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The calculated and observed ionospheric properties during Atmospheric Explorer-C satellite crossings over Millstone Hill

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Abstract-The Atmospheric Explorer-C (AE-C) satellite passed almost directly over the Millstone Hill incoherent scatter radar station on 14 February 1974 and passed within the near vicinity of the station on 15 February 1974. Measurements of ionospheric and atmospheric properties were made simultaneously by the incoherent scatter radar and the AE-C satellite instruments. The incoherent scatter radar measured vertical profiles of the electron and ion temperatures and electron density and these data were used to derive a neutral gas temperature profile. The AE-C satellite measured the electron and ion densities and electron and ion temperatures, neutral gas composition, solar EUV flux, photoelectron spectra, the 6300 Å volume emission rate profile and the distribution of NO along the satellite path. These simultaneous measurements provide a consistent set of data to examine current F-region theory in the daytime ionosphere. We used a time-dependent coupled model of the ionospheric E- and F-region to calculate the ionospheric properties over Millstone Hill at the times of the AE-C crossings and then compared the calculated structure to the observed structure. The results show good agreement between the incoherent scatter radar measurements and the model calculations. There is also good agreement among satellite and incoherent scatter radar measurements and model calculations for the altitude of the satellite crossing, 161 km. The satellite measurements along the orbital path, however, reveal considerable horizontal gradients in the measured ionospheric properties.

1. INTRODUCTION

On 14 February 1974 (on orbit 666) the Atmospheric Explorer-C (AE-C) satellite (DALGARNO et al., 1973) passed within 59 km horizontal distance of the incoherent scatter radar station at Millstone Hill, Massachusetts (42.6° N, 71.5°W), and on 15 February 1974 (on orbit 677) the satellite passed within 254 km horizontal distance of the station. Satellite and incoherent scatter radar data on the ionospheric and atmospheric properties were obtained simultaneously in nearly the same spatial regions and at the same time of day on each satellite pass. The incoherent scatter radar meas-

ured the vertical profiles of electron density and electron and ion temperatures; the vertical distribution of neutral gas temperature and the exospheric temperature were derived from these data. The AE-C satellite measured the electron and ion temperatures, electron and ion densities, neutral gas composition and photoelectron spectra along the satellite path. In addition, measurements were made of the solar EUV and UV fluxes, the vertical profile of the 6300 Å airglow volume emission rate and the vertical profile of the NO number density. These simultaneous satellite and incoherent scatter radar measurements provided a consistent set of data that we used to examine our understanding of E- and F-region theory by comparing the data to model calculations.

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The ionospheric model used to calculate the ionospheric properties over Millstone Hill at the time of the AE-C satellite crossing has been described in detail by ROBLE (1975), STORARSKI et al. (1975) and ROBLE et al. (1976). The model was expanded recently to include auroral processes, as described by ROBLE and REES (1977), and to calculate the distribution of the minor neutral constituents $N(^{2}D)$, $N(^{4}S)$ and NO. The ionospheric model was also modified to include the reaction rates, cross sections, etc., that were derived from previous AE-C studies. The satellite crossings over Millstone Hill occurred near 1130 LT. Therefore, the thermal and chemical properties of the midlatitude daytime ionospheric E- and F-regions were calculated with the model and compared to measurements made by the AE-C satellite and Millstone Hill incoherent scatter radar.

Our results indicate that current ionospheric Eand F-region theory is adequate for describing the gross ionospheric and atmospheric properties over Millstone Hill on the two days when radar and satellite data were obtained simultaneously. Reasonable agreement between the calculated and observed values was obtained for the electron and ion temperatures, electron and ion densities, 6300 Å volume emission rate profile, NO number density profile and photoelectron spectra on both satellite passes of Millstone Hill. The satellite observations show considerable latitudinal gradients in the measured properties, thereby indicating that caution is necessary in the derivation of vertical profiles of ionospheric properties from satellite measurements along the orbital path.

2. IONOSPHERIC MODEL

The ionospheric model used in these calculations is an extension of the model described by ROBLE (1975), STOLARSKI et al. (1975) and ROBLE et al. (1976). A block diagram of the model is shown in Fig. 1. Because most of the processes within the model have been described previously, we only discuss the modifications and additions made for this study. The main addition to the model is the inclusion of minor neutral constituent chemistry and coupling with direct particle precipitation in auroral processes, as discussed by ROBLE and REES (1977).

The ion chemistry in the model was updated from that given in Table 1 of ROBLE (1975) to include many of the processes analyzed by the

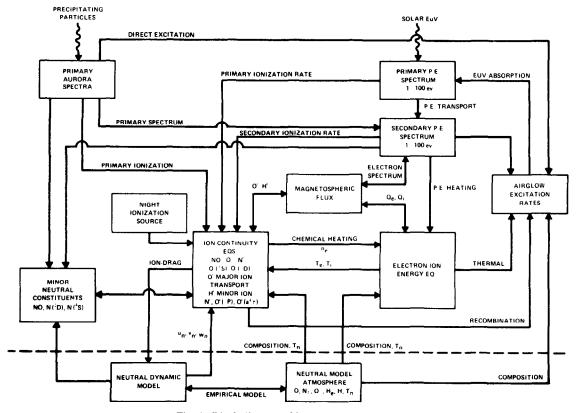


Fig. 1. Block diagram of ionospheric model.

Reaction	Reaction rate (cm ³ s ⁻¹)	Reference
$O_2^+ + N_2 \rightarrow NO^+ + NO$	5×10 ⁻¹⁶	FERGUSON (1974)
$O_2^+ + NO \rightarrow NO^+ + O_2$	4.4×10^{-10}	LINDINGER et al. (1975)
$O_2^+ + N(^2D) \rightarrow N^+ + O_2$	2.5×10^{-10}	Dalgarno (1970)
$O_2^+ + e \rightarrow O + O(^3P, {}^1D, {}^1S)$ [B1]	$1.1 \times 10^{-5} T_{\bullet}^{-0.7}$ for $T_{\bullet} \le 1000 \text{ K}$ $1.6 \times 10^{-7} (T_{\bullet}/300)^{-0.55}; T_{\bullet} > 1000 \text{ K}$	Torr et al. (1976)
$O_2^+ + N(^4S) \rightarrow NO^+ + O$	1.8×10^{-10}	GOLDAN et al. 1966)
$NO^+ + e \rightarrow O + N(^2D, ^4S) [B3]$	$5 \times 10^{-7} (T_e/300)^{-1.0}$ for $T_e > 1000$ K $4.2 \times 10^{-7} (T_e/300)^{-0.85}$; $T_e \le 1000$ K	TORR et al. (1976)
$O^+(^4S) + H \rightarrow H^+ + O$	6.0×10^{-10}	FERGUSON (1974)
$O^+(^4S) + N_2 \rightarrow NO^+ + N(^4S)$	5.0×10^{-13} ; $T_{\rm R} < 1000$ K $4.5 \times 10^{-14} (T_{\rm R}/300)^2$; $T_{\rm R} \ge 1000$ K	LINDINGER et al. (1974)
$O^+(^4S) + O_2 \rightarrow O_2^+ + O$	$2.0 \times 10^{-11} (T_R/300)^{-0.4}; T_R < 1800 \text{ K}$ $1.3 \times 10^{-12} (T_R/300)^{-1.2}; T_R \ge 1800 \text{ K}$	McFarland et al. (1973)
$O^+(^4S) + NO \rightarrow NO^+ + O$	8×10^{-13}	Ferguson (1974)
$O^+(^2D) + N_2 \rightarrow NO^+ + N(^4S)$	5.6×10^{-13} for $T_{\rm R} < 1000$ K	LINDINGER (1974)
$O^+(^2D) + e \rightarrow O^+(^4S) + e$	$7.8 \times 10^{-8} (300/T_{\bullet})^{1/2}$	HENRY et al. (1969)
$O^+(^2D) + N_2 \rightarrow N_2^+ + O$	1.0×10^{-9}	RUTHERFORD and VROOM (1971)
$N_2^+ + O_2 \rightarrow O_2^+ + N_2$	$5.0 \times 10^{-11} (T_R/300)^{-0.8}; T_R < 3560 \text{ K}$	LINDINGER et al. (1974)
N_2^+ + O → NO ⁺ + N(² D, ⁴ S) [B2]	$1.4 \times 10^{-10} (T_R/300)^{-0.44} [1 - 0.07 (T_R/300)^{0.21}]; T_R < 1500 \text{ K}$ $5.2 \times 10^{-11} (T_R/300)^{0.2} [1 - 0.07 (T_R/300)^{0.21}]; T_R \ge 1500 \text{ K}$	McFarland et al. (1974)
$N_2^+ + e \rightarrow N + N(^4S, ^2D)$ [B4]	$1.8 \times 10^{-7} (T_{\rm p}/T_{\rm s})^{-0.4}$	Orsini et al. (1976)
	$1.4 \times 10^{-10} (T_R/300)^{-0.44} \cdot 0.07 (T_R/300)^{0.21}; T_R < 1500 \text{ K}$	McFarland et al. (1974)
$N_2^+ + O \rightarrow O^+ + N_2$	$5.2 \times 10^{-11} (T_R/300)^{0.2} \cdot 0.07 (T_R/300)^{0.21}; T_R \ge 1500 \text{ K}$	
$N_2^+ + NO \rightarrow NO^+ + N_2$	3.3×10^{-10}	FEHSENFELD et al. (1970)
$H^+ + O \rightarrow O^+ + H$	$6.0 \times 10^{-10} \frac{8}{9} \left(\frac{T_1 + T_n/16}{T_n + T_1/16} \right)^{1/2}$	BANKS (1967)
$N(^{2}D) + O_{2} \rightarrow NO + O$	5.0×10^{-12}	LIN and KAUFMAN (1971)
$N(^{2}D) + O \rightarrow N(^{4}S) + O$	$<4.5\times10^{-13}$	ORAN et al. (1975) FREDERICK and RUSCH (1977)
$N(^{2}D) + e \rightarrow N(^{4}S) + e$	$1.0 \times 10^{-9} (T_e/300)^{1/2}$	RUSCH et al. (1975)
$N(^{2}D) + NO \rightarrow N_{2} + O$	7×10^{-11}	BLACK et al. (1969)
$N(^{4}S) + O_{2} \rightarrow NO + O$	$2.4 \times 10^{-11} \exp{(-3975/T_n)}$	WILSON (1967)
$N(^4S) + NO \rightarrow N_2 + O$	$1.5 \times 10^{-12} T_n^{1/2}$	PHILLIPS and SCHIFF (1962)
$NO + h\nu \rightarrow N(^{*}S) + O$	8.3×10 ⁻⁶	STROBEL (1971)
$NO + h\nu \rightarrow NO^+ + e$	6.0×10^{-7}	STROBEL (1971)
$N_2 + e(fast) \rightarrow N(^4S) + N(^2D) [B5]$		WINTERS (1966)
$N(^2D) \rightarrow N(^4S) + h\nu$	1.06×10^{-5}	GASTRANG (1956)

Table 1. Ion and neutral chemistry for midlatitude ionsopheric model

 T_e = electron temperature (K)

 $T_i = \text{ion temperature (K)}$

 $T_v = vibrated temperature (K)$

 T_R = relative reactant kinetic temperature (K)

AE-C team. The ion chemistry currently in the model is summarized in Table 1, along with the various reaction rates. The branching ratios for various chemical reaction rates are given in Table 2.

The minor neutral gas chemistry for $N({}^4S)$, $N({}^2D)$ and NO is also given in Table 1. The diffusion equation that is solved for the minor neutral constituents is the same as that used by STROBEL (1971), RUSCH *et al.* (1975) and ORAN *et al.* (1975). The coupled diffusion equations for $N({}^4S)$, $N({}^2D)$ and NO, with the chemical reactions given in Table 1, are solved simultaneously at each time step according to the technique of HASTINGS and ROBLE (1977).

Although the solar EUV flux is routinely measured on board the AE-C sateilite (HINTEREG-GER, 1976), no direct measurements were made for the particular days under consideration. The solar F10.7 cm flux values are 78.3×10^{-22} and 85.8×10^{-22} W m⁻² Hz⁻¹ for 14 and 15 February respectively (Table 3). For this level of solar activity we use the AE-C solar flux reference spectrum discussed by HINTEREGGER (1975) in our calculations. The Hinteregger reference spectrum was obtained from a rocket flight on 23 April 1974, and the solar flux values are similar to those given in HINTEREGGER (1970). The solar EUV flux measurements available in the AE-C data file show that the values measured before and after 14 and 15

Table 2. Branching ratios for reactions labelled (B-)

(B1)—Production of $O(^1D)$	= 1.0
O(1S)	= 0.04
$O(^{3}P)$) = 0.96
(B2)—Production of $N(^4S)$) = 0.05
N(² D	= 0.95
(B3)-Production of N(4S	0 = 0.26
N(² D	0 = 0.74
(B4)-Production of N(⁴ S) = 1.0
$N(^{2}D)$) = 1.0
(B5)—Production of N(⁴ S) = 1.0
) = 1.0

February are very similar to the AE-C reference spectrum. Therefore, for the period during which solar flux values were not measured, it is reasonable to use the AE-C reference spectrum in our calculations.

We use the satellite measurements of the neutral composition over Millstone Hill at an altitude

161 km for our model calculations. The measurements for both days are given in Table 3. The neutral gas temperature profile is determined from incoherent scatter radar data and is fit to the analytic temperature function specified by BATES (1959) for model calculations. The profile is defined by an exospheric temperature, the temperature at 120 km and the 'shape parameter, S,' that specified the temperature gradient in the lower thermosphere. These values are all given in Table 3 for both days on which the AE-C satellite passed over Millstone Hill.

With the temperature profile thus specified, we determine the vertical distribution of neutral density from the analytic profile of BATES (1959) by requiring that the calculated O and N₂ densities at 161 km agree with the satellite measurements of these constituents. The O₂ number density at 120 km is specified as 5×10^{10} cm⁻³, which gives an

Table 3. Parameters at the time of the AE satellite crossings of Millstone Hill

AE orbit number	666	677
Date	14 Feb. 1974	15 Feb. 1974
Universal time, Hr	1634	1614
EST, Hr	1134	1114
Altitude of crossing, km	161.3	162.4
North latitude, degrees	42.6	42.6
West longitude, degrees	70.7	68.0
Solar zenith angle, degrees	55.4	52.0
Satellite measured values at time of crossing		
N_2 density, cm ⁻³	1.29×10^{10}	1.06×10^{10}
O density, cm^{-3}	1.25×10^{10}	1.18×10^{10}
He density, cm ⁻³	2.48×10^{7}	2.13×10^{7}
Electron Temperature, K	1050	1110
Electron density, cm^{-3}	1.33×10^{5}	1.48×10^{5}
Neutral Temperature, K		702
Millstone Hill determined parameters		
Exospheric temperature, K	790 ± 18	779 ± 13
Temperature at 120 K	350 ± 13	301 ± 11
Shape parameter in lower thermosphere S , $K \text{ cm}^{-1}$	0.053 ± 0.008	0.046 ± 0.006
Electron temperature, K, at satellite crossing	960±20(998)	990 ± 20(966)
Electron density, cm^{-3} ,	$1.85 \pm 0.05 \times 10^{5}$	$2.35 \pm 0.05 \times 10^{5}$
at satellite crossing	(1.50×10^5)	(1.55×10^5)
$N(O^+)/N_{\star}$ at satellite crossing	0.15 ± 0.07	0.36 ± 0.08
O density at satellite crossing, cm^{-3}	$0.7 \pm 0.3 \times 10^{10}$	$1.2 \pm 0.4 \times 10^{10}$
Neutral Temperature, K	737	703
Geophysical parameters		
Solar F10.7	78.3	85.8
Sunspot number	26	46
A _p	12	4
Values in parenthesis give model calculations		

 O_2 to N_2 ratio of 0.12 at an altitude of 200 km; this is consistent with the measurements of NIER *et al.* (1976).

The only other parameters necessary for the model calculation are the boundary conditions for the electron and ion energy equations and the O^+ diffusion equation. The lower boundary of the model is at 80 km, where the electron and ion temperatures are assumed to be equal to the neutral gas temperature. At the lower boundary the time-dependent ion chemistry equations are solved without diffusion and the resulting values are used to specify the ion number densities for the ion transport equations. The upper boundary of the model is set at 500 km, the upper limit of the incoherent scatter radar data. The electron and ion heat fluxes are calculated from the temperature and temperature gradient measurements at 500 km and are used as upper boundary conditions for the electron and ion energy equations. The measured electron density is used as a boundary condition for the O^+ transport equation. The O^+ , NO^+ and O_2^+ densities are determined according to the method of SCHUNK and WALKER (1973). Thus, the only specified parameters for these calculations are the solar EUV flux, the neutral atmospheric properties and the boundary conditions for the plasma energy and transport equations. The initial conditions are specified arbitrarily and the model is run to steady state in the daytime ionosphere, where photochemical time constants are fast. The minor neutral constituent chemistry, however, has a much longer time constant and its distribution can be influenced

by neutral dynamic and transport processes. Realizing these limitations, we nevertheless present steady state calculations for comparison with the measured NO density.

SATELLITE MEASUREMENTS

The satellite geographic longitude, geographic latitude, altitude and local solar time relationship for AE-C orbit 666 on 14 February 1974 are shown in Fig. 2. On the inbound pass, the satellite descended from 400 km near 68° N latitude to 160 km near 45° N latitude. It passed over Millstone Hill at 161 km and ascended to 300 km near 20° N latitude. The total change in latitude was 50° , and the total change in longitude was 60° . The local solar time varied between 0900 and 1300 along the satellite orbit, which covered the altitude range from 400 km to 160 km to 300 km. The orbital parameters of the satellite at the time it crossed the latitude of Millstone Hill are summarized in Table 3.

Several instruments on board the AE-C satellite were used to measure various atmospheric parameters over Millstone Hill. The measurements made by neutral mass spectrometers (NIER *et al.*, 1973; PELZ *et al.*, 1973) were used to specify the neutral composition, ion mass spectrometers (BRINTON *et al.*, 1973) specified ion composition along the satellite path and the retarding-potential analyzer (HANSON *et al.*, 1973) gave the ion temperature and density. A cylindrical electrostatic probe (BRACE *et al.*, 1973) was used to specify the electron density and temperature along the satellite

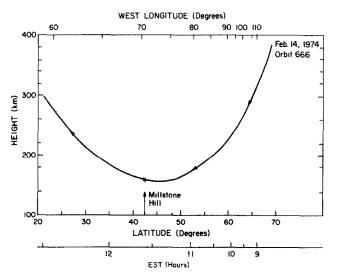


Fig. 2. The longitude, latitude, local solar time and height relationships for the inbound and outbound pass of AE-C satellite orbit 666 on 14 February 1974.

path and the neutral-atmosphere temperature instrument (SPENCER et al., 1973) was used to measure the neutral gas temperature. The solar EUV flux was measured by the EUV spectrophotometer (HINTEREGGER et al., 1973). The 6300 Å volume emission rate profile was determined by the visible airglow experiment (HAYS et al., 1973), and the nitric oxide distribution was obtained from the UV spectrometer measurements (BARTH et al., 1973). Finally, the photo-electron-spectrometer (DOERING et al., 1973) was used to determine the photoelectron spectra at the time of the AE-C crossing over Millstone Hill.

RADAR MEASUREMENTS

The vertically directed UHF (ultra-high frequency; 440 MHz) radar at Millstone Hill (42.6°N, 71.5°W) was used to make incoherent scatter measurements during the overflights of AE-C on orbits 666 and 677. The observations were conducted with the 'two-pulse' program in which pairs of short radar pulses, whose separation is varied, are transmitted during each sweep. A full cycle of operations using basic pulse lengths of 40, 100 and 200 μ s (6, 15 and 30 km height resolution, respectively) required about 30 min and covered the altitude range from 100 to 500 km.

The incoherent scatter signal is processed to determine total power and correlation functions at a number of altitudes. The power measurements permit an altitude profile of electron density to be calculated after normalization to the peak density by using ionosonde records of foF2. If the ion composition is known, both the electron and ion temperatures can be determined directly from the shape of the measured incoherent scatter correlation functions. Direct determination can be made between 115 and 130 km, where the dominant ions are NO^+ and O_2^+ (we assume an ion mass of 31 amu), and also from 225 to 500 km, where O⁺ is the dominant ion. Between 130 and 225 km only two of the three unknowns $(T_e, T_i \text{ and } q =$ $N(O^+)/N_e$ can be determined from the measurements, since there is a transition from molecular ions (q=0) to atomic ions (q=1). Our method of resolving this problem was to construct a model of neutral temperature, T_n , by using only the data outside the transition region, where T_e and T_i are easily determined, and then predicting T_n over the transition region heights from this model. Since $T_i = T_n$ over the transition altitude range, in effect we had determined T_i . The remaining unknowns, T, and q, are obtained directly from the measured correlation functions.

For the T_n model the BATES (1959) form is adopted,

$$T_n(h) = T_{\infty} - (T_{\infty} - T_0) \exp(-s(h - h_0))$$
 (1)

where h is altitude and $h_0 = 120$ km. Under daytime conditions the unknowns T_{∞} , T_0 and s are found by using T_i (i.e. T_n) for altitudes between 115 and 130 km, and T_i , T_e and N_e (together with the heat balance equation for the ion gas) over the altitude range from 225 to 425 km (SALAH and EVANS, 1972). The fitting procedure also allows a fourth parameter to be determined: the concentration of atomic oxygen (assumed to be in diffusive equilibrium) at a reference altitude of 400 km.

Table 3 summarizes results from the 30 min of incoherent scatter data collected during the AE-C passes on orbits 666 and 677. Since the atomic oxygen concentration and ion composition cannot be determined as accurately as the other variables from the incoherent scatter measurements, the values given in Table 3 for these two parameters represent averages over three 30 min runs.

RESULTS

(a) Electron density

The electron density profile measured by the Millstone Hill incoherent scatter radar at the time of the AE-C satellite crossing of Millstone Hill on orbit 666 is shown as the solid line without arrows in Fig. 3. The electron density peak occurs at a value of about 4.5×10^5 cm⁻³ at an altitude of 230 km and is characteristic of a typical daytime profile observed over Millstone Hill for conditions near solar minimum. The electron density profile calculated by the ionospheric model is shown as the dashed line in Fig. 3; its peak is about 5×10^{5} cm⁻³ but at an altitude near 250 km. The difference in altitude probably occurred because the polewarddirected neutral wind flowing in the thermosphere could have been stronger at the time of the measurement than the wind used in the model calculations. In the model the wind at the F2-peak is directed toward the pole with a speed of 25 m s^{-1} at 1130 LT. This is typical of daytime thermospheric winds; however, the results suggest that we would obtain better agreement by increasing the speed of the model thermospheric winds to 50 m s^{-1} . These differences between calculated and observed values are not significant. The AE-C satellite measurements of electron density along the satellite path are shown in Fig. 3 as the solid line with arrows. The electron density measured directly over Millstone Hill is shown as a cross at 161 km. There is

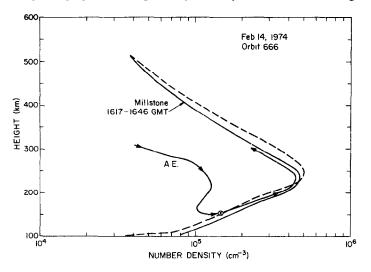


Fig. 3. Electron number density profile over Millstone Hill at the time of the AE-C crossing on orbit 666, 14 February 1974. The solid line is the density profile measured by the incoherent scatter radar, the dashed line gives the density profile calculated by the ionospheric model and the solid line with arrows shows the AE-C satellite measurements along the satellite path.

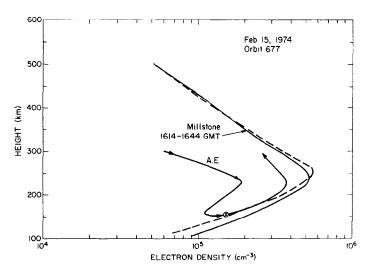


Fig. 4. Same as in Fig. 3, for orbit 677 on 15 February 1974.

considerable latitudinal and longitudinal variation evident in the measurements. However, there is general agreement among the satellite measurements, the incoherent scatter radar measurements and the model calculations of electron density over Millstone Hill at 161 km. Similar agreement among incoherent scatter radar measurements, model calculations and satellite measurements of electron density was obtained for AE-C satellite orbit 677 on 15 February 1974 (Fig. 4). On orbit 677 the satellite measurements again show considerable latitudinal and longitudinal variability along the satellite path in contrast to a vertical profile at a given location.

(b) Ion densities

The ion chemistry profiles of O^+ , O_2^+ and NO^+ for AE-C orbit 66 are shown in Fig. 5. The dashed lines show model calculations for the time of the satellite crossing over Millstone Hill and the solid lines with arrows show satellite measurements of ion density along the satellite path. The breaks in the solid line indicate data gaps along the orbital path. There is reasonable agreement between the

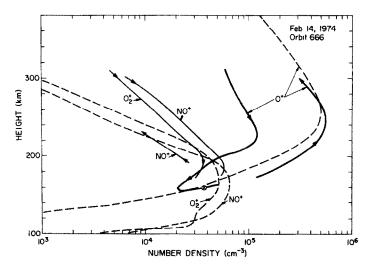


Fig. 5. Ion density profiles calculated over Millstone Hill at the time of the AE-C satellite crossing and ion densities measured along the satellite path. The dashed lines are the ionospheric model calculations and the solid lines with arrows are the satellite measurements along the satellite path. The breaks in the solid lines with arrows indicate regions where satellite data were missing.

calculated and observed O⁺ density at the satellite crossing at 161 km over Millstone Hill. However, there is also considerable latitudinal and longitudinal variation along the satellite path. At a given altitude densities at high latitudes are lower than they are at low latitudes. The O_2^+ and NO^+ ion densities were obtained during only a portion of the orbital pass and measurements were not obtained at the Millstone Hill crossing at 161 km. The apparent discrepancy between the measured and calculated NO⁺ and O_2^+ scale heights is due to latitudinal changes in composition and temperature. When applied point-by-point along the satellite path, the model in fact closely reproduces the satellite measurements. The incoherent scatter measurements give $N(O^+)/N_e = 0.15 \pm 0.07$ at 161 km, compared to values of 0.21 from the model calculation and 0.29 from AE-C measurements.

(c) Electron and ion temperatures

The calculated and observed profiles of electron and ion temperatures and the neutral gas temperature for the AE-C crossing of Millstone Hill on orbit 666 are shown in Fig. 6. The incoherent scatter radar measurements are indicated by crosses with error bars and the solid line is a smooth fit to the data. The dashed lines show model calculations for the same time. The electron and ion heat fluxes at 500 km, determined from the observed values, are used as boundary conditions for the electron

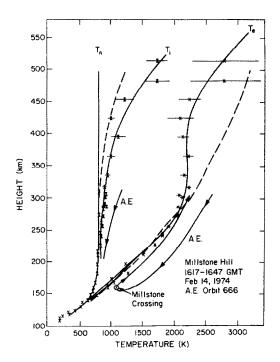


Fig. 6. Electron, ion and neutral gas temperature profiles over Millstone Hill for the time of the AE-C satellite crossing on orbit 666, 14 February 1974. The crosses are incoherent scatter radar measurements and the solid lines are smooth fits to the data. The dashed lines are ionospheric model calculations and the solid line with arrows gives the AE-C satellite measurements along the satellite path. The circled crosses represent the AE-C measurements at 161 km over Millstone Hill.

and ion energy equations. The collisional loss processes in the energy equations are the same as those given in REES and ROBLE (1975) and ROBLE and HASTINGS (1977) except that we used the atomic oxygen fine structure electron loss rate derived by HOEGY (1976) in the calculations. Below 300 km there is good agreement between the electron temperatures determined by the incoherent scatter radar and by the model calculations. The AE-C satellite measurements of the electron and ion temperatures were generally higher than both the incoherent scatter radar measurements and the model calculations. At the only place where overlapping data occur, over Millstone Hill at 161 km, the incoherent scatter radar measured an electron temperature of 960° K, the satellite measured a value of 1050 °K and the model calculations gave 998° K. The spread of 90° K in the measured and calculated electron temperatures is not considered significant since measurements and model calculations were made in an altitude region with large temperature gradients. The satellite measurements of ion temperatures, although not directly measured over Millstone Hill at 161 km, do bracket the incoherent scatter radar observations at higher altitudes with the measurements made by AE-C on its inbound and outbound passes.

Above 300 km, there is a significant difference between the calculated electron temperature and the electron temperature determined by the incoherent scatter radar. Such a difference between calculated and observed electron temperatures also was occasionally noted by BRACE et al. (1976). This effect could be caused by an inaccurate description of the H⁺ density in the topside ionosphere by the ionospheric model for 14 February 1974. If the H⁺ densities were greater than the chemical equilibrium densities calculated by the ionospheric model, the increased H⁺ density would act to decrease the electron temperature and increase the ion temperature through coulomb collision between H⁺ ions and electrons. This would give better agreement between the calculated and observed profiles. However, the shapes of the incoherent scatter correlation functions are markedly changed even if only a small amount of H^+ is present along with the dominant O⁺ ions. For the present measurements, there is no indication of any H^+ below 500 km, within the detection limit of about 2%. The electron and ion temperatures observed on the 30-min runs prior to and after the AE-C pass on orbit 666 show better agreement with model predictions, although some discrepancy remains.

For the next day, 15 February 1974, the calcu-

lated and observed electron temperature profiles above 300 km over Millstone Hill were in better agreement (Fig. 7). It is not clear why the agreement was better on 15 February, but perhaps the H^+ density was different. At these altitudes chemical equilibrium between H^+ and O^+ should prevail, which suggests a change in the neutral hydrogen distribution because the electron density and atomic oxygen density were similar for the two calculations.

The AE-C electron temperature measurements along the satellite path are shown in Fig. 7 as a solid line, with arrows to indicate the inbound and outbound portions of the satellite pass. The satellite measurements show nearly equal temperature gradients for both the inbound and outbound passes, in contrast to the somewhat different gradients observed on orbit 666. The differences were probably due to variations in the heat flow rate of the electron gas from the magnetosphere along the satellite path. There is a spread of 144 K among the electron temperatures observed by the incoherent scatter radar and the AE-C satellite and calculated by the ionospheric model at 162 km, the altitude of the AE-C satellite crossing over Millstone Hill (Table 3). The incoherent scatter radar measured a temperature of 990 K, the AE-C satellite measured 1110 K and the calculated temperature was 966 K. Again, the measurements were made in a region where large temperature gradients exist.

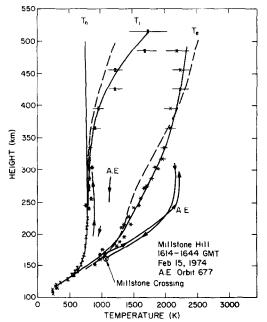


Fig. 7. Same as for Fig. 6, for orbit 677 on 15 February 1974.

(d) Neutral temperature

The neutral gas temperature profile measured by the AE-C satellite is in good agreement with the neutral gas profile determined by the incoherent scatter radar measurements. The satellite measured a temperature of 702 K over Millstone Hill on orbit 677 in contrast to a temperature of 703 K measured by the incoherent scatter radar as given in (Table 3).

(e) OI 6300 Å emission rate

The 6300 Å volume emission rate at the time of the AE-C satellite crossing over Millstone Hill was observed by the visible airglow experiment photometer (HAYS et al., 1973) and is shown in Fig. 8. The dots give the actual data points and the solid line is a smooth fit to the data. The data are an average of about 30 limb scans by the VAE photometer in the vicinity of Millstone Hill. We inverted the data to obtain the volume emission rate profile using the Abel Inversion technique described by HAYS et al. (1977). The 6300 Å volume emission rate was observed to peak at 220 km with a value of 250 photons $cm^{-3}s^{-1}$. The calculated total 6300 Å volume emission rate above Millstone Hill is shown as the dashed line in Fig. 8; and the various components contributing to the total 6300 Å volume emission rate are also shown. The

components are excitation due to dissociative recombination, Schumann-Runge photodissociation, photoelectron impact and thermal electron impact. The various reaction rates, quenching coefficients and excitation cross sections that we used in the analysis were discussed in detail by ROBLE et al. (1976) and HAYS et al. (1977). Above 350 km, thermal electron impact is the major source of 6300 Å excitation. However, its total contribution is only 23 R (Rayleighs). Between 200 and 350 km, photoelectron impact and chemical recombination are the major excitation sources; they contribute a total of 1193 R and 799 R, respectively. And, below 200 km, Schumann-Runge photodissociation is the major excitation source, with a total intensity of 1498 R. The total 6300 Å emission rate calculated at the time of the AE-C satellite crossing of Millstone Hill was 3721 R; the calculated vertical profile of the 6300 Å volume emission rate shows good agreement with the measured values.

(f) Odd nitrogen

The calculated steady state distribution of minor neutral constituents at the time of the AE-C satellite crossing of Millstone Hill is shown in Fig. 9. The chemical reactions and reaction rates used in the calculation are given in Table 1. The AE-C satellite measurements of the NO distribution along

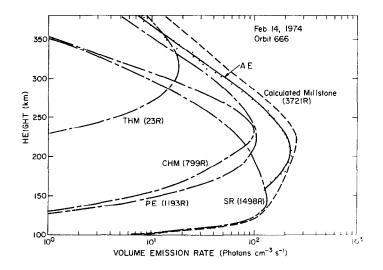


Fig. 8. Calculated and observed 6300 Å volume emission rate. The dots show the 6300 Å volume emission rates observed by the visible airglow experiment aboard the AE-C satellite over Millstone Hill. The solid line is a smooth fit to the data points. The dashed line shows the total 6300 Å volume emission rate calculated by the ionospheric model. The long and short dashed lines show various components of the 6300 Å volume emission rate; THM, excitation of atomic oxygen by thermal electron impact; CHM, dissociative recombination; PE, photoelectron impact; SR, Schumann-Runge photodissociation.

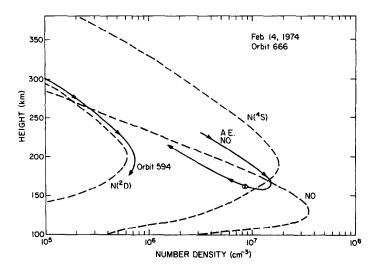


Fig. 9. Nitric oxide distribution measured by the AE-C satellite (solid line) and distribution of the minor neutral constituents $N(^4S)$, $N(^2D)$ and NO calculated by the ionospheric model at the time of the AE-C satellite crossing of Millstone Hill on orbit 666 on 14 February 1974. The cross indicates the measured NO density over Millstone Hill.

the satellite path are shown as a solid line with arrows to indicate the inbound and outbound satellite passes. The NO density measured directly over Millstone Hill is shown as a cross at 160 km and has a value of 9×10^6 cm⁻³. The model calculations gave 1.6×10^7 cm⁻³ for the NO number density at 160 km. The AE-C measurements of NO bracket the calculated values, indicating reasonable agreement between model calculations and measurements from the satellite pass. The measured NO densities are also consistent with the values of NO⁺ shown in Fig. 5, with high values on the inbound and low values on the outbound portion of the orbit. The discrepancies may be due in part to transport by winds and in part to deficiencies in the photochemical theory of minor neutral constituent chemistry. No measurements of NO or $N(^{2}D)$ density were made during AE-C orbit 677. However, an $N(^{2}D)$ density profile was determined from the 5200 Å airglow measurements on AE-C orbit 594 (FREDERICK and RUSCH, 1977) and it is shown in the figure for the sake of comparison. The calculated densities, although obtained for a different but comparable orbit, are somewhat greater than model calculations although the agreement is generally good.

(g) Photoelectron spectrum

Our last comparison was between the measured and calculated photoelectron spectra over Millstone Hill at the time of the AE-C satellite crossing. The calculated and observed photoelectron spectra for orbit 666 on 14 February and orbit 677 on 15 February are shown in Figs. 10 and 11, respectively. The solid line is the calculated total photoelectron flux, which consists of the sum of the

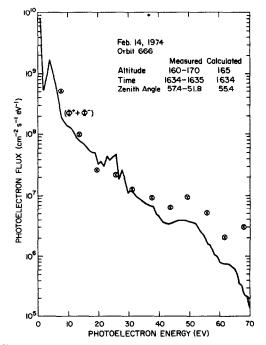


Fig. 10. Photoelectron flux spectra measured by the AE-C satellite during the Millstone Hill crossing on orbit 666, 14 February 1974 (crosses) and calculated by the ionospheric model (solid line).

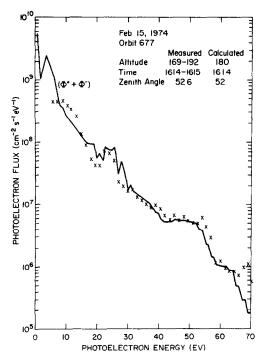


Fig. 11. Same as in Fig. 10, for AE-C satellite orbit 677 on 15 February 1974.

upward- and downward-directed photoelectron streams in the NAGY and BANKS (1970) program. These photoelectron fluxes were assumed to be isotropic for the purpose of comparison to AE-C measurements. On orbit 666 the photoelectron spectrum was calculated at 165 km and on orbit 677 the spectrum was calculated at 180 km. The measurements, shown as crosses in both figures, were obtained over the altitude range shown on the figure insert. On orbit 666 the photoelectron spectrum was obtained at a 6 eV resolution, whereas

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the photoelectron spectrum on orbit 677 was obtained at a higher resolution of 1.5 eV. In both cases the agreement between the calculated and observed values is generally good.

DISCUSSION

The results described above indicate that current F-region theory, as described by model computations, is adequate for describing the measured ionospheric properties in the mid-latitude daytime ionospheric region where photochemical processes dominate. The measurements and calculations apply to a low level of solar activity near solar minimum, with solar F10.7 values of 70–80× 10^{-22} W cm⁻² Hz⁻¹. When used in the ionospheric model, the measured values of the solar EUV flux, as discussed by HINTEREGGER (1975), give a calculated ionospheric structure that generally agrees with observations for the two days that we analyzed.

Our results also show that there is considerable variation in the ionospheric properties along the satellite path and that care should be taken in converting the satellite data into height profiles for comparison with ground-based measurements. Some of the variations are associated with latitudinal gradients of neutral composition and dynamics and some are associated with the satellite's passing through regions (for example, high latitudes) where processes other than solar photochemistry are operating.

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