

The DOE E3SM coupled model version 1: Overview and evaluation at standard resolution

Jean-Christophe Golaz¹, Peter M. Caldwell¹, Luke P. Van Roekel², Mark R. Petersen², Qi Tang¹, Jonathan D. Wolfe², Guta Abeshu³, Valentine Anantharaj⁴, Xylar S. Asay-Davis², David C. Bader¹, Sterling A. Baldwin¹, Gautam Bisht⁵, Peter A. Bogenschutz¹, Marcia Branstetter⁴, Michael A. Brunke⁶, Steven R. Brus², Susannah M. Burrows⁷, Philip J. Cameron-Smith¹, Aaron S. Donahue¹, Michael Deakin^{8,9}, Richard C. Easter⁷, Katherine J. Evans⁴, Yan Feng¹⁰, Mark Flanner¹¹, James G. Foucar⁸, Jeremy G. Fyke^{2,12}, Brian M. Griffin¹³, Cécile Hannay¹⁴, Bryce E. Harrop⁷, Elizabeth C. Hunke², Robert L. Jacob¹⁰, Douglas W. Jacobsen², Nicole Jeffery², Philip W. Jones², Noel D. Keen⁵, Stephen A. Klein¹, Vincent E. Larson¹³, L. Ruby Leung⁷, Hong-Yi Li³, Wuyin Lin¹⁵, William H. Lipscomb¹⁴, Po-Lun Ma⁷, Salil Mahajan⁴, Mathew E. Maltrud², Azamat Mametjanov¹⁰, Julie L. McClean¹⁶, Renata B. McCoy¹, Richard B. Neale¹⁴, Stephen F. Price², Yun Qian⁷, Philip J. Rasch⁷, J. E. Jack Reeves Eyre⁶, William J. Riley⁵, Todd D. Ringler¹⁷, Andrew F. Roberts², Erika L. Roesler⁸, Andrew G. Salinger⁸, Zeshawn Shaheen¹, Xiaoying Shi⁴, Balwinder Singh⁷, Jinyun Tang⁵, Mark A. Taylor⁸, Peter E. Thornton⁴, Adrian K. Turner², Milena Veneziani², Hui Wan⁷, Hailong Wang⁷, Shanlin Wang², Dean N. Williams¹, Philip J. Wolfram², Patrick H. Worley¹⁸, Shaocheng Xie¹, Yang Yang⁷, Jin-Ho Yoon¹⁹, Mark D. Zelinka¹, Charles S. Zender²⁰, Xubin Zeng⁶, Chengzhu Zhang¹, Kai Zhang⁷, Yuying Zhang¹, Xue Zheng¹, Tian Zhou⁷, Qing Zhu⁵

¹Lawrence Livermore National Laboratory, Livermore, CA

²Los Alamos National Laboratory, Los Alamos, NM

³University of Houston, Houston, TX

⁴Oak Ridge National Laboratory, Oak Ridge, TN

⁵Lawrence Berkeley National Laboratory, Berkeley, CA

⁶University of Arizona, Tucson, AZ

⁷Pacific Northwest National Laboratory, Richland, WA

⁸Sandia National Laboratory, Albuquerque, NM

⁹University of British Columbia, Vancouver, BC Canada

¹⁰Applied Numerical Harmonic Analysis

¹¹University of Michigan, Ann Arbor, MI

¹²Howard Engineering Group of Companies, Edison, NJ, Canada

¹³University of Wisconsin – Milwaukee, Milwaukee, WI

¹⁴National Center for Atmospheric Research, Boulder, CO

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¹⁵Brookhaven National Laboratory, Upton, NY

¹⁶Scripps Institution of Oceanography, University of California, San Diego, La Jolla, CA

¹⁷Congressional Science Fellow, U.S. House of Representatives, Washington, DC

¹⁸PHWorley Consulting, Oak Ridge, TN

¹⁹Gwangju Institute of Science and Technology, Gwangju, Korea

²⁰University of California, Irvine, CA

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Key Points:

- Description of E3SMv1, the first version of the U.S. DOE Energy Exascale Earth Sytem Model.
- The performance of E3SMv1 is documented with a set of standard CMIP6 DECK and historical simulations comprising nearly 3000 years.
- E3SMv1 has a high equilibrium climate sensitivity (5.3 K) and strong aerosol-related effective radiative forcing (-1.65 W m^{-2}).

Corresponding author: Chris Golaz, golaz1@llnl.gov

Abstract

This work documents the first version of the U.S. Department of Energy (DOE) new Energy Exascale Earth System Model (E3SMv1). We focus on the standard resolution of the fully-coupled physical model designed to address DOE mission-relevant water cycle questions. Its components include atmosphere and land (110km grid spacing), ocean and sea ice (60km in the mid-latitudes and 30km at the equator and poles), and river transport (55km) models. This base configuration will also serve as a foundation for additional configurations exploring higher horizontal resolution as well as augmented capabilities in the form of biogeochemistry and cryosphere configurations.

The performance of E3SMv1 is evaluated by means of a standard set of Coupled Model Intercomparison Project Phase 6 Diagnosis, Evaluation, and Characterization of Klima (CMIP6 DECK) simulations consisting of a long pre-industrial control, historical simulations (ensembles of fully coupled and prescribed SSTs) as well as idealized CO₂ forcing simulations. The model performs well overall with biases typical of other CMIP-class models, although the simulated Atlantic Meridional Overturning Circulation is weaker than many CMIP-class models. While the E3SMv1 historical ensemble captures the bulk of the observed warming between pre-industrial (1850) and present-day, the trajectory of the warming diverges from observations in the second half of the 20th century with a period of delayed warming followed by an excessive warming trend. Using a two-layer energy balance model, we attribute this divergence to the model's strong aerosol-related effective radiative forcing ($ERF_{\text{ari+aci}} = -1.65 \text{ W m}^{-2}$) and high equilibrium climate sensitivity ($ECS = 5.3 \text{ K}$).

Plain Language Summary

The United States Department of Energy funded the development of a new state-of-the-art Earth system model for research and applications relevant to its mission. The Energy Exascale Earth System Model version 1 (E3SMv1) consists of five interacting components for the global atmosphere, land surface, ocean, sea ice and rivers. Three of these components (ocean, sea ice and river) are new and have not been coupled into an Earth system model previously. The atmosphere and land surface components were created by extending existing components part of the Community Earth System Model, Version 1.

78 E3SMv1's capabilities are demonstrated by performing a set of standardized simu-
79 lation experiments described by the Coupled Model Intercomparison Project Phase 6 Di-
80 agnosis, Evaluation, and Characterization of Klima (CMIP6 DECK) protocol at standard
81 horizontal spatial resolution of approximately 1 degree latitude and longitude. The model
82 reproduces global and regional climate features well compared to observations. Simu-
83 lated warming between 1850 and 2015 matches observations, but the model is too cold
84 by about 0.5°C between 1960 and 1990 and later warms at a rate greater than observed. A
85 thermodynamic analysis of the model's response to greenhouse gas and aerosol radiative
86 affects may explain the reasons for the discrepancy.

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1 Introduction

In 2013, the US Department of Energy (DOE) developed a report summarizing observed long-term trends that, if continued for several decades, would have major impacts on the energy sector [U.S. Department of Energy, 2013]. Among these were regional trends in air and water temperatures, water availability, storms and heavy precipitation, coastal flooding and sea-level rise. The ability to simulate and predict significant, long-term changes in these environmental variables important to energy-sector decisions required capabilities beyond the existing state-of-the-science Earth system models. The Energy Exascale Earth System Model (E3SM) project was conceived to meet this mission need [Bader *et al.*, 2014].

Scientific developments in E3SM are dictated by three science drivers that broadly cover the foundational science for advancing Earth system prediction. Notably, water cycle, biogeochemistry, and cryosphere systems govern variability and changes in water availability and storms, air and river stream temperature, and coastal flooding and sea level rise that are all critical to the energy sector [Leung *et al.*, 2019].

E3SM version 1 (E3SMv1) was branched from the Community Earth System Model (CESM1; Hurrell *et al.* [2013]) but has evolved significantly since. E3SMv1 consists of three coupled modeling systems with varying degrees of sophistication. The present work describes the physical Earth system model that represents water and energy cycles in atmosphere, ocean, sea ice, land and river components. This configuration is aimed at addressing the DOE water cycle science questions relating to interactions between the water cycle and the rest of the human-Earth system on local to global scales, water availability, and water cycle extremes. This physical model also serves as a foundation for two additional configurations: (i) a biogeochemistry configuration with interactive nitrogen and phosphorous for interactions between biogeochemical cycles and other Earth system components and (ii) a cryosphere configuration with added interactive ice-shelf cavities for assessing the impacts of ocean-ice shelf interactions on Antarctic Ice Sheet dynamics and the implications for sea level rise.

The focus here is on the physical model at standard resolution useful for simulations like those specified in Coupled Model Intercomparison Project Phase 6 (CMIP6). This includes a 1° atmosphere and land (equivalent to 110 km at the equator), 0.5° river model (55 km), and an ocean and sea ice with mesh spacing varying between 60 km in

119 the mid-latitudes and 30 km at the equator and poles. A higher resolution configuration
120 with a 0.25° atmosphere and land, 0.125° river model, and ocean and sea ice with mesh
121 spacing between 18 km at the equator and 6 km at the poles (roughly equivalent to a 0.1°
122 resolution) will be documented in a subsequent paper.

123 Given that E3SMv1 features three specific models of components (ocean, sea ice,
124 and river) that have never been used in a coupled Earth System Model and that there were
125 significant developments in the atmosphere and land components, an examination of the
126 model behavior relative to observations and other CMIP class models is needed. We ana-
127 lyze CMIP6 Diagnosis, Evaluation, and Characterization of Klima (DECK) and historical
128 simulations [Eyring *et al.*, 2016] performed with E3SMv1. This allows for a rigorous com-
129 parison of E3SMv1 behavior against observations and many other models. This work will
130 also provide a baseline for all future E3SM developments and experiments.

131 This paper is organized as follows. In Section 2, we present an overview of the
132 model components, with a specific focus on new model components (i.e., those not previ-
133 ously used in coupled Earth system modeling), as well as new developments in the atmo-
134 sphere and land models. In Section 3, we describe the E3SMv1 initialization and spin-up
135 procedure, including model tuning objectives and simulation campaign. Section 4 pro-
136 vides an overview of the pre-industrial control simulation and Section 5 an analysis of the
137 E3SMv1 climate in the historical simulations including short and long-term variability. An
138 extended discussion of E3SMv1's Effective Radiative Forcing (ERF) and climate sensitiv-
139 ity is provided in Section 6. Finally, we offer a summary of E3SMv1 fidelity and discuss
140 future directions for the fully coupled model in Section 7.

141 2 Model Overview

142 E3SM started with a version of the CESM1 [Hurrell *et al.*, 2013, [http://www.
143 cesm.ucar.edu/models/cesm1.0](http://www.cesm.ucar.edu/models/cesm1.0)] from which we developed the fully coupled E3SMv1
144 system. Notable changes between E3SMv1 and CESM1 include:

- 145 • E3SM Atmosphere Model (EAM) component with a spectral-element (SE) dynam-
146 ical core, increased vertical resolution, and substantially revamped physics and the
147 capability of regional grid refinement for multi-resolution simulations.

- 148 • New ocean and sea-ice components based on the Model for Prediction Across Scales
149 (MPAS) framework that uses Spherical Centroidal Voronoi Tessellations (SCVTs)
150 for multi-resolution modeling.
- 151 • The river transport model of CESM1 was replaced by a new river model, Model
152 for Scale Adaptive River Transport (MOSART), for a more physically based repre-
153 sentation of riverine processes.
- 154 • E3SMv1 land model (ELM) is based on the Community Land Model Version 4.5
155 (CLM4.5) with new options for representing soil hydrology and biogeochemistry
156 added to enable analysis of structural uncertainty, with important implications to
157 carbon cycle and climate feedbacks for addressing v1 biogeochemistry questions.

158 The sub-sections below provide a more detailed description of the model compo-
159 nents. Coupling of the new MPAS-Ocean, MPAS-Seaice and MOSART models with EAM
160 and ELM provides E3SMv1 with a unique capability for multi-resolution modeling using
161 unstructured grids in most of its component models. This capability is critical for future
162 simulation campaigns that have a strong regional focus to meet DOE's needs for Earth
163 system modeling in support of energy-sector decisions.

164 The project also built a comprehensive infrastructure for code management, devel-
165 opment, testing, and analysis to enable development of E3SMv1 and future versions at
166 DOE leadership computing centers. Leveraging DOE investments, a flexible framework
167 provides workflow orchestration, provenance capture and management, simulation analysis
168 and visualization, and automated testing and evaluation capabilities (see Appendix C: for
169 a description of some of the analysis tools).

170 **2.1 Atmosphere**

171 The E3SM Atmosphere Model (EAM) is the atmosphere component of E3SMv1. It
172 is a descendant of the Community Atmosphere Model version 5.3 (CAM5.3). EAM uses a
173 spectral element dynamical core at 100 km resolution on a cubed-sphere geometry. Verti-
174 cal resolution was increased from 30 layers, with a top at approximately 40 km, in CAM5
175 to 72 layers with a top at approximately 60 km in EAM. EAM contains many innovations
176 compared to CAM5. While changes in the EAM physics are broadly similar to changes
177 from CAM5.3 to CAM6, many details of tuning, and cloud and aerosols formulations dif-
178 fer in important ways. EAM is described in *Rasch et al.* [2019]. A detailed analysis of its

cloud and convective characteristics and the rationale for model tuning are provided in *Xie et al.* [2018]. Sensitivity of EAM to a number of its adjustable parameters is explored in *Qian et al.* [2018].

Key features of EAMv1 include:

- Implementation of a simplified third-order turbulence closure parameterization (CLUBB; Cloud Layers Unified By Binormals) [*Golaz et al.*, 2002; *Larson and Golaz*, 2005; *Larson*, 2017] that unifies the treatment of planetary boundary layer turbulence, shallow convection, and cloud macrophysics.
- An updated microphysical scheme, the version 2 of *Morrison and Gettelman* [2008] [MG2; *Gettelman et al.*, 2015]. The combination of CLUBB and MG2 enables aerosol-cloud interactions in large-scale and shallow convective clouds to be considered. Significant additional changes to nucleation and ice microphysics have also been incorporated compared to MG2 in CAM6.
- The deep convection scheme of *Zhang and McFarlane* [1995].
- The Modal Aerosol Module (MAM4) with revisions to improve aerosol resuspension, convective transport, aerosol nucleation, scavenging, as well as modifications to sea spray emissions so marine ecosystems can contribute organic matter to aerosols.
- A linearized ozone photochemistry to predict stratospheric ozone changes, which provides an important source of stratospheric variability (Linoz v2, *Hsu and Prather* [2009]).

The computational cost of EAMv1 increased by approximately a factor of four relative to CAM5 due to higher vertical resolution and parameterization complexity, along with a larger number of predicted and transported variables (aerosol species and prognostic snow and rain).

2.2 Ocean

MPAS-Ocean, based on the Model for Prediction Across Scales (MPAS) framework [*Ringler et al.*, 2010], is the ocean component of E3SMv1. MPAS-Ocean uses a mimetic finite volume discretization of the primitive equations and invokes the hydrostatic, incompressible, and Boussinesq approximations on a staggered C-grid [*Arakawa and Lamb*,

1977; Thuburn *et al.*, 2009; Ringler *et al.*, 2013]. MPAS-Ocean grid cells for E3SMv1 simulations are near-hexagons (five or more sides), but the MPAS framework supports cells with any number of sides; the algorithms and code are identical for all cell shapes. The tracer advection scheme is the quasi 3rd-order flux corrected transport (FCT) scheme [Skamarock and Gassmann, 2011] with separate limiting in the horizontal and vertical. The MPAS-Ocean time stepping method is split-explicit, where the barotropic component is subcycled within each baroclinic time step.

The simulations presented here use a z-star vertical coordinate within an arbitrary Lagrangian-Eulerian scheme, where the layer thicknesses of the full column expand and contract with the sea surface height [Petersen *et al.*, 2015; Reckinger *et al.*, 2015]. The prognostic volume-based equation of motion includes surface mass fluxes from the coupler, thus virtual salt fluxes are not needed. Vertical mixing is computed implicitly at the end of each time step using the Community Vertical Mixing project implementation of the K-profile parameterization (KPP) as described by Van Roekel *et al.* [2018] where our configuration of KPP is based on the results of comparison against large eddy simulations.

E3SMv1 standard resolution simulations employ the classic Gent and McWilliams [1990] eddy transport (GM) parameterization. The GM bolus coefficient was tuned, in part, to reduce the transport of heat to depth in the Southern Ocean, to a value of $1800 \text{ m}^2 \text{ s}^{-1}$ for the simulations presented here. The Redi coefficient, which adds diffusion along isopycnal layers, was set to zero for this set of simulations. In the horizontal, bi-harmonic viscosity is used for momentum ($1.2 \times 10^9 \text{ m}^4 \text{ s}^{-2}$). No explicit horizontal tracer diffusivity is included.

2.3 Sea ice

MPAS-Seaice is the sea ice component of E3SMv1. MPAS-Seaice and MPAS-Ocean share identical meshes, but MPAS-Seaice uses a B-grid [Arakawa and Lamb, 1977] with sea-ice concentration, volume, and tracers defined at cell centers and velocity defined at cell vertices. Velocity components at cell vertices are not aligned with the mesh, as in sea-ice models with structured meshes and quadrilateral cells. Instead, the velocity components are aligned with a spherical coordinate system that is locally Cartesian and has its poles on the geographical equator. Velocities are determined by solving the sea-ice momentum equation [Hibler III, 1979; Hunke and Dukowicz, 1997] on cell vertices

240 in an identical manner to the CICE model [Hunke *et al.*, 2015] except for the internal
241 stress term. MPAS-Seaice uses the variational formulation of the internal stress term used
242 by CICE [Hunke and Dukowicz, 2002], but modified to use the non-regular polygons of
243 MPAS meshes. Instead of the bilinear basis functions used by CICE, MPAS-Seaice uses
244 Wachspress basis functions [Dasgupta, 2003], which are integrated with the quadrature
245 rules of Dunavant [1985]. Horizontal transport of ice concentration, volume, and tracers is
246 achieved with an incremental remapping scheme similar to that described in Dukowicz and
247 Baumgardner [2000], Lipscomb and Hunke [2004] and Lipscomb and Ringler [2005] but
248 adapted to MPAS meshes. MPAS-Seaice shares the same column physics code as CICE
249 through the Icepack library [Hunke *et al.*, 2018]. For simulations shown here, MPAS-
250 Seaice uses the “mushy layer” vertical thermodynamics scheme of Turner *et al.* [2013];
251 Turner and Hunke [2015], the level-ice melt pond scheme of Hunke *et al.* [2013], a delta-
252 Eddington shortwave radiation scheme [Briegleb and Light, 2007; Holland *et al.*, 2012], a
253 scheme for transport in thickness space [Lipscomb, 2001], and a representation of mechan-
254 ical redistribution [Lipscomb *et al.*, 2007].

255 Coupling of the sea-ice component to the ocean takes advantage of z-star ocean co-
256 ordinates as described by Campin *et al.* [2008], and is a departure from the coupling of
257 CICE and POP in CESM1. The weight of sea ice contributes to the ocean’s barotropic
258 mode, notably affecting the free surface over continental shelves. In shallow water depths
259 at or less than the floating ice draft, the weight passed to the ocean model is limited to
260 prevent evacuation of the underlying liquid column. When frazil ice forms in the ocean
261 model, the volume of newly formed crystals is passed to the sea ice model with a fixed
262 salinity of 4 PSU, rather than exchanging a freezing potential as in other models. Future
263 versions of E3SM will permit progressive brine drainage to the ocean from the mushy-
264 layer physics used in MPAS-Seaice [Turner and Hunke, 2015]. For E3SMv1, brine drainage
265 occurs internally in MPAS-Seaice for thermodynamic calculations, but for the sake of
266 freshwater coupling, the ocean model only receives mass fluxes back from melted sea ice
267 at the fixed salinity that it originally passed to its cryospheric counterpart (4 PSU). The
268 ocean temperature immediately under the ice is the same as the liquid phase in the lowest
269 layer of the sea ice model, and is not fixed at -1.8°C as is typical of previous generation
270 coupled models [Naughten *et al.*, 2017]. For the current version, we have addressed these
271 long-standing ocean-ice coupling issues identified by the modeling community: explicit
272 sea ice mass and salt exchange, a pressure force of the ice on the ocean, a basal sea ice

273 temperature consistent with the ocean model's equation of state, and resolved inertial os-
274 cillations [Schmidt *et al.*, 2004; Hibler *et al.*, 2006; Lique *et al.*, 2016].

275 **2.4 Land**

276 The E3SM Land Model, version 0 (ELMv0) is the land component in E3SMv1.
277 ELMv0 adopts many of its capabilities from its parent model, the Community Land Model
278 version 4.5 (CLM4.5). Oleson *et al.* [2013] provide a detailed technical description of that
279 parent model, including descriptions of all fundamental equations representing the mod-
280 el's biogeophysical and biogeochemical dynamics. The model describes interactions of
281 the land surface with the near-surface atmosphere, and includes interactions among land
282 sub-systems such as vegetation, soil, snow, groundwater, runoff, urban areas, and managed
283 ecosystems. These interactions are represented through conservation equations of state for
284 energy, water, carbon, and nitrogen, including terms that couple these states and fluxes.
285 While ELMv0 includes prognostic carbon and nitrogen biogeochemistry and dynamic
286 ecosystem structure, the E3SMv1 water cycle experiments do not make use of these ca-
287 pabilities. Instead, these simulations use a static representation of land ecosystem structure
288 with prescribed seasonal changes in vegetation canopies based on a climatology of satel-
289 lite remote sensing data. This simulation mode is referred to as satellite phenology (SP),
290 and is based on methods described by Lawrence and Chase [2007]. In this simulation
291 mode, regardless of long-term climate changes or short-term anomalies, the global dis-
292 tribution and seasonal variation of vegetation structure is constant over the course of the
293 simulations. This includes model representation of vegetation height, timing and amount
294 of displayed leaf area, and vegetation contributions to land surface albedo.

295 Several new developments have been made within ELMv0 since its branch point
296 from CLM4.5. First, an extended representation of the influence of aerosols and black car-
297 bon deposition on snow was introduced, based on data summarized by Liu *et al.* [2012],
298 causing aerosol optical properties within the snowpack to vary as a function of snow grain
299 size and aerosol/ice mixing state [Flanner *et al.*, 2012]. This development also fixed a bug
300 related to calculation of snow grain size following snow layer division. Second, a minor
301 modification was made to reduce the rate of evaporation from area characterized as pervi-
302 ous road under dry conditions. Third, the numerical scheme for calculation of leaf stom-
303 atal conductance was updated to prevent non-physical simulation of negative internal leaf

304 CO₂ concentrations. Fourth, albedo calculation was modified to return a land albedo of
305 1.0 in land cells and time steps when the sun is below the horizon.

306 **2.5 River**

307 E3SMv1 replaced the River Transport Model (RTM) [Branstetter and Erickson III,
308 2003], a linear reservoir routing model used in CESM1 and E3SMv0, with the Model
309 for Scale Adaptive River Transport (MOSART), which uses a physically-based kinematic
310 wave approach for river routing across local, regional and global scales [Li *et al.*, 2013,
311 2015]. In the standard E3SMv1 resolution, MOSART uses a regular latitude-longitude
312 grid with spacing of 0.5°. Surface and subsurface runoff simulated by ELM are mapped
313 from the ELM grid to the 0.5° latitude-longitude grid as input to MOSART, which routes
314 the runoff and provides freshwater input to the ocean model. MOSART does not exchange
315 water with the atmosphere or return water to the land model. MOSART divides each grid
316 cell into three categories of hydrologic units: hillslopes that contribute both surface and
317 subsurface runoff into tributaries, tributaries that discharge into a single main channel,
318 and the main channel that connects the local grid cell with the upstream/downstream grid
319 cells through the river network. Two simplified forms of the one-dimensional Saint-Venant
320 equations are used to represent water flow over hillslopes, in the tributary, or in the main
321 channels. MOSART only routes positive runoff, although spurious negative runoff can be
322 generated occasionally by the land model. Negative runoff is mapped directly from the
323 grid cell where it is generated at any time step to the basin outlet of the corresponding
324 MOSART grid cell. More detailed descriptions of MOSART and its input hydrography
325 data and channel geometry parameters can be found in Li *et al.* [2013]. When driven by
326 runoff simulated by CLM4 with observed meteorological forcing data, streamflow simu-
327 lated by MOSART is shown to reproduce the observed annual, seasonal, and daily flow
328 statistics at over 1600 stations of the world's major rivers reasonably well in terms of the
329 overall bias and the seasonal variation [Li *et al.*, 2015].

330 **2.6 Coupling and performance**

331 E3SMv1 uses component coupling software from the Common Infrastructure for
332 Modeling the Earth (CIME). The top level driver is cpl7 [Craig *et al.*, 2012], which pro-
333 vides a main program for forming the single executable of E3SM, directs the time inte-
334 gration of the coupled model, and performs any necessary interpolation or time averag-

335 ing needed between the components. Communication between the parallel components of
336 E3SM and parallel interpolation is provided by the Model Coupling Toolkit [*Jacob et al.*,
337 2005; *Larson et al.*, 2005] which is included in CIME.

338 cpl7 implements an online-offline method for interpolating values between different
339 grids in E3SMv1. Grid intersections and interpolation weights are calculated with an of-
340 fline tool and then read in at runtime and applied during the coupled integration. cpl7 also
341 allows flux and state variables to be interpolated with different weights. The number of
342 different weight files needed is relatively small because the atmosphere and land models
343 share a grid, as do the ocean and sea-ice models. The interpolation weights for nearly all
344 variables are calculated with the ESMF_RegridWeightGen program from ESMF [*Collins*
345 *et al.*, 2005] using the first-order conservative option. However interpolation weights for
346 state variables from the atmosphere to the ocean/sea ice grids are calculated using Tem-
347 pestRemap [*Ullrich and Taylor*, 2015]. Unlike ESMF, the TempestRemap algorithms have
348 native support for vertex centered finite element grids, such as used by the E3SM atmo-
349 sphere dycore. Future versions of E3SM will use TempestRemap for all mapping weight
350 calculations.

351 cpl7 allows several time sequencing options for the components and flexibility for
352 how often the components communicate with the coupler. E3SMv1 uses "RASM_OPTION1"
353 where the sea ice, land and river runoff models execute simultaneously and in sequence
354 with the atmosphere. The ocean model runs simultaneously with all four of those compo-
355 nents. Model timesteps are as follows:

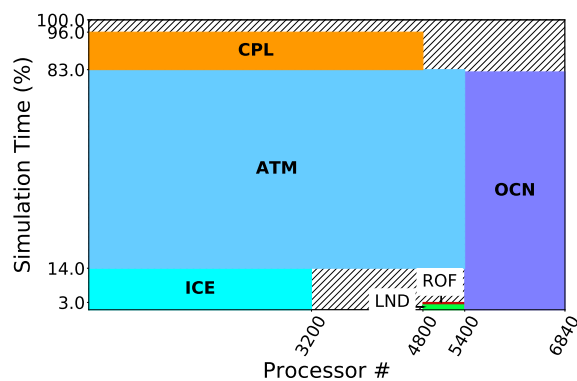
- 356 • The main atmosphere physics timestep is 30 min, but a few parameterizations use
357 a different timestep. CLUBB and MG2 microphysics are substepped together at
358 a timestep of 5 min. Radiation is updated hourly. Several layers of substepping
359 are used by the atmosphere dynamics and tracer transport: the Lagrangian verti-
360 cal discretization uses 15 min timesteps, the horizontal discretization uses 5 min
361 timesteps, and the explicit numerical diffusion uses 100 sec timesteps.
- 362 • Ocean model timestep is 10 min with a barotropic sub-timestep of 40 sec.
- 363 • Sea-ice model timestep is 30 min. The 30 min coupling timestep between the sea-
364 ice and ocean permits transient inertial oscillations in the drift of ice. This can
365 cause instabilities arising from frequent exchange of sea ice weight and sea surface

366 height between sea-ice and ocean. Thus, we damp the sea-surface-height gradient
 367 in the sea-ice momentum equation with a 24-hour Newtonian relaxation constant.

- 368 • Land model timestep is 30 min.
- 369 • River runoff model timestep is 1 hour.

370 The coupling frequency for all components is 30 minutes except the river runoff model
 371 which communicates every 3 hours.

372 Simulations in this work were performed on DOE’s National Energy Research Sci-
 373 entific Computing Center (NERSC) Edison supercomputer ([http://www.nersc.gov/
 374 users/computational-systems/edison](http://www.nersc.gov/users/computational-systems/edison)). Using 285 nodes (each consisting of two 12-
 375 core Intel “Ivy Bridge” processors at 2.4 GHz) with 24 MPI tasks per node, the coupled
 376 model performance averages 10 simulated years per day. Figure 1 illustrates the sequenc-
 377 ing, processor layout and relative cost for all components in E3SMv1.



378 **Figure 1.** Typical processor layout, performance and sequencing for all components in E3SMv1. Compo-
 379 nents include coupler (CPL), atmosphere (ATM), ocean (OCN), sea ice (ICE), land (LND) and river runoff
 380 (ROF). The hatched area represents unoccupied time. OCN runs concurrently on its own nodes and is load-
 381 balanced with ATM+ICE, which run sequentially. Driver overhead is 4%. Other layouts are possible, but we
 382 have found this to be the most robust and efficient.

383 3 Tuning and initialization

384 3.1 Final coupled model tuning

386 Tuning is an integral part of climate model development [e.g., *Hourdin et al.*, 2017;
 387 *Schmidt et al.*, 2017]. While most of the tuning for E3SMv1 was performed at the model

385 **Table 1.** Atmospheric physics parameters that were altered during the development of the coupled model.

Parameter	Initial	Final	Impact	Motivation
zmconv_ke	1.5e-06	5.0e-06	Deep convection rain evaporation	MG2 bug fix 1
so4_sz_thresh_icenuc	7.5e-08	5.0e-08	SO ₄ ice nucleation threshold size	MG2 bug fix 2
clubb_c14	1.30	1.06	CLUBB dissipation for $\overline{u'^2}$ and $\overline{v'^2}$	Final coupled tuning

388 component level, some additional retuning was subsequently required for two reasons: (1)
 389 code errors were discovered in the atmosphere model physics while the coupled system
 390 was being developed and (2) the coupled tuning objectives required an additional level of
 391 adjustment to accommodate biases present in other model components and interactions
 392 between components. Two minor code errors were discovered and addressed in EAMv1:

- 393 1. An incorrectly positioned parenthesis in the code designed to prevent over-depletion
 394 of rain number when multiple microphysics processes act to concurrently deplete
 395 rain drops (<https://github.com/E3SM-Project/E3SM/pull/1599>) led to a
 396 reduction in the rain drop number concentration process rate. Correcting the error
 397 produced a degradation in overall precipitation pattern, approximately compensated
 398 for through increasing the efficiency of convective precipitation evaporation with
 399 zm_conv_ke (see Table 1)
- 400 2. Snow crystal number concentration was updated improperly when snow sublima-
 401 tion occurred (<https://github.com/E3SM-Project/E3SM/pull/1765>). Fixing
 402 the error decreased the ice water path, weakening both LW and SW cloud radiative
 403 effects, which was then countered by a reduction in the sulfate threshold ice nucle-
 404 ation size so4_sz_thresh_icenuc to increase ice number concentration (Table
 405 1).

406 A final tuning of the coupled model with pre-industrial (perpetual 1850) forcings was per-
 407 formed to achieve:

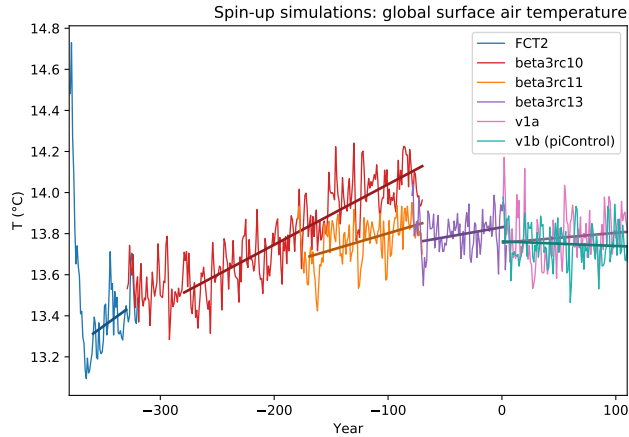
- 408 1. A near-zero long-term average net top-of-atmosphere (TOA) energy flux.
- 409 2. Minimum long-term drift in global mean surface air temperature.
- 410 3. Reasonable absolute global mean surface air temperature.

411 This final tuning required adjustments of less than one W m^{-2} because the E3SM atmo-
412 sphere component was already well tuned from simulations with prescribed SSTs and sea-
413 ice concentrations. To simplify the process, we chose to focus on a single parameter in the
414 CLUBB atmosphere turbulence parameterization, `clubb_c14`, which directly impacts dis-
415 sipation of the horizontal components of the turbulent kinetic energy (TKE). Because of
416 the tight coupling between TKE and boundary layer clouds, `clubb_c14` modulates low-
417 clouds, thus affecting net TOA mostly through SW cloud radiative effects.

418 While most of the coupled tuning was performed under pre-industrial conditions,
419 we also performed a few additional simulations to monitor climate sensitivity and total
420 effective radiative forcing during the developmental phase of E3SMv1. We evaluated sen-
421 sitivity using idealized +4 K simulations [Cess *et al.*, 1989] and one abrupt quadrupling of
422 CO_2 . Total effective radiative forcing was estimated using pairs of atmosphere simulations
423 with identical SSTs but differing forcing (climatological 2000 vs 1850). One historical test
424 simulation was also performed with a near final version of E3SMv1 (“beta3rc10” below).
425 The outcomes of these intermediate simulations were consistent with the final model. A
426 pragmatic, but deliberate, decision was made not to attempt to reduce the high climate
427 sensitivity or reduce the aerosol forcing. We note however that parameters controlling the
428 autoconversion of cloud water to rain were adjusted during the development of the atmo-
429 sphere component, resulting in a reduction in the magnitude of the aerosol forcing by ap-
430 proximately 0.3 W m^{-2} [Rasch *et al.*, 2019].

431 3.2 Spin-up and initialization

432 Since the coupled ocean and sea-ice system take centuries to spin-up, we chose to
433 accelerate the process by simultaneously performing model spin-up and final coupled tun-
434 ing. MPAS-Ocean was initially spun up for one year in stand-alone mode, from the Polar
435 science center Hydrographic Climatology (PHC) dataset [Steele *et al.*, 2001] from rest.
436 During this spin-up, sea surface temperature and salinity were relaxed to an annual mean
437 climatology (from PHC) and the currents were forced by an annual averaged CORE-II
438 normal year wind stress [Large and Yeager, 2009]. This was followed by a ten year Coor-
439 dinated Ocean-ice Reference Experiments (CORE-II), inter-annually varying, forced [Large
440 and Yeager, 2009] ocean and sea-ice simulation. The ocean and sea-ice state at the end of
441 that simulation served as initial conditions for a series of sequential fully coupled simula-
442 tions with pre-industrial forcing (Fig. 2).



443 **Figure 2.** Time evolution of annual global mean surface air temperature for the spin-up simulations. Year 1
 444 marks the first year of CMIP6 DECK *piControl*.

445 The first of these simulations consisted of 50 years with a prototype E3SMv1 con-
 446 figuration (“FCT2”). Due to the change in forcing from present-day CORE-II to pre-
 447 industrial coupled, the simulation experienced a rapid cooling during the first 20 years fol-
 448 lowed by a relative stabilization. Ocean and sea-ice states after 50 years served as initial
 449 conditions for a new coupled simulation (“beta3rc10”) that included bug fixes and associ-
 450 ated retuning described in Table 1, full implementation of the CMIP6 *piControl* forcings,
 451 and a retuning of `cLubb_c14` from 1.3 to 1.2 based on a few previous tuning attempts
 452 with the “FCT2” model configuration. This beta3rc10 configuration was run for a total
 453 of 260 years; it exhibited excessive net influx of energy at TOA and a long-term warm
 454 drift with average temperature approaching present-day values near the end, even with
 455 pre-industrial forcings. As a result, a new configuration (“beta3rc11”) was branched off
 456 after 150 years reusing all components states as initial conditions, but with a slight retun-
 457 ing of `cLubb_c14` (1.2 to 1.1). This new configuration was run for a total of 100 years.
 458 Initially, it cooled slightly from the impact of the retuning, then drifted warm, albeit at a
 459 smaller pace than its predecessor. The process was repeated once again with “beta3rc13”,
 460 further reducing `cLubb_c14` from 1.1 to 1.08. This last simulation was run for a total of
 461 80 years. The drift was smaller than the preceding attempt, but 80 years was too short to
 462 conclude that the drift had been removed entirely.

463 At that point, a decision was made to incorporate a minor code change in the at-
 464 mospheric dynamical core. This change was a workaround for a compiler bug ([https:](https://github.com/NCAR/cesm/issues/17)

465 [//github.com/E3SM-Project/E3SM/pull/1922](https://github.com/E3SM-Project/E3SM/pull/1922)). In the interest of time, two new sim-
 466 ulations were started simultaneously from the end of beta3rc13, one with identical tuning
 467 (“v1a”) and one with a further retuning of the `clubb_c14` parameter from 1.08 to 1.06
 468 (“v1b”). After 110 years, both simulations were examined and “v1b” was selected as the
 469 official E3SMv1 tuning because it exhibited the smaller drift of the two.

470 **Table 2.** Summary of E3SMv1 simulations.

Label	Description	Period	Ens.	Initialization
<i>piControl</i>	Pre-industrial control	500 years	-	Pre-industrial spinup (Sect. 3.2)
<i>1pctCO2</i>	Prescribed 1% yr ⁻¹ CO ₂ in- crease	150 years	1	<i>piControl</i> (101)
<i>abrupt-4xCO2</i>	Abrupt CO ₂ quadrupling	150 years	1	<i>piControl</i> (101)
<i>historical_Hn</i>	Historical	1850-2014	5	<i>piControl</i> (101, 151, 201, 251, 301)
<i>amip_An</i>	Atmosphere with prescribed SSTs and sea-ice concentration	1870-2014	3	<i>historical_Hn</i> (1870)
<i>amip_1850allF_An</i>	Same, but with all forcings held at 1850 values	1870-2014	3	<i>historical_Hn</i> (1870)
<i>amip_1850aeroF_An</i>	Same, but with all aerosol forcings held at 1850 values	1870-2014	3	<i>historical_Hn</i> (1870)

471 3.3 Simulation campaign

472 Table 2 summarizes the E3SMv1 simulation campaign. All simulations were con-
 473 figured to adhere to the CMIP6 DECK specifications [Eyring *et al.*, 2016] as closely as
 474 possible (see Appendix B: for details about input data). *piControl* spans a total of 500
 475 years. Jan 1 of year 101 from *piControl* served as initial conditions for the idealized 1%
 476 yr⁻¹ CO₂ increase (*1pctCO2*) and abrupt CO₂ quadrupling (*abrupt-4xCO2*) simulations, as
 477 well as the first member of the historical simulations (*historical_H1*). Subsequent mem-
 478 bers were branched every 50 years for a total of five ensemble members. AMIP simula-
 479 tions (prescribed SST) were also performed to cover the entire period for which CMIP6
 480 provides surface boundary conditions (1870-2014). Atmosphere and land initial condi-

481 tions for *amip_AI* were taken from year 1870 of *historical_HI*, and similarly for the other
482 members.

483 The entire simulation campaign was performed on NERSC Edison. E3SMv1 did not
484 experience any internal failures during the entire 2930 simulated years. The only failures
485 were system related (generally node failures or file system issues) and the model could
486 always be restarted from the last available set of annual restart files. The source code git
487 hash for the *piControl* simulation was 2e145acf ([https://github.com/E3SM-Project/
488 E3SM/commit/2e145acf](https://github.com/E3SM-Project/E3SM/commit/2e145acf)) and for the remaining simulations was 7de18fc7 ([https:
489 //github.com/E3SM-Project/E3SM/commit/7de18fc7](https://github.com/E3SM-Project/E3SM/commit/7de18fc7)). The difference in source
490 code version arises from code modifications necessary to support transient simulations that
491 were not in place when the *piControl* simulation was started. The newer 7de18fc7 code
492 bit-for-bit (BFB) reproduces 2e145acf for *piControl*. A maintenance branch (maint-1.0;
493 <https://github.com/E3SM-Project/E3SM/tree/maint-1.0>) has also been specifi-
494 cally created to reproduce these simulations. BFB results on NERSC Edison will be main-
495 tained on that branch for as long as the computing environment supports it.

496 **4 Pre-industrial control simulation**

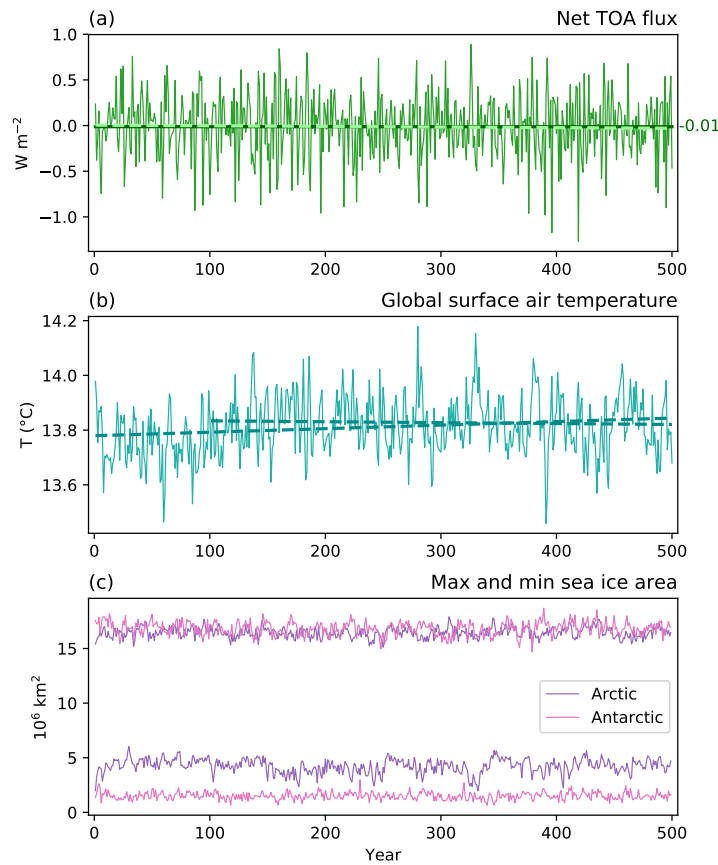
497 The pre-industrial control simulation (*piControl*) is in energy balance at TOA with
498 an average loss of 0.011 W m^{-2} over the course of 500 years (tuning objective 1; Fig. 3a)
499 and almost no long-term linear trend. Among all the model components, the ocean con-
500 stitutes the largest reservoir of heat. It takes up heat at an average rate of 0.016 W m^{-2} ,
501 leaving a small net imbalance of 0.027 W m^{-2} , either from changes in other components
502 or energy non-conservation. Since this imbalance is sufficiently small compared to anthro-
503 pogenic forcings of interest, E3SMv1 can be regarded as essentially conserving energy.
504 We note however that developmental versions of E3SMv1 suffered from much larger im-
505 balances (on the order of 0.5 W m^{-2}). That imbalance was caused by inconsistent defini-
506 tions of energy in the ocean and the atmosphere, with the ocean properly accounting for
507 changes in water heat content with temperature while the atmosphere did not. The 0.5 W
508 m^{-2} imbalance was deemed too large to ignore, but rewriting the atmosphere physics more
509 consistently was impractical due to time constraints. As a result, we decided to incorpo-
510 rate an *ad hoc* correction term. An additional surface (sensible heat) source is introduced
511 to the atmosphere that accounts for the missing energy carried by water molecules leaving
512 the ocean at one temperature, and returning (as condensed water) at another temperature.

513 The apparent flux is small (0.4 W m^{-2}), and applied uniformly at every point globally.
514 The correction doesn't impact the simulated climatology, but its neglect would had led to
515 a long term drift in the global mean temperature of a few tens of degree by century. With
516 this correction, both atmosphere and ocean communicate consistent energy loss and gain
517 with each other (see Appendix A: for more detail).

518 The global mean surface air temperature is very stable over the course of the control
519 simulation (tuning objective 2; Fig. 3b). There is a slight positive temperature trend of
520 $+0.013^\circ\text{C}/\text{century}$ over the entire 500 years, and a much smaller trend of $-0.003^\circ\text{C}/\text{century}$
521 over the last 400 years. We note that all additional simulations were branched from *piControl*
522 in year 101 or later (Table 2), thus ensuring that forcing perturbations are applied atop
523 a non-drifting control simulation. This facilitates subsequent analyses, since it is not nec-
524 essary to correct for a long term drift in the underlying control. Another measure of the
525 stability of the climate in the *piControl* simulation is provided by the annual maximum
526 and minimum sea ice area in Arctic and Antarctic (Fig. 3c). Both hemispheres have little
527 long term drift in their seasonal cycle of sea ice extent.

528 Among 42 CMIP5 models analyzed by *Hawkins and Sutton* [2016], a majority of
529 them have pre-industrial global mean surface temperatures between 13 and 14°C , with
530 some as low at 12°C and as high as 15°C . With a long-term average surface temperature
531 of 13.8°C , E3SMv1 falls within the range of the bulk of CMIP5 models (tuning objec-
532 tive 3). This value is also consistent with estimated warming and the present-day global
533 temperature of $14.0 \pm 0.5^\circ\text{C}$ by *Jones et al.* [1999] for the period 1961-1990 and with lead-
534 ing re-analyses datasets (14.3 to 14.6°C) for the period 1979-2008 [*Hawkins and Sutton,*
535 2016].

540 Water conservation is also important for a global Earth system model. During the
541 early development of the E3SMv1 model, we identified a number of computational prob-
542 lems, leading to water conservation errors in the atmosphere component [*Zhang et al.,*
543 2018]. These problems were all considered, and suitable remedies applied (e.g., borrow-
544 ing from adjacent cells to avoid more drastic non-conservative fixers) so that in the 500
545 year *piControl* simulation, the relative water conservation error (as defined in *Zhang et al.*
546 [2018]) in the atmosphere component is about 0.00226% , equivalent to a computational
547 sea level rise of about $2 \text{ mm century}^{-1}$.



536 **Figure 3.** Time evolution of annual (a) global mean net TOA radiation (positive down), (b) global mean
 537 surface air temperature, and (c) maximum and minimum of total sea ice area for the Arctic and Antarctic in
 538 the *piControl* simulation. Dashed lines in (a) and (b) represent linear trends. The solid straight line in (a) is
 539 the mean TOA energy imbalance of -0.011 W m^{-2} .

548 Additional problems were also identified in the fully coupled system, including non-
549 conservation caused by inconsistent remapping in the exchange of water between model
550 components, excessive storage capacity in the river routing model, missing perched drainage
551 and ponding in the land surface model, and grid definition inconsistencies between dif-
552 ferent model components. Some of these errors caused non-conservation errors in sea
553 level rise as large as tens of meters per century but all were resolved and corrected. Un-
554 fortunately, a smaller non-conservation error was not uncovered until after E3SMv1 was
555 frozen and the DECK simulations were complete. It manifests itself by a steady loss of
556 water in the ocean of $5 \text{ cm century}^{-1}$ in the *piControl* simulation. We note that the water
557 loss does not involve any phase change and therefore does not impact TOA energy bal-
558 ance; however, it does impact ocean heat content. The root cause has since been traced
559 to a nonphysical update of water stored in the unconfined aquifer of urban subgrid land
560 units within the land model. This error, which has been corrected for future versions of
561 E3SM (<https://github.com/E3SM-Project/E3SM/pull/2603>), was partly masked by
562 incomplete and inconsistent water budget checks within the land model and the coupler.
563 These budgets have also been corrected.

564 **5 Historical simulations**

565 **5.1 Atmosphere climatology**

566 We first evaluate the atmosphere in E3SMv1 with climatologies from the final 30
567 years (1985-2014) of simulations. Five ensemble members of the historical coupled exper-
568 iment (H1 to H5) and three ensemble member AMIP simulations (A1 to A3) are evalu-
569 ated. Similar figures are presented in *Rasch et al.* [2019] with comparisons to CAM5.

570 The net TOA radiative flux (positive down) is shown in Figure 4 compared to CERES-
571 EBAF Ed4.0 [Loeb *et al.*, 2009]. Bias patterns are qualitatively similar between the atmo-
572 sphere and coupled simulations, indicating a strong imprinting of atmosphere biases on
573 the coupled system, but the biases are frequently larger in the coupled simulations, where
574 component interactions play a role (RMSE of 9.13 vs 7.80 W m^{-2}). Significant posi-
575 tive biases are seen in the subtropical stratocumulus regions off the west coasts of North
576 America, South America and Africa due to an underestimate of cloudiness. The under-
577 predicted coastal stratocumulus is due to both the use of CLUBB and model re-tuning *Xie*
578 *et al.* [2018]. Southern Oceans are also marked by positive biases due to clouds. The re-

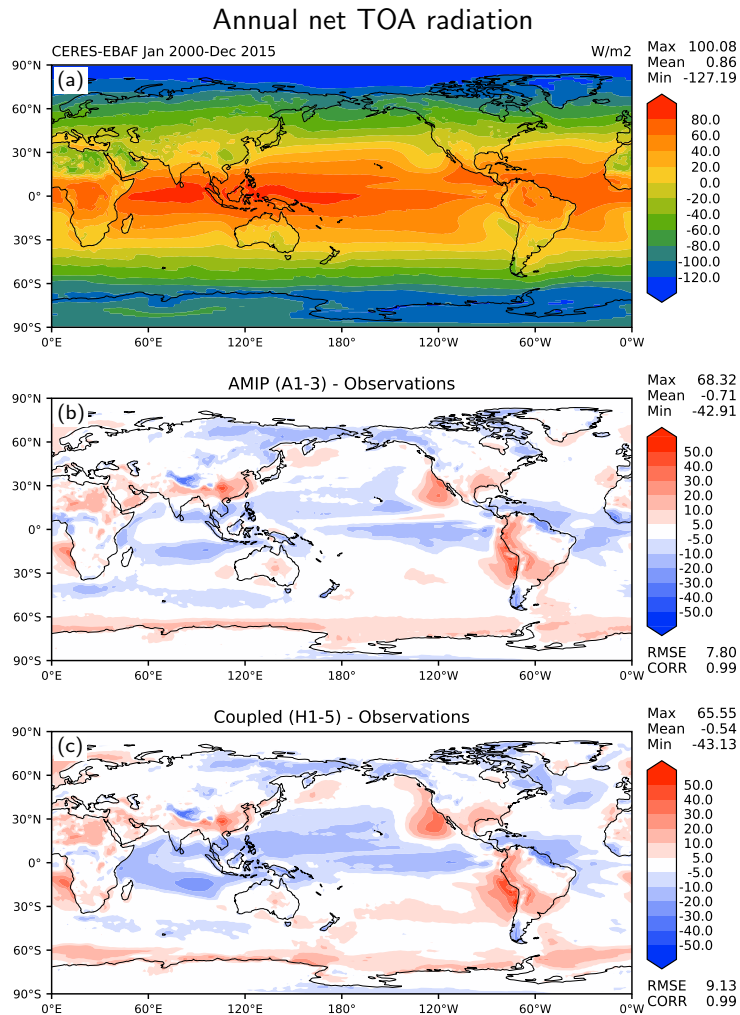
579 maining oceanic regions have either neutral or negative biases. Tropical land masses have
580 generally positive biases, while northern high latitudes suffer from negative ones. The
581 global mean net flux is $+0.32 \text{ W m}^{-2}$ for the coupled simulations and $+0.15 \text{ W m}^{-2}$ for
582 the AMIP simulations, both slightly smaller than the observational estimates of 0.6 W m^{-2}
583 from CERES-EBAF Ed 4.0. [uncertainty range of 0.2 to 1.0; *Wild et al.*, 2012]. We note
584 that the present-day energy imbalance was not a target of model tuning. Instead, we tuned
585 the pre-industrial control simulation to have a net zero global mean TOA flux. The imbal-
586 ance for 1985-2014 emerges from the model evolution in response to historical forcings in
587 the case of the coupled simulations or as a result of imposing present-day SST and sea-ice
588 concentrations (as boundary conditions) for the AMIP simulations.

592 Biases in cloud radiative effects (CRE) at TOA compared to CERES-EBAF Ed4.0
593 [*Loeb et al.*, 2009] are shown in Figure 5 for SW and LW. Many regional biases appar-
594 ent in net TOA (Fig. 4) can be traced to SW cloud biases. LW CRE reveal additional bi-
595 ases, such as a lack of LW cloud trapping in the tropics from high clouds and excessive
596 LW trapping from northern high latitudes clouds. As indicated in *Xie et al.* [2018], high
597 clouds are significantly reduced in the tropical deep convection regions due to the increase
598 of model vertical resolution from 30 levels to 72 levels in EAM, which results in a much
599 weaker LW CRE over these regions. Global mean CRE for both SW and LW are approxi-
600 mately 3 W m^{-2} below the observational estimates.

603 Annual precipitation is depicted in Figure 6 compared to GPCP v2.2 [*Adler et al.*,
604 2003; *Huffman et al.*, 2009]. Simulated global mean precipitation is slightly above 3 mm
605 day^{-1} for both AMIP and coupled simulations, approximately 15% larger than the GPCP
606 estimate. A comprehensive review of 30 currently available global precipitation data sets
607 found large differences in the magnitude of global land annual precipitation estimates
608 [*Sun et al.*, 2018]. In their estimate of the global energy cycle, *Wild et al.* [2012] note
609 that the GPCP estimate may be too low due to systematic underestimations in the satel-
610 lite retrievals [*Trenberth et al.*, 2009; *Stephens et al.*, 2012]. *Wild et al.* [2012] estimated
611 the global mean precipitation at $2.94 \pm 0.17 \text{ mm day}^{-1}$ while the estimate from *Stephens*
612 *et al.* [2012] is $3.04 \pm 0.35 \text{ mm day}^{-1}$. The global mean precipitation from E3SMv1 falls
613 within both estimates.

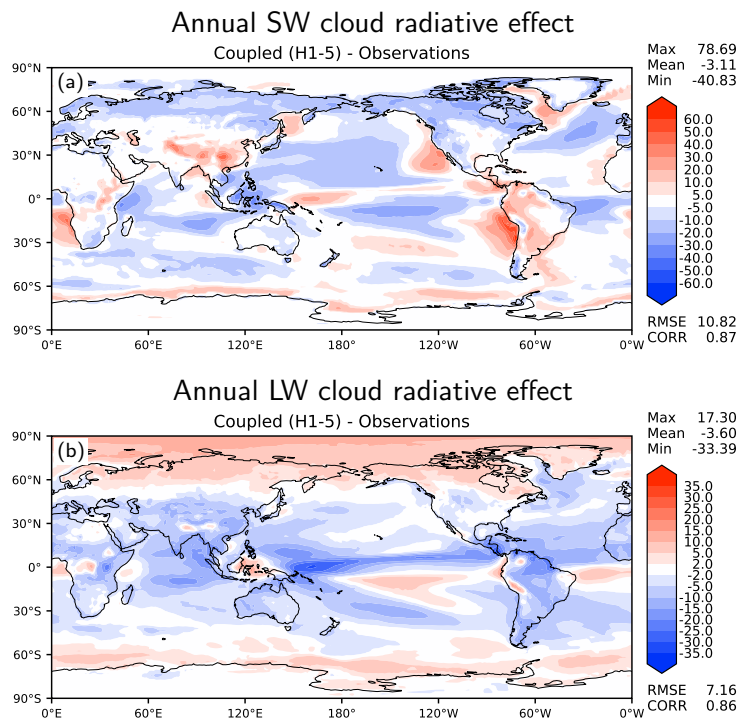
614 Major regional precipitation biases in AMIP simulations (Fig. 6b) include a dry
615 Amazon, excessive precipitation over elevated terrain (e.g. Andes, Tibetan plateau), wet

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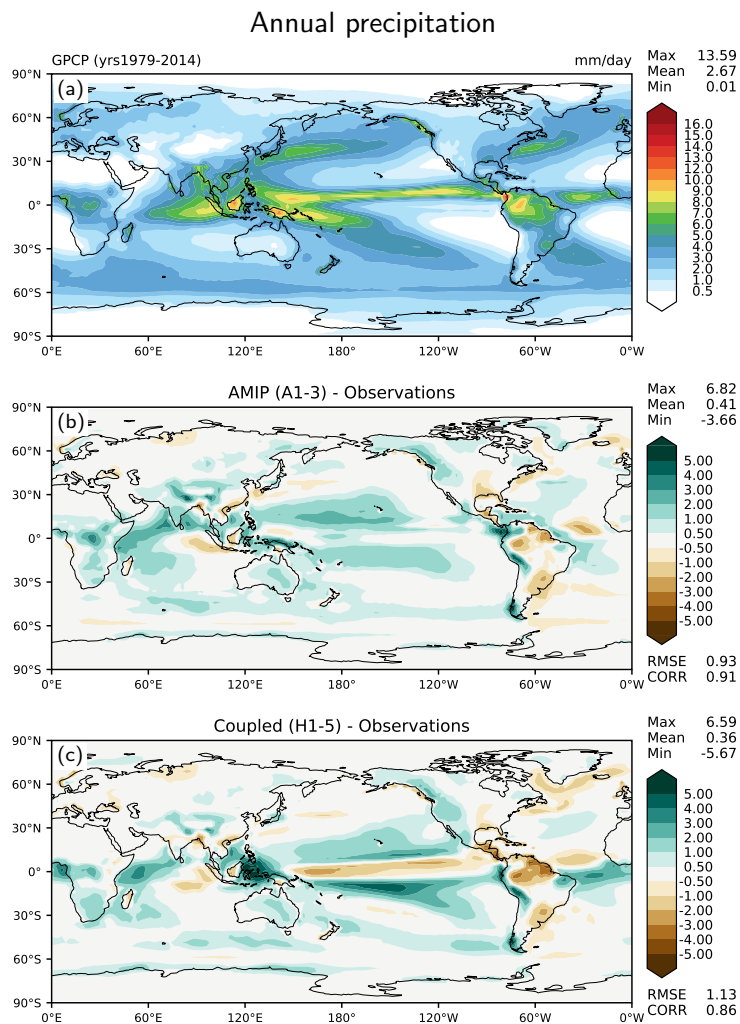
589 **Figure 4.** Annual net top-of-atmosphere (TOA) radiative flux: (a) CERES-EBAF Ed4.0 observational estimate,
 590 (b) model bias from the three AMIP ensemble simulations and (c) model bias from the five ensemble
 591 historical coupled simulations (1985-2014).

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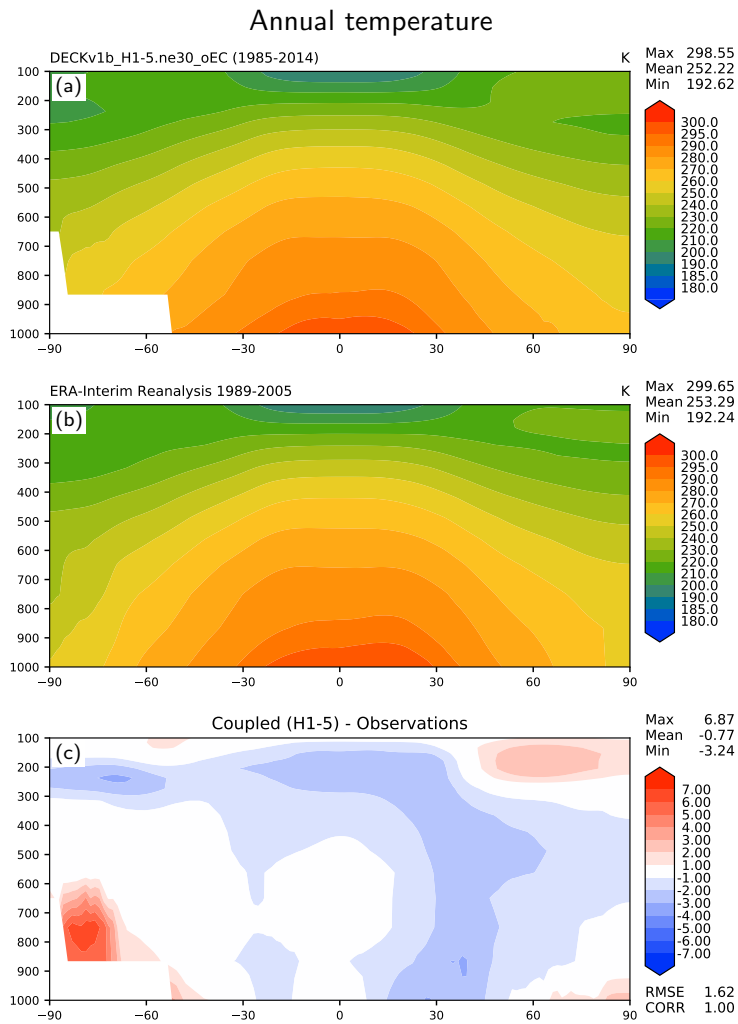
601 **Figure 5.** Annual top-of-atmosphere cloud radiative effect model biases (W m^{-2}) of the five ensemble
602 historical simulations (1985-2014) against CERES-EBAF Ed4.0: (a) SW and (b) LW.

616 biases over tropical Africa and the Indian ocean, and a dry bias over the central US (traced
 617 to the Summer time). Coupled simulations (Fig. 6c) tend to amplify regional biases present
 618 in AMIP simulations (except the central US), as well as develop biases typical of coupled
 619 simulations: double ITCZ and excessive precipitation over the maritime continent. Un-
 620 surprisingly, the RMSE error increases from 0.93 to 1.13 mm day⁻¹ and the correlation
 621 decreases (0.93 vs 0.86) between AMIP and coupled simulations.



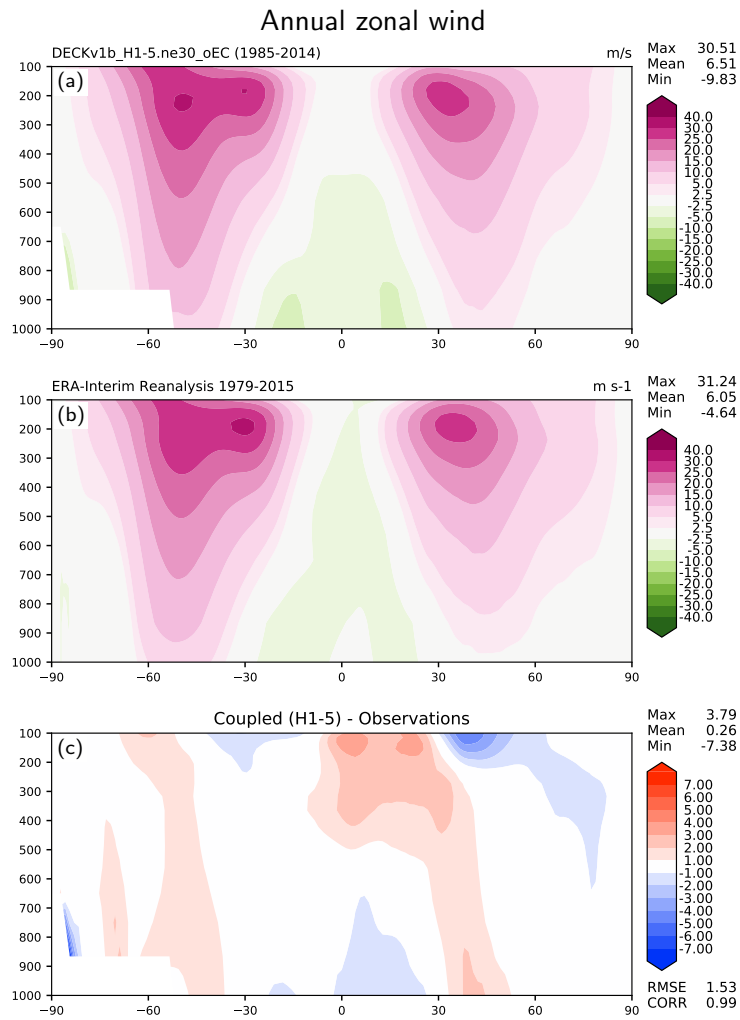
622 **Figure 6.** Annual precipitation rate: (a) GPCP v2.2 observational estimate, (b) model bias from the three
 623 AMIP ensemble simulations and (c) model bias from the five ensemble historical coupled simulations (1985-
 624 2014).

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625 **Figure 7.** Annual zonally averaged temperature: (a) ensemble mean of historical coupled simulations
 626 (1985-2014), (b) ERA-Interim reanalysis, (c) model bias.

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627 **Figure 8.** Annual zonally averaged zonal wind: (a) ensemble mean of historical coupled simulations
 628 (1985-2014), (b) ERA-Interim reanalysis, (c) model bias.

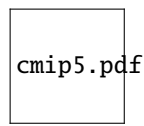
629 Figure 7 shows the annual, zonally averaged temperature for the coupled simulations
630 compared to ERA-Interim reanalysis [Dee *et al.*, 2011]. Overall, the model captures the
631 thermal structure of the atmosphere. There is a tropical cold bias in the upper troposphere
632 which correlates with insufficient tropical high clouds (Fig. 5b). In the northern mid-
633 latitudes, the cold bias extends to the surface because of the colder than observed northern
634 latitudes SSTs (see Section 5.2 below). There is also a significant warm bias in the lower
635 portion of the troposphere in the southern high latitudes. The corresponding zonal wind is
636 shown in Figure 8. The locations of the maximum jet locations are displaced equatorward
637 and the wind magnitude is also too large throughout the mid-latitude troposphere. Tropical
638 easterlies are too weak in the upper troposphere and too strong near the surface.

639 The evaluation above provides only a limited view of the performance of E3SMv1
640 from an atmospheric perspective. For a more exhaustive evaluation, we turn to a com-
641 parison with an ensemble of 45 CMIP5 models using metrics computed with the PCMDI
642 Metrics Package [PMP; Gleckler *et al.*, 2016, 2008]. The comparison covers the period
643 1981-2005 of the historical simulations. The historical and AMIP ensemble members of
644 E3SMv1 were also processed with PMP. Figure 9 shows global RMSE for the CMIP5
645 ensemble using box and whiskers plot, as well as the individual E3SMv1 historical sim-
646 ulations with blue dots and AMIP with red dots. Spatial RMSE against observations col-
647 lected by PMP are shown for nine fields, and each one of them for annual and seasonal
648 averages. Lower values are better. For TOA radiation fields (Fig. 9a-c), E3SMv1 cou-
649 pled (blue) generally falls within the lowest (best) quartile, and is even competitive with
650 some of the best CMIP5 models for certain fields and seasons. We note that we are com-
651 paring a newer model against older ones, so we do not expect this to necessarily hold for
652 CMIP6. For surface variables, precipitation (Fig. 9d), surface air temperature over land
653 (Fig. 9e), and zonal wind stress over ocean (Fig. 9f), E3SMv1's coupled performance is
654 better than the ensemble median and often falls within the lowest quartile, with the ex-
655 ception of surface air temperature over land during DJF and MAM. Precipitation during
656 MAM is also notably worse than other seasons relative to the CMIP5 ensemble. For dy-
657 namical variables, zonal wind at 200 and 850 hPa (Fig. 9g-h) and 500 hPa geopotential
658 height (Fig. 9i), E3SMv1 is generally better than the CMIP5 median, except again for
659 MAM.

660 Unsurprisingly, AMIP simulations (red) perform better than their coupled counter-
661 parts (blue). However, the relative degradation between AMIP and coupled helps attribute

662 sources of errors. For example, the difference is relatively small for the zonal mean wind,
 663 and thus improving overall performance would likely require atmospheric improvements.
 664 On the other hand, surface variables, in particular precipitation and temperature, are much
 665 more strongly affected by the coupled model errors that emerge. The MAM seasonal defi-
 666 ciency also appears to be rooted in coupling errors.

667 In summary, E3SMv1 performs better than the median of the of CMIP5 ensem-
 668 ble for most atmospheric fields and seasons, which helps establish the credibility of the
 669 E3SMv1 simulated climate.



670 **Figure 9.** Comparison of RMSE (1981-2005) of an ensemble of 45 CMIP5 models (box and whiskers
 671 showing 25th and 75 percentiles, minimum and maximum) with the five E3SMv1 historical members (blue
 672 dots) and the first member of the AMIP simulations (ret dots). Spatial RMSE against observations are com-
 673 puted for annual and seasonal averages with the PCMDI Metrics Package [Gleckler *et al.*, 2016]. Fields shown
 674 include TOA net radiation (a), TOA SW and LW cloud radiative effects (b,c), precipitation (d), surface air
 675 temperature over land (e), zonal wind stress over ocean (f), 200 and 850 hPa zonal wind (g,h) and 500 hPa
 676 geopotential height (i).

677 5.2 Ocean climatology

678 An annual climatology (1985-2014) of the E3SMv1 ensemble mean of sea surface
 679 temperature (SST) is shown in Figure 10, compared to the Hadley-NOAA/OI merged data
 680 product [Hurrell *et al.*, 2008] averaged over the same period. Overall, E3SMv1 captures
 681 the observed SST well, with an ensemble mean bias of 0.093°C and an ensemble mean
 682 RMSE of 0.939°C. A few biases do emerge. First, there is a cold bias in the North At-
 683 lantic associated with excessive sea ice in the Labrador Sea. Consistent with this sea ice
 684 bias, these cold SST biases are stronger in the first half of the year than in the second half
 685 (not shown). While the exact cause of the Labrador Sea ice bias is unknown, it is likely
 686 that missing critical heat transports (*e.g.*, from the East and West Greenland currents,
 687 Irminger current, and the Northwest corner), perhaps from low resolution or excessive
 688 horizontal viscosity [Jochum *et al.*, 2008] play an important role. Second, there are warm

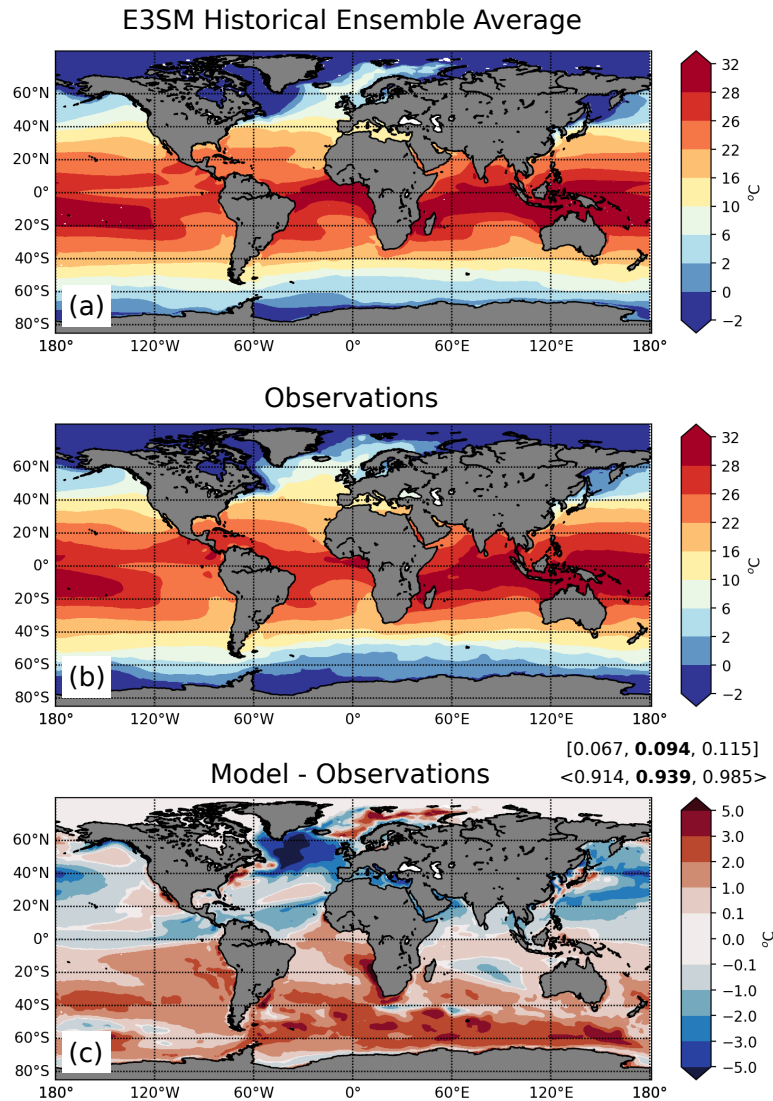
689 biases on the eastern sides of ocean basins, coincident with SW CRE biases discussed in
690 sec. 5.1. Finally, the Southern Ocean is also $\approx 2^\circ\text{C}$ warmer than observed. It is possible
691 this last bias is associated with the relatively large value for the Gent-McWilliams bolus
692 coefficient used, which was chosen to prevent excess heat transport to the deep ocean (see
693 sec. 2.2).

698 The E3SMv1 sea surface salinity (SSS) bias is shown in Fig. 11. The model out-
699 put is compared to the 2011-2014 NASA-Aquarius data [Lagerloef *et al.*, 2015]. Overall
700 the surface ocean is too fresh, with the largest bias in the Labrador Sea. The two minor
701 exceptions to this pattern are the positive salinity bias region over the west Pacific warm
702 pool, which is coincident with the atmospheric precipitation maximum being shifted west-
703 ward relative to observations (Fig. 6) and a positive salinity bias in the Arctic.

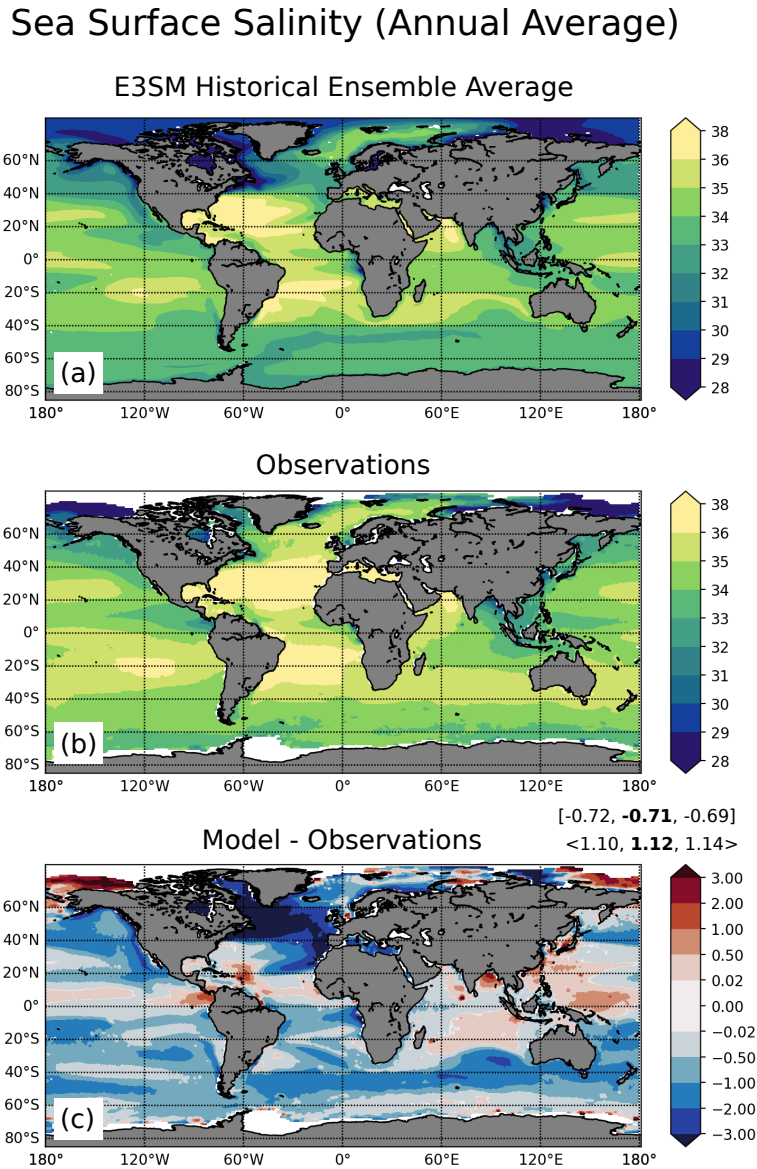
708 The ensemble average Mixed Layer Depth (MLD) annual climatology, based on a
709 critical density threshold ($\sigma_c = 0.03 \text{ kg m}^{-3}$), is presented in Figure 12. The model output
710 is compared to data described by Holte *et al.* [2017]. The globally averaged model MLD
711 is too shallow relative to observations, with the largest bias coincident with the large fresh
712 bias in the North Atlantic. Unlike many other CMIP5 models [e.g., Sallée *et al.*, 2013],
713 the E3SMv1 MLD in the Southern ocean is slightly deeper than observed. This is due
714 to the region of deeper mixed layers in E3SMv1 being broader than observed (Fig. 12a),
715 possibly due to a positive bias in the Southern Ocean wind stress. The maximum MLD
716 simulated by E3SMv1 is still much shallower than observations, which is consistent with
717 other CMIP5 models.

718 Figure 13 shows the E3SMv1 maximum Atlantic Meridional Overturning Circu-
719 lation (AMOC) at 26°N , the site of the RAPID array [Smeed *et al.*, 2017]. The mean
720 AMOC simulated in the historical ensemble (approximately 11 Sv) is weaker than the ob-
721 served mean ($16.9 \pm 3.35 \text{ Sv}$) and also on the weak end of CMIP model AMOC strength
722 (e.g., Cheng *et al.*, 2013 their Fig. 1). There are a number of possible causes for the weak
723 AMOC in E3SMv1. Currently, MPAS-Ocean utilizes a z-star coordinate [Adcroft and
724 Campin, 2004], which is broadly consistent with a traditional z-level coordinate in the
725 deeper ocean. Z-coordinate models are known to experience spurious diapycnal mixing
726 [e.g., Griffies *et al.*, 2000], which could reduce AMOC strength. Second, recent work has
727 shown that Nordic overflows [Wang *et al.*, 2015] and Arctic Freshwater transports [Wang
728 *et al.*, 2018] have a strong impact on AMOC. At low resolution, these processes are poorly

Sea Surface Temperature (Annual Average)

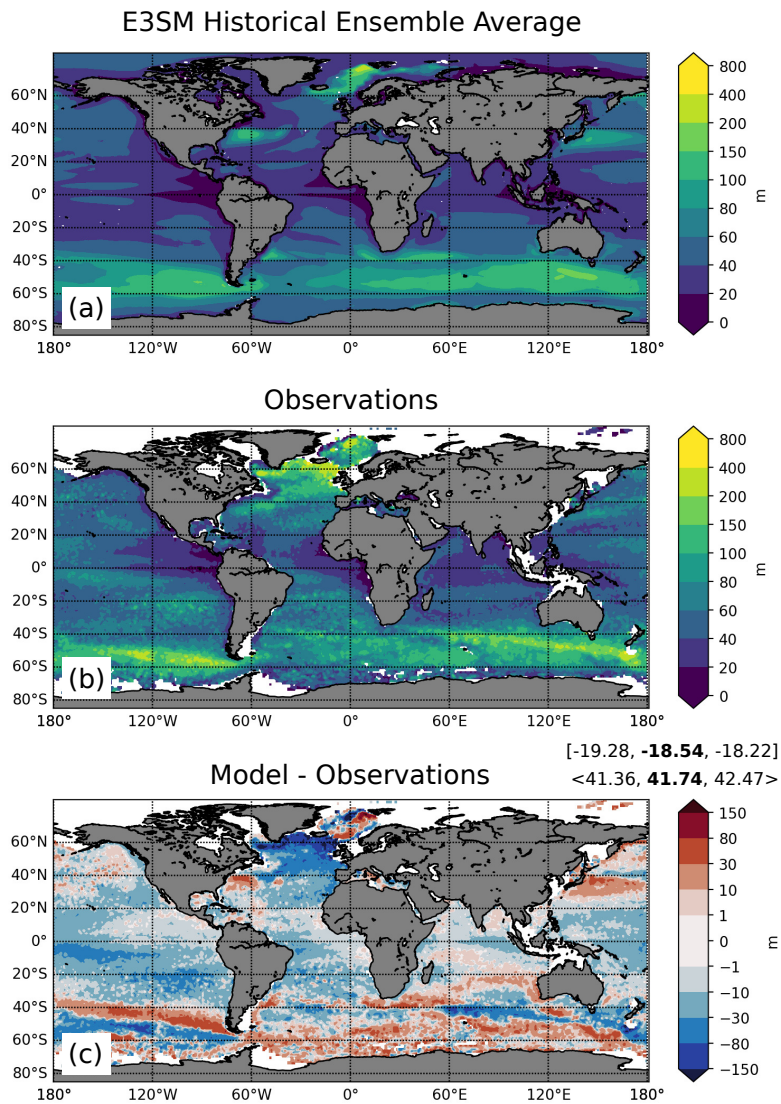


694 **Figure 10.** Annually averaged SST climatology (1985-2014) for (a) E3SMv1 historical simulation (en-
 695 semble mean), (b) Hadley-NOAA/OI merged SST dataset (1985-2014; [Hurrell *et al.*, 2008]), and (c) Model
 696 bias. In the upper right corner of panel (c), the mean bias is shown in square brackets and the RMS error is in
 697 angular brackets. For each error, the min, **mean**, max is listed for the historical ensemble.

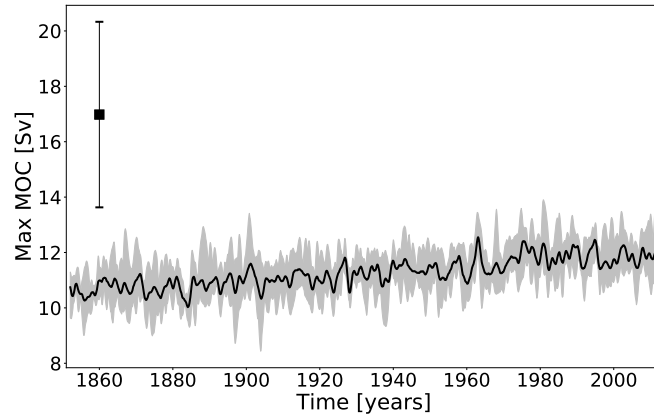


704 **Figure 11.** As in Figure 10 but for Sea Surface Salinity (PSU). The dataset in (b) is from NASA Aquarius
 705 data (averaged 2011-2014).

Mixed Layer Depth (Annual Average)



706 **Figure 12.** As in Figure 10 but for the annual average mixed layer depth. The data in (b) is from *Holte et al.*
707 [2017]. Data has been averaged from 2001-2017.



733 **Figure 13.** Annually averaged maximum Atlantic Meridional Overturning Circulation at 26.5°N below
 734 500m depth. The solid black line represents the ensemble mean and the spread of the shaded gray region
 735 represents the maximum and minimum of the ensemble. The observed AMOC and standard deviation at the
 736 RAPID array is shown by the vertical bar ($16.9 \pm 3.35\text{Sv}$.)

729 represented, where the latter is likely too strong as critical passageways being too wide
 730 (e.g. Davis Strait). Finally, since E3SMv1 exhibits excess sea ice in the Labrador sea, the
 731 simulated MLD (Fig. 12) is reduced, which reduces deep convection and hence AMOC
 732 strength.

737 5.3 Sea ice

738 Sea ice in E3SMv1 is too extensive at the end of winter and too confined in late
 739 summer, reflecting a heightened model seasonality relative to both the Arctic and South-
 740 ern Ocean passive microwave record, regardless of the algorithm used to derive ice area
 741 from these retrievals [e.g., *Cavaliere et al.*, 1996; *Meier et al.*, 2017]. This result is sum-
 742 marized in Figures 14, 15, and 16. Embedded within this heightened seasonality, the melt-
 743 season minimum is delayed in the Northern Hemisphere relative to observations. In the
 744 high north, the heightened winter seasonal extent relative to observations occurs mainly
 745 in the Labrador Sea, as well as the Iceland Sea and Pacific margin of the Sea of Okhotsk,
 746 evident in the historical 5-member ensemble means (Figs. 14a and 14b). Comparison
 747 with satellite-derived albedo, shown for June to August in Figure 17, suggests that surface
 748 radiative processes leading to an albedo bias in the central arctic are unrelated to the phys-
 749 ical exchanges responsible for the Labrador sea bias. As a result, the Labrador Sea bias is

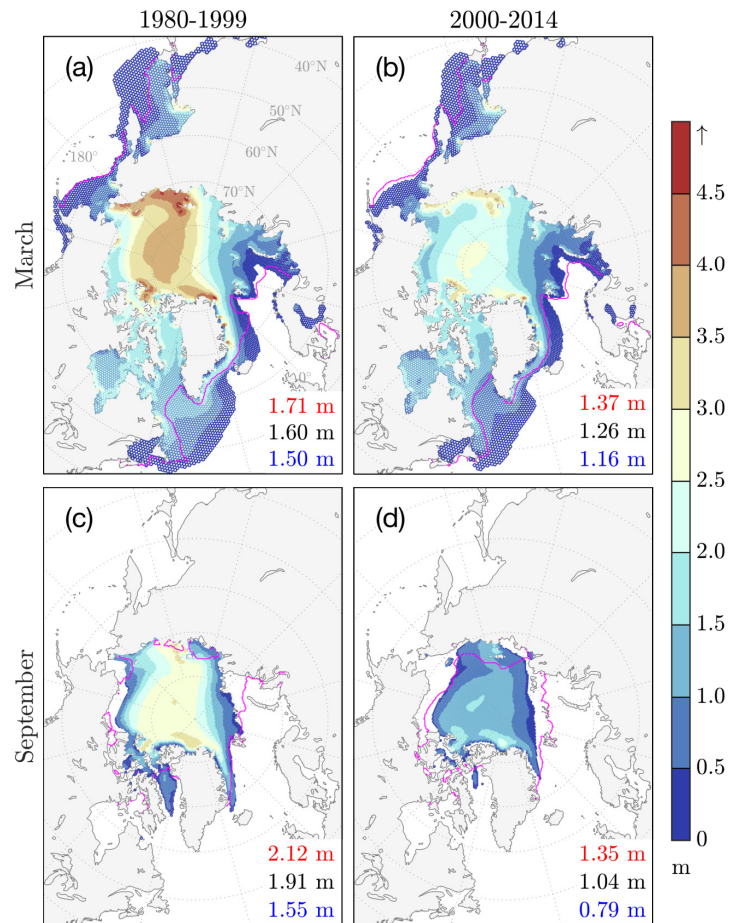
750 unlikely to be attributable to MPAS-Seaice alone, and is more likely a result of coupled
751 interactions and unresolved heat advection in the region.

752 The model possesses a negative climatological trend in Arctic sea ice thickness
753 that is qualitatively consistent with observed basin-wide trends [e.g., *Kwok and Rothrock,*
754 2009], as quantified in Figure 14. The 1980-1999 hemispheric mean spread among the
755 five ensemble members for March (September) decreases from 1.5-1.71 m to 1.16-1.37 m
756 (1.55-2.12 m to 0.79-1.35 m) for the period 2000-2014, which includes the thinnest peri-
757 ods known in the Arctic sea ice thickness record. Still, the spatial ice thickness pattern in
758 the central Arctic does not reflect the predominant build-up of thick ice against the Cana-
759 dian Archipelago seen in observational estimates, including those of *Bourke and McLaren*
760 [1992], *Kwok and Cunningham* [2008], and *Tilling et al.* [2015]. This represents a sec-
761 ond sea ice bias being addressed in ongoing improvements to the polar components of
762 E3SM. There is no significant climatological trend in mean Southern Ocean sea ice thick-
763 ness from the historical simulations, in the ensemble mean or spread (Fig. 15). Extent bias
764 stemming from the heightened polar seasonality of E3SMv1 is consistent across our cho-
765 sen 1980-1999 and 2000-2014 analysis periods, seen in Figure 15. While the model does
766 not replicate the observed climatological trend in austral sea ice extent (Fig. 16), it does
767 simulate the critical decrease in Northern Hemisphere minimum sea ice extent, which has
768 had a significant impact on planetary albedo in the current century and is therefore an im-
769 portant contribution to the energy balance of E3SMv1 as a whole.

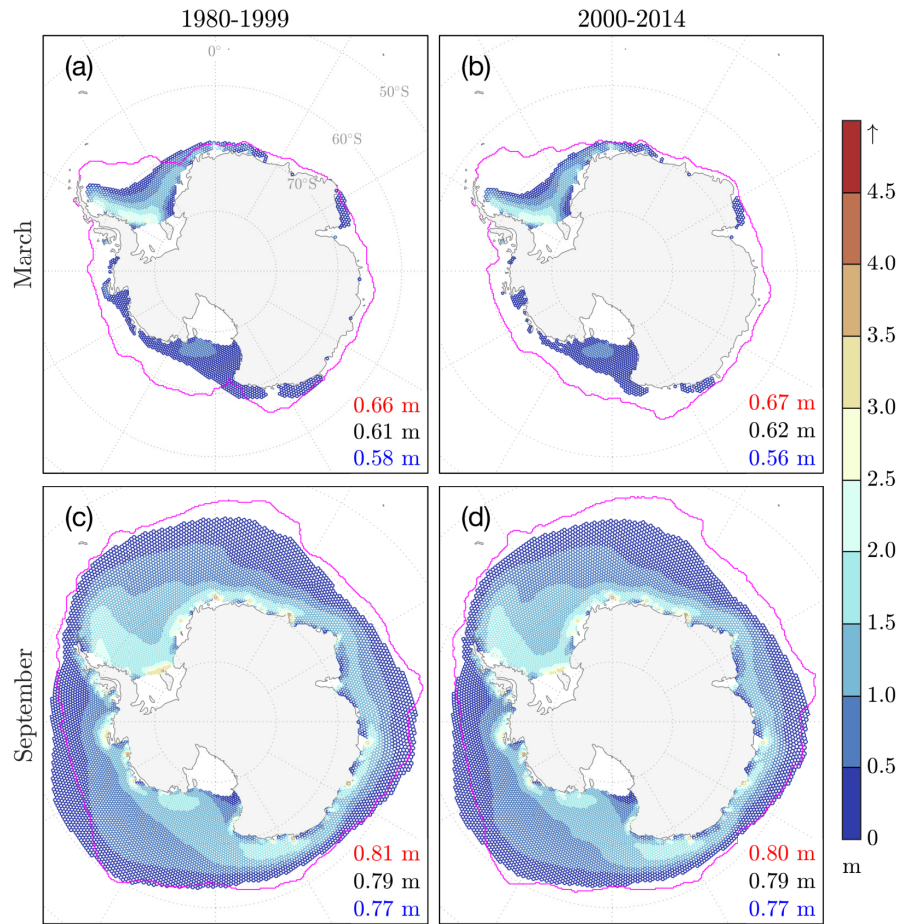
784 5.4 Land and river

791 Figure 18 shows a comparison of the mean annual total (surface and subsurface)
792 runoff simulated by ELM with the composite runoff map from the Global Runoff Data
793 Center (GRDC) [*Fekete and Vörösmarty, 2011*]. Since the GRDC data only provides monthly
794 runoff for 1986-1995, the comparison is shown for the annual total runoff averaged over
795 the same period of the *historical_H1* simulation. ELM captures the general spatial distri-
796 bution of the GRDC runoff, but in relatively arid regions such as Australia and the west-
797 ern U.S., ELM has wet biases, while in the Amazon tropical forest, dry biases are notable.
798 These biases are consistent with the precipitation deficiencies shown in Figure 6.

799 The seasonal cycle of streamflow is an important metric of water availability, as wa-
800 ter deficits resulting from a mismatch in the timing of water supply and water demand

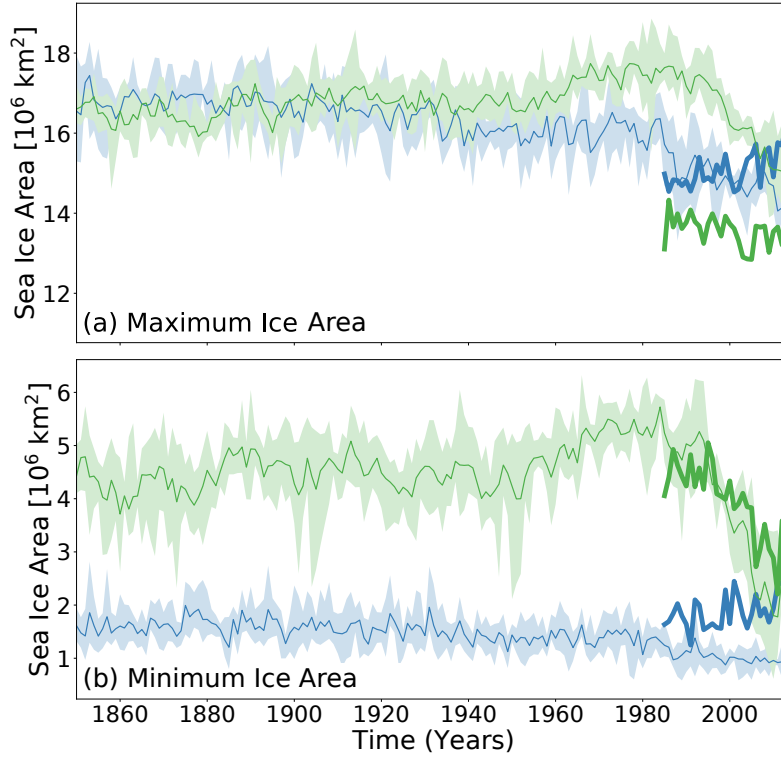


770 **Figure 14.** Northern hemisphere ensemble mean March and September sea ice thickness for two decades
 771 leading up to year 2000 (a and c), and the first 15 years of the 21st century (b and d). Model ice thickness is
 772 truncated at 15% concentration, and magenta represents the *Meier et al.* [2017] NOAA Climate Data Record
 773 (CDR) ice extent for the same averaging periods. Grid density on the polar stereographic projection is indi-
 774 cated by cell translucence. Numbers in the lower right corner of each panel indicate the mean hemispheric ice
 775 thickness for the thinnest ensemble member (blue), the multi-ensemble mean rendered in the map
 776 (black), and the thickest ensemble member (red) for the each period.

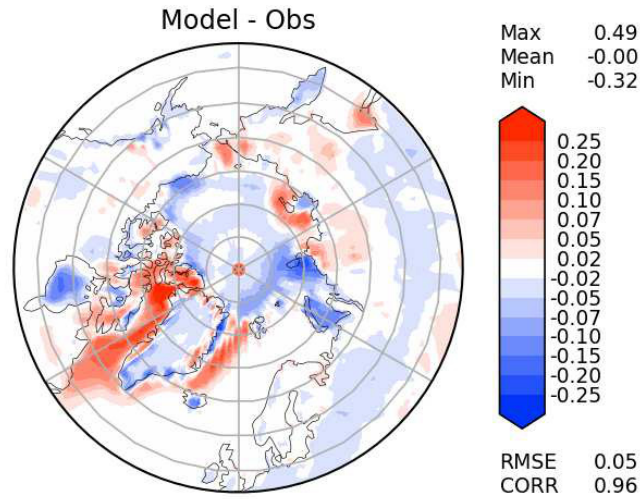


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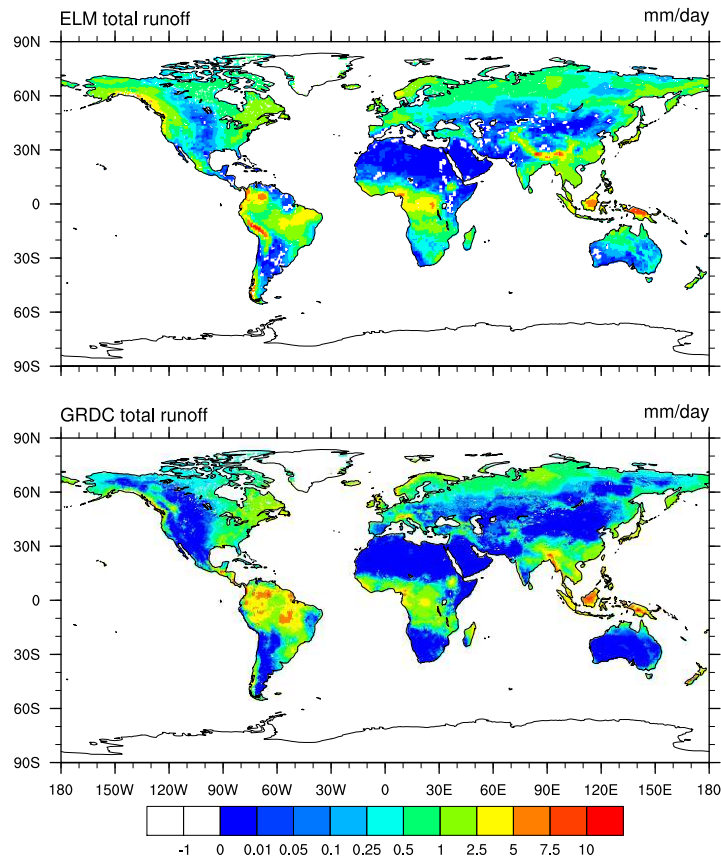
Figure 15. As for figure 14, but for the southern hemisphere.



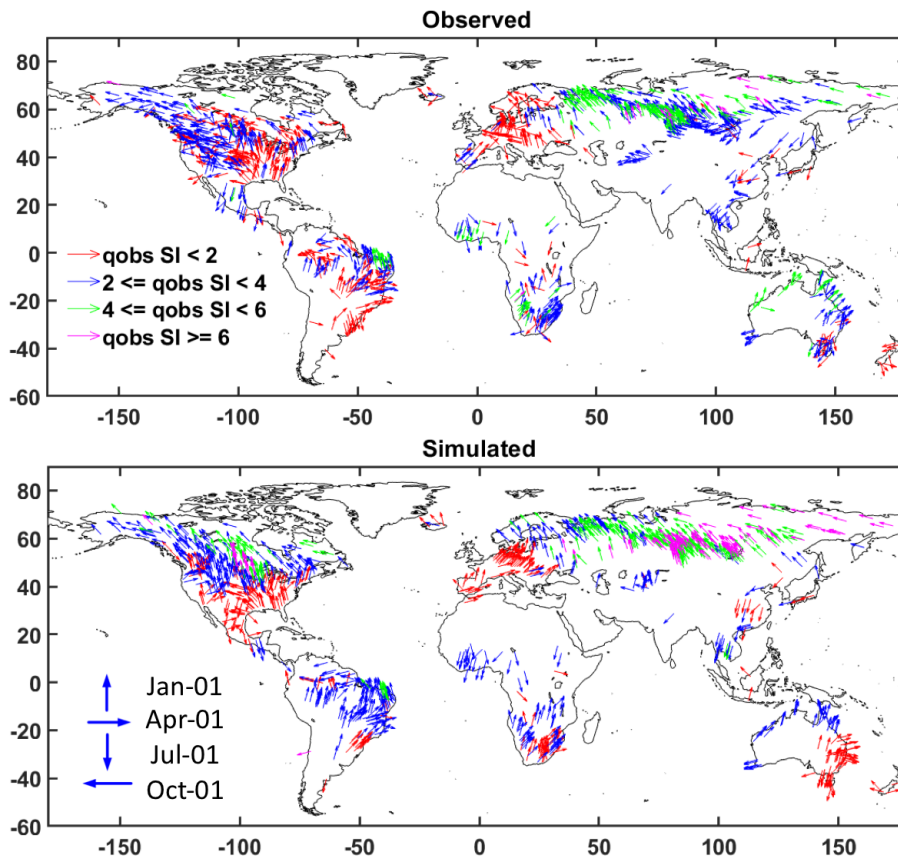
778 **Figure 16.** Maximum (a) and Minimum (b) Sea ice area observed during the year for the Northern (green)
 779 and Southern hemispheres (blue). The shaded bounds represent the E3SMv1 historical ensemble spread, and
 780 the thick colored lines are NASA TEAM observations [Cavalieri *et al.*, 1996]. Model and observational data
 781 are monthly averages.



782 **Figure 17.** Surface albedo bias (E3SMv1 - CERES-EBAF) for northern hemisphere spring, averaged over
 783 1985-2014



785 **Figure 18.** Mean annual total runoff (in mm/day) simulated by ELM (upper panel) and from GRDC (lower
786 panel) averaged between 1986 and 1995.



787 **Figure 19.** Seasonality of observed (upper) and simulated (lower) streamflow at stream gauge stations in
 788 major river basins around the world. The magnitude of seasonality indicated by the seasonality index (SI) is
 789 shown in color and the timing of the peak streamflow is shown by the position of the arrows. SI is calculated
 790 based on monthly streamflow between 1986 and 1995.

801 have important implications for energy and water. Figure 19 compares the seasonal cycle of simulated and observed streamflow at stream gauge stations in major river basins
 802 around the world. A seasonality index (SI) is defined as follows:
 803

$$Pk_i = \frac{12}{n} \sum_{j=1}^n \frac{Q_{ij}}{\sum_{i=1}^{12} Q_{ij}} \quad (1)$$

804

$$SI = \max_i (Pk_i) \quad (2)$$

805 Where Q_{ij} is the monthly streamflow for month i and year j and n is the number of years
 806 in the simulation or observation. With this definition, SI is equal to 1 if the monthly
 807 streamflow is uniformly distributed throughout the year and SI is equal to 12 if stream-
 808 flow only occurs in one month. The seasonality of the simulated streamflow is generally
 809 comparable with that observed in terms of both magnitude and timing. For example, in
 810 North America, streamflow seasonality is stronger in the Northwest with a peak timing
 811 between November and January but in the central and southeastern U.S., streamflow sea-
 812 sonality is weaker with a peak timing generally in spring. In Asia, streamflow generally
 813 peaks in the late summer. Larger biases in seasonality are found in Australia (Murray Dar-
 814 ling River) and Central Asia (Yenisey) where biases in the runoff are also more significant
 815 (Figure 18).

816

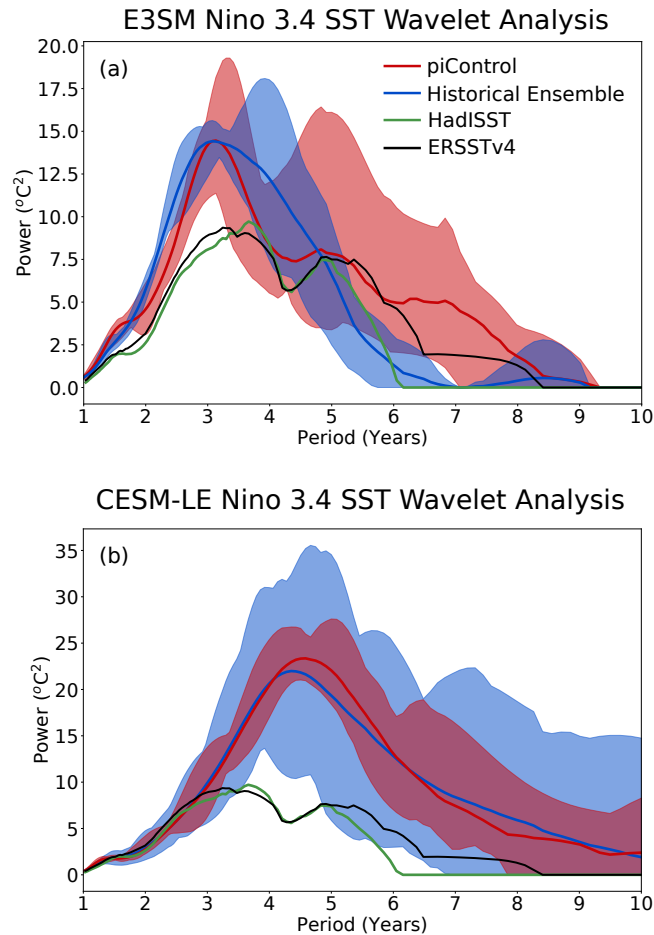
5.5 Variability simulated by E3SMv1

817

In this section we present a variety of metrics to assess E3SMv1 variability.

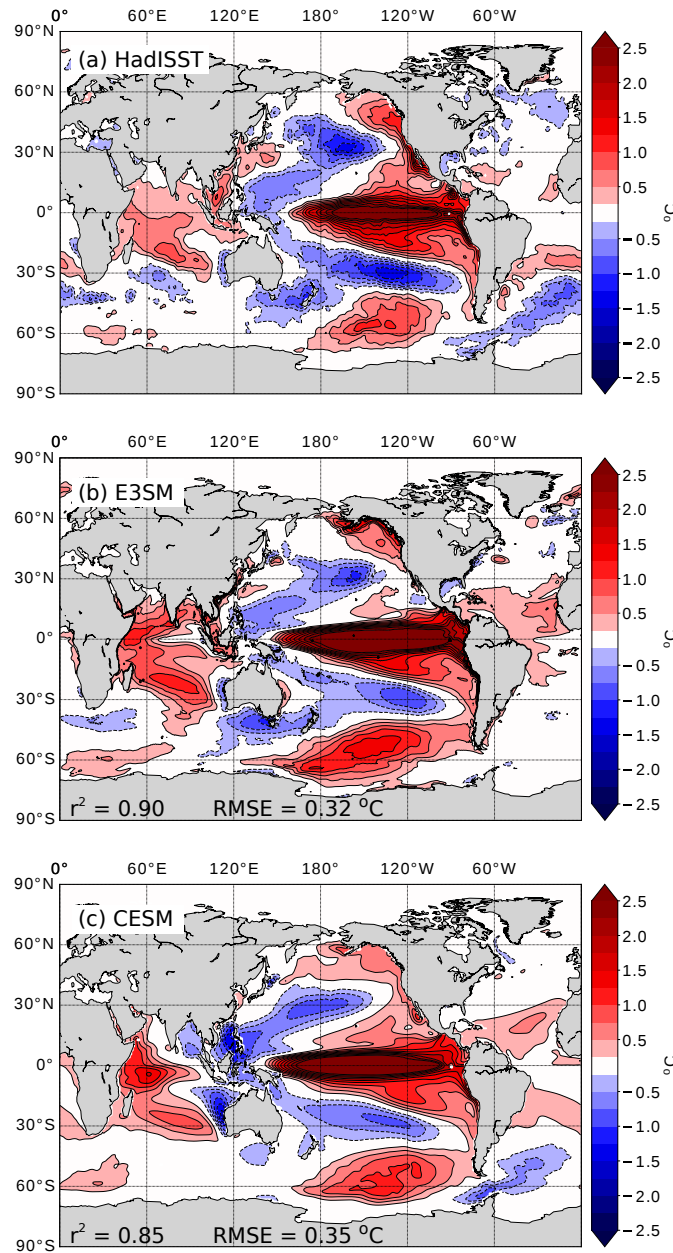
818

819 The long term variability of the El Niño Southern Oscillation (ENSO), as simu-
 820 lated in the *piControl* and historical ensemble simulations is examined via wavelet analysis
 821 [Torrence and Compo, 1998] of Niño 3.4 SST and is shown in Figure 20. In this analy-
 822 sis, the *piControl* simulation is subdivided into five 100-year long sections (as in Steven-
 823 son [2012]). For reference, we also include the spectrum from HadISST [Rayner et al.,
 824 2003] and ERSSTv4 [Huang et al., 2015]. In both the *piControl* and historical ensemble,
 825 E3SMv1 ENSO variability is strong with statistically significant peaks near a three-year
 826 period. We also see a signature of longer term modulation of ENSO variability (6-9 year
 827 period) in the *piControl*. This is consistent with other CMIP models [e.g., Stevenson et al.,
 828 2012; Stevenson, 2012; Wittenberg, 2009] and proxy-data [e.g., McGregor et al., 2013].
 829 However, the modulation is a bit weaker than other models. Relative to the CESMv1
 Large Ensemble [Kay et al., 2015], Figure 20b, where the CESMv1 PI control is subdi-



836 **Figure 20.** ENSO (Nino3.4) variability of the PI control simulation and historical ensemble. The Morlet
 837 wavelet of degree six is used [e.g., Torrence and Compo, 1998] and the maximum and minimum wavelet
 838 power is used for each period. The PI control is subdivided into 100 year intervals to form a five member
 839 ensemble [e.g., Stevenson, 2012]. The solid line represents the mean wavelet power for the ensemble. The
 840 shading bounds the maximum and minimum power that is above the 90% significance threshold. The black
 841 and green lines are two observational data products. (a) E3SMv1 and (b) CESM-LE.

830 provided as in E3SMv1, E3SMv1 variability is slightly closer to observations, but is shifted
 831 strongly to a three year period, whereas CESM-LE is dominant at approximately 4.5 years.
 832 In both models, the dominant ENSO peak remains consistent between the PI control and
 833 historical ensembles. We also note that the spread of ENSO variability in the E3SMv1
 834 historical simulation is much smaller than in the CESM-LE. This is likely due to the small
 835 E3SMv1 ensemble size relative to the 39 member CESM-LE [Newman et al., 2018].



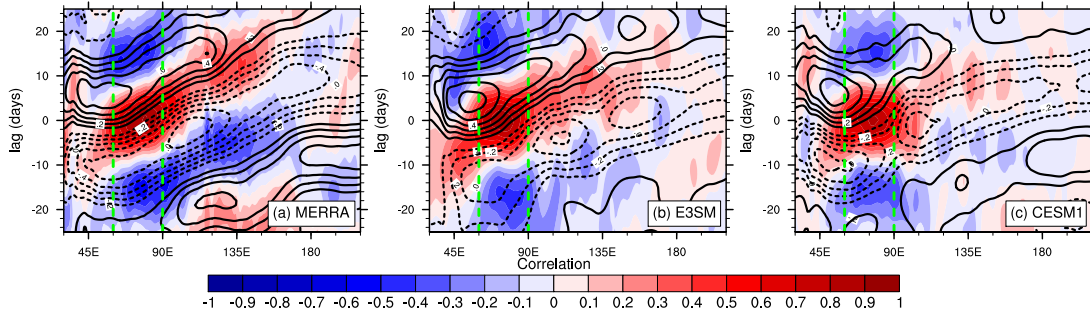
842 **Figure 21.** (a) - (c) Difference of composite El Niño events and composite La Niña events (1950-2015)
 843 for the HadleyISST dataset, the E3SMv1 historical ensemble, the CESM-LE respectively. El Niño events are
 844 defined as periods when the Niño 3.4 SST anomaly exceeds 0.8°C for more than six consecutive months. The
 845 La Niña criterion is Niño 3.4 SST anomaly less than -0.8°C for more than six months (these definitions are
 846 consistent with *Menary et al.* [2018]). When an ENSO event is identified, the SST is averaged from Novem-
 847 ber - March. For model output, every ensemble member contributes to the mean composite. The Pearson
 848 correlation coefficient and RMSE are shown for E3SMv1 and CESM in (b) and (c).

849 A number of CMIP5 models do not well represent the spatial pattern of ENSO vari-
850 ability, in particular its westward extent [*e.g.*, *Van Oldenborgh et al.*, 2005; *Menary et al.*,
851 2018]. Figure 21 shows the difference of a composite of the strongest El Niño events and
852 strongest La Niña events (following [*Menary et al.*, 2018]). Broadly, E3SMv1 reproduces
853 the spatial pattern seen in observations well with a few notable differences. E3SMv1 does
854 not capture the signal along the North American coast, suggesting a bias in the coastally
855 trapped kelvin waves. The westward extent is also larger than observed, but better than
856 seen in CESM-LE. Overall, the comparison with observations is good, with low RMSE
857 and a high correlation coefficient.

858 On sub-seasonal timescales the dominant mode of variability and predictability in
859 the tropics is the Madden Julian Oscillation [*MJO Waliser et al.*, 2003]. The MJO is gen-
860 erally thought to play a role in ENSO initiation [*McPhaden et al.*, 2006], monsoon ac-
861 tive break cycles [*Annamalai and Slingo*, 2001], tropical cyclogenesis [*Sobel and Maloney*,
862 2000] and remote teleconnection effects [*Vitart*, 2017], therefore its accurate simulation
863 is key. The simulation of the MJO in E3SMv1 represents a significant improvement in
864 strength, propagation characteristics and the explained intra-seasonal variance compared to
865 CESM1 (Fig. 22). However, significant biases in Pacific propagation remain. In CESM1,
866 the intra-seasonal propagation is in the wrong direction, westward, in the Indian Ocean.
867 Weak correlations do make it over the Maritime Continent and into the West Pacific, but
868 they are mostly decoupled from the SST signal. In contrast, E3SMv1 has consistent low-
869 level wind propagation coupled in quadrature with the SSTs from the Indian Ocean to the
870 Central Pacific.

875 5.6 Temperature evolution over the historical record

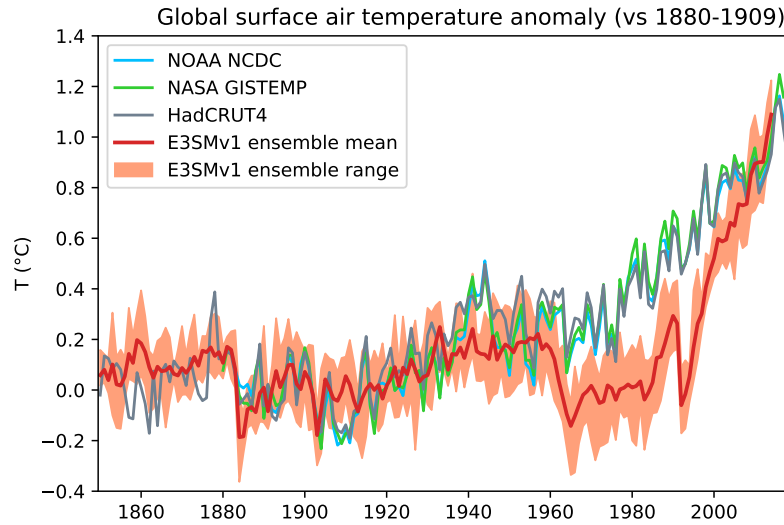
876 The 1850-2014 time evolution of the global mean surface air temperature anomaly
877 from the E3SMv1 historical ensemble is compared against three observational products
878 in Figure 23. Observations include NOAA NCDC [*Smith et al.*, 2008; *Zhang et al.*, 2015],
879 NASA GISTEMP [*Hansen et al.*, 2010; *GISTEMP Team*, 2018], and HadCRUT4 [*Morice*
880 *et al.*, 2012] which are in good agreement with each other. Anomalies are computed with
881 respect to 1880-1909, the earliest 30-year period when data is available from all observa-
882 tional products. The E3SMv1 historical ensemble mean is shown in red and the ensemble
883 minimum and maximum in orange shading. While E3SMv1 captures the bulk of the ob-
884 served warming between the 1850s and 2010s, the trajectory of the warming is at times



871 **Figure 22.** Tropical lag correlation (averaged 10°N - 10°S) of precipitation (colors) and 850-mb zonal wind
 872 (lines) with precipitation in the Indian Ocean region (60°N - 90°E ; shown by the vertical dashed green lines)
 873 for (a) Observed (TRMM for precipitation MERRA for winds), (b) E3SMv1 pre-industrial control and (c)
 874 CESM pre-industrial control. Data used are daily anomalies and band-pass filtered between 20-100 days

885 inconsistent with observations. The ensemble overlaps with observations until the 1950s,
 886 but in the subsequent decades, E3SMv1 departs from observations, first remaining too
 887 cold for several decades before warming up too rapidly starting around year 2000. The
 888 low anomalies in E3SMv1 before 1960 result from a compensation between a downward
 889 trend in the Northern Hemisphere and a positive trend in the Southern Hemisphere (not
 890 shown).

891 We turn to a regional analysis to help elucidate this inconsistency. In Figure 24 we
 892 decompose the SST anomalies into two regions: the Northern and Southern Hemisphere.
 893 The E3SMv1 ensemble range is shaded (blue) and the CESMv1 Large Ensemble (CESM-
 894 LE; [Kay *et al.*, 2015]) mean and range is also plotted (gray). CESM-LE can be regarded
 895 as a proxy for E3SMv0. Model results are compared to observations from the Hadley-
 896 NOAA/OI merged data product [Hurrell *et al.*, 2008] (red). Note here that, unlike in Fig-
 897 ure 23, the anomalies are computed relative to the 1920-1950 period as the CESM data
 898 begins at 1920. In the southern basins, E3SMv1 and CESM represent the evolution of ob-
 899 served SST well, where the E3SMv1 SST is slightly closer to data. However, in the north-
 900 ern hemisphere, the E3SMv1 SST anomalies decrease in the 1950s and warm quickly af-
 901 ter the 1990s, similar to what is observed in the surface air temperature (Figure 29). Fur-
 902 ther, the warming effect is strongest in the North Atlantic. This pattern is not seen in the
 903 data. We discuss possible causes in the next section.



904 **Figure 23.** Time evolution of annual global mean surface air temperature anomalies (with respect to 1880-
 905 1909). Comparison between observations from NOAA NCDC (blue), NASA GISTEMP (green), HadCRUT4
 906 (grey) and E3SMv1 ensemble mean and range (red and orange).

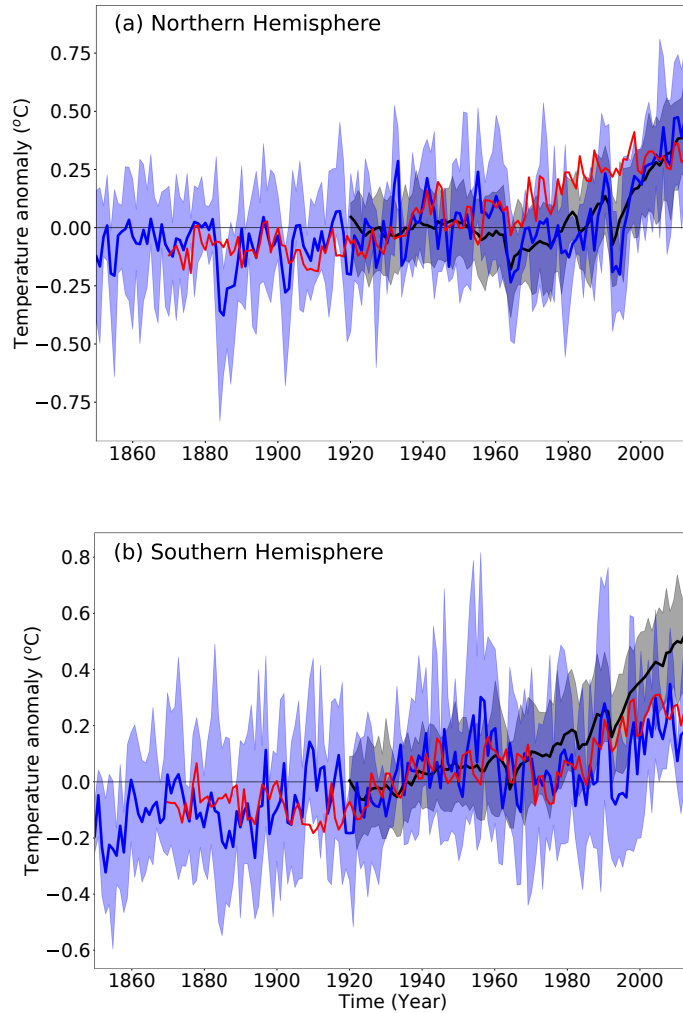
912 6 Radiative forcings and sensitivity

913 We aim to better understand E3SMv1 simulated warming over the historical record
 914 by analyzing the model radiative forcings and sensitivity.

915 6.1 Effective radiative forcing

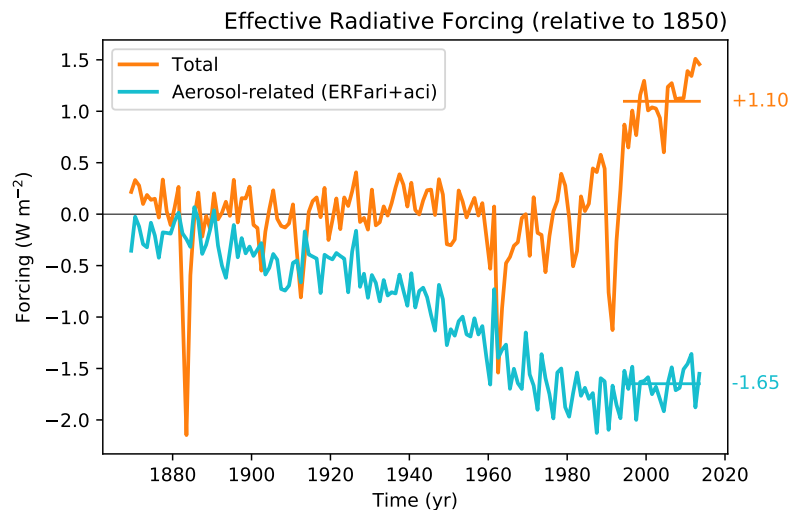
916 Effective radiative forcing (ERF) is the change in net downward TOA radiation due
 917 to observed changes in forcing agents (such as aerosol and greenhouse gases) but not sea
 918 surface temperature (allowing for adjustments in atmospheric temperatures, water vapor
 919 and clouds) [IPCC, 2013, p. 665]. A common approach to estimate ERF is by differ-
 920 encing net TOA fluxes in a pair of atmosphere-only simulations with identical sea sur-
 921 face temperatures and sea ice concentrations but different radiative forcings [e.g Hansen,
 922 2005]. Alternatively, ERF can also be estimated from regression-based approaches [Gre-
 923 gory *et al.*, 2004]. Forster *et al.* [2016] contrast several methodologies to compute ERF
 924 and generally recommend fixed SST methodologies.

925 Total ERF measures the combined effects of GHGs, short-lived gases, aerosols (in-
 926 cluding interactions with clouds), volcanoes, and land-use and land cover changes (LULCC).
 927 We estimate the transient total ERF over the historical record relative to 1850 from pairs



907 **Figure 24.** Regionally averaged SST anomalies from 1850-2014 for E3SMv1 historical ensemble (blue),
 908 Hadley-NOAA/OI merged SST product (red), and the CESMv1 large ensemble *Kay et al.* [2015] (gray). For
 909 the ensemble data, the ensemble maximum and minimum is shown. All anomalies are relative to (1920-
 910 1950), which is the beginning of the CESMv1 record. (a) Northern Hemisphere and (b) Southern Hemi-
 911 sphere.

928 of atmosphere-only simulations. We use the DECK AMIP simulations as reference (*amip_An*)
 929 and perform one additional type of simulation with identical sea surface boundary con-
 930 ditions but with all forcing agents held back at their 1850 values (*amip_1850allF_An*).
 931 This methodology is referred to as *ERF_trans* in the *Forster et al. [2016]* nomenclature.
 932 Separately, we also estimate aerosol-related ERF ($ERF_{ari+aci}$) which measures total an-
 933 thropogenic aerosol effects (including aerosol-radiation interactions, aerosol-cloud interac-
 934 tions, and the effect of light-absorbing particles in snow/ice). This requires a third type of
 935 simulation in which only aerosols and their precursors are held back at their 1850 values
 936 (*amip_1850aeroF_An*). To reduce year-to-year noise in the ERF estimates, we rely on an
 937 ensemble of three members.



938 **Figure 25.** Time evolution of annual global mean total ERF (orange) and aerosol-related ERF (light blue)
 939 derived from three ensemble members.

940 Results are shown in Fig. 25. Aerosol-related ERF (blue) increases in magnitude
 941 from the 1920s to the 1970s as anthropogenic aerosol emissions increase. Improved emis-
 942 sions standards in the 1970s cause aerosol-related ERF to stabilize at a value of -1.65 W m^{-2}
 943 m^{-2} during the last 20 years (1995-2014). This is substantially larger in magnitude than
 944 the IPCC AR5 expert judgment best estimate of -0.9 W m^{-2} [IPCC, 2013, p. 620]. It falls
 945 within the 5 to 95% uncertainty range of -1.9 to -0.1 W m^{-2} , but outside the likely range
 946 of -1.5 to -0.4 W m^{-2} . Strong negative aerosol-related forcing counterbalances warming
 947 from greenhouse gases in E3SMv1, with the net effect that total ERF (orange) hovers near
 948 zero for the first 100 years (except during volcanic eruptions) and only starts to signif-

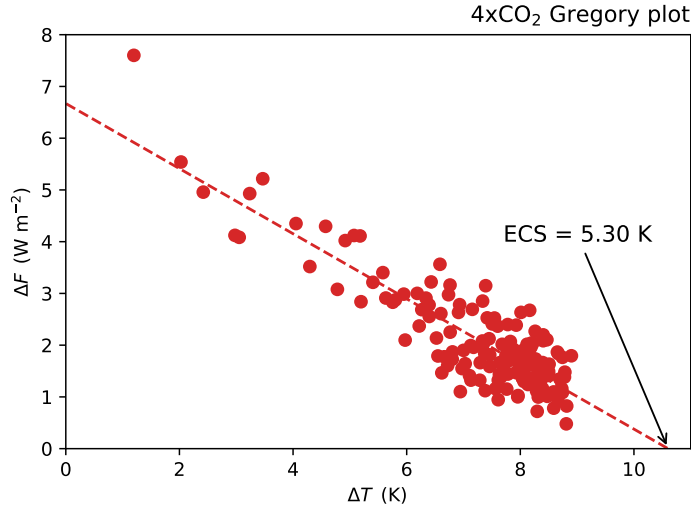
949 icantly become positive once aerosol emissions stabilize in the 1970s. Over the last 20
950 years, the total forcing averages $+1.10 \text{ W m}^{-2}$. The IPCC AR5 expert judgment best es-
951 timate of the total anthropogenic ERF between 1750 and 2011 is $+2.3 \text{ W m}^{-2}$ with an
952 uncertainty range of 1.1 to 3.3 W m^{-2} [IPCC, 2013, p. 696].

953 We note that the treatment of planetary boundary layer turbulence, shallow convec-
954 tion, and cloud macrophysics has been unified by a single parameterization (CLUBB) in
955 E3SMv1. As a result, the aerosol-related ERF now includes a contribution from interac-
956 tions between aerosols and shallow cumulus clouds. In contrast, models whose shallow
957 cumulus is unresponsive to aerosols may be artificially setting part of the aerosol-related
958 ERF to zero.

959 6.2 Sensitivity

960 The DECK simulations include two idealized CO_2 -forcing simulations designed to
961 estimate climate sensitivity at different time horizons. The equilibrium climate sensitivity
962 (ECS) is defined as the equilibrium surface temperature change resulting from a doubling
963 in CO_2 concentrations [e.g. IPCC, 2007]. While it would take thousands of simulated
964 years to run a GCM to equilibrium, ECS is typically approximated from much shorter
965 simulations using the approach of Gregory *et al.* [2004]. This approach takes advantage of
966 the fact that while the responses of global mean surface temperature and TOA energy im-
967 balance to abruptly quadrupling CO_2 are nonlinear in time, the relationship between these
968 variables is usually linear. As a result, ECS can be extrapolated as the surface tempera-
969 ture change associated with zero TOA energy imbalance. This is typically done using 150
970 year “*abrupt-4xCO2*” simulations as demonstrated for E3SMv1 in Figure 26. Because the
971 surface temperature versus TOA energy imbalance slope weakens with time in most mod-
972 els [Armour *et al.*, 2013; Andrews *et al.*, 2015; Ceppi and Gregory, 2017], ECS computed
973 this way is best described as “effective climate sensitivity”. Figure 26 shows that E3SMv1
974 response to abrupt CO_2 quadrupling is relatively linear and produces an ECS of 5.30 K.

978 Figure 27 illustrates the time evolution of annual-average surface air temperature
979 from the E3SMv1 *abrupt4xCO2* simulation (red). Clearly, 150 years is insufficient to
980 achieve equilibrium. This graphic also includes results from another idealized experiment
981 where CO_2 concentration increases by 1% per year (*1pctCO2*; blue). This second simula-
982 tion is useful for computing a shorter-timescale measure of warming called the transient

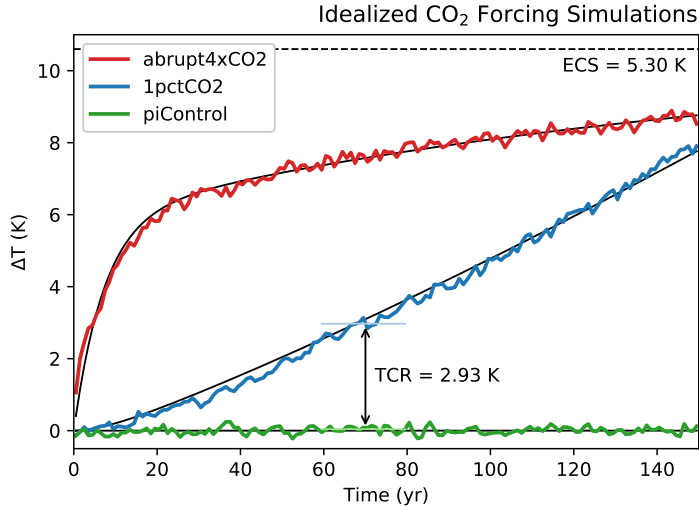


975 **Figure 26.** Global annual mean of surface air temperature change vs. TOA net radiation change for 150
 976 years of *abrupt-4xCO₂* relative to the underlying control simulation. Linear regression is depicted with a
 977 dashed line. Its intersection with the horizontal axis is twice the ECS.

983 climate response (TCR). TCR is defined as the change in surface temperature averaged for
 984 a 20-year period around the time of CO₂ doubling [approximately year 70; *IPCC*, 2007,
 985 p. 629] from a *IpctCO₂* simulation. As such, it depends on both climate sensitivity and
 986 ocean heat uptake rate. For E3SMv1, TCR is 2.93 K from *IpctCO₂*.

991 Both TCR and ECS are on the high side of a compilation of published values [*Knutti*
 992 *et al.*, 2017]. In particular, IPCC AR5 WG1 estimates that ECS is likely (> 66% prob-
 993 ability) between 1.5 and 4.5 K, while TCR is likely between 1 and 2.5 K *IPCC* [2013,
 994 p. 871]. E3SMv1 ECS and TCR are 17% larger than the likely upper bound from IPCC.
 995 While large, these values are below the extremely unlikely (< 5%) upper bounds of 3 K
 996 for TCR and 6 K for ECS.

997 To better understand E3SMv1's high ECS relative to models that took part in CMIP5,
 998 we diagnose 2xCO₂ effective radiative forcing (ERF_{2xCO₂}) and individual radiative feed-
 999 backs from the *abrupt-4xCO₂* simulation. In this particular case, ERF_{2xCO₂} is derived
 1000 by linear regression (a methodology referred to as *ERF_{reg}* in the *Forster et al.* [2016]
 1001 nomenclature). In Figure 28, we show ERFs and radiative feedbacks from 28 CMIP5
 1002 models diagnosed in *Caldwell et al.* [2016], along with those diagnosed following the
 1003 same procedure in E3SMv1. E3SMv1's effective radiative forcing, along with its Planck,

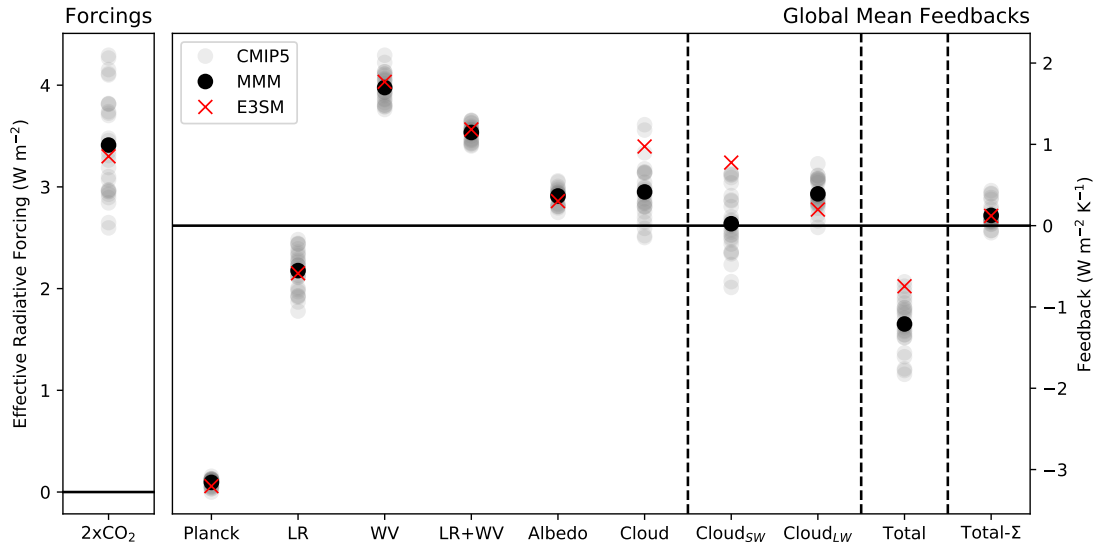


987 **Figure 27.** Time evolution of annual global mean air surface temperature anomalies for the idealized CO₂
 988 forcing simulations *abrupt-4xCO2* (red), *1pctCO2* (blue) and the control simulation (*piControl*; green). Solid
 989 lines are fits obtained with a two-layer energy balance model (discussed in sub-section 6.3). Also depicted are
 990 estimates of ECS and TCR.

1004 lapse rate, water vapor, combined lapse rate plus water vapor, and surface albedo feed-
 1005 backs are all very close to the CMIP5 multi-model mean values. In contrast, E3SMv1's
 1006 net cloud feedback is larger than in all but two CMIP5 models. Although its positive LW
 1007 cloud feedback is slightly smaller than the CMIP5 average, its positive SW cloud feed-
 1008 back is larger than all CMIP5 models. Therefore, E3SMv1's high climate sensitivity
 1009 is solely due to its large positive cloud feedback, which causes its net feedback param-
 1010 eter (which quantifies how strongly the 4xCO₂ forcing is radiatively damped) to be less
 1011 negative than all but two CMIP5 models ("Total" column of Figure 28). A more detailed
 1012 diagnosis of the reasons for E3SMv1's large positive cloud feedback will be reported in a
 1013 subsequent paper.

1026 6.3 Two-layer energy balance model

1027 Having established that E3SMv1 is a high-sensitivity model with a strong aerosol
 1028 forcing, we now explore the degree to which either the sensitivity or the aerosol forcing
 1029 can explain the mismatch in the warming trajectory between E3SMv1 and observations
 1030 (Section 5.6 and Figure 23).



1014 **Figure 28.** Global and annual mean effective radiative forcings (ERFs) and radiative feedbacks derived as
 1015 the y-intercepts and slopes, respectively, of the regression line between TOA radiation anomalies and global
 1016 surface temperature anomalies from 150-year *abrupt-4xCO₂* experiments. ERFs are divided by 2 to express
 1017 them with respect to a doubling of CO₂. The total radiative feedback (Total) is broken down into Planck, lapse
 1018 rate (LR), water vapor (WV), combined lapse rate plus water vapor (LR+WV), surface albedo (Albedo), and
 1019 net cloud (Cloud) components using radiative kernels of *Soden et al.* [2008]. The cloud feedback is further
 1020 broken down into its shortwave (Cloud_{SW}) and longwave (Cloud_{LW}) components. Individual CMIP5 models
 1021 are shown in gray circles, the CMIP5 multi-model mean (MMM) is shown as a black circle, and E3SMv1
 1022 is shown as a red cross. ‘Total’ refers to the net radiative feedback computed directly from TOA fluxes.
 1023 ‘Total-Σ’ refers to the difference between the directly-calculated net feedback and that estimated by summing
 1024 kernel-derived components. This value is near zero for E3SMv1, indicating that errors in the overall radiative
 1025 kernel decomposition of feedbacks are small.

1031 *Stevens* [2015] uses the historical record to constrain the aerosol forcing. He argues
 1032 “that an aerosol forcing less than -1.0 W m^{-2} is very unlikely [because] a more negative
 1033 aerosol forcing would imply that none of the roughly 0.3-K rise in Northern Hemisphere
 1034 surface temperatures during the 100-yr period from 1850 to 1950 could be attributed to
 1035 anthropogenic forcing, which seems implausible.”

1036 Along similar lines of reasoning, *Zhao et al.* [2018] caution against the often seen
 1037 “argument that the twentieth century warming does not strongly constrain either climate
 1038 sensitivity or the strength of aerosol cooling because similar overall warming can result
 1039 from relatively low sensitivity to CO_2 and weak aerosol cooling, or by high sensitivity
 1040 and strong aerosol cooling.” In their words, this argument holds “only to a limited ex-
 1041 tent, because of the likelihood of there having been a peak, or at least a plateau, in aerosol
 1042 forcing in the 1980–1990s. As a result, in order to create the correct overall warming if
 1043 climate sensitivity is high, one requires large enough aerosol forcing to cancel much of the
 1044 warming prior to the 1980s, while after the aerosols peak the high sensitivity and reduc-
 1045 tions in aerosols combine to produce very rapid warming.”

1046 *Held et al.* [2010] demonstrated that the time evolution of GFDL CM2.1 global
 1047 mean warming could be approximated quite realistically using a simple two-layer box
 1048 model driven by a time evolving net radiative forcing. Similar approaches have also been
 1049 used to estimate the climate sensitivity from global mean temperature observations [e.g
 1050 *Padilla et al.*, 2011; *Aldrin et al.*, 2012]. Within the framework of a two-layer energy-
 1051 balance model (EBM), *Geoffroy et al.* [2013] derived analytical solutions for the evolu-
 1052 tion of the global surface temperature in response to idealized forcing scenarios *abrupt-*
 1053 *4xCO2* and *1pctCO2*. Furthermore, using 16 AOGCMs from CMIP5, they demonstrated
 1054 that EBMs calibrated exclusively with *abrupt-4xCO2* data could accurately predict the
 1055 temperature evolution in *1pctCO2* simulations.

1056 The two-layer EBM is defined by the following system of equations:

$$C \frac{dT}{dt} = \mathcal{F} - \lambda T - \gamma(T - T_0) \quad (3)$$

$$C_0 \frac{dT_0}{dt} = \gamma(T - T_0) \quad (4)$$

1057 Prognostic variables are the temperatures of the upper (T) and deep ocean layers (T_0) with
 1058 C and C_0 their respective heat capacities. \mathcal{F} is the total radiative forcing, λ the surface
 1059 feedback parameter and γ the heat exchange coefficient between the upper and deep ocean.

1060 The general solution for the upper (surface) temperature under a time varying forcing is
 1061 given by Eq. (B8) in *Geoffroy et al.* [2013]:

$$T(t) = \frac{a_f}{\lambda\tau_f} \int_0^t \mathcal{F}(t') e^{-(t-t')/\tau_f} dt' + \frac{a_s}{\lambda\tau_s} \int_0^t \mathcal{F}(t') e^{-(t-t')/\tau_s} dt' \quad (5)$$

1062 The solution can be interpreted as the sum of a fast and slow convolution of the forcing
 1063 with exponential decay functions. Table 1 in *Geoffroy et al.* [2013] lists the relationships
 1064 between the weights (a_f , a_s), time scales (τ_f , τ_s), and the parameters characterizing the
 1065 model (Eqs 3 and 4).

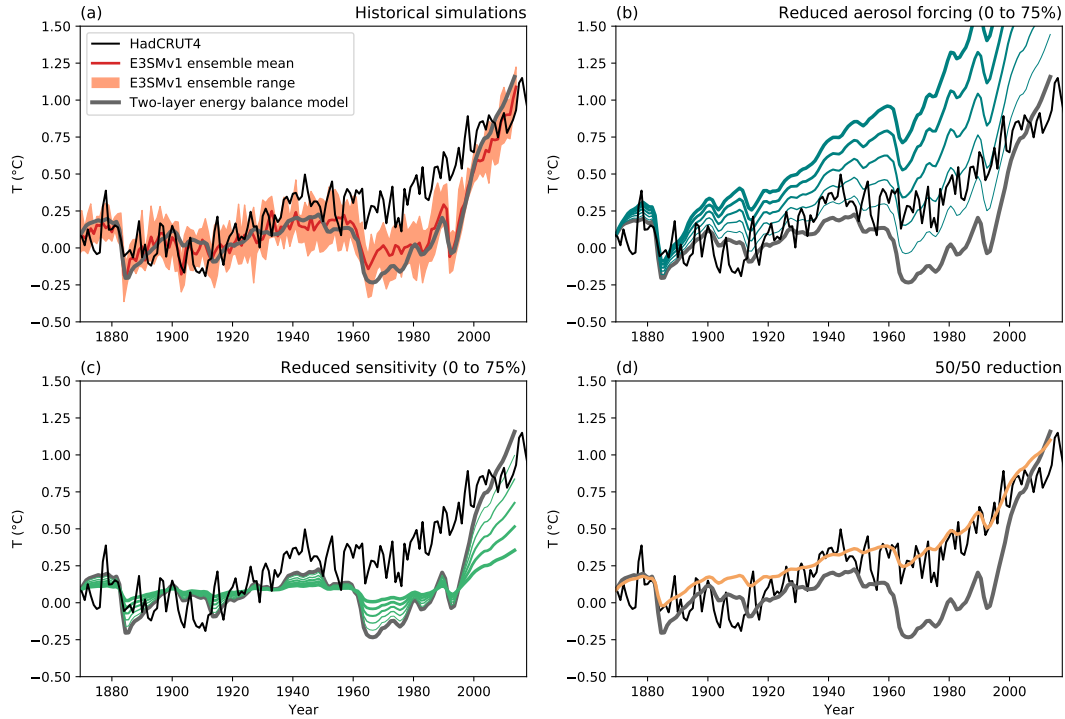
1066 We calibrate the EBM for E3SMv1 using *abrupt-4xCO2* data following the proce-
 1067 dure outlined in *Geoffroy et al.* [2013], except for a small deviation in the second step (p.
 1068 1846) when we average over 20 instead of 10 years to estimate the fast time scale as we
 1069 found that it provides a better fit to *abrupt-4xCO2*. Table 3 lists the parameter values and
 1070 the fits to *abrupt-4xCO2* and *1pctCO2* are shown with solid black lines in Fig. 27. The
 1071 fits clearly demonstrate that the EBM calibrated with *abrupt-4xCO2* can accurately predict
 1072 the behavior of *1pctCO2*. Furthermore, the TCR from the fit (3.07 K) is within 5% of its
 1073 true value (2.93 K), indicating that for E3SMv1, both ECS and TCR could be estimated
 1074 from *abrupt-4xCO2* alone.

1075 **Table 3.** Two-layer EBM parameters calibrated from *abrupt-4xCO2* simulation.

\mathcal{F}	6.671	W m ⁻²
λ	0.629	W m ⁻² K ⁻¹
a_f	0.558	unitless
τ_f	7.263	year
a_s	0.442	unitless
τ_s	160.093	year

1076 We now explore the EBM's ability to reproduce the predicted E3SMv1 warming
 1077 over the historical record following the approach of *Held et al.* [2010]. The time varying
 1078 forcing $\mathcal{F}(t)$ in Eq. (5) is simply the total ERF from Fig. 25 (orange line). The resulting
 1079 temperature is shown with a thick gray line in Fig. 29a. The simple model predictions
 1080 agree much better with modeled rather than observed values, lying within the envelope of
 1081 model ensemble values most of the time. Furthermore, the EBM reproduces – albeit in an

1082 exaggerated fashion – the E3SMv1 behavior of a lack of warming during the 1960-1990s
 1083 followed by an excessive warming trend.



1084 **Figure 29.** Time evolution of annual global mean surface air temperature anomalies. (a): Observations
 1085 (HadCRUT4; black), E3SMv1 ensemble mean and range (red and orange), two-layer energy balance model
 1086 (gray). (b): EBM with reduced aerosol forcing (0% to 75% in 15% increments; blue with increasing line
 1087 thicknesses). (c): EBM with reduced sensitivity (0% to 75%; green with increasing line thicknesses). (d):
 1088 EBM with 50% reduction in both aerosol forcing and sensitivity (brown).

1089 With the credibility of the EBM established, the EBM can also be applied to hypo-
 1090 theoretical scenarios in order to explore the effect of the forcing and sensitivity. New hypo-
 1091 theoretical total forcing with weaker aerosol forcing can be constructed by linear combina-
 1092 tions of the original forcings:

$$\mathcal{F}_{\text{new}} = \mathcal{F}_{\text{tot}} - \alpha_{\text{aero}} \mathcal{F}_{\text{aero}} \quad (6)$$

1093 where \mathcal{F}_{tot} is the orange line and $\mathcal{F}_{\text{aero}}$ the blue line in Fig. 25. Figure 29b explores re-
 1094 ductions in aerosol forcing up to 75% in 15% increments (α_{aero} from 0.0 to 0.75). A 50%
 1095 reduction corresponds approximately to the median IPCC AR5 value of -0.9 W m^{-2} .

1096 Similarly, we can keep the forcing at its original value and reduce the sensitivity

$$\lambda_{\text{new}} = \frac{\lambda}{\alpha_{\text{sensitivity}}} \quad (7)$$

1097 as illustrated in Fig. 29c for a reduction up to 75% ($\alpha_{\text{sensitivity}}$ from 1.0 to 0.25) with no
1098 changes in the fast and slow time scales (τ_f , τ_s) and their corresponding weights (a_f , a_s).

1099 Careful visual inspection of Fig. 29b,c reveals that neither a reduction in aerosol
1100 forcing nor a reduction in sensitivity alone is sufficient to improve the match with the
1101 historical temperature record. Reducing aerosol forcing alone can improve the match up
1102 to the 1980s but not afterwards when the aerosol forcing reaches a plateau and the high
1103 sensitivity causes an excessive warming trend (confirming the argument of *Zhao et al.*
1104 [2018]). Instead, a substantial reduction in both is needed. For example, a 50% reduction
1105 in aerosol forcing and sensitivity (Fig. 29d; brown) matches observations (black) much
1106 better than the original EBM calibrated with E3SMv1 (gray). To the extent that the EBM
1107 is a good proxy for the behavior of the full model, we conclude that improving the tra-
1108 jectory of the historical warming of E3SMv1 would require a substantial reduction in the
1109 magnitude of both aerosol forcing and sensitivity.

1110 7 Conclusion

1111 In this paper, we have described the new E3SMv1 fully-coupled physical model in
1112 its standard resolution configuration. This model is designed to serve as a tool to address
1113 DOE mission-relevant water cycle questions. We have examined the E3SMv1 simulated
1114 climate with a set of experiments from the CMIP6 DECK. Key behaviors and biases are
1115 recapitulated below:

- 1116 • Over the course of the long pre-industrial control simulation, the coupled system
1117 has very little model drift, as evidenced in the net TOA flux, global mean surface
1118 air temperature, and seasonal range in sea ice area (Figure 2).
- 1119 • The present-day climate simulated by an ensemble of historical simulations reveals
1120 that the atmosphere is credible compared to an ensemble of CMIP5 models (Fig-
1121 ure 9) but also subject to biases common to many models, e.g.,
 - 1122 – stratocumulus coverage (Figure 5),
 - 1123 – double ITCZ (Figure 6).

- 1124 • Postively, the atmosphere has a much improved representation of the MJO, in re-
1125 gards to its strength and propagation characteristics (Figure 22).
- 1126 • The ocean also shows biases consistent with lower resolution ocean models, e.g.,
 - 1127 – Gulf stream separation (evident in Figure 10),
 - 1128 – a shallow MLD bias in the SH (Figure 12).
- 1129 • AMOC is weak (Figure 13) and there is large fresh water bias in the North Atlantic
1130 (Figure 11) and accompanying shallow MLD biases (Fig. 12). These biases are
1131 certainly related and the subject of ongoing research with E3SMv1.
- 1132 • The simulated ENSO variability is realistic. It is closer to observations than CESM1
1133 (Figure 20).
- 1134 • E3SM well simulates the spatial pattern associated with ENSO events and is closer
1135 to observations than CESM1 (Figure 21)
- 1136 • Sea ice concentrations are too high in the Labrador sea (Figure 14) and the sea-
1137 sonal growth of ice is delayed and too rapid, relative to observations.
- 1138 • Streamflow simulated by E3SMv1 is consistent with observations in magnitude and
1139 timing, however, the seasonality is too large in a number of regions (Figure 19).
- 1140 • E3SMv1's aerosol-related effective radiative forcing ($ERF_{ari+aci} = -1.65 \text{ W m}^{-2}$),
1141 equilibrium climate sensitivity ($ECS = 5.3 \text{ K}$) and transient climate response (TCR
1142 $= 2.93 \text{ K}$) are larger in magnitude than most CMIP5 models, but fall within previ-
1143 ously published uncertainty bounds. Predictions of large future warming are due to
1144 unusually large positive shortwave cloud feedback.
- 1145 • The coupled climate in the historical ensemble doesn't warm as quickly as observa-
1146 tions between 1960 and 1990, but warms more rapidly thereafter, with an end result
1147 that the E3SMv1 ensemble approaches observations by 2014 (Figure 23).
- 1148 • An analysis with a simple energy balance model reveals that this mismatch is due
1149 to the combination of E3SMv1's strong aerosol-related forcing and high climate
1150 sensitivity (Figure 29).

1151 The climate simulated by E3SMv1 has biases broadly consistent with other climate
1152 class models, but also has improvements in certain regimes (e.g., tropical variability).
1153 These simulations and analysis help to establish the scientific credibility of this new model
1154 and set the stage for future additional analysis of these existing simulations as well as new

1155 simulations with E3SMv1 (e.g., high resolution coupled, future projections, and regionally
1156 refined).

1157 **A: Energy correction term**

1158 The definition of energy is slightly inconsistent between components of E3SMv1, in
1159 particular with respect to the treatment of internal water energy. MPAS-Ocean and sea ice
1160 properly account for energy changes due to water temperature changes, but other compo-
1161 nents currently do not. This inconsistency creates a small spurious energy imbalance be-
1162 tween the atmosphere and ocean due to the fact that water evaporates from the ocean sur-
1163 face at a certain temperature and returns to the ocean as precipitation at a different (lower)
1164 temperature. Globally averaged, the imbalance is less than 0.5 W m^{-2} .

1165 Kirchhoff's equation [Glickman, 2000, p. 432] relates the variation with temperature
1166 of the latent heat of a phase change to the difference between the specific heats of the two
1167 phases. For water, this is written as:

$$\left(\frac{\partial L_v}{\partial T}\right)_p = c_{pv} - c_w \quad (\text{A.1})$$

1168 where L_v is the latent heat of vaporization, c_{pv} is the specific heat at constant pressure of
1169 water vapor, and c_w is the specific heat of liquid water. Like other Earth system models
1170 (Isaac Held, personal communication), the E3SMv1 atmosphere component violates Kirch-
1171 hoff's equation by neglecting variation of L_v with temperature, while at the same time
1172 assuming $c_{pv} \neq c_w$.

1173 The ocean experiences a net cooling because it supplies heat to bring the liquid pre-
1174 cipitation to the temperature of the sea surface, while solid precipitation is first melted
1175 (using heat from the ocean) and then brought to the temperature of the sea surface. How-
1176 ever, there is no corresponding warming term in the atmosphere. The net impact is an off-
1177 set between long term trends in net TOA energy flux and ocean heat content (the largest
1178 heat reservoir of the coupled system).

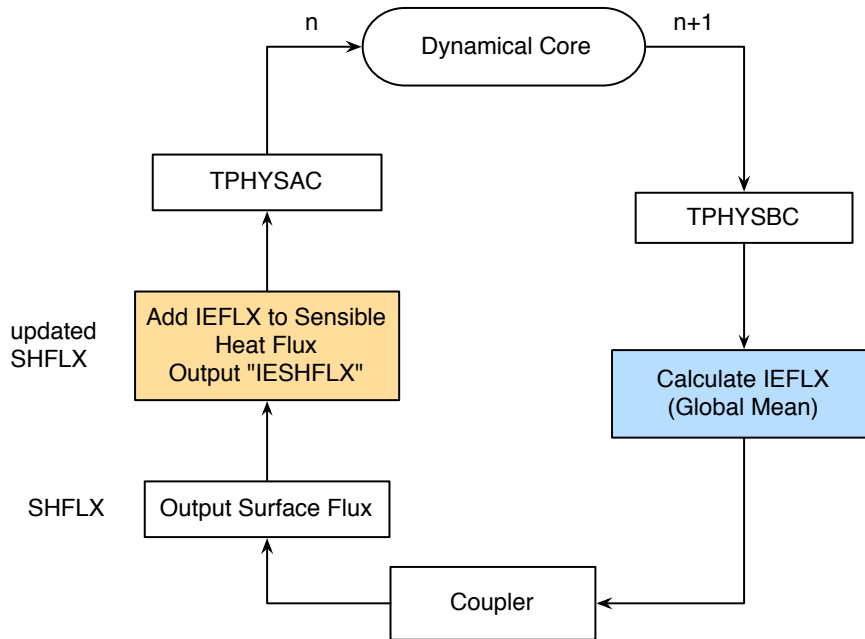
1179 We correct for that imbalance by adding back the missing warming term in the at-
1180 mosphere with an *ad hoc* correction term. To mimic what is done in the ocean, an energy
1181 flux term (IEFLX) is introduced:

$$\text{IEFLX}_i = c_{p_{sw}} * \text{QFLX}_i * T_{surf,i} - c_{p_{sw}} * \text{PRECT}_i * T_{surf,i} \quad (\text{A.2})$$

1182 where i denotes the grid column, cp_{sw} the heat capacity of sea water, QFLX the surface
 1183 moisture flux, PRECT the precipitation flux, and T_{surf} the surface (skin) temperature.
 1184 IEFLX is first calculated for each grid box and then globally averaged and later applied
 1185 as a uniform adjustment to the sensible heat flux at each grid box (Fig. A.1):

$$SHFLX_i = SHFLX_i + \frac{\sum_{i=1}^I (A_i IEFLX_i)}{\sum_{i=1}^I (A_i)}. \quad (A.3)$$

1186 where I denotes the total number of grid columns on the cubed sphere mesh, A_i the grid
 1187 cell area for column i , and SHFLX the sensible heat flux. The calculated global annual
 1188 mean IEFLX is about $+0.4 \text{ W m}^{-2}$. Sensitivity simulations show that using the IEFLX
 1189 correction improves energy conservation consistency between net TOA and ocean heat
 1190 content with minimal impact on the simulated climate.



1191 **Figure A.1.** Flowchart of the IEFLX calculation in E3SMv1. IEFLX shown in this figure is a global mean
 1192 quantity. TPHYSBC and TPHYSAC indicate the model physics parameterizations before (BC) and after (AC)
 1193 the coupler calculation. SHFLX indicates the sensible heat flux.

1194 **B: Input Data**

1195 The E3SMv1 DECK simulations generally followed the input4MIPS datasets, which
1196 are described by *Durack et al.* [2018]. The details of the versions used, and any deviations
1197 from the input4MIPS data are specified in the sections below.

1198 **B.1 Greenhouse Gas Concentrations**

1199 Greenhouse gas (GHG) concentrations were specified using v1.2.0 of input4MIPS
1200 GHG historical concentrations [*Meinshausen and Vogel, 2016; Meinshausen et al., 2017*]
1201 for CO₂, CH₄, N₂O, CFC-12, and CFC-11eq (ie CFC-11 plus all other major halocarbon
1202 species converted to give an equivalent amount of CFC-11 forcing). The data was speci-
1203 fied annually. The model assumed that the concentration applies to January 1, and inter-
1204 polated linearly during the year to the next year's value. To enable time interpolation
1205 until the end of 2014, concentrations for 2015 and 2016 were added by linear extrapola-
1206 tion from 2013 and 2014. The concentrations were assumed to be uniform throughout the
1207 atmosphere (i.e., well mixed).

1208 Ozone concentrations were as specified in sections [B.2](#) and [B.3](#).

1209 **B.2 Tropospheric Aerosols & Related Datasets**

1210 Aerosol concentrations, sizes, and optical properties were provided by the four-mode
1211 Modal Aerosol Module (MAM4, *Liu et al.* [2016]), but modified to include marine organic
1212 aerosols (MOA, *Burrows et al.* [2018]), with the coarse mode extended to include carbona-
1213 ceous aerosols (i.e., BC, POM, MOA and SOA) to treat the resuspension of aerosol parti-
1214 cles from evaporated raindrops more appropriately.

1215 Monthly anthropogenic emissions of aerosols and precursor gases for MAM4 were
1216 specified as follows. SO₂, sulfate, black carbon (BC), primary organic matter (POM), and
1217 second organic aerosol (SOA) gases, were obtained from the CMIP6 (Coupled Model In-
1218 tercomparison Project Phase 6) emission data sets, as described by *Hoesly et al.* [2017,
1219 2018]. Open fire emissions of SO₂, BC, POM and SOA were from the biomass burn-
1220 ing data sets [*van Marle et al., 2016, 2017*] developed for CMIP6. All the CMIP6 emis-
1221 sions had annual data. Biogenic emissions for SOA precursor sources (e.g., isoprene and
1222 monoterpenes) were obtained from the standard MOZART emissions as described by *Em-*

1223 *mons et al.* [2010]. The vertical distribution of SOA precursor sources was prescribed to
1224 mimic the explicit treatment of gas- and particle-phase chemical oxidation of SOA [*Shri-*
1225 *vastava et al.*, 2015]. The inject height of fire emissions, as well as industrial and power
1226 plant emissions, followed the AeroCom (Aerosol Comparisons between Observations and
1227 Models) protocols [*Dentener et al.*, 2006].

1228 The oxidants necessary to calculate secondary aerosol production were read in from
1229 a file, which included O₃, OH, H₂O₂, HO₂, and NO₃. The data was provided for each
1230 month for one year in each decade (1849, 1855, 1865, . . . , 2015). The ozone values were
1231 derived from the input4MIPS Ozone dataset v1.0 [*Hegglin et al.*, 2016]. Concentrations
1232 for the other species were not provided by input4MIPS, so the pre-existing specified-oxidant
1233 data used by MAM4 was used, for which the provenance is not entirely known, but was
1234 probably from a CAM4 based CAM-CHEM transient simulation for IPCC AR5, with an
1235 extension to 2015 by copying 2000 values.

1236 **B.3 Stratospheric Ozone**

1237 We used a prognostic linearized ozone chemistry scheme to calculate stratospheric
1238 ozone using a single tracer (Linoz v2, *Hsu and Prather* [2009]). The linearized chemistry
1239 coefficients were calculated using the GHG concentrations from v1.2.0 of the input4MIPS
1240 GHG historical concentrations with the assistance of Juno Hsu and Michael Prather (pri-
1241 vate communication), with a 3-year lag as simple way to account for the time surface
1242 concentrations take to mix into the stratosphere. The input data was generated for every
1243 month in years spaced 5 years apart (1845, 1850, . . . , 2015).

1244 **B.4 Stratospheric Aerosols**

1245 We modified E3SMv1 to pass the stratospheric aerosol optical properties (extinc-
1246 tion, single scattering albedo, asymmetry factor) from input4MIPS directly to the radiation
1247 routine (after regridding from the input file to the model grid). We used version 3 of the
1248 dataset, which was created to exactly match the E3SMv1 radiation wavelength ranges by
1249 one of the dataset creators, Beiping Luo (private communication).

1250 **B.5 Sea-Surface Temperature & Sea-Ice**

1251 For our AMIP simulations, we used v1.1.3 of the PCMDI sea-surface temperature
1252 and sea-ice fractions [Durack and Taylor, 2017; Taylor et al., 2000], which cover the years
1253 1870 to 2016. E3SM interpolated ocean temperatures and ice fractions between the mid-
1254 month values, hence we used the diddled version of the dataset in which the values were
1255 tweaked so that the monthly-mean values in the model matched the monthly-mean val-
1256 ues in the original dataset. Since the input dataset used the Gregorian calendar and E3SM
1257 uses a 365-day calendar, the model time interpolated on the Gregorian calendar, which
1258 caused Feb 29 data to be skipped and will lead to a small discontinuity at the start of
1259 March 1 during leap-years.

1260 **B.6 Land Use**

1261 Land use, and land use change, files were regrided to the E3SM ne30 grid by George
1262 Hurtt and Ritvik Sahajpal (private communication) using v2.1h of the input4MIPS land
1263 use data [Hurtt et al., 2017].

1264 **B.7 Nitrogen Deposition**

1265 Nitrogen deposition was not used by E3SM for the DECK simulations.

1266 **B.8 Solar Input**

1267 Solar irradiances came from v3.2 of the solar irradiance dataset from input4MIPS
1268 [Matthes et al., 2017a,b]. The data was specified monthly.

1269 **B.9 Orbital Parameters**

1270 Earth's orbital parameters were inadvertently fixed to 1990 values for all DECK sim-
1271 ulations.

1272 **C: Analysis tools**

1273 The analysis of this paper was largely enabled through a suite of diagnostic soft-
1274 ware packages developed in tandem with the model. These packages were designed for the

1275 E3SM development and analysis community to be usable, extensible, and with shareable
1276 results, each with a different scientific focus and goal.

1277 New and improved netCDF Operators (NCO) [Zender, 2008, 2018] were devel-
1278 oped and customized for E3SM analysis, and verified to work on CESM output. These
1279 include a climatology generator and a time-series splitter accessed through the new nc-
1280 climo operator, and a regridded accessed through the new ncremap operator. Each is a
1281 parallelized tool suitable for serial or background-parallel mode execution on personal
1282 laptops and workstations, and background-parallel and MPI-parallel operation on high-
1283 performance computing nodes. These tools are embedded or used as pre-processing steps
1284 for E3SM_Diags, MPAS-Analysis and A-PRIME. Their full documentation is at [http:](http://nco.sf.net/nco.html)
1285 [//nco.sf.net/nco.html](http://nco.sf.net/nco.html).

1286 E3SM_Diags is a modern, Python-based diagnostics package developed to facili-
1287 tate evaluating earth system models. The package includes a set of comprehensive toolk-
1288 its and updated analysis data sets. This software is designed in a flexible, modular, and
1289 object-oriented fashion, enabling users to manipulate different processes in a diagnostics
1290 workflow. Numerous configuration options for metrics computation (i.e., regridding op-
1291 tions) and visualization (i.e., graphical backend, color map, contour levels) are customiz-
1292 able. Built-in functions to generate derived variables and to select diagnostics regions are
1293 supported and can be easily expanded to accommodate earth system models with output
1294 conventions that are CMIP compliant. Modern computer technologies, such as multi-
1295 processing and containerization are applied in the software development, which enhance
1296 the performance and stability of the software. Detailed documentation can be found from
1297 <https://e3sm.org/resources/tools/diagnostic-tools/e3sm-diagnostics>.

1298 MPAS-Analysis (<https://github.com/MPAS-Dev/MPAS-Analysis>) is a Python-
1299 based tool for performing post-processed analysis and plotting of output from E3SM's
1300 ocean and sea-ice components (MPAS-Ocean and MPAS-Seaice, respectively). MPAS-
1301 Analysis uses the NetCDF Operators (NCO) to compute climatologies, extract time series
1302 and remap data sets to common reference grids. Comparisons between simulation results
1303 and a wide variety of observational data sets are supported on both latitude/longitude and
1304 Antarctic stereographic grids. MPAS-Analysis also supports comparisons between E3SM
1305 simulations, allowing users to examine the influence of changing meshes, resolution, pa-
1306 rameters, model physics, etc. Parallelism has been introduced into MPAS-Analysis by

1307 breaking each analysis task into modular subtasks that can be run in parallel to efficiently
1308 produce hundreds of plots. MPAS-Analysis is aware of E3SM namelists, meaning that
1309 tasks are automatically disabled for runs where the necessary output was not produced.
1310 The end result of running MPAS-Analysis is a website with image galleries of all plots,
1311 sorted by component and analysis type, as well as a set of NetCDF files containing the
1312 post-processed data, available for further analysis.

1313 A-PRIME is a *priority* metrics package that is designed to provide a quick,
1314 broad overview of coupled model behavior [Evans *et al.*, 2018]. The target user would ex-
1315 ecute A-PRIME on model data as a run is progressing to determine whether the model
1316 is on track to produce global level expected behavior. It provides a suite of averaged and
1317 time series behavior of the most common variables that drive radiation, dynamical, and
1318 hydrological balance. When there are sufficient simulation years available, it also provides
1319 ENSO metrics. The top-level directory of the software provides a generic script that tar-
1320 gets execution on DOE supercomputers, where E3SM simulations are currently executed
1321 and/or postprocessed. These scripts point to Python postprocessing and visualization mod-
1322 ules for multiple components in the coupled model. The ocean and sea ice modules load
1323 portions of the MPAS-Analysis diagnostics as a submodule.

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1350 The E3SM project, code, simulation configurations, model output, and tools to work
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1352 running E3SM are available at <https://e3sm.org/model/running-e3sm/e3sm-quick-start>.
1353 All model codes may be accessed on the GitHub repository at [https://github.com/
1354 E3SM-Project/E3SM](https://github.com/E3SM-Project/E3SM). Model output data are accessible directly on NERSC or through
1355 the DOE Earth System Grid Federation at [https://esgf-node.llnl.gov/projects/
1356 e3sm](https://esgf-node.llnl.gov/projects/e3sm).

1357 References

- 1358 Adcroft, A., and J.-M. Campin (2004), Rescaled height coordinates for accurate represen-
1359 tation of free-surface flows in ocean circulation models, *Ocean Modelling*, 7, 269–284,
1360 doi:10.1016/j.ocemod.2003.09.003.
- 1361 Adler, R. F., G. J. Huffman, A. Chang, R. Ferraro, P.-P. Xie, J. Janowiak, B. Rudolf,
1362 U. Schneider, S. Curtis, D. Bolvin, A. Gruber, J. Susskind, P. Arkin, and E. Nelkin
1363 (2003), The Version-2 Global Precipitation Climatology Project (GPCP) monthly pre-
1364 cipitation analysis (1979–present), *Journal of Hydrometeorology*, 4(6), 1147–1167, doi:
1365 10.1175/1525-7541(2003)004\$<\$1147:tvGPCP\$>\$2.0.CO;2.
- 1366 Aldrin, M., M. Holden, P. Guttorp, R. B. Skeie, G. Myhre, and T. K. Berntsen (2012),
1367 Bayesian estimation of climate sensitivity based on a simple climate model fitted to ob-
1368 servations of hemispheric temperatures and global ocean heat content, *Environmetrics*,
1369 23(3), 253–271, doi:10.1002/env.2140.

- 1370 Andrews, T., J. M. Gregory, and M. J. Webb (2015), The dependence of radiative forcing
1371 and feedback on evolving patterns of surface temperature change in climate models,
1372 *Journal of Climate*, 28(4), 1630–1648, doi:10.1175/jcli-d-14-00545.1.
- 1373 Annamalai, H., and J. M. Slingo (2001), Active–break cycles: diagnosis of the
1374 intraseasonal variability of the asian summer monsoon, *Climate Dynamics*, 18(1), 85–
1375 102, doi:10.1007/s003820100161.
- 1376 Arakawa, A., and V. R. Lamb (1977), Computational design of the basic dynamical pro-
1377 cesses of the UCLA General Circulation Model, in *General Circulation Models of the*
1378 *Atmosphere, Methods in Computational Physics: Advances in Research and Applications*,
1379 vol. 17, edited by J. Chang, pp. 173–265, Elsevier, doi:10.1016/B978-0-12-460817-7.
1380 50009-4.
- 1381 Armour, K. C., C. M. Bitz, and G. H. Roe (2013), Time-varying climate sensitivity from
1382 regional feedbacks, *J. Climate*, 26, 4518–4534, doi:10.1175/JCLI-D-12-00544.1.
- 1383 Bader, D., W. Collins, R. Jacob, P. J. P. Rasch, M. Taylor, P. Thornton, and D. Williams
1384 (2014), Accelerated Climate Modeling for Energy, U.S. Department of Energy. Re-
1385 trieved from [https://climatemodeling.science.energy.gov/sites/default/
1386 files/publications/acme-project-strategy-plan.pdf](https://climatemodeling.science.energy.gov/sites/default/files/publications/acme-project-strategy-plan.pdf).
- 1387 Bourke, R., and A. McLaren (1992), Contour Mapping of Arctic Basin Ice Draft and
1388 Roughness Parameters, *Journal of Geophysical Research*, 97(C11), 17,715–17,728.
- 1389 Branstetter, M. L., and D. J. Erickson III (2003), Continental runoff dynamics in the Com-
1390 munity Climate System Model 2 (CCSM2) control simulation, *Journal of Geophysical*
1391 *Research*, 108(D17), doi:10.1029/2002jd003212.
- 1392 Briegleb, B. P., and B. Light (2007), A Delta-Eddington multiple scattering parameter-
1393 ization for solar radiation in the sea ice component of the Community Climate Sys-
1394 tem Model, *Tech. Rep. NCAR/TN-472+STR*, National Center for Atmospheric Research,
1395 Boulder, Colorado USA.
- 1396 Burrows, S. M., R. Easter, X. Liu, P. L. Ma, H. Wang, S. M. Elliott, B. Singh, K. Zhang,
1397 and P. J. Rasch (2018), OCEANFILMS sea-spray organic aerosol emissions - Part 1:
1398 implementation and impacts on clouds, *Atmos. Chem. Phys. Discuss.*, pp. 1–27, doi:10.
1399 5194/acp-2018-70.
- 1400 Caldwell, P. M., M. D. Zelinka, K. E. Taylor, and K. Marvel (2016), Quantifying the
1401 sources of intermodel spread in Equilibrium Climate Sensitivity, *Journal of Climate*,
1402 29(2), 513–524, doi:10.1175/jcli-d-15-0352.1.

- 1403 Campin, J.-M., J. Marshall, and D. Ferreira (2008), Sea ice–ocean coupling using a
1404 rescaled vertical coordinate z^* , *Ocean Modelling*, *24*(1-2), 1–14.
- 1405 Cavalieri, D. J., C. L. Parkinson, P. Gloersen, and H. J. Zwally (1996), Sea ice concentra-
1406 tions from Nimbus-7 SMMR and DMSP SSM/I-SSMIS passive microwave data, version
1407 1, doi:10.5067/8GQ8LZQVL0VL.
- 1408 Ceppi, P., and J. M. Gregory (2017), Relationship of tropospheric stability to climate sen-
1409 sitivity and earth’s observed radiation budget, *Proceedings of the National Academy of*
1410 *Sciences*, *114*(50), 13,126–13,131, doi:10.1073/pnas.1714308114.
- 1411 Cess, R. D., G. L. Potter, J. P. Blanchet, G. J. Boer, S. J. Ghan, J. T. Kiehl, H. Le Treut,
1412 Z.-X. Li, X.-Z. Liang, J. F. B. Mitchell, J.-J. Morcrette, D. A. Randall, M. R. Riches,
1413 E. Roeckner, U. Schlese, A. Slingo, K. E. Taylor, W. M. Washington, R. T. Wetherald,
1414 and I. Yagai (1989), Interpretation of cloud-climate feedback as produced by 14 atmo-
1415 spheric general circulation models, *Science*, *245*(4917), 513–516, doi:10.1126/science.
1416 245.4917.513.
- 1417 Cheng, W., J. C. Chiang, and D. Zhang (2013), Atlantic meridional overturning circulation
1418 (amoc) in cmip5 models: Rcp and historical simulations, *Journal of Climate*, *26*(18),
1419 7187–7197.
- 1420 Collins, N., G. Theurich, C. Deluca, M. Suarez, A. Trayanov, V. Balaji, P. Li, W. Yang,
1421 C. Hill, and A. Da Silva (2005), Design and implementation of components in the earth
1422 system modeling framework, *The International Journal of High Performance Computing*
1423 *Applications*, *19*(3), 341–350.
- 1424 Craig, A. P., M. Vertenstein, and R. Jacob (2012), A new flexible coupler for earth system
1425 modeling developed for ccsm4 and cesm1, *The International Journal of High Perfor-*
1426 *mance Computing Applications*, *26*(1), 31–42, doi:10.1177/1094342011428141.
- 1427 Dasgupta, G. (2003), Interpolants within convex polygons: Wachpress’ shape functions, *J.*
1428 *Aerospac. Eng.*, *16*, 1–8, doi:10.1061/(ASCE)0893-1321(2003)16:1(1).
- 1429 Dee, D. P., S. M. Uppala, A. J. Simmons, P. Berrisford, P. Poli, S. Kobayashi, U. Andrae,
1430 M. A. Balmaseda, G. Balsamo, P. Bauer, P. Bechtold, A. C. M. Beljaars, L. van de
1431 Berg, J. Bidlot, N. Bormann, C. Delsol, R. Dragani, M. Fuentes, A. J. Geer, L. Haim-
1432 berger, S. B. Healy, H. Hersbach, E. V. Hólm, L. Isaksen, P. Kállberg, M. Köh-
1433 ler, M. Matricardi, A. P. McNally, B. M. Monge-Sanz, J.-J. Morcrette, B.-K. Park,
1434 C. Peubey, P. de Rosnay, C. Tavolato, J.-N. Thépaut, and F. Vitart (2011), The ERA-
1435 interim reanalysis: configuration and performance of the data assimilation system, *Quar-*

- 1436 *terly Journal of the Royal Meteorological Society*, 137(656), 553–597, doi:10.1002/qj.
1437 828.
- 1438 Dentener, F., S. Kinne, T. Bond, O. Boucher, J. Cofala, S. Generoso, P. Ginoux, S. Gong,
1439 J. J. Hoelzemann, A. Ito, L. Marelli, J. E. Penner, J. P. Putaud, C. Textor, M. Schulz,
1440 G. R. van der Werf, and J. Wilson (2006), Emissions of primary aerosol and precursor
1441 gases in the years 2000 and 1750 prescribed data-sets for AeroCom, *Atmos. Chem.*
1442 *Phys.*, 6(12), 4321–4344, doi:10.5194/acp-6-4321-2006.
- 1443 Dukowicz, J. K., and J. R. Baumgardner (2000), Incremental remapping as a trans-
1444 port/advection algorithm, *Journal of Computational Physics*, 160(1), 318–335, doi:
1445 10.1006/jcph.2000.6465.
- 1446 Dunavant, D. A. (1985), High degree efficient symmetrical Gaussian quadrature rules for
1447 the triangle, *International Journal for Numerical Methods in Engineering*, 21(6), 1129–
1448 1148, doi:10.1002/nme.1620210612.
- 1449 Durack, P. J., and K. E. Taylor (2017), PCMDI AMIP SST and sea-ice boundary condi-
1450 tions version 1.1.3, doi:10.22033/ESGF/input4MIPs.1735.
- 1451 Durack, P. J., K. E. Taylor, V. Eyring, S. K. Ames, T. Hoang, D. Nadeau, C. Doutriaux,
1452 M. Stockhause, and P. J. Gleckler (2018), Toward standardized data sets for climate
1453 model experimentation, *Eos*, 99, doi:10.1029/2018EO101751.
- 1454 Emmons, L. K., S. Walters, P. G. Hess, J. F. Lamarque, G. G. Pfister, D. Fillmore,
1455 C. Granier, A. Guenther, D. Kinnison, T. Laepple, J. Orlando, X. Tie, G. Tyndall,
1456 C. Wiedinmyer, S. L. Baughcum, and S. Kloster (2010), Description and evaluation of
1457 the model for ozone and related chemical tracers, version 4 (MOZART-4), *Geoscientific*
1458 *Model Development*, 3(1), 43–67, doi:10.5194/gmd-3-43-2010.
- 1459 Evans, K. J., S. Mahajan, M. Veneziani, X. S. Asay-Davis, L. P. Van Roekel, M. Branstet-
1460 ter, M. R. Petersen, P. J. Wolfram, C. S. Zender, and S. M. Burrows (2018), ACME
1461 Priority Metrics (A-PRIME), doi:10.5281/zenodo.1311053.
- 1462 Eyring, V., S. Bony, G. A. Meehl, C. A. Senior, B. Stevens, R. J. Stouffer, and K. E. Tay-
1463 lor (2016), Overview of the Coupled Model Intercomparison Project Phase 6 (CMIP6)
1464 experimental design and organization, *Geoscientific Model Development*, 9(5), 1937–
1465 1958, doi:10.5194/gmd-9-1937-2016.
- 1466 Fekete, B. M., and C. J. Vörösmarty (2011), ISLSCP II UNH/GRDC Composite Monthly
1467 Runoff, doi:10.3334/ornladaac/994.

- 1468 Flanner, M. G., X. Liu, C. Zhou, J. E. Penner, and C. Jiao (2012), Enhanced solar energy
1469 absorption by internally-mixed black carbon in snow grains, *Atmospheric Chemistry and*
1470 *Physics*, *12*(10), 4699–4721, doi:10.5194/acp-12-4699-2012.
- 1471 Forster, P. M., T. Richardson, A. C. Maycock, C. J. Smith, B. H. Samset, G. Myhre,
1472 T. Andrews, R. Pincus, and M. Schulz (2016), Recommendations for diagnosing effec-
1473 tive radiative forcing from climate models for CMIP6, *Journal of Geophysical Research:*
1474 *Atmospheres*, *121*(20), 12,460–12,475, doi:10.1002/2016JD025320.
- 1475 Gent, P. R., and J. C. McWilliams (1990), Isopycnal mixing in ocean circulation models,
1476 *Journal of Physical Oceanography*, *20*(1), 150–155.
- 1477 Geoffroy, O., D. Saint-Martin, D. J. L. Olivié, A. Voldoire, G. Bellon, and S. Tytéca
1478 (2013), Transient climate response in a two-layer energy-balance model. Part i: Ana-
1479 lytical solution and parameter calibration using CMIP5 AOGCM experiments, *Journal*
1480 *of Climate*, *26*(6), 1841–1857, doi:10.1175/jcli-d-12-00195.1.
- 1481 Gettelman, A., H. Morrison, S. Santos, P. Bogenschutz, and P. M. Caldwell (2015), Ad-
1482 vanced two-moment bulk microphysics for global models. Part II: Global model so-
1483 lutions and aerosol-cloud interactions, *Journal of Climate*, *28*(3), 1288–1307, doi:
1484 10.1175/jcli-d-14-00103.1.
- 1485 GISTEMP Team (2018), GISS surface temperature analysis (GISTEMP), NASA Goddard
1486 Institute for Space Studies. <https://data.giss.nasa.gov/gistemp>. Dataset accessed 2018-05-
1487 24.
- 1488 Gleckler, P., C. Doutriaux, P. Durack, K. Taylor, Y. Zhang, D. Williams, E. Mason, and
1489 J. Servonnat (2016), A more powerful reality test for climate models, *Eos*, *97*, doi:10.
1490 1029/2016eo051663.
- 1491 Gleckler, P. J., K. E. Taylor, and C. Doutriaux (2008), Performance metrics for climate
1492 models, *Journal of Geophysical Research*, *113*(D6), doi:10.1029/2007jd008972.
- 1493 Glickman, T. S. (2000), *Glossary of Meteorology, Second Edition*, 855 pp., American Me-
1494 teorological Society.
- 1495 Golaz, J.-C., V. E. Larson, and W. R. Cotton (2002), A PDF-based model for boundary
1496 layer clouds. Part I: Method and model description, *Journal of the Atmospheric Sciences*,
1497 *59*(24), 3540–3551, doi:10.1175/1520-0469(2002)059<3540:APBMFB>2.0.CO;2.
- 1498 Gregory, J. M., W. J. Ingram, M. A. Palmer, G. S. Jones, P. A. Stott, R. B. Thorpe,
1499 J. A. Lowe, T. C. Johns, and K. D. Williams (2004), A new method for diagnosing
1500 radiative forcing and climate sensitivity, *Geophysical Research Letters*, *31*(3), doi:

1501 10.1029/2003gl018747.

1502 Griffies, S. M., R. C. Pacanowski, and R. W. Hallberg (2000), Spurious diapycnal mix-
1503 ing associated with advection in az-coordinate ocean model, *Monthly Weather Review*,
1504 *128*(3), 538–564.

1505 Hansen, J. (2005), Efficacy of climate forcings, *Journal of Geophysical Research*,
1506 *110*(D18), doi:10.1029/2005jd005776.

1507 Hansen, J., R. Ruedy, M. Sato, and K. Lo (2010), Global surface temperature change, *Re-*
1508 *views of Geophysics*, *48*(4), doi:10.1029/2010rg000345.

1509 Hawkins, E., and R. Sutton (2016), Connecting climate model projections of global tem-
1510 perature change with the real world, *Bulletin of the American Meteorological Society*,
1511 *97*(6), 963–980, doi:10.1175/bams-d-14-00154.1.

1512 Hegglin, M., D. Kinnison, J.-F. Lamarque, and D. Plummer (2016), CCMI ozone in sup-
1513 port of CMIP6 - version 1.0, doi:10.22033/ESGF/input4MIPs.1115.

1514 Held, I. M., M. Winton, K. Takahashi, T. Delworth, F. Zeng, and G. K. Vallis (2010),
1515 Probing the fast and slow components of global warming by returning abruptly to prein-
1516 dustrial forcing, *Journal of Climate*, *23*(9), 2418–2427, doi:10.1175/2009jcli3466.1.

1517 Hibler, W., A. Roberts, P. Heil, A. Y. Proshutinsky, H. Simmons, and J. Lovick (2006),
1518 Modeling m2 tidal variability in arctic sea-ice drift and deformation, *Annals of Glaciol-*
1519 *ogy*, *44*, 418–428.

1520 Hibler III, W. D. (1979), A dynamic thermodynamic sea ice model, *Journal of Physical*
1521 *Oceanography*, *9*(4), 815–846, doi:10.1175/1520-0485(1979)009\$<\$0815:ADTSIM\$>
1522 \$2.0.CO;2.

1523 Hoesly, R., S. Smith, L. Feng, Z. Klimont, G. Janssens-Maenhout, T. Pitkanen, J. J.
1524 Seibert, L. Vu, R. J. Andres, R. M. Bolt, T. C. Bond, L. Dawidowski, N. Kholod, J.-
1525 i. Kurokawa, M. Li, L. Liu, Z. Lu, M. C. P. Moura, P. R. O'Rourke, and Q. Zhang
1526 (2017), Historical emissions (1750 - 2014) - CEDS - v2017-05-18, doi:10.22033/ESGF/
1527 input4MIPs.1241.

1528 Hoesly, R. M., S. J. Smith, L. Feng, Z. Klimont, G. Janssens-Maenhout, T. Pitkanen,
1529 J. J. Seibert, L. Vu, R. J. Andres, R. M. Bolt, T. C. Bond, L. Dawidowski, N. Kholod,
1530 J. I. Kurokawa, M. Li, L. Liu, Z. Lu, M. C. P. Moura, P. R. O'Rourke, and Q. Zhang
1531 (2018), Historical (1750-2014) anthropogenic emissions of reactive gases and aerosols
1532 from the community emissions data system (CEDS), *Geosci. Model Dev.*, *11*(1), 369–
1533 408, doi:10.5194/gmd-11-369-2018.

- 1534 Holland, M. M., D. A. Bailey, B. P. Briegleb, B. Light, and E. Hunke (2012), Improved
1535 sea ice shortwave radiation physics in CCSM4: The impact of melt ponds and aerosols
1536 on arctic sea ice, *Journal of Climate*, 25(5), 1413–1430, doi:10.1175/JCLI-D-11-00078.
1537 1.
- 1538 Holte, J., L. D. Talley, J. Gilson, and D. Roemmich (2017), An Argo mixed layer climatol-
1539 ogy and database, *Geophysical Research Letters*, 44(11), 5618–5626.
- 1540 Hourdin, F., T. Mauritsen, A. Gettelman, J.-C. Golaz, V. Balaji, Q. Duan, D. Folini, D. Ji,
1541 D. Klocke, Y. Qian, F. Rauser, C. Rio, L. Tomassini, M. Watanabe, and D. Williamson
1542 (2017), The art and science of climate model tuning, *Bulletin of the American Meteorolo-*
1543 *gical Society*, 98(3), 589–602, doi:10.1175/bams-d-15-00135.1.
- 1544 Hsu, J., and M. J. Prather (2009), Stratospheric variability and tropospheric ozone, *Jour-*
1545 *nal of Geophysical Research-Atmospheres*, 114, 15, doi:10.1029/2008jd010942.
- 1546 Huang, B., V. F. Banzon, E. Freeman, J. Lawrimore, W. Liu, T. C. Peterson, T. M. Smith,
1547 P. W. Thorne, S. D. Woodruff, and H.-M. Zhang (2015), Extended Reconstructed Sea
1548 Surface Temperature (ERSST), Version 4., NOAA National Centers for Environmental
1549 Information, doi:10.7289/V5KD1VVF.
- 1550 Huffman, G. J., R. F. Adler, D. T. Bolvin, and G. Gu (2009), Improving the global
1551 precipitation record: GPCP version 2.1, *Geophysical Research Letters*, 36(17), doi:
1552 10.1029/2009gl040000.
- 1553 Hunke, E., R. Allard, D. Bailey, A. Craig, A. Damsgaard, F. Dupont, A. DuVivier,
1554 M. Holland, N. Jeffery, J.-F. Lemieux, C. Newman, A. Roberts, A. Turner, M. Turner,
1555 and M. Winton (2018), CICE-Consortium/Icepack version 1.0.2, doi:10.5281/zenodo.
1556 1213463.
- 1557 Hunke, E. C., and J. K. Dukowicz (1997), An elastic-viscous-plastic model for sea
1558 ice dynamics, *Journal of Physical Oceanography*, 27(9), 1849–1867, doi:10.1175/
1559 1520-0485(1997)027\$<\$1849:AEVPMF\$>\$2.0.CO;2.
- 1560 Hunke, E. C., and J. K. Dukowicz (2002), The elastic-viscous-plastic sea ice dynamics
1561 model in general orthogonal curvilinear coordinates on a sphere – Incorporation of met-
1562 ric terms, *Monthly Weather Review*, 130(7), 1848–1865, doi:10.1175/1520-0493(2002)
1563 130\$<\$1848:TEVPSIS\$>\$2.0.CO;2.
- 1564 Hunke, E. C., D. A. Hebert, and O. Lecomte (2013), Level-ice melt ponds in the Los
1565 Alamos sea ice model, CICE, *Ocean Modelling*, 71, 26–42, doi:10.1016/j.ocemod.2012.
1566 11.008.

- 1567 Hunke, E. C., W. H. Lipscomb, A. K. Turner, N. Jeffery, and S. Elliott (2015), CICE: the
1568 Los Alamos sea ice model documentation and software user's manual version 5.1, *Tech.*
1569 *rep.*, Los Alamos National Laboratory.
- 1570 Hurrell, J. W., J. J. Hack, D. Shea, J. M. Caron, and J. Rosinski (2008), A new sea sur-
1571 face temperature and sea ice boundary dataset for the community atmosphere model,
1572 *Journal of Climate*, 21(19), 5145–5153.
- 1573 Hurrell, J. W., M. M. Holland, P. R. Gent, S. Ghan, J. E. Kay, P. J. Kushner, J.-F. Lamar-
1574 que, W. G. Large, D. Lawrence, K. Lindsay, W. H. Lipscomb, M. C. Long, N. Ma-
1575 howald, D. R. Marsh, R. B. Neale, P. Rasch, S. Vavrus, M. Vertenstein, D. Bader,
1576 W. D. Collins, J. J. Hack, J. Kiehl, and S. Marshall (2013), The community earth sys-
1577 tem model: A framework for collaborative research, *Bulletin of the American Meteoro-*
1578 *logical Society*, 94(9), 1339–1360, doi:10.1175/bams-d-12-00121.1.
- 1579 Hurtt, G., L. Chini, R. Sahajpal, and S. Frohling (2017),
1580 input4MIPs.UofMD.landState.CMIP.UofMD-landState-2-1-h, doi:10.22033/ESGF/
1581 input4MIPs.1127.
- 1582 IPCC (2007), Climate change 2007: The physical science basis, in *Contribution of Work-*
1583 *ing Group I to the Fourth Assessment Report of the Intergovernmental Panel on Climate*
1584 *Change*, edited by S. Solomon, D. Qin, M. Manning, Z. Chen, M. Marquis, K. Acertyt,
1585 M. Tignor, and H. L. Miller, p. 996pp, Cambridge University Press, Cambridge, United
1586 Kingdom and New York, NY, USA.
- 1587 IPCC (2013), *Climate Change 2013: The Physical Science Basis. Contribution of Work-*
1588 *ing Group I to the Fifth Assessment Report of the Intergovernmental Panel on Climate*
1589 *Change*, 1535 pp., Cambridge University Press, Cambridge, United Kingdom and New
1590 York, NY, USA, doi:10.1017/CBO9781107415324.
- 1591 Jacob, R., J. Larson, and E. Ong (2005), M \times n communication and parallel interpolation
1592 in community climate system model version 3 using the model coupling toolkit, *The*
1593 *International Journal of High Performance Computing Applications*, 19(3), 293–307.
- 1594 Jochum, M., G. Danabasoglu, M. Holland, Y.-O. Kwon, and W. Large (2008), Ocean vis-
1595 cosity and climate, *Journal of Geophysical Research: Oceans*, 113(C6).
- 1596 Jones, P. D., M. New, D. E. Parker, S. Martin, and I. G. Rigor (1999), Surface air temper-
1597 ature and its changes over the past 150 years, *Reviews of Geophysics*, 37(2), 173–199,
1598 doi:10.1029/1999rg900002.

- 1599 Kay, J., C. Deser, A. Phillips, A. Mai, C. Hannay, G. Strand, J. Arblaster, S. Bates,
1600 G. Danabasoglu, J. Edwards, et al. (2015), The Community Earth System Model
1601 (CESM) large ensemble project: A community resource for studying climate change
1602 in the presence of internal climate variability, *Bulletin of the American Meteorological*
1603 *Society*, *96*(8), 1333–1349.
- 1604 Knutti, R., M. A. A. Rugenstein, and G. C. Hegerl (2017), Beyond equilibrium climate
1605 sensitivity, *Nature Geoscience*, *10*(10), 727–736, doi:10.1038/ngeo3017.
- 1606 Kwok, R., and G. F. Cunningham (2008), ICESat over Arctic sea ice: Estimation of snow
1607 depth and ice thickness, *Journal of Geophysical Research-Oceans*, *113*(C8), doi:10.1029/
1608 2008jc004753.
- 1609 Kwok, R., and D. A. Rothrock (2009), Decline in Arctic sea ice thickness from submarine
1610 and ICESat records: 1958? @ S2008, *Geophysical Research Letters*, *36*(15), doi:10.
1611 1029/2009gl039035.
- 1612 Lagerloef, G., H.-Y. Kao, T. Meissner, and J. Vazquez (2015), Aquarius salinity
1613 validation analysis; data version 4.0, Aquarius Science Team Technology Doc-
1614 ument, ftp://podaac.jpl.nasa.gov/SalinityDensity/aquarius/docs/v4/AQ-014-PS-
1615 0016_AquariusSalinityDataValidationAnalysis_DatasetVersion4.0and3.0.pdf.
- 1616 Large, W., and S. Yeager (2009), The global climatology of an interannually varying air-
1617 sea flux data set, *Clim. Dynam.*, *33*(2-3), 341–364, doi:10.1007/s00382-008-0441-3.
- 1618 Larson, J., R. Jacob, and E. Ong (2005), The model coupling toolkit: a new Fortran90
1619 toolkit for building multiphysics parallel coupled models, *The International Journal of*
1620 *High Performance Computing Applications*, *19*(3), 277–292.
- 1621 Larson, V. E. (2017), CLUBB-SILHS: A parameterization of subgrid variability in the
1622 atmosphere, arXiv:1711.03675.
- 1623 Larson, V. E., and J.-C. Golaz (2005), Using probability density functions to derive con-
1624 sistent closure relationships among higher-order moments, *Monthly Weather Review*,
1625 *133*(4), 1023–1042, doi:10.1175/MWR2902.1.
- 1626 Lawrence, P. J., and T. N. Chase (2007), Representing a new MODIS consistent land
1627 surface in the Community Land Model (CLM 3.0), *Journal of Geophysical Research*,
1628 *112*(G1), doi:10.1029/2006jg000168.
- 1629 Leung, L. R., D. Bader, M. Taylor, M. R., and D. Koch (2019), An introduction to the En-
1630 ergy Exascale Earth System Model (E3SM) special collection, *Journal of Advances in*
1631 *Modeling Earth Systems*, joint submission to AGU Special Collection on E3SM.

- 1632 Li, H., M. S. Wigmosta, H. Wu, M. Huang, Y. Ke, A. M. Coleman, and L. R. Leung
1633 (2013), A physically based runoff routing model for land surface and earth system mod-
1634 els, *Journal of Hydrometeorology*, *14*(3), 808–828, doi:10.1175/jhm-d-12-015.1.
- 1635 Li, H.-Y., L. R. Leung, A. Getirana, M. Huang, H. Wu, Y. Xu, J. Guo, and N. Voisin
1636 (2015), Evaluating global streamflow simulations by a physically based routing model
1637 coupled with the Community Land Model, *Journal of Hydrometeorology*, *16*(2), 948–
1638 971, doi:10.1175/jhm-d-14-0079.1.
- 1639 Lipscomb, W. H. (2001), Remapping the thickness distribution in sea ice models, *Journal*
1640 *of Geophysical Research: Oceans*, *106*(C7), 13,989–14,000, doi:10.1029/2000JC000518.
- 1641 Lipscomb, W. H., and E. C. Hunke (2004), Modeling sea ice transport using incremental
1642 remapping, *Monthly Weather Review*, *132*(6), 1341–1354, doi:10.1175/1520-0493(2004)
1643 132\$<\$1341:MSITUI\$>\$2.0.CO;2.
- 1644 Lipscomb, W. H., and T. D. Ringler (2005), An incremental remapping transport scheme
1645 on a spherical geodesic grid, *Monthly Weather Review*, *133*(8), 2335–2350, doi:10.1175/
1646 MWR2983.1.
- 1647 Lipscomb, W. H., E. C. Hunke, W. Maslowski, and J. Jakacki (2007), Ridging, strength,
1648 and stability in high-resolution sea ice models, *Journal of Geophysical Research:*
1649 *Oceans*, *112*(C3), doi:10.1029/2005JC003355, C03S91.
- 1650 Lique, C., M. M. Holland, Y. B. Dibike, D. M. Lawrence, and J. A. Screen (2016), Mod-
1651 eling the arctic freshwater system and its integration in the global system: Lessons
1652 learned and future challenges, *Journal of Geophysical Research: Biogeosciences*, *121*(3),
1653 540–566.
- 1654 Liu, X., R. C. Easter, S. J. Ghan, R. Zaveri, P. Rasch, X. Shi, J.-F. Lamarque, A. Gettel-
1655 man, H. Morrison, F. Vitt, A. Conley, S. Park, R. Neale, C. Hannay, A. M. L. Ekman,
1656 P. Hess, N. Mahowald, W. Collins, M. J. Iacono, C. S. Bretherton, M. G. Flanner, and
1657 D. Mitchell (2012), Toward a minimal representation of aerosols in climate models: de-
1658 scription and evaluation in the Community Atmosphere Model CAM5, *Geoscientific*
1659 *Model Development*, *5*(3), 709–739, doi:10.5194/gmd-5-709-2012.
- 1660 Liu, X., P. L. Ma, H. Wang, S. Tilmes, B. Singh, R. C. Easter, S. J. Ghan, and P. J. Rasch
1661 (2016), Description and evaluation of a new four-mode version of the modal aerosol
1662 module (MAM4) within version 5.3 of the community atmosphere model, *Geosci.*
1663 *Model Dev.*, *9*(2), 505–522, doi:10.5194/gmd-9-505-2016.

- 1664 Loeb, N. G., B. A. Wielicki, D. R. Doelling, G. L. Smith, D. F. Keyes, S. Kato,
1665 N. Manalo-Smith, and T. Wong (2009), Toward optimal closure of the earth's top-
1666 of-atmosphere radiation budget, *Journal of Climate*, 22(3), 748–766, doi:10.1175/
1667 2008jcli2637.1.
- 1668 Matthes, K., B. Funke, T. Kruschke, and S. Wahl (2017a), input4MIPs.SOLARIS-
1669 HEPPA.solar.CMIP.SOLARIS-HEPPA-3-2, doi:10.22033/ESGF/input4MIPs.1122.
- 1670 Matthes, K., B. Funke, M. E. Andersson, L. Barnard, J. Beer, P. Charbonneau, M. A. Clil-
1671 verd, T. Dudok de Wit, M. Haberreiter, A. Hendry, C. H. Jackman, M. Kretzschmar,
1672 T. Kruschke, M. Kunze, U. Langematz, D. R. Marsh, A. C. Maycock, S. Misios, C. J.
1673 Rodger, A. A. Scaife, A. Seppälä, M. Shangguan, M. Sinnhuber, K. Tourpali,
1674 I. Usoskin, M. van de Kamp, P. T. Verronen, and S. Versick (2017b), Solar forcing for
1675 CMIP6 (v3.2), *Geosci. Model Dev.*, 10(6), 2247–2302, doi:10.5194/gmd-10-2247-2017.
- 1676 McGregor, H., M. J. Fischer, M. Gagan, D. Fink, S. J. Phipps, H. Wong, and
1677 C. Woodroffe (2013), A weak El Niño/Southern Oscillation with delayed seasonal
1678 growth around 4,300 years ago, *Nature Geoscience*, 6(11), 949.
- 1679 McPhaden, M. J., X. Zhang, H. H. Hendon, and M. C. Wheeler (2006), Large scale dy-
1680 namics and MJO forcing of ENSO variability, *Geophysical Research Letters*, 33(16),
1681 doi:10.1029/2006GL026786.
- 1682 Meier, W., F. Fetterer, M. Savoie, S. Mallory, R. Duerr, and J. Stroeve (2017),
1683 NOAA/NSIDC Climate Data Record of Passive Microwave Sea Ice Concentration, Ver-
1684 sion 3, *Electronic Media*, doi:10.7265/N59P2ZTG.
- 1685 Meinshausen, M., and E. Vogel (2016), input4MIPs.UoM.GHGConcentrations.CMIP.UoM-
1686 CMIP-1-2-0, doi:10.22033/ESGF/input4MIPs.1118.
- 1687 Meinshausen, M., E. Vogel, A. Nauels, K. Lorbacher, N. Meinshausen, D. M. Etheridge,
1688 P. J. Fraser, S. A. Montzka, P. J. Rayner, C. M. Trudinger, P. B. Krummel, U. Bey-
1689 erle, J. G. Canadell, J. S. Daniel, I. G. Enting, R. M. Law, C. R. Lunder, S. O'Doherty,
1690 R. G. Prinn, S. Reimann, M. Rubino, G. J. M. Velders, M. K. Vollmer, R. H. J. Wang,
1691 and R. Weiss (2017), Historical greenhouse gas concentrations for climate modelling
1692 (CMIP6), *Geosci. Model Dev.*, 10(5), 2057–2116, doi:10.5194/gmd-10-2057-2017.
- 1693 Menary, M. B., T. Kuhlbrodt, J. Ridley, M. B. Andrews, O. B. Dimdore-Miles, J. De-
1694 shayes, R. Eade, L. Gray, S. Ineson, J. Mignot, et al. (2018), Preindustrial control sim-
1695 ulations with HadGEM3-GC3. 1 for cmip6, *Journal of Advances in Modeling Earth Sys-*
1696 *tems*.

- 1697 Morice, C. P., J. J. Kennedy, N. A. Rayner, and P. D. Jones (2012), Quantifying uncer-
1698 tainties in global and regional temperature change using an ensemble of observational
1699 estimates: The HadCRUT4 data set, *Journal of Geophysical Research: Atmospheres*,
1700 *117*(D8), n/a–n/a, doi:10.1029/2011jd017187.
- 1701 Morrison, H., and A. Gettelman (2008), A new two-moment bulk stratiform cloud mi-
1702 crophysics scheme in the Community Atmosphere Model, version 3 (CAM3). Part
1703 I: Description and numerical tests, *Journal of Climate*, *21*, 3642–3659, doi:10.1175/
1704 2008JCLI2105.1.
- 1705 Naughten, K. A., B. K. Galton-Fenzi, K. J. Meissner, M. H. England, G. B. Brassing-
1706 ton, F. Colberg, T. Hattermann, and J. B. Debernard (2017), Spurious sea ice formation
1707 caused by oscillatory ocean tracer advection schemes, *Ocean Modelling*, *116*, 108–117,
1708 doi:10.1016/j.ocemod.2017.06.010.
- 1709 Newman, M., A. T. Wittenberg, L. Cheng, G. P. Compo, and C. A. Smith (2018), The
1710 extreme 2015/16 el niño, in the context of historical climate variability and change, *Bul-*
1711 *letin of the American Meteorological Society*, *99*(1), S16–S20.
- 1712 Oleson, K. W., D. M. Lawrence, G. B. Bonan, B. Drewniak, M. Huang, C. D. Koven,
1713 S. Levis, F. Li, W. J. Riley, Z. M. Subin, S. C. Swenson, P. E. Thornton, A. Bozbiyik,
1714 R. Fisher, C. L. Heald, E. Kluzek, J.-F. Lamarque, P. J. Lawrence, L. R. Leung, W. Lip-
1715 scomb, S. Muszala, D. M. Ricciuto, W. Sacks, Y. Sun, J. Tang, and Z.-L. Yang (2013),
1716 Technical description of version 4.5 of the community land model (clm), *Tech. Rep.*
1717 *NCARTN-503+STR*, National Center for Atmospheric Research, Boulder, Colorado,
1718 USA.
- 1719 Padilla, L. E., G. K. Vallis, and C. W. Rowley (2011), Probabilistic estimates of transient
1720 climate sensitivity subject to uncertainty in forcing and natural variability, *Journal of*
1721 *Climate*, *24*(21), 5521–5537, doi:10.1175/2011jcli3989.1.
- 1722 Petersen, M. R., D. W. Jacobsen, T. D. Ringler, M. W. Hecht, and M. E. Maltrud (2015),
1723 Evaluation of the arbitrary Lagrangian–Eulerian vertical coordinate method in the
1724 MPAS-Ocean model, *Ocean Modelling*, *86*, 93–113, doi:10.1016/j.ocemod.2014.12.004.
- 1725 Qian, Y., H. Wan, B. Yang, J.-C. Golaz, B. Harrop, Z. Hou, V. E. Larson, L. R. Le-
1726 ung, G. Lin, W. Lin, P.-L. Ma, H.-Y. Ma, P. Rasch, B. Singh, H. Wang, S. Xie, and
1727 K. Zhang (2018), Parametric sensitivity and uncertainty quantification in the version
1728 1 of E3SM Atmosphere Model based on short Perturbed Parameters Ensemble simula-
1729 tions, *Journal of Geophysical Research: Atmospheres*, doi:10.1029/2018jd028927.

- 1730 Rasch, P. J., S. Xie, P.-L. Ma, W. Lin, H. Wang, Q. Tang, S. M. Burrows, P. Caldwell,
1731 K. Zhang, R. C. Easter, P. Cameron-Smith, B. Singh, H. Wan, J.-C. Golaz, B. E. Har-
1732 rop, E. Roesler, J. Bacmeister, V. E. Larson, K. J. Evans, Y. Qian, M. Taylor, L. Leung,
1733 Y. Zhang, L. Brent, M. Branstetter, C. Hannay, S. Mahajan, A. Mametjanov, R. Neale,
1734 J. H. Richter, J.-H. Yoon, C. S. Zender, D. Bader, M. Flanner, J. G. Foucar, R. Jacob,
1735 N. Keen, S. A. Klein, X. Liu, A. G. Salinger, M. Shrivastava, and Y. Yang (2019), An
1736 Overview of the Atmospheric Component of the Energy Exascale Earth System Model,
1737 *Journal of Advances in Modeling Earth Systems*, joint submission to AGU Special Col-
1738 lection on E3SM.
- 1739 Rayner, N., D. Parker, C. Folland, E. Horton, L. Alexander, and D. Rowell (2003), The
1740 global sea-ice and sea surface temperature (HadISST) data sets, *J. Geophys. Res.*
- 1741 Reckinger, S. M., M. R. Petersen, and S. J. Reckinger (2015), A study of overflow simu-
1742 lations using MPAS-Ocean: Vertical grids, resolution, and viscosity, *Ocean Modelling*,
1743 *96*, 291–313, doi:10.1016/j.ocemod.2015.09.006.
- 1744 Ringler, T., J. Thuburn, J. Klemp, and W. Skamarock (2010), A unified approach to en-
1745 ergy conservation and potential vorticity dynamics for arbitrarily-structured C-grids,
1746 *Journal of Computational Physics*, *229(9)*, 3065–3090, doi:10.1016/j.jcp.2009.12.007.
- 1747 Ringler, T., M. Petersen, R. L. Higdon, D. Jacobsen, P. W. Jones, and M. Maltrud (2013),
1748 A multi-resolution approach to global ocean modeling, *Ocean Modelling*, *69*, 211–232,
1749 doi:10.1016/j.ocemod.2013.04.010.
- 1750 Sallée, J.-B., E. Shuckburgh, N. Bruneau, A. J. Meijers, T. J. Bracegirdle, and Z. Wang
1751 (2013), Assessment of southern ocean mixed-layer depths in cmip5 models: Histori-
1752 cal bias and forcing response, *Journal of Geophysical Research: Oceans*, *118(4)*, 1845–
1753 1862.
- 1754 Schmidt, G. A., C. M. Bitz, U. Mikolajewicz, and L.-B. Tremblay (2004), Ice–ocean
1755 boundary conditions for coupled models, *Ocean Modelling*, *7(1-2)*, 59–74.
- 1756 Schmidt, G. A., D. Bader, L. J. Donner, G. S. Elsaesser, J.-C. Golaz, C. Hannay,
1757 A. Molod, R. B. Neale, and S. Saha (2017), Practice and philosophy of climate model
1758 tuning across six US modeling centers, *Geoscientific Model Development*, *10(9)*, 3207–
1759 3223, doi:10.5194/gmd-10-3207-2017.
- 1760 Shrivastava, M., R. C. Easter, X. Liu, A. Zelenyuk, B. Singh, K. Zhang, P.-L. Ma,
1761 D. Chand, S. Ghan, J. L. Jimenez, Q. Zhang, J. Fast, P. J. Rasch, and P. Tiitta (2015),
1762 Global transformation and fate of SOA: Implications of low-volatility SOA and gas-

- 1763 phase fragmentation reactions, *Journal of Geophysical Research: Atmospheres*, *120*(9),
1764 4169–4195, doi:10.1002/2014JD022563.
- 1765 Skamarock, W. C., and A. Gassmann (2011), Conservative transport schemes for spher-
1766 ical geodesic grids: High-order flux operators for ODE-based time integration, *Mon.*
1767 *Weather Rev.*, *139*, 2962–2975, doi:10.1175/MWR-D-10-05056.1.
- 1768 Smeed, D., G. McCarth, B. Moat, W. Johns, M. Baringer, and C. Meinen (2017), Atlantic
1769 meridional overturning circulation observed by the RAPID-MOCHA-WBTS (RAPID-
1770 Meridional Overturning Circulation and Heatflux Array-Western Boundary Time Series)
1771 array at 26N from 2004 to 2017, British Oceanographic Data Centre - Natural Environ-
1772 ment Research Council, UK, doi:10.5285/5acfd143-1104-7b58-e053-6c86abc0d94b.
- 1773 Smith, T. M., R. W. Reynolds, T. C. Peterson, and J. Lawrimore (2008), Improvements
1774 to NOAA’s historical merged land–ocean surface temperature analysis (1880–2006),
1775 *Journal of Climate*, *21*(10), 2283–2296, doi:10.1175/2007jcli2100.1.
- 1776 Sobel, A. H., and E. D. Maloney (2000), Effect of enso and the mjo on western north
1777 pacific tropical cyclones, *Geophysical Research Letters*, *27*(12), 1739–1742, doi:
1778 10.1029/1999GL011043.
- 1779 Soden, B. J., I. M. Held, R. Colman, K. M. Shell, J. T. Kiehl, and C. A. Shields (2008),
1780 Quantifying climate feedbacks using radiative kernels, *Journal of Climate*, *21*(14),
1781 3504–3520, doi:10.1175/2007jcli2110.1.
- 1782 Steele, M., R. Morley, and W. Ermold (2001), Phc: A global ocean hydrography with a
1783 high-quality arctic ocean, *Journal of Climate*, *14*(9), 2079–2087.
- 1784 Stephens, G. L., J. Li, M. Wild, C. A. Clayson, N. Loeb, S. Kato, T. L’Ecuyer, P. W.
1785 Stackhouse, M. Lebsock, and T. Andrews (2012), An update on earth’s energy bal-
1786 ance in light of the latest global observations, *Nature Geoscience*, *5*(10), 691–696, doi:
1787 10.1038/ngeo1580.
- 1788 Stevens, B. (2015), Rethinking the lower bound on aerosol radiative forcing, *Journal of*
1789 *Climate*, *28*(12), 4794–4819, doi:10.1175/jcli-d-14-00656.1.
- 1790 Stevenson, S. (2012), Significant changes to ENSO strength and impacts in the twenty-first
1791 century: Results from CMIP5, *Geophysical Research Letters*, *39*(17).
- 1792 Stevenson, S., B. Fox-Kemper, M. Jochum, R. Neale, C. Deser, and G. Meehl (2012), Will
1793 there be a significant change to el niño in the twenty-first century?, *Journal of Climate*,
1794 *25*(6), 2129–2145.

- 1795 Sun, Q., C. Miao, Q. Duan, H. Ashouri, S. Sorooshian, and K.-L. Hsu (2018), A review
1796 of global precipitation data sets: Data sources, estimation, and intercomparisons, *Re-*
1797 *views of Geophysics*, 56(1), 79–107, doi:10.1002/2017rg000574.
- 1798 Taylor, K. E., D. Williamson, and F. Zwiers (2000), The sea surface temperature and sea
1799 ice concentration boundary conditions for AMIP II simulations, *PCMDI Report 60*,
1800 <https://pcmdi.llnl.gov/report/pdf/60.pdf>, Program for Climate Model Di-
1801 agnosis and Intercomparison, Lawrence Livermore National Laboratory.
- 1802 Thuburn, J., T. D. Ringler, W. C. Skamarock, and J. B. Klemp (2009), Numerical repre-
1803 sentation of geostrophic modes on arbitrarily structured C-grids, *Journal of Computa-*
1804 *tional Physics*, 228(22), 8321–8335, doi:10.1016/j.jcp.2009.08.006.
- 1805 Tilling, R. L., A. Ridout, A. Shepherd, and D. J. Wingham (2015), Increased Arctic sea
1806 ice volume after anomalously low melting in 2013, *Nature Geoscience*, 8(8), 643–646,
1807 doi:10.1038/ngeo2489.
- 1808 Torrence, C., and G. P. Compo (1998), A practical guide to wavelet analysis, *Bulletin of*
1809 *the American Meteorological society*, 79(1), 61–78.
- 1810 Trenberth, K. E., J. T. Fasullo, and J. Kiehl (2009), Earth’s global energy budget, *Bulletin*
1811 *of the American Meteorological Society*, 90(3), 311–324, doi:10.1175/2008bams2634.1.
- 1812 Turner, A. K., and E. C. Hunke (2015), Impacts of a mushy-layer thermodynamic ap-
1813 proach in global sea-ice simulations using the CICE sea-ice model, *Journal of Geophys-*
1814 *ical Research: Oceans*, 120(2), 1253–1275, doi:10.1002/2014JC010358.
- 1815 Turner, A. K., E. C. Hunke, and C. M. Bitz (2013), Two modes of sea-ice gravity
1816 drainage: A parameterization for large-scale modeling, *Journal of Geophysical Research:*
1817 *Oceans*, 118(5), 2279–2294, doi:10.1002/jgrc.20171.
- 1818 Ullrich, P. A., and M. A. Taylor (2015), Arbitrary-order conservative and consistent
1819 remapping and a theory of linear maps: Part i, *Monthly Weather Review*, 143(6), 2419–
1820 2440.
- 1821 U.S. Department of Energy (2013), U.S. Energy Sector Vulnerabilities to Climate Change
1822 and Extreme Weather, *Tech. Rep. DOE/PI-0013*, [https://www.energy.gov/sites/
1823 prod/files/2013/07/f2/20130710-Energy-Sector-Vulnerabilities-Report.
1824 pdf](https://www.energy.gov/sites/prod/files/2013/07/f2/20130710-Energy-Sector-Vulnerabilities-Report.pdf), U.S. Department of Energy.
- 1825 van Marle, M. J., S. Kloster, B. I. Magi, J. R. Marlon, A.-L. Daniau, R. D. Field, A. Ar-
1826 neth, M. Forrest, S. Hantson, N. M. Kehrwald, W. Knorr, G. Lasslop, F. Li, S. Man-
1827 geon, C. Yue, J. W. Kaiser, and G. R. van der Werf (2016), Biomass burning emissions

- 1828 for cmip6 (v1.2), doi:10.22033/ESGF/input4MIPs.1117.
- 1829 van Marle, M. J. E., S. Kloster, B. I. Magi, J. R. Marlon, A. L. Daniau, R. D. Field,
1830 A. Arneth, M. Forrest, S. Hantson, N. M. Kehrwald, W. Knorr, G. Lasslop, F. Li,
1831 S. Mangeon, C. Yue, J. W. Kaiser, and G. R. van der Werf (2017), Historic global
1832 biomass burning emissions for CMIP6 (BB4CMIP) based on merging satellite observa-
1833 tions with proxies and fire models (1750-2015), *Geosci. Model Dev.*, *10*(9), 3329–3357,
1834 doi:10.5194/gmd-10-3329-2017.
- 1835 Van Oldenborgh, G. J., S. Philip, and M. Collins (2005), El niño in a changing climate: a
1836 multi-model study, *Ocean Science*, *1*(2), 81–95.
- 1837 Van Roekel, L., A. J. Adcroft, G. Danabasoglu, S. M. Griffies, B. Kauffman, W. Large,
1838 M. Levy, B. G. Reichl, T. Ringler, and M. Schmidt (2018), The kpp boundary layer
1839 scheme for the ocean: Revisiting its formulation and benchmarking one-dimensional
1840 simulations relative to les, *Journal of Advances in Modeling Earth Systems*, *0*(0), doi:
1841 10.1029/2018MS001336.
- 1842 Vitart, F. (2017), Madden-Julian oscillation prediction and teleconnections in the s2s
1843 database, *Quarterly Journal of the Royal Meteorological Society*, *143*(706), 2210–2220,
1844 doi:10.1002/qj.3079.
- 1845 Waliser, D. E., K. M. Lau, W. Stern, and C. Jones (2003), Potential predictability of the
1846 madden-Julian oscillation, *Bulletin of the American Meteorological Society*, *84*(1),
1847 33–50, doi:10.1175/BAMS-84-1-33.
- 1848 Wang, H., S. A. Legg, and R. W. Hallberg (2015), Representations of the nordic seas
1849 overflows and their large scale climate impact in coupled models, *Ocean Modelling*, *86*,
1850 76–92.
- 1851 Wang, H., S. Legg, and R. Hallberg (2018), The effect of arctic freshwater pathways on
1852 north atlantic convection and the atlantic meridional overturning circulation, *Journal of*
1853 *Climate*, *31*(13), 5165–5188.
- 1854 Wild, M., D. Folini, C. Schär, N. Loeb, E. G. Dutton, and G. König-Langlo (2012), The
1855 global energy balance from a surface perspective, *Climate Dynamics*, *40*(11-12), 3107–
1856 3134, doi:10.1007/s00382-012-1569-8.
- 1857 Wittenberg, A. T. (2009), Are historical records sufficient to constrain enso simulations?,
1858 *Geophysical Research Letters*, *36*(12).
- 1859 Xie, S., W. Lin, P. J. Rasch, P.-L. Ma, R. Neale, V. E. Larson, Y. Qian, P. A. Bogenschutz,
1860 P. Caldwell, P. Cameron-Smith, J.-C. Golaz, S. Mahajan, B. Singh, Q. Tang, H. Wang,

- 1861 J.-H. Yoon, K. Zhang, and Y. Zhang (2018), Understanding cloud and convective char-
1862 acteristics in version 1 of the E3SM Atmosphere Model, *Journal of Advances in Model-*
1863 *ing Earth Systems*, doi:10.1029/2018ms001350.
- 1864 Zender, C. S. (2008), Analysis of self-describing gridded geoscience data with netCDF
1865 operators (NCO), *Environmental Modelling & Software*, 23(10-11), 1338–1342, doi:10.
1866 1016/j.envsoft.2008.03.004.
- 1867 Zender, C. S. (2018), netCDF Operators (NCO) version 4.7.8, doi:10.5281/zenodo.595745.
- 1868 Zhang, G., and N. A. McFarlane (1995), Sensitivity of climate simulations to the param-
1869 eterization of cumulus convection in the Canadian climate centre general circulation
1870 model, *Atmosphere-Ocean*, 33(3), 407–446, doi:10.1080/07055900.1995.9649539.
- 1871 Zhang, H.-M., B. Huang, J. Lawrimore, M. Menne, and T. M. Smith (2015), NOAA
1872 global surface temperature dataset, version 4.0, doi:10.7289/v5fn144h.
- 1873 Zhang, K., P. J. Rasch, M. A. Taylor, H. Wan, R. Leung, P.-L. Ma, J.-C. Golaz, J. Wolfe,
1874 W. Lin, B. Singh, S. Burrows, J.-H. Yoon, H. Wang, Y. Qian, Q. Tang, P. Caldwell,
1875 and S. Xie (2018), Impact of numerical choices on water conservation in the E3SM
1876 Atmosphere Model version 1 (EAMv1), *Geoscientific Model Development*, 11(5), 1971–
1877 1988, doi:10.5194/gmd-11-1971-2018.
- 1878 Zhao, M., J.-C. Golaz, I. M. Held, H. Guo, V. Balaji, R. Benson, J.-H. Chen, X. Chen,
1879 L. J. Donner, J. P. Dunne, K. Dunne, J. Durachta, S.-M. Fan, S. M. Freidenreich, S. T.
1880 Garner, P. Ginoux, L. M. Harris, L. W. Horowitz, J. P. Krasting, A. R. Langenhorst,
1881 Z. Liang, P. Lin, S.-J. Lin, S. L. Malyshev, E. Mason, P. C. D. Milly, Y. Ming, V. Naik,
1882 F. Paulot, D. Paynter, P. Phillipps, A. Radhakrishnan, V. Ramaswamy, T. Robinson,
1883 D. Schwarzkopf, C. J. Seman, E. Shevliakova, Z. Shen, H. Shin, L. G. Silvers, J. R.
1884 Wilson, M. Winton, A. T. Wittenberg, B. Wyman, and B. Xiang (2018), The GFDL
1885 global atmosphere and land model AM4.0/LM4.0: 2. model description, sensitivity
1886 studies, and tuning strategies, *Journal of Advances in Modeling Earth Systems*, 10(3),
1887 735–769, doi:10.1002/2017ms001209.

Figure 1.

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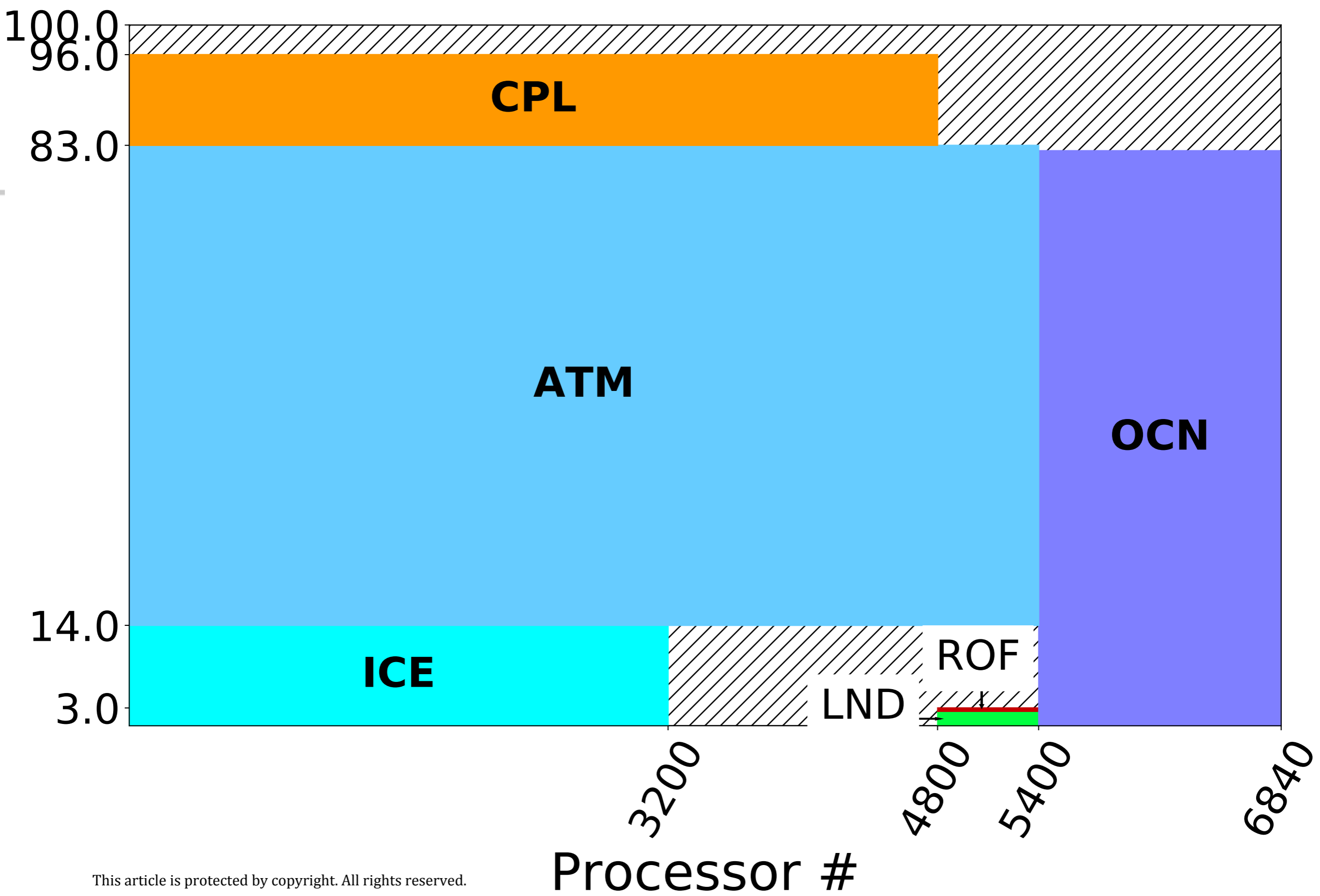


Figure 2.

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Spin-up simulations: global surface air temperature

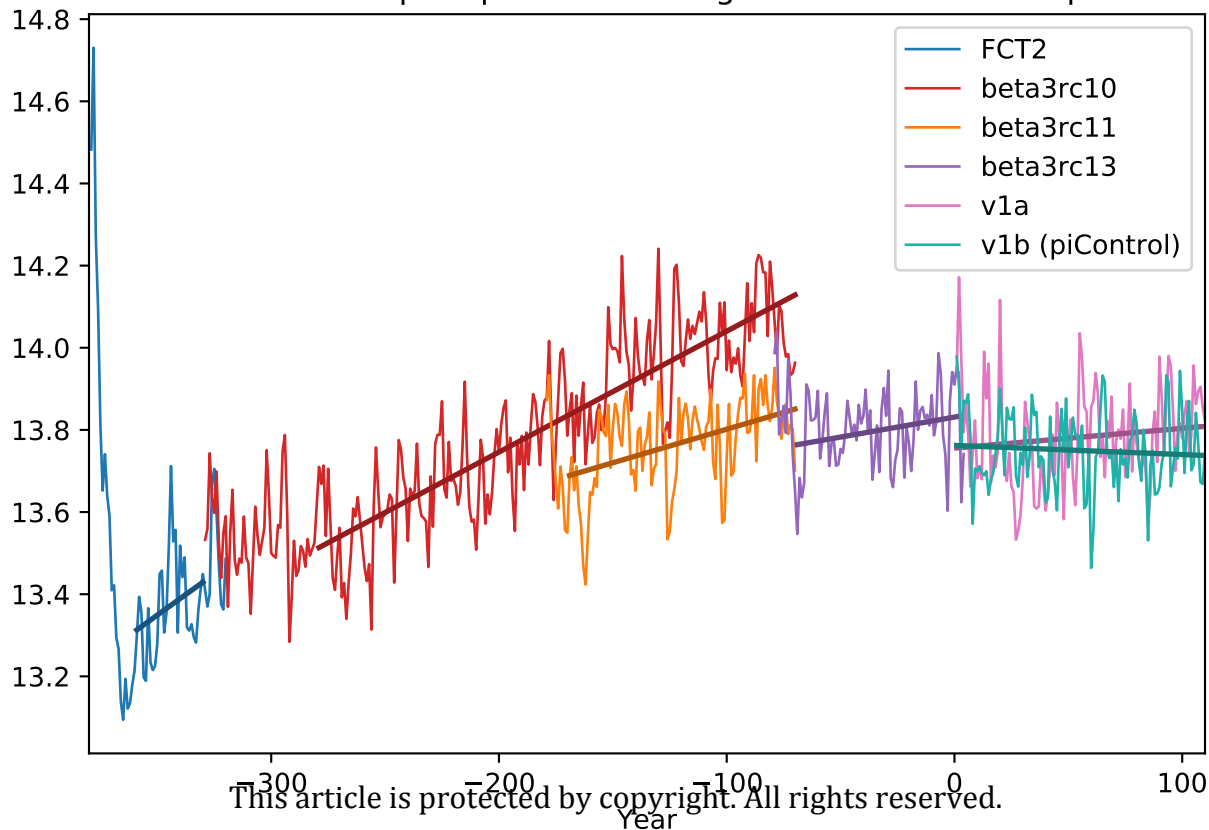


Figure 3.

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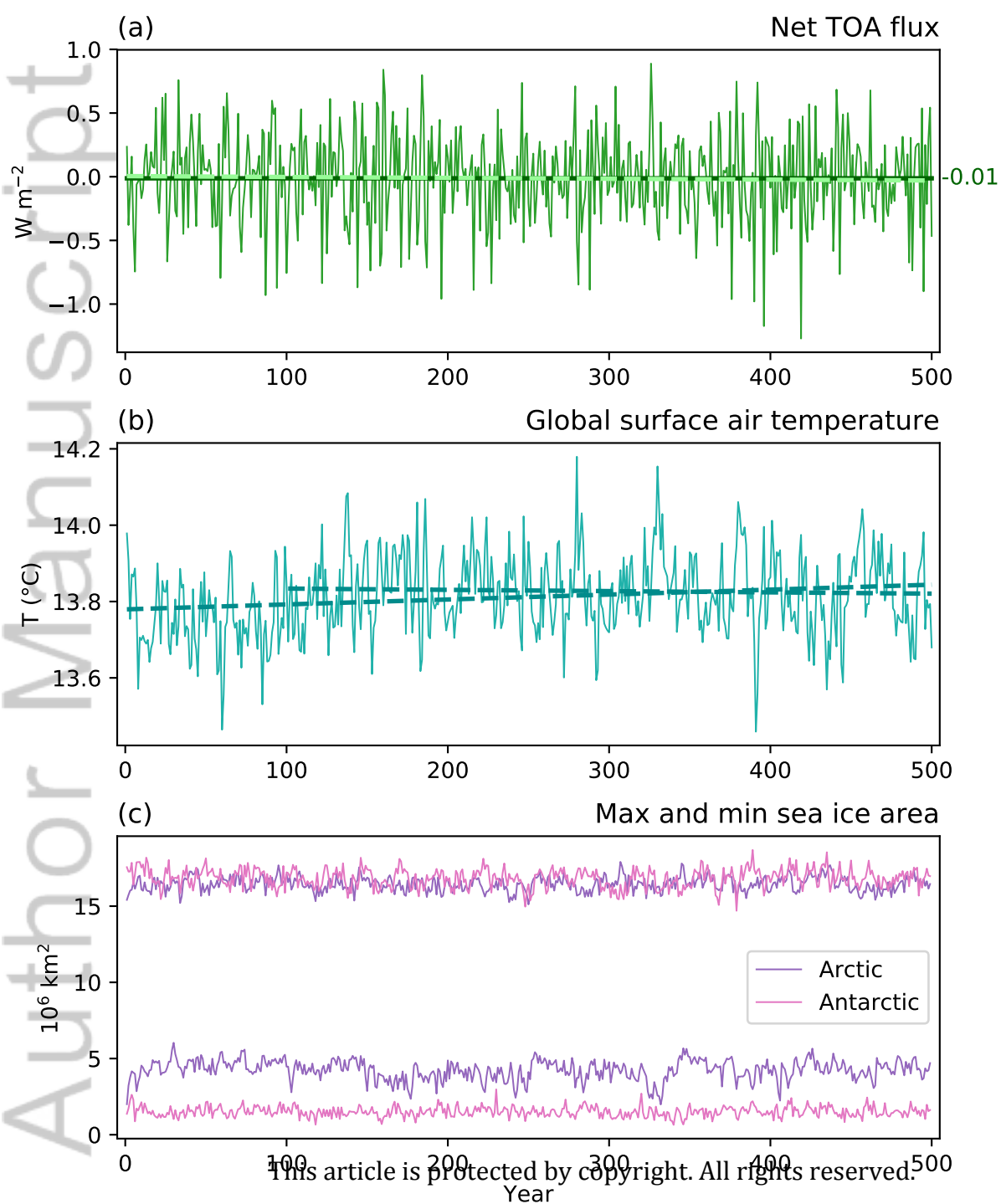


Figure 4.

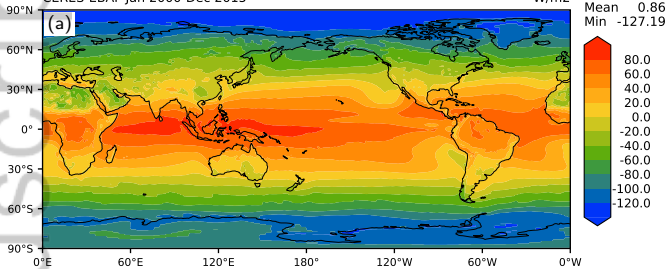
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Annual net TOA radiation

CERES-EBAF Jan 2000-Dec 2015

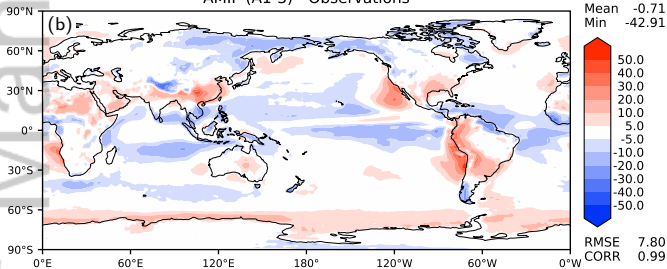
W/m²

Max 100.08
Mean 0.86
Min -127.19



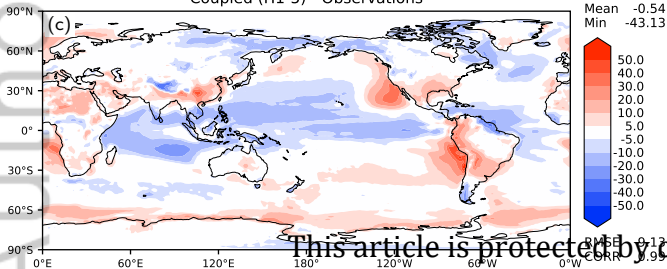
AMIP (A1-3) - Observations

Max 68.32
Mean -0.71
Min -42.91



Coupled (H1-5) - Observations

Max 65.55
Mean -0.54
Min -43.13



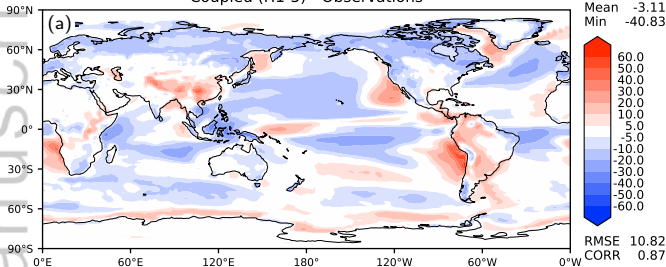
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Figure 5.

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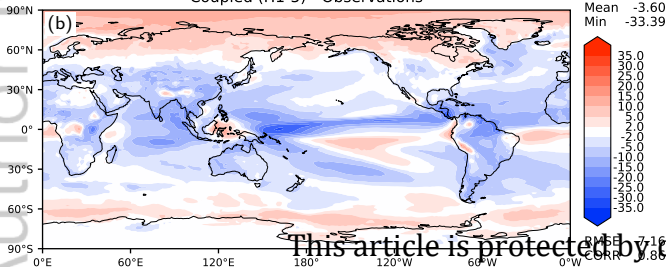
Annual SW cloud radiative effect

Coupled (H1-5) - Observations



Annual LW cloud radiative effect

Coupled (H1-5) - Observations



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Figure 6.

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Annual precipitation

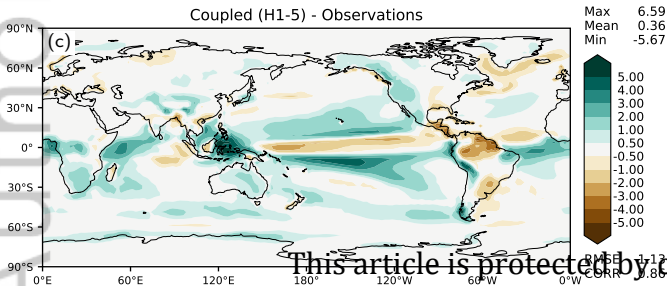
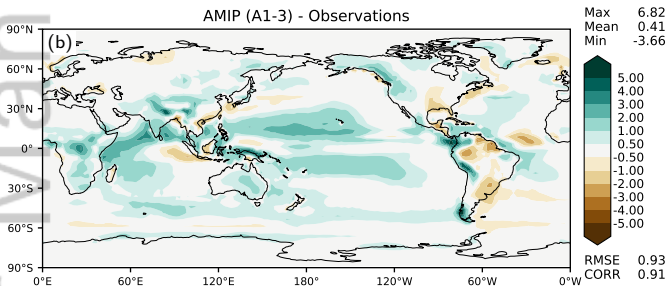
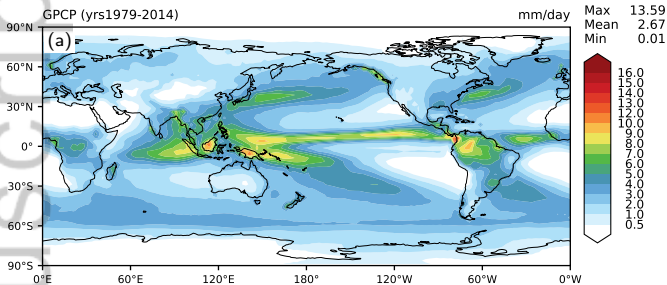


Figure 7.

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Annual temperature

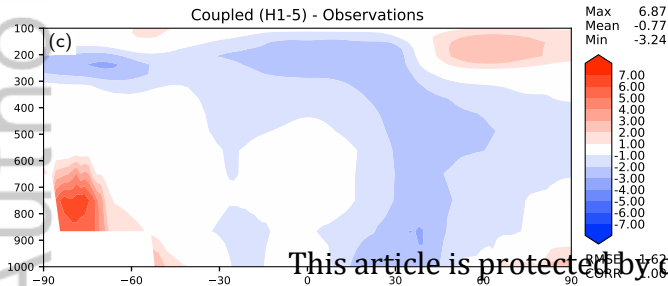
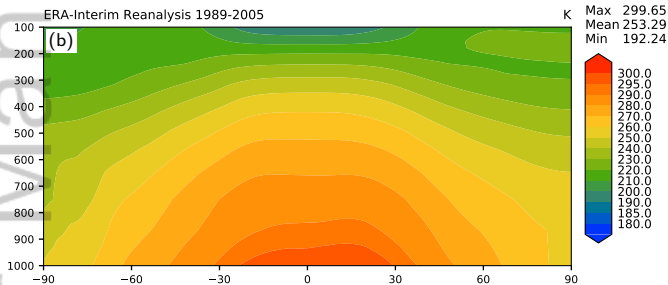
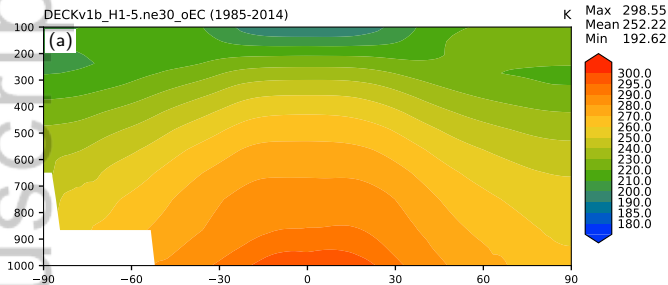


Figure 8.

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Annual zonal wind

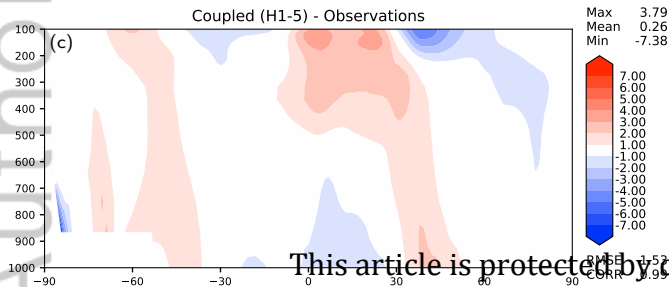
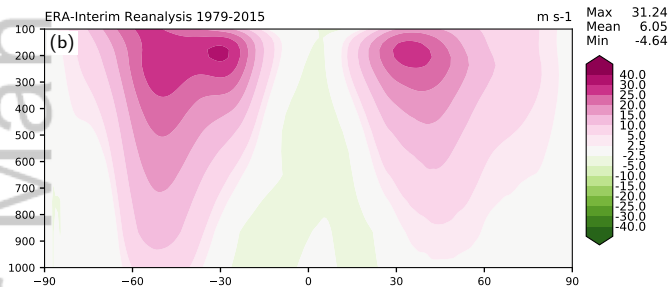
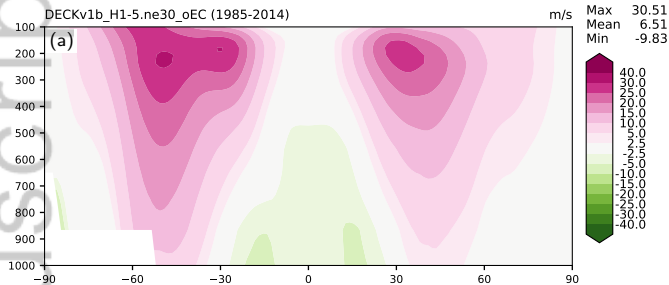
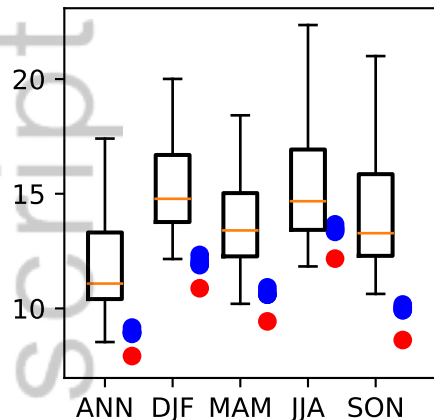


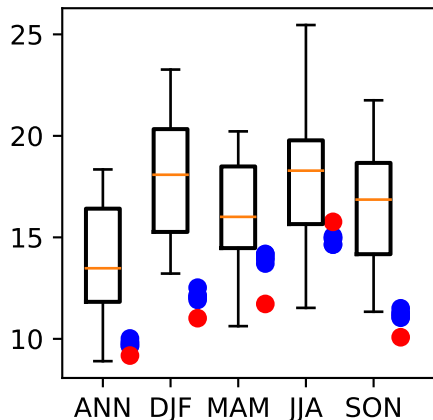
Figure 9.

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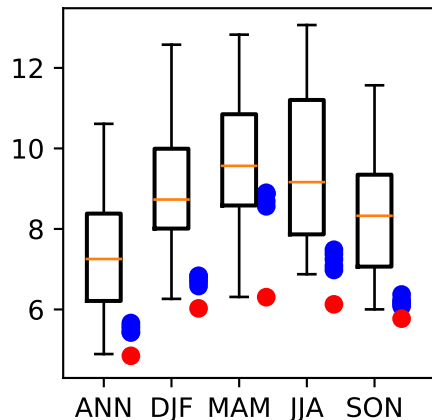
(a) Net TOA (W m^{-2})



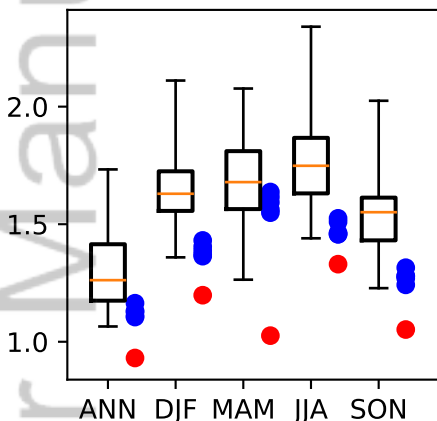
(b) SW CRE (W m^{-2})



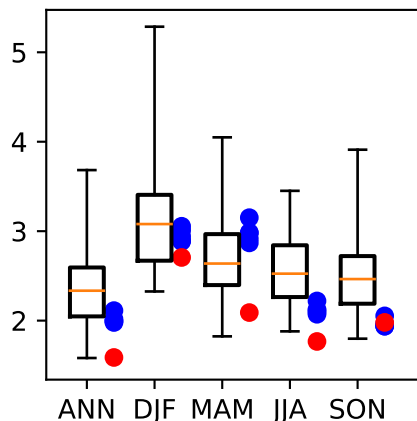
(c) LW CRE (W m^{-2})



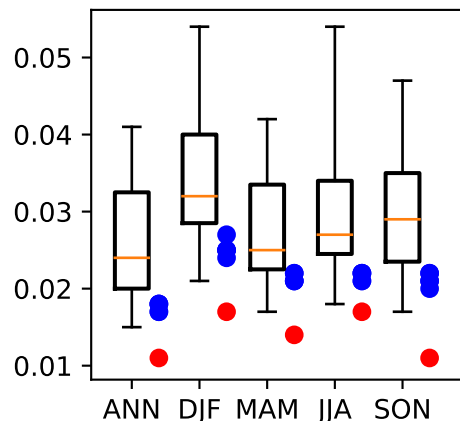
(d) prec (mm day^{-1})



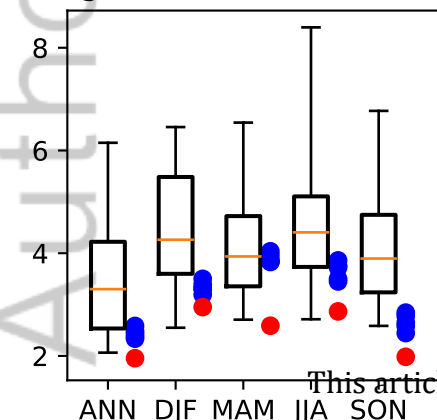
(e) tas (land, K)



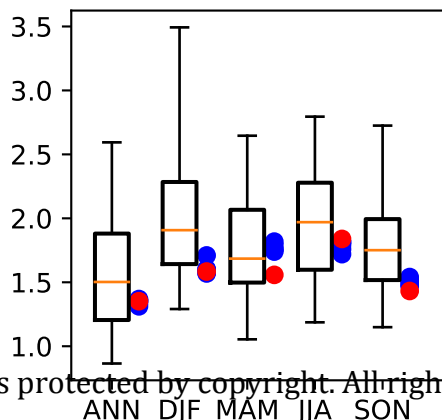
(f) τ_x (ocean, Pa)



(g) u200 (m s^{-1})



(h) u850 (m s^{-1})



(i) Zg-500 (m)

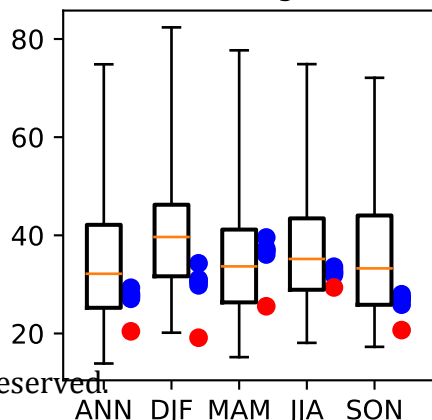
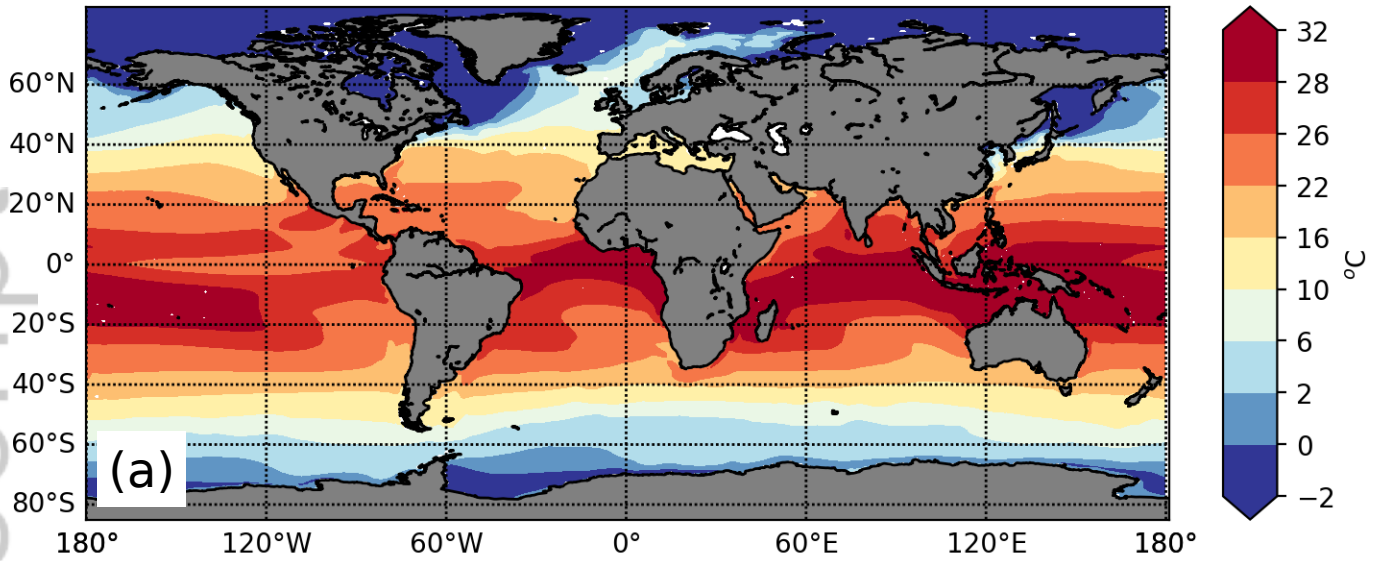


Figure 10.

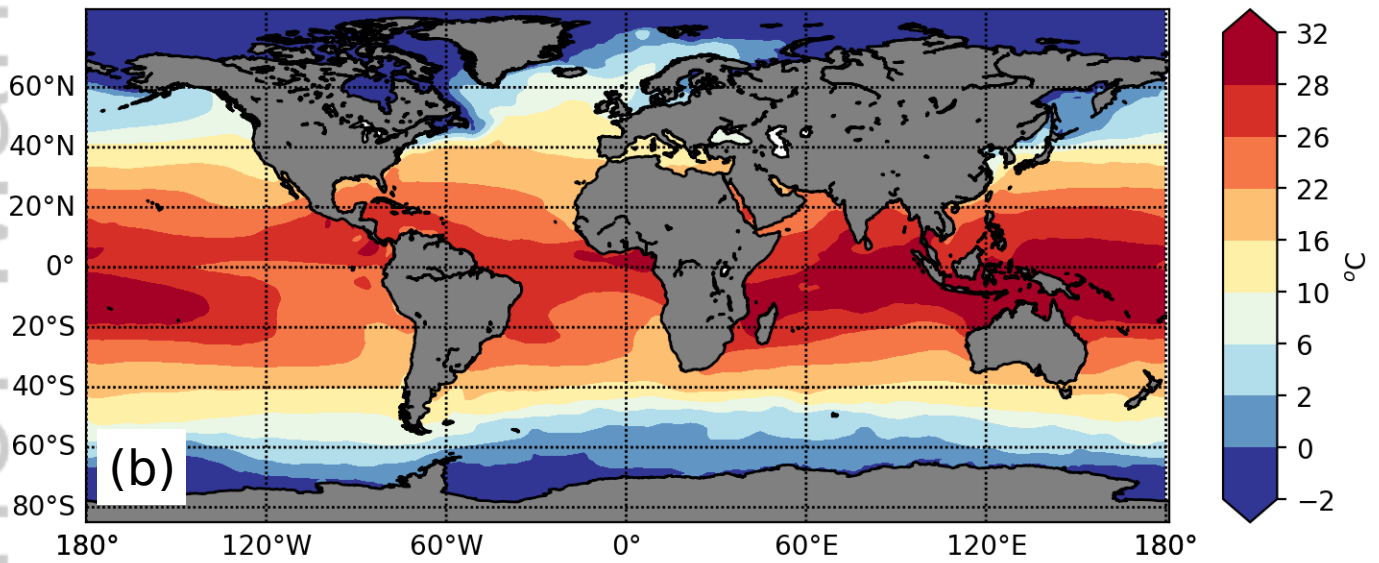
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Sea Surface Temperature (Annual Average)

E3SM Historical Ensemble Average



Observations



Model - Observations

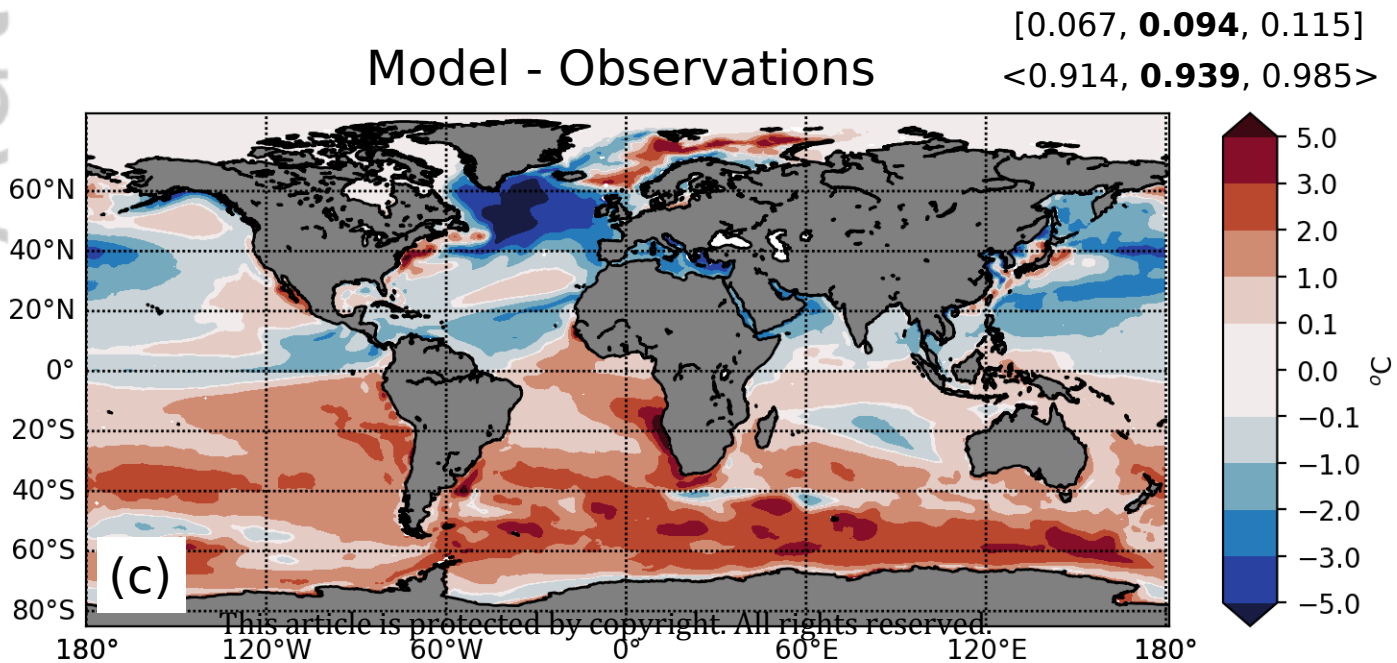
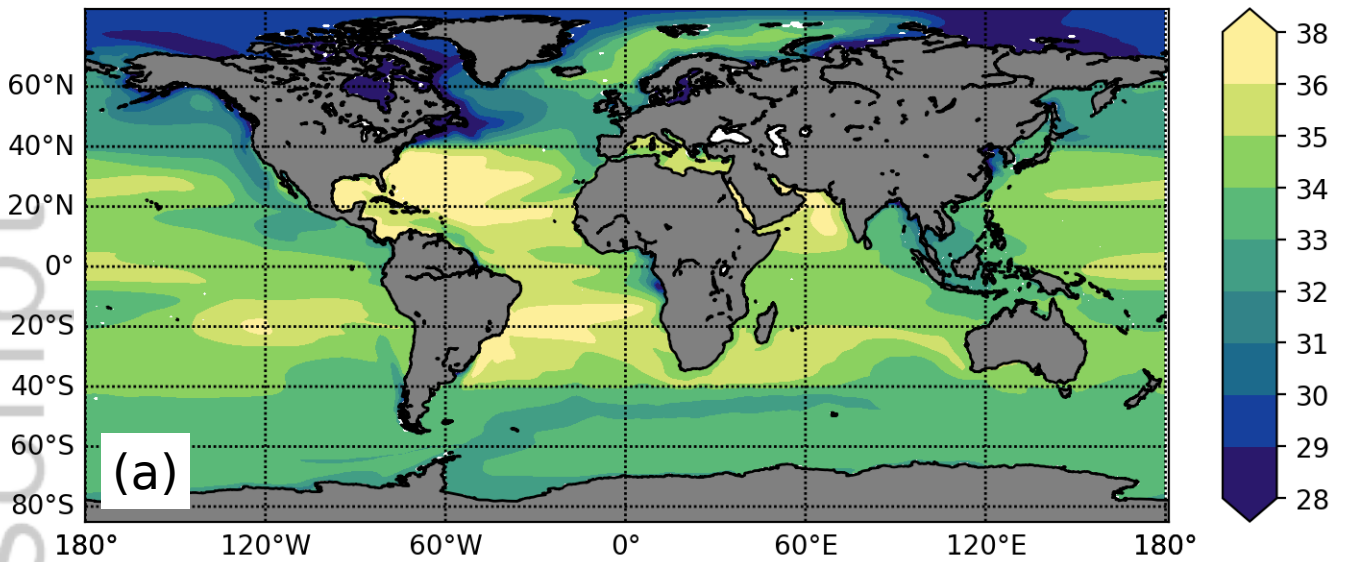


Figure 11.

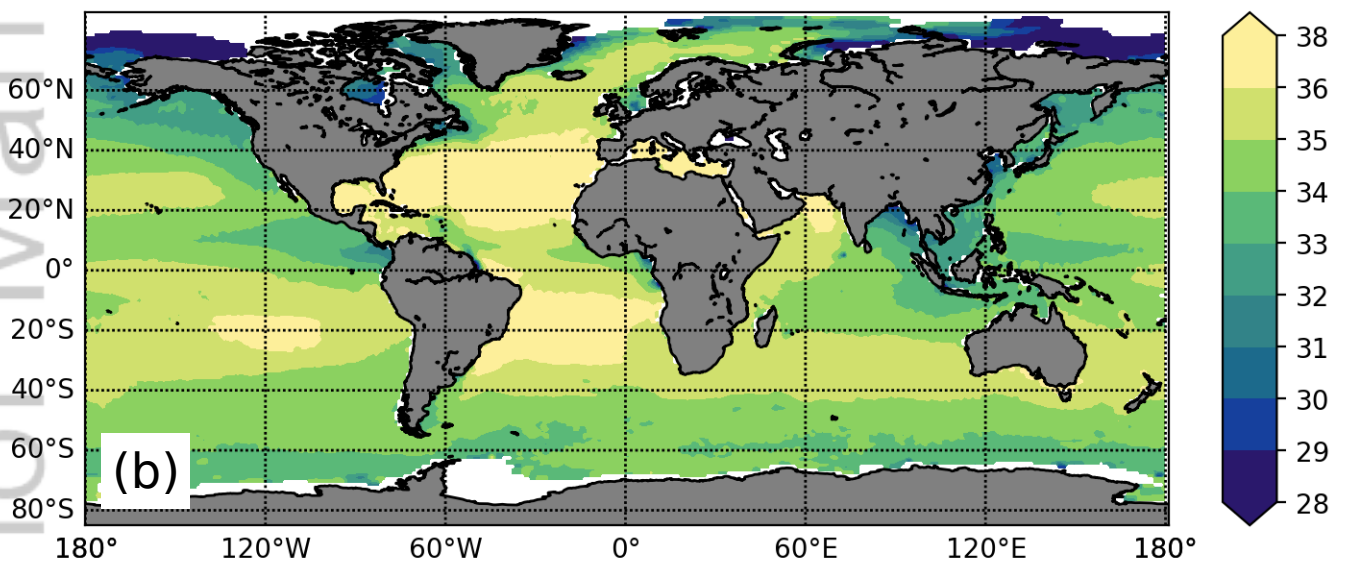
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Sea Surface Salinity (Annual Average)

E3SM Historical Ensemble Average

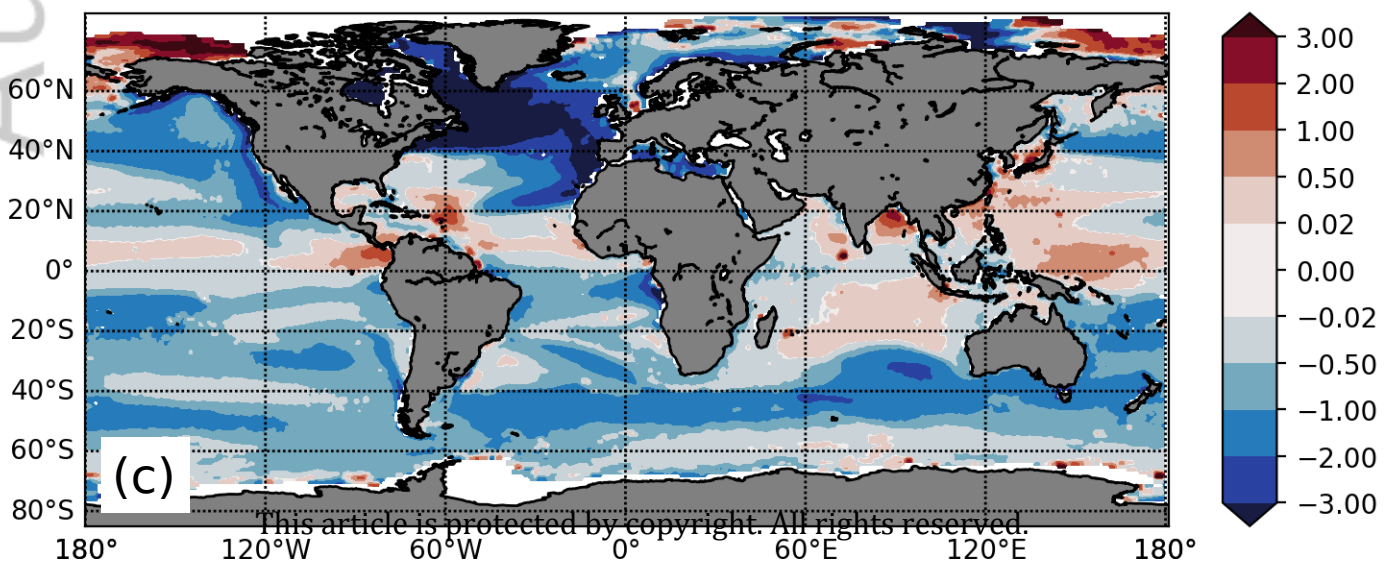


Observations



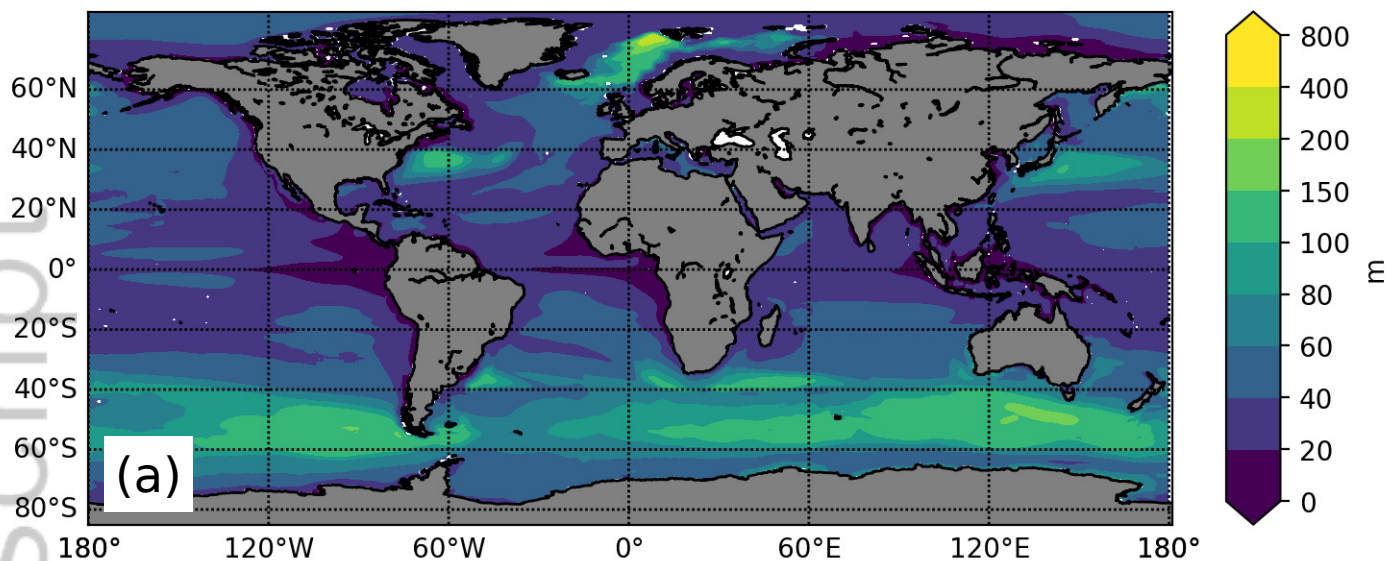
[-0.72, **-0.71**, -0.69]
<1.10, **1.12**, 1.14>

Model - Observations

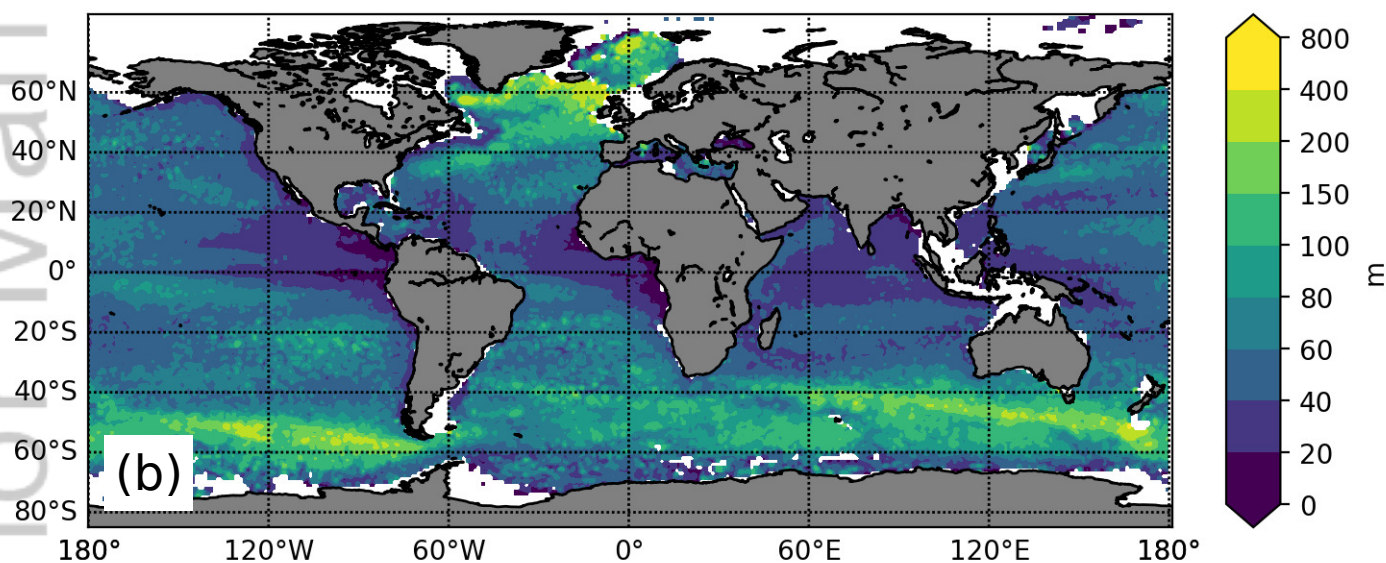


Mixed Layer Depth (Annual Average)

E3SM Historical Ensemble Average



Observations



Model - Observations

[-19.28, -18.54, -18.22]

<41.36, **41.74**, 42.47>

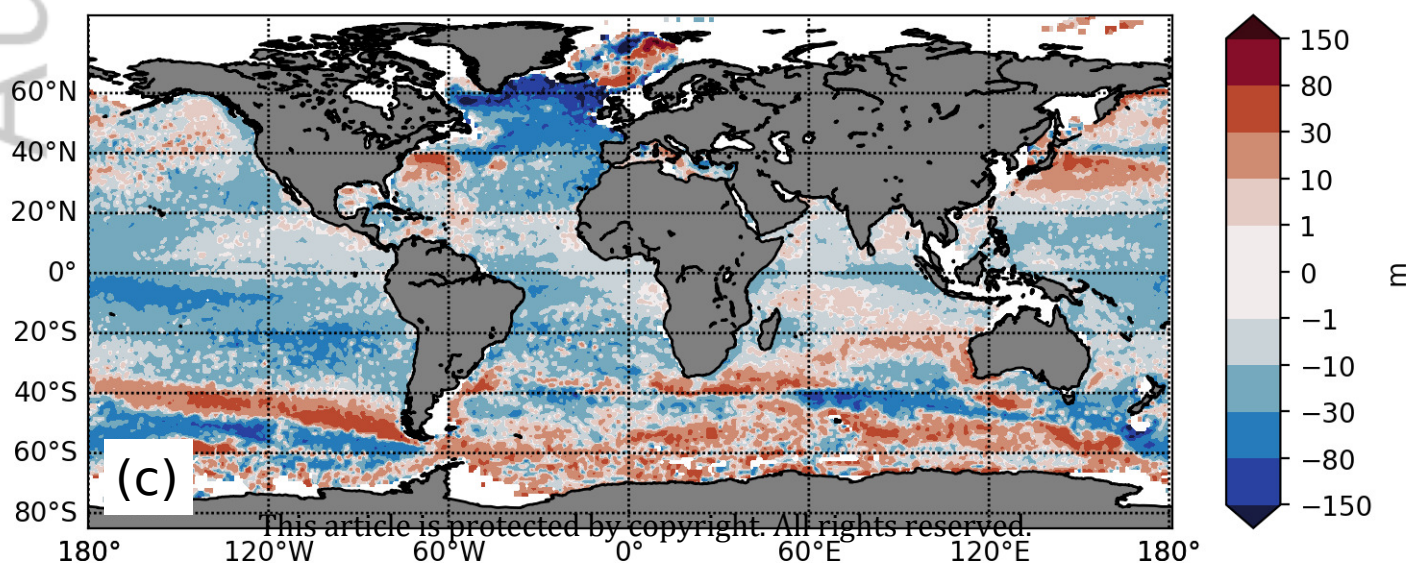
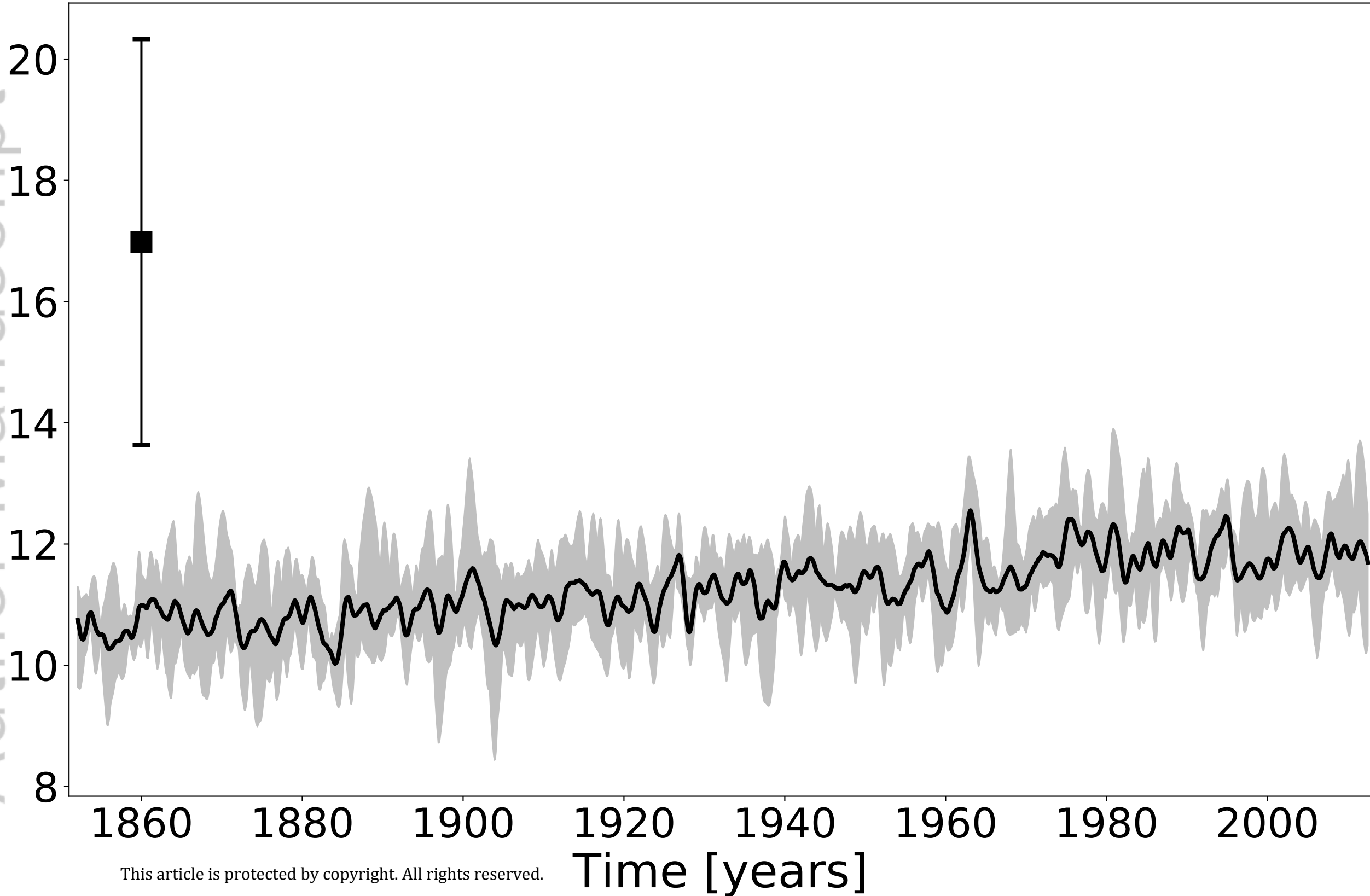
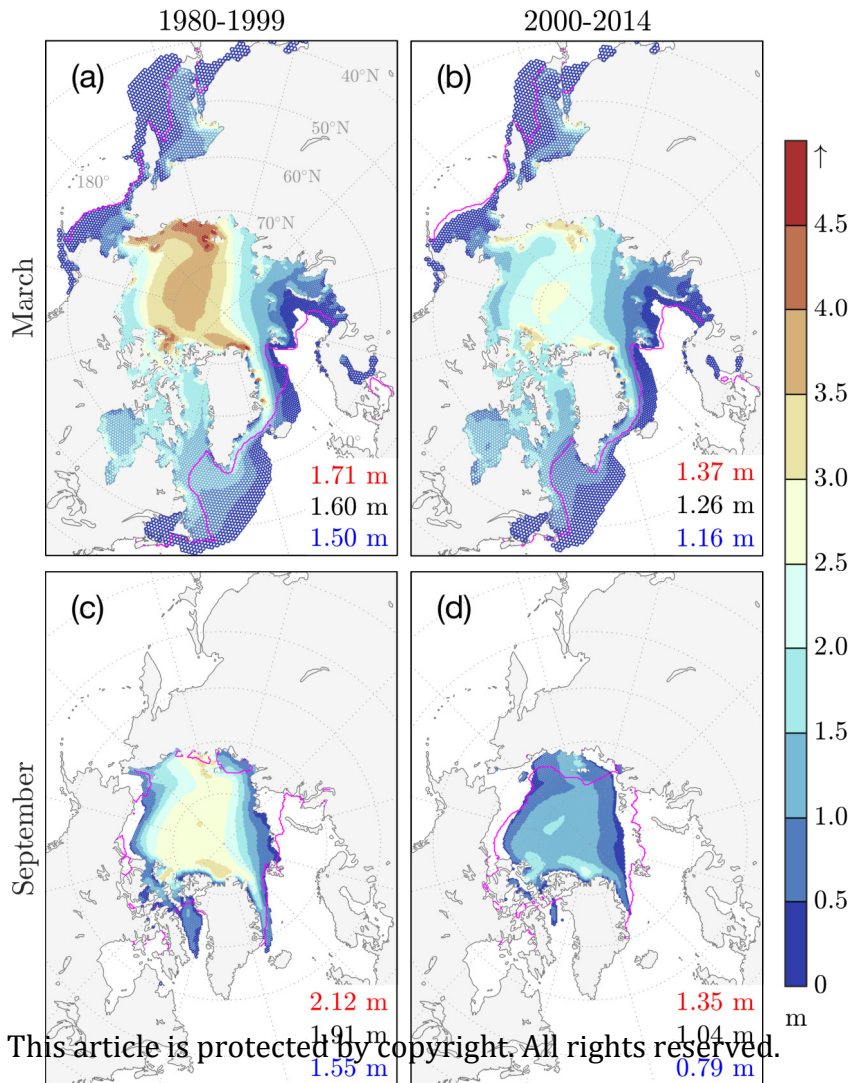


Figure 13.

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Figure 15.

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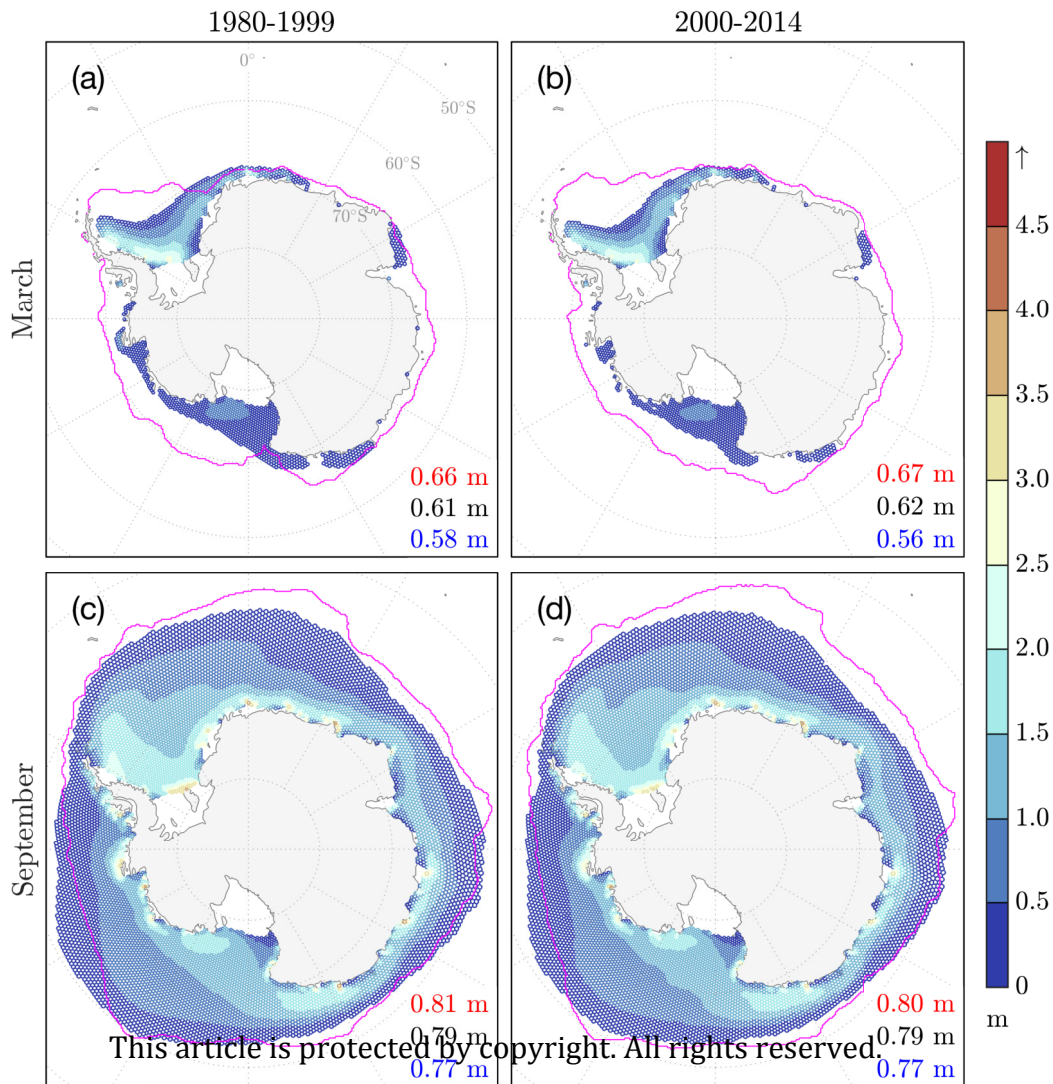


Figure 16.

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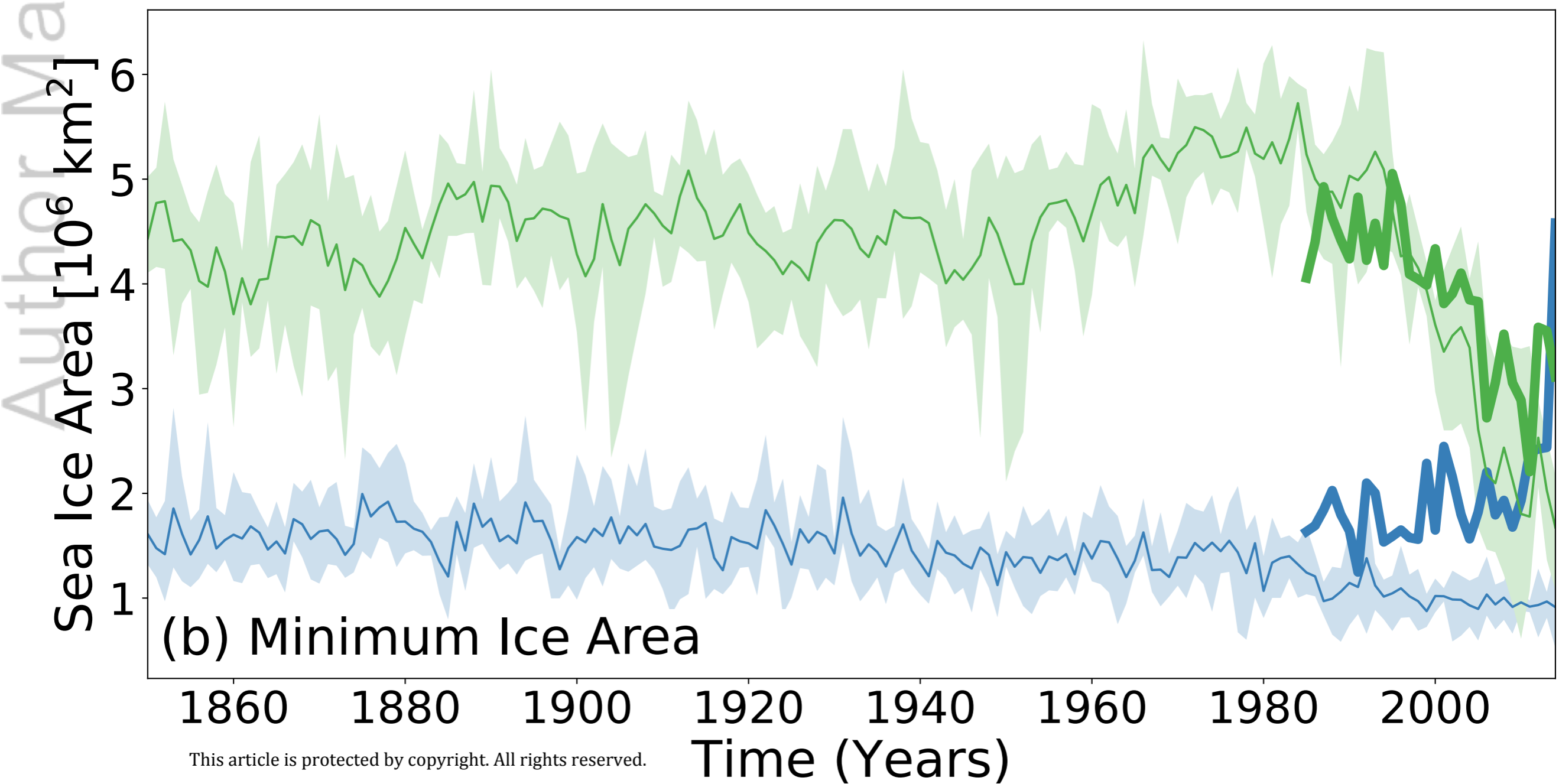
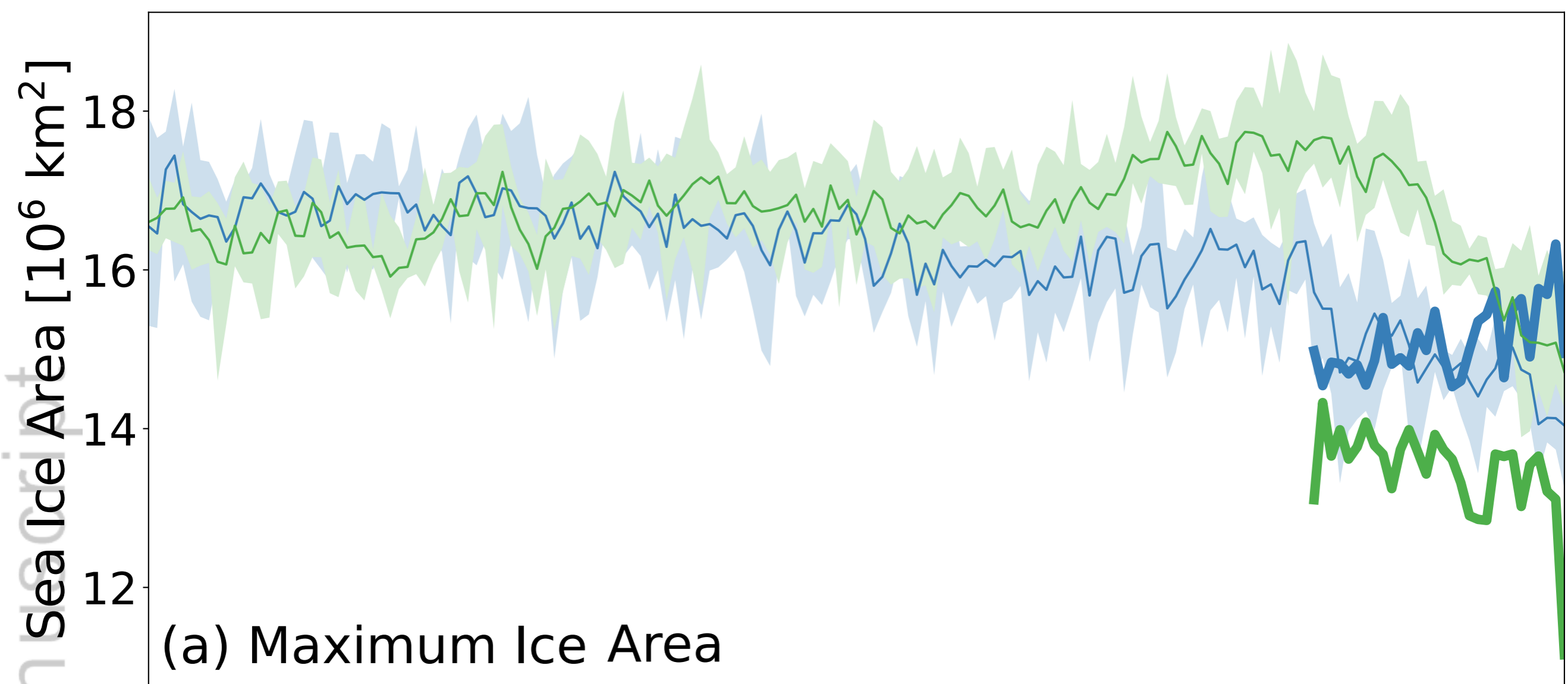
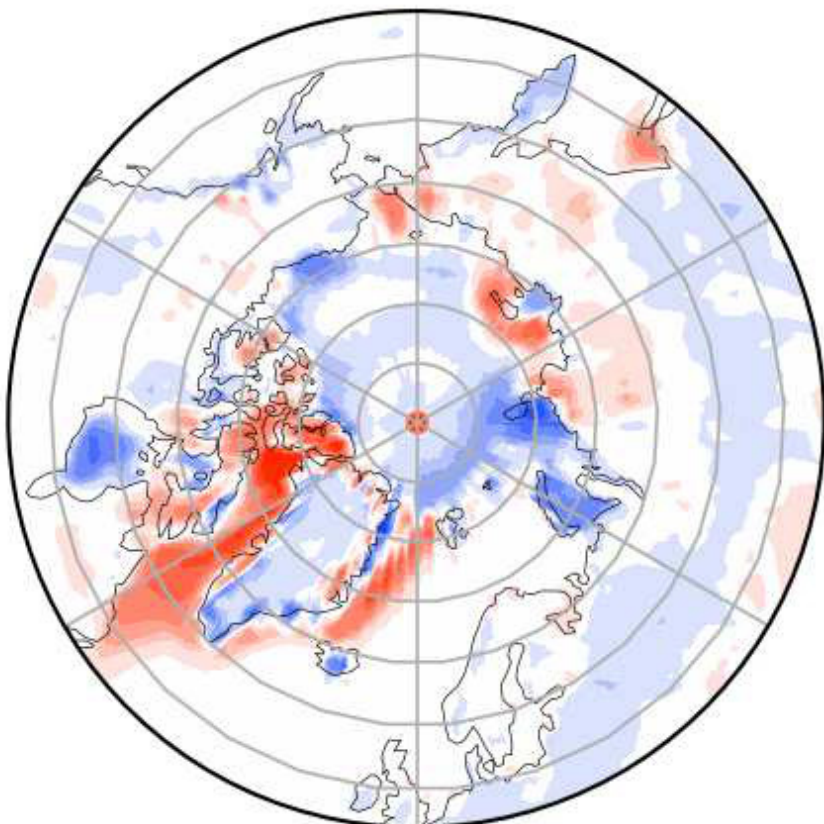


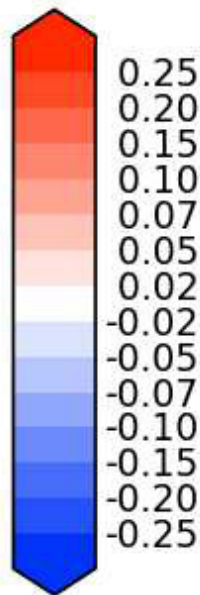
Figure 17.

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Model - Obs



Max 0.49
Mean -0.00
Min -0.32

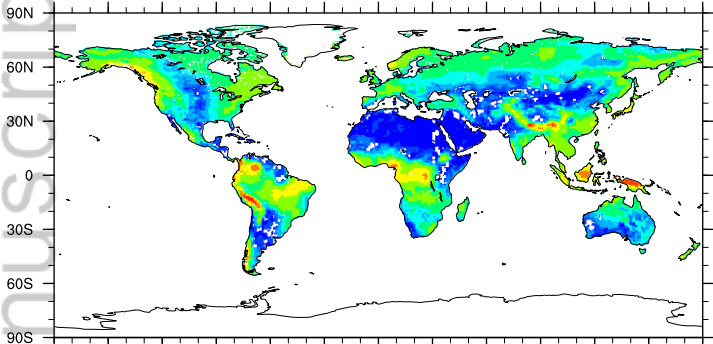


RMSE 0.05
CORR 0.96

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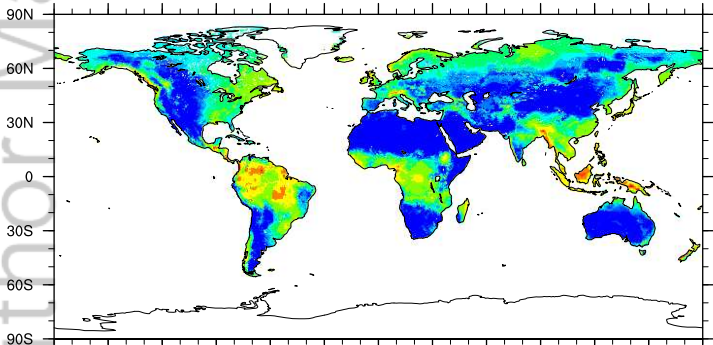
ELM total runoff

mm/day



GRDC total runoff

mm/day

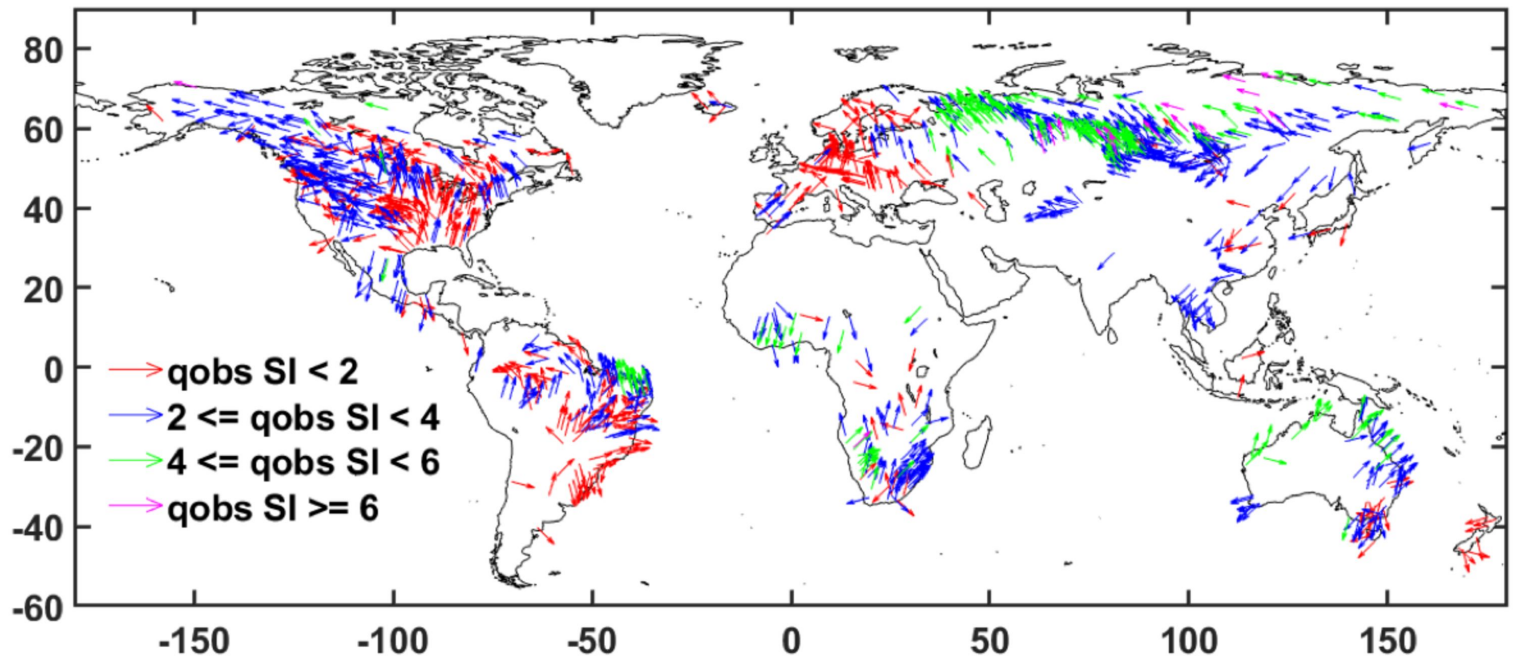


180 150W 120W 90W 60W 30W 0 30E 60E 90E 120E 150E 180

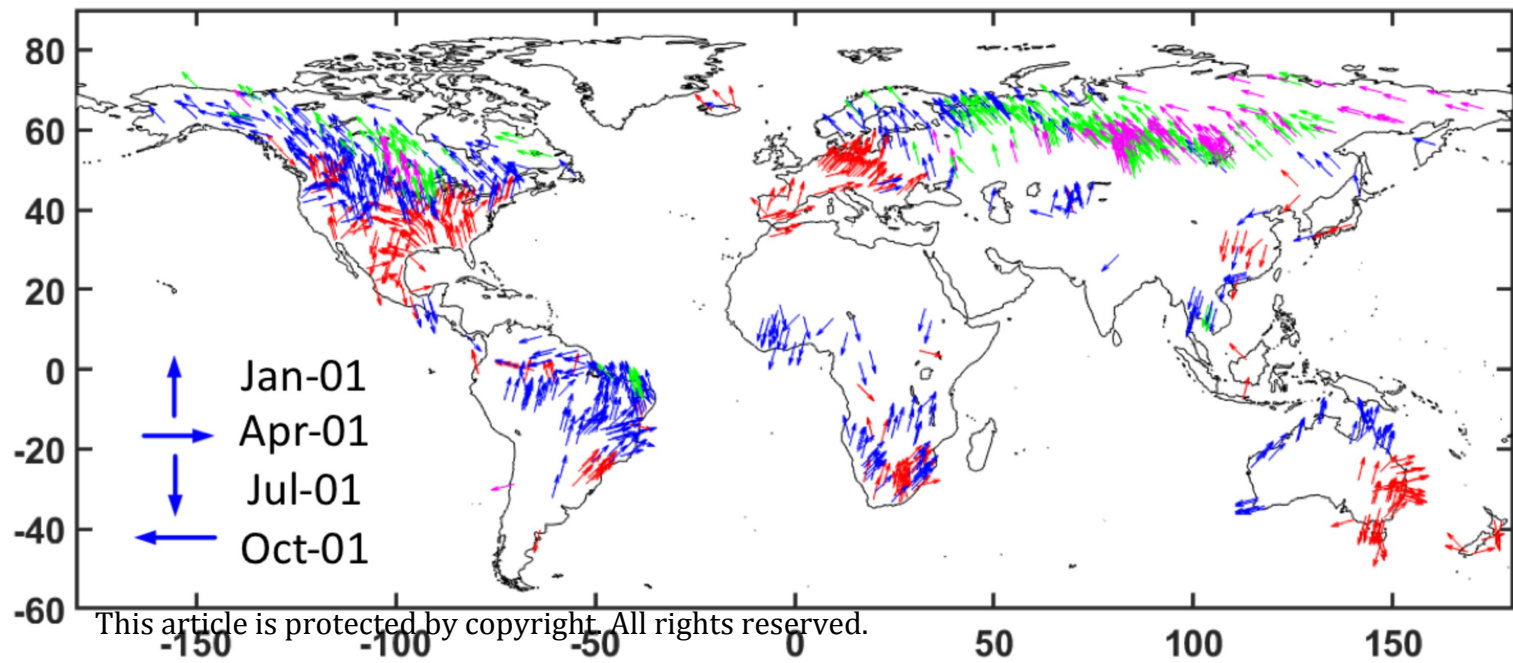


-1 0 0.01 0.05 0.1 0.25 0.5 1 2.5 5 7.5 10

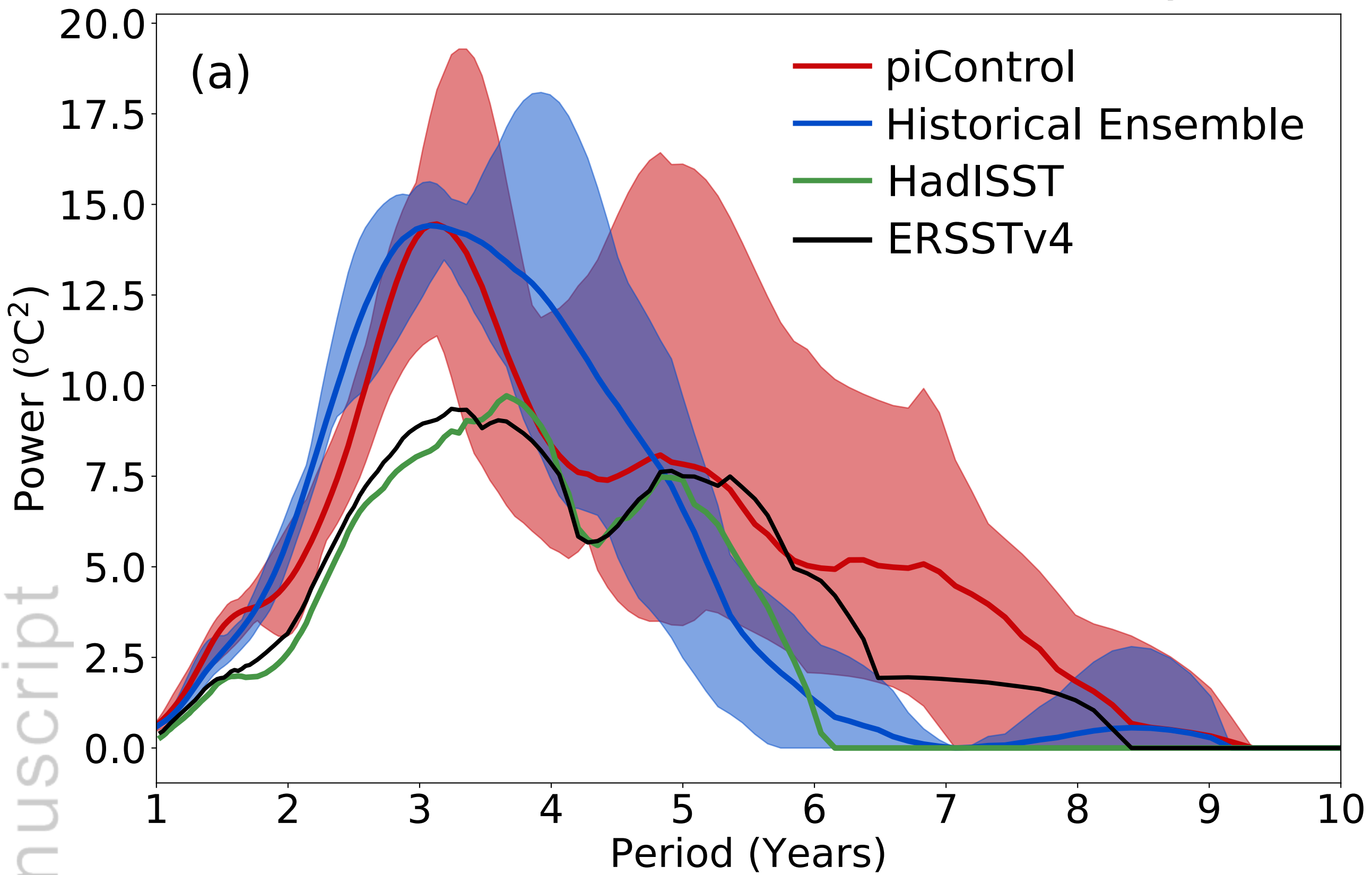
Observed



Simulated



E3SM Nino 3.4 SST Wavelet Analysis



CESM-LE Nino 3.4 SST Wavelet Analysis

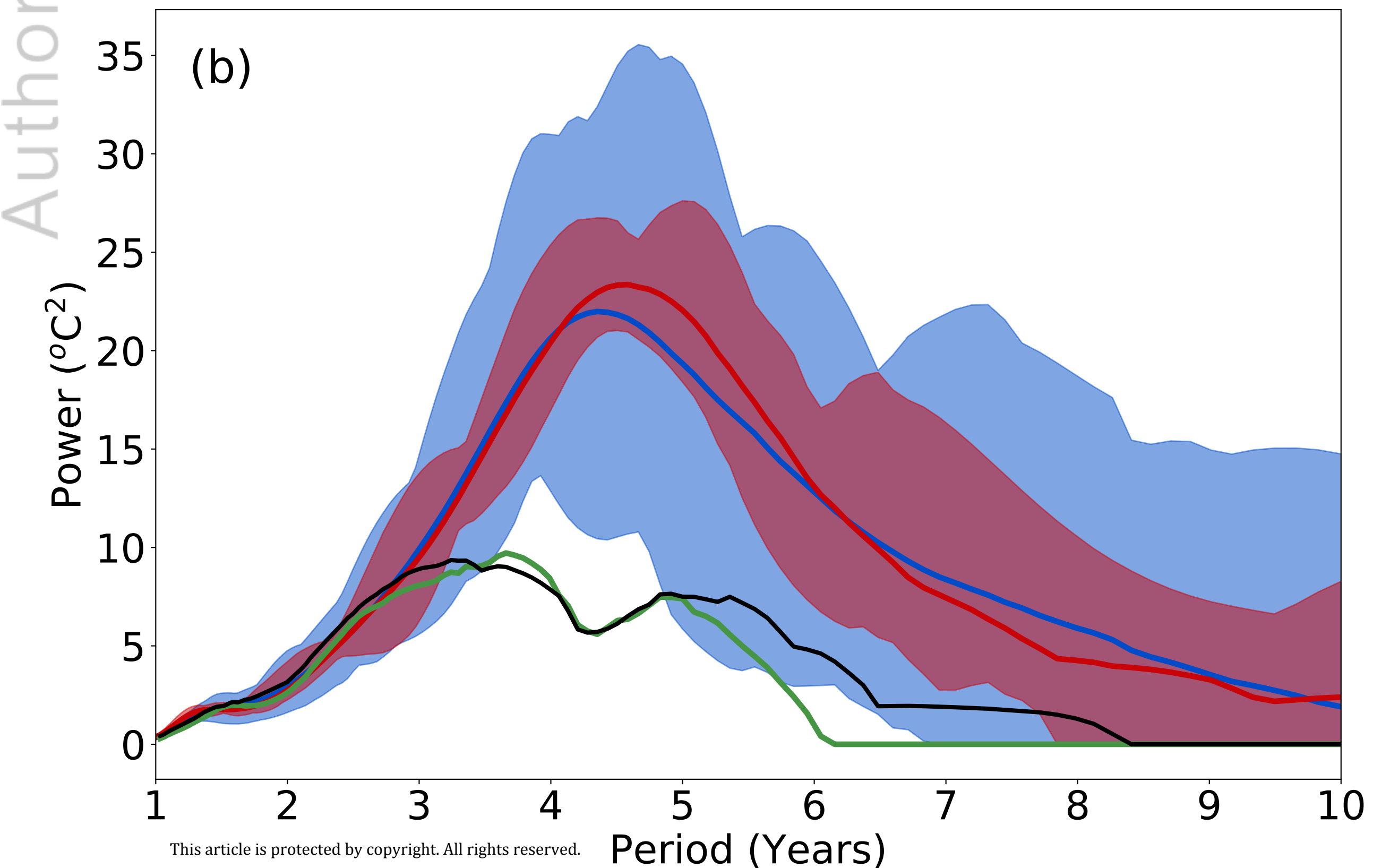
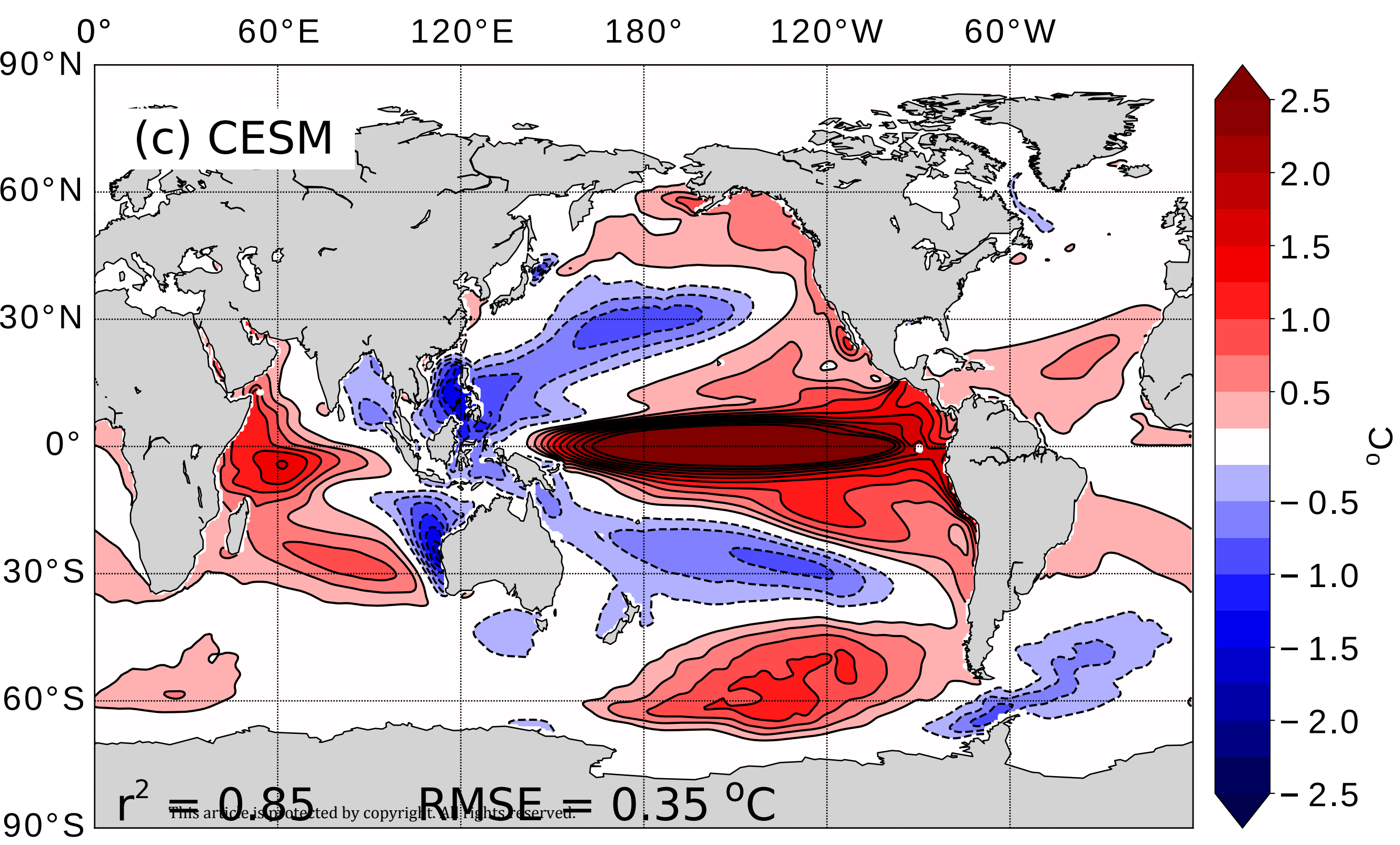
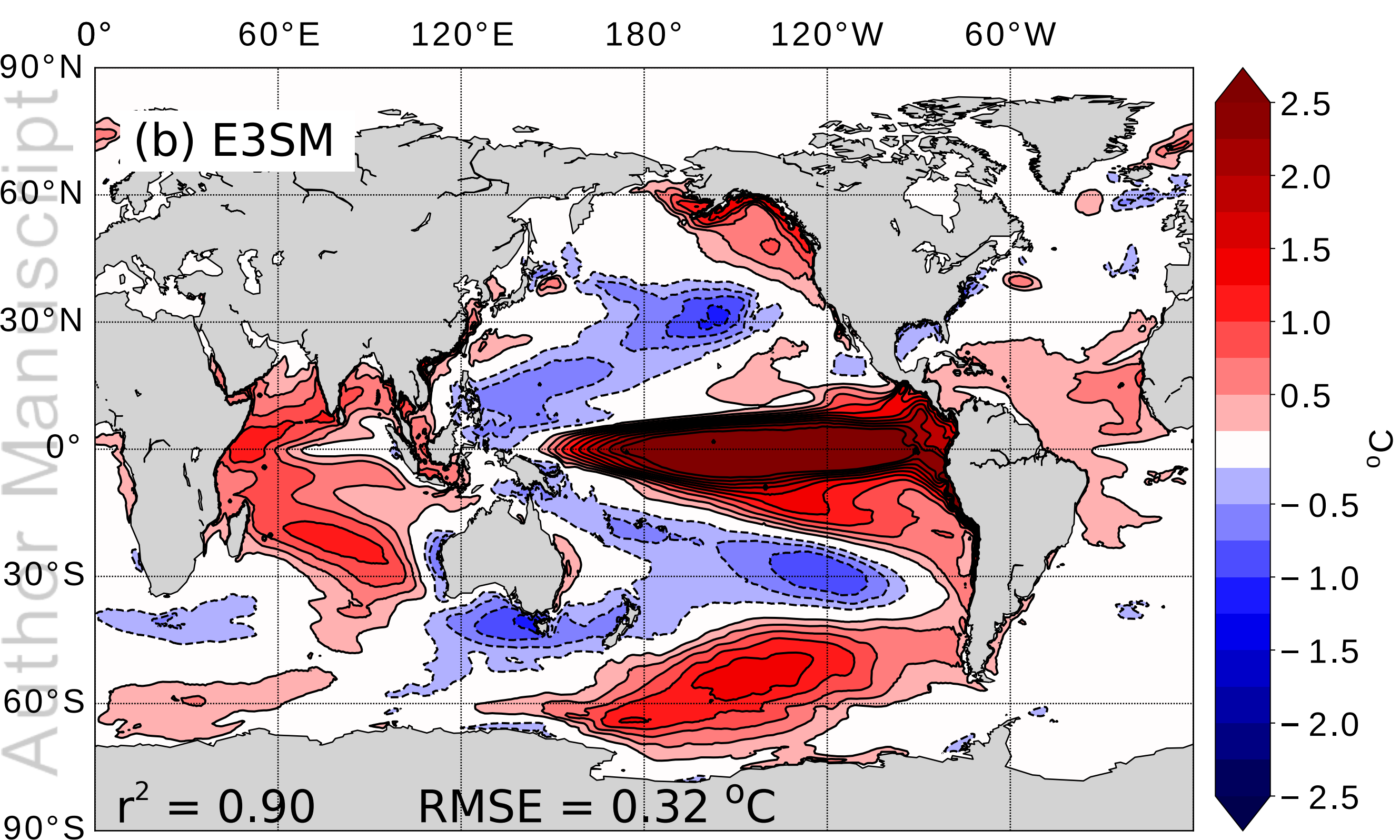
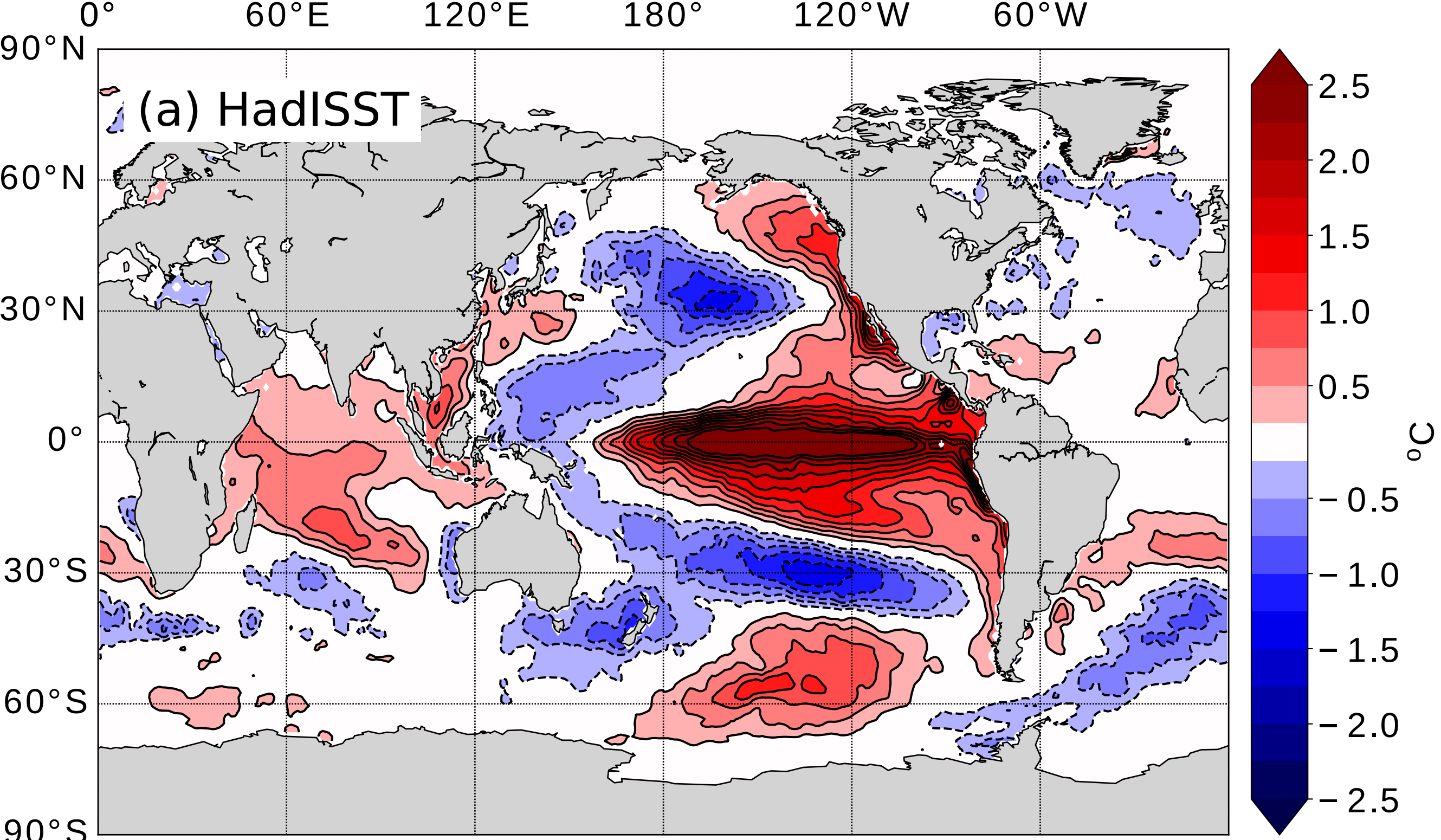


Figure 21.

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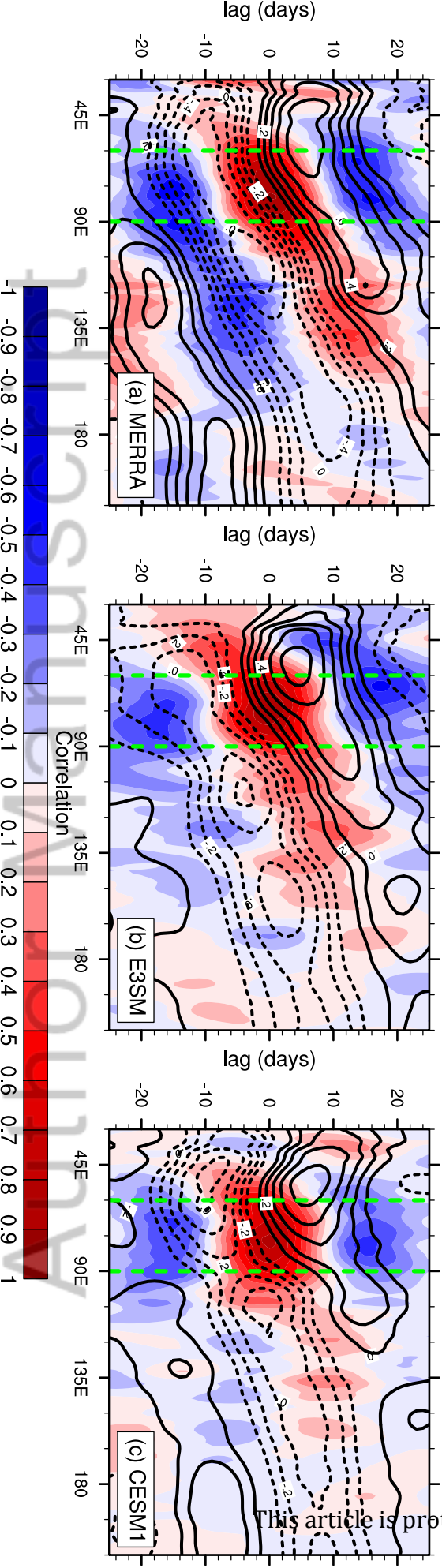
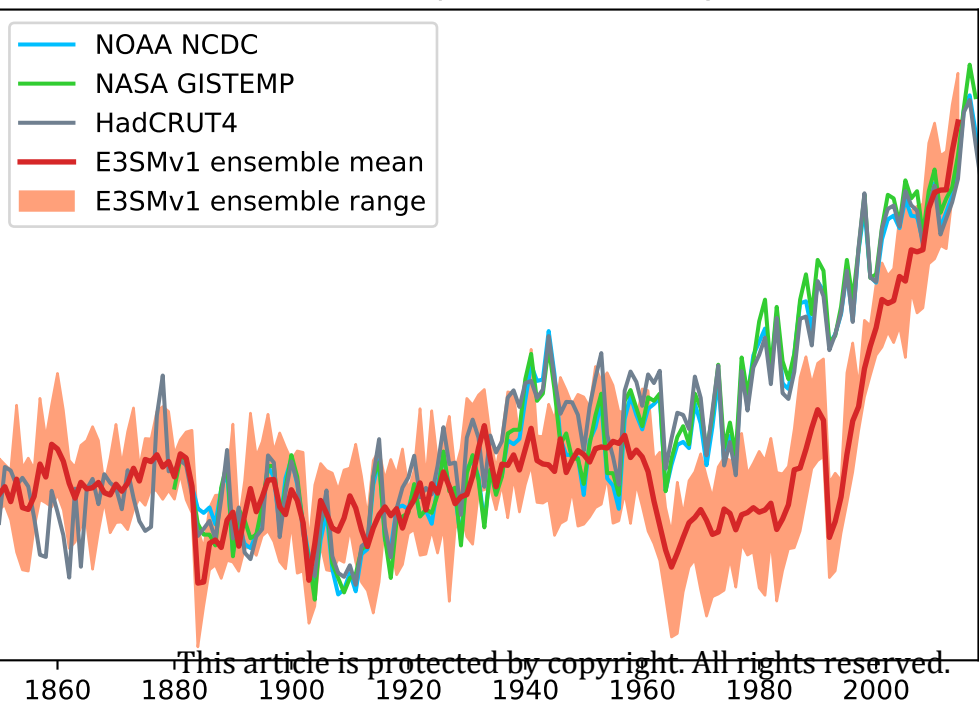


Figure 23.

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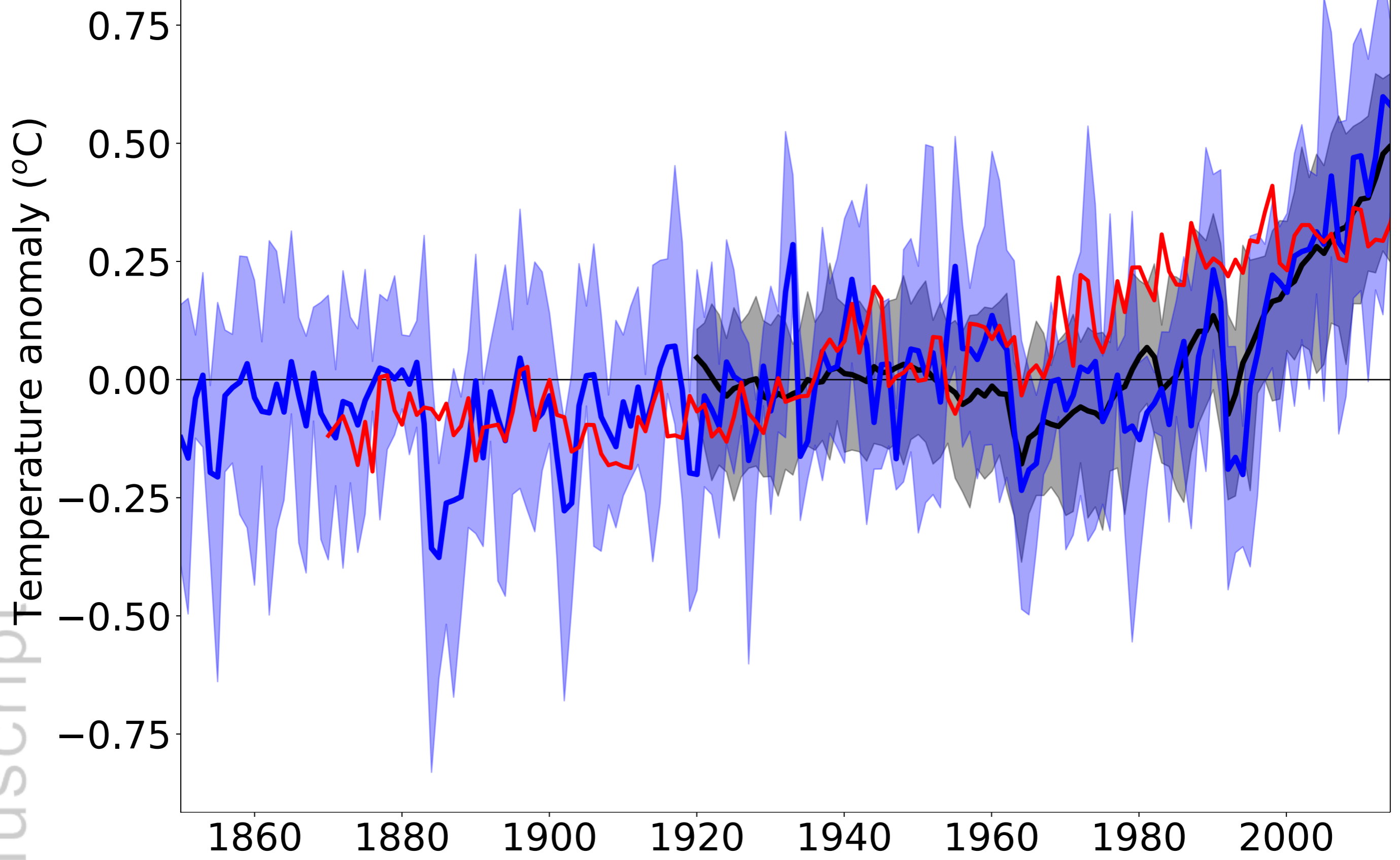
Global surface air temperature anomaly (vs 1880-1909)

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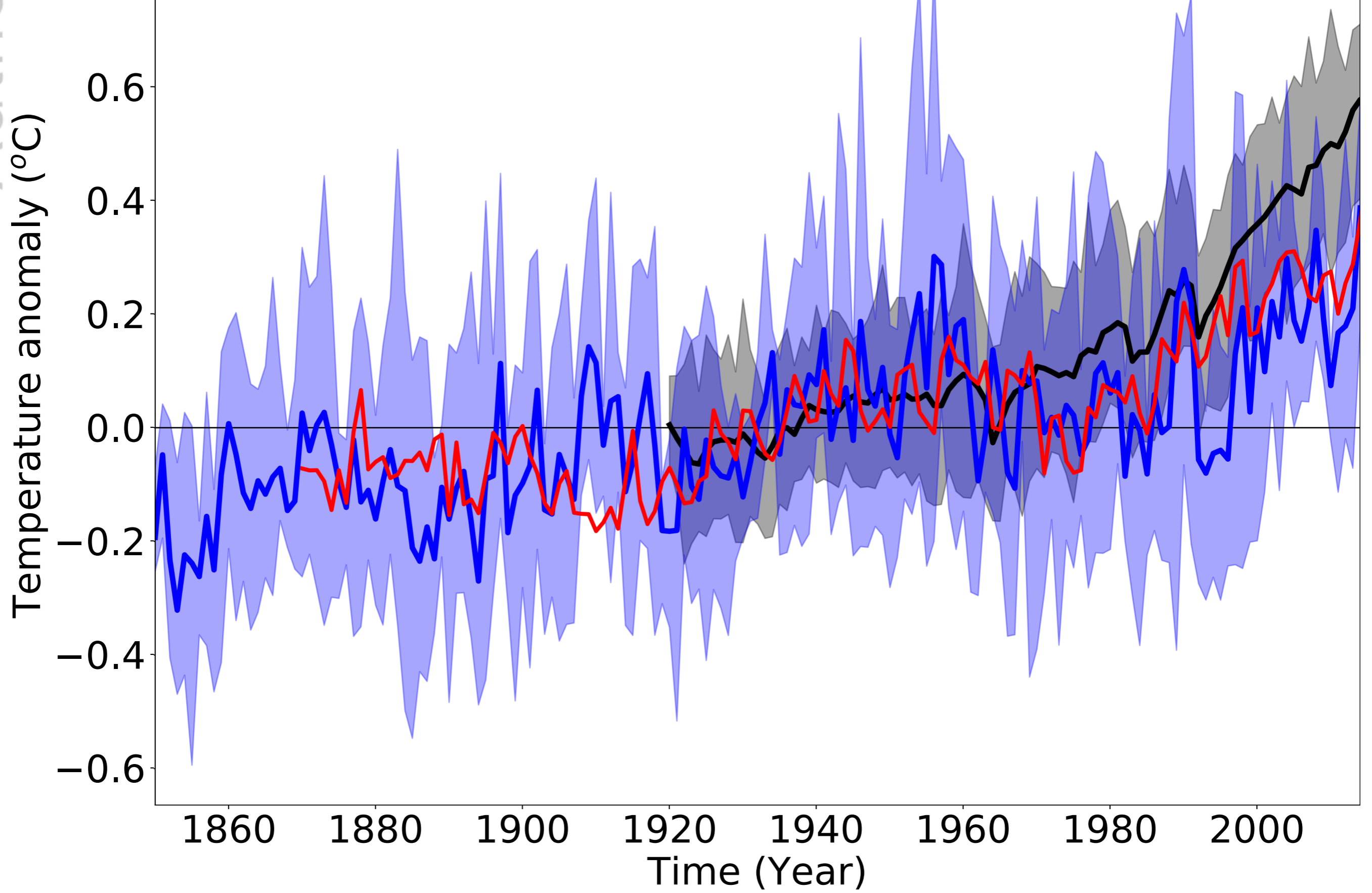


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(a) Northern Hemisphere



(b) Southern Hemisphere



Effective Radiative Forcing (relative to 1850)

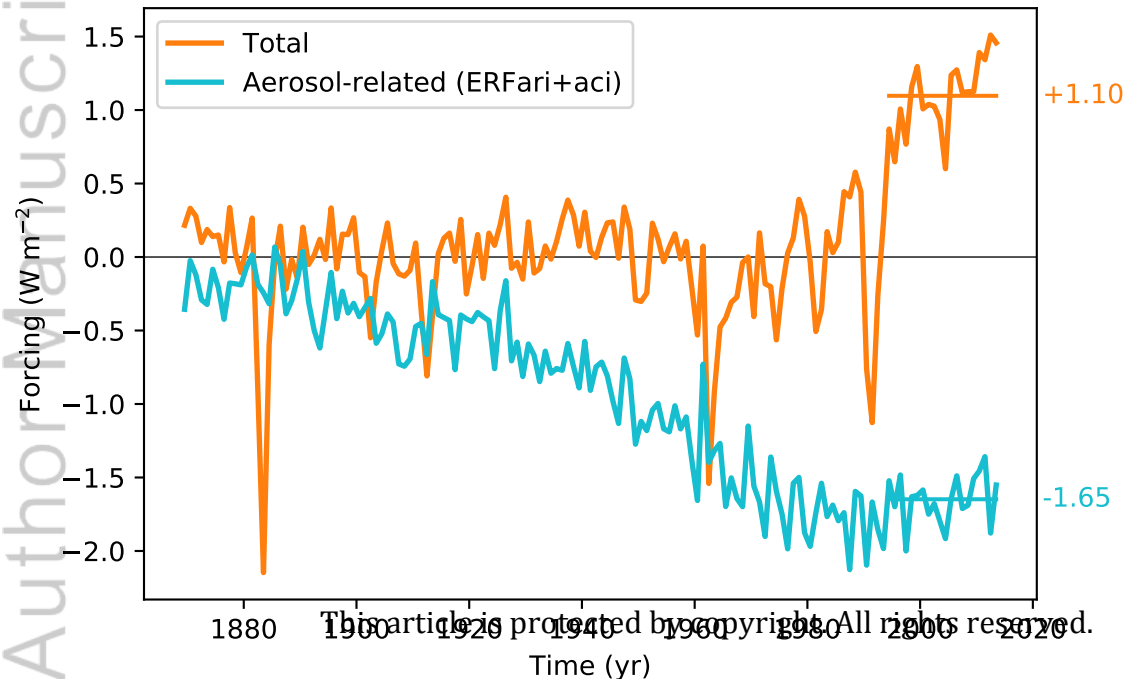
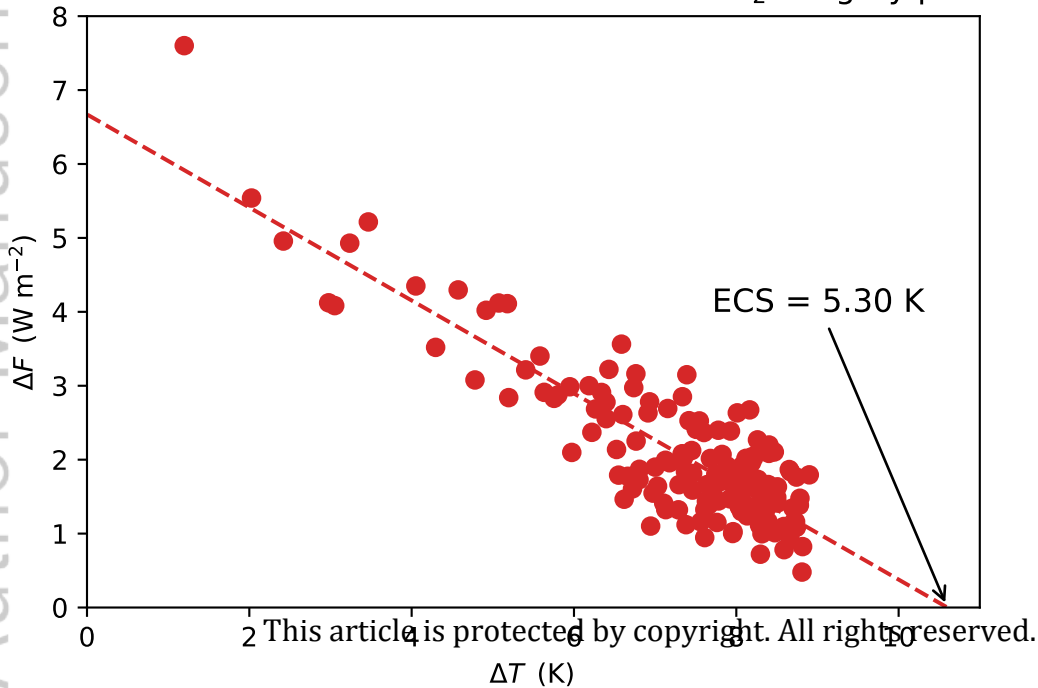
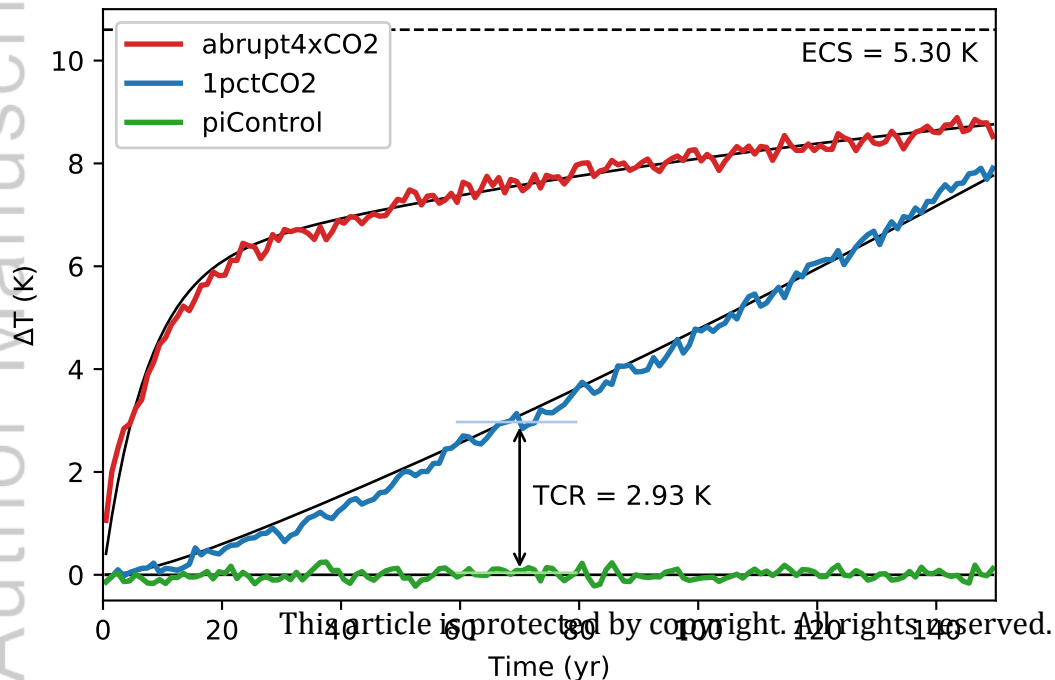


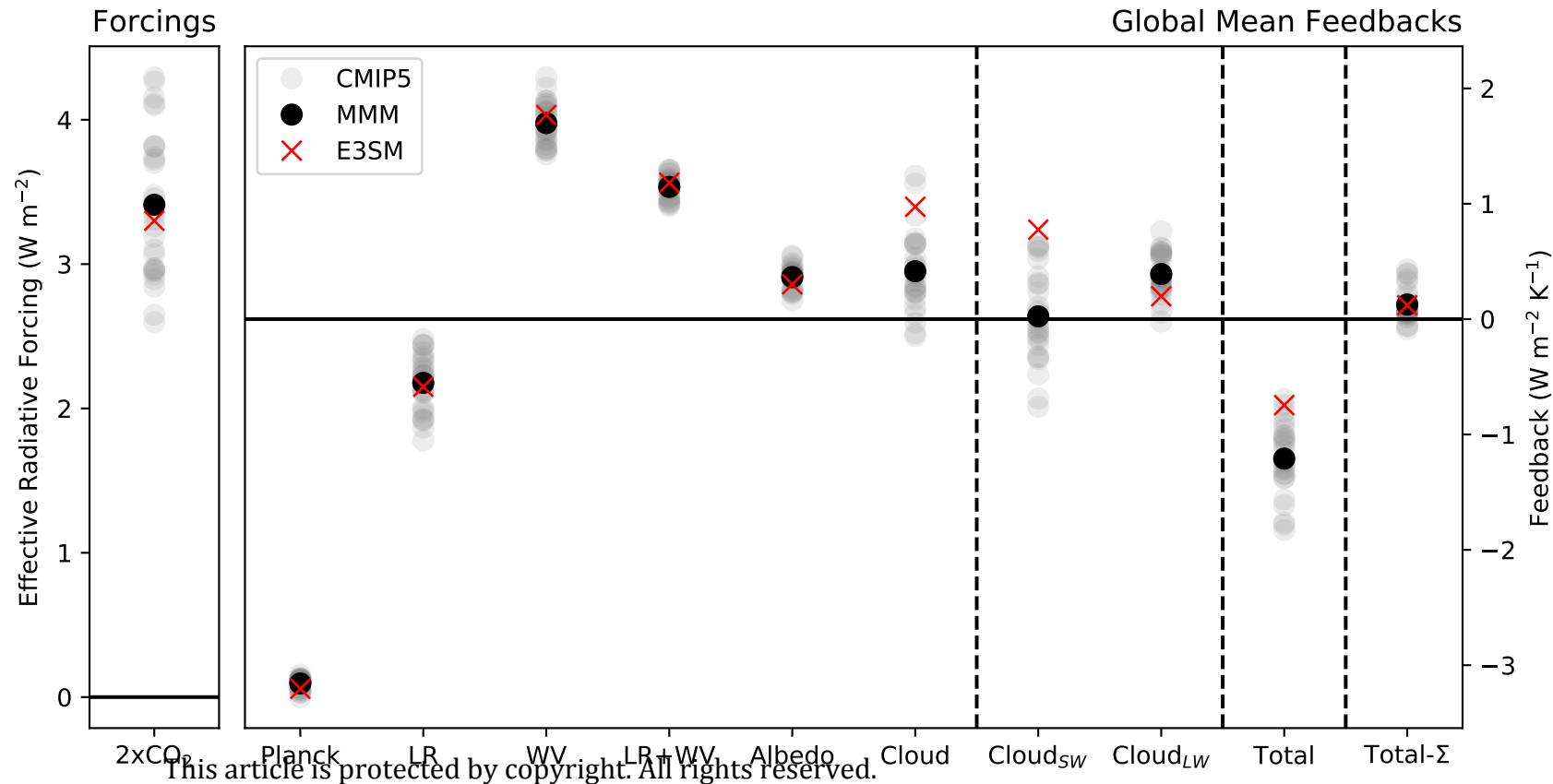
Figure 26.

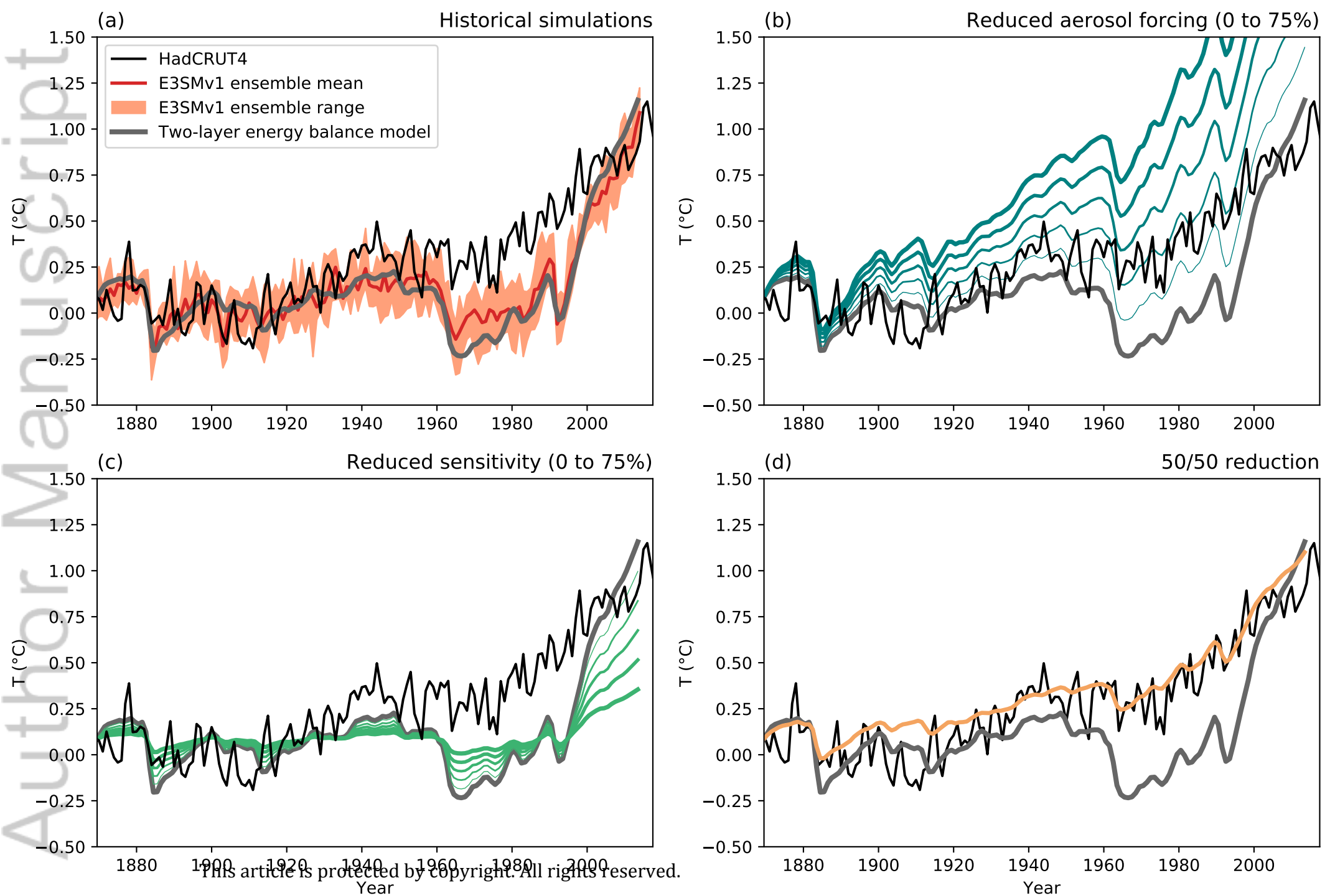
Author Manuscript

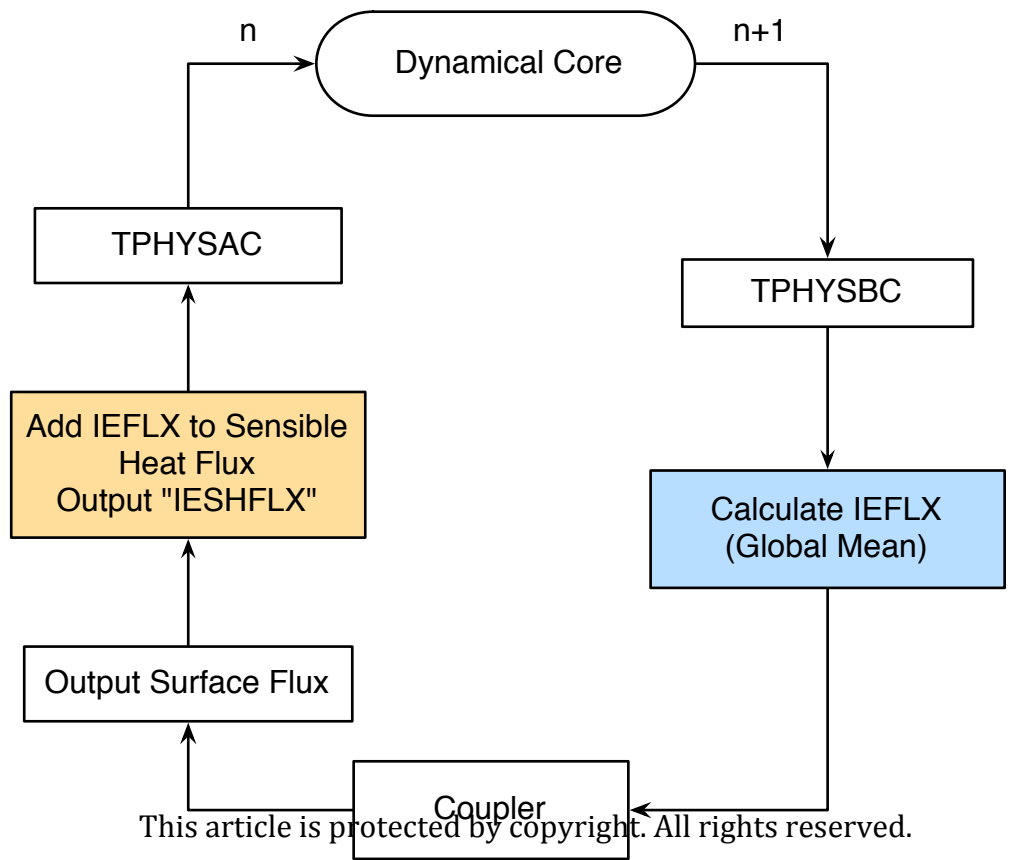
4xCO₂ Gregory plot

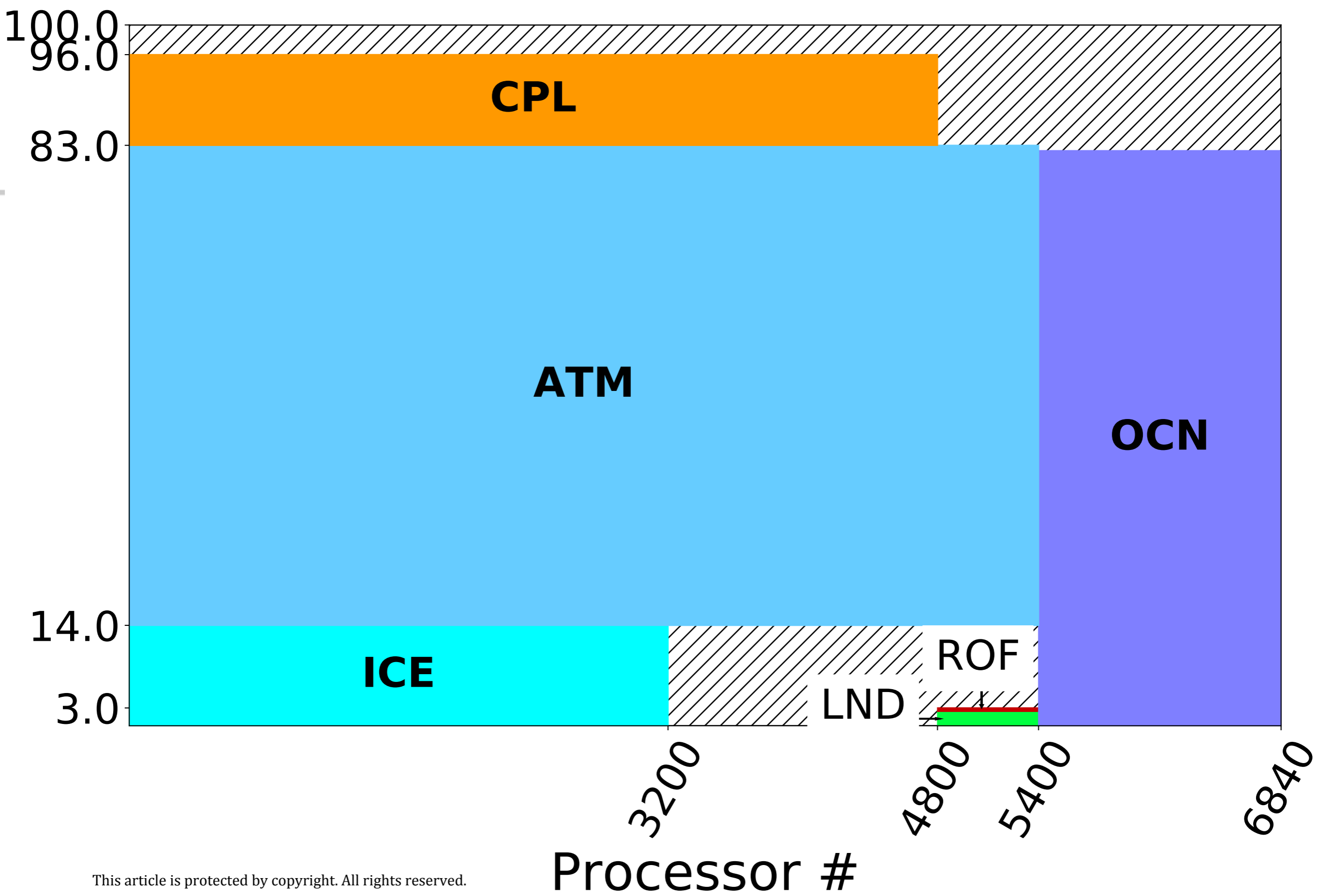
Idealized CO₂ Forcing Simulations



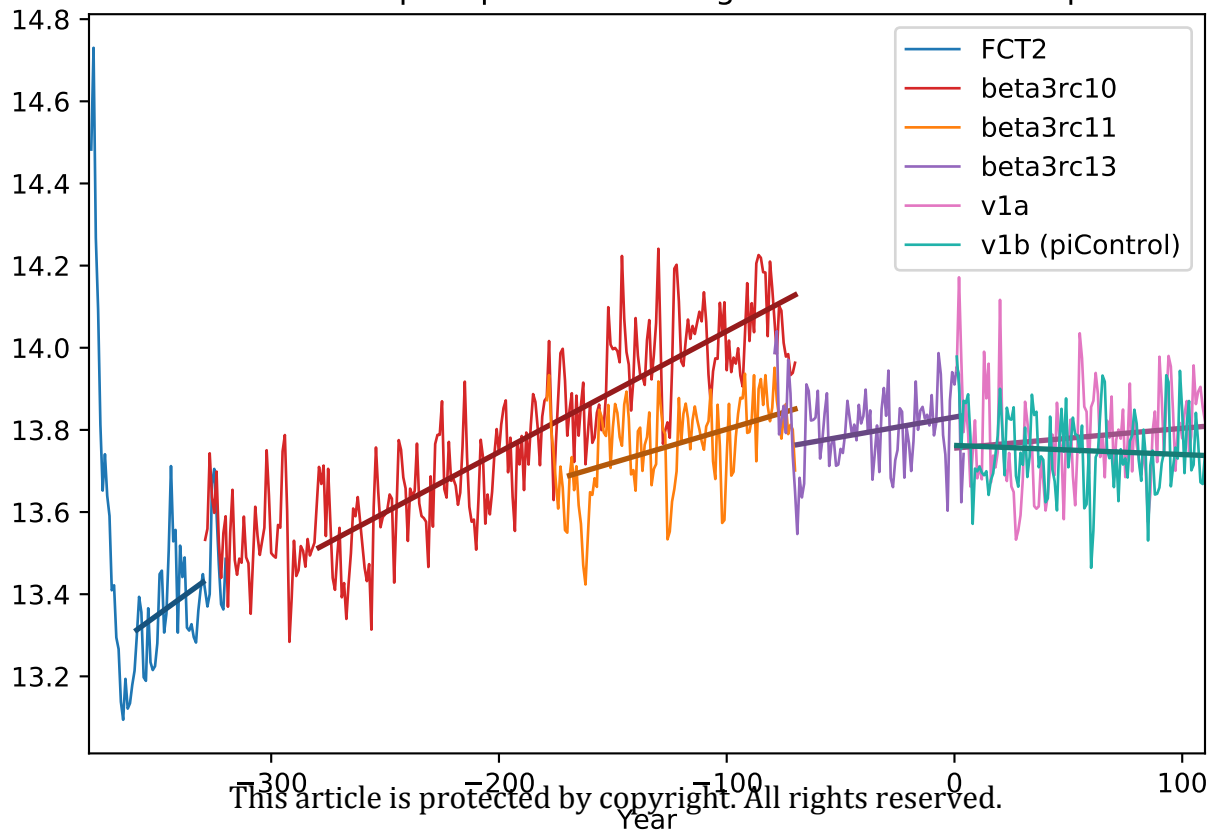


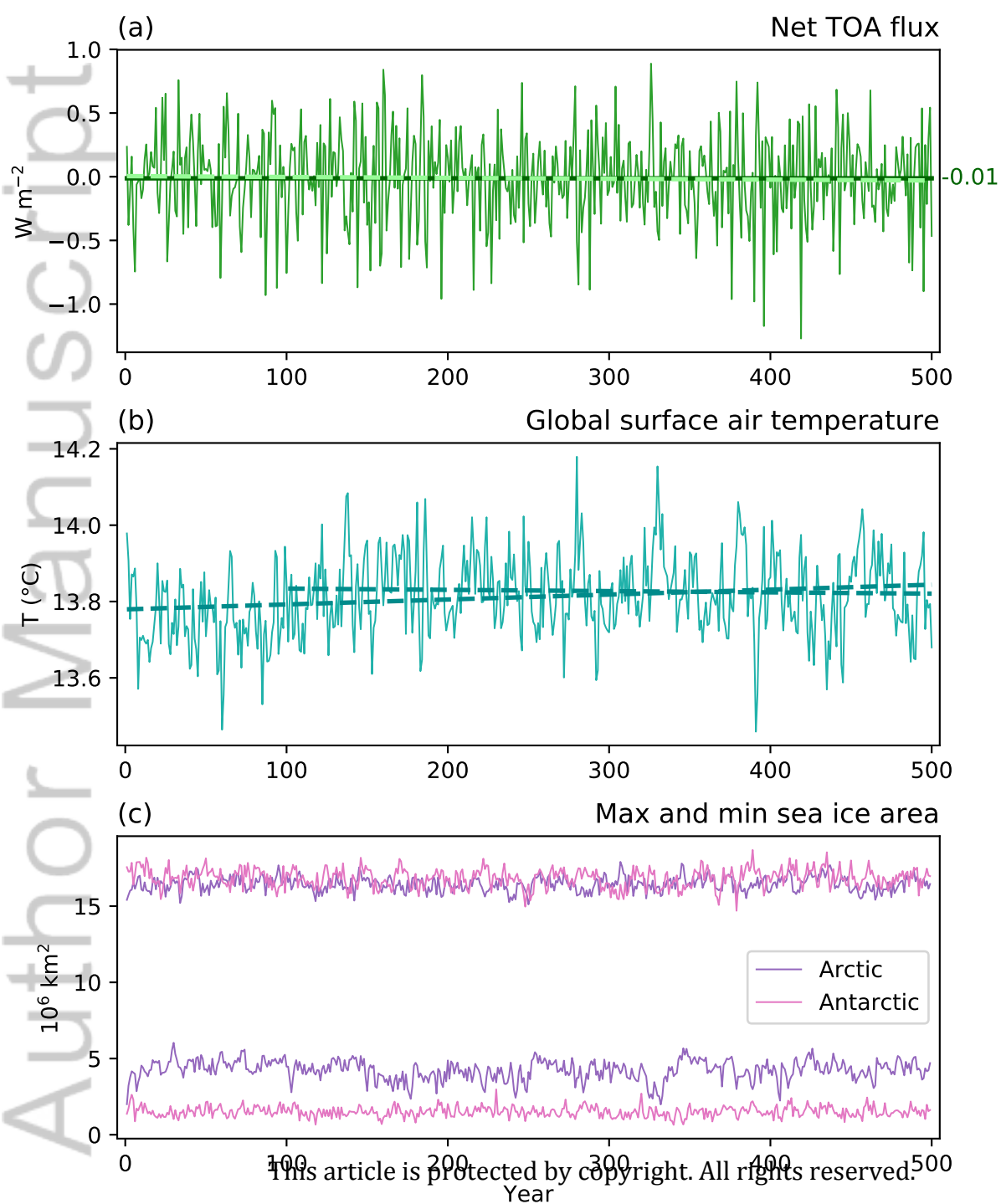






Spin-up simulations: global surface air temperature



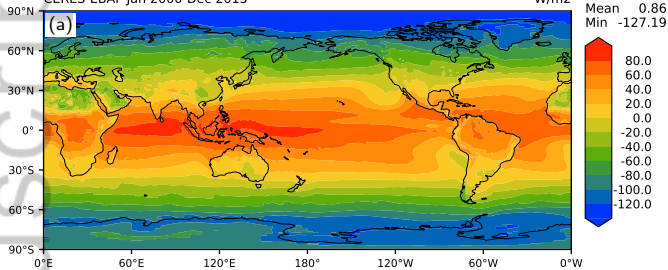


Annual net TOA radiation

CERES-EBAF Jan 2000-Dec 2015

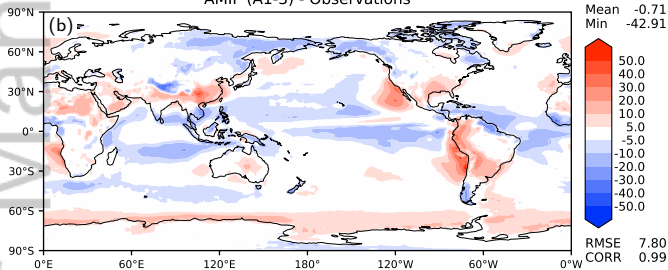
W/m²

Max 100.08
Mean 0.86
Min -127.19



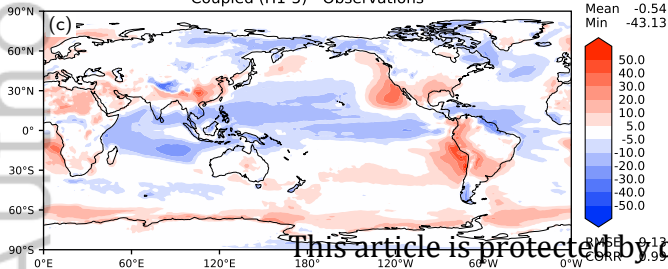
AMIP (A1-3) - Observations

Max 68.32
Mean -0.71
Min -42.91



Coupled (H1-5) - Observations

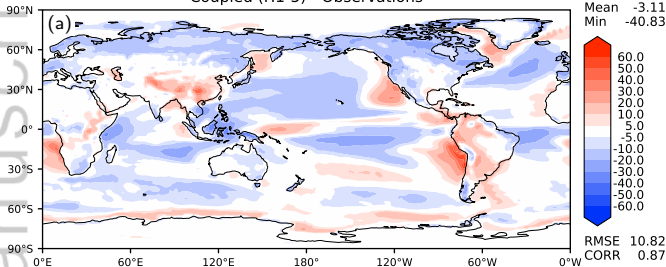
Max 65.55
Mean -0.54
Min -43.13



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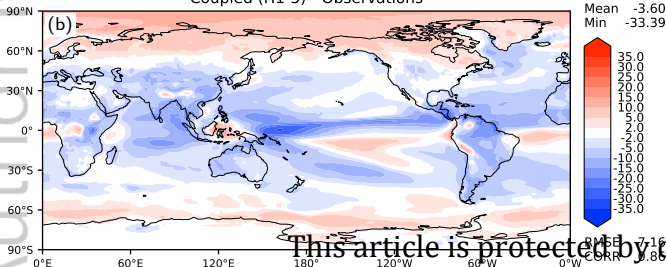
Annual SW cloud radiative effect

Coupled (H1-5) - Observations



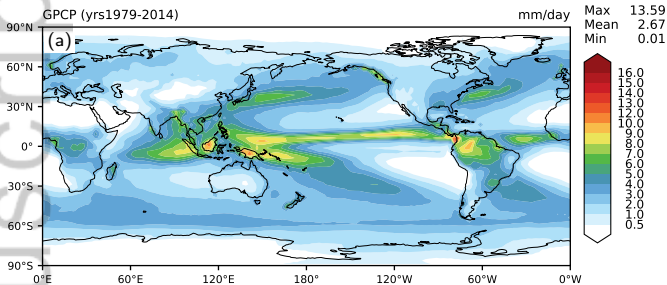
Annual LW cloud radiative effect

Coupled (H1-5) - Observations

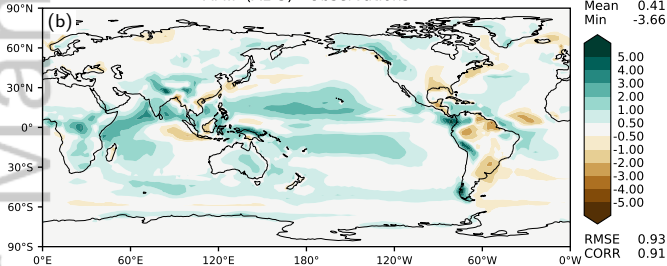


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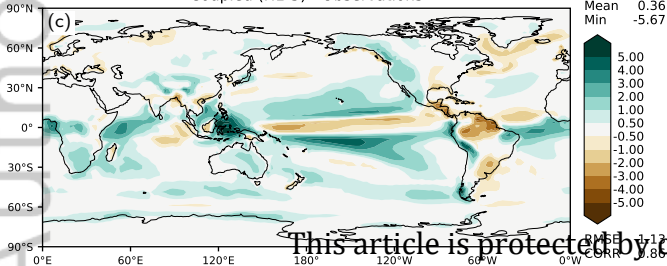
Annual precipitation



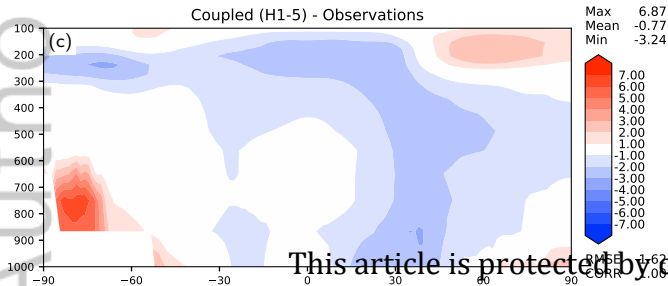
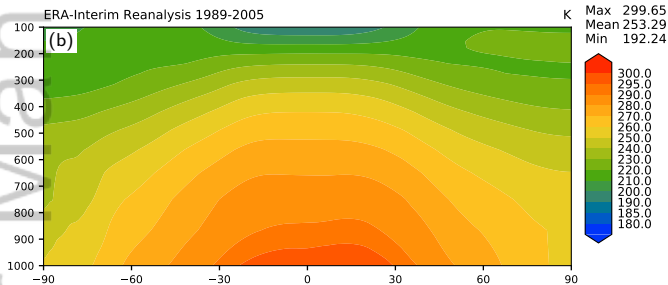
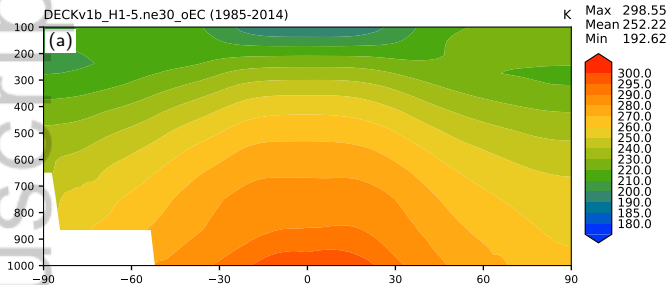
AMIP (A1-3) - Observations



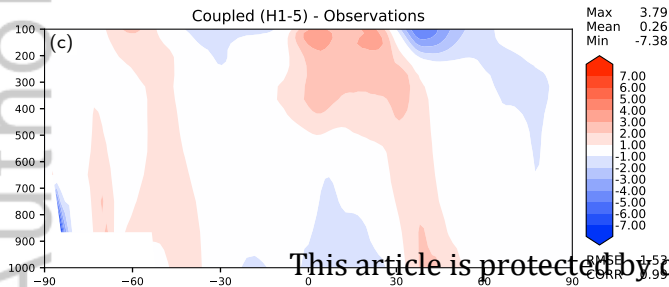
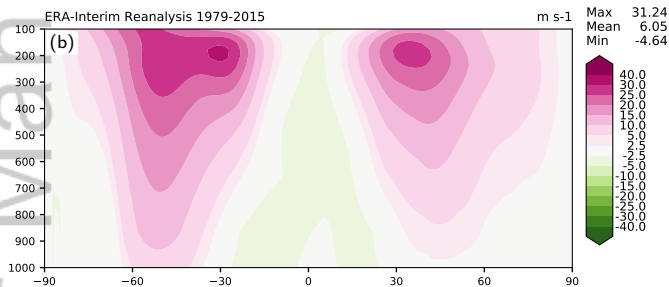
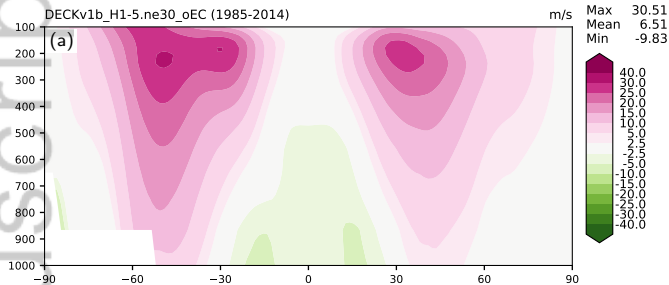
Coupled (H1-5) - Observations



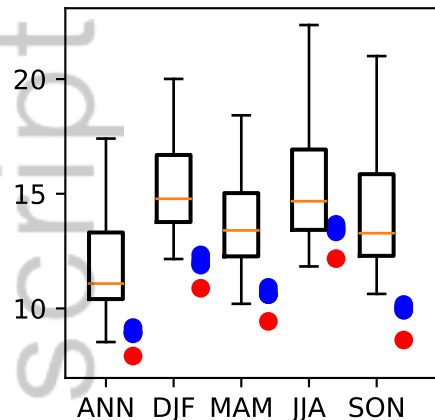
Annual temperature



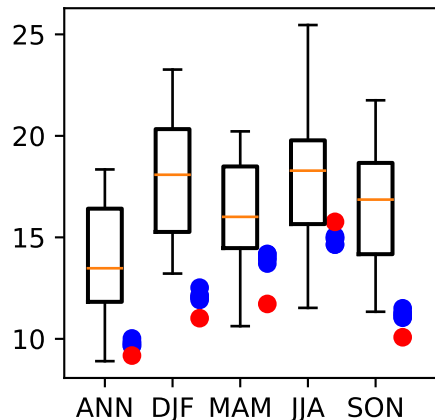
Annual zonal wind



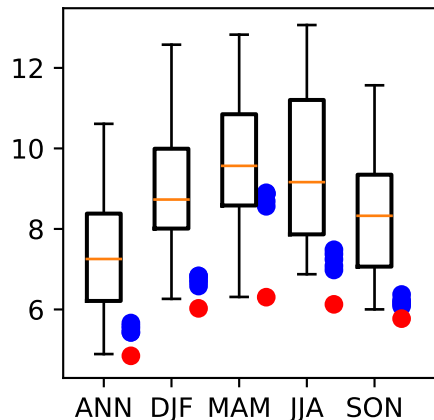
(a) Net TOA (W m^{-2})



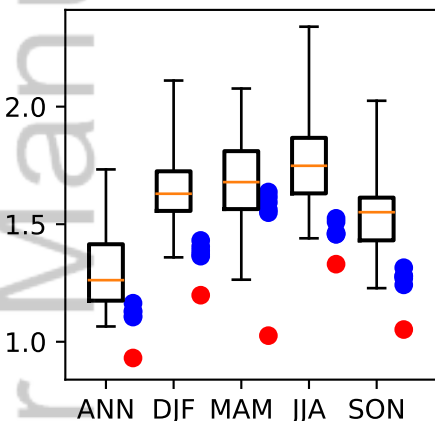
(b) SW CRE (W m^{-2})



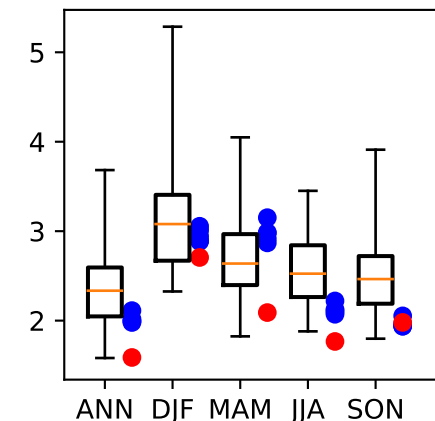
(c) LW CRE (W m^{-2})



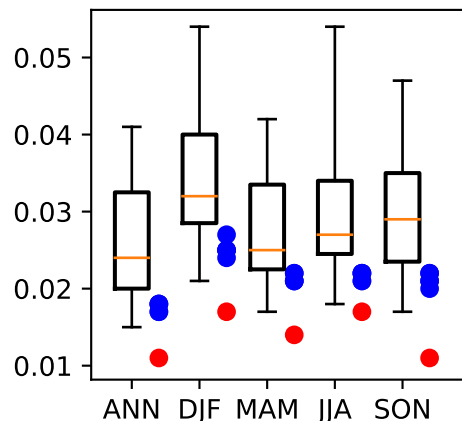
(d) prec (mm day^{-1})



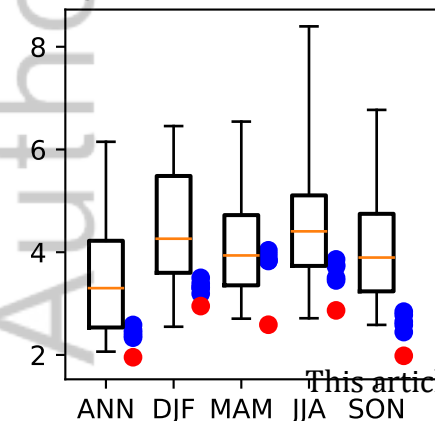
(e) tas (land, K)



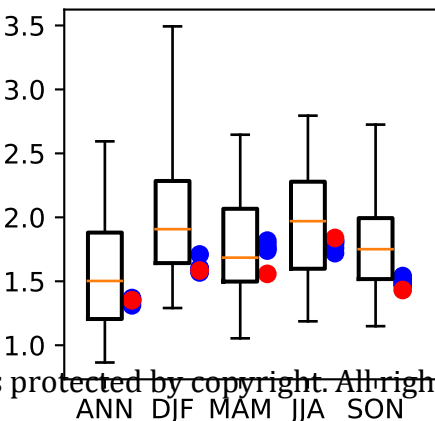
(f) τ_x (ocean, Pa)



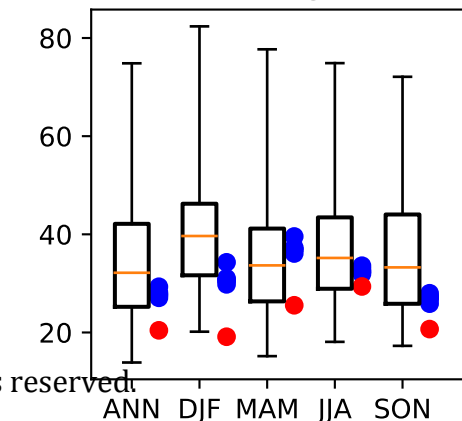
(g) u200 (m s^{-1})



(h) u850 (m s^{-1})

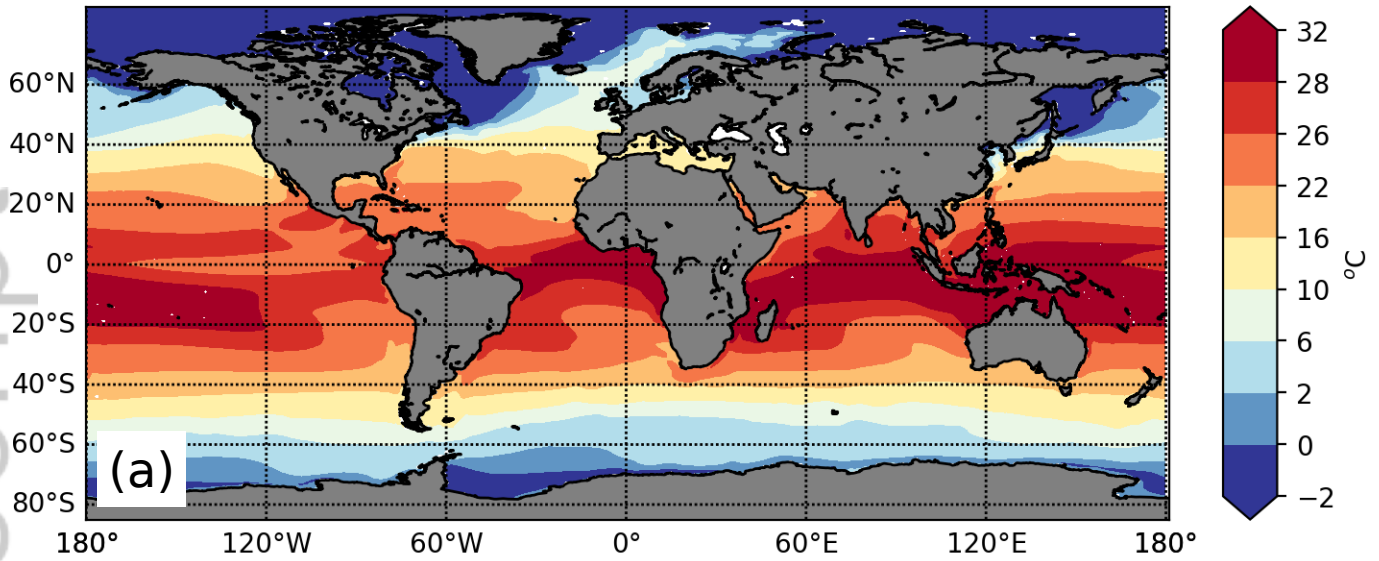


(i) Zg-500 (m)

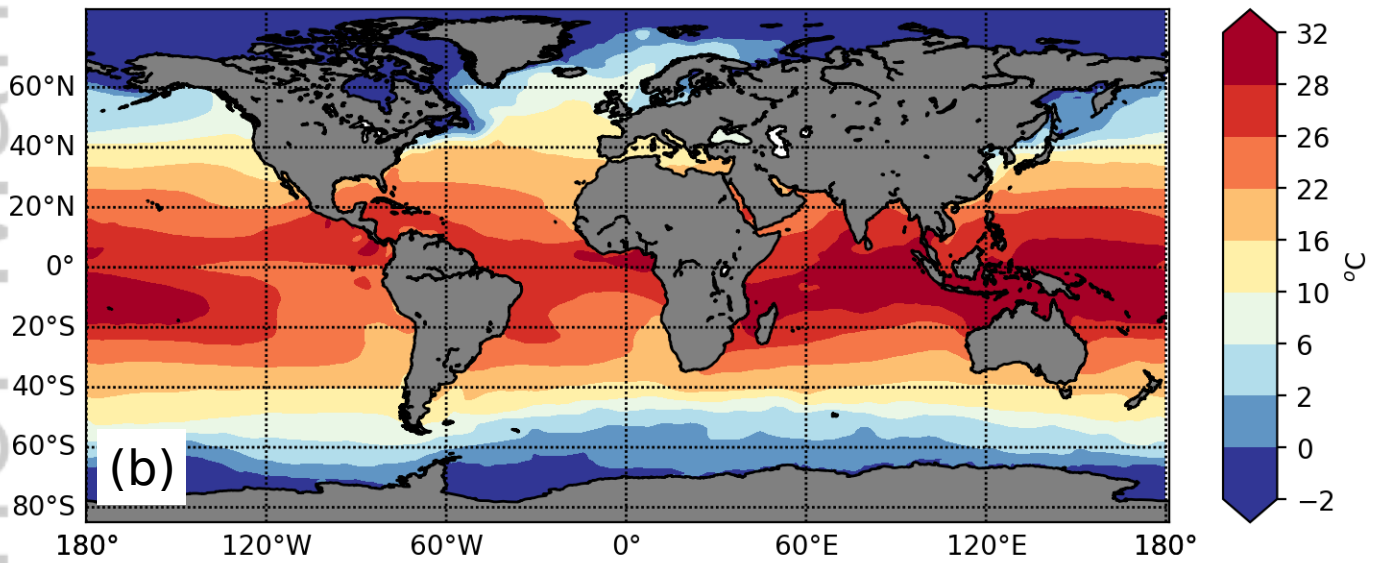


Sea Surface Temperature (Annual Average)

E3SM Historical Ensemble Average



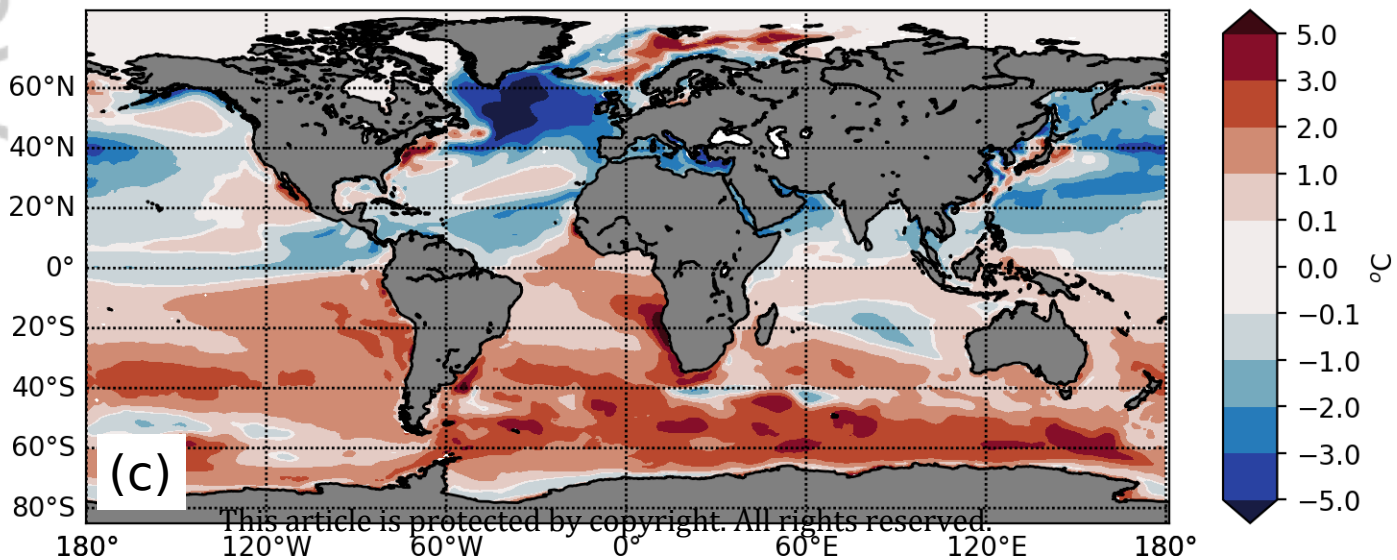
Observations



Model - Observations

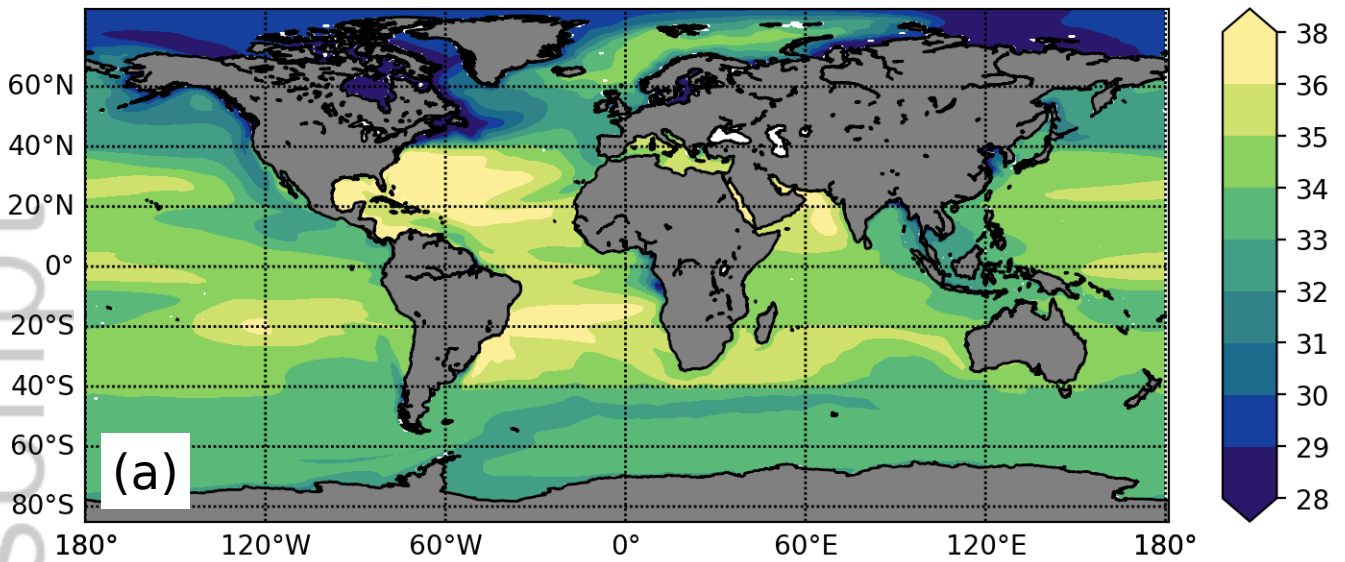
[0.067, **0.094**, 0.115]

<0.914, **0.939**, 0.985>

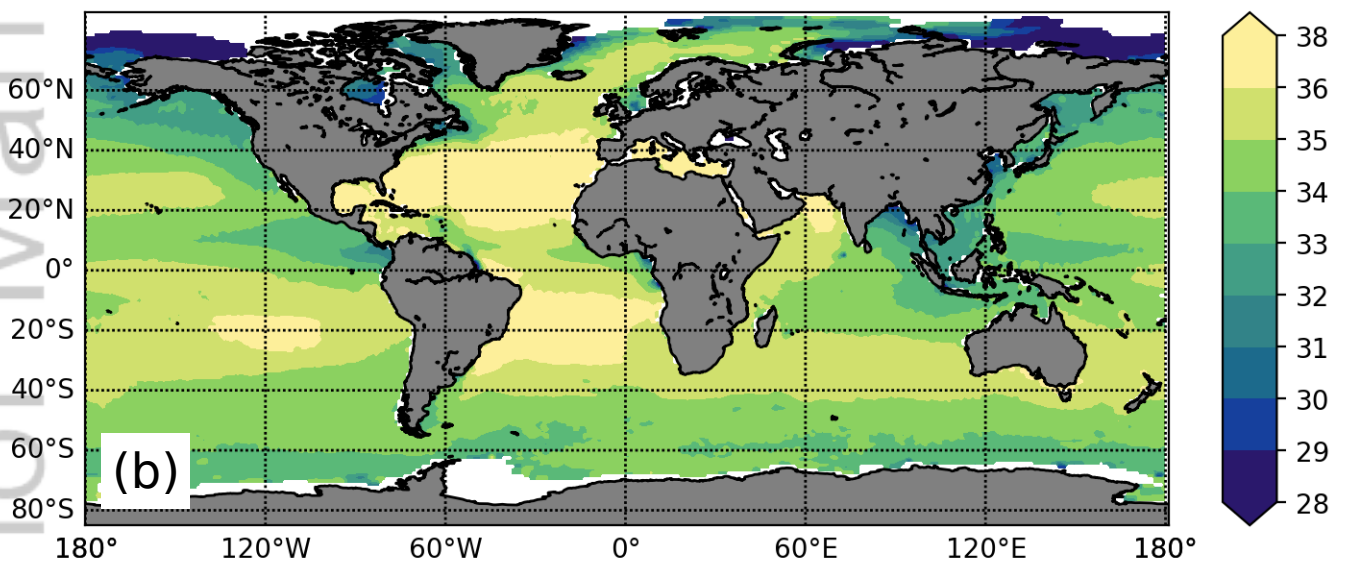


Sea Surface Salinity (Annual Average)

E3SM Historical Ensemble Average

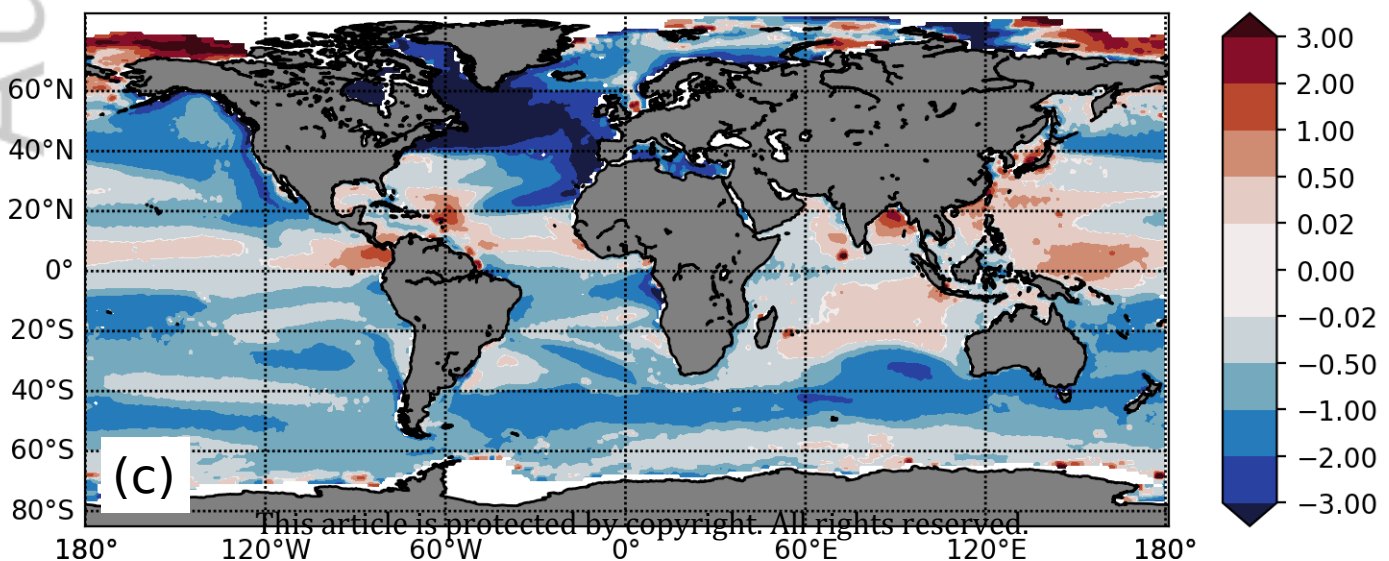


Observations



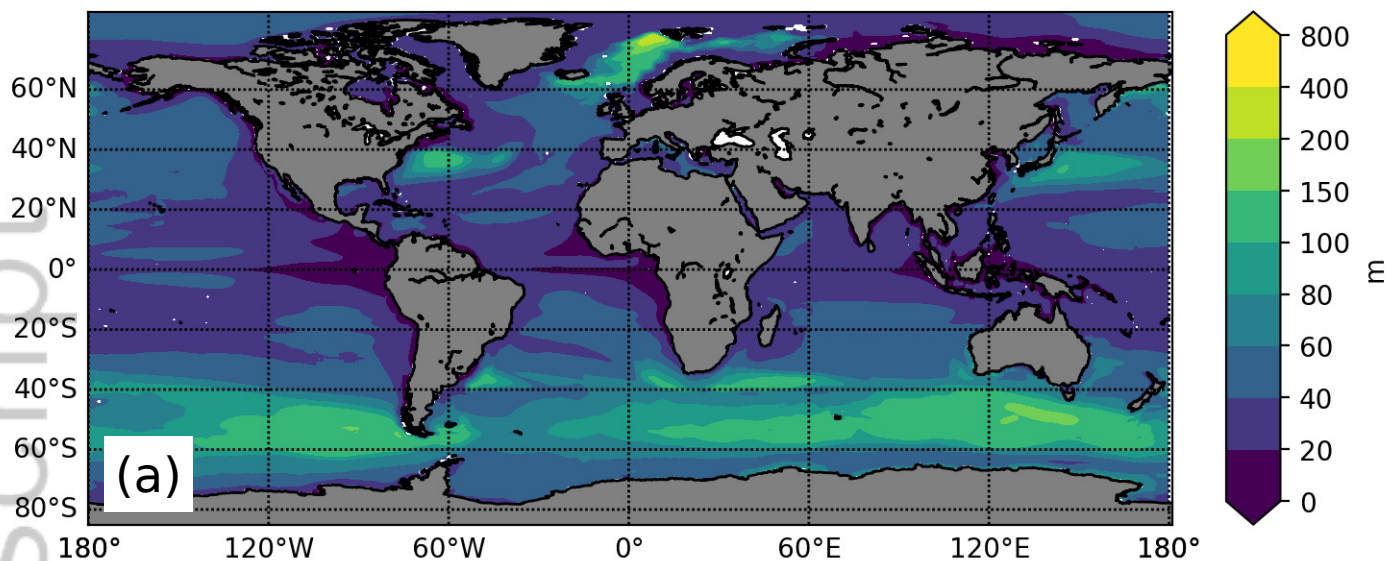
Model - Observations

[-0.72, **-0.71**, -0.69]
<1.10, **1.12**, 1.14>

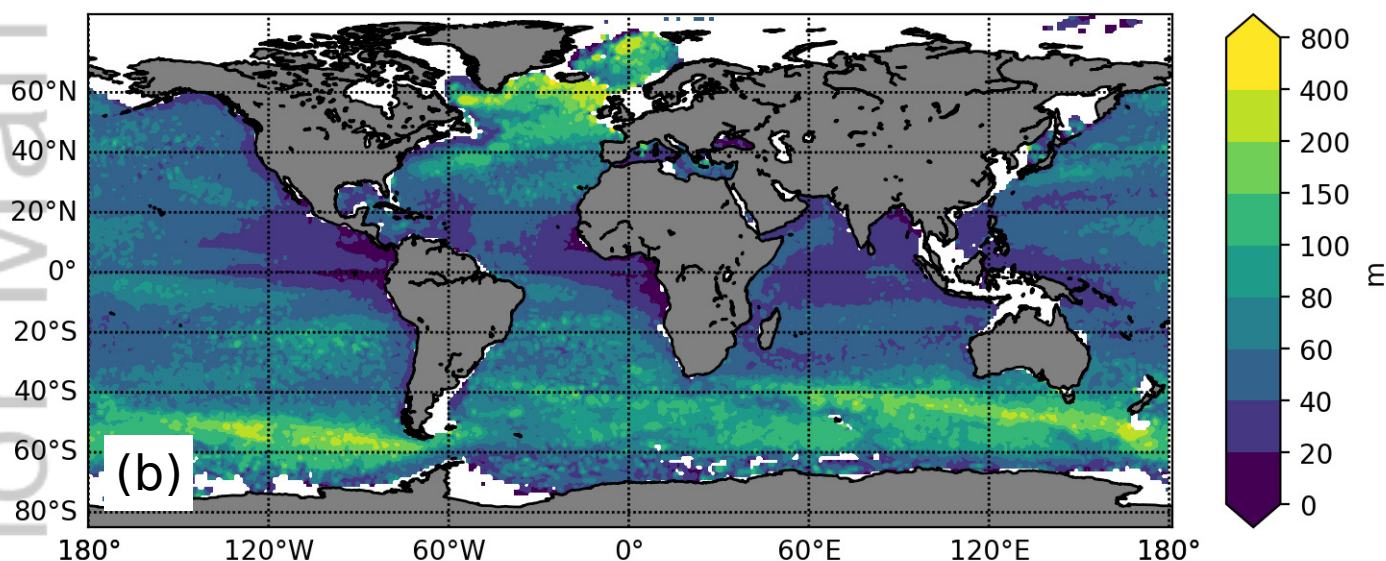


Mixed Layer Depth (Annual Average)

E3SM Historical Ensemble Average



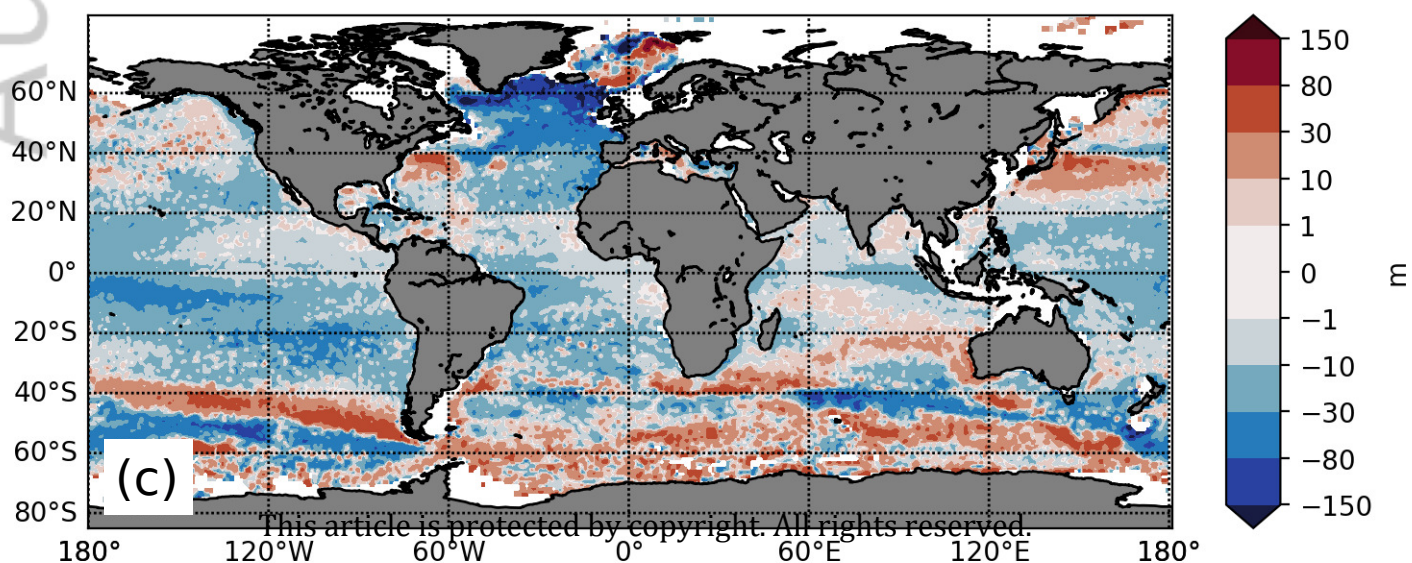
Observations

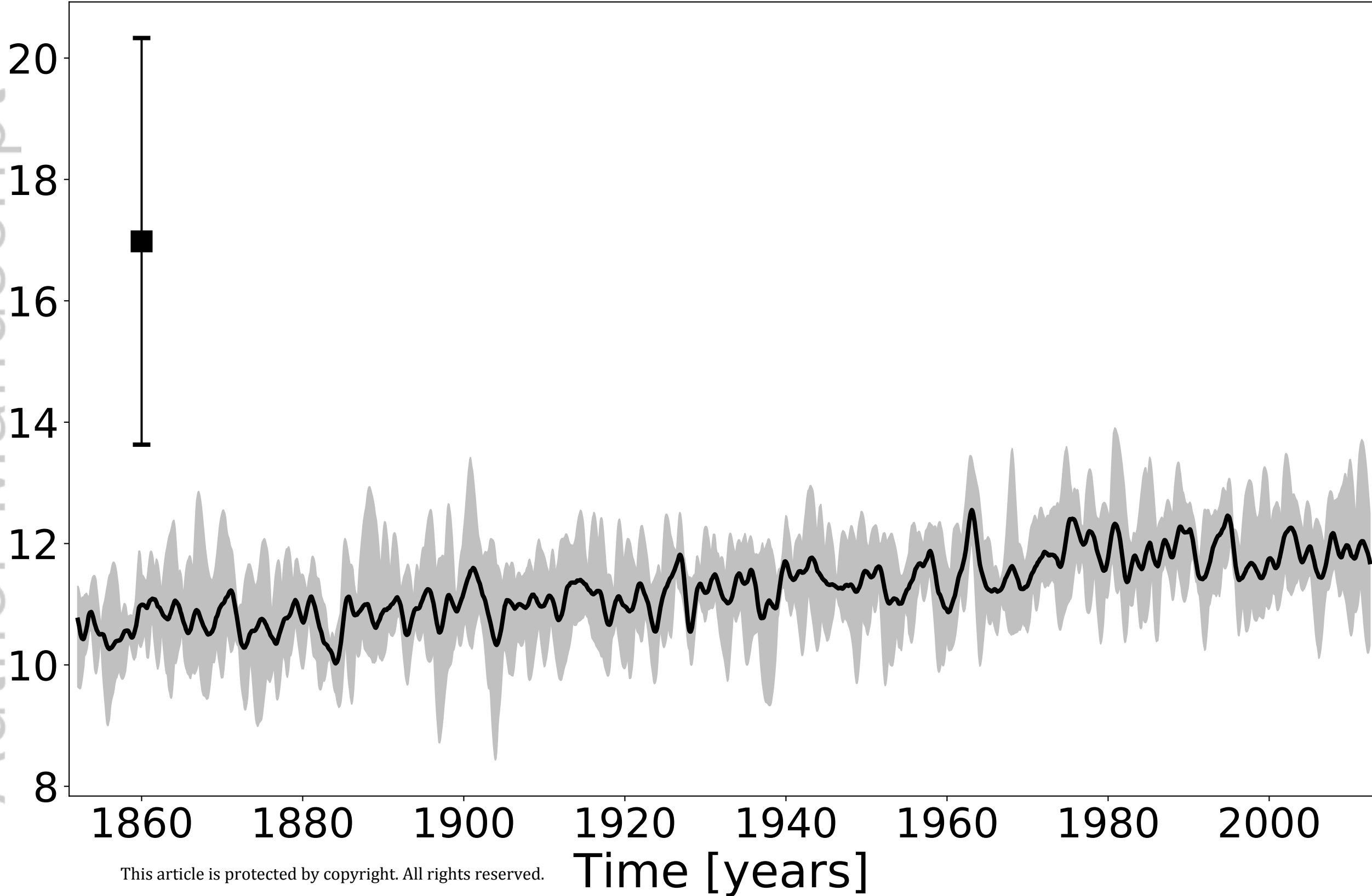


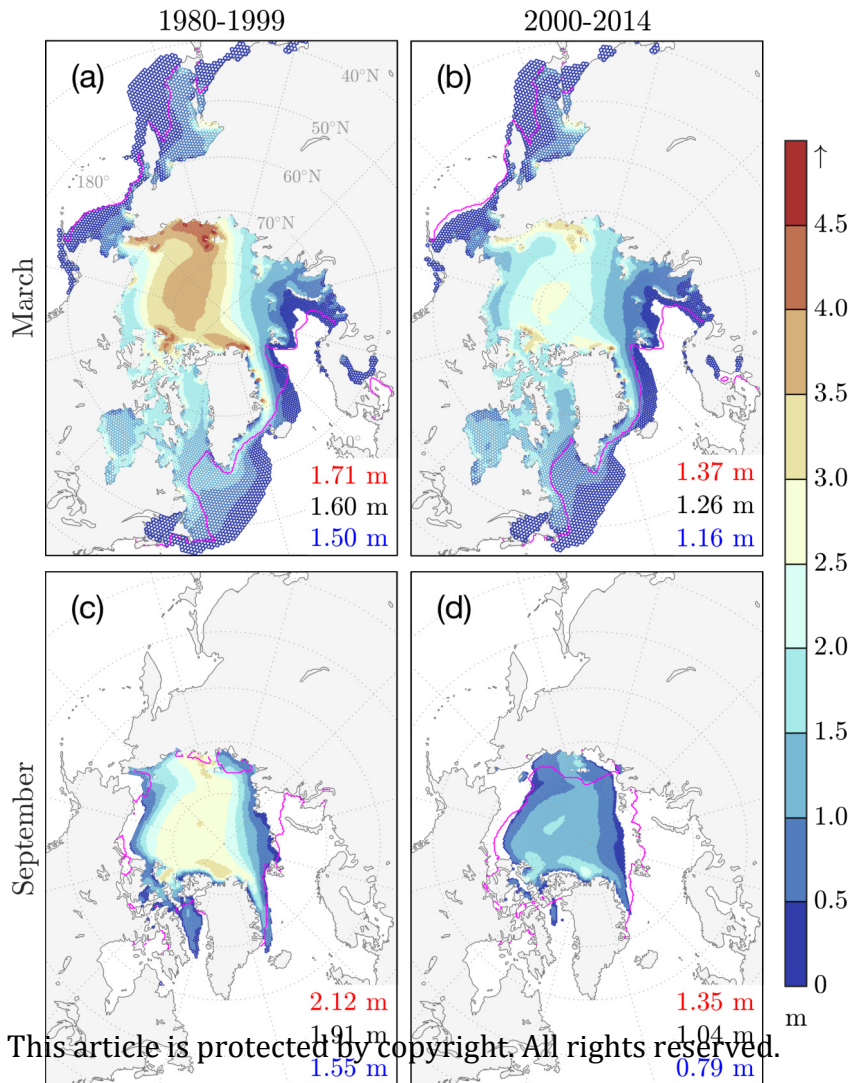
Model - Observations

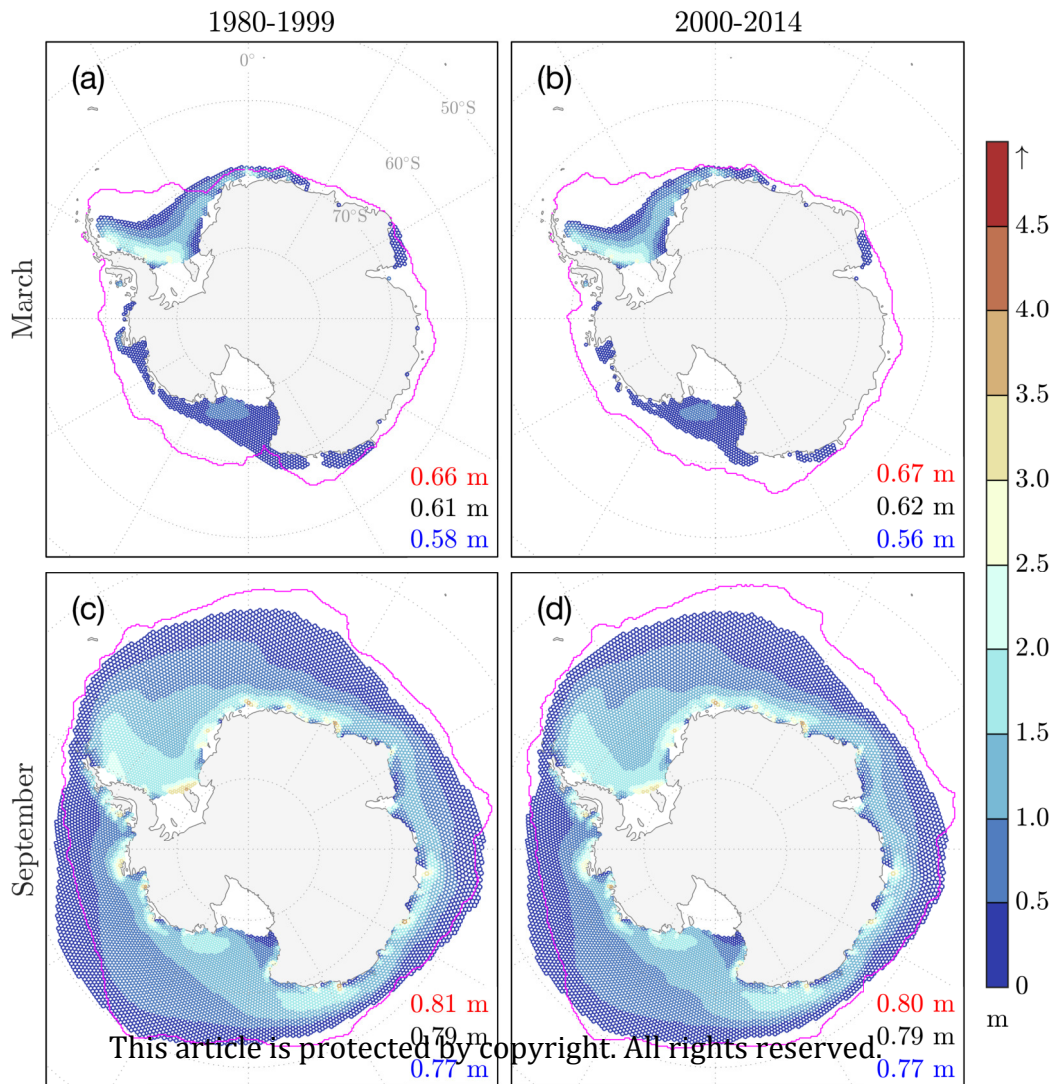
[-19.28, -18.54, -18.22]

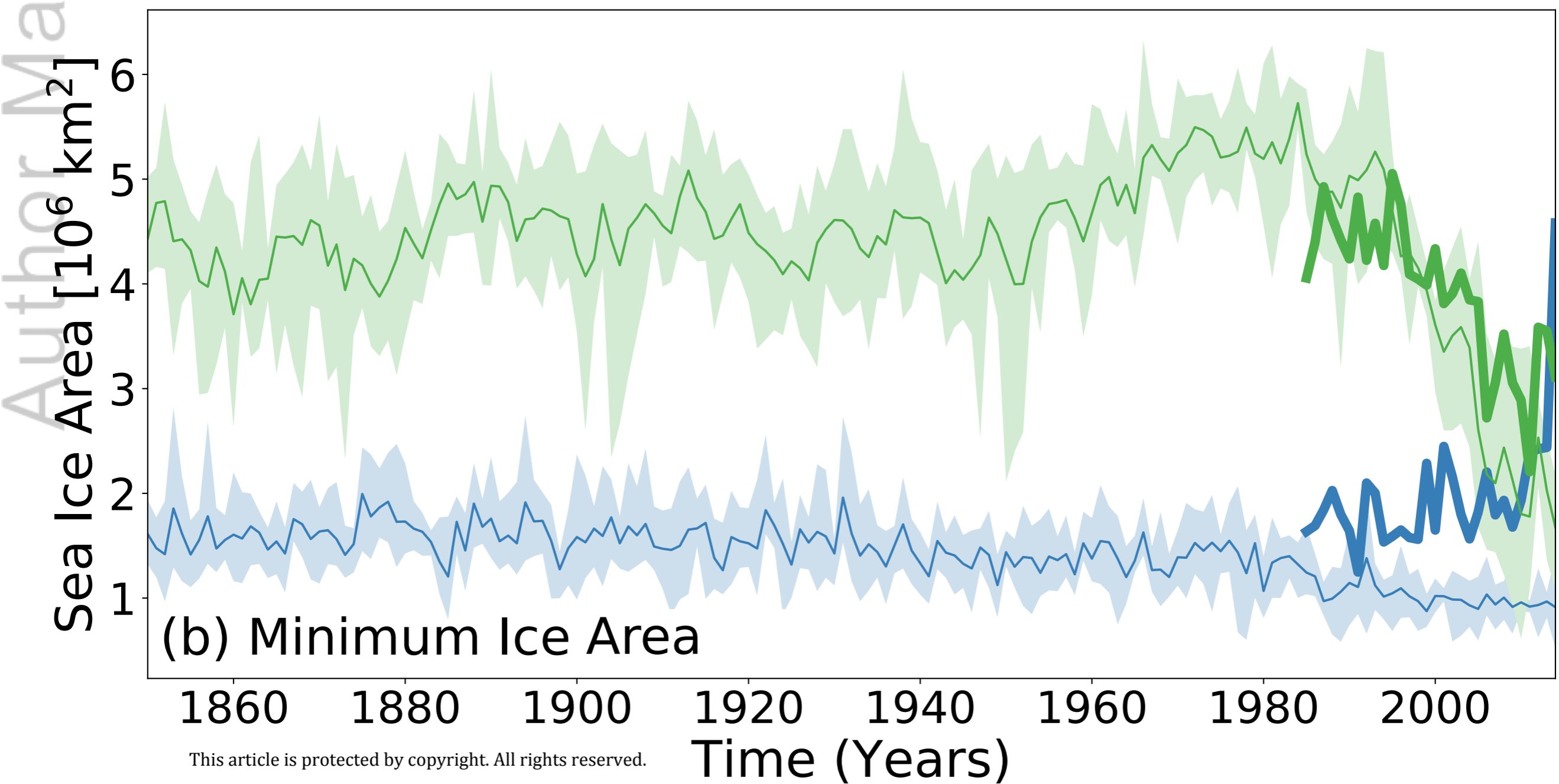
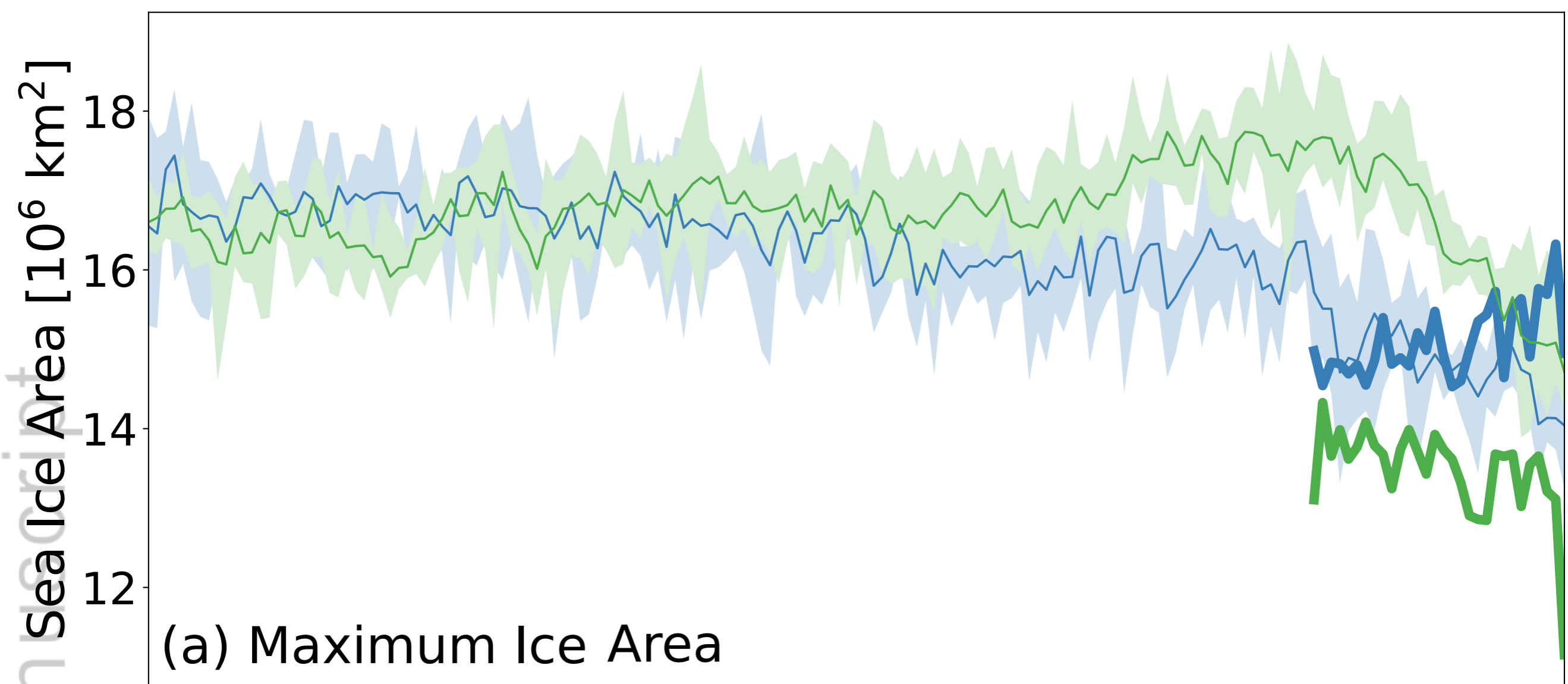
<41.36, **41.74**, 42.47>



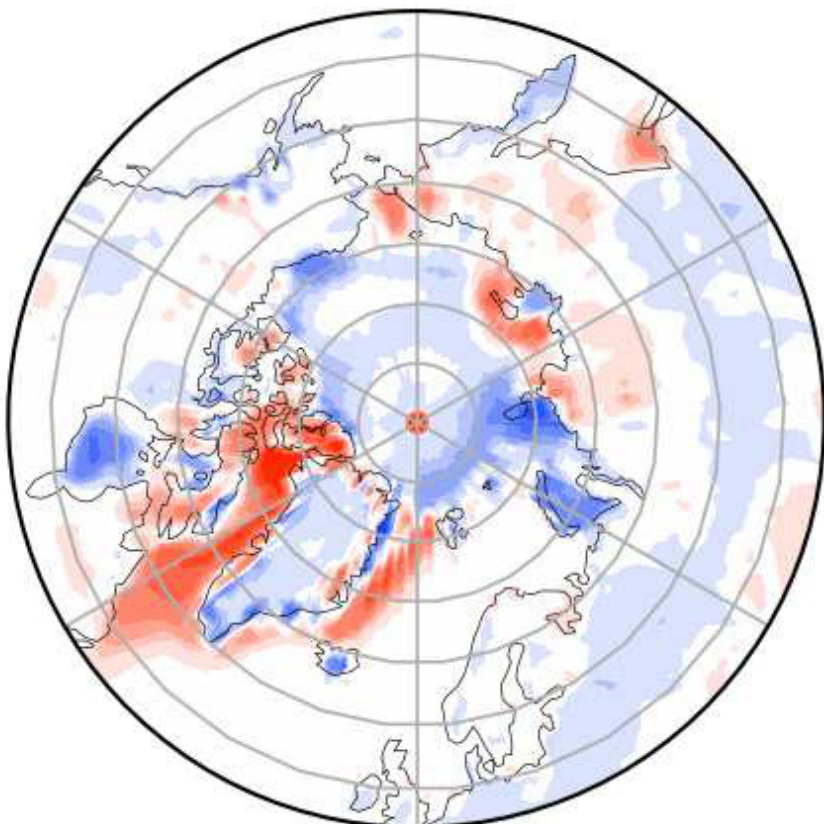




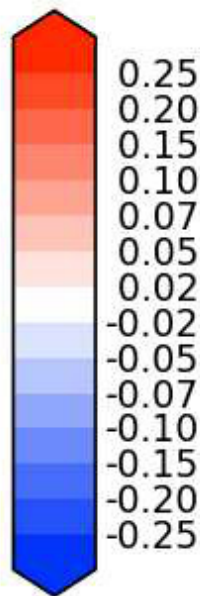




Model - Obs



Max	0.49
Mean	-0.00
Min	-0.32

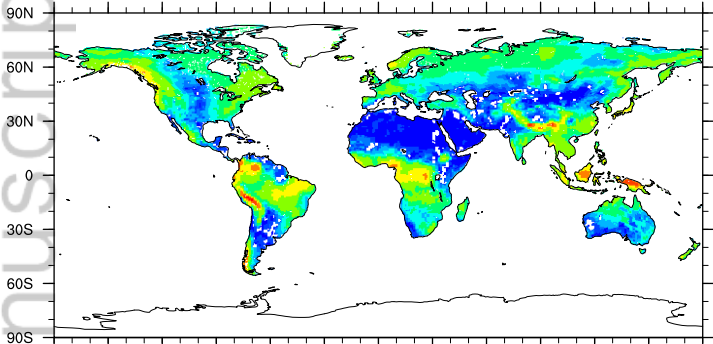


RMSE	0.05
CORR	0.96

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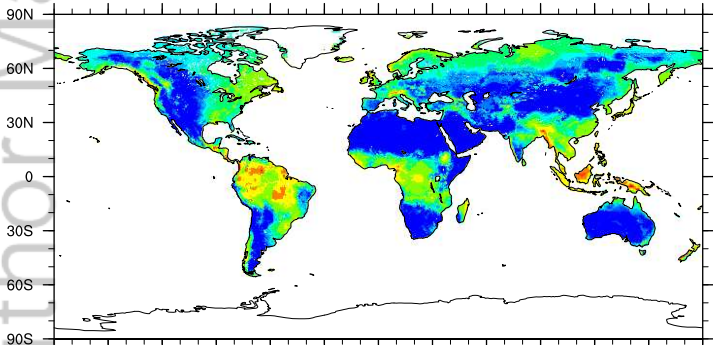
ELM total runoff

mm/day



GRDC total runoff

mm/day

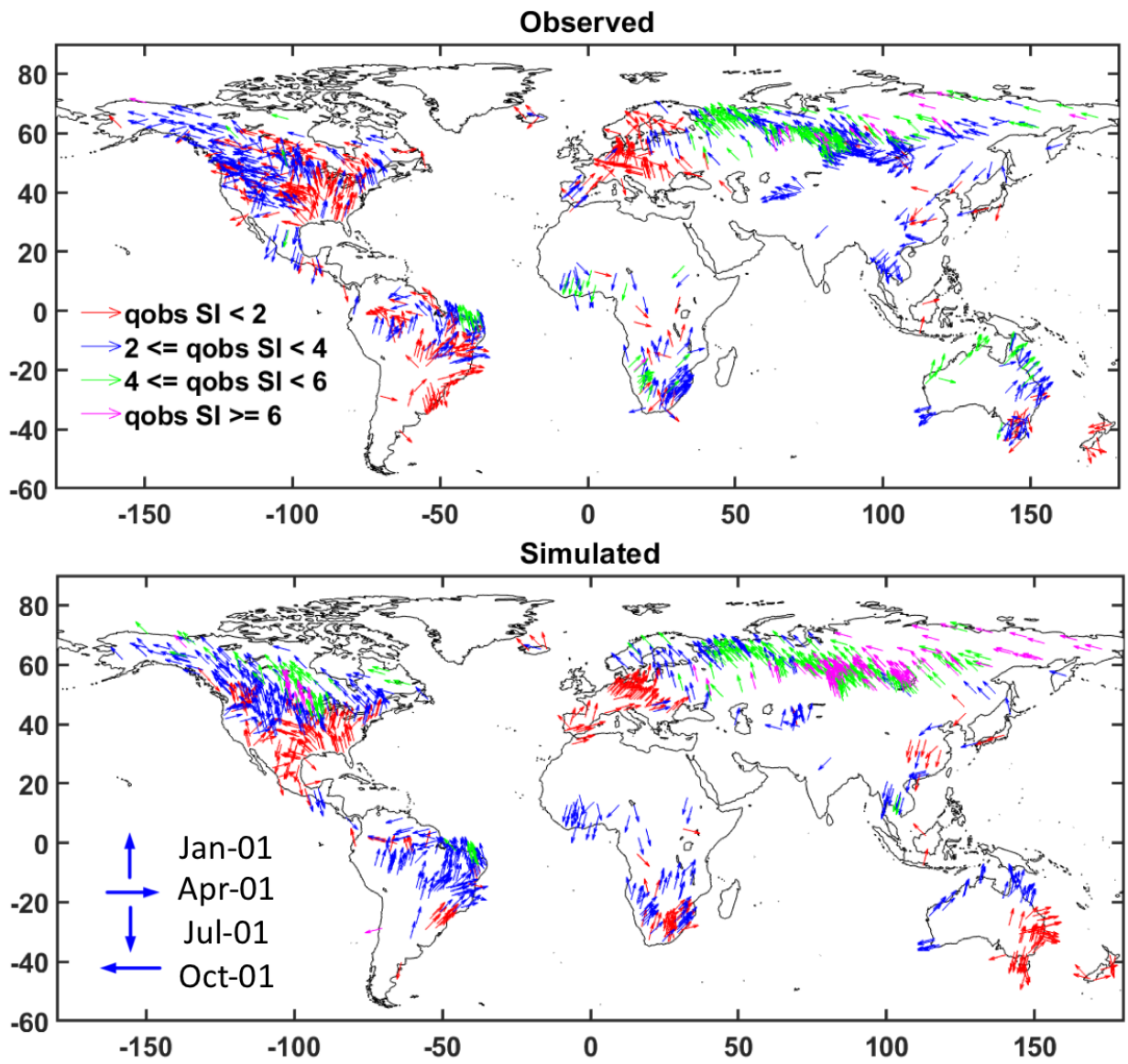


180 150W 120W 90W 60W 30W 0 30E 60E 90E 120E 150E 180



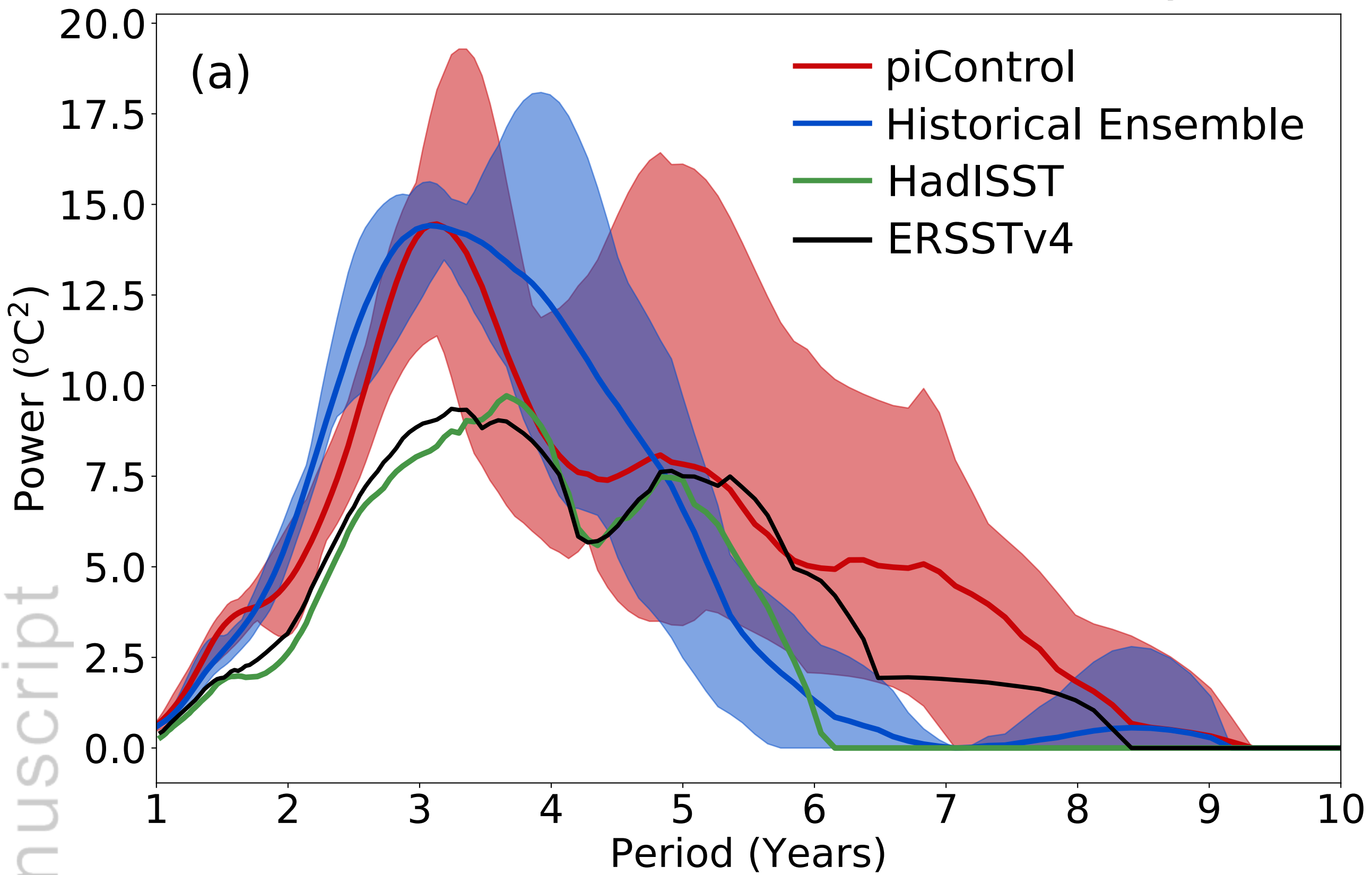
-1 0 0.01 0.05 0.1 0.25 0.5 1 2.5 5 7.5 10

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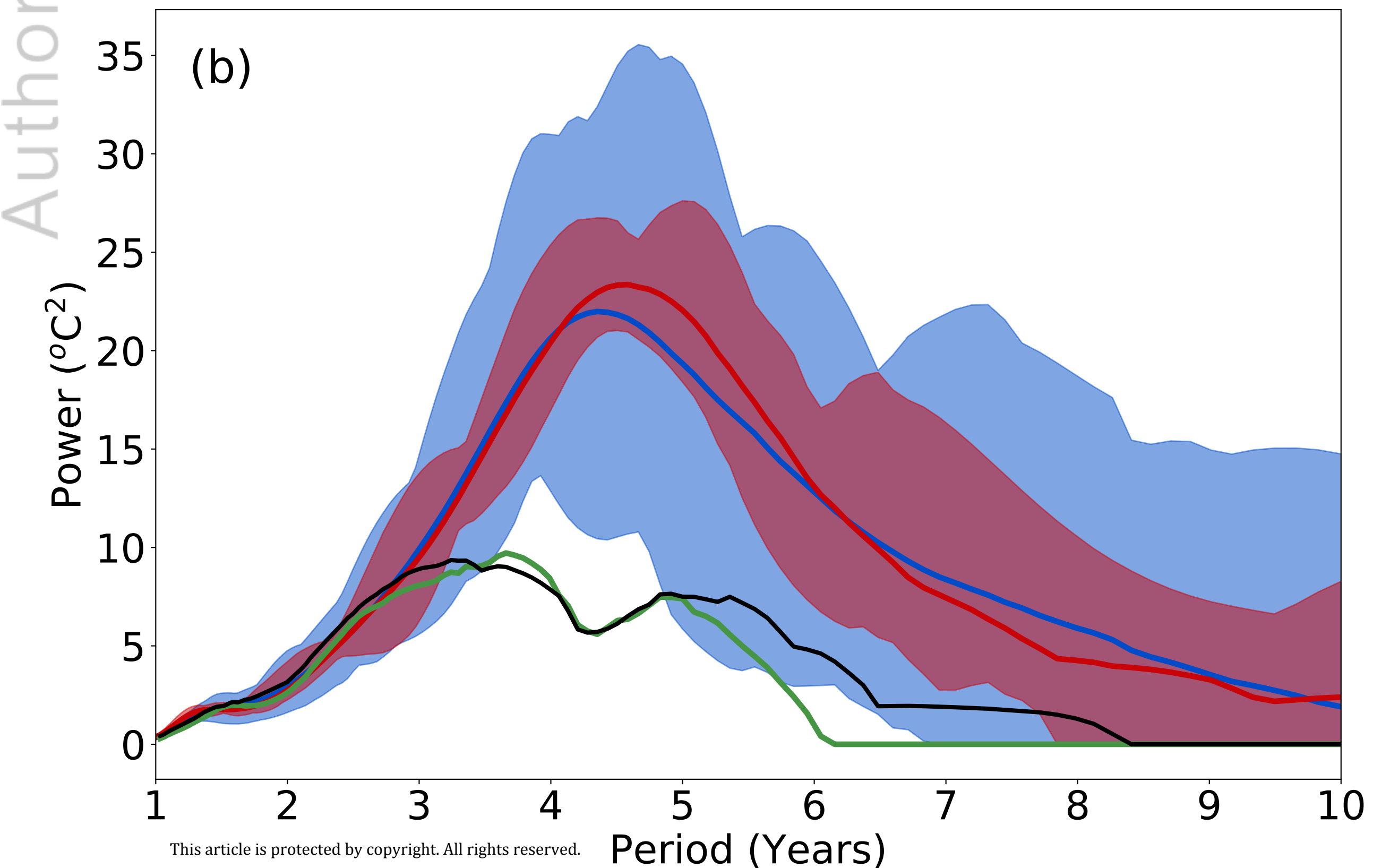


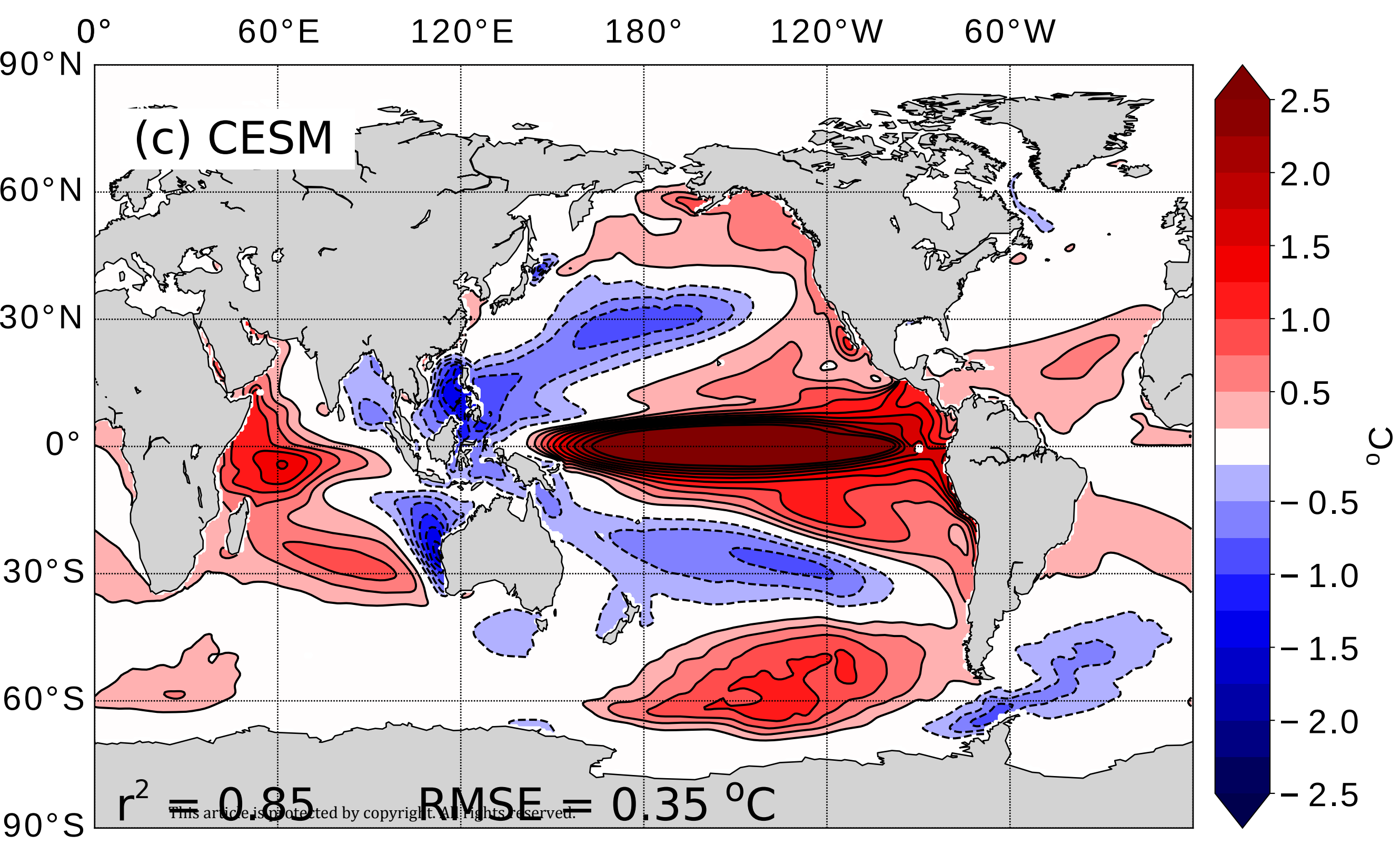
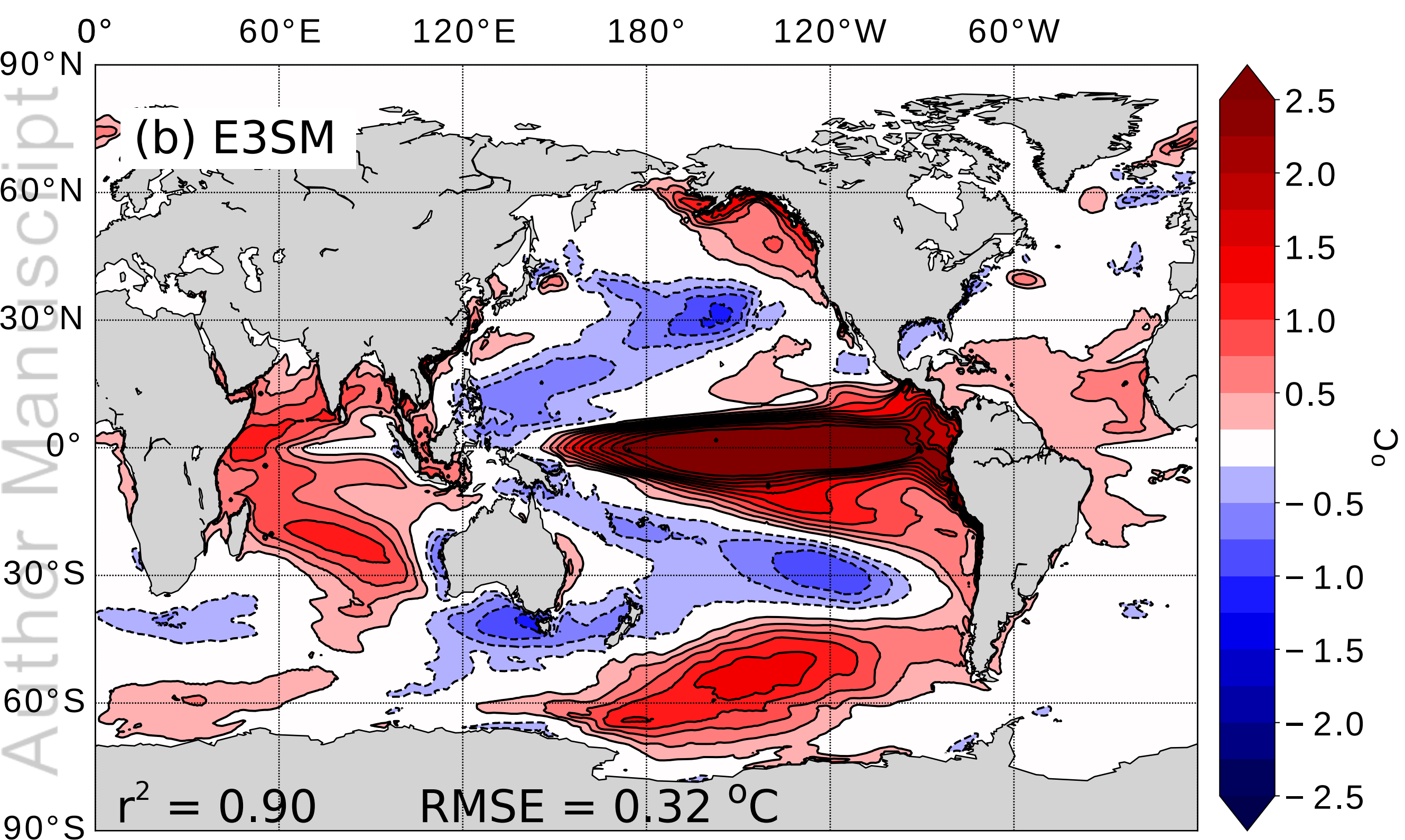
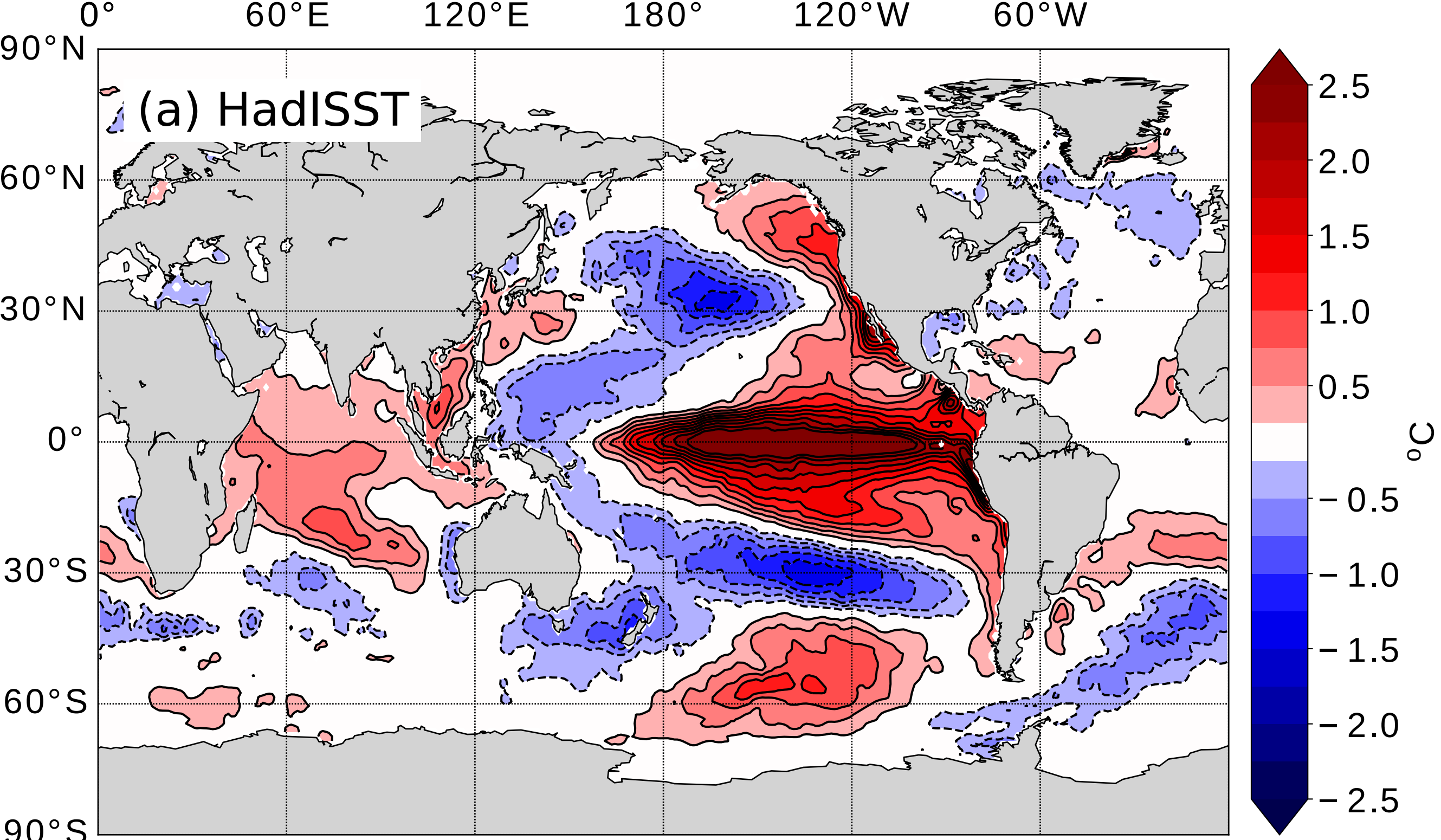
2018MS001603-f19-z-.png

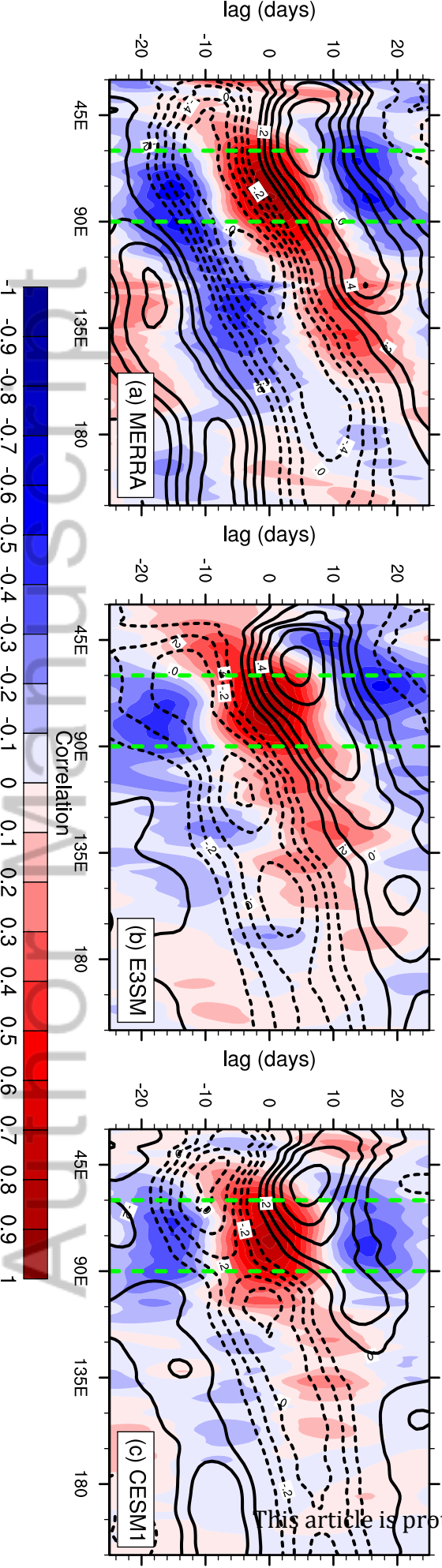
E3SM Nino 3.4 SST Wavelet Analysis



CESM-LE Nino 3.4 SST Wavelet Analysis

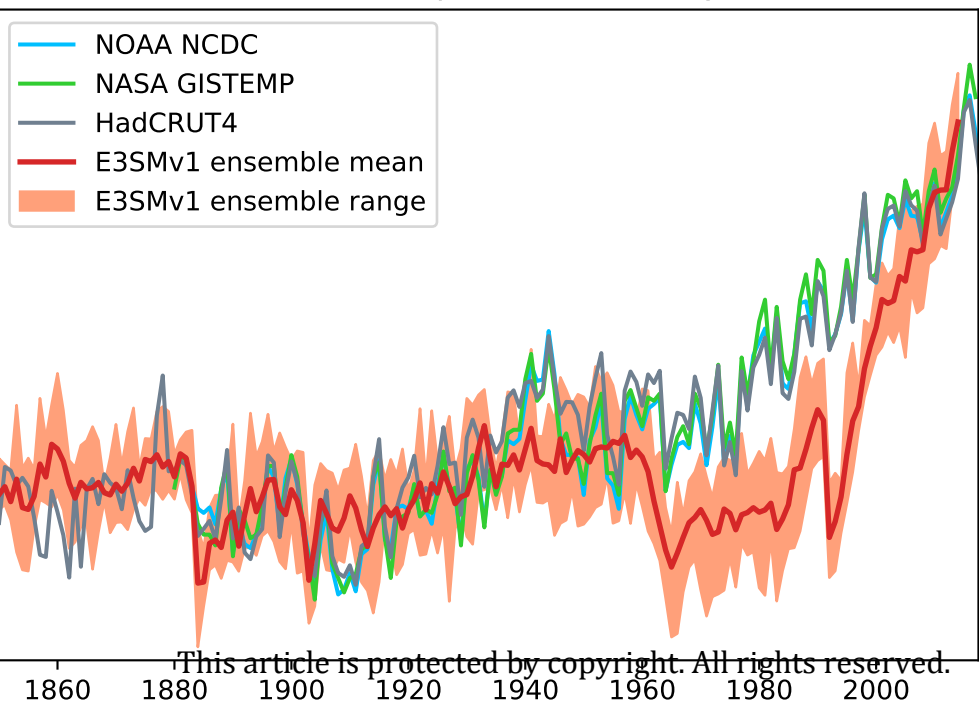






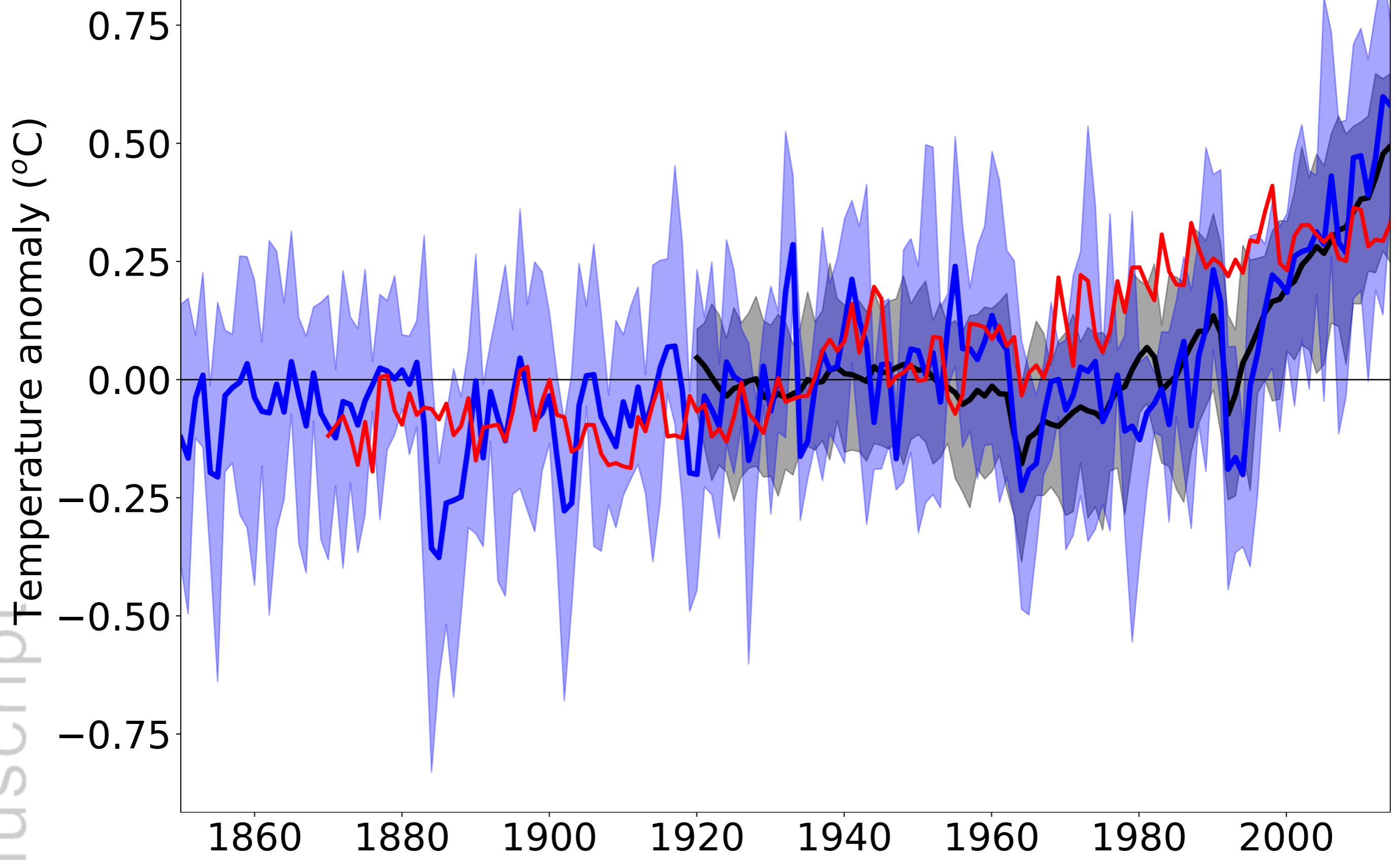
Global surface air temperature anomaly (vs 1880-1909)

Author Manuscript

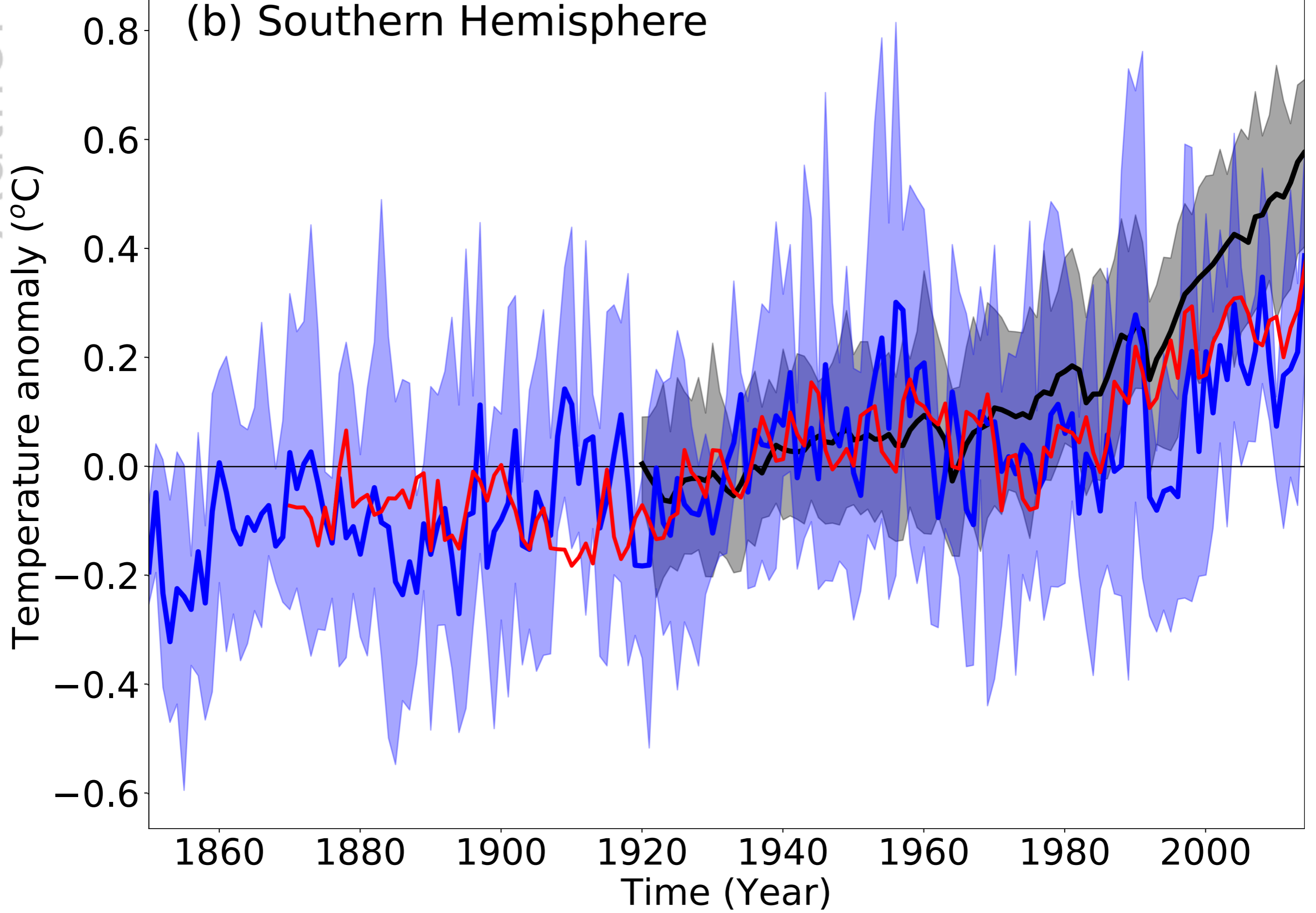


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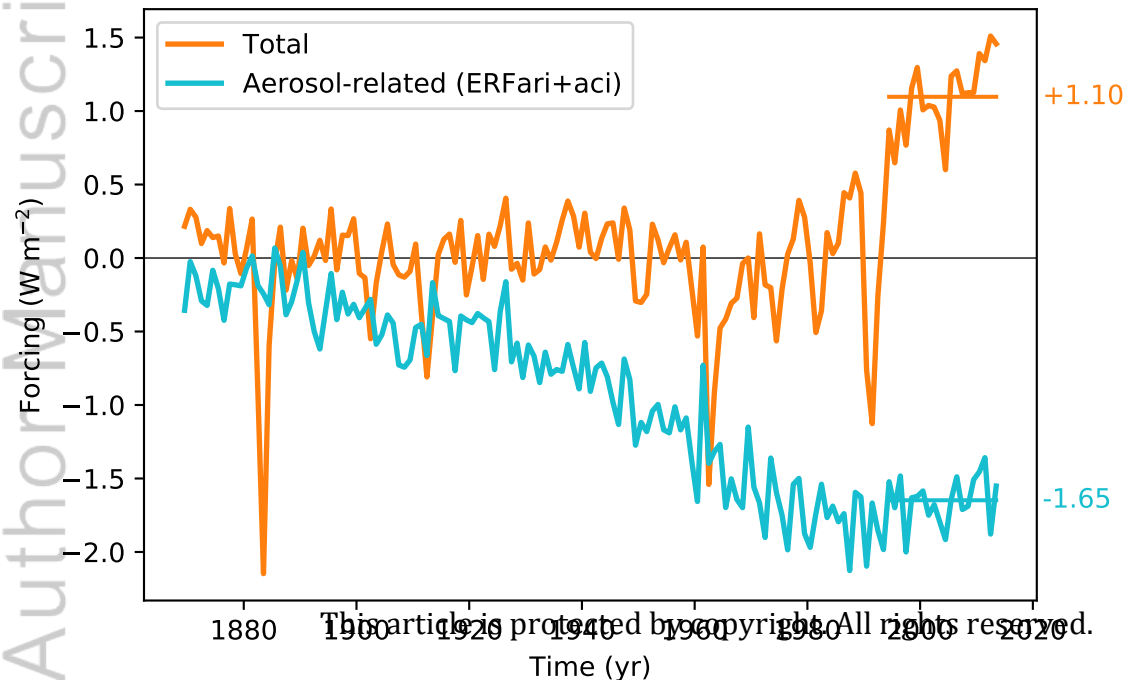
(a) Northern Hemisphere

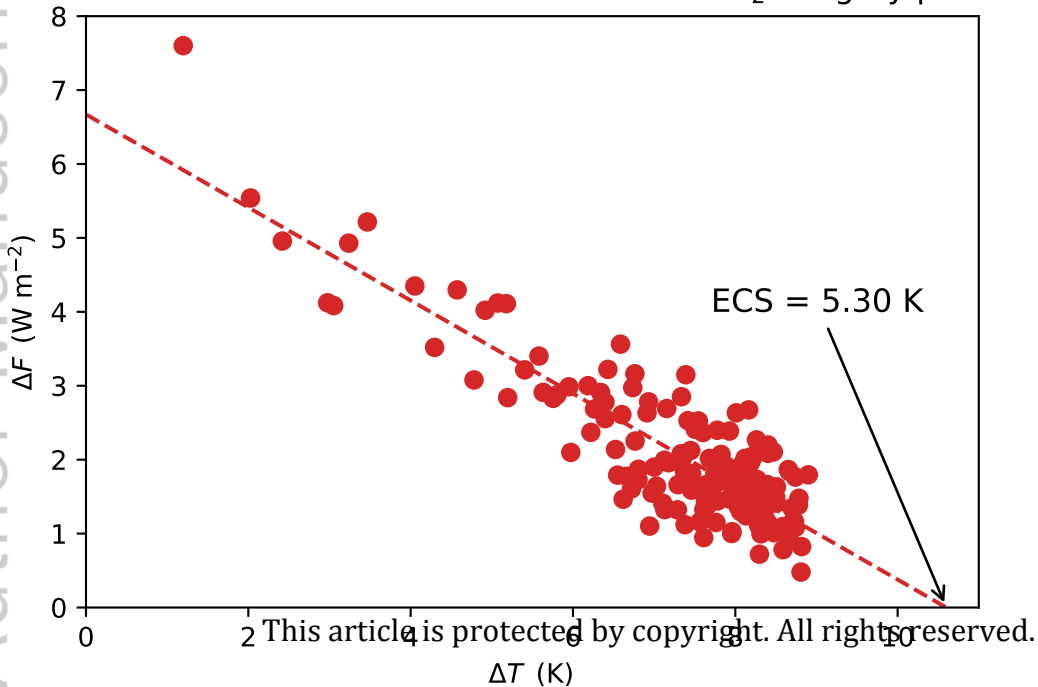


(b) Southern Hemisphere



Effective Radiative Forcing (relative to 1850)



4xCO₂ Gregory plot

Idealized CO₂ Forcing Simulations

