

# ULF Modulations on Plasma Environment and Coherent Waves of Mercury's Magnetosphere: MESSENGER's Observation

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## Key Points:

- ~15 mHz magnetic field pulsations were detected and found to be capable of driving proton flux fluctuations in Mercury's magnetosphere.
- The associated ~1 Hz magnetic waves have ~15 mHz variations in power and compressibility, indicating modulation by the ~15 mHz compressional waves.
- The influence on ~1 Hz waves is possibly attained by plasma flux changes, helping us better interpret its wave mode and character.

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**Abstract**

Ultra low-frequency (ULF) waves are fundamental waves that can energize, transport, and scatter charged particles in planetary magnetospheres. With the measurements from MErcury Surface, Space ENvironment, GEOchemistry, and Ranging (MESSENGER), we investigate the proton flux fluctuations and coherent waves associated with a series of ULF waves on the flanks of Mercury's magnetosphere. The ULF waves are mainly compressional with a frequency of  $\sim 15$  mHz and significantly modulate the intensity of proton flux. The coherent waves accompanied by the ULF waves correspond to a higher frequency ( $\sim 1$  Hz). The wave power and compressibility of the coherent waves vary quasi-periodically with the  $\sim 15$  mHz ULF waves. We conclude that the compressional ULF waves modulate the coherent waves with higher frequency. This modulation might result from the associated periodic proton flux changes and helps us understand the nature of the  $\sim 1$  Hz waves better.

**Plain Language Summary**

This study presents MESSENGER observations of wave-particle interactions on the flank of Mercury's magnetosphere. Clear one-to-one modulation of proton flux by  $\sim 15$  mHz compressional waves is shown in these two cases. Proton gyro-frequency ( $\sim 1$  Hz) waves are observed during the same interval. Their amplitudes and compressibilities are periodic with a frequency of  $\sim 15$  mHz, coinciding with the frequency of the compressional waves. This coincidence indicates that the  $\sim 1$  Hz waves are possibly affected by the  $\sim 15$  mHz waves. The above observational facts demonstrate that ULF waves can modify the plasma environment and further affect the kinetic scale dynamics in the magnetosphere of Mercury. Its importance is further validated by our statistical result, which shows that wave-particle and wave-wave modulation are not rare events.

**1 Introduction**

As the innermost and the smallest planet in the solar system, Mercury has a dipole moment of  $190 \text{ nT} \cdot R_M^3$  with the dipole center shifting northward by  $0.2 R_M$  ( $R_M$  refers to Mercury's radius, which equals to 2440 km) from the geographic center (Alexeev et al., 2010; Anderson et al., 2011, 2012). Due to strong solar wind enforcement, a highly compressed magnetosphere with an average magnetopause subsolar distance of  $\sim 1.45 R_M$  forms (Slavin et al., 2009; Winslow et al., 2013; Philpott et al., 2020). Mercury has a series of magnetospheric structures (e.g. polar cusp, magnetotail plasma sheet, plasma mantle) and substorm activities (e.g., dipolarization, flux rope) that are similar to those of Earth (Raines et al., 2011; Gershman et al., 2014; Sun, Slavin, Fu, Raines, Zong, et al., 2015; DiBraccio et al., 2015; Imber & Slavin, 2017; Zhao et al., 2019). Recent observation study reports that Mercury also has a ring current that consist of  $\sim \text{keV}$  proton within the magnetosphere (Zhao et al., 2022). Apart from these structures and activities, magnetohydrodynamic and plasma waves have also been reported and investigated in previous studies (e.g., Sundberg et al., 2012; Boardsen et al., 2012, 2015; Kim et al., 2016).

Ultra-low frequency (ULF) pulsations are one of the most common waves in space plasma (Jacobs & Westphal, 1964). They play essential roles in both the magnetosphere and solar wind. These waves are efficient mediums for energy transport (e.g., fast mode magnetosonic waves at the arrival of interplanetary shock (Q.-G. Zong et al., 2009)). Fluctuating electromagnetic fields can directly change the energy and momentum of charged particles and result in modulation, energization, transport, and loss of charged particles (Elkington et al., 2003; Q.-G. Zong et al., 2009; Loto'aniu et al., 2010; Yue et al., 2016; Q. Zong et al., 2017; Z.-Y. Liu et al., 2020). These changes in charged particles may further influence the amplitudes and propagations of the original waves or excite other types of waves (Kim et al., 2016; S. Liu et al., 2019).

66 ULF pulsations near Mercury have been investigated by the observation of Mariner  
 67 10 and MESSENGER. The earliest study can be traced back to a Mariner 10 observa-  
 68 tion of  $\sim 0.5$  Hz linearly polarized magnetic field fluctuations with both compressional  
 69 and transverse modes (C. T. Russell, 1989). The similar type of waves is further inves-  
 70 tigated through MESSENGER observations (e.g., Boardsen et al., 2012, 2015; Kim et  
 71 al., 2016). Based on two years of observations from MESSENGER, Boardsen et al. (2012)  
 72 statistically analyzed and quantified the characteristics of these waves. The frequency  
 73 ranges from 0.4 Hz to 5 Hz, approaching the local proton gyro-frequency. These waves  
 74 are more likely to occur at the duskside rather than at the dawnside. They have the max-  
 75 imum occurrence rate near the radial distance of  $\sim 1.4 R_M$  around the magnetic equa-  
 76 tor. These waves tend to be transverse in the higher magnetic latitudes and tend to be  
 77 compressional near the magnetic equator. These transverse dominant waves (ratio of trans-  
 78 verse power to total power  $> 0.7$ ) are linearly polarized at higher magnetic latitudes. Some  
 79 of these waves have up to at least fourth harmonics.

80 Several studies proposed that these  $\sim 1$  Hz waves may be generated via ion-ion hy-  
 81 brid resonance since the contribution of heavy ions in Mercury's magnetosphere is sig-  
 82 nificant (e.g., Kim et al., 2016). In addition, the observed compressional component can  
 83 be converted from the transverse component via mode conversion. However, from MES-  
 84 SENGER's observation, the compressional mode is always dominant ( $\delta B_{\parallel}^2 > \delta B_{\perp}^2$  for 75%  
 85 wave events), which does not support the ion-ion hybrid resonance hypothesis (Boardsen  
 86 et al., 2012; Kim et al., 2016). Such compressible waves that concentrated near the mag-  
 87 netic equator were interpreted as another mode, electromagnetic ion Bernstein mode waves  
 88 (Boardsen et al., 2015). The electromagnetic ion Bernstein mode waves exhibit large com-  
 89 pressibility in high beta plasma (Denton et al., 2010; Boardsen et al., 2015). Moreover,  
 90 ray tracing simulation (Rönmark & André, 1991; Boardsen et al., 2015) shows that these  
 91 electromagnetic ion Bernstein mode waves propagate between the magnetic latitude of  
 92  $\pm 12^\circ$ , which explains the latitudinal difference in wave compressibility.

93 Long-period (minute-period,  $\sim 15$  mHz, or  $\sim 60$  s) waves were also examined from  
 94 MESSENGER's measurements. Boardsen et al. (2010) and Sundberg et al. (2012) first  
 95 reported minute-period Kelvin-Helmholtz (KH) pulsations at Mercury's magnetopause  
 96 with comprehensive analysis of both the magnetic field and plasma. These KH waves  
 97 are found to drive large amplitude  $\sim 15$  mHz waves inside the magnetosphere (Liljeblad  
 98 et al., 2016). James et al. (2019) proposed that some of these  $\sim 15$  mHz waves are gen-  
 99 erated via field line resonance because they have a dominant toroidal component and re-  
 100 versal of handedness. In a study on the nightside plasma sheet, MESSENGER observed  
 101 series of plasma waves with period of 10 to 20 s, which are near circularly polarized in  
 102 the high magnetic latitude and are compressional waves near the magnetic equator (Sun,  
 103 Slavin, Fu, Raines, Sundberg, et al., 2015).

104 However, the influence of these magnetospheric waves on the plasma has not been  
 105 well investigated in Mercury's magnetosphere, which was partly due to the field-of-view  
 106 (FOV) limitation of MESSENGER's ion measurements. There is still no measurements  
 107 on the direct modulations between waves of different frequencies. In this study, we re-  
 108 port two cases of proton flux fluctuations modulated by  $\sim 15$  mHz compressional waves  
 109 observed by the MESSENGER spacecraft. Coherent  $\sim 1$  Hz waves are also present si-  
 110 multaneously in these events. Their power and compressibility are correlated with the  
 111  $\sim 15$  mHz waves. Moreover, our statistical research verifies that the  $\sim 15$  mHz wave mod-  
 112 ulations on proton flux and high-frequency waves are commonly observed and notewor-  
 113 thy. The paper is organized as follows: First, we will show two case studies: including  
 114 an overview of the wave events and detailed analyses from both magnetic field and par-  
 115 ticle views. Then follow the statistics on the  $\sim 15$  mHz wave modulations covering two  
 116 Mercury years. The possible wave mode and explanations for the observed modulation  
 117 are discussed in the following section. Finally, we will summarize the observations and  
 118 inferences of this study.

## 2 Result

### 2.1 Case I: 19th March, 2014

On 19th March, 2014, MESSENGER's magnetometer (MAG) (Anderson et al., 2007) detected periodic magnetic field fluctuations with a period of  $\sim 64$  s (i.e.  $\sim 15$  mHz) in the duskside magnetosphere of Mercury. The first two panels in Figure 1 show the 46 eV to 13.3 keV energy spectrum and pitch angle distribution of protons measured by the fast imaging plasma spectrometer (FIPS) with 10 s resolution (Andrews et al., 2007). Figure 1c shows the magnetic field vectors with 20 Hz resolution, and Figure 1d shows the magnetic latitude and magnetic local time of MESSENGER. The average energy spectrum during the time interval of interest (UT 10:30 - UT 10:42) is shown as red crosses in Figure 1e. Figure 1f displays the average pitch angle distribution of proton within this interval. MESSENGER's trajectory is presented in Figures 1g and 1h. The time interval of interest is marked by the gray shadow region and bold gray line in Figures 1a-1d and 1g-1h, respectively. The coordinate system used here is the aberrated Mercury Solar Magnetospheric coordinate system (aMSM). In this coordinate system, the X-axis is antiparallel to the solar wind inflow direction (assuming that the solar wind speed is 400 km/s and radial outward). The Z-axis points to Mercury's geographic north pole. The orientation of the Y-axis can be determined by their cross product (i.e.  $\vec{e}_y = \vec{e}_z \times \vec{e}_x$ ). The coordinate system's center deviates from Mercury's geographical center  $0.2 R_M$  northwardly due to the off-centered dipole field (Anderson et al., 2011).

During this interval, MESSENGER moves equatorward from Northern Hemisphere within the meridian plane of  $\sim 18$  h local time. The magnetic field direction changes from downward to northward, and the strength decreases gradually as the spacecraft moves toward the equator. The increasing signature in the proton flux coincides with the previous statistics about the background proton distribution (Zhao et al., 2022). The absence of the protons with a pitch angle of  $\sim 180^\circ$  can be seen in Figure 2b, while the conjugate proton with a pitch angle of  $\sim 0^\circ$  is not measurable due to the limited field of view of FIPS. The  $\sim 64$  s periodic fluctuations with  $\sim 10$  nT amplitude appear in the magnetic field observations between UT 10:30 and UT 10:40. Such  $\sim 64$  s periodic variation also exists in the proton energy spectrum of Figure 1a and pitch angle spectrum of Figure 1b. The magnetic field and proton flux fluctuate irregularly after UT 10:48, implying that the spacecraft enters the magnetosheath. As the proton energy spectrum and pitch angle spectrum do not vary much during this interval, we averaged the proton energy spectrum (Figure 1e) and pitch angle spectrum (Figure 1f) within the time of interest to improve signal to noise ratio. The average energy spectrum shown in Figure 1e greatly exceeds the one-count level except for the low-energy range ( $< 0.12$  keV, the left gray dashed line) and high-energy tail ( $> 4.7$  keV, the right gray dashed line), which means that the proton distributions with energies between the low- and high-energy ends are reliable. The average pitch angle spectrum shown in Figure 1f presents an anisotropic pitch angle distribution. The parallel ( $0^\circ$ - $10^\circ$ ) and anti-parallel ( $150^\circ$ – $180^\circ$ ) pitch angle bins are highly uncertain (the standard error of the mean flux exceeds the 50% of the mean flux) and shown as red bins in the integrated pitch angle distribution.

The proton energy spectrum during the interval with substantial wave activity is presented in Figure 2a. The integrated proton flux within the reliable energy range (0.12 keV to 4.7 keV) is presented as the solid black line, revealing an obvious periodic variation. Pitch angle distributions shown in Figure 2b demonstrates that the periodic variation is almost pitch-angle independent. The magnetic pulsation signals, after subtracting a moving-averaged B field with a window of 300 s, are shown in Figure 2c. Here, the residual magnetic field vectors are transformed into the field-aligned coordinate system. In this coordinate system, the compressional component ( $\delta B_{\parallel}$ ) is parallel to the local magnetic field, and the poloidal and toroidal components ( $\delta B_{po}$  and  $\delta B_{to}$ ) point outward and eastward along the normal plane of the local magnetic field, respectively. The pulsations consist of compressional, toroidal, and poloidal modes. Panel d presents the wavelet spec-

172 trum of the detrended compressional mode signal. The governing component has a pe-  
 173 riod of  $\sim 64$  s, and the secondary component locates around the period of  $\sim 36$  s, which  
 174 is close to the second harmonic period. The power of  $\sim 36$  s fluctuations increases rapidly  
 175 at UT 10:40 as the  $\sim 64$  s component gradually decreases. Here, we focus on the obser-  
 176 vation of monochromatic  $\sim 64$  s waves with clear proton modulation before UT 10:40.

177 In addition,  $\sim 1$  Hz waves are detected during this interval. These waves are an-  
 178 alyzed by the singular value decomposition technique described in Santolík et al. (2003)  
 179 and Boardsen et al. (2015). Figures 3a and 3b present the power spectrum density de-  
 180 rived from wavelet analysis. The solid white line in each panel indicates the proton gyro-  
 181 frequency. The wave power is mainly concentrated at constant frequencies between 0.5  
 182 Hz and 2.0 Hz, close to the proton gyro-frequency. The transverse components are the  
 183 dominant components, and the compressional component becomes comparable as the  
 184 spacecraft approaches the magnetic equator. Meanwhile, the total wave power increases.  
 185 Ellipticity and coherence analyses are presented in Figures 3c and 3d. The low elliptici-  
 186 ty and high coherence characteristics of these waves are consistent with the observational  
 187 characteristics of the waves investigated in Boardsen et al. (2015). The waveforms de-  
 188 trended from the 5 s moving average magnetic field are presented in Figure 3e, show-  
 189 ing a difference in wave power at different latitudes. Apart from the latitudinal varia-  
 190 tions, periodic wave power and components can also be seen, which are likely associated  
 191 with the  $\sim 15$  mHz waves. Here we select three time slices with large amplitude  $\sim 1$  Hz  
 192 waves and show their zoomed-in waveforms in the bottom three panels (Figures 3f-h).  
 193 These waveforms demonstrate that the  $\sim 1$  Hz waves have different amplitudes, compress-  
 194 ibility, and frequency bandwidth at different times.

195 Here, we focus on the five periods of  $\sim 15$  mHz waves from UT 10:31:30 to UT 10:39:00  
 196 with clear modulation of the  $\sim 1$  Hz waves. Figures 4a, 4b, and 4c are adapted from Fig-  
 197 ures 2a, 2c, and 3e to show the proton energy spectrum, the  $\sim 15$  mHz wave signal, and  
 198 the  $\sim 1$  Hz wave signal. Integral wave power within 0.5 Hz to 2 Hz is presented in Fig-  
 199 ure 4d. The compressional and transverse mode power time series are plotted as solid  
 200 red and blue lines, respectively. We use compressibility (i.e., the ratio of compressional  
 201 mode and total power) to quantify the wave mode, as shown in Figure 4e. There are one-  
 202 to-one correspondences between the  $\sim 15$  mHz compressional mode wave dips (Figure  
 203 4b) and  $\sim 1$  Hz wave compressibility peaks. In addition, the integral power of the  $\sim 1$  Hz  
 204 waves (Figure 4d) also seems to be modulated by the  $\sim 15$  mHz compressional waves, al-  
 205 though their one-to-one correspondences are not very distinct. The integral wave power  
 206 and compressibility are Fourier transformed to explore their temporal variation. The re-  
 207 sults shown in Figure 4f demonstrate that they vary with a common period of  $\sim 64$  s,  
 208 coinciding with the frequency of long-period  $\delta B_{\parallel}$  and 0.12 keV to 4.7 keV proton flux  
 209 variations observed during the same interval. Additional power density peaks at approx-  
 210 imately  $\sim 1/36$  Hz are shown in the dynamic spectrum of the  $\sim 1$  Hz wave power and com-  
 211 pressibility.

## 212 2.2 Case II: 4th July, 2013

213 On 4th July 2013, MESSENGER detected another ULF waves-proton modulation  
 214 event (Figure 5, in the same format as Figure 1) in the dawnside flank region. Appar-  
 215 ent magnetic field perturbation can be seen between UT 00:22 and UT 00:30 (Figure 5c).  
 216 We adopted the same method used in Case I to analyze this event. The proton energy  
 217 spectrum (Figure 5a) and pitch angle spectrum (Figure 5b) do not present abrupt changes  
 218 during this interval. The averaged proton energy spectrum presented in Figure 5e re-  
 219 veals a significant count within 0.2 keV and 8.4 keV and maximum counts around  $\sim 1$   
 220 keV. And the averaged pitch angle spectrum shown in Figure 5f indicates an even larger  
 221 anisotropy than the first case. The background magnetic field is about 100 nT and dom-  
 222 inated by the northward component. This event is observed at no more than  $15^{\circ}$  mag-

netic latitudes (Figure 5d). As the spacecraft moves out of the equator (Figures 5g,h), the observed proton flux decrease, and the magnetic field intensity increase.

The proton modulation analysis is presented in Figure 6. Panels a, b, c, and d show the zoom-in proton energy spectrum, pitch angle spectrum, detrended magnetic field, and compressional wavelet spectrum. Five one-to-one correspondences between compressional ULF waves and integral proton fluxes are observed in this event. The magnetic field fluctuations have amplitudes  $\sim 5$  nT and are mainly contributed by the compressional component. The wavelet spectrum suggests the wave periods are 90 s and 55 s (11 and 18 mHz,  $\sim 15$  mHz) at the magnetic equator and higher magnetic latitude. Compared to the first case, the ULF wave amplitudes are lower, and the transverse components are negligible to the compressional component. The proton flux fluctuations are also less significant.

Proton-gyro frequency waves are observed within the same duration. Figures 7a and 7b show the compressional and transverse power spectra within this duration, respectively. The proton gyro-frequency is shown as the solid white line in both panels. The second harmonic waves around  $\sim 2$ -3 Hz are revealed between UT 00:25-00:30. Ellipticity and coherence spectrum are presented in the following c, d panels. The detrended waveforms shown in Figure 7e suggest both the overall increasing trend and periodic fluctuations of the compressional and transverse wave amplitudes as the spacecraft moves poleward. Also, both waves tend to be monochromatic at higher magnetic latitudes and broadband around the equator. These waves present constant central frequency, low ellipticity, and high coherence. These characteristics are coincident with the first case. Figures 7f, g, and h present three exemplified waveforms during this event. Unlike the first case, the compressional mode waves are always dominant.

Figure 8 shows the modulation relationship between the compressional long-period ULF waves, proton flux, and short-period wave powers. It should be noted that the long-period ULF waves do not modulate the  $\sim 1$  Hz wave compressibility in this case. The lack of compressibility modulation indicates that long-period ULF waves might not be sufficient to drive perturbation of proton flux and high-frequency waves every time. A statistical investigation is required to confirm the long-period ULF waves' common influence on particles and high-frequency waves.

### 2.3 Statistical Result

The above cases reveal similar observational characteristics: The long-period ULF waves modulate proton flux and short-period ULF wave power. To understand the importance and role of the ULF waves, we furtherly investigate another 142 long-period ULF wave cases from 1st July 2013 to 30th September 2013 and from 1st January 2014 to 31st March 2014 (consisting of 546 MESSENGER orbits,  $\sim 2$  Mercury years). These ULF waves have periods between 30 - 120 s, durations longer than three wave periods, and an amplitude larger than 1 nT. The case list is shown in the Supporting Information(SI). Among the 142 cases, 63 present modulation (at least three wave periods with one-to-one correspondences). 31 of the 63 proton-modulation cases and 16 of the 79 non-proton-modulation cases reveal high-frequency waves with power modulated by the long period ULF waves. These results imply that the long-period ULF waves-proton flux modulation and long-period ULF waves-short period ULF wave modulation occur with considerable frequencies in Mercury's magnetosphere.

## 3 Discussion and Summary

Although we cannot determine the exact wave modes of the  $\sim 1$  Hz waves and  $\sim 15$  mHz waves, our observations provide some pieces of evidence to diagnose the wave property.

272 The  $\sim 15$  mHz compressional waves play the most crucial role in this event. Their  
 273 frequency is much lower than the proton gyro-frequency and  $\sim 2$  times lower than the  
 274  $\text{Na}^+$  gyro-frequency. Compared to the frequency ( $\sim$  several Hz) of ion-ion hybrid res-  
 275 onance and electromagnetic ion Bernstein mode waves (Kim et al., 2016; Boardsen et  
 276 al., 2015), the frequency of  $\sim 15$  mHz is much lower. Therefore, these waves are more likely  
 277 magnetohydrodynamic waves than kinetic plasma waves. According to the out-of-phase  
 278 correlation between proton flux and compressional component of magnetic field, slow mode  
 279 magnetosonic waves and drift mirror mode waves are two possible candidates for the ob-  
 280 served  $\sim 15$  mHz waves. For slow magnetosonic waves, the thermal pressure gradient force  
 281 and magnetic pressure gradient force are out-of-phase coupled and contribute to the restor-  
 282 ing force together, naturally explaining the one-to-one correspondence of proton flux. The  
 283 associated transverse mode can be regarded as Alfvén waves coupled with magnetosonic  
 284 waves, as in the terrestrial magnetosphere. However, in the low-beta magnetosphere ( $\beta < 1$ ),  
 285 the slow magnetosonic waves have a phase speed on the same order of ion thermal speed,  
 286 which indicates that the waves will be damped via Landau damping (Southwood & Hughes,  
 287 1983). The observed long-life waves do not agree with the theoretical expectation. One  
 288 possible explanation is that long-lasting drivers, such as KH waves at the magnetopause,  
 289 continuously excite the  $\sim 15$  mHz waves (Liljeblad et al., 2016). For the mirror mode (or  
 290 drift mirror waves), the particle is modulated via mirror effects. The modulation efficiency  
 291 of the mirror mode waves depends on the anisotropy of plasma and the loss cone size of  
 292 the magnetic bottle. This interpretation can be supported by the observed anisotropies  
 293 in Figures 1f and 5f while further quantitative analysis is still required. Other candidates,  
 294 including heavy ion cyclotron waves, are not considered here due to insufficient obser-  
 295 vation. Apart from compressional waves, the transverse components are also notable.  
 296 However, they are highly coherent the compressional waves. And, the one-to-one cor-  
 297 respondences between the transverse modes and proton flux are not as compatible as that  
 298 between the compressional mode and proton flux. So, we cannot determine what roles  
 299 do they play in this event.

300 Compared to the ULF wave-particle modulation on Earth, the large amplitude (rel-  
 301 ative to the ambient field) long-period compressional ULF waves are capable of causing  
 302 a more substantial mirror effect. This effect could be signified by the large ratio of pro-  
 303 ton flux maximum and minimum. On the other hand, the relatively large amplitude also  
 304 suggests the presence of the non-linear growth stage of the wave development. The ap-  
 305 parent asymmetric wave peaks and dips in Case I could be one of the pieces of evidence  
 306 for the above inference. Besides, resonant modulations (e.g., bounce resonance, drift res-  
 307 onance) are essential in explaining the wave-particle modulation on Earth. Nevertheless,  
 308 there is no clear energy dependence or pitch angle dependence of the modulation shown  
 309 in the present observation.

310 The  $\sim 1$  Hz waves observed in these events have similar characteristics to the waves  
 311 reported by C. Russell et al. (1988) and Boardsen et al. (2009). Previous interpretations  
 312 of these waves include ion-ion hybrid waves (for the transverse dominant waves) and elec-  
 313 tromagnetic ion Bernstein mode waves (for the compressional dominant waves at lower  
 314 magnetic latitudes). In this case, both compressional dominant and transverse dominant  
 315 waves are observed between  $10^\circ$  S and  $40^\circ$  N magnetic latitude. So, the observed waves  
 316 are possible a mixture of two kinds of waves. The compressibility and power of these waves  
 317 are overall higher near the magnetic equator than at higher latitudes, coinciding with  
 318 the previous statistics and theoretical interpretations. Apart from the latitudinal differ-  
 319 ence, the  $\sim 1$  Hz wave power and compressibility have an additional periodic variation  
 320 with periods of long-period ULF waves. This periodicity suggests the  $\sim 1$  Hz waves may  
 321 be modulated by the long period compressional ULF waves via the changes in plasma  
 322 flux (or pressure) since the plasma beta may influence the growth of plasma instability  
 323 and mode conversion of waves. The unchanged  $\sim 1$  Hz wave frequency implies that the  
 324 excitation and modulation of the  $\sim 1$  Hz waves might not occur locally, unlike the mod-  
 325 ulation on proton flux. This possibility is also supported by the statistical fact that the

326 proton flux modulation and high-frequency wave modulation are not always observed  
 327 together. However, these deductions still have large uncertainty based on the existing  
 328 observations. Several parameters (e.g., heavy ion density, plasma beta, ion perpendic-  
 329 ular and parallel temperatures) are also crucial to the excitation of ion-ion hybrid waves  
 330 and ion Bernstein mode waves. They may also be determinative in the wave-wave mod-  
 331 ulation and are not considered at the present stage (Denton et al., 2010; Kim et al., 2016).  
 332 The charged particle and electromagnetic field instruments onboard BepiColombo are  
 333 also expected to unveil nature of these waves furtherly (Benkhoff et al., 2021)

334 This paper can be summarized as follows:

- 335 1. Our observations confirm that  $\sim 15$  mHz compressional waves can modulate pro-  
 336 ton flux significantly.
- 337 2. Large amplitudes out-of-phase one-to-one correlations between the magnetic field  
 338 intensity and proton flux indicates that ULF waves control the plasma environment.
- 339 3. The  $\sim 1$  Hz coherent wave power and compressibility observed during the same  
 340 interval are modulated by  $\sim 15$  mHz compressional waves. This modulation may be im-  
 341 plemented via the changes in plasma flux, and it helps us better understand the ubiq-  
 342 uitous  $\sim 1$  Hz waves in Mercury's magnetosphere.

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 351 [-3-CDR-CALIBRATED-V1.0](https://pds-ppi.igpp.ucla.edu/search/view/?f=yes&id=pds://PPI/MESS-E\_\V\_H\_SW-MAG-3-CDR-CALIBRATED-V1.0) and Fast Imaging Plasma Spectrometer: [https://pds-ppi](https://pds-ppi.igpp.ucla.edu/search/view/?f=yes&id=pds://PPI/MESS-E\_\V\_H\_SW-EPPS-3-FIPS-DDR-V2.0)  
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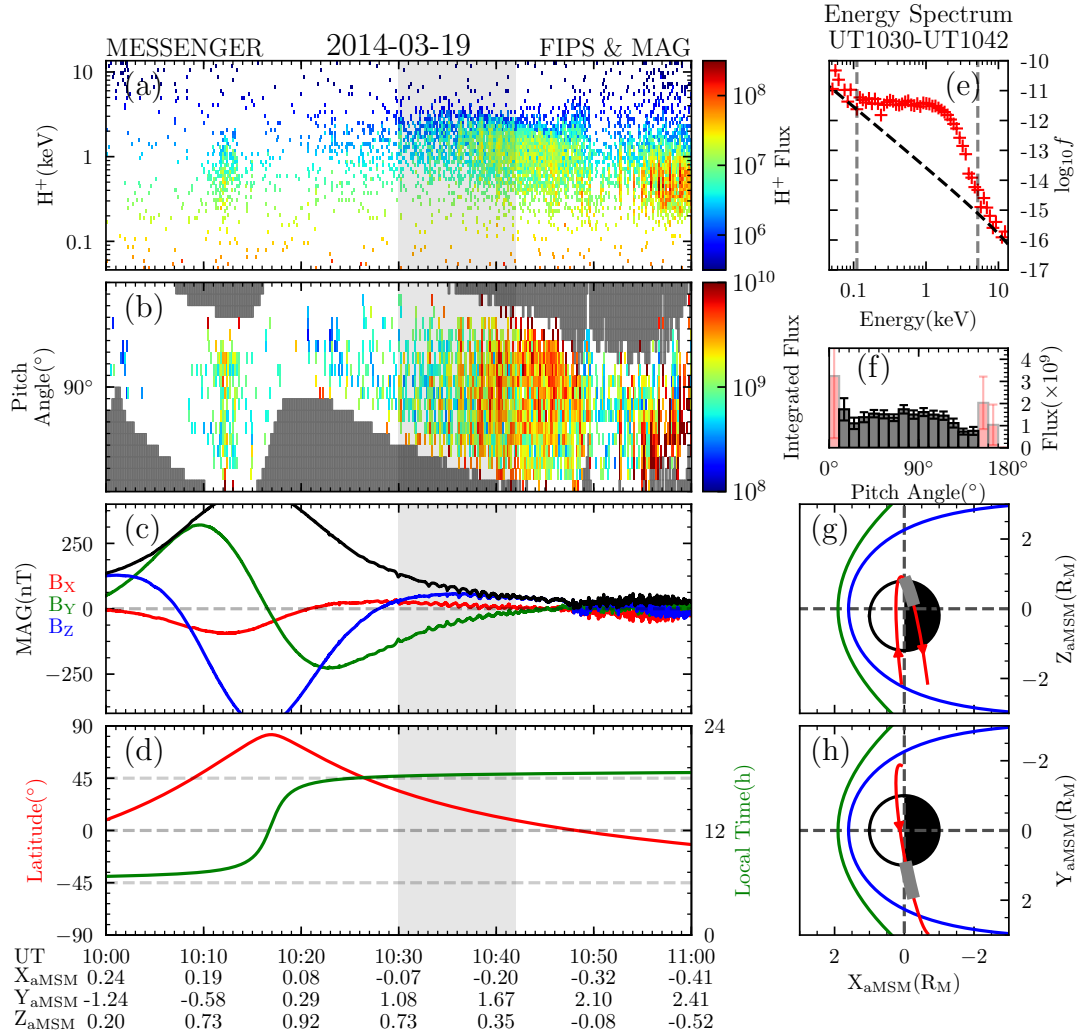
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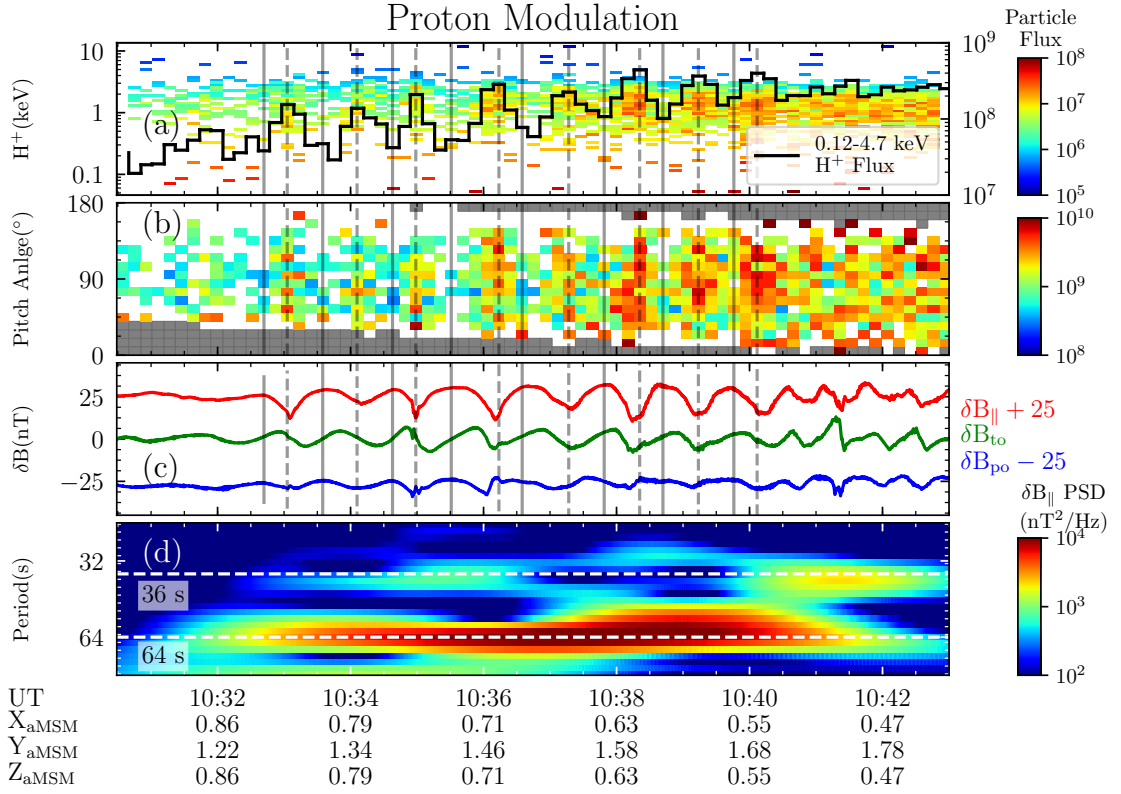
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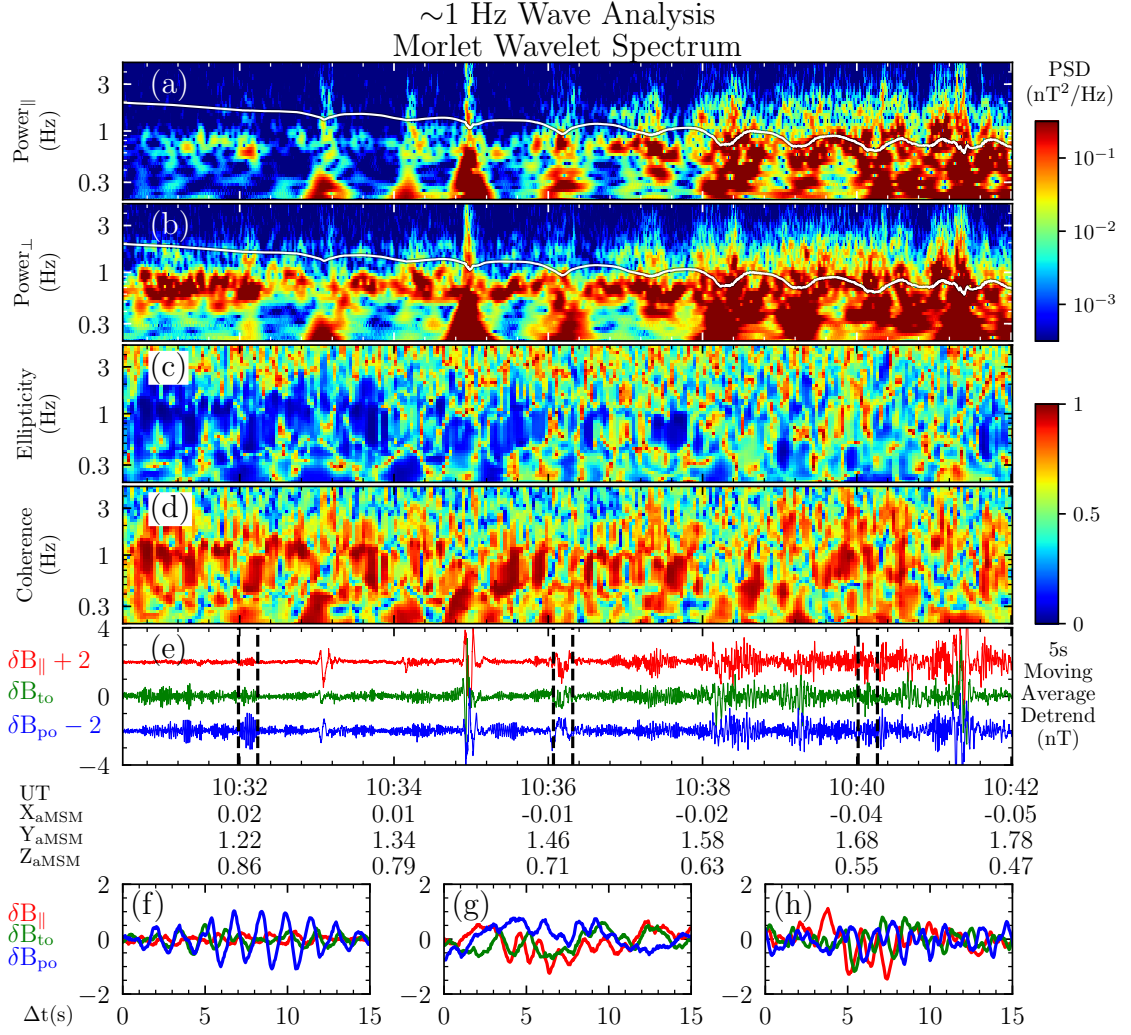
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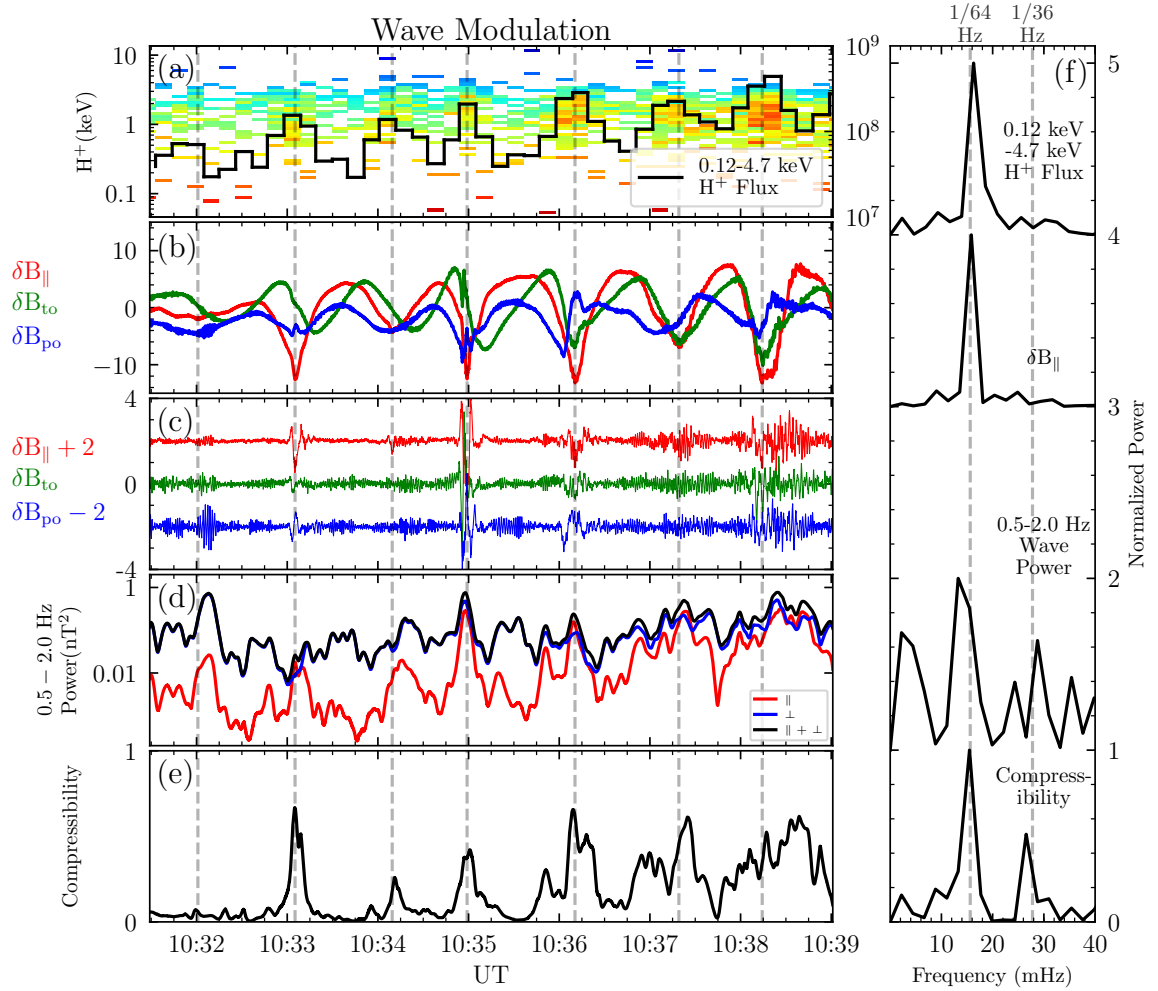
**Figure 1.** An overview of the ultra-low frequency pulsations event on 19th March, 2014 observed by MESSENGER’s MAG and FIPS. (a) Differential number flux (in the unit of  $\text{cm}^{-2} \cdot \text{s}^{-1} \cdot (\text{keV}/e)^{-1} \cdot \text{sr}^{-1}$ ) of protons measured by FIPS. (b) Pitch angle distribution (in the unit of  $\text{cm}^{-2} \cdot \text{s}^{-1} \cdot \text{sr}^{-1}$ ). The pitch angle bins out of the field-of-view are filled with dark gray. (c) Magnetic field components (solid red, green, and blue lines for  $B_x$ ,  $B_y$  and  $B_z$ ) and intensity (solid black line). (d) Magnetic latitude and local time of the spacecraft. (e) The average energy spectrum of protons within the duration from UT 10:30 to UT 10:42. (f) Average pitch angle spectrum of protons within the duration from UT 10:30 to UT 10:42. Errorbars here are the standard errors of the mean. Pitch angle bins with large relative uncertainty ( $>50\%$ ) are marked with red color. (g, h) Trajectories of the spacecraft in the XY, XZ planes. The interval of interest is marked by the light gray shaded area in panels a-d and bold gray line in panels f-g. Universal time and spacecraft location in aMSM coordinate are presented at the bottom of the figure.



**Figure 2.** Proton modulation by  $\sim 15$  mHz compressional ULF waves. (a) The differential number flux of the proton (in the unit of  $\text{cm}^{-2} \cdot \text{s}^{-1} \cdot (\text{keV}/e)^{-1} \cdot \text{sr}^{-1}$ ). The integrated flux of 0.12 keV - 4.71 keV protons is presented as the solid black line (in the unit of  $\text{cm}^{-2} \cdot \text{s}^{-1} \cdot \text{sr}^{-1}$ ). (b) Pitch angle distribution (in the unit of  $\text{cm}^{-2} \cdot \text{s}^{-1} \cdot \text{sr}^{-1}$ ). (c) The magnetic field detrended from the 300 s moving average in the field-aligned coordinate system. (d) Wavelet spectrum of the compressional mode magnetic field pulsations. The dashed white line indicates the periods of 36 s and 64 s.

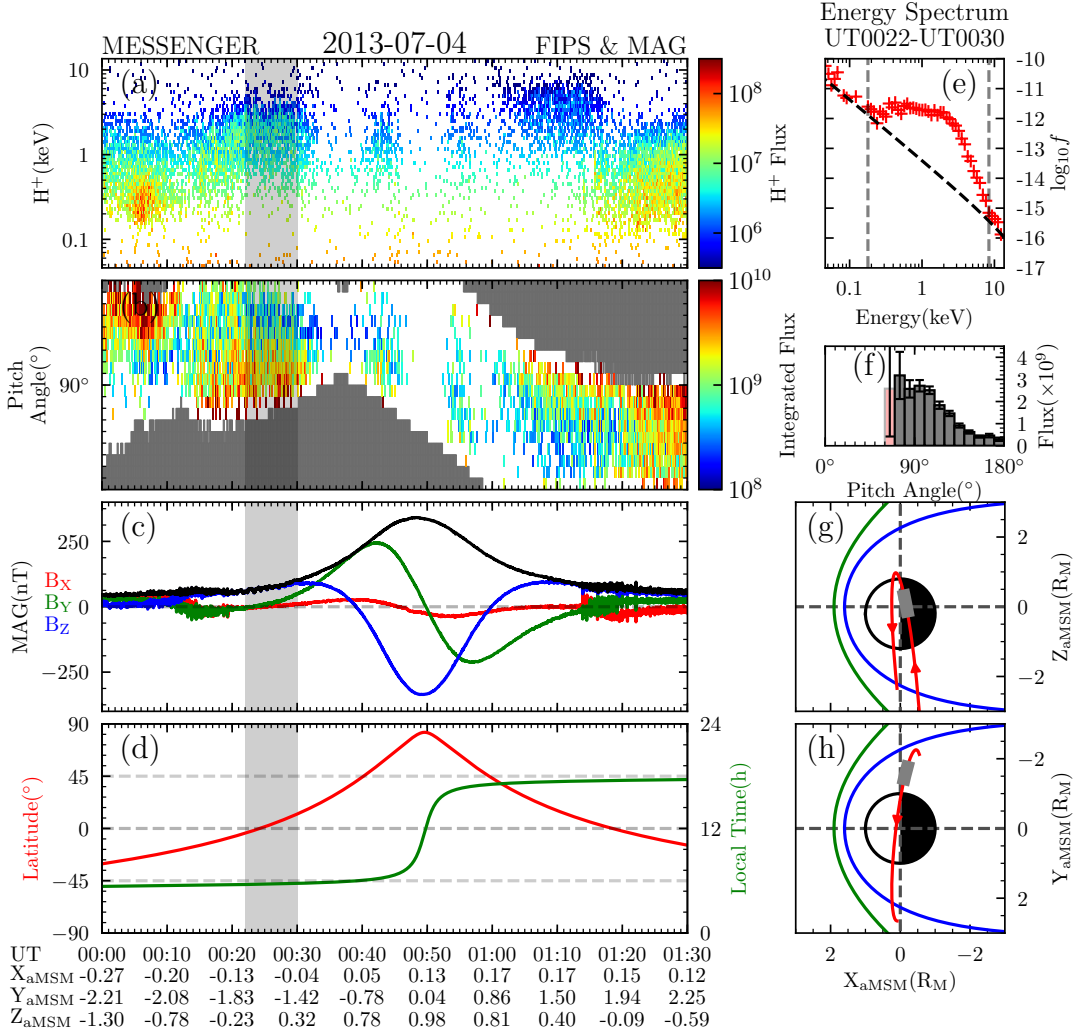


**Figure 3.**  $\sim 1$  Hz coherent waves analysis. (a) Power spectrum of the compressional component, (b) and the transverse component. (c) Ellipticity. (d) Coherence. (e) The magnetic field waveform detrended from a 5 s moving average. (f, g, h) Three slices of  $\sim 1$  Hz coherent waves with a detrended magnetic field. The time ranges of these three slices are marked as dashed black lines in panel e.

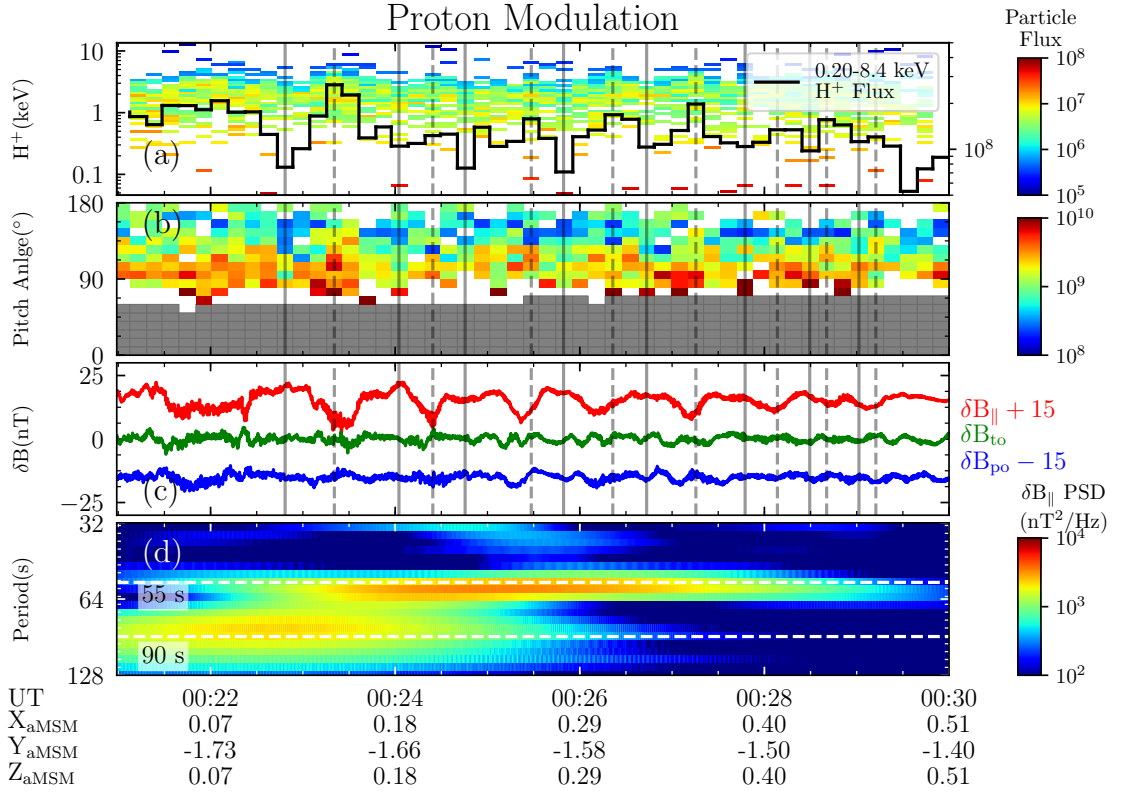


**Figure 4.** The one-to-one correspondence of  $\sim 1$  Hz waves and  $\sim 15$  mHz waves. (a) Proton spectrum adapted from Figure 2a. The solid black line displays the 0.12 keV to 4.7 keV proton flux. (b) The waveform of detrended  $\sim 15$  mHz waves. (c) The waveform of detrended  $\sim 1$  Hz waves. (d) 0.5-2.0 Hz integral wave power. (e) The ratio between 0.5-2.0 Hz integral compressional wave power and 0.5-2.0 Hz integral total wave power. (f) Discrete Fourier transform of 0.12 keV to 4.7 keV proton flux (solid black line in panel a),  $\sim 15$  mHz compressional waves (solid red line in panel b), integral wave power (solid black line in panel d), and compressibility (solid black line in panel e) of the  $\sim 1$  Hz waves.

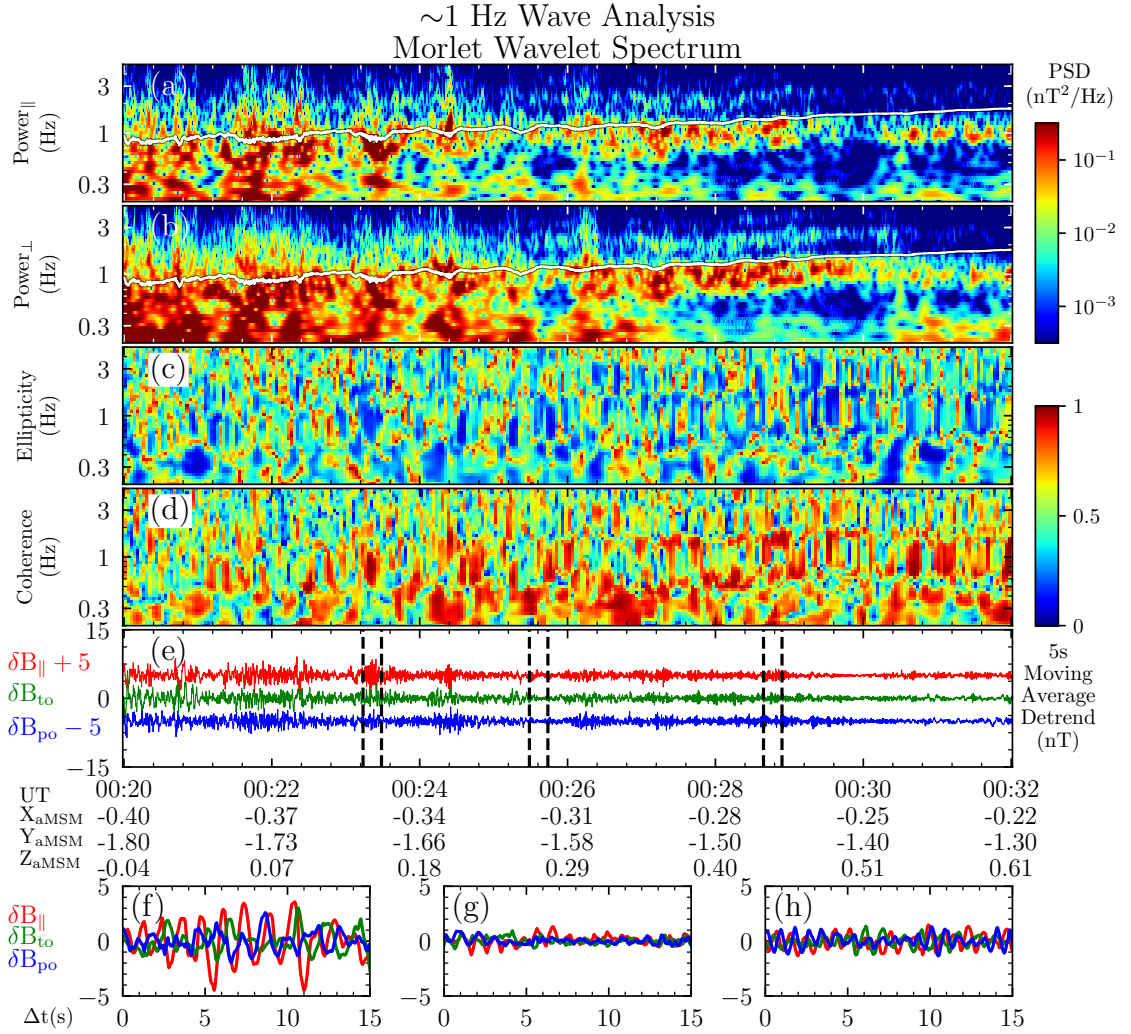




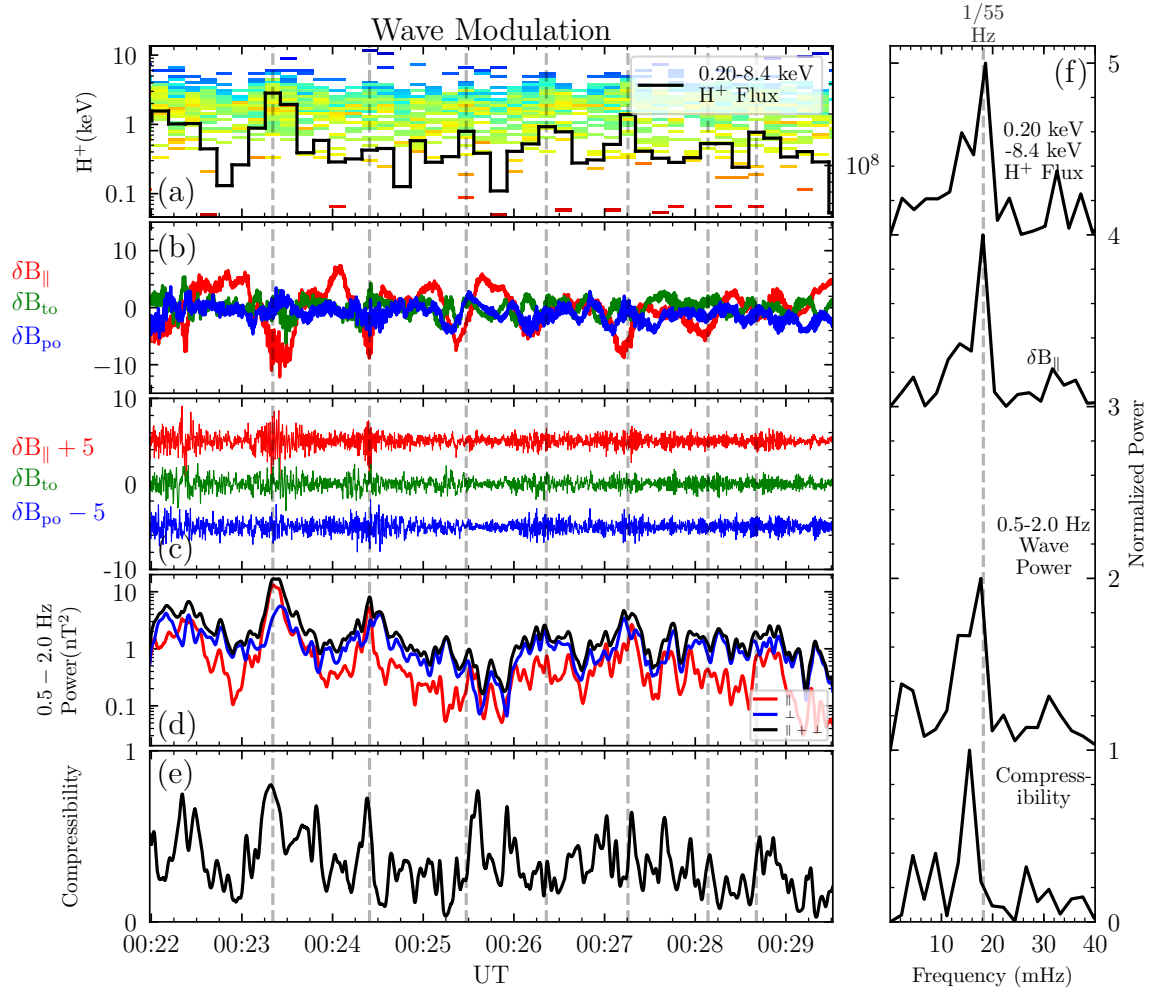
**Figure 5.** An overview of the ultra-low frequency pulsations event on 4th July, 2013 observed by MESSENGER’s MAG and FIPS. (a) Differential number flux (in the unit of  $\text{cm}^{-2} \cdot \text{s}^{-1} \cdot (\text{keV}/e)^{-1} \cdot \text{sr}^{-1}$ ) of protons measured by FIPS. (b) Pitch angle distribution (in the unit of  $\text{cm}^{-2} \cdot \text{s}^{-1} \cdot \text{sr}^{-1}$ ). The pitch angle bins out of the field-of-view are filled with dark gray. (c) Magnetic field components (solid red, green, and blue lines for  $B_x$ ,  $B_y$  and  $B_z$ ) and intensity (solid black line). (d) Magnetic latitude and local time of the spacecraft. (e) The average energy spectrum of protons within the duration from UT 00:22 to UT 00:30. (f) Average pitch angle spectrum of protons within the duration from UT 00:22 to UT 00:30. Errorbars here are the standard errors of the mean. Pitch angle bins with large relative uncertainty ( $>50\%$ ) are marked with red color. (g, h) Trajectories of the spacecraft in the XY, XZ planes. The interval of interest is marked by the light gray shaded area in panels a-d and gray red line in panels f-g. Universal time and spacecraft location in aMSM coordinate are presented at the bottom of the figure.



**Figure 6.** Proton modulation by  $\sim 15$  mHz compressional ULF waves. (a) The differential number flux of the proton (in the unit of  $\text{cm}^{-2} \cdot \text{s}^{-1} \cdot (\text{keV}/e)^{-1} \cdot \text{sr}^{-1}$ ). The integrated flux of 0.20 keV - 8.4 keV protons is presented as the solid black line (in the unit of  $\text{cm}^{-2} \cdot \text{s}^{-1} \cdot \text{sr}^{-1}$ ). (b) Pitch angle distribution (in the unit of  $\text{cm}^{-2} \cdot \text{s}^{-1} \cdot \text{sr}^{-1}$ ). (c) The magnetic field detrended from the 300 s moving average in the field-aligned coordinate system. (d) Wavelet spectrum of the compressional mode magnetic field pulsations. The dashed white line indicates the periods of 55 s and 90 s.



**Figure 7.**  $\sim 1$  Hz coherent waves analysis. (a) Power spectrum of the compressional component, (b) and the transverse component. (c) Ellipticity. (d) Coherence. (e) The magnetic field waveform detrended from a 5 s moving average. (f, g, h) Three slices of  $\sim 1$  Hz coherent waves with a detrended magnetic field. The time ranges of these three slices are marked as dashed black lines in panel e.



**Figure 8.** The one-to-one correspondence of  $\sim 1$  Hz waves and  $\sim 15$  mHz waves. (a) Proton spectrum adapted from Figure 6a. The solid black line displays the 0.20 keV to 8.4 keV proton flux. (b) The waveform of detrended  $\sim 15$  mHz waves. (c) The waveform of detrended  $\sim 1$  Hz waves. (d) 0.5-2.0 Hz integral wave power. (e) The ratio between 0.5-2.0 Hz integral compressional wave power and 0.5-2.0 Hz integral total wave power. (f) Discrete Fourier transform of 0.20 keV to 8.4 keV proton flux (solid black line in panel a),  $\sim 15$  mHz compressional waves (solid red line in panel b), integral wave power (solid black line in panel d), and compressibility (solid black line in panel e) of the  $\sim 1$  Hz waves.