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Supporting Information

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Unlocking the Single-Domain Epitaxy of Halide Perovskites

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Supporting Information

Unlocking the Single-Domain Epitaxy of Halide Perovskites

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Movies S1



Figure S1. Schematic crystal structure of cubic CsSnBr₃. (A) monolayer and (B) bilayer (i.e., 1 unit cell). Green spheres are Cs; gray spheres are Sn; red spheres are Br. For cubic CsSnBr₃, the lattice constant is 5.8 Å; therefore, the monolayer (ML) and bilayer (BL) thicknesses are defined as a/2 (2.9 Å) and a (5.8 Å), respectively.



Figure S2. RHEED patterns of substrate and precursors. (A) NaCl along the [110], (B) CsBr and (C) SnBr₂ with 22 Å thickness.



Figure S3. Rotation dependent RHEED patterns. Streaky RHEED patterns of (A-D) cubic phase CsSnBr₃ taken from different rotation angles showing that single-domain epitaxial layer is formed on the substrates. RHEED patterns of roughened (E-H) tetragonal phase CsSn₂Br₅ taken from different rotation angles.



Figure S4. Film adhesion of the epitaxial growth of CsSnBr₃ on NaCl. (A) Epitaxial CsSnBr₃ sample before application of Scotch tape; (B) after attachment on the film surface; and (C) after peeling off.



Figure S5. RHEED patterns of growth of CsSnBr₃ on Ge and InP. (A) Ge single crystalline substrate along the [100]; (B) CsSnBr₃ grown on Ge; (C) Ge single crystalline substrate along the [100] pre-treated with HCl acid etching for 30s; (D) CsSnBr₃ grown on Ge from (C); (E) InP single crystalline substrate along the [100]; (F) CsSnBr₃ grown on InP.



Figure S6. Epitaxy lift-off and regrowth. (A) Epitaxial lift-off procedure: the epitaxial CsSnBr₃ grown on NaCl with a gold layer (~300 Å) on the top is rapidly immersed into liquid nitrogen and then rapidly transferred into diethyl ether; Cu tape is then pressed onto the surface and then slowly peeled which results in separation of the epitaxial film from the substrate; (B-D) RHEED patterns for the epitaxial re-growth of CsSnBr₃ after epitaxial lift-off.



Figure S7. RHEED oscillations. Specular RHEED intensity recorded during $CsSnBr_3$ epitaxial growth at 1:1 stoichiometry on NaCl at (A) 0.28 Å/s and (B) 0.14 Å/s. The oscillation period is 5 s and 10 s for (A) and (B), respectively, and corresponds to

thickness of a half monolayer. (C) Cross-section SEM image used for calibration of the growth rate.



Figure S8. Investigation of interface between epitaxial CsSnBr₃ and NaCl. (A) Crosssection TEM images of NaCl/CsSnBr₃ (~25 nm). The area marked with a white frame in (A) is enlarged and shown in Figure 2a. (B) Same area shown in Figure 2b with black arrows marking dislocations.



Figure S9. Top-view SEM images of epitaxial CsSnBr₃ on NaCl. The scale bar in (A) is 200 nm, in (B) is 1 μ m. The small dots shown on the surface are Pt nanoparticles deposited to avoid charging on the surface.

Identification of CsSn₂Br₅

CsSn₂Br₅ has a bulk tetragonal structure with lattice constants of a = 8.48 Å and c = 15.28 Å.^[1] The *d*-spacings along the substrate normal and along in-plane axes parallel to the NaCl [110] are 7.58±0.12 Å and 3.77±0.05 Å, respectively. The RHEED pattern of the tetragonal phase along NaCl [110] is consistent with the simulated SAED pattern of CsSn₂Br₅ along the [210] direction (shown below), i.e. (001)[110]_{NaCl}//(002)[210]_{CsSn²Br⁵}. Therefore, the measured lattice constants of the tetragonal CsSn₂Br₅ are $a = 8.43\pm0.11$ Å and $c = 15.16\pm0.24$ Å, which is within error of the value reported in Ref^[1].



Figure S10. RHEED pattern, simulated SAED and crystal structure of tetragonal CsSn₂Br₅. (A) RHEED pattern of the sample grown at the ratio of 0.25:1 (CsBr:SnBr₂) collected along the [110] direction of NaCl. (B) Simulated SAED pattern of CsSn₂Br₅ along the [210] direction. The calculated *d*-spacings of (002) and (210) are 7.63 Å and 3.79 Å, respectively, which are consistent with the values calculated from the RHEED pattern (7.58±0.12 Å and 3.77±0.05 Å). Schematic illustration of the crystal structure of CsSn₂Br₅ viewed along (C) *a*-axis and (D) *c*-axis. Green spheres are Cs; gray spheres are Sn; red spheres are Br.

0.5:1 (CsBr:SnBr₂)



Figure S11. *In-situ* real-time monitoring of the phase transition. A phase transition from the cubic to tetragonal phase occurs when the deposition ratio of CsBr to SnBr₂ is 0.5:1 after 1-2 monolayers. Note that while the pattern for the tetragonal phase appears monoclinic, it is actually a rotated tetragonal phase as shown in Fig. S4 and the diffraction spots are therefore not along primary axes.



Figure S12. RHEED oscillations monitored during growth process. (A) RHEED pattern with the monitored intensity area highlighted with the red circle and the corresponding (B) RHEED intensity profile with time.



Figure S13. Crystal structure characterization of the two epitaxial phases. (A) XRD patterns of NaCl (blue curve) and samples grown at different ratios of CsBr:SnBr₂: 0.25:1 (black curve) and 1:1 (red curve).



Figure S14. Calculated XRD patterns. (A) cubic CsSnBr₃ and (B) tetragonal CsSn₂Br₅.



Figure S15. XRD patterns of samples grown at different precursor ratios. (A) 0.5:1 and (B) 1.5:1 of CsBr to SnBr₂. The insets show the appearance of samples prepared at 0.5:1 and 1.5:1 ratio, respectively. It shows both phases (CsSnBr₃ and CsSn₂Br₅) occur when the sample is prepared at 0.5:1 ratio.



Figure S16. XPS spectra of samples grown at different precursor ratios. All the spectra were taken at the top surfaces of the epitaxial film. From the sensitivity factors and the peak area of binding energy of different elements (Cs, Sn, Br), the elemental ratio is obtained.

XPS of CsSn₂Br₅ sputtered by Ar⁺ ion

When prepared with a 0.25:1 ratio, the atomic concentration of Cs is much lower than that of Sn and Br. However, after Ar^+ sputtering of the top surface, the atomic concentration of Cs increases, and the elemental ratio is close to stoichiometric CsSn₂Br₅ (1:2:5 for Cs:Sn:Br) (as shown in fig. S17 and Table S3). This suggests that the Cs vacancies concentrated at the interface are likely eliminated as the growth proceeds or subsequently concentrated as the growth is halted.



Figure S17. XPS spectra of $CsSn_2Br_5$ **after** Ar^+ **ion sputtering**. It should be noted that Sn^{2+} is partially reduced by Ar^+ during sputtering (1.5 mins), resulting in the Sn3d peak splitting; however, this does not change the molar ratios calculated by integrating the peak area of different elements divided with sensitivity factors.



Figure S18. Epitaxial growth of CsSnBr₃ on pseudomorphic interlayers of NaCl-NaBr. (A-E) RHEED patterns of growth of pseudomorphic interlayers of NaCl-NaBr on NaCl substrates. The ratio of NaBr: NaCl is gradually increased to reduce the misfit between substrates and interlayers. Note that when the ratio of NaBr: NaCl reaches 1:1, the lattice constant of interlayers is 5.81 Å perfectly matching that of CsSnBr₃ (5.80 Å). During the growth, the lattice constant calculated from RHEED patterns shows gradual increase along with increasing the ratio of NaBr: NaCl. (F-I) RHEED patterns of growth of CsSnBr₃ on pseudomorphic interlayers of NaCl-NaBr.



Figure S19. XRD patterns of alloyed NaCl-NaBr pseudomorphic interlayers and CsSnBr₃. (A) Comparison of NaCl substrate alone and alloyed NaCl-NaBr layers. (B) Enlarged view of (A) shows a shoulder peak arises in the sample of alloyed NaCl-NaBr pseudomorphic interlayers on NaCl substrates. (C) Enlarged view shows the peak splitting that reflects the change of lattice constant from NaCl substrate, alloyed NaCl-NaBr, and CsSnBr₃.



Figure S20. Absorption spectra of two phases and NaCl substrate. (A) $CsSnBr_3$ of varying well thickness and (B) $CsSn_2Br_5$ (black curve) and NaCl (blue curve). The spectra are converted from (1-Transmission) and shifted for clarity.



Figure S21. Calculated bandgap as a function of lattice parameter. The bandgap of CsSnBr₃ decreases substantially with a decrease of lattice parameter.



Figure S22. DFT calculation using the PBE functional. Band structure, density of states (DOS) and projected density of states (PDOS) of (A) CsSnBr₃ and (B) CsSn₂Br₅.



Figure S23. Solar cell device from epitaxial lift-off. AFM image of (A) single-domain epitaxial lift-off film, and (B) amorphous film. (C) Device architecture of the photovoltaic cells. (D) *I-V* curve of devices fabricated with the single domain epitaxy film and amorphous film, respectively, showing nearly an order of magnitude more photocurrent and double the voltage for single domain films versus the amorphous film.



Figure S24. Multilayer quantum well fabrication. RHEED patterns of (a) freshly cleaved NaCl along [110] direction, (b) NaCl/CsSnBr₃(~10 nm)/NaCl(1.5 nm), (c) NaCl/CsSnBr₃(~10 nm)/NaCl(1.5 nm) /CsSnBr₃(~10 nm)/NaCl(1.5 nm),(d) NaCl/CsSnBr₃(~10 nm)/NaCl(1.5 nm) /CsSnBr₃(~10 nm)/NaCl(1.5 nm) /CsSnBr₃(~10 nm)/NaCl(1.5 nm). It indicates that with well controlled growth, no obvious change occurs even after growing three pairs of CsSnBr₃(~10 nm)/NaCl(1.5 nm). This provides opportunity for fabrication multi-junction quantum wells.



Figure S25. PL spectra of quantum well samples CsSnBr₃/CsSn₂Br₅ with various well widths: 20 nm (black curve), 40 nm (red curve), 80 nm (blue curve), 100 nm (magenta curve). PL of quantum well samples CsSnBr₃/NaCl as a comparison: 80 nm with NaCl (violet curve), and 100 nm with NaCl (olive curve).

Bulk Crystal Structure		CsSnBr ₃	CsSn ₂ Br ₅
		Pm3m, <i>a=b=c=</i> 5.80 Å	I4/mcm, <i>a=b=</i> 8.48 Å, <i>c=</i> 15.28 Å
	0.25:1	-	Observed
	0.5:1	Observed $< 2 \text{ ML}$	Observed $> 3 \text{ ML}$
Precursor Ratio of	1:1	Observed	-
CsBr to SnBr ₂	1.5:1	-	-
	Along NaCl [110]	[110]	[210]
Orientation	Along NaCl [100]	[100]	[310]
	Along NaCl [001]	[001]	[002]
Misfit		-2.8 %	4.9 %

Table S1. Lattice constants, film orientation and misfit of two phases.

Table S2. Elemental ratio of as-deposited films obtained from XPS data.

Dragurgar Datio	Real Ratio		
Precuisor Ratio	Cs	Sn	Br
0.25:1	1.0 ± 0.1	10.0 ± 1.0	13.3 ± 1.3
0.5:1	1.0 ± 0.1	3.0 ± 0.3	5.0 ± 0.5
1:1	1.0 ± 0.1	1.3 ± 0.1	2.9 ± 0.3

Table S3. Elemental ratio of as-deposited films obtained from XPS data collected after 1.5 min Ar^+ ion sputtering.

Precursor Ratio	Real Ratio		
	Cs	Sn	Br
0.25:1	1.0 ± 0.1	4.0 ± 0.4	5.5 ± 0.6

Materials	PBE band gaps (eV)	HSE06 band gaps (eV)	Experimental value (eV)
CsSnBr ₃	0.40	0.84	1.83±0.02
CsSn ₂ Br ₅	2.33	3.12	3.34±0.04

Table S4. Calculated band gaps of CsSnBr₃ and CsSn₂Br₅ using the DFT-PBE and DFT-HSE06 methods.

Table S5. Emission energy of quantum well CsSnBr₃/NaCl with various well width.

Quantum Well Width (nm)	Emission Peak (nm)	Emission Energy (eV)
5	654	1.896
10	664	1.867
20	669	1.854
40	673	1.842
80	684	1.813
100	685	1.810

Table S6. Emission energy of quantum well CsSnBr₃/CsSn₂Br₅ with various well width.

Quantum Well Width (nm)	Emission Peak (nm)	Emission Energy (eV)
20	665	1.865
40	672	1.845
80	673	1.842
100	677	1.832
100 (NaCl)	681	1.821

Movie S1. *In-situ* real-time monitoring of the phase transition from cubic to tetragonal phase.

The phase transition from cubic to tetragonal phase arises after deposition is proceeding for around 20 s, which is interpreted as the film thickness of around 1-2 monolayers.

Supplementary Text

Computational Details

All the DFT calculations are performed using the Vienna Ab initio Simulation package (VASP) with the implemented projector augmented wave (PAW) potentials. All electronic self-consistent energy calculations converge within the accuracy of 10^{-5} eV between each electronic iteration. Gaussian smearing with the smearing width of 0.05 eV is used to treat bands with partial occupancies. For the cubic unit cell, an energy cutoff of the plane wave basis set (E_{CUT}) equal to 300eV and a Γ -centered 8×8×8 k-point mesh are used for both PBE and HSE06 calculation. For the tetragonal unit cell, using a gammacentered $4 \times 4 \times 2$ k-point mesh (i.e. k-point spacing of 0.03 Å⁻¹) was found to be suitable. For HSE06 calculations, we use $E_{CUT} = 250 \text{eV}$, which yields an error of about 10 meV but reduces the computational cost. For PBE and HSE06 band structure calculations, paths connecting high-symmetry k-points are divided into 10-15 k-points. The high symmetry points in the plots are defined in Ref.[2]. For PBE DOS calculations, a finer grid of k-point that has the separation between k-points of around 0.01 Å⁻¹ is used, i.e. Γ centered 16×16×16 and 11×11×6 grids are used for the CsSn₂Br₅ and CsSnBr₃ DOS calculations, respectively. For HSE06 DOS calculations, we use the gamma-centered $8 \times 8 \times 8$ and $4 \times 4 \times 2$ k-point meshes for CsSnBr₃ and CsSn₂Br₅, respectively. These grids result in a k-point spacing of around 0.02 - 0.03 Å⁻¹, which is sufficient for DOS calculations.

Fitting the Quantum Well data

The emission energy of a quantum well is described by the Brus equation as:^[3]

$$E_{g}^{well}(L_{z}) = E_{g}^{0} + \Delta E(L_{z}) = E_{g}^{0} + \frac{\hbar^{2}\pi^{2}}{2m^{*}L_{z}} - \frac{1.8q^{2}}{4\pi\varepsilon_{r}\varepsilon_{o}L_{z}}$$
(1)

where $E_g^{\ 0}$ is the bulk band gap, ΔE is the confinement energy of both electrons and holes, \hbar is reduced Planck constant, L_z is the thickness of the quantum well, and m^* is the reduced mass that can be obtained from the effective masses of the electron (m_e^*) and the hole (m_h^*) as $\frac{1}{m^*} = \frac{1}{m_e^*} + \frac{1}{m_h^*}$, q is the charge of electron, ε_r is the relative permittivity, and ε_0 is the vacuum permittivity. Since the exciton binding energy of CsSnBr₃ has been reported to be less than 1 meV^[4] the size dependence of the quantum well bandgap can be expressed as:^[5]

$$E_{total} = E_g^0 + \Delta E = E_g^0 + \frac{\hbar^2 \pi^2}{2m^* L_z^2}$$
(2)

The value of m^* obtained by fitting the PL data of quantum well PL data can be extracted and then used for calculation of Bohr radius of this material by using Equation (3):^[6]

$$a_{B} = \frac{4\pi\hbar^{2} \varepsilon_{o}}{e^{2}} \left(\frac{1}{m_{e}^{*}} + \frac{1}{m_{h}^{*}}\right)$$
(3)

where a_B is the Bohr radius, e is the electron charge, ε_o is the vacuum permittivity, and ε is the dielectric constant of the semiconductor which has been reported for CsSnBr₃ to be 32.4.^[4]

References:

[1] I. Abrahams, D. Z. Demetriou, R. T. Kroemer, H. Taylor, M. Motevalli, *Journal of Solid State Chemistry* 2001, **160**, 382.

[2] W. Setyawan, S. Curtarolo, *Computational Materials Science* 2010, **49**, 299.

[3] L. E. Brus, The Journal of Chemical Physics 1983, 79, 5566; L. Brus, *The Journal of Physical Chemistry* 1986, **90**, 2555.

[4] L.-Y. Huang, W. R. Lambrecht, *Physical Review B* 2013, **88**, 165203.

[5] D. A. Miller, *Quantum Dynamics of Simple Systems*, ed. G.-L. Oppo, SM Barnett, E. Riis, and M. Wilkinson (Institute of Physics, London, 1996) 1996, **239**.

[6] R. Koole, E. Groeneveld, D. Vanmaekelbergh, A. Meijerink, C. de Mello Donegá, in *Nanoparticles: Workhorses of Nanoscience*, (Ed: C. de Mello Donegá), Springer Berlin Heidelberg, Berlin, Heidelberg 2014, **13**.