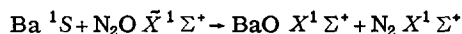
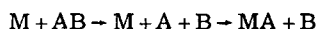


is that our original diagram² focusses on the transition state region. Our Ba+N₂O correlation diagram, augmented by local Ba+O+N₂ correlation and energetic arguments, predicts that the adiabatic (spin conserving)

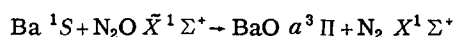


surface may not be the lowest energy surface in the transition region. The augmented correlation diagram was drawn in an attempt to clarify an apparently paradoxical situation by suggesting that a smaller activation energy permits a spin nonconserving reaction to occur with larger yield than a spin conserving reaction.² Augmented correlation diagrams of the type



may prove to be useful for understanding reactions which, as in the Ba+N₂O reaction, involve transfer of an initially weakly bound atom.

Several mechanisms have been postulated in an attempt to explain the high BaO *a*³Π yield from the Ba+N₂O reaction. These include direct but nonadiabatic formation of BaO *a*³Π from ground state reactants²; direct adiabatic formation of BaO *a*³Π from metastable Ba ³D which is formed by a precursor reaction⁵; direct adiabatic formation of BaO *X*¹Σ⁺ (*v* >> 0) from ground state reactants followed by collision induced transitions into *a*³Π^{3b}; collision induced surface crossing in a long lived Ba-N₂O complex.⁵ If only because of its simplicity we continue to prefer the direct nonadiabatic mechanism



and do not believe that this mechanism is ruled out by the correlation arguments of Husain and Wiesenfeld.⁵ It remains to be determined whether the total Ba+N₂O reactive cross section for ground state reactants,⁹ measured under single collision conditions, is too large to be consistent with a 20–100% yield of products formed via a spin nonconserving reaction which should be well described by the weak spin-orbit coupling limit.

¹C. R. Jones and H. P. Broida, *J. Chem. Phys.* **60**, 4369 (1974).

²R. W. Field, C. R. Jones, and H. P. Broida, *J. Chem. Phys.* **60**, 4377 (1974).

³(a) D. J. Eckstrom, S. A. Edelstein, and S. W. Benson, *J. Chem. Phys.* **60**, 2930 (1974); (b) S. A. Edelstein, D. J. Eckstrom, and D. L. Huestis, Status Report 2 (March 1974) for SRI Project PYU-3190.

⁴C. J. Hsu, W. D. Krugh, and H. B. Palmer, *J. Chem. Phys.* **60**, 5118 (1974).

⁵D. Husain and J. R. Wiesenfeld, *J. Chem. Phys.* **62**, 2010 (1975), accompanying Comment.

⁶R. W. Field, *J. Chem. Phys.* **60**, 2400 (1974).

⁷Semiempirical estimates of diagonal and off-diagonal spin-orbit matrix elements are made by replacing molecular orbitals of the dominant electronic configuration by a dominant atomic orbital and using the observed spin-orbit splitting of an appropriate atomic state to estimate the radial part of the spin-orbit matrix element.

⁸A. G. Gaydon, *Dissociation Energies and Spectra of Diatomic Molecules* (Chapman and Hall, London, 1968), Appendix 1.

⁹An upper limit of 3 × 10⁻¹⁶ cm² is obtained by C. D. Jonah, R. N. Zare, and Ch. Ottinger [*J. Chem. Phys.* **56**, 263 (1972)], who measure reactive plus nonreactive attenuation of a Ba atomic beam by N₂O.

ERRATA

Erratum: Crystal-field model study of the xenon hexafluoride molecule. I. Energy levels and molecular geometry [*J. Chem. Phys.* **60**, 3901 (1974)]

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The calculated *t*_{1u} bending frequency reported as 109 cm⁻¹ on pp. 3908 and 3913 is in error. The zero point energy is also in error. The correct frequency is 240 cm⁻¹ using a harmonic oscillator approximation matching the curvature of the potential energy at the C_{4v} radial minimum. Alternatively a fit of the potential $V(x) = -ax^2 + bx^4$ to the position and depth of the C_{4v} radial minimum yields¹ a spacing of 319 cm⁻¹ between the first two levels (ignoring inversion splittings) and a zero point energy of 163 cm⁻¹. The reduced mass of 8.7 × 10⁻²³ g molecule⁻¹ = 52.4 amu molecule⁻¹ used for this mode is $4m_F m_{Xe} / (m_F + m_{Xe})$ (*m*_F and *m*_{Xe} are fluorine and xenon masses), an expression which is eight times larger than that obtained using Wilson's high frequency approximation² for the separation of the *t*_{1u} bending and stretching modes. However, our treatment is equivalent to the latter, as it is

used with a molecular force constant eight times larger than that associated with the *t*_{1u} symmetry coordinate in Eq. (27) on p. 3906. The *t*_{1u} displacement is taken as the bond length times the change in a F-Xe-F angle. For the *t*_{2g} mode our mass is 4*m*_F rather than $\frac{1}{4}m_F$ and is used with a force constant 16 times larger than that associated with the *t*_{2g} symmetry coordinate in Eq. (27). The displacement for this mode is the bond length times one-half of the change in a F-Xe-F angle. For the *t*_{2u} mode our mass is 4*m*_F rather than $\frac{1}{2}m_F$ and is used with a force constant eight times larger than that associated with the *t*_{2u} symmetry coordinate in Eq. (27).

The pseudorotational frequency 2*B*_g estimated as 0.61 cm⁻¹ on p. 3908 is in error by a factor of eight and should be 4.9 cm⁻¹. In addition, the value of Δ*E*_{5g-5p}

for $\text{Te}(+4)$ in Table III should be 11.10 eV instead of 16.10 eV, the $I(s, 0; p, 0)$ matrix element in Eq. (20) should contain the factor R_k^{-3} instead of R_k^{-1} , the two references to Eq. (35) in the second column of p. 3911 should be instead to Eq. (40), and the subscript on Q_5

in Eq. (59) for $V_{445}(C_{2v})$ should be 110 instead of 100.

¹S. Y. Wang and L. L. Lohr, Jr., *J. Chem. Phys.* 61, 4110 (1974).

²E. B. Wilson, Jr., *J. Chem. Phys.* 9, 76 (1941).

Erratum: Rates of collision-induced emission from metastable $\text{O}(^1S)$ atoms [J. Chem. Phys. 61, 1118 (1974)]

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Table II and associated discussion should be corrected to show that the apparent rate constant for collisionally induced emission from $\text{O}(^1S)$ in xenon ($k_g + k_c$) is 25% of the total deactivation rate constant (k_d), rather than 40% as asserted in the paper.

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Erratum: Some theoretical aspects of chemically-induced dynamic nuclear polarization [J. Chem. Phys. 61, 1517 (1974)]

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Printer's errors were made in Eq. (3.8) and (3.6). They should read

$$P^\infty(\text{R.I.})/\mathcal{F}(\text{R.I.}) = -P_{h=0}^\infty(S) = +P_{h=0}^\infty(T) . \quad (3.6)$$

$$\Lambda = k\tau_1/(1 + k\tau_1) . \quad (3.8)$$