

ALPHA-CLUSTER PICKUP REACTIONS ON EVEN Sn AND Te NUCLEI*

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Spectra (Fig.1) for $\text{Te}(d, {}^6\text{Li})\text{Sn}$ and $\text{Sn}(d, {}^6\text{Li})\text{Cd}$ have been obtained for $E_d=33$ MeV at BNL. Complete angular distributions were measured for ${}^{122}\text{Te}(d, {}^6\text{Li}){}^{118}\text{Sn}$. The g.s. spectroscopic strength closely follows that for (p,t) and (t,p)¹ demonstrating the close correlation^{2,3} between 2- and 4-nucleon transfer reactions. For excited states, however, the (d, ${}^6\text{Li}$) strength sometimes exceeds that of (p,t) considerably. The proton pairing-vibration state in ${}^{118}\text{Sn}$ at 1758 keV, for example, is much stronger in (d, ${}^6\text{Li}$). For states with $J>0$ the increased strength is due to coherent combinations (L_π, L_ν) [$\vec{L}_\pi + \vec{L}_\nu = \vec{L} = \vec{J}$], whereas in (p,t) only $L_\nu = J$ is possible.

Semi-microscopic calculations have been performed for selected transitions using RCS⁴ and pairing wave functions. The equations of

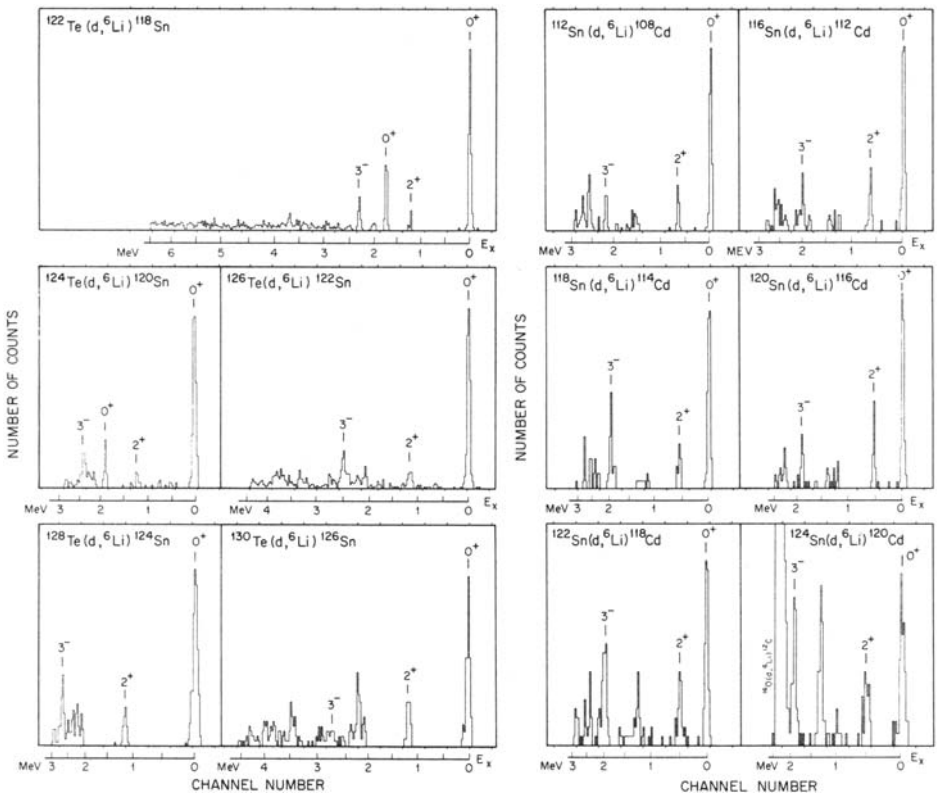


Fig.1. Spectra of ${}^6\text{Li}$ particles from (d, ${}^6\text{Li}$) reactions at $\theta_{\text{lab}}=16^\circ$.

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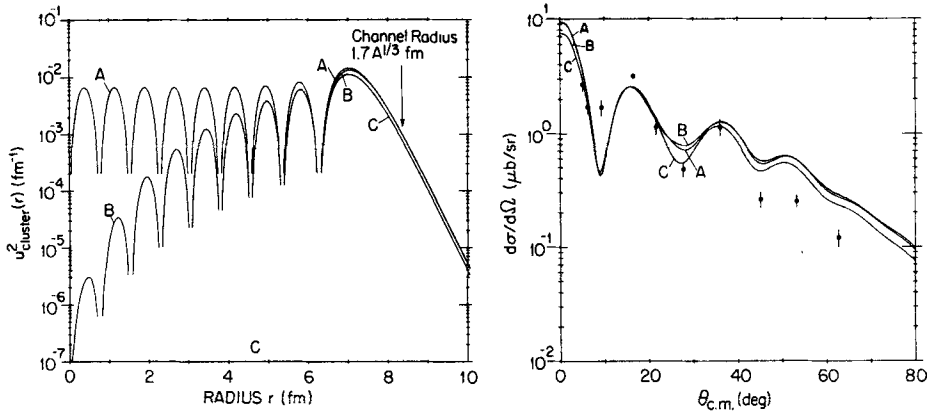


Fig.2. Cluster wave functions and angular distributions for ^{122}Te ($d, {}^6\text{Li}$) $^{118}\text{Sn}(\text{g.s.})$. A: Macroscopic; B: Microscopic; c: Macroscopic with $u^2(r) \equiv 0$ inside $r=6.3$ fm.

Kurath and Towner⁵ were employed with α -parentage amplitudes in SU_3 notation and wave functions generated in Woods-Saxon potential wells. About 60 contributions have to be considered for 0^+ states, about 200-300 for 2^+ and 3^- states. The interaction $V_{d\alpha}$ was not treated microscopically. Fig.2 shows the microscopic cluster wave function for $^{118}\text{Sn} + \alpha$. In the exterior region it is practically identical to the macroscopic function. Enhancement factors $\epsilon = \sigma(\text{exp})/\sigma(\text{calc}) = 3.1$ and ≈ 1.1 for the g.s. transitions to ^{118}Sn and ^{114}Cd were found. The 0^+ pairing-vibration state in ^{118}Sn has $\epsilon \approx 0.6$ in excellent agreement with the strength observed in $({}^3\text{He}, n)$.⁶ The weak 0^+ state at 2057 keV is well described ($\epsilon \approx 1$) as a 2-quasi-particle state. Transitions to 2^+ and 3^- states are well described ($\epsilon \approx 1$) assuming only contributions from $(L_\pi, L_\nu) = (0, J)$ and $(J, 0)$. One might expect, though, that for $J > 3$ additional contributions become important. The dependence on N for the observed $2^+, 3^-, 4^+, 5^-$ and 7^- strengths is at least qualitatively understood in terms of the BCS and pairing wave functions. The pronounced increase in 7^- strength, for example, results from the stretched $(2d_{3/2} 1h_{11/2})_7$ configuration. The drastic decrease for the 0^+ proton pairing-vibration states is not fully understood.

Using the known α -decay of ^{148}Sm , earlier $^{148}\text{Sm}(d, {}^6\text{Li})^{144}\text{Nd}$ data³ make it possible to deduce the normalization constant N needed in zero-range DWBA calculations. Depending on the optical parameter sets used, $N = 3.8 \pm 25\%$. Finite-range calculations are in good agreement with this value.

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