

1 **The Ionospheric Source of Magnetospheric Plasma is**
2 **Not a Black-Box Input for Global Models**

D. T. Welling,¹ M. W. Liemohn,¹

Corresponding author: D. T. Welling, Department of Climate and Space, University of Michigan, 2455 Hayward St., Ann Arbor, Michigan 48109, USA. (dwelling@umich.edu)

¹Department of Climate and Space,
University of Michigan, Ann Arbor,
Michigan, USA.

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3 **Abstract.** Including ionospheric outflow in global magnetohydrodynamic
4 models of near-Earth outer space has become an important step towards un-
5 derstanding the role of this plasma source in the magnetosphere. Of the ex-
6 isting approaches, however, few tie the outflowing particle fluxes to magne-
7 tospheric conditions in a self-consistent manner. Doing so opens the magnetosphere-
8 ionosphere system to non-linear mass-energy feedback loops, profoundly chang-
9 ing the behavior of the M-I system. Based on these new results, it is time
10 for the community eschew treating ionospheric outflow as a simple black-box
11 source of magnetospheric plasma.

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1. Introduction

12 The ionosphere is not merely a black-box electrodynamic boundary condition on the
13 magnetosphere-ionosphere system; rather, it is an integral part of the tightly-coupled,
14 non-linear solar wind-magnetosphere-ionosphere system. The Birkeland currents flowing
15 from the magnetosphere close through the ionosphere, creating a large scale electric field.
16 This electric field controls magnetic convection and particle drifts in the magnetosphere,
17 playing a profound role in global dynamics. The interplay between the two creates a
18 system that is non-linear and tightly-coupled. This relationship is widely accepted as fact
19 and the use of these adjectives should be neither controversial nor surprising.

20 This line of reasoning should apply to ionospheric outflow as well. It is now widely
21 accepted that ionospheric outflow is an important, if not dominant, source of magneto-
22 spheric plasma [Chappell *et al.*, 1987]. Similar to the case of ionospheric electrodynamic,
23 the characteristics of the outflowing plasma is tied to the energy input from both the
24 magnetosphere and solar wind (in addition to the above reviews, see *Yau and André*
25 [1997]). From these two statements alone, it would be expected that ionospheric outflow
26 is a *tightly-coupled* and *non-linear* part of the solar-magnetosphere-ionosphere system and
27 should be treated as such by the research community.

28 A review of global numerical modeling literature, however, shows that outflow is treated
29 very differently than electrodynamic. Most global-level studies neglect outflow outright;
30 only a fraction of those that include the ionospheric plasma source connect it to magneto-
31 spheric conditions in a causal, self-consistent manner. In contrast, it is exceedingly rare to

32 find a simulation-based study of the global system that does not include a self-consistently
33 calculated convection electric field.

34 Those studies that do include the ionospheric source of magnetospheric plasma are
35 finding that it has a profound impact on global dynamics. The subset of this minority
36 that includes outflow causally are discovering that previously unobserved mass-energy
37 feedback loops can develop between the ionosphere and magnetosphere. Based on this
38 small but growing body of work, it is time for the numerical modeling community to
39 accept that ionospheric outflow is not merely a black-box input to the magnetosphere,
40 but a critical part of the tightly-coupled, non-linear magnetosphere-ionosphere system.

2. Outflow and the Magnetosphere

41 The case for this argument begins with observations of heavy ions in the magnetosphere,
42 a clear sign of plasma of ionospheric source. The report of plasma sheet oxygen by *Shelley*
43 *et al.* [1974] marked the first of many observations of ionospheric plasma in the plasma
44 sheet [*Lennartsson and Shelley*, 1986; *Nosé et al.*, 2003; *Denton et al.*, 2005; *Mouikis*
45 *et al.*, 2010] and inner magnetosphere [*Sharp et al.*, 1985; *Daglis et al.*, 1999; *Nosé*, 2005;
46 *Kronberg et al.*, 2012; *Nosé et al.*, 2015]. Multiple reviews now cover this topic in depth
47 [*Hultqvist et al.*, 1999; *Yau and André*, 1997; *Yau et al.*, 2007; *Kronberg et al.*, 2014;
48 *Welling et al.*, 2015]. Two themes emerge in these works: the strength of the ionospheric
49 source varies with geomagnetic activity and this source plays a critical role in the storm
50 time magnetosphere.

51 Many studies, both observational and numerical, have tied variability of outflow to
52 different energy sources. Solar energy deposition is tied to increases in outflow, both in
53 terms of increased photoionization of neutrals (via correlation with F10.7 solar radio flux

54 [e.g., *Young et al.*, 1982; *Cully et al.*, 2003]) and production of photoelectrons on the day
55 side [e.g., *Khazanov et al.*, 1997; *Tam et al.*, 1998; *Kitamura et al.*, 2011]. On the day
56 side, especially the cusp region, Alfvénic Poynting flux and soft precipitating electron
57 fluxes are associated with increases in outflow into the magnetosphere [e.g., *Barghouthi*
58 *et al.*, 1998; *Barakat et al.*, 1998; *Barakat and Schunk*, 2001; *Strangeway et al.*, 2005].
59 Impulses in the solar wind dynamic pressure have also been tied to bursts of outflow
60 [*Moore et al.*, 1999]. On the night side auroral zone, outflow is tied to various measures of
61 magnetospheric activity [*Yau et al.*, 1985] and have been directly tied to magnetospheric
62 substorms [e.g., *Øieroset et al.*, 1999; *Wilson et al.*, 2004]. *Gombosi and Nagy* [1989]
63 showed that transient oxygen outflows can be tied to changes in field-aligned current
64 strength. The reviews previously listed cover these topics in more depth.

65 Connecting these lines of research gives us the possibility of tight coupling between
66 ionospheric outflow and magnetospheric activity. Clearly, outflow is important to mag-
67 netospheric dynamics. Conversely, many energy inputs to ionospheric ions, though ulti-
68 mately of solar origin, are filtered through the magnetosphere. As one changes, we expect
69 the other to evolve in response.

3. Outflow and Global Modeling

70 Though evidence concerning the importance of outflow in the magnetosphere was
71 mounting, global magnetohydrodynamic (MHD) models of the magnetosphere were re-
72 luctant to include this source. The first model to purposely include this source was that
73 of *Winglee* [1998], which relied on simple inner boundary conditions to provide an iono-
74 sphere population. Noting intent is necessary as other modelers had been including inner
75 boundary sources via the same mechanism [*Siscoe et al.*, 2001; *Walker et al.*, 2003], but

76 only recognized this after the work of *Winglee* [1998] and did not leverage this feature
77 until some time later [*Zhang et al.*, 2007; *Welling and Ridley*, 2010; *Welling and Liemohn*,
78 2014]. The work of *Winglee* [1998] was an important step for the community because,
79 though particle tracing models have previously explored ionospheric particles in the mag-
80 netosphere [e.g., *Delcourt et al.*, 1989; *Perroomian and Ashour-Abdalla*, 1996; *Huddleston*
81 *et al.*, 2005; *Moore et al.*, 2007], global MHD provided the opportunity to do so in a self-
82 consistent manner (i.e., where the ionospheric population can act on the fields through
83 which it drifts). This breakthrough inclusion led to a series of successful follow up studies
84 that included both ionospheric oxygen and hydrogen [*Winglee*, 2000; *Winglee et al.*, 2002;
85 *Harnett et al.*, 2008], yet no other modeling group followed suit for a decade.

86 The slow pace to adopt outflow stands in stark contrast to the situation concerning
87 ionospheric electrodynamics during the early MHD era [*Brecht*, 1985]. The first 3D MHD
88 simulations of the Earth's magnetosphere [*Brecht et al.*, 1981; *LeBoeuf et al.*, 1981; *Wu*
89 *et al.*, 1981] neglected any connection with the ionosphere (but frequently noted its abs-
90 cence). Only three years later, rudimentary attempts to connect to the ionosphere were
91 being made [*Fedder and Lyon*, 1984; *Ogino and Walker*, 1984]. Second generation MHD
92 models, such as the Block Adaptive Tree Solar wind Roe type Upwind Scheme (BATS-R-
93 US, *Gombosi et al.* [1994]), did not publish Earth simulation results until a full treatment
94 of ionospheric electrodynamics was installed [e.g., *Gombosi et al.*, 1998].

95 Eventually, however, the rest of the MHD community began to realize the benefits
96 of including both light and heavy ionospheric sources. The results were immediately
97 enlightening. The importance of ionospheric plasma to the central plasma sheet and
98 ring current, a result previously explored with observations and particle tracing methods,

99 was confirmed in the global models, but the impacts of this outflow on magnetospheric
100 dynamics could now be explored [*Glocer et al.*, 2009a, b; *Brambles et al.*, 2010; *Welling and*
101 *Ridley*, 2010; *Welling et al.*, 2011]. Controlled, idealized experiments were constructed
102 to test the impact of specific populations, including outflow from specific regions [*Garcia*
103 *et al.*, 2010; *Yu and Ridley*, 2013a] and outflow with different characteristics [*Wiltberger*
104 *et al.*, 2010; *Brambles et al.*, 2010; *Yu and Ridley*, 2013b] in isolation from other regions.
105 The explosion of investigations into the role of outflow in the global system was just
106 starting.

107 Intertwined within the flurry of new results were hints that non-linear feedback loops
108 could manifest. The first hint came via the ubiquitous result that ionospheric outflow could
109 strongly affect cross polar cap potential [*Winglee et al.*, 2002; *Glocer et al.*, 2009a; *Brambles*
110 *et al.*, 2010; *Wiltberger et al.*, 2010; *Welling et al.*, 2011], indicating that mass loading of
111 the magnetosphere is somehow capable of changing its electromagnetic coupling with the
112 ionosphere. The exact nature of this effect remains unresolved [*Welling and Zaharia*,
113 2012]. Further, there were profound effects upon tail dynamics, including the potential
114 for heavy ion outflow to trigger substorms [*Wiltberger et al.*, 2010], alter substorm intensity
115 and onset timing [*Welling et al.*, 2016], and impact the development of Kelvin-Helmholtz
116 stability criteria along the flanks of the magnetosphere. It was also found that as outflow
117 from one region of the ionosphere changed the shape and dynamics of the magnetosphere,
118 the role of outflow from other regions would be changed [*Welling et al.*, 2016], obfuscating
119 results from previous studies. The complex role that ionospheric outflow plays in the
120 global system was being realized.

4. Mass-Energy Feedback

121 In all of these studies, only a handful closed the ionosphere-magnetosphere-ionosphere
122 loop by including a terrestrial plasma source that is *causal*, i.e., one that responds dynami-
123 cally to solar and magnetospheric drivers. Two methods are currently employed. The first
124 is an empirical method that leverages the work of *Strangeway et al.* [2005] to tie Poynting
125 flux, as calculated by the MHD model, to O⁺ outflow rates [*Gagne, 2005; Damiano et al.,*
126 *2010; Brambles et al., 2010*]. The second is to use first-principles-based modeling via the
127 Polar Wind Outflow Model (PWOM [*Glocer et al., 2007, 2012*]), to simulate outflow dy-
128 namics using energy inputs from the magnetosphere [*Glocer et al., 2009a, b*]. Both have
129 been used in a number of studies, as listed above.

130 It was with the empirical approach that the first clear mass-energy feedback loop was
131 discovered within the confines of global models. *Brambles et al.* [2011] found that, using
132 constant solar wind drivers, heavy ion outflow as provided by the adapted *Strangeway et al.*
133 [2005] formula could drive global sawtooth oscillations [e.g., *Huang, 2003; Henderson,*
134 *2004*] in the tail. The mechanism responsible was a chain of O⁺ outflow that would
135 mass-load and destabilize the tail, which in turn would drive a temporary increase in
136 Poynting flux into the inner boundary, therefore releasing a new burst of O⁺ [*Ouellette*
137 *et al., 2013*]. *Brambles et al.* [2013] further explored these dynamics for different categories
138 of storms. These publications marked the first true identification of a possible ionosphere-
139 magnetosphere feedback loop that relied entirely on the inclusion of outflow as a coupling
140 mechanism.

141 The community was intensely interested in the ionosphere-magnetosphere connection
142 illustrated in *Brambles et al.* [2011]. First, it was a potential solution to the mystery

143 of sawteeth oscillations, which are yet to be fully understood. Statistical studies were
144 immediately undertaken to test the veracity of this result, finding that O^+ in the plasma
145 sheet appears to be a necessary but not sufficient condition for sawteeth events [*Liao*
146 *et al.*, 2014]. Second, *Brambles et al.* [2011] was the first to demonstrate that ionospheric
147 outflow can develop non-linear feedback loops with the magnetosphere if causality with
148 the global system is considered. This idea was expanded upon by *Moore et al.* [2014], who
149 chronicled the different ways that outflow may be interacting with the magnetosphere.

150 Following this new found excitement of outflow was the identification of another po-
151 tential feedback loop [*Welling et al.*, 2015]. Using the first-principle-based PWOM model
152 with global MHD and a dedicated ring current model, it was found that region 2 Birkeland
153 currents could increase the amount of outflowing O^+ via a previously known mechanism
154 [*Gombosi and Nagy*, 1989]. The increase in outflow mass loaded the ring current and
155 increased energy density of the ring current significantly. This would, in turn, increase
156 the magnitude of O^+ outflow associated with the region 2 currents. When compared to a
157 similar simulation where the Birkeland currents from the ring current were not allowed to
158 close to the ionosphere (a “one-way” coupled simulation), the “two-way coupled” simula-
159 tion (i.e., the feedback loop is closed via the region 2 currents) yielded a six-fold increase
160 in outflowing O^+ , a ring current dominated by oxygen, and far more realistic behavior in
161 terms of D_{ST} index and Birkeland current distribution.

162 Figure 1 illustrates the dramatic change in outflow when the feedback loop is allowed to
163 close. The solid lines show the total fluence, or outflowing flux integrated over the polar
164 cap, over the duration of the storm when the feedback loop is broken (the simulation is
165 only one-way coupled to the ring current). Colors indicate O^+ (green), H^+ (orange), or

total fluence (black) values. When the feedback loop is closed (dashed lines), total outflow increases by about a factor of two owing almost completely to the increase in oxygen. This line plot clearly demonstrates the necessity of using a ionospheric source of plasma that is self-consistent with the state of the magnetosphere.

5. Forward to Outflow Self-Consistency

Brambles et al. [2011] and *Welling et al.* [2015], two illustrations of outflow-magnetosphere feedback loops, are not without faults. Both are, at best, a limited representation of the true system. Each of the outflow models include only a limited set of acceleration physics (*Strangeway et al.* [2005] representing wave-particle acceleration, PWOM representing ambipolar acceleration). Additionally, coupling from the magnetosphere into the ionosphere is limited. Assumptions must be made to obtain precipitating fluxes from the MHD models, and outflowing fluxes are very sensitive to these. Further, the extent to which different magnetospheric processes are being accurately captured, such as reconnection in the tail, is not well understood. It is reasonable to question if these feedback loops exist in the real system.

However, at a minimum, these studies have demonstrated the need to include two-way coupled outflow in global simulations. In each case, the inclusion of self-consistent outflow fundamentally changed the behavior of both the magnetosphere and of the spatial and temporal dynamics of the ionospheric plasma source. By continuing to ignore the possibility of such feedback loops and treating outflow as a black-box boundary condition, researchers risk severely limiting the capabilities of their models and misinterpreting the results. Now is the time for the community to embrace the challenge of advancing numerical models to include self-consistent ionospheric outflow.

188 Accepting this challenge will require redoubled effort to improve the physics capabilities
189 in our models. Outflow model physics must be expanded to include as many acceleration
190 mechanisms as possible. Currently, very few models have this capability and only one
191 has been coupled to a global model [*Welling et al.*, 2016], albeit in a simple, one-way
192 manner. Additionally, the coupling between outflow and magnetosphere models must be
193 expanded so that all possible energy coupling paths can be included. Our current capabil-
194 ities available for two-way coupling are, at present, singularly focused on either capturing
195 wave-particle interactions or obtaining the ambipolar acceleration. This will likely re-
196 quire capabilities beyond multi-fluid MHD that provide additional information about the
197 precipitating populations and parallel acceleration mechanisms at higher altitudes (e.g.,
198 *Glocer* [2016]). Finally, we need to start considering the problem from the bottom up by
199 capturing densities and upflows from the thermosphere upwards. This will yield a more
200 accurate and dynamic limit on the supply of ions from below. These efforts will require
201 broad commitment from ionospheric and magnetospheric scientists as well as computa-
202 tional specialists. However, the scientific return on such an investment will be deeply
203 beneficial to our understanding of the ionosphere-magnetosphere-heliosphere system.

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Figure 1. Hydrogen, oxygen, and total fluence (orange, green, and black lines, respectively) taken at the interface between the PWOM and BATS-R-US during the 1-way coupled simulation (solid lines) and the 2-way coupled simulation (dashed lines). Reproduced from *Welling et al.* [2015]

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Total Fluence: 1-Way vs. 2-Way Coupling

