Ultra-high temperatures from oxygen isotope thermometry of a coesite-sanidine grospydite

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Abstract. The oxygen isotope compositions of coesite, sanidine, kyanite, clinopyroxene and garnet were measured in an ultra-high pressure-temperature grospydite from the Roberts Victor kimberlite, South Africa. The δ^{18} O values (per mil v. SMOW) of each phase and (1 σ) are as follows: coesite, 8.62 (0.31); sanidine, 8.31 (0.02); kyanite, 7.98 (0.08); pyroxene, 7.63 (0.11); garnet, 7.53 (0.03). In situ analyses of the coesite with the laser extraction system are $\delta^{18}O = 9.35$ (0.08), n = 4, demonstrating that the coesite is homogeneous. The coesite has partially inverted to polycrystalline quartz and the pyroxene is extensively altered during uplift. The larger scatter for the mineral separate coesite and pyroxene data may be due to partial reequilibration between the decompression-related breakdown products of these two phases. The anomalously high δ^{18} O value of the grospydite $(\delta^{18}O_{wholerock} = 7.7\%)$ is consistent with altered oceanic crust as a source rock. Temperature estimates from a linear regression of all the data to three different published calibrations correspond to an equilibrium temperature of 1310 ± 80 °C. The calculated isotopic pressure effect is to lower these estimates by about 40 °C at 40 kb. The estimated temperature based on Al-Si disorder in sanidine is 1200 + 100 °C and that from Fe-Mg exchange thermometry between garnet and clinopyroxene is 1100 ± 50 °C. Given the large errors associated with thermometry at such high temperatures, it is concluded that the xenolith equilibrated at 1200+100 °C. Pressure estimates are 45 ± 5 kb, based on dilution of the univariant equilibria albite = jadeite + coesite and 2 kyanite + 3 diopside = grossular + pyrope + 2 coesite. Zoning in the outer 20 μ m of the feldspar from Ab_{0.8} to Ab₁₆ indicates rapid decompression to 25 kb or less. The isotopic temperature estimates are the highest ever obtained and combined with the high degree of Al-Si disorder in sanidine require rapid cooling from ultra-high temperatures. It is inferred that the xenolith was sampled at the time of equilibration, providing a point on the upper Cretaceous geotherm in the mantle below South Africa.

Introduction

Oxygen isotope thermometry has been applied successfully to a wide range of rock types. The advantages of oxygen isotope thermometry compared to other geothermometric techniques is that it usually requires no pressure correction and can be used over a broad temperature range (e.g., O'Neil 1986). Unfortunately, oxygen isotope thermometry is not suitable for all high temperature rocks. Oxygen isotope fractionations between coexisting phases in many high temperature igneous rocks are relatively insensitive to temperature variations. In slowly cooled metamorphic rocks, significant retrograde oxygen isotope exchange occurs by the process of intracrystalline oxygen diffusion (Bottinga and Javoy 1975; Giletti 1986; Sharp et al. 1988).

The highest temperature estimates from oxygen isotope thermometry are from rapidly cooled igneous rocks (Onuma et al. 1970, Taylor and Epstein 1970; Anderson et al. 1971). Mantle xenoliths are ideal candidates for oxygen isotope thermometry because they meet the "rapidly-cooled" criterion for preserving peak temperatures. Unfortunately, the mafic, undersaturated compositions of nearly all mantle xenoliths are not suited to oxygen isotope thermometry, as the fractionations among common upper mantle phases (pyroxene, garnet, olivine) are very small (O'Neil 1986). In this study, the δ^{18} O values of mineral separates from an unusual coesite-bearing grospydite were measured in order to estimate temperatures of formation. This mantle nodule is well suited for oxygen isotope thermometry because the fractionations between the coexisting phases (coesite, sanidine, kyanite, garnet and pyroxene) are relatively large, and the rock is thought to have cooled extremely rapidly under anhydrous conditions (Smyth and Hatton 1977).

Pressure-temperature estimates from mantle xenoliths give the most direct indication of the sub-crustal geothermal gradient (e.g., Boyd 1973; Boyd and Nixon 1973; Finnerty and Boyd 1984, 1987). The δ^{18} O values of mantle samples have been measured (Garlick et al. 1971;

Kyser et al. 1981, 1982; McGregor and Manton 1986; Ongley et al. 1987; Shervais et al. 1988; Caporuscio 1990), but these data have not been used for thermometry except in conjunction with other cation-based thermometers (Kyser et al. 1981). The measured oxygen isotope values of the coesite-sanidine-bearing eclogite (Smyth and Hatton 1977) should provide a check on the cation-based thermometers; any evidence of isotopic disequilibrium may indicate partial reequilibration during relatively slow cooling.

Sample description

A detailed description of the coesite-bearing nodule analyzed in this study (sample SRV 1) has been made by Smyth and Hatton (1977) and is only briefly described here. The sample comes from the Roberts Victor kimberlite near Kimberley, South Africa, and was found by Smyth and Hatton as a 5 kg tabular nodule. The sample consists of an anhedral matrix of clinopyroxene, with subhedral garnet, coesite and sanidine, euhedral kyanite, and very minor pyrrhotite. Approximate abundances of each phase are as follows: pyroxene, 56%; garnet, 28%; kyanite, 9%; coesite, 6%; sanidine, <1%. The matrix pyroxene is extremely aluminous (18 wt.% Al₂O₃) with significant amounts of jadeite (49 mol%) and some substitution of Ca-Tschermaks (6 mol%) and Ca-Eskola (6 mol%) components (Table 1, Fig. 1a). The pyroxene is light green to white in hand sample and is partially altered due to decompression-related exsolution of Ca-Tschermaks component (Smyth and Hatton 1977) (Fig. 1a). The garnets have an average composition of Gr₄₇Py₂₈Al₂₅ (Table 1), are up to 5 mm in diameter and show almost no alteration. Colorless coesite grains are up to 3 mm in diameter, with ubiquitous alteration rims of polycrystalline quartz (Fig. 1b, c) and have >99.9 wt.% SiO₂. The nearly colorless sani-

Tab	le	1.	Microprol	be analyses	of se	lected pha	ases in	sample	SRV	/-	1
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	Garnet ^a	Omphacite ^b	Sanidine ^c		
			(Core)	(Rim)	
SiO ₂	39.84	56.13	64.06	64.36	
TiO ₂ ^d	0.17	0.07	0.00	0.00	
$Al_2 \bar{O}_3$	23.20	17.95	18.83	19.08	
FeO ^e	12.19	1.84	0.01	0.02	
MgO	7.64	5.89	0.00	0.00	
MnO ^d	0.29	0.00	0.00	0.00	
CaO	17.82	11.15	0.01	0.04	
K ₂ O	0.01	0.18	17.37	14.60	
Na ₂ O	0.07	7.33	0.10	2.00	
Sum	101.22	100.53	100.38	100.11	
Si	2.941	1.940	2.949	2.940	
Ti	0.009	0.002	0.000	0.000	
Al	2.018	0.731	1.021	1.028	
Fe	0.752	0.053	0.000	0.001	
Mg	0.841	0.304	0.000	0.000	
Mn	0.018	0.000	0.000	0.000	
Ca	1.409	0.413	0.000	0.002	
K	0.001	0.008	1.020	0.851	
Na	0.010	0.491	0.009	0.177	
vacancy	0	0.058	0	0	

^a Normalized to 8 cations; ^b normalized to 6 oxygens for vacancy calculation; ^c normalized to 5 cations; ^d data for this element from Smyth and Hatton (1977); ^c all iron assumed as ferrous



Fig. 1. a Photomicrograph of pyroxene showing the decompressionrelated breakdown of the high-pressure Ca-Tschermak and Ca-Eskola components. Scale bar = 300 μ m. b-c Coesite with ubiquitous inversion rim to quartz in plane-polarized (b) and cross-polarized light (c). Other minerals: pyroxene (groundmass), kyanite (*left*), garnet (*top-left*). Scale bar = 300 μ m



Fig. 2. Zoning profile in sanidine measured with the electron microprobe

dine is very highly disordered (Scambos et al. 1987); it is chemically zoned with a core composition $Or_{99}Ab_1$ and a rim composition in the outer 20 µm of $Or_{83}Ab_{17}$ (Table 1; Fig. 2). Kyanite grains are up to 3 mm across, light blue in color with no alteration, and are nearly pure Al_2SiO_5 . Previous pressure-temperature estimates for the nodule are 48–50 kb and 1050–1125 °C (Wohletz and Smyth 1984; McCormick and Smyth 1986).

Analytical methods

Oxygen isotope compositions were determined with the laser based extraction method of Sharp (1990) on mineral separates and by in situ measurements on coesite. For mineral separates, a small sample was crushed, ultrasonically cleaned in distilled water, and sieved to between 180-250 µm mesh. Preliminary separation of the samples was made with heavy liquids. The sample was treated with methylene iodide (d = 3.31 - 3.32) to separate garnet and kyanite from the coesite, sanidine and pyroxene. Sanidine was further concentrated by floating in bromoform (d=2.814-2.818). The coesiteand sanidine-concentrates were passed through a magnetic separator; coesite and sanidine were then handpicked from their respective concentrates and treated for 30 s in hydrofluoric acid to weakly etch the sanidine. This allowed for the coesite and sanidine to be distinguished optically. A final handpicking of each concentrate resulted in a purity for both the sanidine and coesite separates that approached 100%. Pure garnet and kyanite separates were obtained by simple handpicking of the coarse-grained heavy concentrate. The purity of all concentrates was checked under crossed polars in a petrographic microscope.

The in situ analyses were made on a 200 μ m thick, doubly polished plate (Fig. 3). Coesite grains were first identified with EDS spectroscopy on an SEM at the University of Lausanne before laser analysis. A detailed description of the laser extraction technique is given by Sharp (1990, 1992) and is only briefly described here. Mineral separate analyses were made on about 1 mg of material. Fifteen to eighteen samples were loaded simultaneously on a nickel sample holder. A minimum of three laboratory quartz standards (relative to NBS-28 = 9.5‰) were analyzed in every block of 15–18 samples; all phases were analyzed at least twice. A 20 W CO₂ laser was used for heating, with a BrF₅ oxidizing reagent at 0.1 bar. All mineral phases could be reacted to 100% completion with the laser extraction technique. The in situ analyses on coesite were made by a series of pulses of the focused laser at 6 W power; each pulse duration was 50–80 ms. Pulsing continued until a hole



Fig. 3. Backscatter electron image of a 200 μ m thick plate of the coesite-bearing grospydite. The hole in the coesite grain (dark grey) constitutes a single oxygen isotope analysis with the in situ laser method (Sharp 1990). Scale bar = 300 μ m. Light grey matrix, pyroxene; white grain, garnet; black grain, kyanite

was bored through the thin plate. The gas was converted to CO_2 via combustion of a heated graphite rod with a platinum catalyst and introduced directly into the mass spectrometer.

Major element analyses were made with a Cameca SX50 microprobe at the University of Lausanne. Analyses were made with a tightly focused beam rastered over a $25 \,\mu\text{m}^2$ area at 15 kV and 10 nA sample current. The analytical standards jadeite, diopside, and orthoclase were used for calibration. The PAP correction program supplied by Cameca was used for data reduction. Mineral normalizations were made with the aid of the program MINFILE (Afifi and Essene 1988).

Results

The δ^{18} O values of all mineral separates and in situ analyses are given in Tables 2 and 3. The standard deviations for the data from mineral separates of garnet, kya-

 Table 2. Oxygen isotope ratios for all mineral separates

Mineral	Run number	δ^{18} O (SMOW)	Sample size (mg)
Garnet	S058-7	7.58	1.46
	S058-8	7.49	0.90
	S074-5	7.52	1.52
	S097-5	7.52	1.75
	S097-6	7.56	1.83
	Avg. (1σ)	7.53 (0.03)	,
Pyroxene	S097-7	7.71	1.78/
5	S097-8	7.48	1.30
	S098-6	7.55	1.82
	S098-7	7.76	1.28
	Avg. (1σ)	7.63 (0.11)	
Kvanite	S058-9	7.89	1.38
2	S058-10	8.11	0.88
	S074-6	7.99	1.03
	S097-9	7.99	1.46
	S097-10	7.90	1.80
	Avg. (1σ)	7.98 (0.08)	
Coesite	S097-3	14.38ª	0.60
	S097-4	8.36	0.52
	S098-4	9.05	0.80
	S099-3	8.77	1.02
	S099-4	8.28	0.76
	Avg- (1σ)	8.62 (0.31)	
Sanidine	S098-5	8.33	1.39
	S099-5	8.28	1.03
	Avg. (1 σ)	8.31 (0.02)	
Whole rock		7.7 ^b	

^a Excluded from average; ^b calculated based on mineral modes of Smyth and Hatton (1977)

Table 3. Oxygen isotope values of in situ analyses of coesite. Each analysis is a single 200 μ m diameter hole

Run number	δ^{18} O (SMOW)	μ moles CO ₂
S096-2	9.38	0.9
S096-4	9.44	0.8
S096-5	9.34	2.0
S096-6	9.22	1.6
Avg (1σ)	9.35 (0.08)	

nite and sanidine are all very small, indicating a high degree of homogeneity. The pyroxene data have only a slightly larger range in δ^{18} O values, which may be related to the extensive alteration and/or retrograde isotopic exchange. The δ^{18} O values of the coesite have the largest spread. The δ^{18} O values range from 8.28-9.05‰, excluding an anomalous value of 14.38‰. This large range in δ^{18} O may be due to either (1) real sample heterogeneities, (2) partial reequilibration of the coesite during retrogression or (3) contamination by other phases. The degree of sample contamination for coesite can be assessed visually with the laser extraction system. The fluorination products of SiO₂ are all gases (Br₂, BrF_3 , SiF_4); any residual solid fluoride in the sample chamber indicates sample contamination. In all cases, very minor or no residues were observed. Assuming an unreasonably high sample contamination of 20% by the isotopically lightest phase in the rock, the shift in the δ^{18} O would still be <0.2‰. The variations in the δ^{18} O of coesite, therefore, cannot be explained exclusively by contamination.

Potential isotopic heterogeneity of the coesite was assessed by making in situ analyses of coesite. The high spatial resolution ($\sim 200 \,\mu m$) of the in situ technique (Fig. 3) allows for analyses to be made on pure, uncontaminated material. All in situ δ^{18} O values are 9.35 $\pm 0.08\%$ (Table 3), demonstrating that the coesite is isotopically homogeneous. The in situ δ^{18} O values are slightly higher than those obtained on mineral separates. This ¹⁸O enrichment may be due to an isotopic shift associated with the in situ method for coesite. There is no shift associated with in situ analyses of quartz plates (Sharp 1992; Kirschner et al. 1992), but other refractory phases, such as garnet, do exhibit an isotopic shift (Sharp 1992). In contrast to the in situ data, no isotopic shift has been observed on mineral separate data with the laser extraction system (Sharp 1990, 1992); therefore, the mineral separate data are accepted as the true isotopic composition. The in situ data are used to demonstrate that the coesite is isotopically homogeneous.

A possible explanation for the large range in measured coesite δ^{18} O values is retrograde exchange with the host pyroxene. The coesite has a ubiquitous inversion rim to polycrystalline quartz (Fig. 1b, c), and the pyroxene is partially altered, presumably due to the instability of the Ca-Eskola, Ca-Tschermaks and jadeite components during decompression (Fig. 1a). If oxygen isotope exchange between the inverting coesite and the decomposing pyroxene occurred during decompression, there could be a significant shift in the δ^{18} O values of the SiO₂ polymorphs. Even slight retrograde deformation can cause significant isotopic exchange (e.g., Sharp et al. 1988). If 20% of the coesite inverted to quartz and exchanged with the surrounding pyroxene at temperatures of about 700 °C, the δ^{18} O of the quartz would be 1.8‰ higher than the coesite, and the sum of the two would be about 0.4% higher than the unaltered coesite. An isotopic fractionation between the two silica polymorphs is predicted based on electrostatic site potential calculations for oxygen (Smyth 1989). At 1000 °C, the calculated fractionation between quartz and coesite

is 0.2‰, based on the data of Smyth (1989) and the equation of Smyth and Clayton (1988). This difference will be as large as 1‰ at 40 kb (see following discussion). Preferential analysis of either the polycrystalline quartz or the coesite will lead to a large scatter in the isotopic data for the SiO₂ polymorph.

There are several lines of evidence supporting the hypothesis for retrograde exchange. First, the pyroxene has a larger standard deviation than the unaltered phases garnet, sanidine and kyanite. Partial exchange with the inverting coesite would explain this small degreee of heterogeneity. Secondly, analyses of the δ^{18} O of the coesite in the same sample by Jeffrey Deen measured in the laboratory of R.O. Rye, U S G S, Denver with the conventional fluorination technique (Clayton and Mayeda 1963) are in the range 9.5 to 10.2‰. Very low yields were obtained due to the extreme difficulty of fluorinating coesite. The high δ^{18} O values obtained conventionally probably reflect a preferential reaction of the ¹⁸O-enriched polycrystalline quartz as compared to the more refractory coesite.

Oxygen isotope thermometry

Temperature estimates based on the isotopic fractionations between all phases are given in Table 4 and shown graphically in Fig. 4a–d. Four calibrations were used: (1) Bottinga and Javoy (1975); (2) the incremental method of Richter and Hoernes (1988) with the modified data set of Hoffbauer et al. (1992); (3) Smyth and Clayton (1988) using the oxygen site potentials of Smyth (1989); and (4) experimentally determined calibrations of Lichtenstein and Hoernes (1992) for quartz-garnet and Chiba et al. (1989) for all other phases. The fractionation factor for albite was used in place of sanidine for the temperature estimates with the fractionation factors of Bottinga and Javoy (1975) and Chiba et al. (1989). This substitution is justified as there is no oxygen isotope fractionation associated with K-Na exchange in feldspar (O'Neil and Taylor 1967; T Chacko, personal communication 1992). A correction for the jadeite component in clinopyroxene based on the temperature coefficient of fractionation a_{Di-Jd} relationship of Matthews et al. (1983) was made for the temperature estimates made with the fractionation factors of Bottinga and Javoy (1975) and Chiba et al. (1989). The fractionation factors with the technique of Hoffbauer et al. (1992) were kindly calculated by Dr. R. Hoffbauer for the exact mineral compositions given in Table 1. A fractionation factor intermediate to diopside and jadeite was used for the Smyth and Clayton calibration.

Oxygen isotope fractionations between any two phases are related by the equation:

$$\Delta_{i-j} = a_{ij} \times 10^6 / (T^2) + b_{ij} (T \text{ in K})$$
(1)

where Δ_{i-j} is the difference in the isotopic compositions of phases *i* and *j*, a_{ij} is the temperature coefficient of the fractionation equation for the mineral pair i-j and b_{ij} is the isotopic fractionation between mineral pair i-j

Table 4. Oxygen isotope thermometric estimates (°C). Abbreviations: Coe, coesite; Sd, sanidine; Ky, kyanite; Di, diopside; Om, omphacite; Gt, garnet

	Sd	Ку	Di	Om	Gt
Experime	ental calibra	ations ^{a, b, c}			
Coe Sd Ky Di Om	1470	1370 1280	1400 1360 1430	1240 1120 950	1340 1290 1290 620 2140
Bottinga	and Javoy	(1975) ^{a, b}			
Coe Sd Ky	1500		1370 1310 _	1210 1060 -	1350 1290 –
Di Om					1180 2390
Smyth an	nd Clayton	(1988) ^d			
Coe Sd Ky Di Om	1170	1550 1850	1280 1330 580	$\begin{array}{c} 1150\\ 1150\\ \infty\end{array}$	1330 1390 950 1780 2580
Hoffbaue	er et al. (199	2) ^{a, d}			
Coe Sd Ky Di Om	1520	1840 2100	$\begin{array}{c} 1370\\ 1300\\ \infty\end{array}$	$\begin{array}{c} 1230\\ 1080\\ \infty\end{array}$	1450 1430 650 2000 2940

^a The fractionation factor for quartz was used in place of coesite. ^b The fractionation factor for albite was used in place of sanidine; the fractionation factor for omphacite was assumed to be midway between that of diopside and jadeite based on the diopside-jadeite calibration of Matthews et al. (1983). ^c All fractionation factors from Chiba et al. (1989) except quartz-garnet from Lichtenstein and Hoernes (1992), and kyanite (empirical) from this study. ^d Fractionation factor for garnet corrected for solid-solution

at infinite temperature (Clayton and Epstein 1961; Bottinga and Javoy 1973). For the phases and calibrations considered in this study, b=0, and will be left out of the following equations. A temperature estimate for all coexisting minerals in a rock can be made with the isotherm method of Javoy et al. (1970) by rearranging Eq. 1 to a simple linear form

$$\Delta_{i-k} = a_{ik} \times 10^6 / (T^2) (T \text{ in } \mathbf{K}), \tag{2}$$

where all phases *i* are treated in terms of one of the phase *k* (the choice of *k* is arbitrary), and b=0. All data will plot on a Δ_{i-j} versus a_{i-j} diagram as a straight line with a slope of T^{-2} if all phases are in equilibrium. Equation 2 requires the assumptions that the phase *i* has a single value and the statistical best fit of Δ_{i-j} versus a_{i-j} must pass through zero (i.e., Δ_{k-k} and $a_{k-k}=0$). Equation 2 can be generalized to incorporate every measurement of each phase so that temperatures can be calculated with a linear best fit to *all* of the measurements by the following equation:

$$\delta_{ij} = a_{jk} \times 10^6 / (T^2) + C (T \text{ in } \text{K})$$
(3)

Table 5. Empirical estimation of the a(Qz-Ky) and a(Ky-Gt) temperature coefficient from published data on quartz – kyanite – garnet triplets

Reference	(1)	(2)	(2)	(3)	(4)	(4)	
δ^{18} O (Quartz)	0.00	8.70	8.40	14.00	9.41	7.83	
$\delta^{18}O$ (Kyanite)	-4.70	6.40	6.10	10.90	7.71	6.76	
$\delta^{18}O$ (Garnet)	_	5.10	5.10	9.50	5.17	4.79	
T(Qz-Gt)	430ª	615	650	520	545	690	
a(Qz-Ky)	2.32	1.81	1.97	1.95	1.13	1.00	
a(Ky-Gt)	0.37	1.02	0.86	0.88	1.70	1.83	
Reference	(5)	(5)	(5)	(5)	(5)	(5)	
δ^{18} O (Quartz)	23.05	10.60	10.35		18.55	14.00	
$\delta^{18}O$ (Kyanite)	18.50	8.00	4.70	7.15	15.50	10.15	
δ^{18} O (Garnet)	18.60	6.10	6.00	5.55	14.60	10.15	
T(Qz-Gt)	525	520	535	670 ^b	575	585	
a(Qz-Ky)	2.89	1.64	3.68	1.42	2.19	2.83	
a(Ky-Gt)	-0.06	1.19	-0.85	1.41	0.64	0.00	

^a Equilibration temperature determined from quartz-muscovite fractionation. ^b Equilibration temperature determined from garnetilmenite fractionation. References: (1) Taylor et al. 1963; (2) Garlick and Epstein 1965; (3) Hoernes and Hoffer 1979; (4) G. Frueh Green, ETH, Zurich, unpublished data; (5) Hoernes and Friedrichsen 1980

where δ_{ij} is the *i*th measurement of phase *j*, a_{jk} is the temperature constant of phase *j* relative to a reference mineral *k* (garnet was used in this study, although the choice of a reference mineral is arbitrary) and *C* is a constant. A least squares fit of all δ_{ij} measurements with respect to the corresponding a_{jk} constants yields a slope equal to T^{-2} with a *y* intercept equal to *C*. Equation 3 is more general than Eq. 2. Equation 2 forces the *y* intercept to be zero, while Eq. 3 eliminates this constraint. Graphical representations of Eq. 3 and temperature estimates based on the four different isotope calibrations are shown in Fig. 4a–d.

Temperature estimates and regression uncertainties using is method are as follows: Bottinga and Javoy (1975), 1320 (\pm 90) °C; experimental calibrations, 1310 (\pm 80) °C; Smyth and Clayton (1988), 1310 (\pm 90) °C; Hoffbauer et al. (1992), 1400 (\pm 50) °C. The kyanite data were excluded from all regressions.

Fractionation factors involving kyanite have not been estimated by Bottinga and Javoy (1975) or measured experimentally, but Smyth (1989) and Hoffbauer et al. (1992) provide kyanite fractionation data. In these two calibrations, the data for kyanite lie well off the isotherm defined by the other phases (Fig. 4c, d). The temperature coefficient of fractionation *a* for kyanite (relative to either quartz or garnet) was determined empirically in this study using published δ^{18} O values of quartz-kyanitegarnet assemblages (Table 5). Equilibration temperatures for these quartz-kyanite-garnet triplets were determined from the measured Δ_{gt-qz} and the calibration of Lichtenstein and Hoernes (1992). The garnet-kyanite or quartz-kyanite temperature coefficient of fractionation *a* could then be calculated with Eq. 1.

The average a_{ky-gt} in Table 5 is 1.1 ± 0.5 (excluding negative values). The large standard deviation of this



Fig. 4a–d. Graphical presentation of temperature estimates by Eq. 3. The isotopic composition (δ^{18} O value_{SMOW}) of every measurement for each mineral is plotted against the temperature coefficient *a* (from Eq. 1) for that mineral relative to garnet. The slope

value is almost certainly due to analytical errors, because kyanite is so difficult to fluorinate with conventional extraction techniques (in contrast, kyanite is easily fluorinated with the high-temperature laser extraction method). The temperature estimate for the grospydite SRV 1 using the value 1.1 for temperature coefficient a_{ky-gt} is 1290 °C, in agreement with other fractionations, although the uncertainties in the estimate for the temperature coefficient *a* may make this agreement fortuitous.

The polymorphic transition guartz-coesite may impose an additional complication to the oxygen isotope temperature estimates of the coesite-bearing grospydite. In the four calibrations applied in this study, only that of Smyth and Clayton (1988) with the electrostatic site potential data for coesite of Smyth (1989) provides fractionation factors in place of coesite. In all the others, quartz was used in place of coesite. The difference in the average site potentials for quartz and coesite is 0.45 V, corresponding to an a_{qz-coe} of 0.35 and a Δ_{qz-coe} of 0.14 at 1300 °C. Although this difference is small, an $a_{\text{coe}-\text{garn}}$ value of 2.48 is obtained when the $a_{\text{qz-coe}}$ value of 0.35 is subtracted from the experimentally determined $a_{qz-garn}$ value of 2.83. Application of this $a_{coe-garn}$ coefficient will result in temperature estimates that are 100 °C lower than those based on the experimental quartz-garnet fractionation. Preliminary measurements of Δ_{qz-coe}

of the best fit line to all the data is related to the temperature by Eq. 3. **a** From experimental calibrations (Chiba et al. 1989; Lichtenstein and Hoernes 1992); **b** Bottinga and Javoy (1975); **c** Smyth and Clayton (1988); **d** Hoffbauer et al. (1992)

by Sharp and P. Ulmer (ETH, Zurich) following the procedure of Clayton et al. (1989) do not support the theoretical fractionation indicated by the electrostatic data. In addition, no quartz-coesite fractionation is predicted by the increment method (Hoffbauer et al. 1992). It is therefore reasonable to use the quartz fractionation factors for coesite.

Temperature estimates from Al-Si disorder in feldspar

The very high temperature estimates from the stable isotope data are supported by the extraordinarily high degree of Al-Si disorder in the sanidine (Scambos et al. 1987). The sanidine exceeds the upper limits of Al-Si disorder that have been measured on one atmosphere, experimentally disordered samples. The degree of Al-Si disorder in feldspars increases with temperature, so that an estimate of disorder in a crystal defines a minimum temperature of equilibration. The unit cell parameters of Scambos et al. (1987) are a=8.595(3) Å, b=13.028(5) Å and c=7.175(2) Å. A measure of the Al-Si disorder can be calculated from the relation (Hovis 1986)

$$Z = -138.575 + 19.3153 c_k \tag{4}$$

where $c_k = c_{obs} + 0.0381 (1 - N_{Or})$ and the degree of disorder Z is defined as $2[N_{A1(T1)} - N_{A1(T2)}]$. The disorder term Z is 0.02 by this technique. Alternatively, Z may be estimated from the T-O distance technique of Kroll and Ribbe (1983) with the structure refinement of Scambos et al. (1987); the Z from this method is 0.06. Temperature estimates, corrected for pressure, are made from the following relation of Hovis (1974):

$$T(\mathbf{K}) = \frac{\Delta E^{0} + P \Delta V^{0}}{\Delta S^{0} - 2R \operatorname{ath} \left[\frac{4Z}{3 + Z^{2}} \right]}$$
(5)

where

ath
$$(x) = 0.5 \ln [(1+x)/(1-x)].$$
 (6)

In Eq. 5, $\Delta E^0 = -22200 \text{ J/mol}$, $\Delta S^0 = -15.5 \text{ J/deg·mol}$, $\Delta V^0 = -0.054 \text{ J/bar·mol}$; *P* in bar. Temperature estimates from Eq. 5 for Z = 0.02 and Z = 0.06 are 1220 and 1140 °C at 30 kb and 1310 and 1220 °C at 55 kb, respectively.

Geothermobarometric estimates from mineral equilibria

Various univariant equilibria involving the phases garnet, clinopyroxene, feldspar, coesite and kyanite are applicable to the sample SRV-1. The equilibria considered in this study are the following:

$$NaAlSi_{3}O_{8} = NaAlSi_{2}O_{6} + SiO_{2}$$

albite jadeite coesite (7)

$$\begin{array}{c} Fe_{3}Al_{2}Si_{3}O_{12} + 3CaMgSi_{2}O_{6} = Mg_{3}Al_{2}Si_{3}O_{12} + 3CaFeSi_{2}O_{6} \\ almandine & diopside & pyrope & hedenbergite \end{array}$$
(8)

 $2 \text{Al}_2 \text{SiO}_5 + 3 \text{CaMgSi}_2 \text{O}_6$

$$= 2 \operatorname{SiO}_2 + \operatorname{Mg}_3 \operatorname{Al}_2 \operatorname{Si}_3 \operatorname{O}_{12} + \operatorname{Ca}_3 \operatorname{Al}_2 \operatorname{Si}_3 \operatorname{O}_{12}$$

coesite pyrope grossular (9)

$$2Al_{2}SiO_{5} + 3CaMgSi_{2}O_{6} + 2NaAlSi_{2}O_{6}$$

kyanite diopside jadeite
$$= 2NaAlSi_{3}O_{8} + Mg_{3}Al_{2}Si_{3}O_{12} + Ca_{3}Al_{2}Si_{3}O_{12}$$

albite pyrope grossular (10)

In order to use these equilibria for accurate thermobarometric estimates, it is necessary to adjust the equilibrium curves of each reaction for the solid solutions present in the phases clinopyroxene, garnet and feldspar. Activity models for each of these phases are essential for accurate pressure-temperature estimates.

Activity models

The activity of NaAlSi₃O₈ in sanidine was determined with the ternary feldspar model of Fuhrman and Lindsley (1988) with the updated excess volume terms (W_v) of Lindsley and Nekvasil (1989). The pressure effect on the activity-composition (a-X) relations is very large in the pressure range experienced by the grospydites (Fig. 5). The NaAlSi₃O₈ activity at X(NaAlSi₃O₈)=0.01



Fig. 5. Activity-composition relation for activity of $NaAlSi_3O_8$ in sanidine for various pressures with parameters of Lindsley and Nekvasil (1989). In the dilute range considered in this study (*inner box*) pressure has a strong effect on the activity-composition relations



Fig. 6. Activity-composition relation for NaAlSi₃O₈ in sanidine at the critical solvus conditions defined by Thompson and Waldbaum (1969). The *horizontal line* in each curve satisfies the thermodynamic conditions that a solvus exists at that given pressure and temperature. P-T as follows (P in kb; T in °C): a - 0.001, 650; b - 10, 780; c - 20, 910; d - 30, 1020; e - 40, 1170

increases from 0.01 at 1 bar to 0.09 at 50 kb (1200 °C). The validity of this large pressure correction was evaluated by plotting the (a-X) relation for alkali feldspar at the critical point of the solvus from 0.001-40 kb (Thompson and Waldbaum 1969) (Fig. 6). A horizontal slope in the a-X relations at any pressure and temperature indicates that the solvus is reached. At all P-Tvalues plotted in Fig. 6, a horizontal slope occurs, in agreement with the solvus data of Thompson and Waldbaum (1969). This substantiates the large pressure effect indicated by the W_v of Lindsley and Nekvasil (1989), although extrapolations to 40 kb should theoretically include a W term for compressibility (Lindsley, personal communication 1992).

Garnet activities were estimated from the quaternary mixing model of Berman (1990). The calculated activities of Mg₃Al₂Si₃O₁₂, Fe₃Al₂Si₃O₁₂ and Ca₃Al₂Si₃O₁₂ in garnet based on earlier a - X models (e.g., Ganguly and Saxena 1984; Anovitz and Essene 1987) are more ideal and yield slightly lower pressures and temperatures based on the calculated equilibria of reactions 8–10. Substitution of a Wohl ternary interaction parameter C₁₂₃ = -66.4 kJ/mol (Ganguly and Saxena 1984) in place of the C₁₂₃=0 assumed by Berman (1990) will shift the location of reactions 8–10 less than 20 °C and 2 kb.

The clinopyroxene composition in sample SRV-1 (Table 1) can be approximated by the major components jadeite (NaAlSi₂O₆) and diopside (CaMgSi₂O₆), with minor hedenbergite (CaFeSi₂O₆), Ca-Eskola molecule $(Ca_{0.5}\square_{0.5}AlSi_2O_6)$, and Ca-Tschermaks component (CaAl₂SiO₆). The quaternary system jadeite-diopside-hedenbergite-Ca-Eskola adequately describes the pyroxene, but a - X relations involving Ca-Eskola are limited to the Ca-Eskola – Ca-Tschermaks join (Gasparik 1984), so that quaternary mixing terms in the above system could not be calculated. The activities of NaAlSi₂O₆, $CaMgSi_2O_6$, and $CaFeSi_2O_6$ in clinopyroxene (cpx) were calculated using the jadeite-diopside-hedenbergite ternary of Anovitz (1991) based on the jadeite-diopside binary relations of Cohen (1986) and assuming ideal mixing between diopside and hedenbergite. The activity coefficient (γ) for each end-member was calculated at the pressure and temperature of interest by normalizing $X_{\text{jadeite}} + X_{\text{diopside}} + X_{\text{hedenbergite}} = 1$. The activities of each component were then estimated as follows: $a(NaAlSi_2O_6)$ in $\operatorname{cpx} = (X_{\operatorname{Na}}) \cdot (X_{\operatorname{Al}}) \cdot (X_{\operatorname{Si}})^2 \cdot (\gamma_{\operatorname{Jd}}), \quad a(\operatorname{CaMgSi}_2O_6 \quad \text{in}$ $cpx) = (X_{Ca}) \cdot (X_{Mg}) \cdot (X_{Si})^2 \cdot (\gamma_{Di})$ and $a(CaFeSi_2O_6 \text{ in } cpx) = (X_{Ca}) \cdot (X_{Fe}) \cdot (X_{Si})^2 \cdot (\gamma_{Di})$ where all cations (including vacancies) are normalized to 4.

Pressure-temperature estimates

The *P*-*T* locations of reactions 7–10 were calculated with the computer program THERMOCALC (Powell and Holland 1988) for the thermodynamic data set of Holland and Powell (1990) corrected for reduced activities outlined already. Pressure estimates from the intersection of reactions 7, 9 and 10 are 44 kb at 1110 °C (Fig. 7). The feldspar is zoned in the outer 50 microns due to new growth of more sodic feldspar by the net-transfer reactions 7 and 10 (Fig. 2) and is consistent with pressures of 25–30 kb (Fig. 7).

The garnet-clinopyroxene exchange reaction 8 is the only reaction considered that is strongly temperature sensitive. The calculated equilibria for reaction 8 corrected for solid solution occurs at 1000 °C at 43 kb, far lower than estimates from either the isotopic data or Al-Si disorder data from sanidine. The calculated P-T position of this reaction is very sensitive to small changes in the free energy of the end-member phases. Small uncertainties in the thermodynamic data or a-X rela-



Fig. 7. Pressure-temperature diagram of selected equilibria. The equilibria coesite = quartz from Bohlen and Boettcher (1982); graphite = diamond from Kennedy and Kennedy (1976). Melting relations from Green and Ringwood (1967), Ito and Kennedy (1971) and Lambert and Wyllie (1972). The Fe-Mg garnet-clinopyroxene exchange reaction (*E&G*) was calculated with the formula from Ellis and Green (1979). All other equilibria calculated with the computer program THERMOCALC (Powell and Holland 1988) and the thermodynamic data set of Holland and Powell (1990) corrected for solid solution. Abbreviations: $ab_{0.8}$, albite ($X_{albite} = 0.008$); ab_{16} , albite ($X_{albite} = 0.16$); alm, almandine; coe, coesite; di, diopside; dm, diamond; gph, graphite; gr, grossular; hed, hedenbergite; jd, jadeite; ky, kyanite; py, pyrope

tions for any of the phases in reaction 8 can dramatically shift the calculated P - T location of this reaction. More reliable temperature estimates can be obtained with experimental calibrations of the Fe-Mg garnet-clinopyroxene exchange thermometer (e.g., Råheim and Green 1974; Ellis and Green 1979; Pattison and Newton 1989; see references in Fonarev et al. 1991). Temperature estimates with the Ellis and Green calibration are 1120 °C, in excellent agreement with the invariant point of reactions 7, 9 and 10, but lower than the estimates from either oxygen isotope thermometry or Al-Si disorder in sanidine. Temperature estimates based on Powell (1985) and Krogh (1988) are in agreement with those from Ellis and Green (1979), while the calibration of Pattison and Newton (1989) yields far lower temperatures. Pattison and Newton (1989) state that their calibration "does not address the problem of K_D versus T in eclogites with omphacite of jadeite contents approaching 50%". In a comparison of the different thermometers, Fonarev et al. (1991) suggesting using the Ellis and Green calibration for samples above 1200 °C. Overall, the pressure, temperature and composition range of the Ellis and Green experiments more closely match the xenolith composition and metamorphic conditions than those of Pattison and Newton.

Discussion

Pressure estimates from the reaction 7, 9 and 10 are greater than 40 kb. The effect of pressure on oxygen isotope thermometry is thought to be insignificant for crustal pressures (Hoering 1961; Clayton et al. 1974) and is routinely ignored. There have not been, however, any previous oxygen isotope thermometric studies of samples equilibrated at such high pressures. It is therefore worthwhile to estimate any potential pressure contribution to the measured oxygen isotope fractionations in sample SRV-1.

Effect of pressure on oxygen isotope fractionations

Clayton et al. (1974) and Clayton (1981) estimated that the $\Delta V_{\text{reaction}}$ for a typical solid-solid reaction is on the order of 25×10^{-3} J/mole·kb. A rough estimate of the effect of pressure on a typical oxygen isotope exchange reaction can be estimated from the relationship

$$\frac{d(1000 \ln K)}{dP} = 1000 \frac{P\Delta V}{RT}$$
(11)

At temperatures of 1100 °C and 40kb, the isotopic shift due to pressure will be less than 0.1‰, assuming a $\Delta V_{\text{reaction}}$ of 25×10^{-3} J/mole kb. Such a pressure effect is nearly insignificant when compared to the analytical uncertainties associated with isotopic measurements, and no method of calculating the relative pressure corrections for each mineral is available by this method.

A second method for estimating the pressure effects on oxygen isotope fractionation between minerals can be made by comparing the electrostatic site potentials of oxygen in minerals at high pressure. High pressure structures have been determined at room temperature for quartz, coesite, diopside, and alusite and pyrope (references in Table 6). Unfortunately, high pressure structure data were not found for kyanite, albite or sanidine. We used these data to compute electrostatic potentials for the average oxygen sites at the pressures of the structure refinements and interpolated the values to 40 kb. These results are given in Table 6 along with the 1 atm values. The small increase in the average oxygen site potentials with pressure for most of the phases will only result in a relative change of less than 0.1‰ at 40 kb. The effect of pressure is such that at 50 kb, the calculated temperatures with the 1 bar temperature coefficients will be about 40 °C too high.

Table 6. Electrostatic site potentials, (V) for average oxygen sites at 1 atm and 40 kb

Mineral	1 atm	40 kb	Ref ^a
Quartz	30.816	30.710	1
Coesite	30.350	30.408	2
Pyrope	27.070	27.281	3
Gr ₄ , Py ₂₈ Alm ₂₃	26.852	n.d.	
Diopside	27.328	27.444	4
Andalusite	28.143	28.268	5
Kvanite	27.67	n.d.	
Sanidine	29.550	n.d.	

^a Ref. (all 1 atm data from Smyth (1989). High pressure data from (1) Levien et al. 1980; (2) Levien and Prewitt 1981a; (3) Hazen and Finger 1978; (4) Levien and Prewitt 1981b; (5) Ralph et al. 1984; n.d. not determined

For quartz, the oxygen site potential *decreases* with increasing pressure. This is consistent with a slight increase in the Si-O distance with pressure resulting from a decrease in the Si-O-Si angle and resulting decrease in the Si-Si next-nearest neighbor distance.

Pressure effects on isotopic substitutions have been considered in more detail by Polyakov and Kharlashina (1989) and Heybey and Wand (1989). A method to calculate the volume change for an isotopic substitution ($\Delta V'$) has been proposed (Heybey and Wand 1989) such that

$$\Delta V' = -R(\Gamma)/T \text{ (cm}^3/\text{mol)}$$
(12)

where

$$\Gamma = 2\gamma A o/K \quad (K^2/bar) \tag{13}$$

The effect of pressure on the isotopic fractionation between phase i and j can be denoted as

$$d(1000 \ln \alpha_{i-i})/dP = (\Gamma_i - \Gamma_i)(P - 1)/T^2$$
(14)

where $\gamma = \text{Grüneisen parameter}$, $\kappa = \text{bulk modulus (bars)}$, T = temperature (K), R = gas constant (83.143 cm³ · bar/mol·K) and Ao is related to the reduced partition function by the relation

$$1000 \ln Q = A \, o/T^2. \tag{15}$$

Table 7. Physical constants of mineral phases required for the calculation of the pressure effect on isotopic fractionation between phases. Data applicable to Eqs. 14 and 18 in text

Phase	k (mbar)	Ref.	γ (calc)	$Ao \times 10^{-6}$ (K ²)	Ref.	Cp^1 (J/mol·K)	V ⁰ (J/bar)	$\alpha \times 10^{6}$ (°C)	Ref.	Γ K ² ·mol/bar	qz-min (at	coe-min 1200 °C, 4	py-min 0 kb)
Ouartz	0.37	1	0.64	11.7	10	44.59	2.27	34	16	40.5			
Coesite	1.14	2	0.41	11.7	11	45.40	2.06	8	16	8.5	0.59		
Pyrope	1.75	3	1.16	8.9	12	325.50	11.32	19	16	11.8	0.53	0.06	
Diopside	1.13	4	1.08	9.0	10	166.50	6.62	24	16	17.3	0.43	-0.16	-0.10
Sanidine	0.67	5	0.50	10.8	10	204.50	10.9	14	16	16.1	0.45	-0.14	-0.08
Kvanite	0.43	6	0.17	10.0	13	121.70	4.41	11	16	8.0	0.60	0.01	0.07
Jadeite	1.43	7	1.08	8.0	14	160.00	6.03	20	16	12.1	0.52	-0.06	-0.01
Calcite	0.73	8	0.29	11.4	10	83.47	3.69	9	17	9.1	0.58	-0.01	0.05
Rutile	2.16	9	1.84	6.9	15	55.10	1.88	25	17	11.8	0.53	0.06	-0.00

Ref. (1) Bass et al. 1981; (2) Weidner and Carleton 1977; (3) Leitner et al. 1980; (4) Levien et al. 1979; (5) Angel et al. 1988; (6) Brace et al. 1969; (7) Kandelin and Weidner 1988; (8) Birch 1966; (9) Hazen and Finger 1981; (10) Clayton and Kieffer 1991; (11) assumed equal to quartz; (12) Lichtenstein and Hoernes 1992; (13) this study; (14) diopside value corrected to jadeite based on relation of Matthews et al. 1983; (15) Agrinier 1991; (16) Skinner 1966; (17) calculated from Berman 1987



Fig. 8. Combined pressure-temperature estimates from mineral equilibria, Al-Si disorder in sanidine and oxygen isotope thermometry. The errors associated with the oxygen isotope and Al-Si disorder sanidine temperature estimates are shown by the *shaded areas* for each. The pressure-correction for the oxygen isotope data yield a slight negative slope. Abbreviations for mineral equilibria as in Fig. 7, Z = disorder term in sanidine

Equation 14 can be modified using the identity:

 $\gamma = \kappa_s \beta V^0 / C p \tag{16}$

to yield

 $\Gamma = 2\beta V^0 A o/C p \tag{17}$

and finally

 $d(1000 \ln \alpha_{i-j})/dP = 2(\beta_i V_i^0 A o_i/C p_i - \beta_j V_i^0 A o_j/C p_j)(P-1)/T^2. (18)$

where κ_s is the adiabatic zero pressure bulk modulus, β is the thermal expansivity, V^0 is the molar volume and Cp is the molar heat capacity at constant pressure. A compilation of the necessary constants is given in Table 7. The fractionation between quartz and other minerals due to the pressure increase from 0 to 40 kb at 1200 °C is 0.4–0.7‰. This is a significant effect that would drastically alter the calculated oxygen isotope temperature estimates. The stable polymorph of SiO₂ in sample SRV-1 is not quartz, however, but coesite. The pressure-induced fractionation between coesite and all other phases is less than 0.2‰. Such an effect is small but sufficient to reduce the calculated temperatures by about 50 °C at 50 kb (Fig. 8). This pressure effect is similar to the 40 °C temperature reduction at 50 kb based on the electrostatic site potential data. The pressure effect estimates are certainly not without large errors, but they may offer an explanation in part, for the discrepancy between the cation-based and oxygen isotope thermometers. For high pressure rocks in the quartz stability field, the pressure effect could be significant and should be considered.

Pressure-temperature estimates for the coesite-bearing grospydite

The complete set of pressure-temperature estimates from all methods are presented in Fig. 8. The temperature estimates based on the oxygen isotope fractionations between coesite, sanidine, clinopyroxene and garnet are 1310 ± 80 °C (1 bar). Temperature estimates are 50 °C lower if a pressure correction is applied to the data. These later temperatures are in agreement with those predicted from Al-Si disorder in feldspar. Temperature estimates from Fe-Mg exchange between garnet and clinopyroxene are 1100 °C. The discrepancy between the isotope and cation exchange thermometry requires that either there is an error in one of the estimates or that the different thermometers equilibrated at different temperatures.

The three thermometric calibrations (a, b, c) for oxygen isotope thermometry have been derived by a combination of empirical, theoretical and experimental techniques. All give the same results, and there is no reason to suppose that the calibrations are in error. The temperature estimates based on Eq. 3 have an uncertainty of \pm 80 to 90 °C. At very high temperatures, total isotopic variations are small, so that the real uncertainty of the isotopic measurements may be close to double this value. It is possible that the isotopic data are, in fact, compatible with the cation exchange thermometry. The coesite data have the largest scatter, and they strongly influence the linear regression used for temperature estimates. If the predicted temperature coefficient of fractionation $a_{az-coe} = 0.35$ from the data of Smyth (1989) is correct, then the temperature estimates will be 1200 °C instead of 1300 °C. Furthermore, if the highest $\delta^{18}O_{\text{coesite}}$ =9.05% is actually the correct value, then the linear regression of all data by Eq. 3 will be 1100 °C, in agreement with the cation exchange thermometry.

The cation exchange thermometry estimates also have large uncertainties. The experimental calibration of Ellis and Green (1979) is valid only for Ca-poor garnets. For calcium-rich garnets, the garnet-clinopyroxene exchange thermometer is over extrapolated and may have a large error (Essene 1989). A further possibility for the discrepancy between the garnet-clinopyroxene exchange thermometer and the isotope thermometers is that there has been partial reequilibration of the former during retrogression. In this case, while the calibrations for both the isotope and cation-exchange thermometers may be correct, the garnet-clinopyroxene Fe-Mg exchange thermometer is recording cooling, rather than peak, metamorphic temperatures. A best overall estimate for equilibration temperatures is considered to be 1200 + 100 °C. This temperature is intermediate to the oxygen isotope estimates and garnet-clinopyroxene estimates, and equal to the estimate based on the Al-Si disorder in sanidine.

Pressure-temperature-time paths

The peak temperatures and pressures of formation for sample SRV-1 are 1200 ± 100 °C and 45 ± 5 kb. The retrograde zoning profile in the sanidine (Fig. 2) indicates formation of a more Na-rich feldspar from jadeite and coesite (or quartz) by reactions 7 and 10 at a pressure of 25–30 kb. The coesite-quartz inversion also occurs in this pressure range. This decompression must have been 368

close to isothermal. Otherwise, a Fe-Mg zoning profile would have developed between garnet and clinopyroxene and the sanidine would have become more ordered. Experimental determinations of the anhydrous solidus of quartz eclogite (Green and Ringwood 1967; Ito and Kennedy 1971) are well above 1300 °C at 40 kb and about 1250 °C at 25 kb (Fig. 7). Melting experiments of quartz eclogite with excess H₂O place the liquidus at approximately 1100 °C at 40 kb (Lambert and Wyllie 1972; Stern et al. 1975), well below the temperatures recorded by sample SRV-1 (Fig. 7). Rapid decompression of the coesite – sanidine eclogite does not require partial melting as long as the sample was free of H₂O.

The preservation of disordered sanidine in the coesitesanidine grospydite is difficult to explain unless the xenolith was sampled while at or near equilibrium. A similar conclusion has also been reached for some deep crustal xenoliths (Hayob et al. 1989), although the inference is controversial (Harley 1990; Harte and Barnicoat 1990 vs. Hayob et al. 1990). In the case of the Roberts Victor kimberlites, the inference yields a point on the upper Cretaceous geotherm in Southern Africa. The P-T conditions from sample SRV-1 lie on the geotherm of Carswell and Gibb (1987) for Lesotho garnet peridotites, but are at somewhat higher temperatures than that obtained with a typical continental shield geotherm calculated for a 40 mW/m² surface heat flow (Pollack and Chapman 1977). It may be dangerous to generally rely on xenoliths transported in kimberlites for average geothermal conditions because the portions of the mantle that produce kimberlites may be anomalously heated by a plume (Irving 1976). Nonetheless, we propose that the 45 ± 5 kb and 1200 + 100 °C recorded in the xenolith represent conditions extant in the upper mantle shortly before the grospydite became entrained in the kimberlite eruptive.

Source rock and tectonic implications

The process of grospydite and eclogite formation is enigmatic. The unusual chemistry of grospydites, in particular, and eclogites, in general, has been explained in terms of high pressure fractional crystallization from primary mantle melts or melts of subducted material (e.g., Lappin 1978; Smyth et al. 1989) or as prograde metamorphic products of subducted oceanic crust (e.g., Jagoutz et al. 1984; MacGregor and Manton 1986; Shervais et al. 1988).

Oxygen isotope compositions of mantle xenoliths provide some of the most indisputable evidence for subducted material being the ultimate source, but do not shed light on whether the formation process was high pressure accumulation of hyper-aluminous clinopyroxene (Smyth et al. 1989) or low pressure accumulation of plagioclase (MacGregor and Manton 1986; Shervais et al. 1988). The δ^{18} O value of uncontaminated mantle has a maximum range of 5 to 7‰ (Kyser et al. 1982; Kyser 1986). Samples lying outside this range are best explained by recycling of oceanic crust. Previous measurements of the δ^{18} O values of grospydites are in the range of pristine mantle (Garlick et al. 1971; MacGregor and Manton 1986 – sample no. 53). The calculated whole rock δ^{18} O value of 7.7‰ for this coesite-bearing grospydite is outside the primary mantle range and is consistent with its derivation from subducted altered oceanic basalt. Pressure estimates for the grospydite SRV-1 correspond to equilibration depth of 150–160 km; subduction of oceanic crust has therefore occurred into the asthenosphere.

The oxygen isotope thermometric estimates in this study are the highest measured to date. The preservation of such high temperatures is due to the extremely refractory behavior of the grospydite mineral assemblage combined with a rapid quenching from peak temperatures. Oxygen isotope thermometry of refractory phases has previously been limited by the difficulty in fluorination of such materials (e.g., Kyser 1986). The limited and scattered data on kyanite confirm this problem. Recent advances in laser-based analytical techniques allow for the analysis of extremely refractory phases. Further oxygen isotope geochemical studies of ultra-refractory phases. such as garnet, olivine, coesite, kyanite, corundum and staurolite are expected to preserve peak temperature chemistries even when other phases have partially reset. Stable isotope thermometry may prove, in fact, to be one of the most robust thermometers in high temperature systems, as long as the appropriate mineralogy and mineral phases are chosen for analysis.

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