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FRACTAL EXCITON FUSION: SIMULATIONS ON LONG-RANGE
PERCOLATION CLUSTERS*
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Computer simulations of exciton fusion on percolating clusters were performed for interaction distances $n=2,3$, and 4. The data yield slopes less than those expected from superuniversailty ( $f=2 / \dot{\circ}$ ) at interaction distances above 2 , when the probability of successively longer hops falls by a factor of $1 / 10$.

Exciton fusion kinetics in isotopic mixed crystals have been modeled by Monte Carlo random walks on random binary lattices, giving for both two- and three-dimensional systems of reacting walkers an effective spectral dimension $d_{s}=2 f$, where $f=2 / 3$ at the critical percolation concentration $\left(C_{C}\right), f=1$ at $100 \%$ concentration (classical behavior), and a "crossover" from fractal to classical regime in between. Fusion occurs when two walkers occupy the same site simultaneously. The integrated rate equation for fusion is: $\rho(t)^{-1}-p(0)^{-1}=K_{0} t^{f}$, where $\rho(t)$ is the walker density at time $t$, and $k_{0}$ is a constant. The simulations have here been extended to walks with jumps longer than nearest-neighbor ${ }^{2}$ (see Fig. 1). The probability assigned to a long range jump is a function of the number of ways an exciton may arrive there and the assigned probability for an interaction range. Static percolation gave the following: $C_{c}=0.282, n=2$;
$C_{C}=0.156, n=3$; and $C_{c}=0.099, n=4$.
The slopes of the curves in Fig. $2(n=2)$, Fig. $3(n=3)$, and Fig. $4(n=4)$ correspond to $f$. Simulations close to $C_{c}$ were run ---
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FIGURE 1
Four "diamonds" corresponding to $n=1,2,3$, and 4 , nested around a walker $X$. The walker may choose any diamond with assigned probability. The numbers are the number of routes to a site.


FIGURE 2
$n=2, \quad \log \left[1 / \rho-1 / \rho_{0}\right]$
vs. Log Time plotted. 5000 step random walks on largest cluster of 250,000 site square lattices. $10 \%$ initial concentration of walkers. 100 runs averaged. See text for explanation.
somewhat above $C_{C}$, to ensure percolation. Cluster Multiple 4
Labelling Technique was used to find the largest cluster. The three sets of three curves in Fig. 2 are for $n=2$ with $1 / 10$ ratio between diamonds, 1/2, and $1 / 1$, top to bottom. (Top and bottom sets are shifted 1 log cycle up and down for clarity.) Interaction range $n=2$ seems to conform to the asymptotic superuniversality hypothesis at criticality. At $29 \%$, there is a trend to higher $f$ when all sites are equally probable and a trend to lower $f$ when the ratio is $1 / 10$.

In Figs. $3 \& 4$ the ratio between diamonds was $1 / 10$. These data show a trend in $f$ toward lower values as the interaction range increases. This may be due to walkers in inaccessible regions not having taken sufficiently many long hops to meet other walkers and annihilate. Indeed, the $20 \%$ curve on the $n=4$ graph might be viewed as if it were an $n=2$ or $n=3$ curve near percolation--as if the longer hops were not contributing. Over the course of 100,000 steps, a walker can only perform 100
hops of four lattice spacings ( $n=4$ ). Therefore the walks should be extended to much longer times to assure that asymptotic $f$ values (for superuniversality) are obtained. Similar behavior was observed in single walker simulations.

We note that recent experiments for triplet excitons in naphthalene-perdeuteronaphthalene indicate that "classical" behavior becomes "fractal" behavior between $16 \%$ and $12 \%$ concentrations, corresponding roughly to $n=4$.


FIGURE 3
$\mathrm{n}=3 . \quad 100,000$ step random walks. Top: $30 \%, 8$ runs. Bottom: $16.5 \%, 8$ runs.


FIGURE 4

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