Thermodynamics of the monohydrogen difluorides VI. Heat capacities of RbHF₂ and CsHF₂ from 300 to 530 K, transitions and melting*

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The heat capacities of RbHF₂ and CsHF₂ were determined from 5 to 533 K, and hence into the molten-salt region by adiabatic calorimetry. The thermodynamics of the transitions were determined. The entropy of transition ΔS_t of the II-to-I transition in RbHF₂ is $(5.59 \pm 0.04) \text{ cal}_{th} \text{ K}^{-1} \text{ mol}^{-1}$ and the corresponding entropy of melting ΔS_m is $(2.31 \pm 0.03) \text{ cal}_{th} \text{ K}^{-1} \text{ mol}^{-1}$. Both entropy values are comparable to those in the isostructural compound KHF₂. For CsHF₂, ΔS_t (III-to-II), ΔS_t (II-to-I), and ΔS_m are (1.74 ± 0.07) , (2.20 ± 0.08) , and $(1.62 \pm 0.09) \text{ cal}_{th} \text{ K}^{-1} \text{ mol}^{-1}$. The occurrence of the II-to-I transition was first reported on the basis of these studies. At 500 K the standard entropies S° are (47.91 \pm 0.05) and $(50.63 \pm 0.05) \text{ cal}_{th} \text{ K}^{-1} \text{ mol}^{-1}$ for RbHF₂ and CsHF₂.

1. Introduction

The structure and the phase transitions of the alkali-metal hydrogen diffuorides have attracted considerable attention from both theorists and experimentalists. Although initially interest was focused mainly on the single-minimum potential-energy curves for the proton in the strong hydrogen bond of the FHF⁻ ion, originally established by thermophysical calorimetry^(2,3) and later confirmed by other techniques,⁽⁴⁻⁷⁾ more recently interest has centered on the structural and ionic orientational disorder of these interesting ionic crystals consisting of spherical cations and cylindrical (linear) anions and the comparison and correlation of their thermophysical properties and phase behavior with those of the azides.⁽⁸⁻¹²⁾ Interest in the nature of the hydrogen bond in alkali and related substances continues unabated.⁽¹³⁻¹⁶⁾

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From the isostructural (tetragonal) phases of these alkali difluorides⁽¹⁷⁾ transitions to phases of KHF_2 ,^(3, 18) of RbHF_2 ,⁽¹⁹⁾ and of $\text{CsHF}_2^{(19, 20)}$ stable at higher temperatures are known. The higher-temperature phases of KHF_2 and RbHF_2 are isostructural (NaCl-type); that of CsHF_2 is of the CsCl-type structure.

In the ambient-temperature phases the FHF⁻ ion is in an ordered array in the plane perpendicular to the tetragonal axis.⁽¹⁷⁾ The work of Kruh *et al.*⁽¹⁹⁾ has shown that in the higher-temperature phases the FHF⁻ anions are aligned along the cube diagonals for the potassium and rubidium compounds and along the Cartesian axes for the cesium salt. A randomly oriented hydrogen difluoride ion rather than a freely rotating ion is characteristic of the ionic motion in these salts. Unfortunately, extensive thermophysical investigation has been hitherto achieved only on the transition in crystalline KHF₂.⁽¹⁸⁾

Structural differences between the high-temperature phases of KHF_2 and $CsHF_2$ introduce differences in the thermodynamic characteristics of transition and melting. Hence, the present study of the phase transitions and melting of $RbHF_2$ and $CsHF_2$ by precise heat-capacity measurements clarifies the behavior of the rod-shaped anions at the transitions involved and extends the low-temperature heat-capacity studies presented elsewhere.⁽²¹⁾

2. Experimental

SAMPLE PROVENANCE

The high-purity samples synthesized by Burney and Westrum⁽²¹⁾ for low-temperature calorimetry has been stored in Teflon and were utilized for this study. Full details of the preparation are reported in their paper and purities higher than 99.92 moles per cent are claimed. Fractional melting studies of this investigation are consistent with purities of 99.9 moles per cent if account is taken of the decomposition pressure of hydrogen fluoride.

CALORIMETRIC MEASUREMENTS

All measurements were taken in the silver calorimeter of laboratory designation W-22-0 in the Mark IV adiabatic calorimetric thermostat described elsewhere.⁽²²⁾ This apparatus is provided with an electronic shield-control system consisting of three channels of recording circuitry with proportional, rate, and reset modes. These controls maintain the temperature difference between calorimeter and adiabatic shield within 1 mK and reduce heat interchange with the calorimeter to a negligible amount. Platinum resistance thermometer A-5, in the re-entrant well of the calorimeter, was calibrated by the National Bureau of Standards against the IPTS-48. All measurements of mass, time, potential, current, and resistance are referred to standards maintained by the National Bureau of Standards.

The heat capacity of the empty calorimeter was determined separately; its volume is 83.77 cm³. The heat capacities of $RbHF_2$ and $CsHF_2$ represented 72 per cent and 76 per cent of the total measured heat capacity of the sample plus calorimeter. Helium gas (7.2 kPa at 300 K) was added to facilitate thermal equilibration.

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3. Results

The experimental heat capacities for RbHF₂ and CsHF₂ are presented in table 1 in chronological sequence so that the temperature increments used in the measurements may usually be deduced from the differences in the adjacent (mean) temperatures. These results have been adjusted for curvature, *i.e.* for the difference between $\Delta H/\Delta T$ and the corresponding derivative. The results in the transition and melting regions are presented in figures 1 to 3. Comparisons of enthalpy-type determinations with values obtained by integration of the smoothed curve in both transition and normal heat-capacity regions are shown in tables 2 and 3.

The smoothed heat capacities and the thermodynamic functions derived from them are presented in table 4 at selected temperatures. The smoothed heat capacities were obtained by a digital-computer program and checked by comparison with largescale plots. The integrations also were performed by digital computer. For the entropies at 298.15 K, the results of Burney and Westrum⁽²¹⁾ were used. The highertemperature functions have a standard deviation of less than 0.1 per cent. Adjustment

$\langle T \rangle \over K$	$\frac{C_p}{\operatorname{cal_{th}} K^{-1} \operatorname{mol}^{-1}}$	$\frac{\langle T \rangle}{K} = \frac{1}{c}$	$\frac{C_p}{\operatorname{cal_{th}} K^{-1} \operatorname{mol}^{-1}}$	$\frac{\langle T \rangle}{K}$	$\frac{C_p}{\operatorname{cal_{th}} \mathrm{K}^{-1} \operatorname{mol}^{-1}}$	$\frac{\langle T \rangle}{K}$ c	$\frac{C_p}{\text{al}_{\text{th}}\text{K}^{-1}\text{mol}^{-1}}$
	Series I	448.91	94.7	331.44	3900	320.28	22.58
303.83	21.18	450.08	398.9	331.80	134.6	325.99	23.27
322.80	22.92	450.36	2330	333.72	23.21		
329.40	75.0	450.37	6290	336.84	23.33	9	eries X
331.44	445	450.39	2240			AH.(II-t	
336.09	269	450.65	272.9	S	Series VI	Δ11t(11-((AH detn G
345.78	23.77	451.45	123.5	322.80	22.92	452.90	319
355.86	23.99	452.29	225.4	$\Delta H_{t}(III)$	-to-II) detn. D	454 17	64 55
365.80	24.35	452.75	371.8				04,00
375.63	24.64	455 .08	433.5	S	eries VII		YT
385.36	24.95	455.68	31.17	320.27	22.56	56	ries Al
394.98	25.27	460.86	27.78	$\Delta H_{\rm t}({\rm III})$	-to-II) detn. E	$\Delta H_{t}(\Pi - t)$	5-1) detn. H
405.05	25.59	466.25	27.86			452.07	240
415.56	25.90			Se	eries VIII	464.07	27.77
423.94	26.33	Se	eries III	433.65	27.64	4/4.93	27.90
434.06	27.52	431.10	26.81	$\Delta H_{\rm t}({\rm II})$	to-I) detn. F	485.56	28,09
443.61	35.66	$\Delta H_{\rm t}({\rm II-to})$	o-I) detn. A	450.41	901	496.35	28.08
449.19	189.7	$\Delta H_{\rm m}$ det	n. B	450.53	242	507.06	28.22
450.25	4132			450.84	101.8	517.74	28.32
451.28	197	Se	eries IV	451.40	121.5	528.40	28.44
452.76	353	$[\Delta H_{\rm t} ({ m II}-$	to-I)	451.86	164.2		
457.95	34.2	+	$\Delta H_{\rm m}$] detn. C	452.20	216.9	Se	ries XII
468.20	27.75					306.20	21.24
		S	eries V		Series IX	307.49	21.34
	Series II	331.13	165.2	292.96	20.44	308.78	21.47
428.72	26.84	331.41	5422	298.4 4	20.87	310.07	21.60
441.78	31.11	331.42	19000	303.86	21.21	311.35	21.59
446.09	42.18	331.43	13940	309.21	21.53	312.63	21.79

TABLE 1.	Experimental	heat capa	icity deter	minations	on	RbHF ₂	and	CsHF ₂
		(cal	h = 4.184	I J)				

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$\langle T \rangle$	C_p	$\langle T \rangle$	C_p	$\langle T \rangle$	C_p	$\langle T \rangle$	C _p
K	calth K ⁻¹ mol ⁻¹	K	$cal_{th} K^{-1} mol^{-1}$	ĸ	$\operatorname{cal_{th}} K^{-1} \operatorname{mol}^{-1}$	ĸ	$cal_{th} K^{-1} mol^{-1}$
	······································	R	ubidium hydrogen	difluorid	e (RbHF ₂)		
	Series I	397.65	22.23	444.71	3610	522.66	28.71
307.88	19.33	406.85	22.91	448.32	113.4	530.16	28.79
318.30	19.59	415.90	22.99	S	eries VIII	S	eries XIII
328.60	19.86		Series V	415.90	22.99	481.02	214
338.55	20.10	429.90	23.84	392.11	21.82	481.25	271
348.62	20.40	438.21	27.44	398.34	22.20	481.47	324
358.56	20.68	442.99	136.8	401.36	22.40	481.68	334
368.37	20.99	444.12	422	404.33	22.76	481.84	329
378.07	21.27	444.40	823	407.31	23.01	481.99	368
398.87	22.37	444.55	1410	410.28	22.68	482.14	406
409.22	23.12	444.64	1750	413.24	22.71	482.26	451
419.41	23.24	444.71	2430	416.22	23.05	482.38	541
429.45	24.07	444.77	3250	419.19	23.18	482.47	600
	Sarias II	444.81	4980		Sarias IV	482.56	643
127 71	26 10	444.85	4130		n R	482.64	719
437.14	20.19	444.88	5364	Δm_t uci		482.72	707
AAA 26	633	444.91	11000		Series X	482.80	591
111.20	011	444.93	4420	$\Delta H_{\rm t}$ det	in. C	48 2.9 9	186.3
AAA 5A	1380	444 .9 6	937	462.15	23.05	483.88	27.79
444 63	1650	444 .9 8	4390	470.91	24.50	S	eries XIV
444.60	2320	445.03	2770	477.31	56.70	473 07	24.93
AAA 75	2270	445.74	104.5	480.30	154.6	AH. de	tn D
444.75	2210	447.59	22.68	481.57	313	Ling ut	
	Series III	450.54	22.68	482.28	476	3	eries XV
444 76	1211		1	483.15	229	404.39	23.13
444.83	2890	204.00	Series VI	487.48	28.99	4/1.28	24.89 Am E
444.87	4320	394.23	21.98	495.03	28.30	$\Delta H_{\rm m}$ de	an, E
444.89	5690	401.08	22.22	5	Series XI	Se	eries XVI
444.91	6350	404.00	22.56	477.74	61.40	$\Delta H_{\rm m}$ de	tn. F
444.93	6680	407.02	22.78	480.75	212	Se	ries XVII
444.96	5320	409.90	22.75	481.89	368	$\Delta H_{\rm t}$ (II-t	o-I) detn. G
444.99	4980	412.93	22.70	482.67	434	445.26	535
445.03	3230	412.88	22.79	486.35	40.07		1 3232777
448.32	22.82	410.03	22.92	S	eries XII	Se	ries XVIII
		421.75	23.10	492.40	28.19	$\Delta H_{t}(11-$	(0-1) detn. H
1	Series IV	S	eries VII	500.01	28.32	447.84	23.03
378.70	21.32	$\Delta H_{\rm t}({\rm II}-$	to-I) detn. A	507.60	28.39	S	eries XIX
388.25	21.70	444.09	84.6	515.16	28.48	$\Delta H_{t}(II-$	to-I) detn. J

TABLE 1-continued

has not been made for isotope-mixing nor for nuclear-spin contributions to $\{S^{\circ}(T) - S^{\circ}(0)\}$ or to $-\{G^{\circ}(T) - H^{\circ}(0)\}/T$; hence, these values are practical ones for use in chemical thermodynamic calculations.

4. Discussion

TRANSITION AND MELTING OF RbHF₂

Corresponding $RbHF_2$ and KHF_2 phases are isostructural and both compounds have phase transitions from the tetragonal structure (phase II) with the space group



FIGURE 1. The heat capacity of RbHF₂ near the II-to-I and melting transitions.



FIGURE 2. The heat capacity CsHF2 near the III-to-II transition.



FIGURE 3. The heat capacity of CsHF₂ near the II-to-I and melting transitions.

	No. of	T_2	T_1	$\underline{H^{\circ}(T_2)}-\underline{H^{\circ}(T_1)}$	$H^{\circ}(452 \text{ K}) - H^{\circ}(420 \text{ K})$
	detns.	K	ĸ	cal _{th} mol ⁻¹	cal _{th} mol ⁻¹
TRANSITION I	I-to-I				
Series V	20	451.53	425.51	3103	3249 ª
Series VII	4	451.82	424.95	3114	3236
Series XVIII	2	449.34	437.66	2740	3228
Series XIX	1	464.83	434.39	3183	3230
			Average	value $(H^{\circ}(452 \text{ K}) - 1)$	H°(420 K)}: 3232 ± 2
		La	attice contri	bution {H°(452 K) -	<i>H</i> °(420 K)}: 746 ± 4
		$\Delta H_{\rm t}$ (II-t	o-I)/cal _{th} m	$ol^{-1} = 2486 \pm 5$	
	Δ.	St(II-to-I)/	$cal_{th} K^{-1} m$	$101^{-1} = 5.59 \pm 0.4$	
		T_{t}	(II-to-I)/K =	= 444.93 ± 0.03	

TABLE 2. Thermodynamics of transition II-to-I and melting of $RbHF_{\rm 2}$ (cal_{th} = 4.184 J)

	No. of detns.	$\frac{T_2}{K}$	$\frac{T_1}{K}$	$\frac{H^{\circ}(T_2) - H^{\circ}(T_1)}{\operatorname{cal_{th}} \mathrm{mol}^{-1}}$	$\frac{H^{\circ}(496 \text{ K}) - H^{\circ}(460 \text{ K})}{\text{cal}_{\text{th}} \text{ mol}^{-1}}$
MELTING ^b					
Series X	8	491.22	457.66	1941	2023
Series XIV	2	491.22	470.54	1634	2016
Series XV	2	490.18	467.05	1692	2020
Average value { $H^{\circ}(496 \text{ K}) - H^{\circ}(460 \text{ K})$ } Lattice contribution { $H^{\circ}(496 \text{ K}) - H^{\circ}(460 \text{ K})$ } $\Delta H_{m}/cal_{tb} \text{ mol}^{-1} = 1117 \pm 5$ $\Delta S_{m}/cal_{tb} \text{ K}^{-1} \text{ mol}^{-1} = 2.31 \pm 0.03$ $T_{m}/\text{K} = 482.67 \pm 0.04$					H°(460 K)}: 2020 ± 3 H°(460 K)}: 903 ± 4

TABLE 2-continued

 ^a Rejected from average because of incertainties in the drift corrections.
 ^b Melting determinations include "bump" 481.1 to 481.6 K which has 32.4 cal_{th} mol⁻¹ of excess enthalpy.

TABLE 3. Thermodynamics of transitions III-to-II and II-to-I and melting for CsHF2 $(cal_{th} = 4.184 \text{ J})$

	No. of	$\frac{T_1}{K}$	T_2	$\frac{H^{\circ}(T_2)-H^{\circ}(T_1)}{2}$	$\frac{H^{\circ}(340 \text{ K}) - H^{\circ}(315 \text{ K})}{100 \text{ mol}^{-1}}$			
		<u> </u>	<u> </u>	caith mor	calth mor			
TRANSITION	III-to-II							
Series I	4	318.15	340.72	1099.2	1151.7			
Series VI	2	316.76	337.25	1044.8	1149.9			
Series VII	2	316.35	337.12	1051.5	1150.2			
	ΔΙ	La La (III-to-l	Average ttice contrib	value $\{H^{\circ}(340 \text{ K}) - H^{\circ}(340 \text{ K}) - H^{\circ}(340 \text{ K}) - H^{-1} = 580.3 + 2$	<i>H</i> °(315 K)}: 1150.6 ± 0.7 H°(315 K)}: 570.3 ± 2			
	$\Delta S_{t}(II)$	I-to-II)/ca	h K ⁻¹ mol	$^{-1} = 1.74 \pm 0.07$				
			T _t (III-to-II)	$/K = 331.42 \pm 0.02$				
	No. of	<i>T</i> ₁	T2	$H^{\circ}(T_2) - H^{\circ}(T_1)$	$H^{\circ}(470 \text{ K}) - H^{\circ}(420 \text{ K})$			
	detns.	ĸ	ĸ	cal _{th} mol ⁻¹	cal _{th} mol ⁻¹			
TRANSITION	II-to-I AND	MELTIN	ſG		······································			
Series I	9	418.77	473.72	3122.3	2998.3ª			
Series II	15	439.34	468.96	2437.7	2999.1ª			
Series III	3	428.30	460.45	2509.1	2995.8			
Series IV	3	419.88	460.16	2721.9	2995.7			
	Average value $\{H^{\circ}(470 \text{ K}) - H^{\circ}(420 \text{ K})\}$: 2995.8 \pm 0 Lattice contribution $\{H^{\circ}(470 \text{ K}) - H^{\circ}(420 \text{ K})\}$: 1349 \pm 4 $\{\Delta H_{*}(\text{II-to-I}) + \Delta H_{m}\}$ (calls mol ⁻¹ = 1646 + 4							
(TRANSITION	II-to-I)							
	-	$\Delta H_{\rm t}$ (II-to	-I)/cal _{th} mo	$d^{-1} = 991 \pm 4$				
	ΔS_{t}	(II-to-I)/ca	llth K ⁻¹ mo	$1^{-1} = 2.20 \pm 0.08$				
			$T_{\rm t}$ (II-to-I)/	$K = 450.36 \pm 0.02$				
(MELTING)								
		Δ	$H_{\rm m}/{\rm cal_{th}}~{\rm mc}$	$d^{-1} = 655 \pm 4$				
		$\Delta S_{\rm m}/{\rm cal}$	™ K ⁻¹ mol ⁻ T_/K	$^{-1} = 1.62 \pm 0.09$ = 453.20 + 0.08				
			- 17/1					

^a Rejected from average because of magnitude (12 cal_{th} mol⁻¹) of drift corrections.

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TABLE 4. Thermodynamic properties of RbHF2 and CsHF2

<u>T</u>	<u></u>	$\underline{S^{\circ}(T) - S^{\circ}(0)}$	$\frac{H^{\circ}(T)-H^{\circ}(0)}{1-H^{\circ}(0)}$	$- \{G^{\circ}(T) - H^{\circ}(0)\}/T$
ĸ	$\operatorname{cal}_{\operatorname{th}} \mathrm{K}^{-1} \operatorname{mol}^{-1}$	$cal_{th} K^{-1} mol^{-1}$	cal_{th} mol ⁻¹	$cal_{th} K^{-1} mol^{-1}$
		Rubidium hydrogen d	ifluoride (RbHF ₂)	
CRYSTA	AL II		(110111 2)	
298.15	19.03	28.70	3932	15.51
300	19.07	28.82	3967	15.59
325	19.65	30.36	4451	16.67
350	20.36	31.85	4951	17.70
375	21.34	33.28	5472	18.69
400	22.35	34.69	6018	19.65
425	23.58	36.08	6590	20.57
444.93 ª	(11000)	37.16	7060	21.29
CRYSTA	L I (Transition II-to	-I)		
444.93 ª	(11000)	42.73	9545	21.29
475	26.47	44.25	10245	22.68
482.67 ª	(720)	44.60	10415	23.02
LIQUID				
482.67 °	(720)	46.91	11531	23.02
500	28.29	47.91	12020	23.87
525	28.75	49.30	12732	25.05
		Cesium hydrogen dif	luoride (CsHF2)	
CRYSTA	L III			
298.15	20.83	32.30*	4199°	18.22 *
300	20.93	32.43	4238	18.30
325	23.08	34.18	4784	19.46
331.42 °	(≈ 20000)	34.63	4929	19.76
CRYSTA	L II (Transition III-	to-II)		
331.42 ª	(≈ 20000)	36.38	5509	19.76
350	23.85	37.66	5946	20.67
375	24.63	39.31	6552	21.84
400	25.42	40.92	7178	22.97
425	26.32	42.50	7823	24.09
450.36 ª	(≈ 8000)	44.03	8497	25.16
CRYSTA	L I (Transition II-to	-I)		
450.36 °	(≈ 8000)	46.23	9488	25.16
453.2 ª	(≈ 600)́	46.38	9566	25.27
LIQUID	(Melting)			
453.2 ª	(≈ 600)	48.00	10222	25.27
475	27.91	49.19	10828	26.39
500	28.14	50.63	11529	27.57
525	28.43	51.94	12236	28.63

 $(cal_{th} = 4.184 J)$

^a Assuming the transitions involved to be entirely isothermal. ^b These values have been revised on better knowledge of the heat capacity near 300 K but differ from those reported in reference (20) by less than 0.05 per cent.

 D_{4b}^{18} -I 4/mcm stable at 300 K into a NaCl-type cubic structure (phase I), O_{b}^{5} -Fm3m,⁽¹⁹⁾ stable at higher temperatures until melting occurs. For KHF₂ Phase I is a soft translucent wax-like solid at 469.8 K. The entropy increment of 5.69 cal_{th} K⁻¹ mol⁻¹ for the II-to-I transition is nearly twice that of melting.⁽¹⁷⁾†

The temperatures of transition and melting of RbHF_2 were determined by the present study as $T_t = 444.93$ K and $T_m = 482.67$ K, both of which differ considerably from the values reported by Kruh *et al.*,⁽¹⁹⁾ ($T_t = 449$ K and $T_m = 461$ K, determined by using a polarizing microscope equipped with Kofler hot stage). Prideaux and Webb⁽²³⁾ reported T_m as 477 to 478 K. Very recently, White and Pistorius⁽⁹⁾ found the values of T_t and T_m at atmospheric pressure as $(448_2 \pm 1)$ K and (476 ± 1) K, respectively, with a high-pressure d.t.a. apparatus. These discrepancies are probably due in part to the purity of the samples and in part to the methods of determination of the transition and melting temperatures.

Entropies of transition and melting of $RbHF_2$ (5.59 and 2.31 cal_{th} K⁻¹ mol⁻¹, respectively) are comparable to well known values for KHF₂ (i.e. 5.69 and 3.09 $\operatorname{cal}_{th} \mathrm{K}^{-1} \operatorname{mol}^{-1}$.⁽¹⁸⁾ The ΔV_t for the transition of RbHF₂ is calculated from the X-ray data⁽¹⁹⁾ as 7.53 cm³ mol⁻¹, which also corresponds with that of KHF₂, calculated as 5.6 cm³ mol⁻¹ from p, V data, $^{(10)}$ 4.2 cm³ mol⁻¹ from X-ray diffraction data,⁽¹⁹⁾ and determined directly by dilatometry as (4.2 ± 0.1) cm³ mol⁻¹ in this laboratory.⁽²⁷⁾ The large entropy and volume increments accompanying the transitions of KHF₂ and RbHF₂ are not easy to explain. Kruh et al.⁽¹⁹⁾ suggested that in phase I each hydrogen difluoride ion would randomly occupy positions along the body diagonals of the unit cell. The expected disordering entropy increment would be close to $R \ln 4$ (2.75 cal_{th} K⁻¹ mol⁻¹). The experimental values of entropy increments, 5.69 and 5.59 cal_{th} K^{-1} mol⁻¹ for KHF₂ and RbHF₂, respectively, are much larger. The large entropy increments are not easily explained but may accompany the structural changes as well as the considerable volume increments which is larger for the former $(4.2 \text{ cm}^3 \text{ mol}^{-1} \text{ for KHF}_2)^{(19,27)}$ and 3.8 cm³ mol⁻¹ for RbHF₂)⁽¹⁰⁾ as is the ΔS_t . This suggests that phase I has additional degrees of orientational freedom, that the structural change itself is enhancing the transitional entropy, and that other contributions are involved, e.g. $(\partial S/\partial V)_T$.

TRANSITION AND MELTING OF CsHF₂

Although it has been known that $CsHF_2$ has a transition near 330 K,^(19,20) the present investigation shows that, in addition to this first transition at 331.42 K, a second sharp heat-capacity anomaly occurs at 450.36 K only 2.8 K below the melting temperature of 453.2 K. The two heat-capacity anomalies are nearly separate and thus clearly indicate the existence of a second transition, phase II \rightarrow phase I, in CsHF₂.

The lower temperature (III-to-II) transition was first reported by Windsor and Cady⁽²⁰⁾ at 330 K and confirmed by Kruh *et al.*⁽¹⁹⁾ but at 334 K. In an extensive high-temperature X-ray diffraction study, they concluded that phase III of CsHF₂ was isostructural with phase II of KHF₂ and transforms to the CsCl-type, cubic structure with the space group O_h^2 -Pm3m.

† Throughout this paper $cal_{th} = 4.184$ J.

The structural features of the CsHF₂ (III-to-II) transition are very similar to those of univalent azides such as RbN₃, CsN₃, T1N₃, etc.,^(11,12) which have linear N₃⁻ anions instead of HF₂⁻ anions. Furthermore, the entropy increment of 1.75 cal_{th} K⁻¹ mol⁻¹ for this CsHF₂ transition as obtained in the present study, is in good accord with the entropy increments of 1.9 and 1.8 cal_{th} K⁻¹ mol⁻¹ reported for RbN₃ and CsN₃.⁽⁹⁾ The small ΔV_t is calculated as 0.59 cm³ mol⁻¹ from the X-ray data and is also comparable to the values for these azides: $\Delta V_t(RbN_3) = 0.5$; $\Delta V_t(CsN_3) = 0.3$ cm³ mol⁻¹.⁽⁸⁾ Thus, taking into account both thermodynamic and X-ray diffraction data, the nature of the III-to-II transition of CsHF₂ can be regarded as essentially the same as that of the transition in these azides. Mueller and Joebstl⁽²⁴⁾ pointed out that the entropy increment at these azide transitions, assuming completely random orientation of the N₃⁻ anions parallel to the edges of the unit cell for the hightemperature phase, is expected to be $R \ln 3$ (2.18 cal_{th} K⁻¹ mol⁻¹).⁽²⁴⁾ Experimental values of ΔS_t would be less than $R \ln 3$ if the length of the linear anion is large enough relative to the unit-cell dimensions to make collinear alignment of two or three ions

Substance	Phase change	$\frac{T_{t}}{K}$	$\frac{\Delta H_{\rm t}}{{\rm cal_{th}mol^{-1}}}$	$\frac{\Delta S_{\rm t}}{{\rm cal_{th}K^{-1}mol^{-1}}}$	$\frac{\Delta V_{\rm t}}{\rm cm^3 \ mol^{-1}}$
KHF ₂	II-to-I	469.8ª	2671 ª	5.69ª	4.2°
_		± 0.05	± 14	± 0.06	4.2 ^{<i>d</i>}
RbHF ₂	II-to-I	444.93	2486	5.59	[3.8]*
-		\pm 0.03	± 5	± 0.04 [5.49] ^b	[]
CsHF ₂	III-to-II	331.42	580.3	1.74	[0.24] ^b
		\pm 0.02	± 2	± 0.07 [3.10] ^b	[1.0] < 0.59°
CsHF ₂	II-to-I	450.36	991	2.20	י ך3.2
		± 0.02	± 4	± 0.08 [2.0] ^b	to 4.9
Substance	Phase	$T_{\rm m}$	$\Delta H_{ m m}$	ΔS_{m}	$\Delta V_{\rm m}$
Substance	change	K	cal _{th} mol ⁻¹	$cal_{th} K^{-1} mol^{-1}$	cm ³ mol ⁻¹
KHF ₂	melting	511.8ª ±0.1	1582ª ± 6	3.09 ^{<i>a</i>} ± 0.05	[0.8] ′
RbHF	melting	482.67	1117	[3.09] ^o 2 31	[0 96] 0
		\pm 0.04	± 5	± 0.03 [2.80] ^b	[0120]
CsHF2	melting	453.2 ± 0.08	665 ± 4	1.62 ± 0.09 [1.60] *	[1.9]

TABLE 5. Transition and melting thermodynamics for some alkali dihydrogen fluorides $(cal_{th} = 4.184 \text{ J})$

^a Reference 18. ^b Values in brackets [], are from high-pressure phase studies summarized in reference 10. ^c Reference 19. ^d Reference 27. ^e At 160 kPa, reference 10. ^f At 220 kPa, reference 10.

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HEAT CAPACITIES OF RbHF₂ AND CsHF₂

in adjacent cells improbable for steric reasons—as is also the case for the alkali azides. A calculation of the correlation energy for succino-nitrile has been made by Descamps and Coulon, $^{(25, 26)}$ but the CsHF₂ case has not been evaluated.

Our detection of the II-to-I transition in $CsHF_2$ near 450 K was subsequently confirmed by White and Pistorius⁽⁹⁾ by d.t.a. They found that this transition occurs at atmospheric pressure, but that at 1.3 MPa and higher pressures only melting was observed. Although their measured T_t (II-to-I) is lower (447.4 K) and that of melting (459 K) is higher than the present values, other features are essentially the same.

White and Pistorius⁽⁹⁾ estimated the volume change associated with this transition from our ΔS_t and ΔS_m of 2.20 and 1.45 cal_{th} K⁻¹ mol⁻¹, respectively, and the trend of the melting curve with pressure, to be between 3.2 and 4.9 cm³ mol⁻¹. The large ΔS_t and ΔV_t suggest that a part of the increments is due to structural randomness. Many CsCl-type crystals undergo transition into a NaCl-type structure at high temperatures and/or low pressures, with a decrease in coordination from 8- to 6-fold. Direct structural analysis of this phase by high-temperature X-ray diffraction is an obvious desideratum, but the temperature range in which CsHF₂ phase I exists is so narrow the X-ray study will not be convenient. Hence, we assume following the arguments of White and Pistorius⁽¹⁰⁾ based on the large ΔV_t (II-to-I) that phase I of CsHF₂ may have an NaCl-type cubic structure with the space group O⁵_h-Fm3m, as is characteristic of phase I KHF₂ and RbHF₂.

It will be noticed in table 5 that the calorimetrically determined ΔS_t and ΔS_m values are typically in good accord with the less accurately determined quantities from high-pressure phase-equilibrium studies. Although one finds increasing evidence of interest in the comparisons between the linear anionic orientational disorder and thermophysical parameters in the alkali azides and the hydrogen difluorides^(9, 10-12) and in the related thermodynamic characteristics, pressures, and mechanisms of the decomposition processes in the alkali (and alkaline earth) hydrogen difluorides⁽²⁸⁾ much work remains to be done before a comprehensive appreciation of these compounds is achieved.

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