

THE UNIVERSITY OF MICHIGAN  
INDUSTRY PROGRAM OF THE COLLEGE OF ENGINEERING

ADSORPTION OF NITROGEN-METHANE ON  
LINDE MOLECULAR SIEVE TYPE 5A

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A dissertation submitted in partial fulfillment  
of the requirements for the degree of  
Doctor of Philosophy in The  
University of Michigan  
1961

March, 1961

IP-502



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## ACKNOWLEDGEMENT

The author wishes to take this opportunity to thank the many people who have assisted in the completion of this thesis.

Among these are Dr. T. L. Thomas of the Linde Company, who kindly supplied the Molecular Sieve and whose advice was most helpful; Mr. R. L. Porter of Autoclave Engineers for his as well as his associates' valuable time in aiding to solve some valve problems; Mr. Foglesong of The Babcock and Wilcox Company for his help in obtaining material for the adsorption vessel; The Kolene Corporation for Tufftriding some vital parts and the Dow Chemical Company for its gift of the necessary Styrofoam insulation.

Special thanks are due to Messers. Cleatis Bolen, John Wurster, William Hines, and Peter Severn, who assisted in building the equipment and were always ready with helpful suggestions. Mr. Frank Drogoz also deserves many thanks as do his assistants, Richard Naragon and David Williams for the help in the mass spectrometer analysis. Mr. Jude Sommerfeld assisted with the surface area determinations. The staff of the University Computing Center and the use of the IBM 704 Computer materially lightened the task of reducing the experimental data to a useful form.

To his committee the author desires to extend a special note of thanks for their patience, ever open doors and encouragement.

The Industry Program of the College of Engineering was of great assistance in preparing the final manuscript and in arranging for its presentation in this form.

The encouragement given me by my wife, as well as her aid in preparing the final draft was most helpful.

Finally, to my parents, for making this task possible, and for their continuing encouragement and guidance, I dedicate this work.



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## NOMENCLATURE

a	constant in single component Langmuir equation; milligram moles/gm. atm.
a'	constant in multicomponent Langmuir equation; milligram moles/gm. atm.
b	constant in single component Langmuir equation; l/atm.
b'	constant in multicomponent Langmuir equation; l/atm.
b''	constant in multicomponent Langmuir equation; l/atm.
b <sub>v</sub>	van der Waals' constant; cc./gm. mole
E	Schay's mixing coefficient
f	fugacity
k	constant in Freundlich equation; milligram moles/gm.
N	amount adsorbed; milligram moles/gm.
N'	amount adsorbed as pure component under the same pressure as the component partial pressure in the gas phase mixture.
n	constant in Freundlich equation
P	pressure; atm.
R	universal gas constant
T	temperature °Kelvin
v	volume adsorbed
x	mole fraction adsorbed phase
y	mole fraction gas phase
γ	constant
ε	adsorption potential; calories/gm. mole
ρ	density; gm. moles/cc.

## Subscripts

cr	critical
F	final
FS	free space (adsorption cell)
I	initial
i	i <sup>th</sup> component
m	maximum
s	at saturation conditions

## ABSTRACT

The adsorption of nitrogen and methane and their mixtures on Linde Type 5A Molecular Sieve has been studied in a batch system from  $-140^{\circ}\text{C}$  to room temperature and at pressures up to 85 atmospheres. The amount adsorbed at equilibrium is reported for the pure components. Data indicating the equilibrium gas and adsorbate compositions, as well as the total amount adsorbed, are presented for mixtures of nitrogen and methane.

Constants for the Langmuir equation, over the range of pressure and temperature studied, have been obtained for pure component adsorption. The Freundlich equation constants are also presented for the adsorption of pure nitrogen and methane and for total adsorption of mixtures as a function of temperature. The equations with the appropriate constants represent the capacity of the Molecular Sieve for pure components within five per cent and for mixtures within seven per cent of the total amount adsorbed.

Individual component adsorption from mixtures may be expressed as a function of the pressure of the components in the gas phase in equilibrium with the adsorbate. Constants for the system studied are presented as a function of temperature between  $100^{\circ}\text{Kelvin}$  and  $300^{\circ}\text{Kelvin}$ .

Limited data at room temperature for some other adsorbents, such as silica gel and some activated carbons, indicate that Molecular Sieves have a higher capacity for methane and nitrogen. A study made on another activated charcoal which had a slightly greater capacity shows that available surface area is a critical factor in determining capacity.

## INTRODUCTION

The present study was undertaken to analyze the adsorption system nitrogen - methane - Linde Molecular Sieve Type 5A\*. A method of predicting the adsorbate loading from both pure and mixed gas streams was to be developed. If possible, a model for the adsorption mechanism was to be postulated.

The nitrogen - methane system is of ever growing interest because these two gases appear together in many industrial applications. In addition, the present program to conserve helium will require removal of nitrogen and methane from that gas in varying quantities, for subsequent recovery.

The study, as undertaken, included a wide range of temperatures from room temperature to below the critical of the two gases as well as the pressure range of industrial interest. Temperatures varied from 123° Kelvin to 295° Kelvin and pressures ranged from essentially one atmosphere to as high as 50 atmospheres. The range of compositions at the higher temperatures varied from pure methane to pure nitrogen, while at the lower temperatures, only one or two mixtures were evaluated in addition to the pure components.

For purpose of completeness, a discussion of the various adsorption theories has been included, followed by a discussion of the experimental equipment. The experimental method and results are described and evaluated together with a series of conclusions and recommendations.

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\* Registered Trade Mark

## DISCUSSION OF THEORY

### Pure Component Adsorption Isotherms

Several basic approaches to physical adsorption equilibrium have been presented since the phenomenon was first observed by Scheele in 1773.<sup>(13)</sup> These may roughly be distinguished as the kinetic and the thermodynamic approaches. This distinction is somewhat artificial as the isotherms based on the kinetic approach, first advanced by Langmuir,<sup>(33,34)</sup> can be readily developed from thermodynamic or statistical principles. In the present context an isotherm defines the amount adsorbed as a function of pressure at constant temperature.

Langmuir first derived the isotherm which bears his name:

$$N = \frac{N_m b P}{1 + b P} \quad (1)$$

or

$$N = \frac{a P}{1 + b P} \quad (1a)$$

from kinetic principles<sup>(34)</sup>. This derivation requires that arbitrary forward and reverse rates be balanced. Fowler and Guggenheim<sup>(20)</sup> have derived the Langmuir isotherm on statistical grounds, considering that the adsorbed molecules do not interact with each other and had no translational motion and were thus confined to localized sites, or points of active adsorption. If one examines the Langmuir isotherm, one can readily see that at high pressures  $N$  approaches  $N_m$  which corresponds to the amount of adsorbate in a complete monolayer.

While the Langmuir isotherm fitted a great deal of experimental data, the three basic assumptions:

1. Uniform sites with no interaction between adsorbate molecules,
2. No translational motion on the surface,
3. The maximum adsorption corresponds to a complete monolayer,

placed severe restrictions on the theory. Brunauer, Emmett and Teller<sup>(15)</sup> in their well-known "B E T" equation made the first significant advance when they considered multilayers. The "B E T" equation takes several forms, of which

$$\frac{N}{N_m} = \frac{c P/P_s}{(1-P/P_s)[1 + (c-1)P/P_s]} \quad (2)$$

where the number of layers, or stacks of molecules above the solid surface are finite, is the most widely used form. Several modifications of the original "B E T" equation were proposed by Brunauer et al.<sup>(15)</sup>, where the number of layers are limited, and Brunauer and Demming<sup>(14)</sup> for the case of adsorption on capillary walls. Joyner<sup>(31)</sup> has discussed in detail a method of determining the number of adsorption layers in the complete "B E T" equation by graphical methods.

The original derivation of the "B E T" isotherm suffered from the same limitation that Langmuir's derivation did, in that it was based on kinetic equilibria. Hill<sup>(24)</sup> and others have derived the equation from statistical principles, based on a model allowing vertical but no horizontal interaction between molecules. The "B E T" equation still suffers from the same basic deficiencies as the Langmuir isotherm: no interaction of adsorbate molecules and no lateral motion of the adsorbate molecules in any given layer. More recently, a multi-molecular adsorption model, allowing for lateral interaction, has been suggested by Lee<sup>(35)</sup>.

Unfortunately, no experimental work has been done to confirm the postulated isotherm.

Peierls<sup>(46)</sup> and Wang<sup>(61)</sup> derived the Langmuir equation from statistical principles for adsorbed atoms, exhibiting attractive and repulsive interaction respectively. Schay<sup>(56)</sup> has derived Langmuir's isotherm on statistical grounds in the case of a mobile monolayer, where the adsorbed molecules exhibit a covolume in the adsorbed state, thus removing the last serious theoretical restriction.

The Freundlich isotherm<sup>(21)</sup>

$$N = KP^{1/n} \quad (3)$$

has been derived from the Langmuir isotherm by Zeldowitsh<sup>(63)</sup>. By assuming that the adsorption surface is heterogeneous, the Langmuir equation may be written as

$$N = \sum_i \frac{a_i P}{P + c_i} \quad (4)$$

where  $c_i$  depends mainly on the heat of adsorption on the adsorption spaces and  $a_i$  is a function of the number of these spaces. If "c" is said to vary continuously over the surface and the distribution function of "a" assumed the form

$$a(c) = Ac^{1/n-1} \quad (5)$$

then Zeldowitsh arrives at the Freundlich isotherm.

The isotherm may also be derived for multi-molecular adsorption, as was shown by Baly<sup>(3)</sup>, who assumed that each adsorbed layer obeys a Langmuir isotherm with different constants, and that each succeeding layer has a slightly lower heat of adsorption than the one below. These assumptions lead to the Freundlich isotherm for moderate pressures.

More recently, Sips<sup>(57,58)</sup> has shown that a generalized Freundlich isotherm:

$$N = (P/(P + a))^c \quad 0 < c < 1 \quad (6)$$

may be obtained by assuming a localized adsorption model with no interaction. Equation (6) reduces to the Freundlich equation for small "P", and when "c" is equal to one the Langmuir equation results. Unfortunately, this theory does not appear to fit the experimental data. Sips' assumptions of localized sites and no molecular interaction may be necessary from a mathematical point of view, but present a simplified picture of the physical process, thus leaving much to be desired.

In an excellent review of the integral inversion approach used by Sips, Honig<sup>(28)</sup> indicates that the isotherm is very insensitive to changes in the chosen distribution function. In addition, all attempts to obtain an isotherm by this method require a model with no molecular interaction. He further states that many of the most useful isotherms do not satisfy all of the mathematical restrictions found to be necessary in applying this method of analysis. The isotherm suggested by Sips is, as he states, the simplest satisfying all the necessary restrictions. A real physical restriction, that at very low pressures, the amount adsorbed is proportional to the pressure is not satisfied except in the case of "c" equal to one. It would seem, then, that Sips' model is too simple to explain the total adsorption isotherm. Real improvements in obtaining valid adsorption isotherms must come from advanced or improved models of the adsorption phenomena.



Redlich and Peterson<sup>(52)</sup> recently proposed an isotherm

$$N = aP/(1 + cP^\gamma) \quad (7)$$

where  $0 < \gamma < 1$

which fits much of the experimental data well. It further reduces to the phenomenologically sound, limiting Langmuir equation

$$N = ap \quad (8)$$

at low pressure. At moderate pressure the equation reduces to the Freundlich isotherm

$$N = \frac{a}{c} p^{1-\gamma} \quad (9)$$

which appears to be sound in this range. The equation, however, breaks down as the pressure becomes very high.

At present the kinetic-statistical approach, while indicating correct avenues, has, due to the limitations imposed by the simplified models, not fully explained the adsorption process.

About the same time that Langmuir first advanced his kinetic adsorption model, Polanyi suggested a model based essentially on thermodynamic principles. As first proposed by Polanyi<sup>(47,48,49,50)</sup>, the theory assumed that long range forces controlled the adsorption. This basic assumption was later modified by Polanyi and Goldman<sup>(22)</sup>, so that the theory would conform with the concept of molecular forces.

Basically, the Potential Theory of Adsorption assumes that work must be done to bring the molecules from the gas phase into the adsorbed phase, and that this work,  $\epsilon$ , is work of compression:

$$\epsilon = \int_{p_g}^{p_a} V dp \quad (10)$$

Then, if any layer "i" is considered, one can obtain the work required,  $\epsilon_i$ , which is called its potential, and the equipotential surfaces will vary from a maximum of  $\epsilon_0$ , the potential of the adsorbent-adsorbate interface, to zero, the potential at the gas-adsorbate interface. Each equipotential surface will enclose a given volume,  $v_i$ , around the adsorbent. Thus it was shown that the volume adsorbed and the work expended in the adsorption process are related. Polanyi denoted this relationship

$$\epsilon_i = \phi(v_i) \quad (11)$$

as the characteristic curve dependent only upon the adsorbent-adsorbate system.

Although Polanyi first applied the theory, Berényi<sup>(7,8)</sup> improved the method considerably. Both found that the computation of the characteristic curve, although fundamentally the same over the whole range of interest, required slight alterations near the critical point and above the critical temperature.

Well below the critical temperature, the characteristic curve may be computed, if one assumes an ideal gas model; 1) by integrating the compressive work integral between the limits of the system and saturation pressures respectively; thus:

$$\epsilon_i = \int_P^{P_s} VdP \quad (10)$$

or

$$\epsilon_i = RT \ln(P_s/P) \quad (10a)$$

and 2) if one assumes an incompressible liquid, the adsorbed volume

becomes:

$$v_i = N_i/\rho_s \quad (12)$$

where  $\rho_s$  is taken as the density of the saturated liquid at the adsorption temperature. Berenyi suggested corrective terms for both the potential integral and the density near or above the critical temperature, so that near the critical:

$$\epsilon = RT \ln (P_s/P) + \mu \quad (13)$$

$$v = N/\rho_s \quad (12)$$

and above the critical,

$$\epsilon = RT \ln \frac{0.14T}{Pb_v} \quad (14)$$

where  $b_v$  is van der Waals' constant.

The model as originally suggested has been used effectively; more recently modifications have been introduced, in particular for the conditions near or above the critical to remove the trial and error part of the solution for the correction terms originally introduced by Berényi. Nikolaev and Dubinin<sup>(43)</sup> have proposed that above the critical the potential be computed by:

$$\epsilon = RT \ln \tau^2 (P_{cr}/P) \quad (15)$$

where

$$\tau = T/T_{cr} \quad (16)$$

and the adsorbate volume:

$$v = Nb_v .$$

They obtain this result by showing at the critical temperature the characteristic curve obtained using the normal method of evaluating  $\epsilon$ ;

$$\epsilon = RT_{cr} \ln (P_{cr}/P) \quad (17)$$

and using Berenyi's modification above the critical

$$\epsilon = RT_{cr} \ln \mu (P_{cr}/P) \quad (18)$$

will coincide when

$$\mu = (P_{cr} b_v / T_{cr}) \tau \quad (18a)$$

In addition, Nikolaev and Dubinin<sup>(43,17)</sup> suggests that above the normal boiling point a density:

$$\rho^* = \rho_s - ((\rho_s - \rho_v) / (t_{cr} - t_s))(t - t_s) \quad (19)$$

where  $\rho_v$  is equal to van der Waals' constant, would better represent the density of the adsorbate. This density then will take into account the existence of the adsorbate in a compressed state due to the attractive forces of the adsorbent.

The deviation from ideality of the gas in equilibrium with the adsorbate should influence the adsorption potential. Although Polanyi in his original work suggests this, he did not develop the idea. Lewis and co-workers<sup>(27)</sup>, in correlating some adsorption isotherms of hydrocarbons at moderate pressures, found that the deviations were sufficient to warrant introducing corrections for the deviations from ideality. They suggest the use of fugacity, the most reasonable approach, instead of pressure. Thus,

$$\epsilon = RT \ln (f_s/f) \quad (20)$$

where, from their calculations,  $f_s$  is the fugacity at the saturation pressure, obtained by using the generalized  $f/p$  plots. Above the critical

$f_s$  was set equal to the fugacity at the pseudo-saturation pressure as determined by extrapolating the usual vapor pressure curve:

$$\ln P_s = (-A/T) + B \quad (21)$$

where "A" and "B" are constants.

Lewis and his co-workers suggested that the density of the adsorbate was equal to the saturated liquid density at the adsorption pressure. It would seem that their analysis is quite satisfactory except for their assumption regarding density, as liquid density is principally a function of the system temperature.

A combination of the work of Dubinin and co-workers with that of Lewis should give a better representation of the adsorption potential. Thus, the use of fugacities, to represent the driving force, evaluated at the system pressure and saturation or the critical condition, together with the density evaluated by Dubinin's method should give a good characteristic curve; that is a single function

$$\epsilon = \phi(v) \quad (22)$$

independent of temperature.

Although the Potential Theory does not clearly explain the adsorption phenomena, it does present a satisfactory explanation for the gross process. The process above the usual critical point is identical in nature to that below the critical point in that a dense phase forms. This phase, while not a liquid as normally envisioned, does possess many of the same properties. The theory does have the very fine property of generalizing the system from one isotherm. Lack of knowledge regarding the dense or adsorbed phase limits the general applicability of the

theory at present. Hopefully, added knowledge in this area will make this theory more useful.

The deterministic approach, illustrated by Hill, Schay and many others may lead to a number of models and isotherm equations. At present, no totally satisfactory model has been suggested to explain in particular the adsorption phenomena at moderate pressures. Most probably, this is the result of the many simplifications necessary to enable one to treat the phenomena from the statistical viewpoint and still obtain mathematical models which may be solved analytically. It may be that by the application of numerical methods at some future date more complex models, which allow for all types of molecular interactions, may be proposed and a more definitive test for determining the applicability of specific models at moderate pressures may be found.

#### Adsorption From Mixtures

Much of the work done to date on adsorption from gaseous mixtures falls into the category of dehumidification. Some work has been done on removal of trace impurities from gaseous streams. In this area two works, that of Johnson<sup>(29)</sup> and Hiza<sup>(28)</sup> are of some interest as they deal with the removal of trace amounts of nitrogen and methane by silicon gel adsorption from a high purity hydrogen stream.

A good review of the ~~early~~ work in mixed adsorption is given by Brunauer.<sup>(13)</sup> Very little work has been done in this area. The first work of any significance is that of Markham and Benton<sup>(40)</sup>, which extend Langmuir's isotherm to mixtures. Assuming the same model as Langmuir, the component isotherms resulting from the solution of the

rate equations (two adsorption and two desorption) are:

$$N_1 = \frac{N_{1m} b_1 P_1}{1 + b_2 P_2 + b_1 P_1} \quad (23a)$$

$$N_2 = \frac{N_{2m} b_2 P_2}{1 + b_2 P_2 + b_1 P_1} \quad (23b)$$

Equations (23a) and (23b) indicate a mutual decrease in the adsorption of both components at a given partial pressure. This has been shown to be true in general; the equation does not take into account the possibility of interaction between the two components in the adsorbed state.

The relationship appears to hold well at low concentrations, but breaks down at the higher concentrations reinforcing the problems caused by interaction. Schay et al.<sup>(57)</sup> have developed essentially the same equation as Markham and Benton:

$$N_1 = N_{1m} \frac{a_1 P_1 / E_1}{1 + \sum_j (a_j P_j / E_j)} \quad (24)$$

where  $E_j$  is a function of the interaction of the molecular species in the adsorbed state and is defined by Schay as:

$$\ln E_k = \left[ \sum_j n_j \beta_j - \beta_k \frac{\sum_j n_j}{\sum_j n_j \beta_j} \right] \left[ \frac{A}{\sum_j n_j \beta_j} \ln \left( 1 - \frac{\sum_j n_j \beta_j}{A} \right) + 1 \right]$$

$\beta$  = two dimensional residual molecular volume

$A$  = adsorbent area

$n$  = number of molecules adsorbed at the asymptotic value of the isotherm.

For a binary component, if  $E_1 = E_2 = 1$ , the result is identical to that of Markham and Benton. In addition, Schay and co-workers have shown that the

empirical relationship, for a binary mixture, suggested by Williams<sup>(62)</sup> and observed by Lewis and Gilliland<sup>(36)</sup>

$$\frac{N_1}{N'_1} + \frac{N_2}{N'_2} = 1 \quad (25)$$

can be shown to be:

$$\frac{N_1}{N'_1} + \frac{N_2}{N'_2} = 1 + \frac{\frac{1}{P} \left( \frac{P_1}{E_1} + \frac{P - P_1}{E_2} \right) - 1}{1 + \frac{a_1}{E_1} P_1 + \frac{a_2}{E_2} (P - P_1)} \quad (26)$$

for their model. Thus Equation (25) is strictly correct for

$$E_1 = E_2 = 1$$

which indicates essentially no interaction of the species in the adsorbed state. In practice, Schay has found that  $E_1$  and  $E_2$  each differ slightly from unity, one being greater the other less and Lewis' observation is still valid in practice.

Redlich and Peterson<sup>(52)</sup> have suggested that their equation for pure components (Equation 8) is also valid for mixtures if it is modified in a manner analogous to the Markham modification of Langmuir's isotherm for multi-component adsorption. This results in:

$$N_1 = \frac{a_1 P_1}{1 + c_1 \frac{P_1^{\gamma_1}}{1} + c_2 \frac{P_1^{\gamma_2}}{2}} \quad (27)$$

which does not account for any interactions if  $a_i$  and  $c_i$  are identical values for the mixed isotherm as for the pure component. A modification for interaction of the molecular species, similar to that suggested by Schay, can be introduced so that the form of the equation would be:



$$N_i = \frac{a_i P_i / E_i}{1 + \sum_j (c_j P_j^2 / E_j)} \quad (28)$$

Very little work has been carried out in the area of co-adsorption of mixtures so that the phenomena of competing adsorption has not been studied except as an extension of pure component adsorption. At the present time it would appear that very little can or should be said in this area until the phenomenon of pure component physical adsorption has been more clearly defined.

A detailed sub-microscopic study of the solid adsorbent surface and the surface forces must be undertaken before a clear picture of the adsorption phenomena can be developed. In addition the molecular state of the adsorbate, both for pure and multicomponent systems, requires examination. Until the adsorption phenomena is clearly established, a semi-empirical approach such as that discussed above, must suffice for engineering purposes.

## EQUIPMENT

### Adsorption Apparatuses

There are basically three methods of measuring adsorption equilibria, the continuous once-through flow system, the recirculating flow system and the closed system. Each system or method has inherent advantages and disadvantages.

The continuous flow system requires careful control of the flow rates and continuous stream analysis. Both of the requirements can easily present experimental difficulties as they must often be determined and controlled within very small tolerance limits. A modification, used by Linde<sup>(51)</sup>, eliminates these difficulties to some extent. In this modified flow system, adsorbate is passed over the adsorbent bed until flow rates at the entrance and exit of the cell are equal. The cell effluent is then sampled and the cell taken off the line. The adsorbate is desorbed; the desorbed material is collected and its mass and composition determined. Knowing the gas phase composition and volume, the adsorbate volume and composition may be determined by material balance. Both the once-through and modified once-through flow systems require large amounts of adsorbate gas and a great deal of coolant at low temperature. They do, however, enable the experimenter to obtain kinetic data, as well as equilibrium data.

The recirculating flow system, such as the one described by Lewis and co-workers<sup>(37)</sup>, is quite satisfactory for equilibrium data but can present problems as there is a greater possibility of leaks. As it recirculates the gas, a uniform system will be achieved and by appropriate

pressure measurements and knowing the free space volume one can determine the amount adsorbed. As it has the additional feature of a large gas space not in contact with the adsorbent, it is easier to control the gas phase equilibrium composition and even the pressure by adding small increments of gas as the adsorption process proceeds. This system requires mechanical agitation, or rocking, which can present problems when considered in respect to a system that must be maintained at low temperatures and high pressures. In addition, the entire system must be leak-free which in itself can be a difficult problem under severe operating requirements and constant movement. This system is probably best adapted to operating pressures below ten atmospheres and temperatures in the ambient region.

By far, in many respects, the simplest system is the closed system in which a measured amount of sample is admitted to the equilibrium vessel and in the case of the mixture, the gas phase is sampled when equilibrium is reached. Then, knowing the amount of each constituent in the gas phase, the equilibrium adsorbate can be determined by material balance. Several variations of this type of apparatus have been used. At low pressures and particularly for pure components, a gravimetric adsorption apparatus, as typified by the sorption balance of McBain and Bakr<sup>(38)</sup>, is very useful in that the amount adsorbed can be determined directly. Sawyer, Josefowitz and Othmer<sup>(29)</sup> have made modifications to the spring and general assembly which make the apparatus more durable. At high pressure, modifications of this balance have been developed by McBain and Britton<sup>(39)</sup> and Morris and Mass<sup>(42)</sup>.

Volumetric apparatuses for low pressure work are used frequently. The simplest apparatus of this type is probably that of Pease<sup>(44)</sup> and is

similar to ones used by Homfray, Titoff and others. Coolidge<sup>(16)</sup> developed an apparatus of this type which had no stopworks. A high pressure volumetric adsorption apparatus which was suitable for work over a wide temperature range was built by Antropff and co-workers.<sup>(1)</sup> This apparatus is excellent for pure component adsorption studies at high pressure and was used to determine the adsorption isotherms of Nitrogen and Argon. The equipment included a mercury pressure transmitting device to keep the gas volume constant. At low temperatures a temperature gradient in the gas in the pressure transmitter could present difficulties. In addition, there was some gas volume which never came in contact with the adsorbent; this presents no difficulties when working with pure components but makes it impossible to use this apparatus for multicomponent adsorption studies.

#### The Present Apparatus

A schematic diagram of the system used in the present study can be found in diagram I. There are two large feed tanks (A), size 1A cylinders, which contain pure component. These are connected to the main feed reservoir (B), which is a high pressure vessel fabricated of 2 1/2-in. extra heavy-duty pipe and has a volume of about 940 cc. The feed reservoir is connected by a quarter inch high pressure line to the feed pressure gauge (C), and the distribution block. From the distribution block, lines lead to the sampling connection (D) and the vacuum and expansion system (E). A third line leads through the sampling connection (K), through the low temperature cell block valve (L), to the adsorption cell (F) which is discussed in greater detail below. The vacuum and

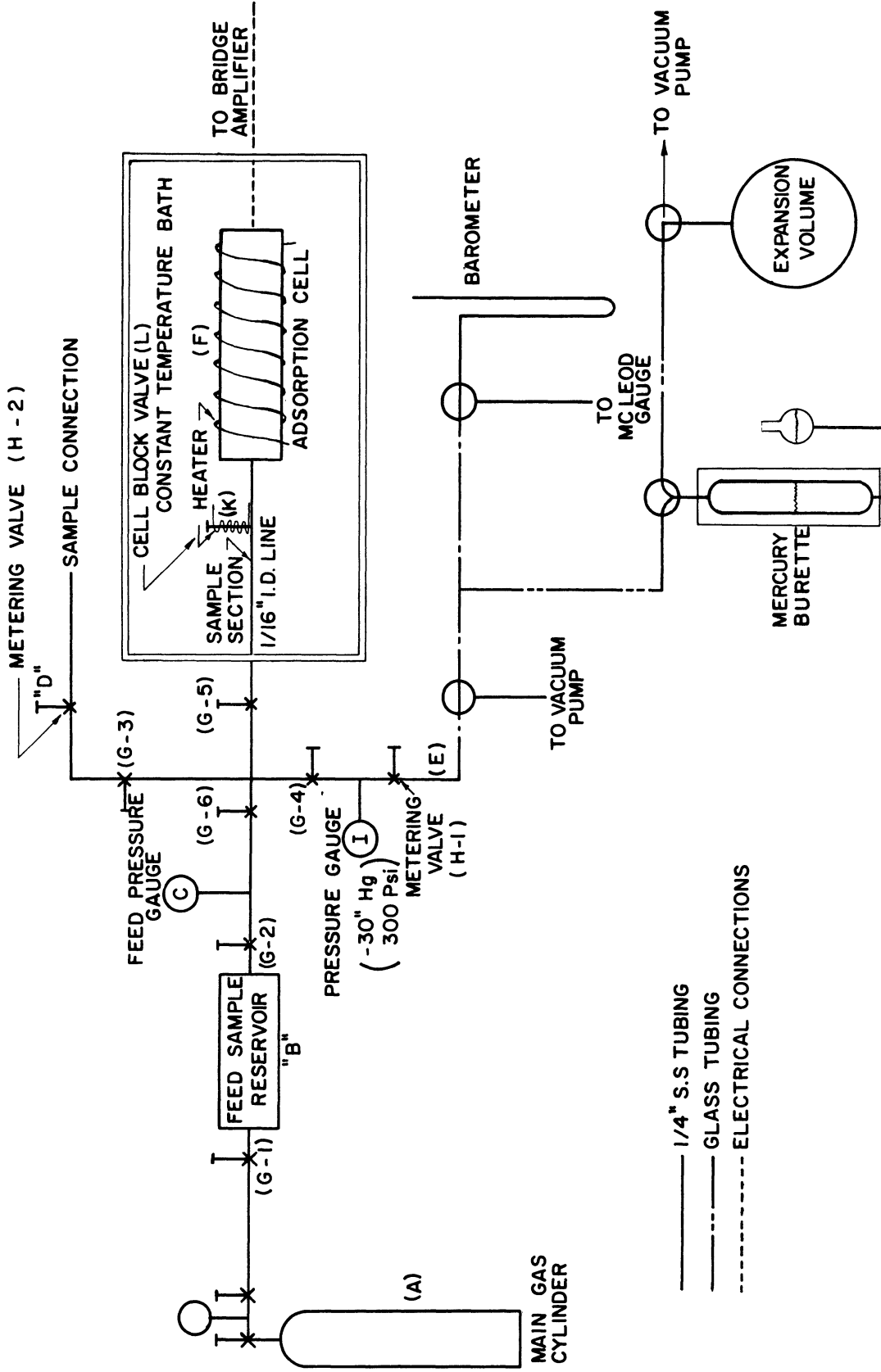


Figure 1. Flow Diagram of Adsorption Apparatus.



Figure 2. Top View of Closed Bath (showing protective steel shield).

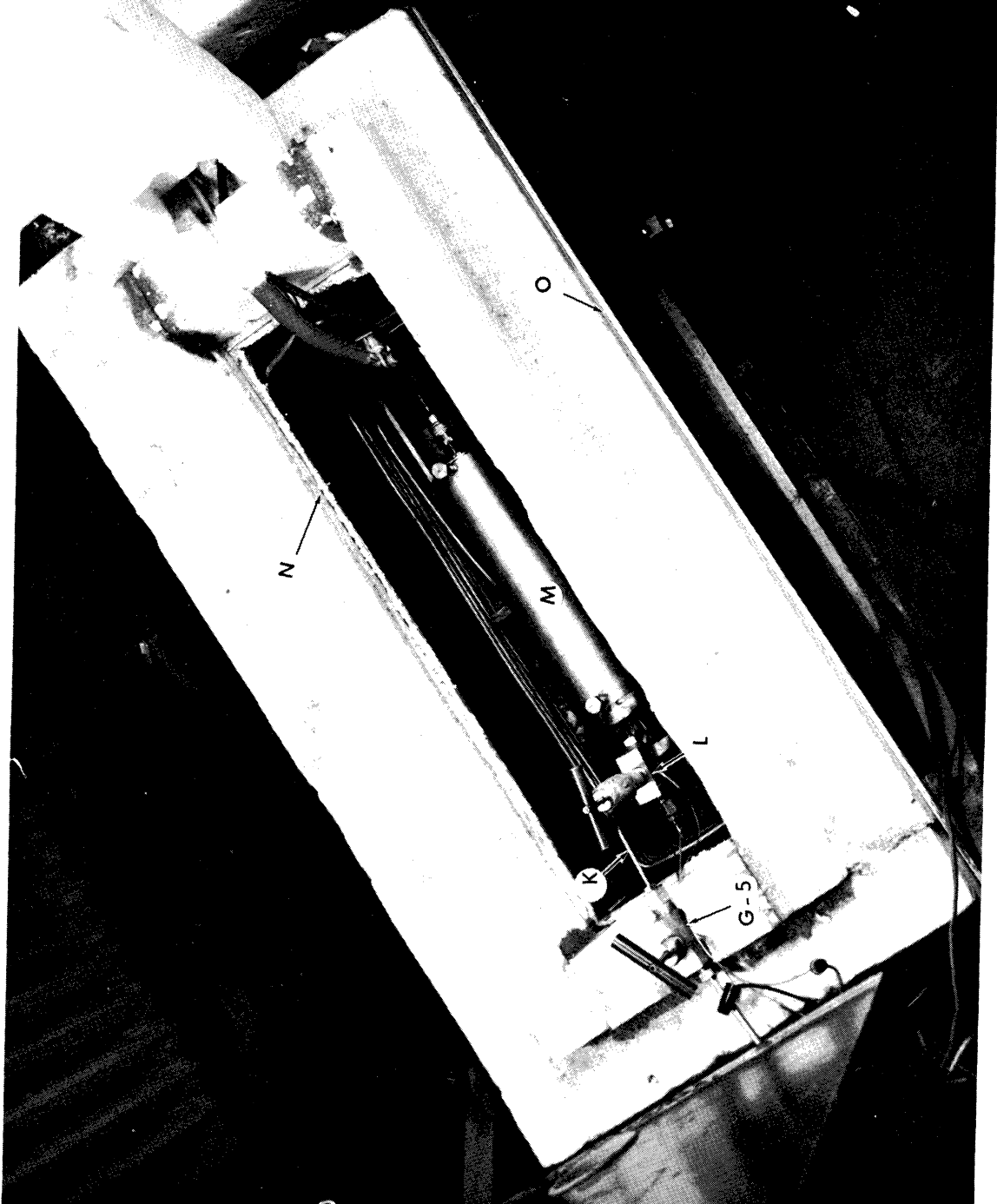


Figure 3. Open Constant Temperature Bath.

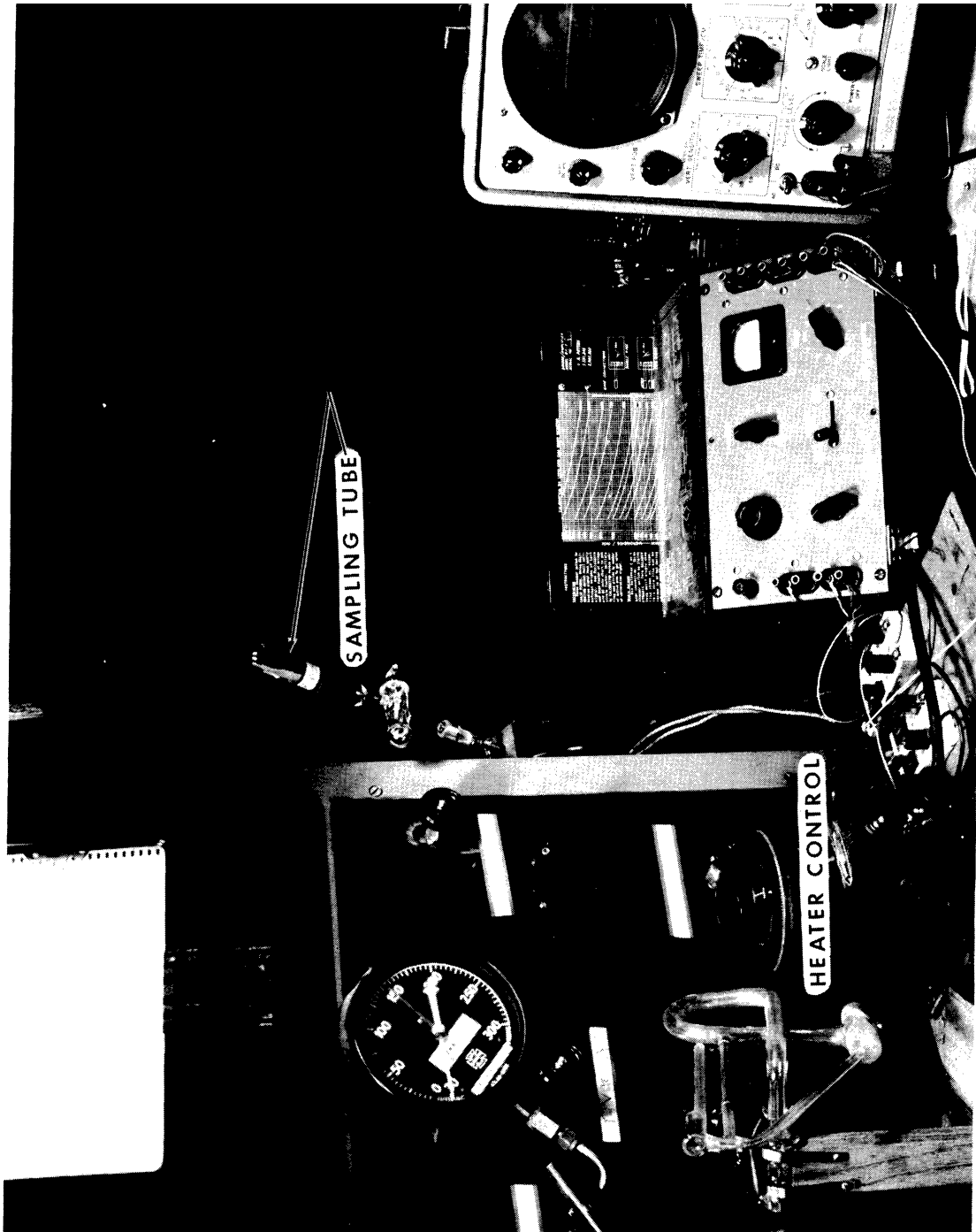


Figure 4. Instrument and Control Section.



expansion systems are made of glass and are protected against high pressure surges by standard Autoclave 10,000 psi. valves (G3, G4) and Hooke metering valves (H1, H2) in series. A low pressure gauge (I) is placed in between the block and metering valves on the vacuum line as an additional precaution. The entire system is at room temperature except for the adsorption cell and cell sampling line (K) which are in the constant temperature bath.

### The Adsorption Cell

The major modifications from other similar system were made to the adsorption cell. The cell, shown in detail in Figure 3, was made of 316 stainless steel and designed according to the ASME-API Code for Unfired Pressure Vessels.<sup>(2)</sup> Detailed design calculations may be found in Appendix D.

The two ends were machined from bar stock and the center cylinder is a piece of 1 1/2-in. schedule 180 seamless tubing donated by Babcock and Wilcox. The ends were joined to the cylinder using 308 stainless steel welding rod in order to prevent the formation of delta ferrite. The vessel was then x-rayed to insure the soundness of the welds. In order to release any stresses produced during welding or machining, the entire vessel was heat treated as 750-800°F for two hours and then furnace cooled.

The quarter inch tubing connects to the main block valve (L) and is fitted with a fine stainless steel screen to prevent particulate matter from fouling the valve. The thermocouple well is standard 1/8 inch high pressure tubing. The 18 mm. threaded aperture serves both for filling the vessel with adsorbent and as the connector for a

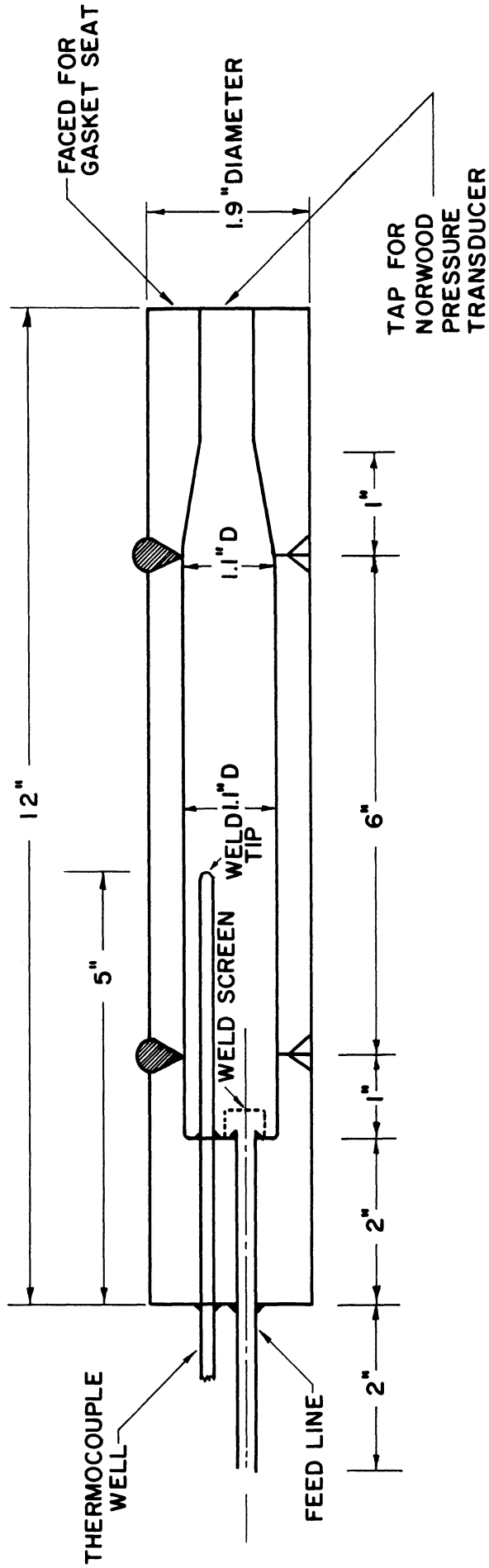


Figure 5. Adsorption Cell in Cross Section, (F).

pressure transducer used for determining the pressure in the cell. This end is faced for a gasket and a teflon gasket is used. Thirty foot-pounds torque is sufficient to seal the gasket and insure a tight system.

The use of a pressure transducer eliminates the need for any long connections where gas not in contact with the adsorbent can collect. In order to further reduce any "dead" space the quarter inch outside diameter connecting line is filled with clean copper wire.

The cell sampling connection (K) is a length of 1/8 inch outside diameter tubing, also filled with copper wire so that the sample size can be kept to a minimum.

Thus, the main feature of this system is the low "dead" space, assuring good contact between the entire gas sample and the adsorbent.

#### The Pressure Transducer

Pressure transducers are basically strain gauge devices. Several manufacturers make a variety of these devices and the Norwood Controls Transducer Model 101, of which Figure 6 is a schematic diagram, was picked because of its ease of operation, good response and the wide choice of pressure ranges available.

The pressure signal is received by a diaphragm which compresses a strain tube which supports the center of the diaphragm. Any changes in the strain tube are immediately detected by two strain gauges, one wound circumferentially and the other longitudinally. The gauges stretch and shorten respectively with an increase in pressure.

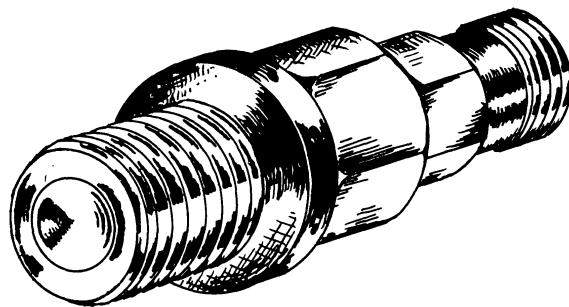
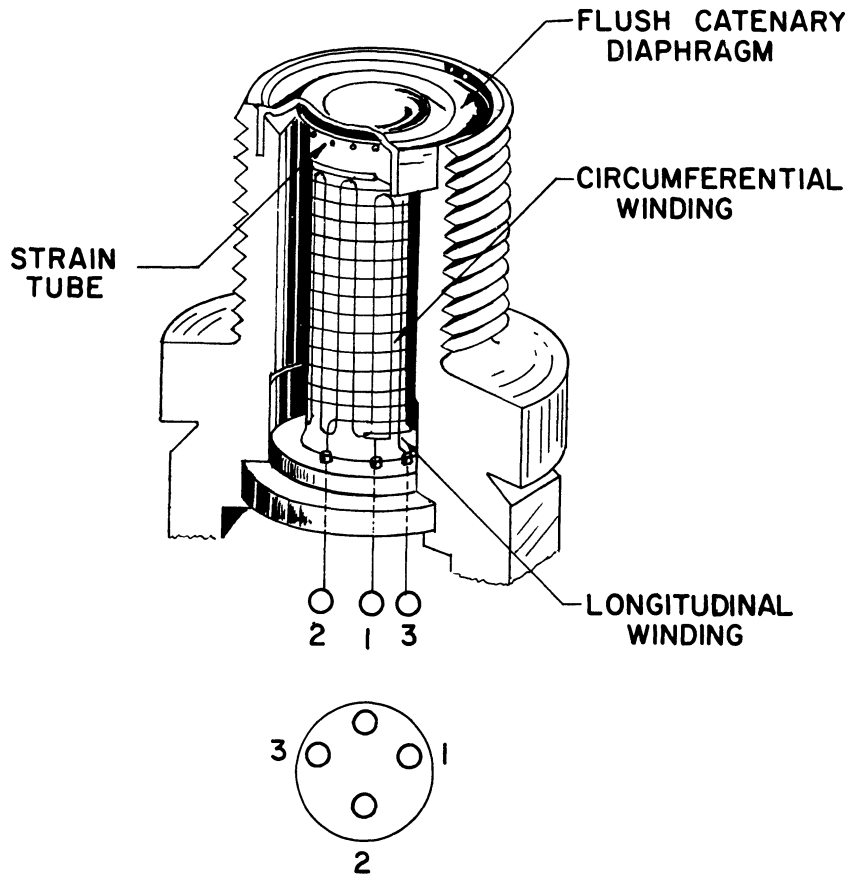


Figure 6. Schematic Diagram of Norwood Pressure Transducer.  
Courtesy of Norwood Unit - Controls Division  
American Standard Corporation.

The output produced by changes in the strain gauge is transmitted to an Ellis Bridge Amplifier which contains the other elements of the bridge circuit as well as an amplifier. The output signal from the amplifier is proportional to the induced strain and is read on a Model 130B Hewlett Packard Oscilloscope. The complete circuit including amplifier and scope are shown schematically in Figure 7.

TABLE I

NORWOOD PRESSURE TRANSDUCER SPECIFICATIONS

		<u>-15 to 500 psig Transducer</u>	<u>-15 to 2000 psig Transducer</u>
Accuracy	$\pm 1\%$ of rated pressure	<u>+50 psi</u>	<u>+200 psi</u>
Linearity	$\pm 1\%$ over entire range		
Resolution	0.1% of rated pressure	5 psi	20 psi
Repeatability	1/4% of full scale	12.5 psi	50 psi
Temperature Effect	0.02%/°F may be balanced out at any given temperature		

The specifications listed in Table I and the calibrations found in Appendix B, clearly show that these transducers are acceptable as secondary pressure measuring devices, if used with a linear amplifier and sensitive display device such as the oscilloscope mentioned previously. There is, however, a very definite need to choose the range of the transducer carefully so that the rated pressure is not exceeded by fifty per cent and that the pressure to be measured is greater than ten per cent of the rated pressure. Within these bounds the transducer will perform most satisfactorily.

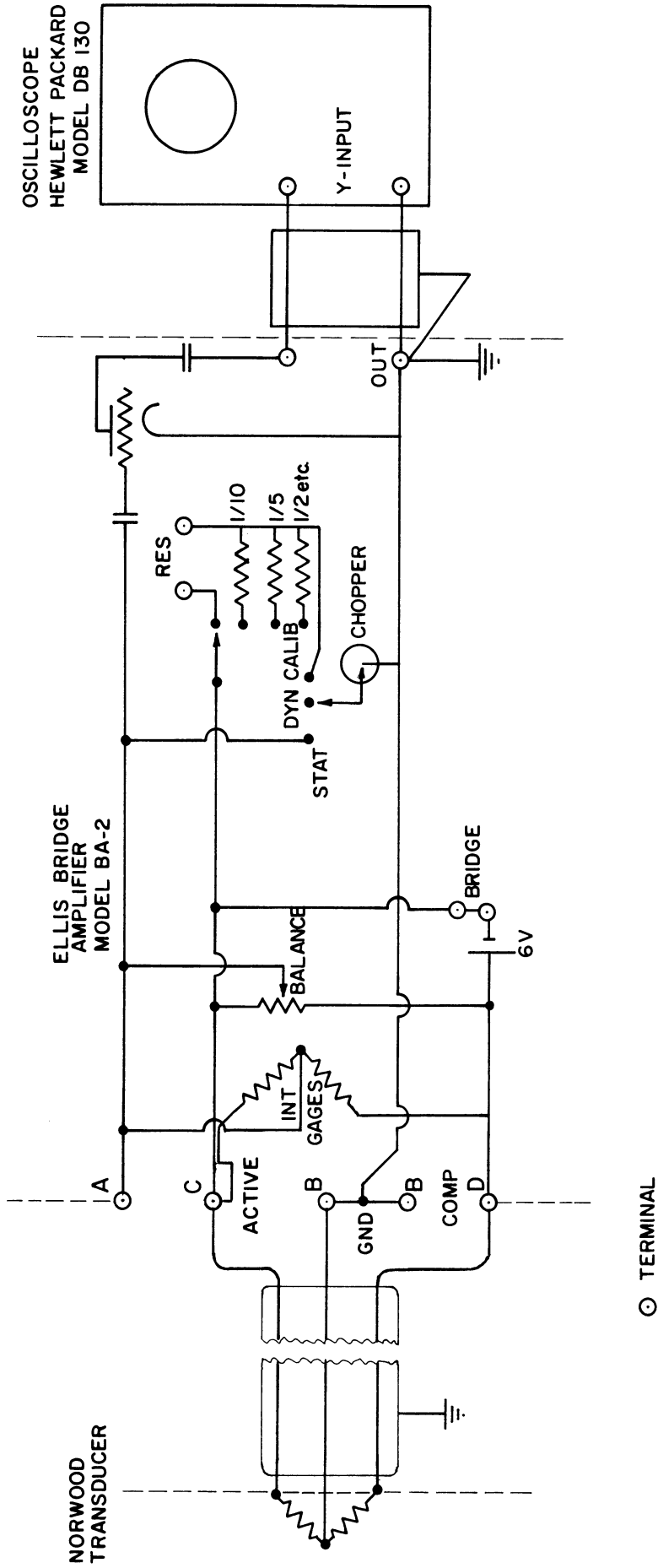


Figure 7. Simplified Schematic of Electrical Measuring Circuit.

In order to cover the desired range of pressures adequately, two transducers were used, a 2000 psi transducer and a 500 psi transducer.

#### Material Selection

As low temperature operation was anticipated, all high pressure lines, vessels and connectors were made of austenitic stainless steel. All welds were made with 308 rods to prevent delta-ferrite formation. Wherever possible, standard parts of 316 stainless steel were purchased. Special stainless steel packing washers and packing nuts, rather than the standard bronze parts, were used in the valves which were exposed to low temperature. All valve packings used in valves at low temperature were made of Dilecto, a glass impregnated with teflon. This material is dimensionally more stable than teflon and proved an excellent packing. Some difficulty was encountered at low temperature in the main block valve (L). This is discussed in some detail below.

The inner liner of the constant temperature bath is made of an 1/8 inch copper plate. The cooling coils were also constructed of copper. Around the copper tank there are six inches of Styrofoam insulation.

#### Valve Operation at Low Temperature

The operation of valves at extreme temperatures can present a problem. The basic problems encountered are the differences in thermal expansion of various metals and packings and the lack of good, low temperature lubricants.

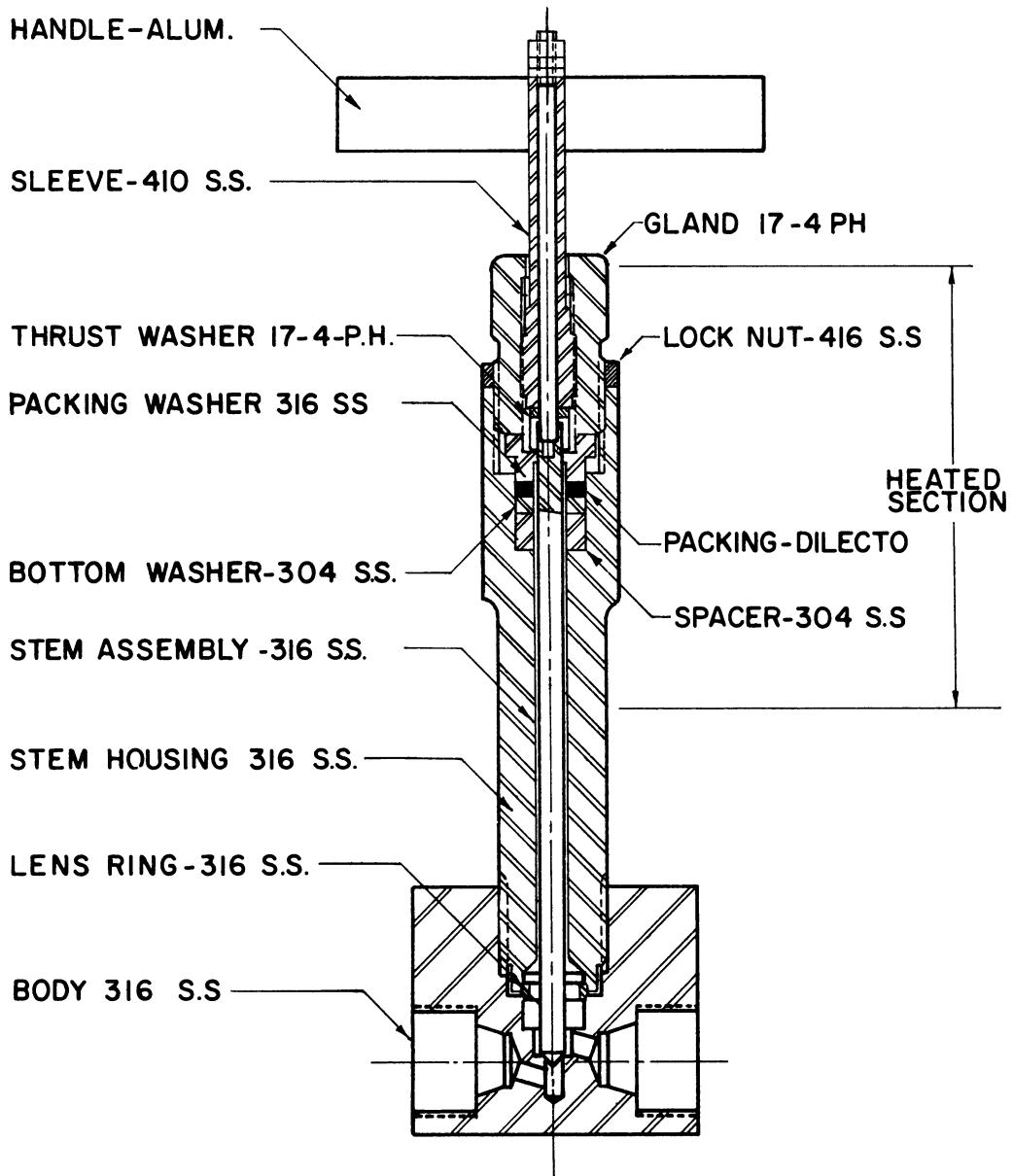


Figure 8. Cross Section of Main Block Valve (L).  
Courtesy of Autoclave Engineers Inc.



A standard Autoclave high pressure valve with stainless steel used throughout proved unsatisfactory, possibly because of the incorrect choice of material as seizing occurred between the stem (420 stainless steel) and the bottom washer made of 17-4PH Armco steel.

A high temperature valve, Figure 8, was inserted and with several modifications worked well throughout the tests. This valve has a two piece stem, the sleeve of which is made of 410 stainless steel. At low temperatures, where it is very critical that all contacting parts be clean and dry, this sleeve tends to gall in the packing nut threads. In order to eliminate this problem the stem must be treated so that its surface has properties different from those of the packing nut. This may be accomplished by either nitriding the surface or putting on a surface film which can withstand low temperatures. Both methods were tried, the former by the Tufftride\* process, and the latter with a coating of teflon about 0.002 inches thick. Both methods appeared to work well, although no extensive tests were made. During the course of the tests, the Tufftrided stem was used and the long stem was wound with a heater to keep the valve packing and gland nut at temperatures no lower than  $-100^{\circ}\text{C}$ .

#### Temperature Measurement and Control

Temperature in the cell was measured with a thermocouple made of matched 30 gauge copper and constantan wire. A continuous record of the cell temperature was made and the recorded reading was checked periodically with a Leeds and Northrup Semi-Precision Potentiometer.

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\* Trademark of the Kolene Corporation.

Another copper-constantan thermocouple was attached to the outside of the cell. The output from this thermocouple was recorded by a Micromax recorder-controller. The control element was an on-off switch controller Skinner three way valve. This valve, when on, allowed gaseous nitrogen, under pressure, to force liquid nitrogen into the constant temperature bath. With the aid of a 192 watt heating element wound around the cell, and controlled by a Variac, temperatures could easily be kept at  $\pm 0.5^{\circ}\text{C}$  within the cell.

The temperature of the reservoir was measured with a copper-constantan thermocouple and Leeds and Northrup student type potentiometer. The temperature of the expansion volume and gas burette were determined with standard thermometers.

Temperatures taken with the thermocouples could be determined to  $\pm 0.5^{\circ}\text{C}$  under the most severe conditions,  $-192^{\circ}\text{C}$  and within  $\pm 0.1^{\circ}\text{C}$  at ambient conditions. The thermometers could be read to  $\pm 0.2^{\circ}\text{C}$ .

#### Safety Considerations

Whenever experimental work is carried out at high pressures and extreme temperatures the safety of both personnel and equipment becomes an important consideration.

The adsorption cell, after welding, x-raying and stress relieving was hydraulically tested to 8,500 psi which was about one and a half times the design pressure. The entire high pressure system was then hydraulically tested at 7500 psi. In addition to the hydraulic testing, several safety shields were provided as can be seen in Figure 3. The inner shield (M) is a 2 1/2-inch schedule 80 stainless steel pipe

with the end near the transducer closed, to provide a shield should the transducer fail. A 1/8 inch copper bath (N) provided additional protection. One quarter inch steel plate (O) served as a final shield with space for the gas to escape between the hinged cover and the fixed tank. The temperature of this shield was checked with the bath at liquid nitrogen temperature (-196°C) in order to insure that its temperature was above the brittle fracture temperature, about 30°F. The temperature was about 65°F with a room temperature of about 72°F.

A blow out disk was provided in the high pressure lines and was rated at 3,000 psi or one and a half times the maximum normal operating pressure.

In order to protect the glass vacuum system and sample bottles, a block valve (G) and a metering valve (H) were placed in series and the two were never opened together if the pressure upstream from the block valve exceeded two atmospheres.

These precautions have proved sufficient; to date no accident has occurred.

#### Leakage in the System

In a static system it is important that the test cell in particular be leak-free. In order to determine the leakage, if any, in the test cell, a series of pressure tests were conducted over a twenty-four hour period each. In addition, tests at vacuum were conducted in order to determine if any leakage at low pressure was present. The cell was also tested using the pressure method with a helium detector.

The high pressure tests were conducted at 2,000 psi. It was found that no visible leakage, as determined by readings on the pressure gauge, occurred in a twenty-four hour period. This would mean that any losses would be less than one part in a thousand during that period of test. This was deemed sufficient to consider the system free of any leakage. Vacuum tests were conducted on the cell. The cell was evacuated to a pressure of ten microns and was then closed. The rest of the system was kept under a vacuum of ten microns and after a period of eight, sixteen and twenty-four hours the pressure in the cell was determined. The pressure in the cell was found to be no greater than twenty microns after a twenty-hour period. This was considered sufficient to consider the cell free of any leaks.

The helium leak detector was used to find any possible small leaks. In this particular test the cell was filled with helium to a pressure of 2,000 psi and all joints were examined with the probe from the helium leak detector. No leakage indications were discovered.

The loading system was also checked for any leaks. This was done in a manner similar to the one used for the test cell. The duration of the tests, however, was only two and four hours. No leakage was found over that period of time as determined by readings taken on a pressure gauge. The vacuum system was tested for possible leaks and the leakage in this part of the system was found to be less than one micron per hour. This was considered satisfactory for the purposes of the experiment.

### Calibration of Equipment

All vessels and lines, with the exception of the large expansion volume, were calibrated using the volume expansion method. Helium, as well as dry air, were used as the calibrating gases except in the free space determination in which only helium was used. The large expansion volume was calibrated by filling it with a measured amount of water. The calibrated volumes are listed in Table II. The effect of pressure and temperature on the cell volume is negligible. Complete calculations may be found in Appendix D.

Several pressure gauges were used in the study in order to cover adequately the range of pressures studied. The two feed pressure gauges, 0-800 psi and 0-2000 psi and the -30 to 300 psig were calibrated against a dead weight tester. The calibration curves may be found in Appendix B. The calibration data was subjected to regression analysis using a program developed for the Bendix G-15 computer after the method of Milne<sup>(47)</sup>. A second order fit represented the data very well.

The pressure transducers, 2000 psi and 500 psi, were calibrated in the system against the pressure gauges. These data were also fitted with a regression curve of second order. The calibrations may be found in Appendix B. Excellent reproducibility was achieved at a given temperature range and although they are compensated for changes in temperature over part of the range, a change in the calibrations occurred over the range of interest. No drift in the calibrations occurred with time and it appeared, therefore, that the slight decrease in battery power had no effect on the readings.

TABLE II  
CALIBRATED VOLUMES

<u>Volume Calibrated</u>	<u>Volume CC.</u>
Constant Volume Reservoir	923 $\pm$ 3.cc
Reservoir Connections	38.2 $\pm$ 1.cc
Central Distribution "T"	8.63 $\pm$ 0.1cc
Cell Sampling Line Section	2.46 $\pm$ 0.05cc
Vacuum Safety Gauge and Connections	12.7 $\pm$ 1.cc
Glass Lines (To top of barometer)	127.7 $\pm$ 5.cc
Barometer	0.196 $\pm$ 0.004cc
Expansion Volume	11,780cc $\pm$ 50.cc
Adsorption Cell-Filled-Free Space	
Runs 1-107	89.7 $\pm$ 0.7cc
Runs 107 -	88.8 $\pm$ 1.0cc

The thermocouple used to determine the cell temperature was checked against a nitrogen-filled pentane thermometer at the ice point, in dry ice, and in liquid nitrogen. At all three temperatures the readings checked within  $0.5^{\circ}\text{C}$ . As this was well within the accuracy of the equipment the temperatures recorded by the thermocouple were assumed to be accurate.

## EXPERIMENTAL METHOD

### Adsorbent

The adsorbent is a calcium substituted Aluminosilicate commercially known as Linde Molecular Sieve 5A. This material is similar to the natural zeolites, chabazite and analcite. Natural materials of this type and related synthetic zeolites were studied by Barrer and co-workers.<sup>(4)</sup>

The sieve used in this study has been studied in some detail by Breck and co-workers of the Linde Company. The basic chemical composition of the hydrated crystal is  $\text{Ca}_6(\text{AlO}_2)_{12}(\text{SiO}_2)_{12} \cdot 3\text{H}_2\text{O}$ .<sup>(12,53)</sup> The basic structure of the unit cell has been described as a framework of alumina ( $\text{AlO}_4$ ) and silica ( $\text{SiO}_4$ ) tetrahedra linked at their apexes, providing a large interior cavity linked to other cavities by eight sided and six sided passages. The edge of the tetrahedra form eight membered oxygen rings with a diameter of 5.0 to 5.6 $\overset{\circ}{\text{A}}$ . In addition there are small six membered rings of approximately 3 $\overset{\circ}{\text{A}}$  in diameter leading from the interior cavity.

The sieve is supplied in the form of 1/16 inch or 1/8 inch pellets and in this study the former were used. These pellets are about 80 per cent actual crystal. The remainder is binder. The sieve has a surface area of 500 square meters per gram (see Appendix C) as measured by the "BET" method and agrees with the data of Breck (12).

The adsorbent will rapidly adsorb water vapor upon exposure to the atmosphere. It is therefore necessary to activate the sieve before any equilibrium determinations are made. The recommended activation conditions are 350°C at one atmosphere or lower temperatures at



lower pressures. The sieve used in these experiments was activated for a minimum of 12 hours at 175°C and at a pressure of  $10^{-2}$  mm.Hg. Between individual equilibrium determinations, the adsorbent was held for 4 hours at 175°C and  $10^{-2}$  mm Hg after nominally complete desorption had been achieved.

In order to determine the effect of regeneration temperature and pressure, the surface area was determined after desorbing samples of sieve at 300°C and  $10^{-6}$  mm Hg. and at 125°C and  $10^{-2}$  mm.Hg for eight hours. There was no significant difference; the specific surfaces of the samples (See Appendix C) were within three per cent of each other which is well within the accuracy of the method. In addition, a sample of the adsorbent which had been regenerated in the cell at the normal desorption conditions, 175°C,  $10^{-2}$  mm.Hg for a minimum of 4 hours, was desorbed for 12 hours at 350°C. The weight difference was 1.45 grams per 100 grams of adsorbent, equivalent to about five per cent of the water capacity of the sieve.

#### Gases

The Methane used in this study was obtained through the courtesy of Phillips Petroleum Company in standard, seventy cubic cylinders at 2000 psi. The gas received was pure grade, 99 plus per cent methane. A number of analyses of the methane were made during the course of the study with the mass spectrometer. A typical analysis can be found in Table III.

The nitrogen used was prepurified grade nitrogen purchased from the Matheson Company in standard 220 cubic foot cylinders at 2000 psi. This gas is sold with a specification of 99.99 per cent nitrogen

TABLE III  
ANALYSES OF PURE FEED GASES

<u>Component</u>	<u>Per Cent</u>	
METHANE		
Methane	98.6	98.7
Nitrogen	1.0	0.9
C <sub>2</sub> <sup>+</sup>	0.2	0.2
CO <sub>2</sub>	0.2	0.2
H <sub>2</sub> O	None Evident*	
NITROGEN		
N <sub>2</sub>	99.9+	
H <sub>2</sub> O	None Evident*	
O <sub>2</sub>	None Evident**	

\* No. 18 peak on mass spectrometer

\*\* No. 32 peak on mass spectrometer

content and a maximum dew point of  $-96^{\circ}\text{C}$ . Each cylinder was sampled and a mass spectrometer analysis made; a typical analysis can be found in Table III.

As both gases were essentially pure and free of water, they were used as received.

#### Experimental Method

Of the two basic experimental methods, dynamic or flow and static or batch, the latter is in many ways the easier to execute and was chosen as the method to obtain the equilibrium data. Besides avoiding some of the experimental difficulties, such as keeping flow rates constant and measuring both flow rates and composition accurately for the entire run, the static method is more economical in regard to material consumption.

After activation of the adsorbent which is discussed above, any number of adsorption-desorption cycles can be made. The gas mixture to be adsorbed is mixed in the constant volume reservoir. The mixture is allowed to come to equilibrium for 24 hours. This was found sufficient to provide a feed mixture of constant composition in the reservoir.

Before any run the adsorption cell is brought to the adsorption temperature and the initial pressure must be less than 0.02 mm. mercury. The pressure and temperature of the reservoir are then determined. The pressure transducer circuit is balanced so that at full vacuum the oscilloscope display shows a balanced bridge circuit.

The main block valve (G-4) and metering valve (H-1) to the vacuum system are closed and the main block valve from the reservoir (G-6)

is opened. The cell is loaded with the gas mixture and the main cell valve (L) is closed. The pressure and temperature of the reservoir are again determined and a sample of the feed gas is taken for analysis.

The system is allowed to come to equilibrium. Temperature is recorded continuously during the run. At the beginning of any run, as the gas is admitted and adsorbed, the system temperature is above the equilibrium temperature. The system temperature, depending upon the conditions, will be as much as 50°C above the equilibrium temperature. Within 15 minutes, the temperature dropped to within five degrees of the equilibrium temperature. The equilibrium temperature is reached within one hour.

Periodic pressure readings are made during the run. The pressure decreases rapidly during the first fifteen minutes. The pressure normally will decrease to the equilibrium pressure within the first hour. In order to insure that equilibrium is achieved, at least eight hours is allowed. A series of tests over a longer period of time indicated that essentially no change in the gas phase occurred after four hours, the time the first sample was taken. After eight hours, two pressure readings taken half an hour apart, at constant temperature, must be identical before the system is said to be at equilibrium. At equilibrium, the cell pressure and temperature are recorded and a gas sample is taken for analysis.

The adsorbate is desorbed into the large, 11,780 cc., expansion chamber and the pressure, temperature, and composition determined. Desorption for purposes of material balance is said to be complete when the equilibrium pressure is less than 1 mm. Hg. It requires between three and six desorption steps to obtain essentially complete desorption. In

order to speed this process, temperature cycling is used with a maximum temperature of 170°C. After the equilibrium pressure is below 1 mm., the system is evacuated further to an equilibrium pressure of 0.01 mm. Hg. at 170°C.

### Sampling

For ease of construction, only one line connects to the adsorption vessel. A small, approximately 2.5 cc. sampling section exists between the main cell valve (L) and the first block valve (G-5). The main distribution system then connects with the sample outlet and the glass sampling tube connection.

In order to obtain a gas phase sample, a 2.5 cc. purge sample is first taken. The line is then evacuated and the sampling section block valve (G-5) is closed. The sample is then obtained by opening the main cell valve (L). It is expanded into the glass sample tube. It is imperative that a true gas sample be obtained. The total sample is about five per cent of the total gas phase. In several tests several samples were taken in rapid succession. As indicated in Table IV, there is no difference in samples taken in rapid order; thus it can be assumed that the sampling technique gives a representative sample of the gas phase. On the other hand, if some time, say 20 minutes, elapses between samples, a new equilibrium is being approached. The rapid sampling is also accompanied by a temporary drop in pressure which increases again to a new equilibrium value with time; this supports the opinion that the gas phase sample is a representative sample of the gas phase only.

The samples were analyzed on a Consolidated Electrodynamics Corporation Model 21-103B Mass Spectrometer. This instrument is said

TABLE IV

PAIRED GAS SAMPLES FROM ADSORPTION CELL TAKEN AS A  
FUNCTION OF TIME TO DETERMINE SAMPLING VALIDITY

Sample Number	Time	Mole Fractions		Deviation from Aver. for CH <sub>4</sub> Analysis
		CH <sub>4</sub>	N <sub>2</sub>	
P-289	8:00 AM	.6512	.3488	- 0.0049
P-290	8:20 AM	.6561	.3439	0.0049
P-318	9:55 AM	.1182	.8818	0.0007
P-319	10:00 AM	.1169	.8831	- 0.0007
P-340	7:20 AM	.5030	.4970	0.0010
P-341	7:25 AM	.5010	.4990	- 0.0010
P-266	2:15 PM	.1459	.8541	0.0000
P-267	2:18 PM	.1459	.8541	0.0000
-	-	-	-	-

TABLE V

REPLICATE ANALYSES OF A SINGLE SAMPLE

Sample No.	Analysis Mole Fraction		% Deviation From Average for CH <sub>4</sub> Analysis
	CH <sub>4</sub>	N <sub>2</sub>	
P-570a	0.1731	0.8269	- 0.34
P-570b	0.1735	0.8265	- 0.12
P-570c	0.1744	0.8256	0.40
Mean	0.1737		
-	-	-	-

TABLE VI

REPLICATE SAMPLES  
FROM A UNIFORM MIXTURE

Sample No.	Analysis Mole Fraction		% Deviation From Average for CH <sub>4</sub> Analysis
	CH <sub>4</sub>	N <sub>2</sub>	
P-428	0.7100	0.2900	- 0.14
P-431	0.7112	0.2888	0.03
P-435	0.7113	0.2887	0.04
P-438	0.7115	0.2885	0.07
Mean	0.7110		
St'd Deviation	= 5.87 x 10 <sup>-4</sup>		
-	-	-	-

to be accurate to less than one-tenth of one per cent but replicate analysis of the same sample, found in Table V, indicate that the precision of 0.2 per cent can be expected with certainty. To determine the accuracy of analysis, a series of replicate samples were taken and analyzed. The results may be found in Table VI. This data also indicates that any given analysis is good to  $\pm 0.002$  mole fraction or 0.2 per cent.

#### Method of Calculating Adsorption Equilibrium

The adsorption equilibrium cannot be directly calculated from the experimental measurements, pressure, temperature, gas volume and composition. The amount actually adsorbed is the difference between the amount of gas admitted to the cell and the gas remaining in the free space of the adsorption cell at equilibrium and may be represented for a given component by:

$$N_i = (N_{I,i} - N_{F,i}) - N_{FS,i} \quad (29)$$

Equation (29) may be rewritten in terms of volume, density and compositions:

$$N_i = (V_I \rho_I - V_F \rho_F) y_{I,i} - V_{FS} \rho_{FS} y_i \quad (30)$$

but the density is some function of pressure, temperature and composition:

$$\rho = f(P, T, y) \quad (31)$$

In the regions of low pressure and high reduced temperatures the ideal gas law will represent the P-V-T data adequately. At high pressures and lower temperatures the gas behavior cannot be adequately represented by the ideal gas law and some correction must be made. Therefore, a more complex equation of state or a compressibility factor must be used. Of

these two methods, the equation of state lends itself to computer solution and was used here. Several equations of state have been developed to represent the experimental P-V-T data. Of these, the most adequate is the Benedict, Webb, Rubin equation, which is an empirical equation developed for hydrocarbons using eight constants<sup>(6,7)</sup> and is of the form:

$$P = RT\rho + C_1\rho^2 + C_2\rho^3 + C_3\rho^6 \quad (32)$$

where

$$C_1 = B_0RT - A_0 - C_0/T^2 \quad (32a)$$

$$C_2 = (bRT - a) + C/T^2 [(1 + \gamma\rho^2)/e^{\gamma\rho^2}] \quad (32b)$$

$$C_3 = a\alpha \quad (32c)$$

The B-W-R equation as developed originally was intended to predict vapor-liquid equilibria of hydrocarbon mixtures. In the original work no constants for nitrogen were published. Bloomer and co-workers fitted a modified B-W-R equation of state for nitrogen.<sup>(10)</sup> The equation added two constants to better represent the data, changing the first and second virial coefficients, so that:

$$C_1 = B_0RT - A_0 - C_0/T^2 - D_0/T^4 \quad (32d)$$

$$C_2 = bRT - a + \left(\frac{C}{T^2} + \frac{\delta}{T^4}\right) [(1 + \gamma\rho^2)/e^{\gamma\rho^2}] \quad (32e)$$

In later work, based on experimental data of Sault<sup>(54)</sup>, Darby,<sup>(17)</sup> Pace,<sup>(44)</sup> and Keyes and Burks<sup>(34)</sup> for nitrogen-methane mixtures, Ellington and co-workers<sup>(19)</sup> recommended the following combinatorial rules for the ten-constant modified B-W-R equation:

$$B_{Om} = \sum_i x_i B_{Oi} \quad (32f1)$$

$$C_{Om} = \left(\sum_i x_i C_{Oi}^{1/2}\right)^2 \quad (32f2)$$



$$D_{om} = \left( \sum_i x_i D_{oi}^{1/2} \right)^2 \quad (32f3)$$

$$a_m = \left( \sum_i x_i a_i^{1/3} \right)^3 \quad (32f4)$$

$$b_m = \left( \sum_i x_i b_i^{1/3} \right)^3 \quad (32f5)$$

$$c_m = \left( \sum_i x_i c_i^{1/3} \right)^3 \quad (32f6)$$

$$\alpha_m = \left( \sum_i x_i \alpha_i^{1/3} \right)^3 \quad (32f7)$$

$$\gamma_m = \left( \sum_i x_i \delta_i^{1/3} \right)^3 \quad (32f8)$$

with a special modification for the constant  $A_o$  :

$$A_{om} = \left( \sum_i x_i A_{oi}^{1/2} \right)^2 - 0.1000 (X_{N_2} \times X_{CH_4}). \quad (32f9)$$

The equation, as modified, represents data with an average deviation under one per cent at pressures up to 1500 psi and temperatures as low as  $-135^\circ\text{C}$ .

As this equation is explicit in pressure, it is necessary to solve for the gas density by an iterative procedure. The Newton-Raphson method suits itself ideally to obtaining a solution rapidly in this case. In addition, the equation may be modified so that the fugacity of a gas may be computed.

Once the density is determined, the amount of each constituent in each phase can be readily determined, using relationship (30). The relative volatility can then be obtained from its definition:

$$\alpha = (y_{N_2}/X_{N_2}) / (y_{CH_4}/X_{CH_4}) \quad (33)$$

As part of the calculations necessary to obtain the equilibrium loading are iterative and lengthy, the entire calculation was programmed for the IBM 704 computer utilizing the MAD translator. The logical flow sheet, together with a sample calculation may be found in Appendix D.

#### Errors in Measurement

Two basic methods for error analysis exist. The first and perhaps most often used is statistical in nature and depends on the basic assumption that the measurements are Gaussian-distributed about the true value. In order to get a good measure of the errors, which are propagated in the results, a number of values of the result are necessary at a given parameter or it is necessary to assume a model for the data.

The second method of evaluating the effect of errors in measurement on the results is to determine the effect of the maximum possible errors in measurement on the calculated results. This method will bound the maximum error and in all but a very few cases the error in the finally calculated result will be measurably less than the maximum indicated from this type of analysis.

The second or maximum propagated error method has been used in analyzing the effect of errors in measurements on the results. In determining equilibrium values it is important that the degree of accuracy of the results be known with some certainty. This could be achieved by obtaining data for a number of runs at a given set of conditions, but it is often uneconomical to carry out this number of experiments. A more realistic approach, therefore, is to make only a single run at most conditions and obtain results over a wider range of the test parameters.

It is then best to make the error analysis by the second method unless one can assume a model which suits itself readily to a statistical regression analysis.

The maximum errors in the various measured variables are listed in Table VII. The maximum errors were determined by checking the values read against primary standard values or secondary standard values. In addition, the errors in reading due to parallax and other factors were taken into consideration.

In order to determine the effect of the maximum errors in measurements, the values used in a number of calculations covering the spectrum of results and basic data were altered. These altered data were then used to compute results using the basic computer program discussed above. The results obtained were then compared with the results obtained using the original data.

The conclusions of the error analysis may be found in Tables VIII and IX. The major effect on the calculated equilibria ensues from inaccuracies in the measurement of the feed pressures. This occurs because a relatively small net pressure is the difference between two rather large pressure values; thus small errors in the initial or final feed cell pressure can materially effect the total cell loading. The calculations indicate that the other errors in measurement propagate errors which are quite reasonable. The errors resulting from inaccuracies in the feed pressure determination should be well below the maximum as both pressures are read with the indicated pressure decreasing and within five minutes of each other. Any reading errors, therefore, should be in the same direction.

TABLE VII

## ACCURACY OF MEASURING DEVICES

<u>Measuring Device</u>	<u>Accuracy</u>	<u>% Error of Max Reading</u>	<u>Min Reading</u>
2000 psi Pressure Gauge (500-2000 psi)	<u>+ 5 psi</u>	0.25	1.0
800 psi Pressure Gauge (100-800 psi)	<u>+ 2 psi</u>	0.25	2.0
2000 psi Pressure Transducer (600-2000 psi)	<u>+ 20 psi</u>	1.0	3.33
500 psi Pressure Transducer (700-100 psi)	<u>+ 5 psi</u>	0.7	5.0
Temperature (Thermocouple) (123°K-300°K)	<u>+ 0.5°K</u>	0.5	0.16
Temperature (Thermometer) (300°K)	<u>+ 0.5°K</u>	0.16	0.16
Mass Spectrometer Composition Analysis (Mole Frac.)	<u>+ 0.002</u>		

TABLE VIII

MAXIMUM ERROR PROPAGATED IN PURE COMPONENT  
EQUILIBRIUM DATA

<u>Measurement</u>	<u>% Changed</u>	<u>Milligram Moles Adsorbed</u>	<u>% Deviation</u>
Original Run (Run 26)		0.266	0
Feed Pressure	0.5	0.258	3
Feed Volume	1	0.265	0.5
Cell Pressure	3	0.260	2
Cell Temperature	0.2	0.267	0.5
Cell Volume	1	0.265	0.5

TABLE IX

MAXIMUM ERROR PROPAGATED IN  
BINARY EQUILIBRIUM DATA

<u>Measurement</u>	<u>% Changed</u>	<u>Milligram Moles CH<sub>4</sub> Adsorbed</u>	<u>% Deviation</u>
Original (Run 108)		0.0348	
Feed Pressure	2	0.0338	3.0
Feed Volume	1	0.0356	1.8
Cell Pressure	4.0	0.0350	0.6
Cell Pressure	4.0	*	0.7
Cell Composition	1.3	0.0349	0.3
Cell Volume	1.0		
Cell Pressure, Composition and Temperature	4.0 1.3 0.16	0.0350	0.57

\* Total adsorption

The amount adsorbed should be within three per cent of the true value under all but the most unusual cases but even in these cases the value will be correct to plus or minus six per cent. The values of relative volatility in all cases are within six per cent of the true value.

## EXPERIMENTAL RESULTS

### Pure Component Adsorption

Pure component isotherms were obtained for methane at three temperatures and over a wide pressure range. Well above the critical temperature of methane at 295°K, nine equilibrium determinations were made up to 50 atmospheres or 760 psia. At 195°K, close to the critical temperature, nine determinations covered a pressure range of about 30 to 450 psia. Below the critical temperature, at 175°K, the equilibrium loading was evaluated at four pressures covering the range of pressures up to the saturation pressure. A tabulation of these results may be found in Table X. As discussed in detail below, the maximum error in any individual determination of the amount adsorbed is plus or minus four per cent. A statistical regression analysis of the data indicates that the isotherms are accurate to plus or minus five per cent 95 per cent of the time at the pressure extremes. The isotherms may be found in Figure 9.

Twenty-one values of pure nitrogen equilibrium loading, or amount adsorbed, were obtained. These may be found in Table XI. The isotherms were determined at 295°K, 195°K and below the critical temperature of nitrogen, at 123°K. At the higher temperatures, data was obtained at pressures up to 1300 psi and 900 psi respectively. Below the critical, the amount adsorbed was determined up to the saturation pressure. The 95 per cent confidence limits on the extremes of the adsorption isotherms, which may be found in Figure 10, are plus or minus five per cent.

The experimental data from which the equilibrium calculations were made may be found in Appendix A.

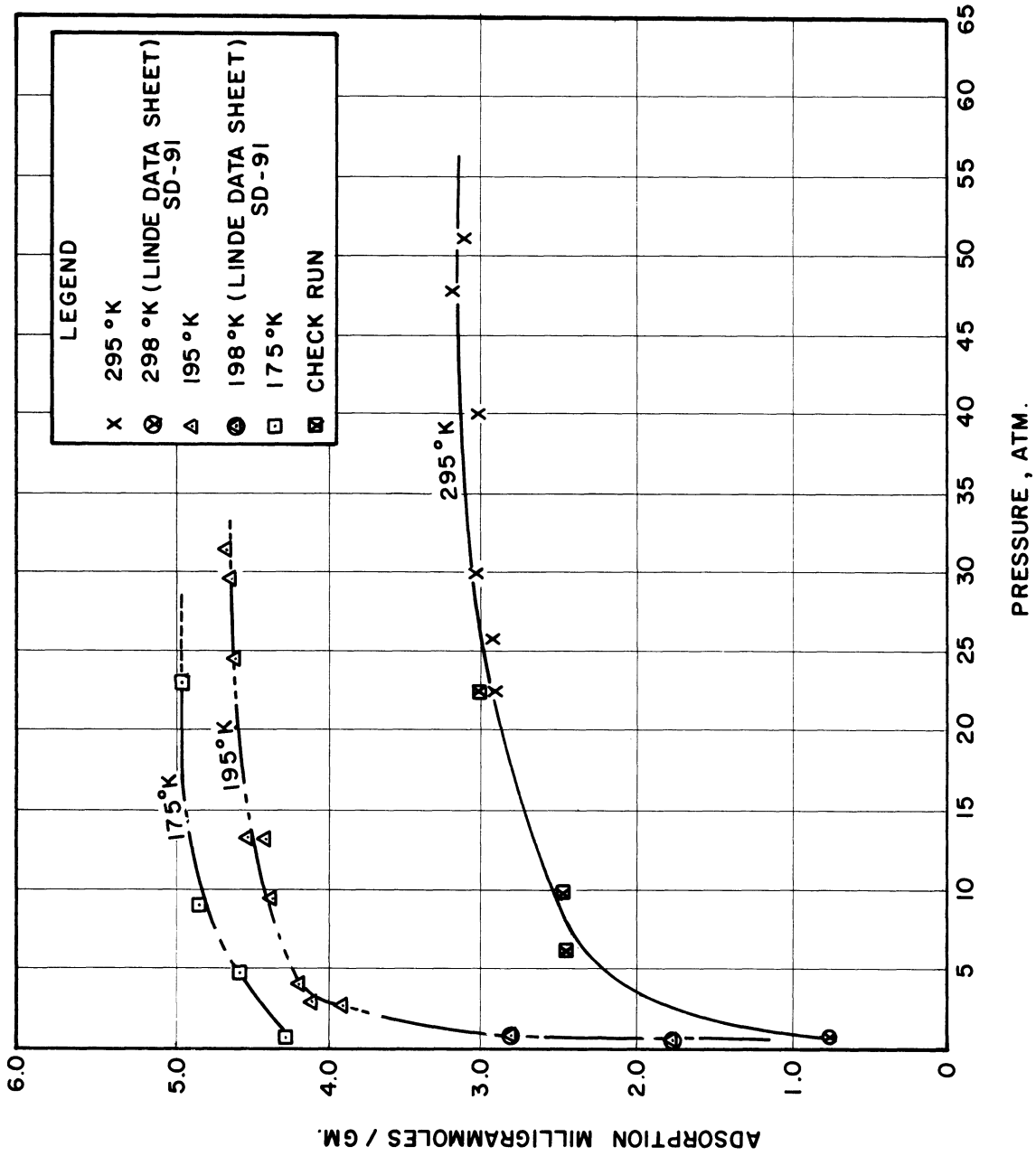


Figure 9. Methane Adsorption Isotherms on 5A Molecular Sieve.



TABLE X

## PURE METHANE ADSORPTION

<u>Run No.</u>	<u>Pressure ATM.</u>	<u>Temperature °K</u>	<u>Amount Adsorbed Milligram Moles/GM</u>
Average Temperature: 295°K			
115	51.3	297.0	3.12
25-1	47.63	293.0	3.24
26	40.15	297.0	3.01
27	29.94	296.0	3.05
28	25.86	295.5	2.91
122	22.5	296.0	3.02
25-2	22.46	293.0	2.94
124	9.78	296.5	2.46
123	6.40	296.0	2.52
Average Temperature: 195°K			
114	31.30	195.5	4.70
52	29.70	196.0	4.61
56	24.84	194.0	4.56
58	13.13	195.5	4.39
92	9.25	193.0	4.34
53	9.08	195.5	4.66
57	3.54	195.5	4.17
93	2.93	193.0	4.12
59	2.72	195.0	3.92
Average Temperature: 175°K			
78	22.52	175.5	4.95
79-1	8.98	174.5	4.83
79-2	4.87	174.0	4.62
80	0.50	175.0	4.25

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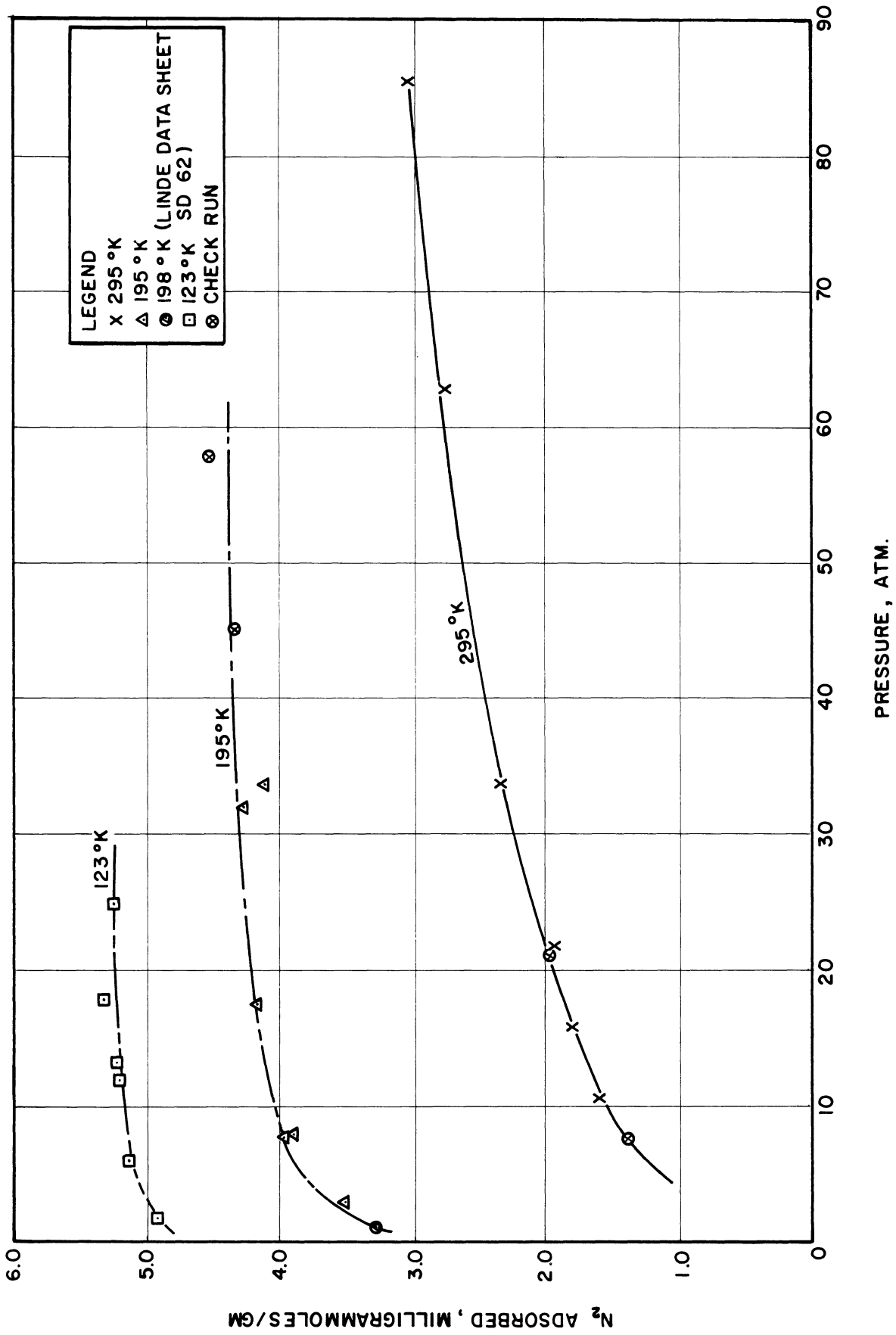


Figure 10. Nitrogen Adsorption Isotherms on 5A Molecular Sieve.

TABLE XI

PURE NITROGEN ADSORPTION

<u>Run No.</u>	<u>Pressure ATM.</u>	<u>Temperature °K</u>	<u>Amount Adsorbed Milligram Moles/GM</u>
Average Temperature: 295°K			
29	85.4	294.0	3.06
30	62.9	295.5	2.74
31	33.68	295.0	2.32
32	21.57	295.5	1.92
120	21.3	295.5	1.95
33	15.65	296.0	1.80
103	10.55	296.0	1.66
121	6.78	295.0	1.39
Average Temperature: 195°K			
112	57.84	195.5	4.54
113	44.91	195.5	4.35
71	33.48	195.0	4.12
51	31.98	195.0	4.31
72	17.56	194.0	4.16
73	7.72	194.0	3.90
74	2.86	194.0	3.54
Average Temperature: 123°K			
84	24.82	123.0	5.24
85	17.60	123.5	5.30
87	13.34	123.0	5.23
86	11.64	123.0	5.19
88	5.78	123.0	5.14
89	1.59	123.0	4.93

- - - - -

Lewis<sup>(27)</sup> reported very limited data on methane adsorption on Davison Silica Gel and Columbia G Activated Carbon. The data indicate that the Molecular Sieve has a greater capacity for methane. Antropff<sup>(1)</sup> studied the adsorption of nitrogen on a Bayer-Werk Leverkusen activated charcoal. The data show that his particular adsorbent has a slightly higher capacity for nitrogen; for example, at 0°C and 27 atmospheres it adsorbs 2.7 milligram moles per gram as opposed to the 2.2 milligram moles per gram capacity of the Molecular Sieve at 22°C and 27 atmospheres pressure.

#### Binary Adsorption

Adsorption equilibria for mixtures of methane and nitrogen were evaluated at the four temperatures at which pure isotherms were determined, 295°K, 195°K, 175°K and 123°K. At the higher two temperatures, equilibrium adsorption was determined at pressures as high as 1350 psi or 90 atmospheres and gas phase compositions ranging between ten and ninety per cent. Results for the sixty runs may be found in Tables XII and XIII. Just below the critical temperature of pure methane, at 175°K, the equilibrium loading increased from 4.75 to 5.04 milligram moles per gram of adsorbent with an increase of pressure from 7.6 to 30 atmospheres. The equilibrium amounts of each component adsorbed may be found in Table XIV.

At 123°K, near the critical temperature of pure nitrogen, the quantity adsorbed was determined up to 300 psia for an equilibrium gas mixture containing five per cent methane.

The maximum pressure investigated was limited at the higher two temperatures by the cylinder pressure without any additional compression. At the lower two temperatures the dew point pressure was the upper pressure limit.

TABLE XII

MIXED ADSORPTION AT  
295°K AVERAGE TEMPERATURE

Run No.	Pressure ATM.	Mole Fraction CH <sub>4</sub> Gas Phase	Amount Adsorbed Milligram Moles/GM			Relative Volatility
			N <sub>T</sub>	N <sub>CH<sub>4</sub></sub>	N <sub>N<sub>2</sub></sub>	
13-1	89.82	0.793	3.27	2.61	0.68	2.2
13-2	89.82	0.802	3.09	2.48	0.61	2.31
13-3	83.02	0.787	3.30	2.60	0.70	2.08
13-6	57.16	0.790	3.33	2.63	0.70	2.16
13-9	42.19	0.793	3.22	2.55	0.67	2.14
13-11	30.28	0.818	3.19	2.61	0.58	2.46
14-1	76.21	0.639	3.45	2.66	0.79	1.91
14-2	74.85	0.633	3.45	2.68	0.77	2.02
15-1	28.58	0.593	2.70	2.04	0.66	2.12
15-2	27.90	0.544	2.70	2.10	0.60	2.94
16-2	14.29	0.590	2.36	1.74	0.62	1.94
17-2	90.84	0.161	3.08	0.87	2.21	2.06
17-4	61.92	0.161	3.16	0.88	2.28	2.02
17-5	59.98	0.168	3.18	0.87	2.31	1.87
18	47.29	0.155	2.86	0.74	2.12	1.91
19	22.80	0.146	2.05	0.53	1.53	2.03
20	74.85	0.705	3.51	2.94	0.57	2.15
21	56.14	0.700	3.00	2.50	0.50	2.14
22	34.02	0.675	2.83	2.29	0.54	2.05
23	24.50	0.667	2.38	1.94	0.44	2.19
24	15.31	0.656	2.02	1.62	0.40	2.15
104	30.62	0.271	2.86	1.28	1.58	2.16
105	21.71	0.274	2.56	1.10	1.46	2.00
106	9.12	0.270	1.89	1.78	1.11	1.90
107	21.07	0.150	2.41	0.60	1.81	1.89
108	8.39	0.149	1.75	0.39	1.36	1.66
109	35.40	0.171	2.90	0.86	2.04	2.03
110	19.64	0.173	2.31	0.67	1.64	1.94
111	12.83	0.174	2.13	0.60	1.53	1.85

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TABLE XIII

MIXED ADSORPTION AT  
195°K AVERAGE TEMPERATURE

Run No.	Pressure ATM.	Mole Fraction $\text{CH}_4$ Gas Phase	Amount Adsorbed			Relative Volatility
			Milligram $\text{N}_2$	Moles/GM $\text{N}_{\text{CH}_4}$	Moles/GM $\text{N}_{\text{N}_2}$	
34-1	84.04	0.150	4.44	1.55	2.89	3.03
34-2	81.65	0.153	4.54	1.55	2.99	2.86
35	46.61	0.120	4.50	1.40	3.10	3.30
36	28.92	0.110	4.34	1.26	3.08	3.29
37	16.47	0.112	4.18	1.18	3.00	2.93
39	47.63	0.606	4.29	3.53	0.76	3.04
40	28.58	0.556	4.25	3.45	0.80	3.44
41	11.36	0.503	3.50	2.64	0.85	3.36
43	37.43	0.405	4.51	3.08	1.43	3.20
44	16.33	0.406	4.28	2.98	1.30	3.40
45	76.89	0.476	4.80	3.20	1.60	2.21
46	36.74	0.611	4.30	2.78	1.52	2.88
48	18.17	0.451	4.18	2.98	1.20	3.02
61	6.51	0.314	4.10	2.22	1.88	2.58
62	1.97	0.334	3.75	1.97	1.78	2.21
64	20.48	0.121	4.23	1.18	3.05	2.81
65	7.59	0.126	4.07	1.03	3.04	2.36
66	1.38	0.133	3.42	0.83	2.59	2.69
67	31.20	0.546	4.75	3.72	1.03	2.71
68	16.40	0.531	4.96	3.74	1.22	2.70
69	7.96	0.508	4.36	3.21	1.15	2.69
70	0.95	0.542	3.54	2.53	1.01	2.34
94	33.4	0.201	4.33	1.89	2.44	3.05
95	13.27	0.176	4.19	1.59	2.60	2.87
96	4.76	0.188	3.92	1.40	2.52	2.42
97	28.58	0.357	4.63	2.91	1.72	3.05
98	19.53	0.334	4.44	2.67	1.77	3.03
99	6.98	0.326	4.13	2.30	1.83	2.74
100	24.36	0.687	4.60	3.96	0.64	2.76
101	13.98	0.658	4.55	3.86	0.69	2.92
102	4.49	0.700	4.21	3.46	0.75	1.95

TABLE XIV

MIXED ADSORPTION AT  
175°K AVERAGE TEMPERATURE

Run No.	Pressure ATM.	Mole Fraction CH <sub>4</sub> Gas Phase	Amount Adsorbed Milligram Moles/GM			Relative Volatility
			N <sub>T</sub>	N <sub>CH<sub>4</sub></sub>	N <sub>N<sub>2</sub></sub>	
75	30.55	0.807	5.04	4.07	0.97	3.04
76	7.60	0.754	4.75	3.59	1.16	3.18
118	17.08	0.786	4.90	3.85	1.05	3.10

TABLE XV

MIXED ADSORPTION AT  
123°K AVERAGE TEMPERATURE

Run No.	Pressure ATM.	Mole Fraction CH <sub>4</sub> Gas Phase	Amount Adsorbed Milligram Moles/GM			Relative Volatility
			N <sub>T</sub>	N <sub>CH<sub>4</sub></sub>	N <sub>N<sub>2</sub></sub>	
81	11.80	0.050	5.03	1.07	3.96	5.20
82	1.27	0.050	5.02	0.92	4.10	4.30
90	18.98	0.051	5.16	1.28	3.88	6.10
91	5.99	0.047	5.07	1.03	4.04	5.02

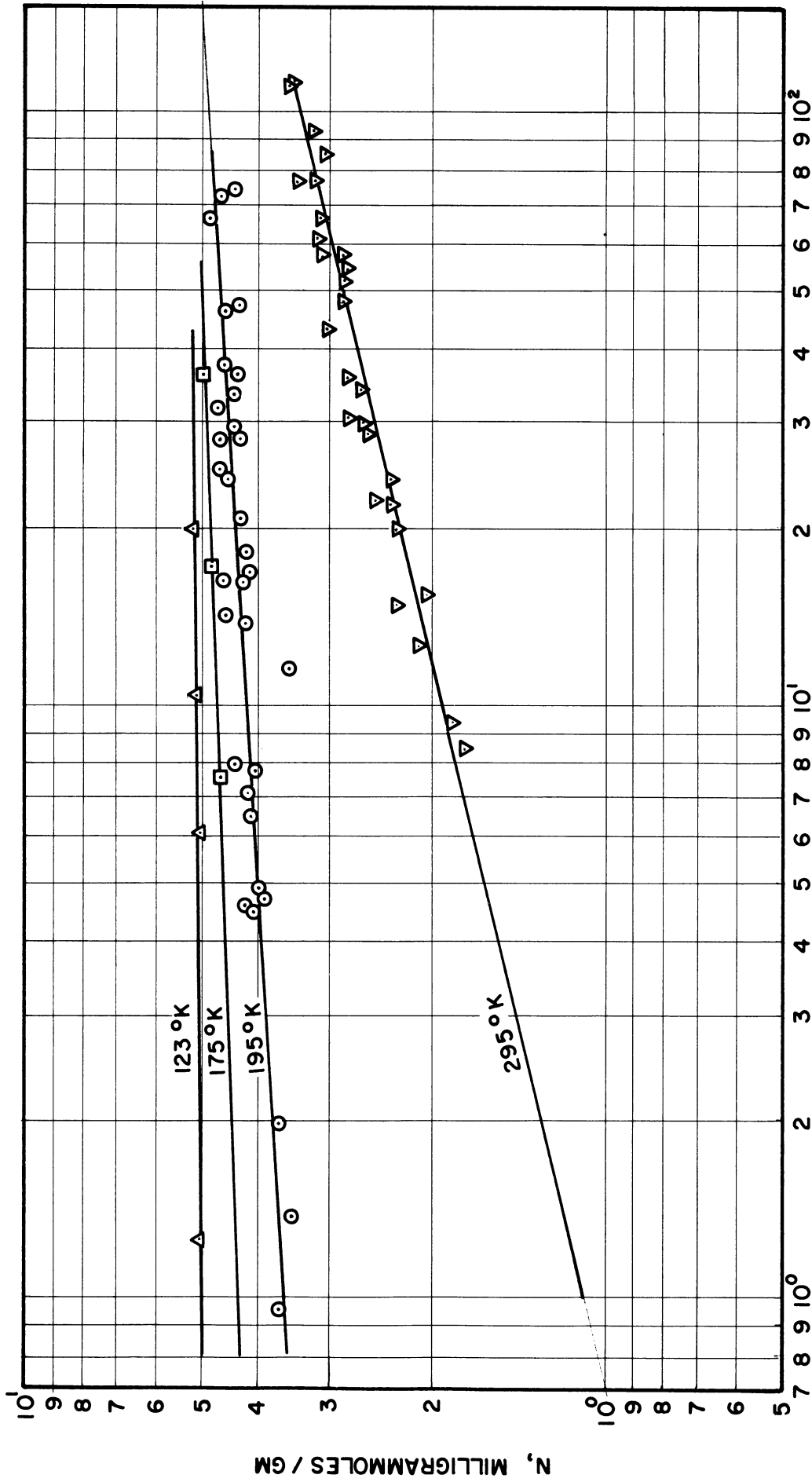
The individual determinations of the binary adsorption equilibria have a maximum error of plus or minus six per cent. It is apparent from Figure 11, which shows the total adsorption for the binary system, that the total equilibrium amount adsorbed is independent of composition. The 95 per cent confidence of these isotherms at their extreme is seven per cent.

TABLE XVI

CONFIDENCE LIMITS ON EQUILIBRIUM VALUES

	95 Per Cent Confidence on Isotherms at Pressure Extremes
Methane	5%
Nitrogen	5%
Methane-Nitrogen Mixtures	7%





PRESSURE, ATM.

Figure 11. Mixed Adsorption as a Function of Pressure and Temperature.

## DISCUSSION OF RESULTS

### Pure Component Adsorption Correlations

A number of methods of correlating the data were tried in order to extend the usefulness of the data over the entire range of interest.

For pure component adsorption, both methane and nitrogen, the data was correlated using both the Freundlich

$$N = KP^{1/n} \quad (3)$$

and Langmuir

$$N = \frac{aP}{1 + bP} \quad (1a)$$

relationships as previously discussed. The constants for the correlations were obtained by the method of least-squares which is well detailed in standard references on numerical methods such as Milne.<sup>(47)</sup>

A plot of the experimental results for the pure components after Freundlich may be found in Figure 12. A Langmuir plot of the results has not been included but this would be rather easily constructed by plotting "P/N" in atmospheres per milligram mole per gram versus "P" in atmospheres. The constants, valid between 100 and 300° Kelvin, for the Freundlich equation are shown as a function of temperature in Figure 13, while those necessary to complete the Langmuir relationship may be found in Figure 14. Analytical relations for the Langmuir constants as a function of temperature will also be found with the graphical representation. The dependence of the constants on temperature is quite similar to that found by Fowler and Guggenheim<sup>(20)</sup> in their statistical development of the Langmuir equation.

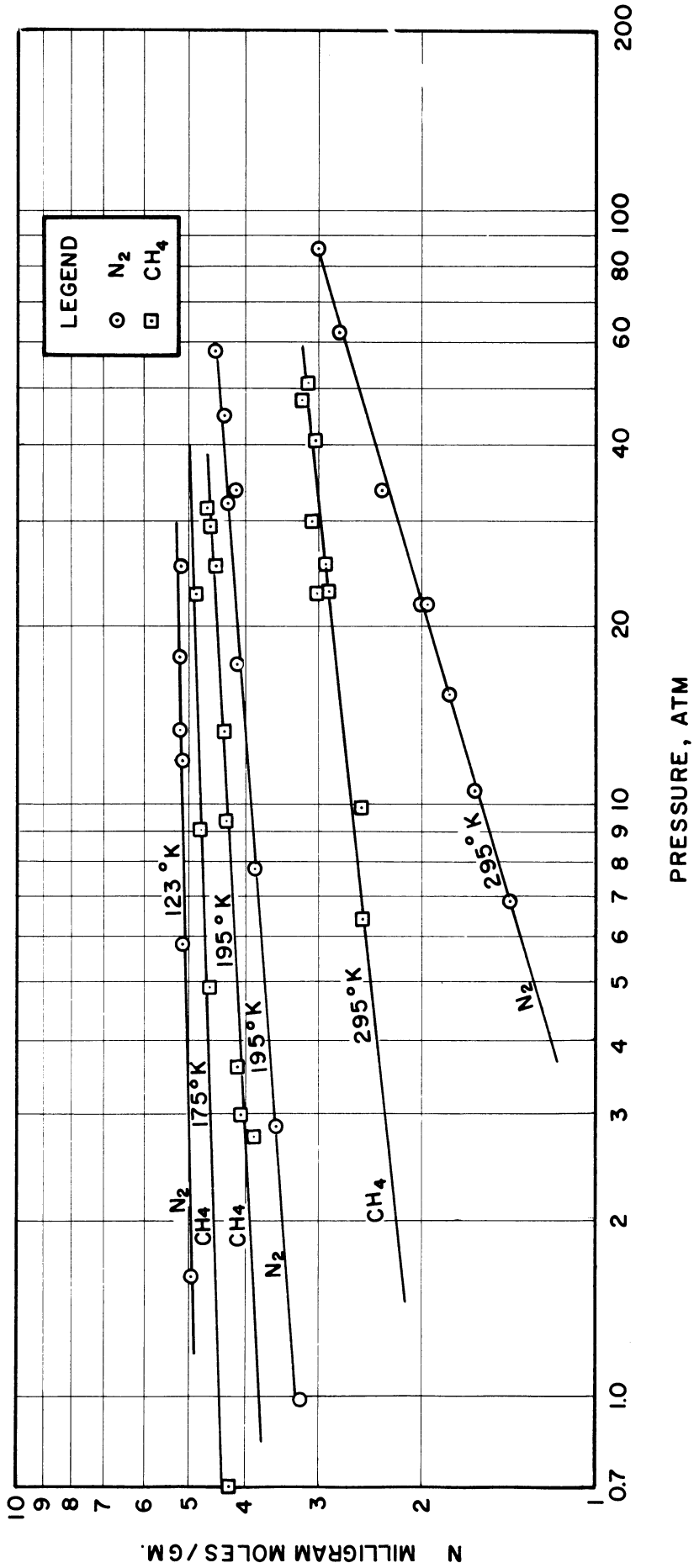


Figure 12. Pure Component Adsorption on 5A Molecular Sieve.

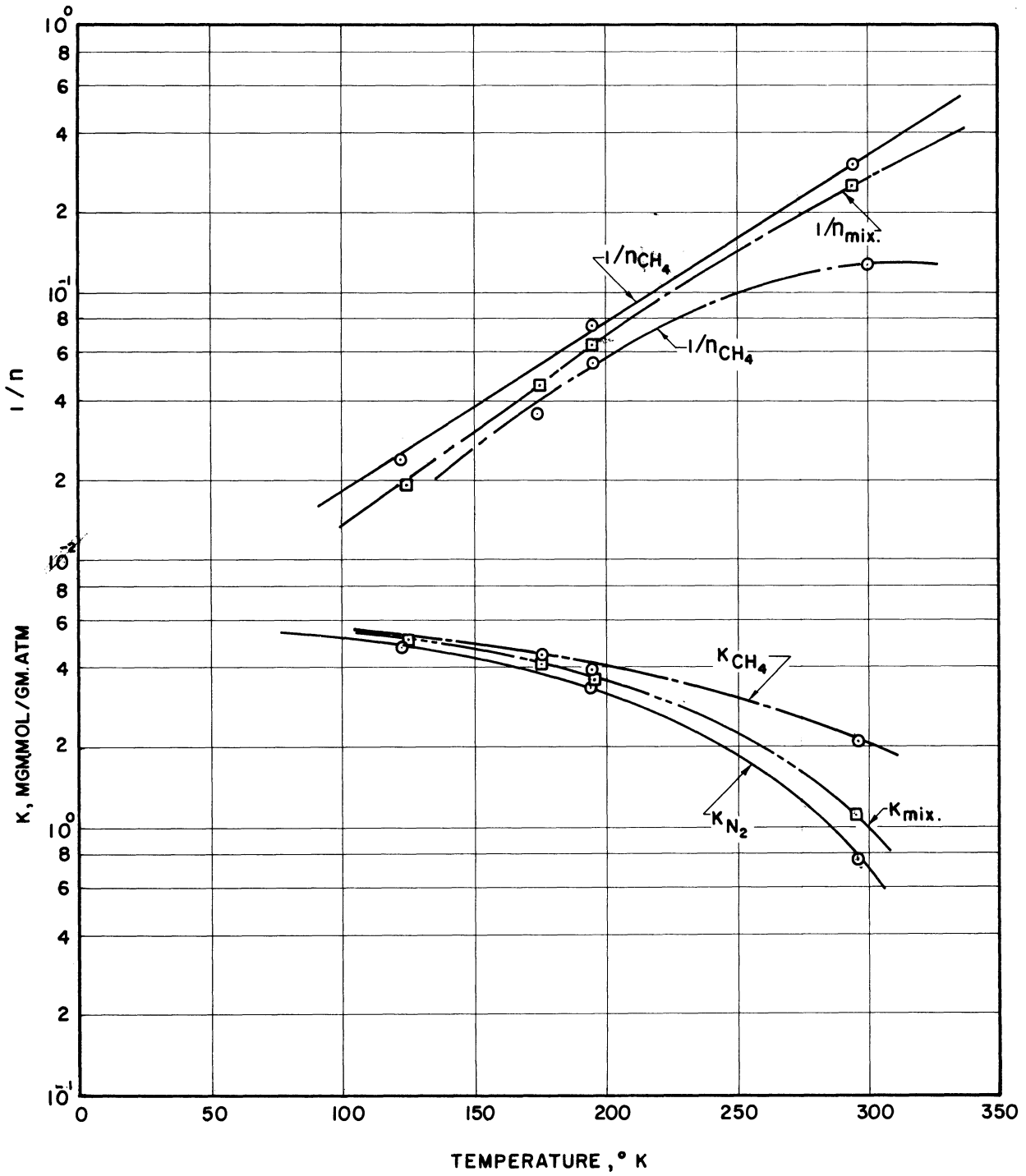


Figure 13. Constants for the Freundlich Equation for Nitrogen and Methane and Their Mixtures (P in ATM.)

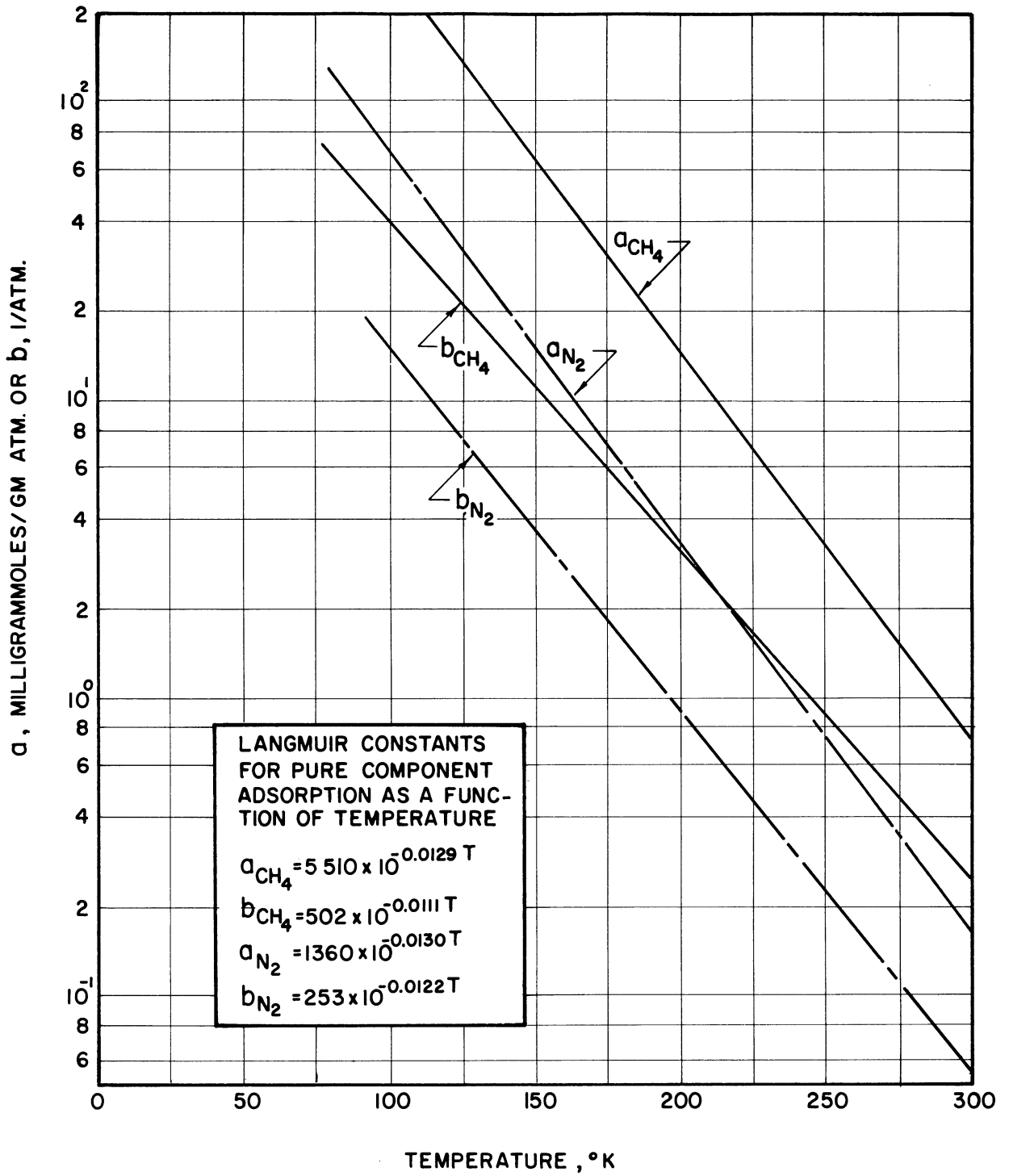


Figure 14. Langmuir Constants for Pure Component Adsorption.

Both relations fit the experimental data rather well with an average correlation coefficient for the logarithmic form of the Freundlich relationship equal to 0.96 for nitrogen and 0.92 for methane. The standard error in the logarithm of the amount adsorbed is less than 0.16 milligram moles per gram (95 per cent confidence) when  $N$  is equal to 4 milligram moles per gram. In the case of the Langmuir relationship, the average correlation coefficient is 0.99 plus and a standard error is less than 0.05 which represents a 95 per cent confidence limit of  $\pm 0.1$  in  $N$ , the amount adsorbed.

That both relationships represent the data well is not surprising. Trapnell<sup>(60)</sup> indicated that, with the proper choice of constants, either the Langmuir or the Freundlich relationship will represent adsorption data in the middle pressure range adequately.

In addition to the Langmuir and Freundlich correlations, two correlations based on the Polanyi Potential Theory were attempted. The first is the one first proposed by Lewis and co-workers<sup>(37)</sup> where  $N/\rho_S$ , the volume adsorbed, is plotted as a function of  $T_{p_S} \ln f_S/f$ , a measure of the adsorption potential. Here,  $f_S$ , is the fugacity at the saturation pressure, as determined from an extrapolated Clausius-Clapeyron vapor pressure equation, calculated with the help of the Benedict-Webb-Rubin fugacity relationship. The gas phase fugacity,  $f$ , was also evaluated using the B-W-R relationships. The density,  $\rho_S$ , which is the important correlating factor, was evaluated at the temperature equal to the temperature of saturation at the equilibrium pressure. The actual density data for methane is due to Bloomer and Parent<sup>(10)</sup> and for nitrogen is due to Bloomer and Rao.<sup>(11)</sup>

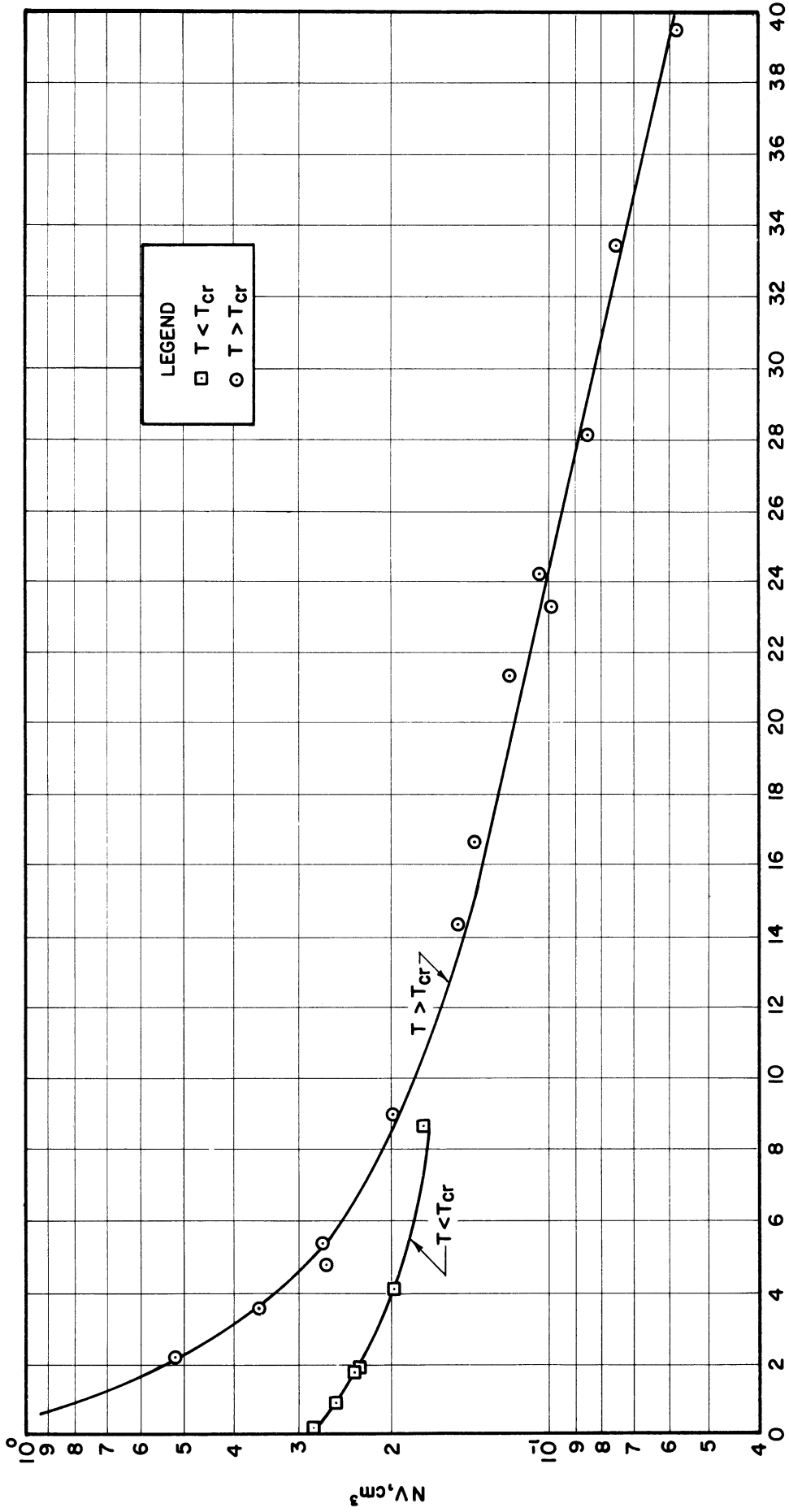


Figure 15. Volume of Nitrogen Adsorbed as a Function of the Adsorption Driving Force.

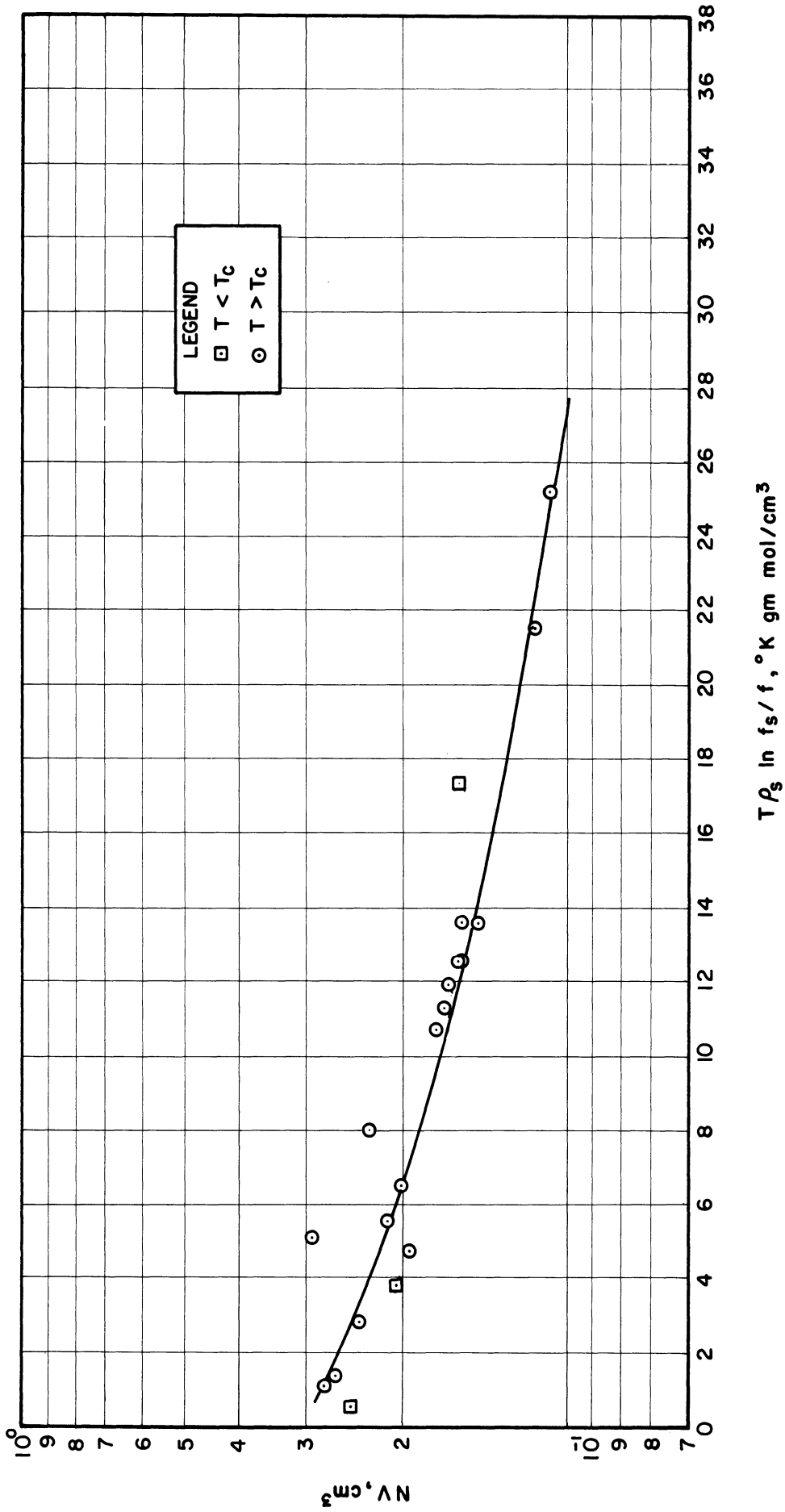


Figure 16. Volume of Methane Adsorbed as a Function of the Adsorption Driving Force.



The correlation for pure nitrogen, as shown in Figure 15, indicates that a good correlation is obtained above the critical point independent of temperature; however, a different correlation appears to exist of adsorption below the critical temperature. This difference could easily be due to an incorrect choice for the density correlating variable. The methane correlation, Figure 16, by this method appears to be quite good for values both above and slightly below the critical with the exception of two points which lie noticeably above the correlating line. These points are obtained from equilibrium data obtained at 40 and 50 atmospheres where the present method of evaluating the adsorbed state density may not predict this density correctly. This same observation can be made about the single point lying well above the nitrogen correlating line.

#### Correlations for Total Adsorption from a Mixture

A correlation similar to the one detailed in the previous paragraph was attempted for the total amount adsorbed from mixtures of methane and nitrogen. The density in this case was assumed to be the molal average density of the two components in the adsorbed state, using the densities of pure components at their respective saturation temperatures as a basis for obtaining the average. It can readily be seen from Figure 17 that no single correlation could be achieved by this method. The correct choice of density again seems to be the major problem.

The total amount adsorbed from a mixture of nitrogen and methane can also be predicted by either a Freundlich or Langmuir type relationship. The Freundlich equation appears to represent the data better; the necessary constants may be found as a function of temperature in Figure 13. The

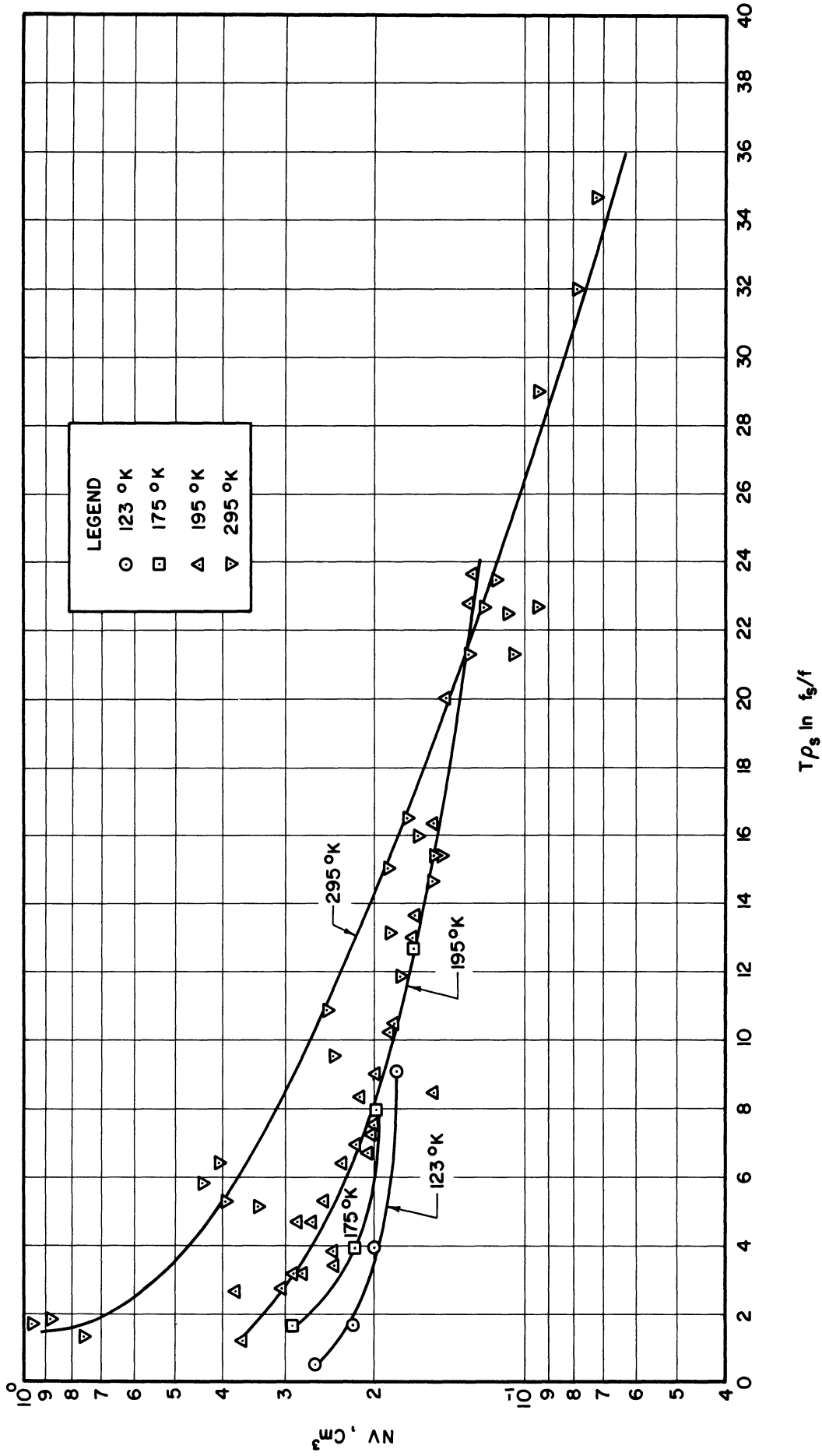


Figure 17. Volume of Nitrogen-Methane Mixtures Adsorbed as a Function of the Adsorption Driving Force.

plots of the total amount adsorbed as a function of pressure and temperature, as found in Figure 11, show that the total adsorption is independent of the composition of the binary gas in equilibrium with the adsorbate. This is not entirely unexpected at the higher temperatures where both gases are well above their critical temperatures. Near the critical point of methane, this behavior is still present and can best be explained by assuming that both gases compete for a given amount of adsorption space which limits the number of molecules adsorbed under a given set of conditions. This assumption would then follow for lower temperatures. More data over a wider range of pressures and compositions at these temperatures would be helpful in confirming this theory.

Another correlating method based on the work of Berényi and Polanyi is that advanced by Dubinin and co-workers. This method has been previously discussed in some detail. Although this method appears to be fundamentally sound, no correlation was possible using Dubinin's methods of evaluating adsorbed phase density as a function of the system temperature, normal boiling point density and van der Waals' density.

#### Adsorption of Components from a Mixture

Methods for predicting the adsorption of one component from a mixture have not received as much attention as pure component adsorption. The basic work on this area is that of Markham and Benton, which has been previously discussed, and which was examined by Schay. A correlation by these methods proved quite successful. The constants for the equation

$$N_{N_2} = \frac{a_{N_2} P_{N_2}}{1 + b_{N_2} P_{N_2} + b_{CH_4} P_{CH_4}} \quad (23a)$$

and

$$N_{\text{CH}_4} = \frac{a_{\text{CH}_4}^i P_{\text{CH}_4}}{1 + b_{\text{N}_2}'' P_{\text{N}_2} + b_{\text{CH}_4}^i P_{\text{CH}_4}} \quad (23b)$$

are plotted as a function of temperature in Figure 18. It is of some interest that the "a" and "b" constants take on different values from those found in pure component adsorption. This indicates, in conformity with the ideas advanced by Schay, that the adsorbate is not an ideal solution and that certain interactions take place in the adsorbed phase. The constants were obtained using the standard least squares technique and represent the data within plus or minus seven per cent. No simple interaction coefficients, as suggested by Schay were obtained.

A correlation to predict individual component adsorption from a mixture using a Freundlich type relationship

$$\log N_1 = \log K + 1/n_1 \log y_1 P + 1/n_2 \log y_1 \quad (34)$$

was attempted. Although at given temperatures a fair degree of correlation could be obtained, it was not possible to obtain a good correlation of the coefficients as a function of temperature. This may in part be due to the lack of data at intermediate temperatures and a better picture might be obtained with a greater amount of data below the critical temperature. However, the nature of the equation makes it difficult to deduce any fundamental relationship between pure component coefficients and mixed coefficients making it difficult to predict the behavior of the equation constants when mixtures are considered.

From the observation of Lewis and co-workers and the theoretical proof of Schay that

$$\sum \frac{N_i}{N_i^0} = 1 \quad (35)$$

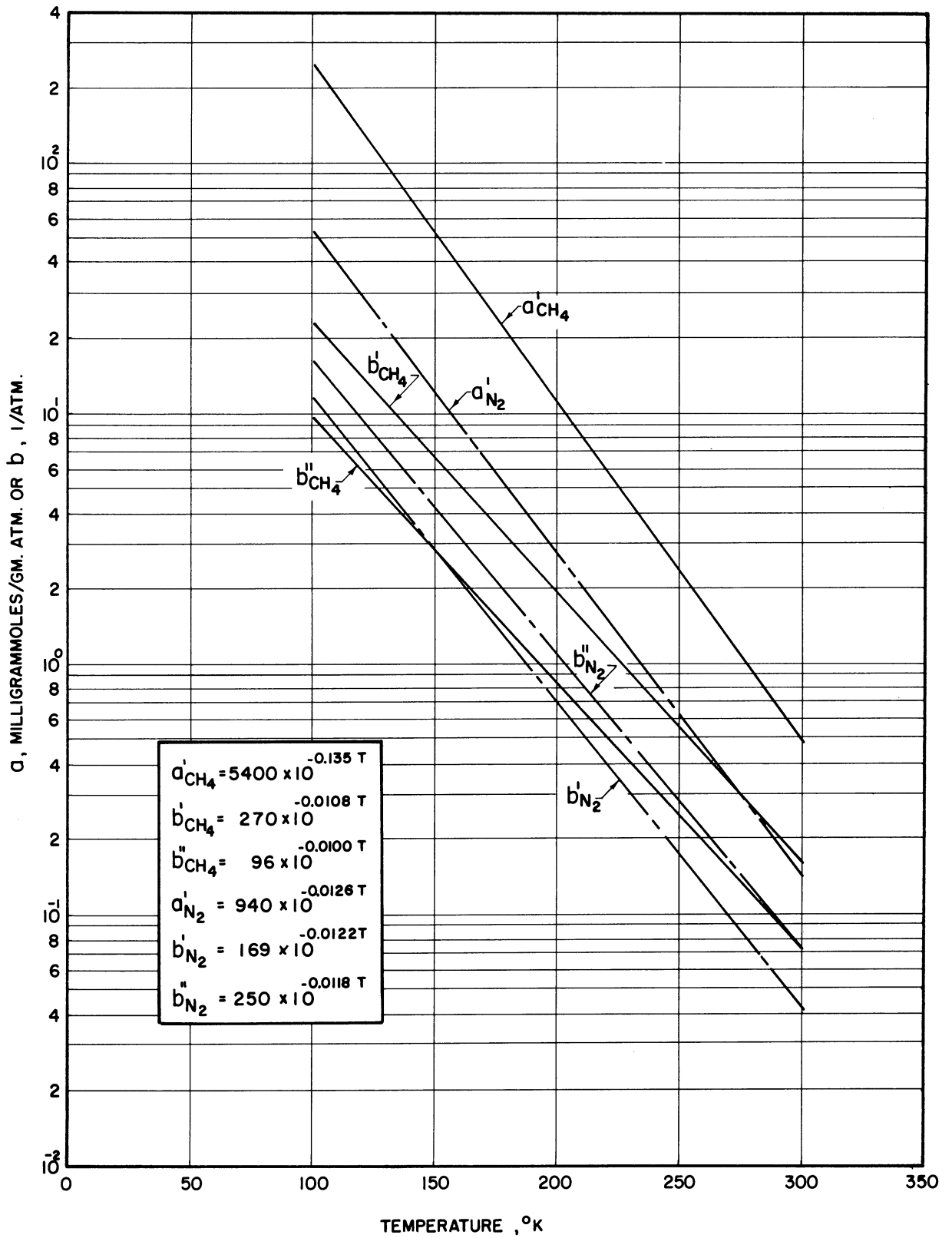


Figure 18. Langmuir Constants for Adsorption From Nitrogen-Methane Mixtures.

for ideal mixtures, a plot of  $N_{N_2}/N_{N_2}^i$  as a function of  $N_{CH_4}/N_{CH_4}^i$  should result in a straight line through the values 0, 1 and 1, 0. If the mixture is non-ideal then the data will lie on one side of the line, which is the case here, as seen in Figure 19.

The relative volatility, as shown in Figure 20, is independent of composition. At the higher pressures above the critical temperature, pressure appears to have little or no effect on the relative volatility. Below the critical point of the mixture, however, as is indicated by the values at 123°K, there is a marked increase in relative volatility with increased pressure. The relative volatility, together with the total loading, will yield the equilibrium loading from a mixture.

#### Computing Adsorption Loading from the Correlations

It is possible to calculate the amount of each constituent adsorbed from a mixture knowing the temperature, the total pressure and the composition of the equilibrium gas.

As outlined in the sample calculation below, the total amount adsorbed from a mixture can be determined by the Freundlich equation with the appropriate constants, as found in Figure 12. The relative volatility, obtained from Figure 20, is independent of composition over the range of pressures of interest. From the gas phase data and the values of the relative volatility and total loading, the amount of each component may now be computed.

These above calculations may be checked by determining the amount of pure component adsorbed at the system temperature and a pressure equal to its partial pressure in the mixture. These values, together with

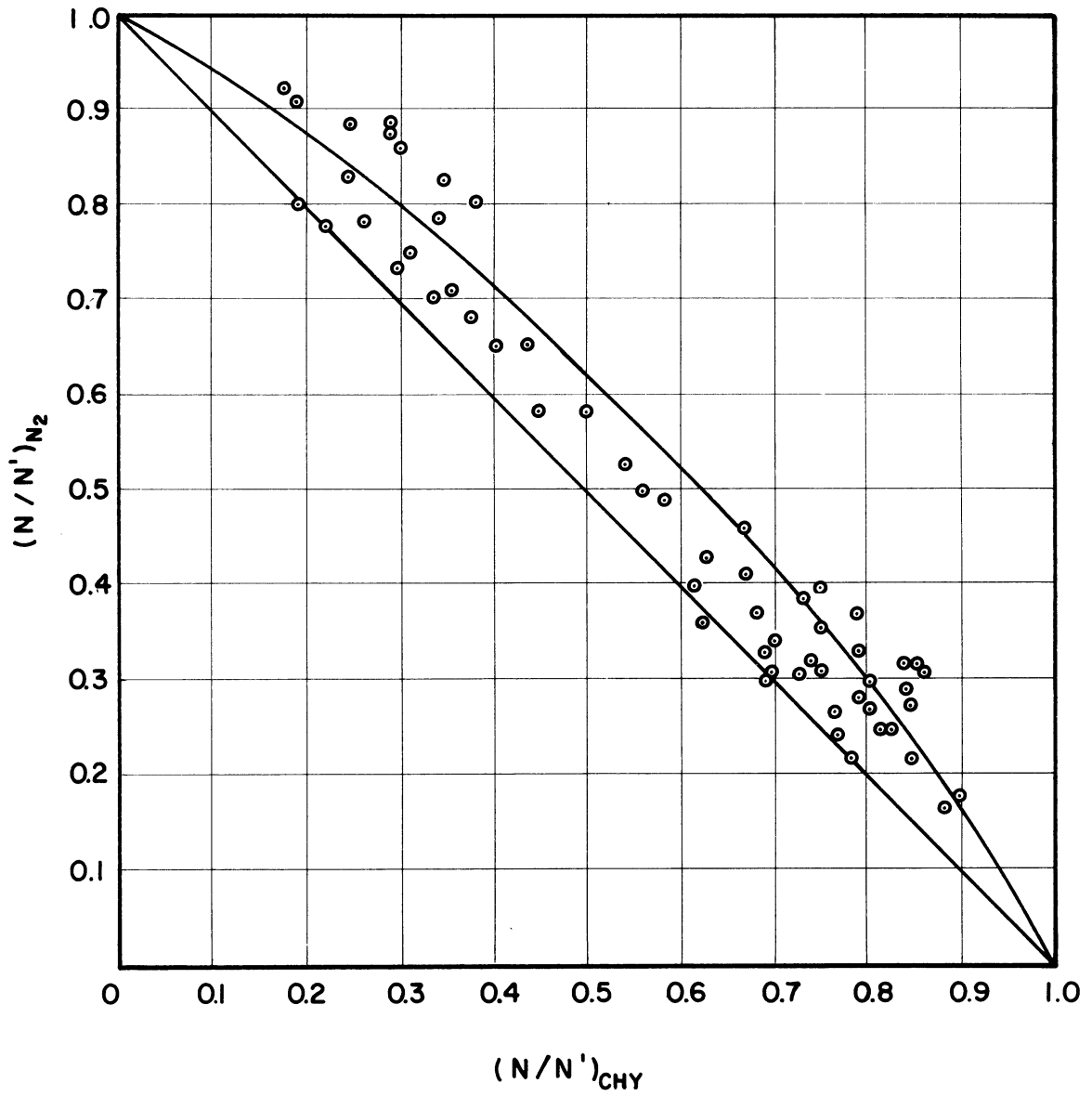
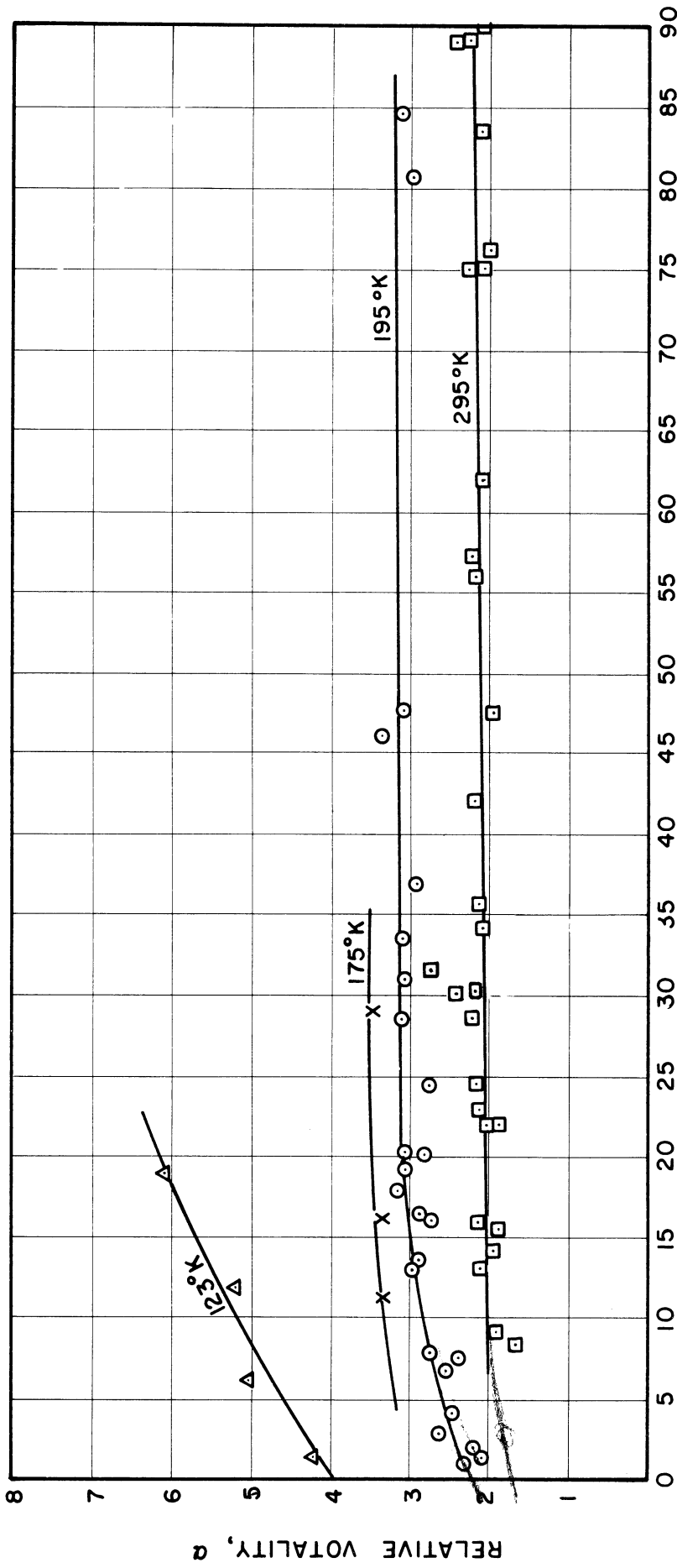


Figure 19. Total Adsorption Correlation on Molecular Sieve 5A.



RELATIVE VOLATILITY, α

Figure 20. Relative Volatility as a Function of Pressure.  
 $(\alpha = (y_{N_2}/X_{N_2}) / (y_{CH_4}/X_{CH_4}))$ .

195°K

295°K

175°K

123°K

PRESSURE, ATM.



the amount of component adsorbed, may be used to evaluate the dimensionless factor  $N/N'$ . The values of this dimensionless factor for methane and nitrogen should satisfy the restriction placed upon them in Figure 19, a plot of the sum of the values of  $(N/N')$  for methane and nitrogen.

The individual component Langmuir equations will also give individual component capacities from a mixture.

Sample Calculation of Binary Adsorption

$$P = 15 \text{ atm.}$$

$$y_{\text{CH}_4} = 0.6$$

$$P_{\text{CH}_4} = 9 \text{ atm.}$$

$$P_{\text{N}_2} = 6 \text{ atm.}$$

From Figure 20.

$$\alpha = 2.0 = (0.4/x_{\text{N}_2}) / (0.6/(1-x_{\text{N}_2}))$$

then

$$x_{\text{N}_2} = 0.25$$

From Figure 11.

$$N_T = 2.2 \text{ milligram moles/gm.}$$

then

$$\begin{aligned} N_{\text{N}_2} &= 2.2 \times 0.25 \\ &= 0.55 \text{ milligram moles/gm.} \end{aligned}$$

and

$$N_{\text{CH}_4} = 1.65 \text{ milligram moles/gm.}$$

This can now be checked by using Figure 19.

From Figure 12.

$$N_{N_2} = 1.35 \text{ milligram moles/gm at 6 atm.}$$

$$N_{CH_4} = 2.6 \text{ milligram moles/gm at 9 atm.}$$

therefore

$$(N/N')_{N_2} = 0.41$$

$$(N/N')_{CH_4} = 0.64$$

and

$$\frac{(N/N')_{N_2} + (N/N')_{CH_4}}{\quad}$$

satisfies the relationship illustrated in Figure 12.

The Langmuir relationships predict

$$N_{N_2} = 0.59 \text{ milligram moles/gm.}$$

$$N_{CH_4} = 1.71 \text{ milligram moles/gm.}$$

One or more points were checked at each temperature using both methods and they yielded results that were good to plus or minus ten per cent as indicated by the three examples tabulated below.

TABLE XVII  
ACTUAL AND PREDICTED BINARY ADSORPTION

Run	Experimental					Predicted			
	P atm.	T °K	$Y_{CH_4}$	$N_{CH_4}$ mgm.	$N_{N_2}$ moles/gm.	Relative Volatility Method		Langmuir Equation	
						$N_{CH_4}$	$N_{N_2}$	$N_{CH_4}$	$N_{N_2}$
16	14.29	295	0.590	1.74	0.62	1.65	0.55	1.71	0.59
64	20.48	195	0.121	1.18	3.05	1.30	3.05	1.31	3.00
81	11.80	123	0.050	1.07	3.96	1.12	3.98	0.72	4.02

At the lowest temperature, 123°K, the Langmuir equation, with the predicted constants, gives low values for the amount of methane adsorbed. Data over a wider range of pressure and composition at this temperature should be evaluated before any definite conclusions regarding these predicted values are made.

#### An Adsorption Model

Many authors, as discussed previously, have suggested specific models for adsorption as a basis for theoretical development of various correlating equations. It would appear, however, that in light of the present knowledge and development, a modelistic approach based on data in the mid-pressure range is not justified. This is especially true since either the Langmuir or Freundlich relation can be developed from any of a number of models.

The Langmuir equation suggests a mono-layer. In the present studies, the maximum amount adsorbed would cover approximately 95 per cent of the surface area (as measured by the "B E T" method). Many authors, among them Joyner<sup>(31)</sup> and Brunauer,<sup>(14)</sup> suggest that a second layer begins to form when the first layer is 35 per cent filled. In the case of definite porous structures of the type under study here, these assumptions are of little significance. Until a more detailed sub-microscopic study of the adsorption phenomena on sieve type adsorbents can be made, it appears unwise to draw any conclusions as to the method by which the adsorption proceeds.

## CONCLUSIONS AND RECOMMENDATIONS

The present study was undertaken to obtain adsorption data for the binary system methane-nitrogen on Molecular Sieves.

For pure component adsorption, both the Langmuir and Freundlich equations will predict the adsorbent loading well, to the saturation pressure below the critical temperature and to about 90 atmospheres above the critical point, over a wide range of temperature. A comparison of the results with limited data previously published by Lewis indicates that the Molecular Sieve has a greater capacity for methane than do either Davison Silica Gel or Columbia G Activated Carbon. Limited data obtained by von Antropff on an activated charcoal indicate that the capacity of the carbon for nitrogen is slightly higher than the capacity of the Molecular Sieve.

A Polanyi type correlation, as those suggested by Dubinin or Lewis and co-workers, would be very useful from a theoretical as well as practical point. Unfortunately, the data did not correlate well using these methods. This appears to be due to the lack of a satisfactory method of predicting the density of the adsorbed state. Therefore, this method of correlation must wait for more fundamental studies than the one attempted here.

For mixtures, for practical purposes, the use of the total loading together with the relative volatility will give sufficiently accurate results. Somewhat more accurate results may be obtained using the individual Langmuir isotherms with proper coefficients as proposed in the section on discussion of results. A study of the constants supports the ideas advanced by Schay regarding interaction of the components in the adsorbed state which

changes the values of the isotherm constants from those of the pure component. No simple interaction coefficients as those suggested by Schay could be obtained from a study of the present data. This suggests that more complex models must be assumed or that the interaction is more complex than can be analytically treated at this time.

Much fundamental work remains to be done here before a general method of predicting adsorbate loading from mixtures from pure component adsorption data can be advanced. The present study indicates that there is equal competition for adsorption surface since the total amount adsorbed is independent of the equilibrium gas phase composition. It is quite possible that there is some interference between the molecular species.

~~This future work must include a better relationship for the~~ phase behavior of mixtures in both the gas and liquid state, a microscopic study of the adsorbate-adsorbent interface, and the phase behavior of the adsorbate in its adsorbed state.

In the present study at low temperatures, the composition range was severely limited as it was difficult to obtain good equilibrium data below one atmosphere due to limitations imposed by the equipment. The study of this system at low pressures and over the whole composition range is recommended.

From the present study it would appear that at the lower temperatures adsorption does not enjoy any advantage as a separative method for nitrogen from methane over vapor-liquid contacting, since the relative volatilities of both methods are about equal and the adsorbent must be

regenerated constantly. At higher temperatures, or for final scrubbing of a stream of nitrogen and methane, there may be applications where Molecular Sieve type adsorbents will prove to be quite useful.

## APPENDIX A

### EXPERIMENTAL DATA

The basic experimental data was converted to absolute pressure and temperature before being used in computations. The mass spectrometer analyses of the samples were reduced to mole fractions. These data, together with the applicable volumes, were then coded on cards which were read by the computer program and served as the basis for computation.

The experimental data as reduced for the computer program may be found in Table XIX. In order to amplify the listing, a sample set of data with the necessary dimensions added may be found in Table XVIII. Each line corresponds to a data card and at the **right of the line or card** will be found an identification "Run 16". It will be noted that the first line of each series contains only a series of digits on the left. These identify the data set or run to the program as well as the number of samples taken plus descriptions made. The next two data cards or lines of data contain the information regarding the condition of the feed cell before (PI, TI, VI) and after (PF, TF, VF) loading as well as the composition of the gas in the cell ( $YI(CH_4)$ ,  $YI(N_2)$ ). The cards following either refer to samples taken at equilibrium or to desorption. They may easily be distinguished as the equilibrium samples will have a three-digit number (i.e., 002) preceding the pertinent data, while the desorption data are preceded by a two digit number preceded by a minus sign. In addition, no cell pressure (PC) or cell temperature (TC) is recorded for desorption.

The data cards containing equilibrium data, contain the following information in order from left to right:

- 1) a three digit number identifying the sample,

- 2) the cell pressure (PC) in psia,
- 3) the cell temperature (TC) in °K,
- 4) the pressure of the sample (PS) in psia,
- 5) the temperature of the sample (TS) in °K,
- 6) the sample volume (VS) in cc,
- 7) the mole fraction methane (YC(CH<sub>4</sub>)) in the sample,
- 8) the mole fraction nitrogen (YC(N<sub>2</sub>)) in the sample.

In the data cards which contain desorption information, as previously stated, no cell data will be found. The pressure, temperature and volume of the desorbate, followed by the composition where it was possible to obtain a sample or an assumed value in case no experimental value was available, are recorded.

It will be noted that a number of runs were eliminated from consideration. Several of these were eliminated because the amount desorbed indicated a leak in the system. One run was not used because the liquefaction pressure had been exceeded. In addition, two pure component runs were not used because they did not agree with the remainder of the data as well as the check runs made to check these values.

Runs 116 and 117 were blank runs made to check the adsorptive capacity of the empty cell.



TABLE XVIII

SAMPLE OF EXPERIMENTAL DATA INPUT FOR COMPUTER

Run No.	No. of Data Pts.	PI	TI	PF	TF	YI <sub>CH<sub>4</sub></sub>	YI <sub>N<sub>2</sub></sub>	Run
016	004	415.0 psia	297.0 °K	317.0 psia	297.0 °K	0.7062	0.2938	Run 16
VI		948.8cc.	VF	959.9cc.				
				89.72cc.				
		PC	TC	PS	TS	VS	YC <sub>CH<sub>4</sub></sub>	YC <sub>N<sub>2</sub></sub>
001	210.0 psia	292.75 °K	29.40 psia	295.2 °K	23.19cc.	0.6612	0.3388	Run 16
002	210.0	292.75	29.4	296.3	23.79	0.5903	0.4097	Run 16
- 03			7.85	294.1	11771.0	1.0	0.0	Run 16
- 04			1.82	293.0	179.6	1.0	0.0	Run 16

TABLE XIX

EXPERIMENTAL DATA

(See Table XVIII for Explanation  
of Format)







042003											
1810.0	305.7	1398.0	305.7	0.5093	0.4907					RUN42	RUN42
948.8	959.9	89.72									
001 580.0	194.45	79.0	296.5	23.79	0.4370	0.5629				RUN42	RUN42
002 565.0	194.45	77.0	296.5	23.79	0.4369	0.5531				RUN42	RUN42
-03		16.881	297.6	11774.7	0.4427	0.5573				RUN42	RUN42
043005										RUN43	RUN43
1295.0	301.25	1055.0	301.25	0.5759	0.4241					RUN43	RUN43
948.8	959.9	89.72									
001 550.0	191.9	72.0	297.1	23.79	0.4048	0.5952				RUN43	RUN43
002 540.0	191.9	70.0	297.1	23.79	0.4054	0.5946				RUN43	RUN43
-03		9.668	298.2	11774.3	0.4176	0.5824				RUN43	RUN43
-04		7.696	298.8	11773.2	0.7011	0.2989				RUN43	RUN43
-05		1.025	297.6	11770.8	1.0	0.0				RUN43	RUN43
044004										RUN44	RUN44
948.00	302.75	772.0	302.75	0.6360	0.3640					RUN44	RUN44
948.8	959.9	89.72									
001 240.0	194.45	22.0	299.9	23.79	0.4059	0.5941				RUN44	RUN44
-02		7.429	296.0	11772.8	0.4866	0.5134				RUN44	RUN44
-03		5.279	296.0	11772.0	0.7479	0.2521				RUN44	RUN44
-04		0.812	295.4	11769.6	1.0	0.0				RUN44	RUN44
045006										RUN45	RUN45
1867.0	301.5	1455.0	301.5	0.5499	0.4511					RUN45	RUN45
948.8	959.9	89.72									
001 1130.0	192.5	160.0	297.1	23.79	0.4764	0.5236				RUN45	RUN45
-02		19.530	296.5	11779.6	0.4494	0.5506				RUN45	RUN45
-03		8.160	296.5	11773.6	0.6571	0.3429				RUN45	RUN45
-04		2.340	296.5	11770.5	0.75	0.25				RUN45	RUN45
-05		0.367	297.6	11770.0	1.0	0.0				RUN45	RUN45
-06		0.048	297.6	11770.0	1.0	0.0				RUN45	RUN45
046005										RUN46	RUN46
915.0	302.75	673.0	298.25	0.5466	0.4534					RUN46	RUN46
948.8	959.9	89.72									
001 540.0	194.95	70.0	295.4	23.79	0.3887	0.6113				RUN46	RUN46
-02		9.630	296.5	11774.3	0.3954	0.6046				RUN46	RUN46
-03		6.652	296.5	11772.7	0.6537	0.3463				RUN46	RUN46
-04		2.224	296.5	11770.0	0.75	0.25				RUN46	RUN46
-05		0.155	296.5	11769.0	1.0	0.0				RUN46	RUN46
048005										RUN48	RUN48
577.0	298.25	402.0	298.25	0.6522	0.3478					RUN48	RUN48
948.8	959.9	89.72									
001 267.0	192.5	35.0	298.8	23.79	0.4512	0.5488				RUN48	RUN48
-02		10.538	299.9	11774.8	0.5479	0.4521				RUN48	RUN48
-03		4.041	299.3	11771.3	0.8133	0.1869				RUN48	RUN48
-04		1.006	299.9	11769.7	1.0	0.0				RUN48	RUN48
-05		0.048	299.9	11769.0	1.0	0.0				RUN48	RUN48
049005										RUN49	RUN49
1593.0	298.75	1240.0	298.75	0.0007	0.9993					RUN49	RUN49
948.8	959.9	89.72									
001 1145.0	192.9	67.0	298.8	49.6	0.0007	0.9993				RUN49	RUN49
-02		15.237	301.5	11777.2	0.0009	0.9991				RUN49	RUN49
-03		9.978	299.9	11774.4	0.0018	0.9982				RUN49	RUN49
-04		0.309	299.9	11769.3	0.0	1.0				RUN49	RUN49
-05		0.010	299.9	11769.0	0.0	1.0				RUN49	RUN49
050004										RUN50	RUN50
1213.0	305.0	940.0	305.0	0.0006	0.9994					RUN50	RUN50
948.8	959.9	89.72									
001 850.0	193.9	99.0	300.9	23.79	0.0016	0.9984				RUN50	RUN50
-02		15.798	299.9	11779.6	0.0013	0.9987				RUN50	RUN50
-03		4.738	300.4	11771.8	0.0	1.0				RUN50	RUN50
-04		0.512	299.9	11769.5	0.0	1.0				RUN50	RUN50
051003										RUN51	RUN51
825.0	298.0	600.0	298.0	0.0009	0.9991					RUN51	RUN51
948.8	959.9	89.72									
001 470.0	194.7	12.820	299.3	11775.9	0.0009	0.9991				RUN51	RUN51
-02		4.670	298.8	11771.6	0.0	1.0				RUN51	RUN51
-03		0.058	298.8	11769.0	0.0	1.0				RUN51	RUN51
052004										RUN52	RUN52
750.4	298.2	541.9	298.2	0.9920	0.0080					RUN52	RUN52
961.1	972.2	89.72									
001 436.4	195.9	60.0	302.0	23.79	0.9897	0.0103				RUN52	RUN52
-02		7.677	299.8	11783.0	0.9842	0.0158				RUN52	RUN52
-03		10.364	298.7	11758.7	0.9940	0.0060				RUN52	RUN52
-04		0.590	298.7	11733.5	1.0	0.0				RUN52	RUN52
053004										RUN53	RUN53
508.0	298.6	345.0	298.6	0.9900	0.0100					RUN53	RUN53
961.1	972.2	89.72									
001 133.46	195.6	20.0	301.5	23.79	0.9860	0.0140				RUN53	RUN53
-02		4.892	299.8	11786.1	0.9884	0.0116				RUN53	RUN53
-03		9.127	300.4	11758.0	1.0	0.0				RUN53	RUN53
-04		0.300	300.9	11753.0	1.0	0.0				RUN53	RUN53
056005										RUN56	RUN56
687.0	296.9	493.5	296.9	0.9904	0.0096					RUN56	RUN56
961.1	972.2	89.72									
001 365.0	193.8	47.0	299.3	23.79	0.9883	0.0118				RUN56	RUN56
-02		6.710	299.3	11782.5	0.9862	0.0138				RUN56	RUN56
-03		10.055	299.3	11783.0	1.0	0.0				RUN56	RUN56
-04		1.006	299.8	11779.0	1.0	0.0				RUN56	RUN56
-05		0.029	299.8	11779.0	1.0	0.0				RUN56	RUN56

057004								
470.0	297.4	333.0	297.4	0.9887	0.0113		RUN57	
961.1	972.2	89.72					RUN57	
001 52.0	195.5	20.0	299.8	23.79	0.9814	0.0186	RUN57	
-02		11.254	299.8	11784.9	0.9919	0.0081	RUN57	
-03		0.986	298.7	11779.0	1.0	0.0	RUN57	
-04		0.029	298.7	11779.0	1.0	0.0	RUN57	
058006							RUN58	
299.0	296.2	172.8	296.2	0.9895	0.0105		RUN58	
961.1	972.2	89.72					RUN58	
-01		337.3	296.2	-972.2	0.9914	0.0086	RUN58	
-02		296.9	296.2	972.2	0.9914	0.0086	RUN58	
003 193.0	195.5	25.0	300.4	23.79	0.9849	0.0151	RUN58	
-04		11.873	298.7	11780.0	0.9894	0.0106	RUN58	
-05		2.881	299.3	11780.0	1.0	0.0	RUN58	
-06		0.211	299.3	11780.0	1.0	0.0	RUN58	
059003							RUN59	
236.4	298.4	102.2	298.4	0.9797	0.0203		RUN59	
961.1	972.2	89.72					RUN59	
001 40.0	194.9	40.0	300.9	23.79	0.9830	0.0170	RUN59	
-02		10.287	300.4	11780.0	0.9918	0.0082	RUN59	
-03		0.850	300.4	11780.0	1.0	0.0	RUN59	
060005							RUN60	
770.5	298.16	556.0	798.16	0.5241	0.4759		RUN60	
961.1	972.2	89.72					RUN60	
001 485.0	195.8	60.0	298.2	23.79	0.3517	0.6483	RUN60	
-02		7.793	299.3	11780.0	0.3572	0.6428	RUN60	
-03		9.108	299.8	11780.0	0.500	0.500	RUN60	
-04		1.122	299.3	11780.0	1.0	0.0	RUN60	
-05		0.072	299.3	11780.0	1.0	0.0	RUN60	
060004							RUN60A	
537.9	295.8	369.6	296.2	0.5205	0.4795		RUN60A	
961.1	972.2	89.72					RUN60A	
001 178.0	194.4	20.0	298.7	23.79	0.5066	0.4934	RUN60A	
-02		13.014	299.8	11780.0	0.3060	0.6940	RUN60A	
-03		0.870	299.8	11780.0	1.0	0.0	RUN60A	
-04		0.012	299.8	11780.0	1.0	0.0	RUN60A	
061004							RUN61	
349.4	296.9	203.1	296.9	0.5208	0.4792		RUN61	
961.1	972.2	89.72					RUN61	
001 95.6	195.5	13.26	297.0	23.79	0.3139	0.6861	RUN61	
-02		10.887	293.2	11780.0	0.4959	0.5041	RUN61	
-03		1.305	297.6	11780.0	0.500	0.500	RUN61	
-04		0.064	297.6	11780.0	1.0	0.0	RUN61	
062004							RUN62	
193.0	296.2	66.8	296.2	0.5195	0.4805		RUN62	
961.1	972.2	89.72					RUN62	
001 29.0	192.8	24.0	295.4	23.79	0.3344	0.6656	RUN62	
-02		9.707	299.8	11780.0	0.5065	0.4935	RUN62	
-03		0.831	299.3	11780.0	1.0	0.0	RUN62	
-04		0.041	299.3	11780.0	1.0	0.0	RUN62	
063005							RUN63	
381.2	296.2	559.0	296.2	0.2411	0.7589		RUN63	
961.1	972.2	89.72					RUN63	
001 550.0	200.9	65.0	295.4	23.79	0.1382	0.8618	RUN63	
-02		8.373	297.6	11780.0	0.1382	0.8618	RUN63	
-03		8.412	298.7	11780.0	0.3078	0.6922	RUN63	
-04		0.928	298.7	11780.0	1.0	0.0	RUN63	
-05		0.039	298.7	11780.0	1.0	0.0	RUN63	
064005							RUN64	
533.8	296.2	349.4	296.4	0.2405	0.7595		RUN64	
961.1	972.2	89.72					RUN64	
001 301.0	192.8	37.0	294.3	23.79	0.1214	0.8786	RUN64	
-02		12.105	298.7	11780.0	0.1962	0.8038	RUN64	
-03		2.630	298.7	11780.0	0.500	0.500	RUN64	
-04		0.261	298.7	11780.0	1.0	0.0	RUN64	
-05		0.014	298.7	11780.0	1.0	0.0	RUN64	
065005							RUN65	
336.2	296.9	186.0	296.9	0.2402	0.7598		RUN65	
961.1	972.2	89.72					RUN65	
001 111.5	192.8	14.7	296.5	23.79	0.1261	0.8739	RUN65	
002 106.0	192.8	14.7	296.5	23.79	0.1263	0.8737	RUN65	
-03		11.505	299.8	11780.0	0.2274	0.7726	RUN65	
-04		0.870	298.7	11780.0	0.500	0.500	RUN65	
-05		0.021	299.8	11780.0	1.0	0.0	RUN65	
066004							RUN66	
177.9	298.66	61.8	298.66	0.2395	0.7605		RUN66	
961.1	972.2	89.72					RUN66	
001 20.3	190.5	8.35	298.2	23.79	0.1326	0.8674	RUN66	
-02		8.798	301.5	11780.0	0.2282	0.7718	RUN66	
-03		0.657	301.5	11780.0	0.500	0.500	RUN66	
-04		0.029	301.5	11780.0	1.0	0.0	RUN66	
067005							RUN67	
755.4	297.91	535.8	297.91	0.7100	0.2900		RUN67	
961.1	972.2	89.72					RUN67	
001 458.5	193.83	60.0	294.3	23.79	0.5638	0.4362	RUN67	
-02		8.257	299.8	11780.0	0.5740	0.4260	RUN67	
-03		9.475	298.7	11780.0	0.8185	0.1815	RUN67	
-04		0.503	298.7	11780.0	1.0	0.0	RUN67	
-05		0.029	298.7	11780.0	1.0	0.0	RUN67	
068005							RUN68	
515.7	298.16	323.2	298.16	0.7122	0.2878		RUN68	
961.1	972.2	89.72					RUN68	
001 241.0	192.84	28.0	296.5	23.79	0.5312	0.4688	RUN68	
-02		12.086	299.8	11780.0	0.6803	0.3196	RUN68	
-03		2.765	299.8	11780.0	0.7500	0.2500	RUN68	
-04		0.193	299.8	11780.0	1.00	0.0	RUN68	
-05		0.014	299.8	11780.0	1.0	0.0	RUN68	

069005											RUN69
311.0	300.66	152.0		300.66	0.7113	0.2887					RUN69
961.1	972.2	89.72									
001 117.0	192.84	16.0	299.8	23.79	0.5081	0.4919					RUN69
-02		11.119	300.9	11780.0	0.6802	0.3198					RUN69
-03		1.992	302.0	11780.0	0.7500	0.2500					RUN69
-04		0.213	302.0	11780.0	1.0	0.0					RUN69
-05		0.014	302.0	11780.0	1.0	0.0					RUN69
070004											RUN70
145.6	300.91	26.5		300.91	0.7115	0.2885					RUN70
961.1	972.2	89.72									
001 14.0	194.66	8.84	299.3	23.79	0.5415	0.4585					RUN70
-02		8.856	302.0	11780.0	0.6977	0.3023					RUN70
-03		1.006	301.5	11780.0	0.7500	0.2500					RUN70
-04		0.014	301.5	11780.0	1.0	0.0					RUN70
071004											RUN71
762.5	297.9	543.0		297.9	0.0	1.0					RUN71
961.1	972.2	89.72									
001 492.0	192.16	56.0	297.6	23.79	0.0	1.0					RUN71
-02		16.243	298.7	11780.0	0.0	1.0					RUN71
-03		1.102	299.8	11780.0	0.0	1.0					RUN71
-04		0.021	299.8	11780.0	0.0	1.0					RUN71
072004											RUN72
520.4	300.28	339.3		300.28	0.0	1.0					RUN72
961.1	972.2	89.72									
001 258.0	192.84	30.0	298.2	23.79	0.0	1.0					RUN72
-02		12.878	300.4	11780.0	0.0	1.0					RUN72
-03		1.568	302.0	11780.0	0.0	1.0					RUN72
-04		0.073	302.0	11780.0	0.0	1.0					RUN72
073004											RUN73
323.2	300.0	175.2		300.0	0.0	1.0					RUN73
961.1	972.2	89.72									
001 113.5	192.84	30.0	297.6	23.79	0.0	1.0					RUN73
-02		10.771	300.9	11780.0	0.0	1.0					RUN73
-03		1.431	302.0	11780.0	0.0	1.0					RUN73
-04		0.039	302.0	11780.0	0.0	1.0					RUN73
074004											RUN74
168.7	298.75	45.0		298.75	0.0	1.0					RUN 74
961.1	972.2	89.72									
001 42.0	192.84	25.0	298.7	23.79	0.0	1.0					RUN74
-02		9.444	302.0	11780.0	0.0	1.0					RUN74
-03		0.870	302.0	11780.0	0.0	1.0					RUN74
-04		0.019	302.0	11780.0	0.0	1.0					RUN74
075005											RUN 75
742.3	300.4	498.6		300.4	0.7260	0.2740					RUN 75
961.1	972.2	89.72									
001 449.0	175.7	57.0	297.0	23.75	0.5791	0.4209					RUN 75
-02		12.105	298.7	11780.0	0.6278	0.3723					RUN 75
-03		6.226	298.7	11780.0	0.8760	0.1240					RUN 75
-04		1.044	298.7	11780.0	1.0	0.0					RUN 75
-05		0.07	298.7	11780.0	1.0	0.0					RUN 75
076004											RUN 76
328.2	300.4	155.7		300.4	0.7257	0.2743					RUN 76
961.1	972.2	89.72									
001 111.7	174.87	15.7	297.6	23.79	0.4908	0.5092					RUN 76
-02		12.447	301.5	11780.0	0.6949	0.3051					RUN 76
-03		1.799	302.1	11780.0	1.0	0.0					RUN 76
-04		0.126	302.1	11780.0	1.0	0.0					RUN 76
077005											RUN 77
749.4	301.2	458.1		301.2	0.9990	0.0010					RUN 77
961.1	972.2	89.72									
001 420.0	173.2	60.0	297.6	23.79	0.9963	0.0014					RUN 77
-02		13.652	299.8	11780.0	0.9984	0.0016					RUN 77
-03		10.925	299.8	11780.0	1.0	0.0					RUN 77
-04		1.856	300.4	11780.0	1.0	0.0					RUN 77
-05		0.102	300.9	11780.0	1.0	0.0					RUN 77
078004											RUN 78
584.2	301.2	366.5		301.2	0.9995	0.0005					RUN 78
961.1	972.2	89.72									
001 331.0	175.2	48.0	299.8	23.79	0.9978	0.0022					RUN 78
-02		14.232	302.1	11780.0	0.9931	0.0069					RUN 78
-03		3.887	303.2	11780.0	1.0	0.0					RUN 78
-04		0.435	303.2	11780.0	1.0	0.0					RUN 78
079004											RUN 79
348.4	299.2	172.8		299.2	0.9931	0.0069					RUN 79
961.1	972.2	89.72									
001 132.0	174.2	17.7	298.7	23.79	0.9883	0.0117					RUN 79
-02		12.337	302.1	11780.0	0.9928	0.0072					RUN 79
-03		2.398	302.1	11780.0	1.0	0.0					RUN 79
-04		0.261	302.1	11780.0	1.0	0.0					RUN 79
080004											RUN80
165.8	300.4	24.4		300.4	0.9924	0.0076					RUN80
961.1	986.9	89.72									
001 7.4	173.2	7.37	298.7	23.79	0.9862	0.0138					RUN80
-02		10.345	302.6	11754.0	1.0	0.0					RUN80
-03		1.276	300.9	11754.0	1.0	0.0					RUN80
-04		0.048	300.9	11754.0	1.0	0.0					RUN80
081006											RUN81
421.0	301.0	206.1		301.0	0.1779	0.8221					RUN81
961.1	972.2	89.72									
-01		27.0	299.8	23.79	0.1766	0.8234					RUN81
002 173.4	123.2	24.0	300.9	23.79	0.0495	0.9505					RUN81
-03		24.0	300.9	2.3	0.0500	0.9500					RUN81
-04		13.748	302.6	11780.0	0.1348	0.8652					RUN81
-05		3.539	300.9	11780.0	0.0	1.0					RUN81
-06		0.261	298.7	11780.0	0.0	1.0					RUN81



082005									RUN 82
197.06	302.2	24.48	302.2	0.1805	0.8195				RUN 82
961.2	984.9	89.72							
001 18.7	122.0	18.38	299.8	23.79	0.0498	0.9502			RUN 82
-02		12.240	303.2	11780.0	0.1694	0.8306			RUN 82
-03		1.315	303.2	11780.0	0.0	1.0			RUN 82
-04		0.097	303.2	11780.0	0.0	1.0			RUN 82
-05		0.010	303.2	11780.0	0.0	1.0			RUN 82
084007									RUN 84
724.2	296.7	419.96	296.7	0.0	1.0				RUN 84
961.1	972.2	89.72							
001 364.7	123.0	100.0	298.0	0.0	0.0	1.0			RUN 84
002 374.8	123.0	100.0	298.0	0.0	0.0	1.0			RUN 84
003 369.8	122.0	41.0	296.5	23.79	0.0	1.0			RUN 84
-04		18.525	299.8	11780.0	0.0	1.0			RUN 84
-05		5.434	298.7	11780.0	0.0	1.0			RUN 84
-06		1.740	298.7	11780.0	0.0	1.0			RUN 84
-07		0.386	298.7	11780.0	0.0	1.0			RUN 84
085004									RUN 85
515.71	298.2	263.65	298.2	0.0	1.0				RUN 85
961.1	972.2	89.72							
001 258.7	123.0	36.56	297.6	23.79	0.0	1.0			RUN 85
-02		17.694	300.9	11780.0	0.0	1.0			RUN 85
-03		1.798	300.9	11780.0	0.0	1.0			RUN 85
-04		0.126	300.9	11780.0	0.0	1.0			RUN 85
086005									RUN 86
404.84	299.7	185.96	299.7	0.0	1.0				RUN 86
961.1	972.2	89.72							
001 171.0	123.0	100.0	298.0	0.0	0.0	1.0			RUN 86
002 173.0	123.0	101.0	298.7	23.79	0.0	1.0			RUN 86
-03		16.301	300.9	11780.0	0.0	1.0			RUN 86
-04		1.160	300.9	11780.0	0.0	1.0			RUN 86
-05		0.019	300.9	11780.0	0.0	1.0			RUN 86
087006									RUN 87
262.5	299.4	37.3	299.4	0.0	1.0				RUN 87
961.2	984.9	89.72							
001 196.0	125.5	100.0	298.0	0.0	0.0	1.0			RUN 87
002 176.0	123.0	100.0	298.0	0.0	0.0	1.0			RUN 87
003 196.0	123.0	30.0	304.1	23.79	0.0	1.0			RUN 87
-04		16.378	301.5	11780.0	0.0	1.0			RUN 87
-05		1.470	301.5	11780.0	0.0	1.0			RUN 87
-06		0.048	301.5	11780.0	0.0	1.0			RUN 87
088005									RUN 88
350.4	303.0	155.7	303.0	0.0	1.0				RUN 88
961.2	972.2	89.72							
001 87.0	125.5	100.0	298.0	0.0	0.0	1.0			RUN 88
002 85.0	123.0	49.0	296.5	23.79	0.0	1.0			RUN 88
-03		14.483	298.7	11780.0	0.0	1.0			RUN 88
-04		1.102	299.8	11780.0	0.0	1.0			RUN 88
-05		0.029	299.8	11780.0	0.0	1.0			RUN 88
089005									RUN 89
257.6	298.0	89.1	298.0	0.0	1.0				RUN 89
961.2	972.2	89.72							
001 23.3	123.0	100.0	298.0	0.0	0.0	1.0			RUN 89
002 23.3	122.0	21.0	296.5	23.79	0.0	1.0			RUN 89
-03		12.047	300.4	11780.0	0.0	1.0			RUN 89
-04		1.682	298.7	11780.0	0.0	1.0			RUN 89
-05		0.087	298.7	11780.0	0.0	1.0			RUN 89
090006									RUN 90
574.15	298.7	319.13	298.7	0.1818	0.8182				RUN 90
961.2	972.2	89.72							
001 279.0	123.2	20.0	298.0	0.0	0.0511	0.9488			RUN 90
002 269.0	122.7	20.0	298.0	0.0	0.0511	0.9488			RUN 90
003 274.0	123.2	36.0	294.3	23.79	0.0511	0.9488			RUN 90
-04		19.279	299.3	11780.0	0.1643	0.8357			RUN 90
-05		0.889	299.3	11780.0	0.0	1.0			RUN 90
-06		0.042	299.3	11780.0	0.0	1.0			RUN 90
091007									RUN 91
297.6	299.0	110.0	299.0	0.1855	0.8145				RUN 91
961.2	972.2	89.72							
001 84.0	124.2	20.0	298.0	0.0	0.0466	0.9534			RUN 91
002 88.0	124.2	20.0	298.0	0.0	0.0466	0.9534			RUN 91
003 88.0	124.0	80.0	297.0	23.79	0.0466	0.9534			RUN 91
-04		13.804	298.7	11780.0	0.1607	0.8393			RUN 91
-05		1.624	298.7	11780.0	0.0	1.0			RUN 91
-06		0.077	298.7	11780.0	0.0	1.0			RUN 91
-07		0.004	298.7	11780.0	0.0	1.0			RUN 91
092007									RUN 92
411.0	298.2	255.0	298.2	0.9904	0.0096				RUN 92
961.2	972.2	89.72							
001 166.0	200.2	20.0	298.0	0.0	0.9860	0.0140			RUN 92
002 136.0	192.2	20.0	298.0	0.0	0.9860	0.0140			RUN 92
003 161.0	200.2	115.0	297.0	23.79	0.9860	0.0140			RUN 92
-04		11.879	299.3	11780.0	0.9917	0.0083			RUN 92
-05		1.760	299.3	11780.0	1.0	0.0			RUN 92
-06		0.155	299.3	11780.0	1.0	0.0			RUN 92
-07		0.014	299.3	11780.0	1.0	0.0			RUN 92

093005												RUN 93
233.0	299.7	92.6	299.7	0.9895	0.0105							RUN 93
961.2	972.2	89.72										RUN 93
001 43.0	192.9	20.0	298.0	0.0	0.9855	0.0145						RUN 93
002 37.0	191.5	33.0	298.2	23.79	0.9855	0.0145						RUN 93
-03		10.887	300.4	11780.0	0.9910	0.0090						RUN 93
-04		0.967	300.9	11780.0	1.0	0.0						RUN 93
-05		0.039	300.9	11780.0	1.0	0.0						RUN 93
094004												RUN 94
751.4	302.2	529.8	302.2	0.3528	0.6472							RUN 94
961.2	972.2	89.72										RUN 94
001 490.0	196.9	52.0	300.9	23.79	0.2010	0.7990						RUN 94
-02		16.746	303.2	11780.0	0.3528	0.6472						RUN 94
-03		0.677	302.6	11780.0	0.0	1.0						RUN 94
-04		0.029	302.6	11780.0	0.0	1.0						RUN 94
095004												RUN 95
413.0	302.2	244.5	302.2	0.3448	0.6552							RUN 95
961.1	972.2	89.72										RUN 95
001 195.0	195.9	43.0	298.7	23.79	0.1762	0.8238						RUN 95
-02		12.898	300.9	11780.0	0.3310	0.6690						RUN 95
-03		0.657	300.9	11780.0	1.0	0.0						RUN 95
-04		0.029	300.9	11780.0	1.0	0.0						RUN 95
096004												RUN 96
232.4	300.7	92.1	300.7	0.3457	0.6543							RUN 96
961.1	972.2	89.72										RUN 96
001 70.0	194.2	69.0	296.5	23.79	0.1876	0.8124						RUN 96
-02		10.248	300.4	11780.0	0.3317	0.6683						RUN 96
-03		0.851	300.4	11780.0	1.0	0.0						RUN 96
-04		0.048	300.4	11780.0	1.0	0.0						RUN 96
097004												RUN 97
744.3	300.7	529.8	300.7	0.5455	0.4545							RUN 97
961.1	972.2	89.72										RUN 97
001 420.0	194.7	102.0	298.7	23.79	0.3566	0.6434						RUN 97
-02		15.585	299.8	11780.0	0.5250	0.4750						RUN 97
-03		1.276	300.4	11780.0	1.0	0.0						RUN 97
-04		0.058	300.4	11780.0	1.0	0.0						RUN 97
098005												RUN 98
505.6	302.2	317.1	302.2	0.5405	0.4595							RUN 98
961.1	972.2	89.72										RUN 98
001 287.0	195.2	66.0	298.2	23.79	0.3335	0.6665						RUN 98
-02		13.768	299.8	11780.0	0.5202	0.4798						RUN 98
-03		1.257	299.8	11780.0	1.0	0.0						RUN 98
-04		0.081	299.8	11780.0	1.0	0.0						RUN 98
-05		0.0097	299.8	11780.0	1.0	0.0						RUN 98
099004												RUN 99
277.5	301.0	126.6	301.0	0.5340	0.4660							RUN 99
961.1	972.2	89.72										RUN 99
001 102.6	195.2	94.0	297.6	23.79	0.3257	0.6743						RUN 99
-02		11.428	299.3	11780.0	0.5289	0.4711						RUN 99
-03		0.851	299.3	11780.0	1.0	0.0						RUN 99
-04		0.035	299.3	11780.0	1.0	0.0						RUN 99
100005												RUN100
660.0	299.2	463.0	299.2	0.8123	0.1877							RUN100
961.1	972.2	89.72										RUN100
001 358.0	195.2	81.0	297.0	23.79	0.6873	0.3127						RUN100
-02		15.276	299.8	11780.0	0.8062	0.1938						RUN100
-03		1.102	299.8	11780.0	1.0	0.0						RUN100
-04		0.058	299.8	11780.0	1.0	0.0						RUN100
-05		0.0039	299.8	11780.0	1.0	0.0						RUN100
101005												RUN101
432.0	299.0	257.6	299.0	0.8156	0.1844							RUN101
961.1	972.2	89.72										RUN101
001 205.5	194.7	67.0	298.7	23.79	0.6562	0.3418						RUN101
-02		12.956	300.4	11780.0	0.7842	0.2158						RUN101
-03		1.392	300.4	11780.0	1.0	0.0						RUN101
-04		0.073	300.4	11780.0	1.0	0.0						RUN101
-05		0.015	300.4	11780.0	1.0	0.0						RUN101
102005												RUN102
245.5	299.7	99.2	299.7	0.8138	0.1862							RUN102
961.1	972.2	89.72										RUN102
001 66.0	196.2	51.7	298.7	23.79	0.7003	0.2997						RUN102
-02		10.829	301.5	11780.0	0.8053	0.1942						RUN102
-03		1.044	301.5	11780.0	1.0	0.0						RUN102
-04		0.056	301.5	11780.0	1.0	0.0						RUN102
-05		0.0097	301.5	11780.0	1.0	0.0						RUN102
103004												RUN103
235.4	301.2	160.7	301.2	0.0349	0.9651							RUN103
961.1	972.2	89.72										RUN103
001 155.0	295.9	132.5	297.6	23.79	0.0007	0.9993						RUN103
-02		5.356	300.9	11780.0	0.0	1.0						RUN103
-03		0.213	300.9	11780.0	0.0	1.0						RUN103
-04		0.014	300.9	11780.0	0.0	1.0						RUN103
104005												RUN104
643.6	301.2	502.6	301.2	0.3900	0.6100							RUN104
961.1	972.2	89.72										RUN104
001 488.0	300.9	50.7	299.8	23.79	0.2801	0.7199						RUN104
002 450.0	295.7	82.0	302.1	23.79	0.2712	0.7288						RUN104
-03		10.016	305.2	11780.0	0.3828	0.6172						RUN104
-04		0.348	305.2	11780.0	1.0	0.0						RUN104
-05		0.019	305.2	11780.0	1.0	0.0						RUN104



116004									RUN116
1585.0	295.4	1243.0	295.4	0.0	1.0				RUN116
961.7	972.7	123.6							RUN116
001 1290.0	191.2	15.0	297.6	0.0	0.0	1.0			RUN116
-02		24.152	299.6	11780.0	0.0	1.0			RUN116
-03		0.425	297.6	11780.0	0.0	1.0			RUN116
-04		0.019	297.6	11780.0	0.0	1.0			RUN116
117004									RUN117
885.0	295.7	698.0	295.7	1.0	0.0				RUN117
961.7	972.7	123.6							RUN117
001 625.0	192.5	15.0	297.0	0.0	1.0	0.0			RUN1 7
-02		17.171	297.0	11780.0	1.0	0.0			RUN117
-03		0.309	297.0	11780.0	1.0	0.0			RUN117
-04		0.019	297.0	11780.0	1.0	0.0			RUN117
118004									RUN118
525.0	300.0	328.0	300.0	0.7360	0.2640				RUN118
961.2	972.2	89.7							RUN118
001 251.0	175.0	30.0	297.0	23.7	0.5489	0.4511			RUN118
-02		12.5	298.0	11780.0	0.6298	0.3703			RUN118
-03		3.5	298.0	11780.0	1.0	0.0			RUN118
-04		1.5	298.0	11780.0	1.0	0.0			RUN118
120004									RUN 120
469.0	298.2	370.0	298.2	0.0	1.0				RUN 120
961.7	972.7	88.91							RUN 120
001 312.0	298.4	27.0	300.9	23.79	0.0	1.0			RUN 120
-02		7.851	297.7	11780.0	0.0	1.0			RUN 120
-03		0.715	297.7	11780.0	0.0	1.0			RUN 120
-04		0.019	297.7	11780.0	0.0	1.0			RUN 120
121006									RUN121
174.0	296.7	117.0	296.7	0.0	1.000				RUN121
961.7	972.7	88.91							RUN121
001 99.6	297.9	10.0	302.6	0.0	0.0	1.000			RUN121
002 94.0	295.9	10.0	302.6	0.0	0.0	1.00			RUN121
003 99.6	297.2	90.0	302.6	26.0	0.0	1.0			RUN121
-04		4.351	298.8	11780.0	0.0	1.0			RUN121
-05		0.193	298.8	11780.0	0.0	1.0			RUN121
-06		0.004	299.3	11780.0	0.0	1.0			RUN121
122005									RUN 122
553.0	298.2	429.0	298.2	1.0	0.0				RUN 122
961.7	972.7	88.91							RUN 122
001 330.8	295.9	39.0	300.0	0.0	1.0	0.0			RUN 122
002 326.0	293.4	39.0	299.3	23.79	1.0	0.0			RUN 122
-03		10.596	298.8	11780.0	1.0	0.0			RUN 122
-04		0.406	298.8	11780.0	1.0	0.0			RUN 122
-05		0.014	298.8	11780.0	1.0	0.0			RUN 122
123004									RUN 123
217.00	299.7	125.0	299.7	1.0	0.0				RUN 123
961.7	972.7	88.91							RUN 123
001 94.0	295.7	85.0	298.7	23.79	1.0	0.0			RUN 123
-02		6.052	297.6	11780.0	1.0	0.0			RUN 123
-03		0.367	297.6	11780.0	1.0	0.0			RUN 123
-04		0.014	297.6	11780.0	1.0	0.0			RUN 123
124005									RUN124
265.0	293.9	170.0	293.9	1.0	0.0				RUN124
961.2	972.7	88.91							RUN124
001 115.0	296.4	30.0	297.61	0.0	1.0	0.0			RUN124
002 145.0	296.4	30.0	297.61	23.79	1.0	0.0			RUN124
-03		7.31	297.1	11780.0	1.0	0.0			RUN124
-04		0.406	297.1	11780.0	1.0	0.0			RUN124
-05		0.029	297.1	11780.0	1.0	0.0			RUN124

APPENDIX B  
CALIBRATIONS

TABLE XX

## THERMOCOUPLE CALIBRATION DATA

Reference Point	Reference Temperature °C	Average* EMF. MV.	Thermocouple Reading	
			Average* Temp. °C (38 Calib)	Average Deviation
Boiling H <sub>2</sub> O (739.5 mm Hg)	99.2	4.25	99.4	0.2
Ice Point	0	0.0	0	0.0
CO <sub>2</sub> Sublimation Point	-78.5	-2.71	-78.25	0.25
Boiling O <sub>2</sub>	-182.96	-5.25	-182.5	0.46
Boiling N <sub>2</sub>	-195.8	-5.47	-195.5	0.30

\* 5 readings each at each point.

TABLE XXI

CALIBRATION DATA FOR PRESSURE GAUGES

Dead Weight Tester PSIG	Gauge Readings PSIG	
	Up	Down
Gauge C-2-522		
50		75
100	120	125
200	220	225
300	325	325
400	420	425
500	520	525
750	800	805
1000	1050	1055
1100	1160	1165
1250	1305	1305
1350	1410	1415
1500	1585	1585
1750	1820	
Gauge C-2-464		
20	27	26
30	36	37
40	48	50
60	67	69
80	85	88
110	117	117
160	165	165
210	215	216
260	265	266
310	316	318
410	415	416
510	518	519
560	570	
Gauge C-2-82		
20	26	26
40	48	47
80	86	85
110	116	117
210	215	

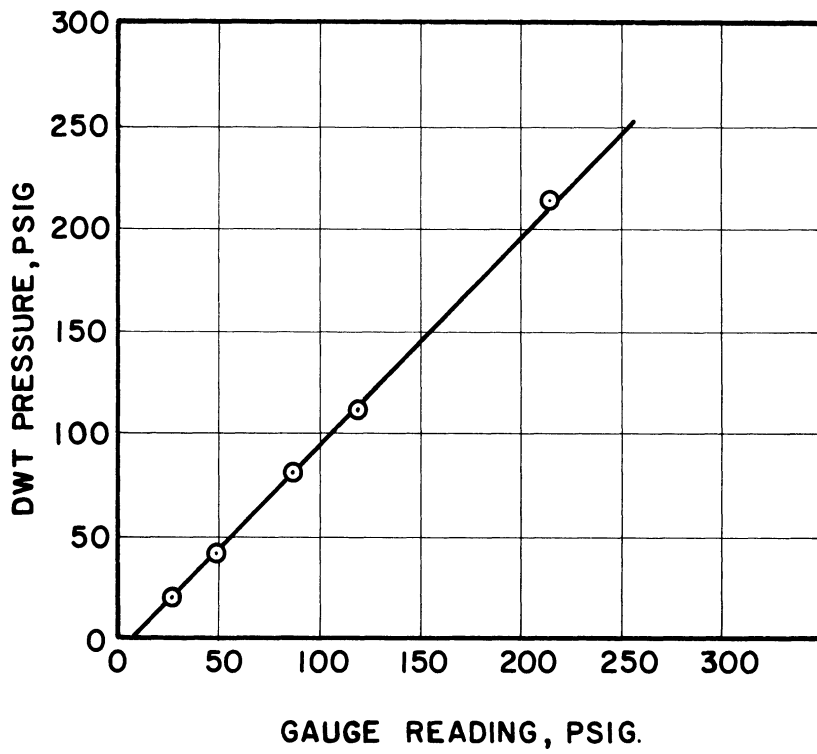


Figure 21. Calibration of 300 Psi Pressure Gauge (C-2-82).



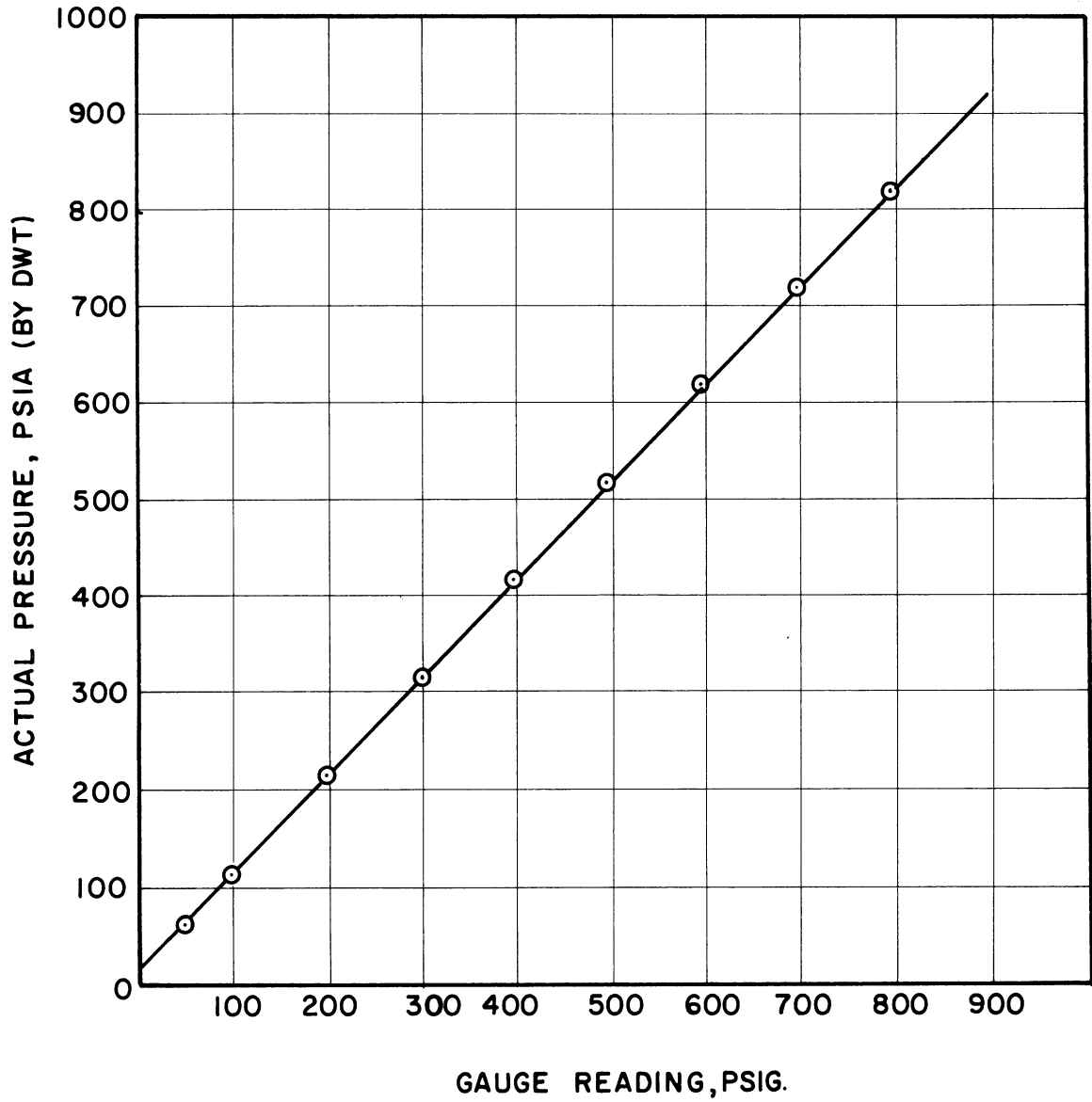


Figure 22. Calibration for 800 Psi Pressure Gauge (No. C-2-464).

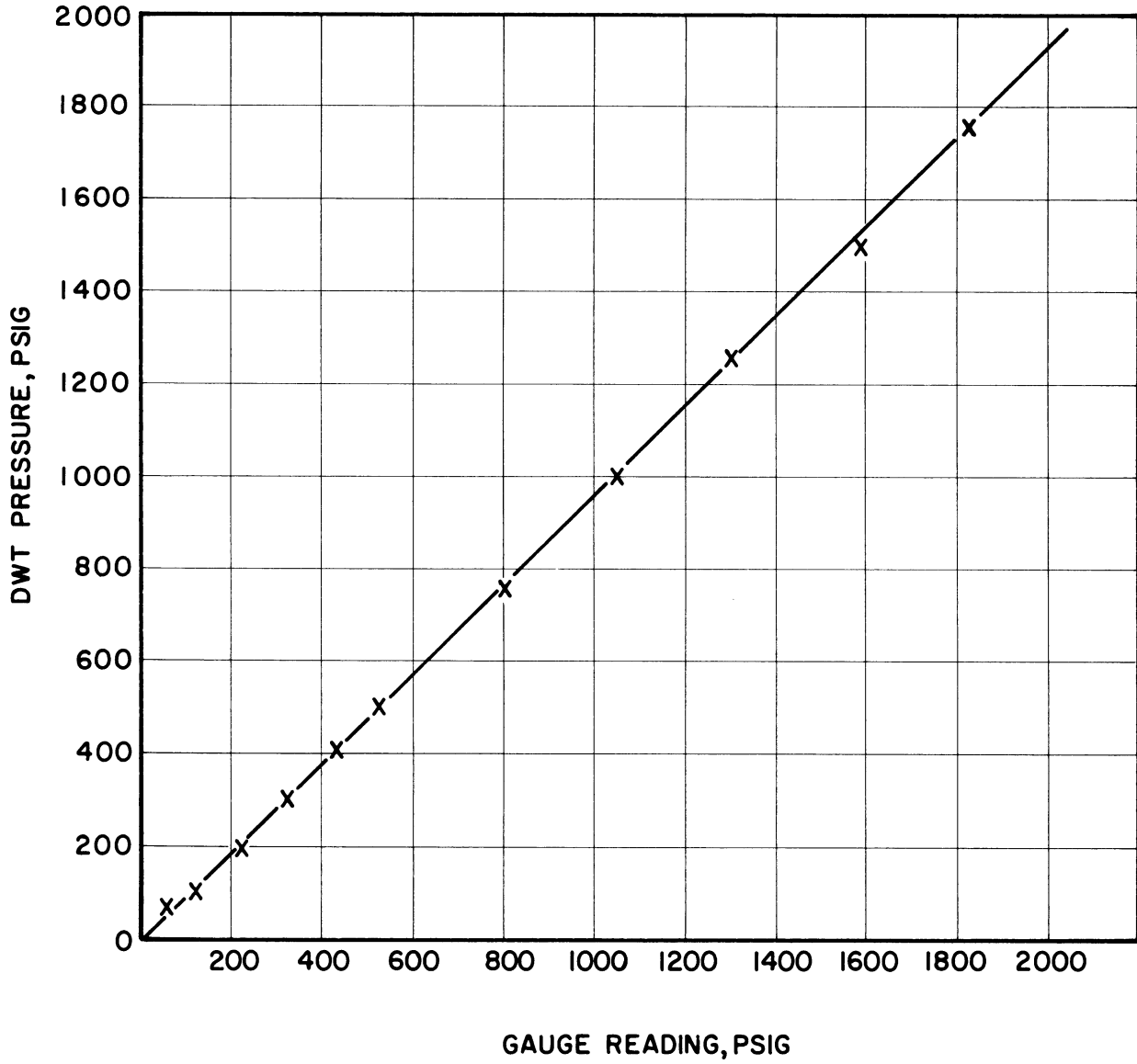


Figure 23. Calibration for 2000 Psi Pressure Gauge (No. C-2-522).

TABLE XXII

CALIBRATION DATA FOR 2000 PSI  
PRESSURE TRANSDUCER  
(Serial No. 110)

Temperature: 26°C

Reference Pressure: 0.01mm Hg.

<u>GAUGE READING</u> <u>PSIA</u>	<u>TRANSDUCER READING</u> <u>MV</u>
290	7.75
1215	36.5
1305	39.0
955	28.5
760	22.5
570	16.75
475	13.5
375	10.5
205	5.0
260	6.25
450	12.5
290	7.75

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TABLE XXIII  
CALIBRATION OF 2000 PSI  
PRESSURE TRANSDUCER  
(Serial No. 110)

Temperature: 79.6°C      Reference Pressure: 0.01mm Hg

<u>GAUGE READING</u> PSIA	<u>TRANSDUCER READING</u> MV
1435	39.0
1140	30.8
1055	27.5
932	25.0
741	20.0
612	16.0
360	9.2
245	6.0
162	4.0
87	2.0
1100	29.5
913	24.8
815	22.0
665	18.0
550	14.4
430	11.2
330	8.4
220	5.4
180	4.5
325	8.4
450	11.8
573	15.0
680	18.0
645	17.0
965	27.0
1419	39.2
1320	36.5
1050	29.0
945	26.0
680	21.0
620	16.7
455	12.2
240	6.0
180	4.6

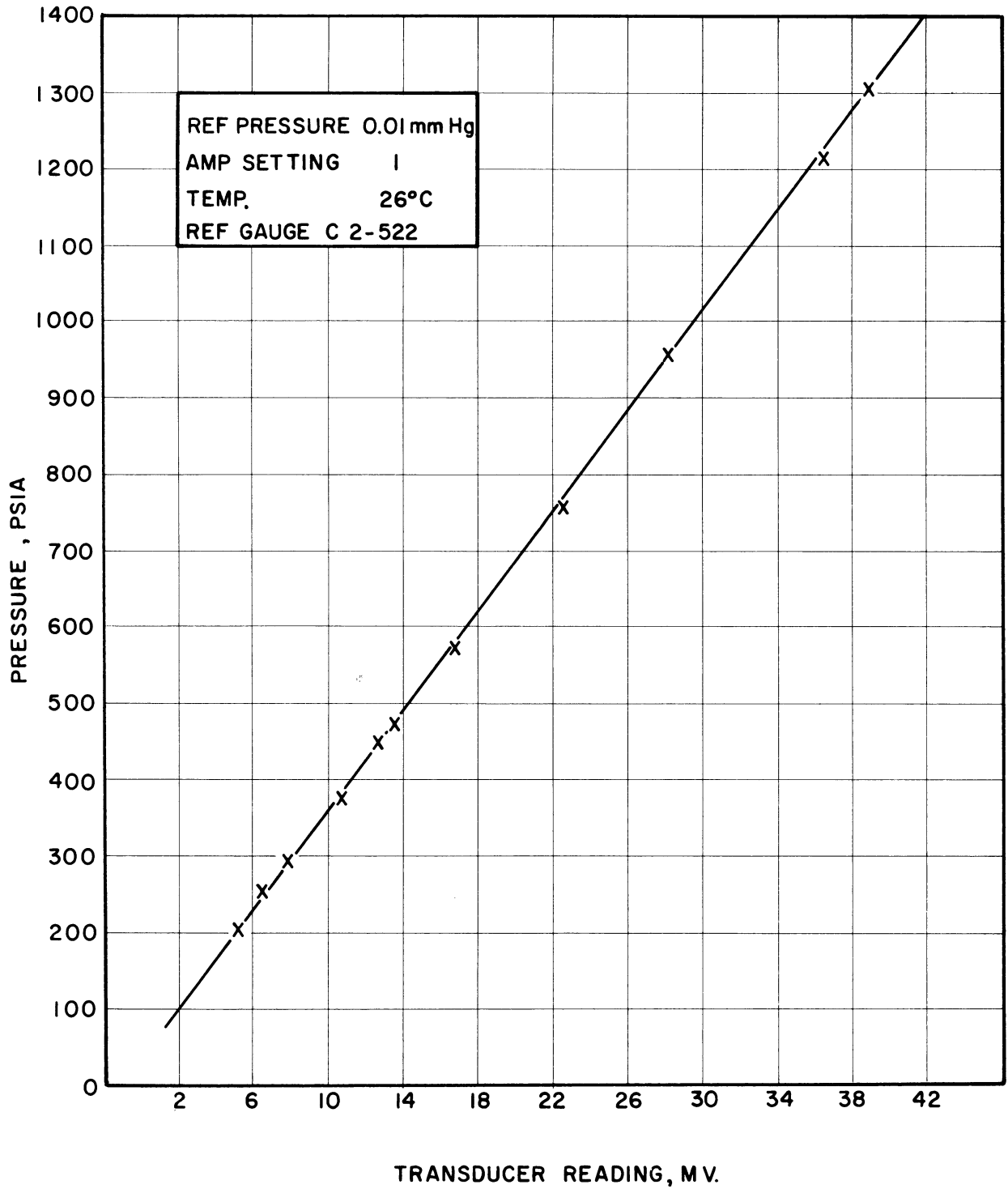


Figure 24. Calibration for 2000 Psi Pressure Transducer (Serial No. 110).

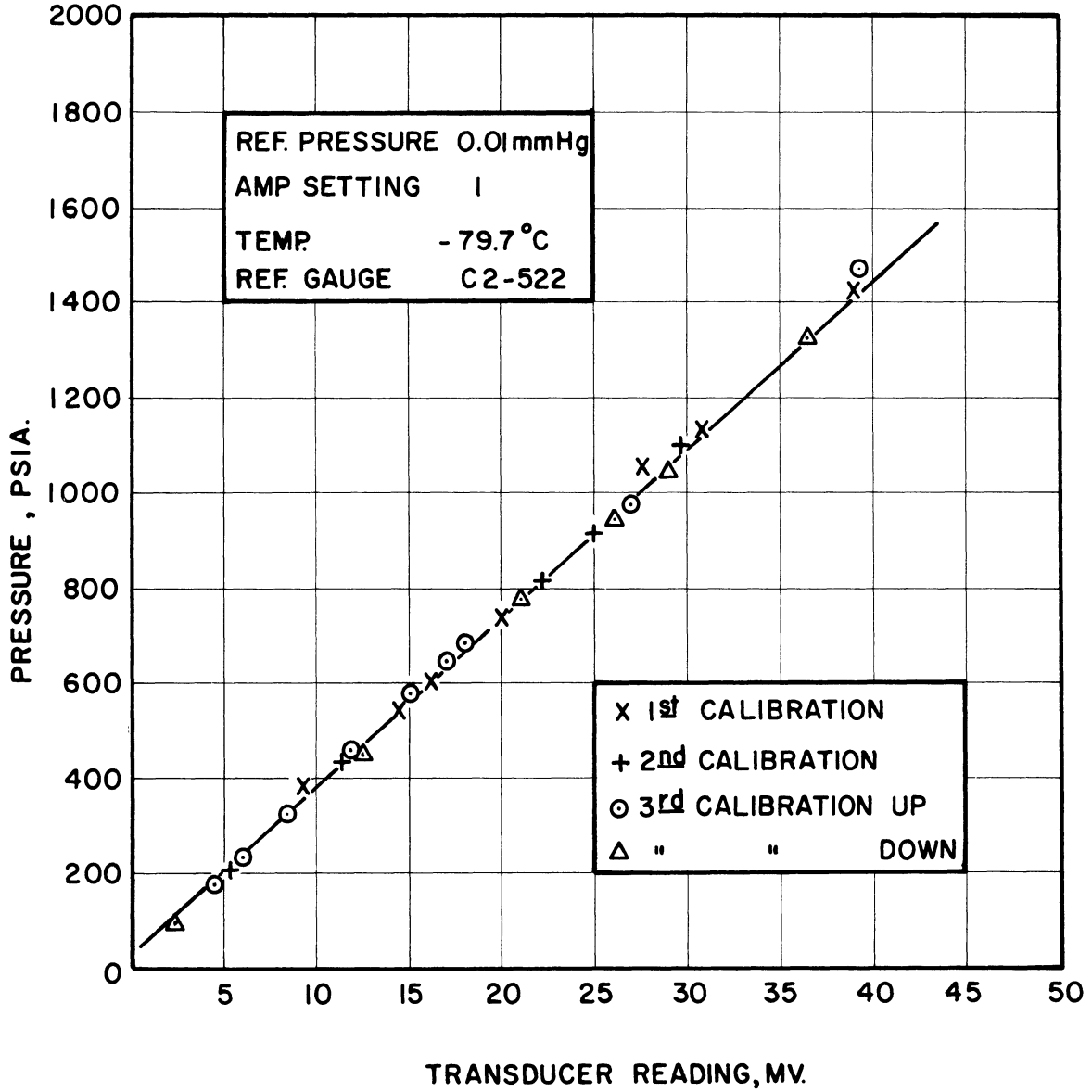


Figure 25. Calibration for 2000 Psi Pressure Transducer (Serial No. 110).

TABLE XXIV

CALIBRATION DATA FOR  
500 PSI PRESSURE TRANSDUCER  
(Serial No. 60)

Gauge Reading PSIA	Transducer Reading MV
Temperature: -79.7°C	Reference Pressure: 0.01mm Hg
115	13.5
218	24.0
321	33.5
423	43.5
523	53.0
621	63.0
656	67.0
566	57.0
519	52.5
420	44.0
368	39.0
320	33.5
216	23.9
150	17.7
119	14.0
67	8.4
421	44.0
115	13.2
Temperature: 20°C	Reference Pressure: 0.01mm Hg
138	16
256	25.5
254	25.0
419	45.0
455	49.0
322	35.0
218	24.0
138	15.6
67	8.0

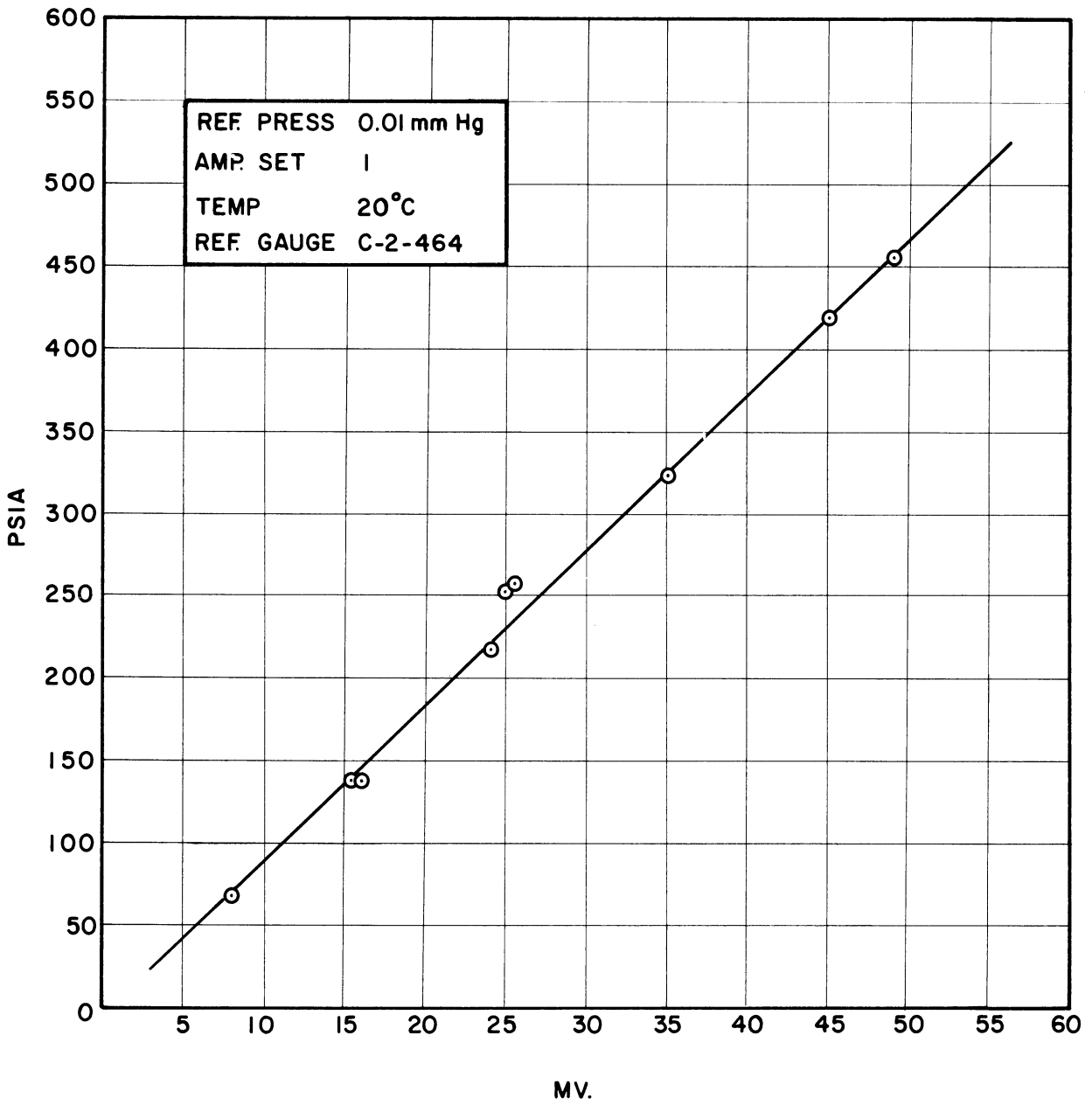


Figure 26. Calibration for 500 Psi Pressure Transducer. (Serial No. 60)



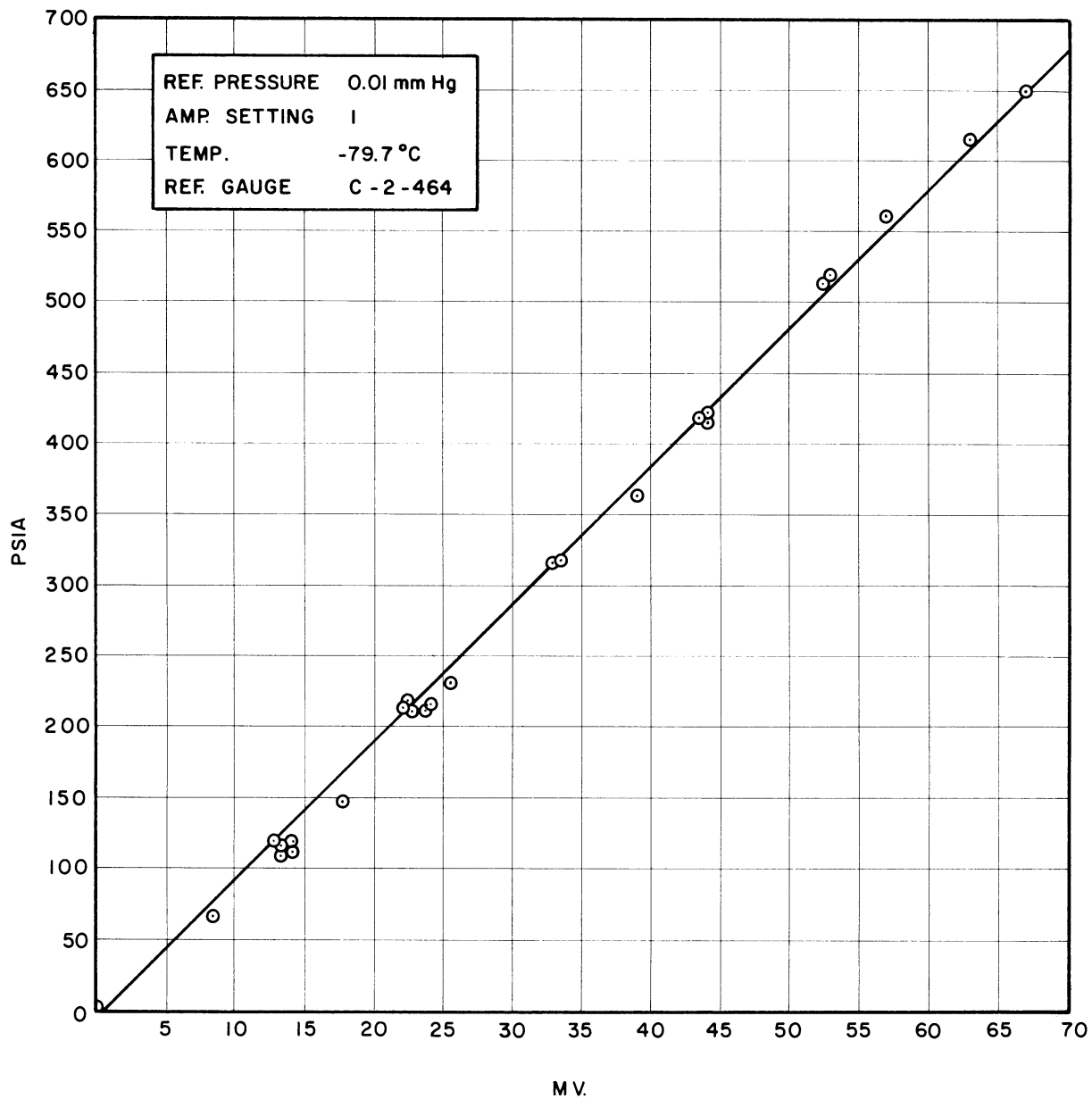


Figure 27. Calibration for 500 Psi Pressure Transducer.  
(Serial No. 60)

## APPENDIX C

### CALCULATION OF SURFACE AREA

Two determinations of the surface area of the adsorbent were made using the standard "BET" method. The first was made after activating the adsorbent for 12 hours at 350°C under a vacuum of  $10^{-6}$  mm. of Hg. For the second determination, the sample was activated at 125°C for 12 hours at  $3 \times 10^{-2}$  mm. Hg. pressure. The results were essentially identical being 494 and 512 square meters per gram for the respective "BET" determinations. The calculated data may be found in Tables XXV and XXVI and Figures 28 and 29.

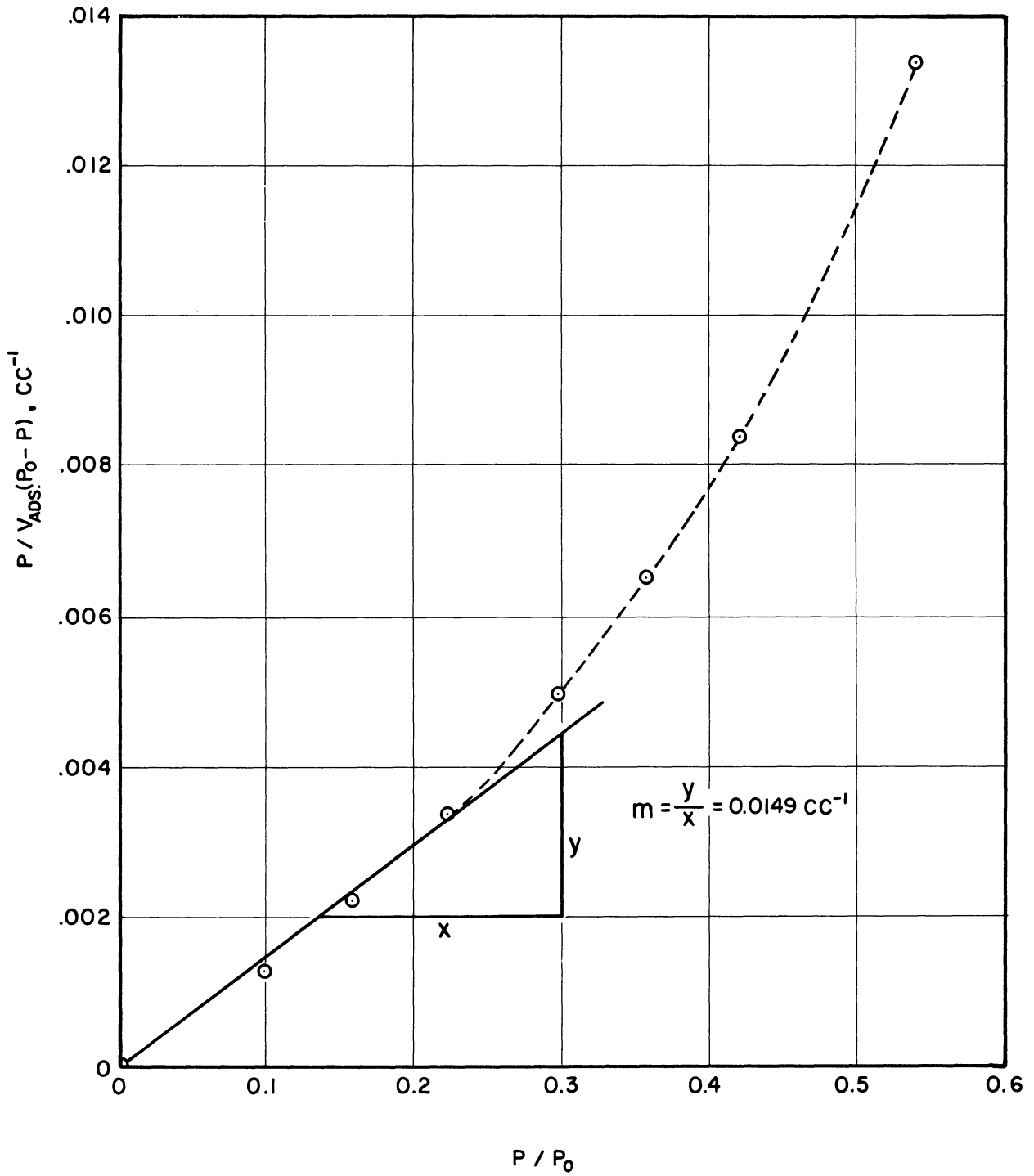


Figure 28. "B E T" Plot for Surface Area  
Determination 1  
Sample - 5A Molecular Sieve

TABLE XXVI

CALCULATION OF BET SURFACE AREA DETERMINATION  
DETERMINATION 2

Activation Temperature = 125°C  
 Activation Pressure =  $3 \times 10^{-2}$  mm Hg  
 Weight of Sample = 0.3023 gms.  
 Average Adsorption Temperature = 193.5°C

DATA

<u>Readings</u>	<u>Pressure mm Hg.</u>	<u>P-Po mm Hg</u>	<u>V Ads ccSTP</u>	<u>P/V Ads(P-Po)</u>	<u>P/Po</u>
1	23.0	712.3	39.5	0.00083	0.031
2	32.0	703.3	39.54	0.00115	0.044
3	41.0	694.3	39.77	0.00148	0.056
4	50.0	685.3	40.02	0.00182	0.068
5	59.5	675.8	40.25	0.00219	0.081
6	78.5	656.8	40.53	0.00295	0.107
7	100.5	634.8	40.68	0.00389	0.137
8	147.5	587.8	41.20	0.00609	0.201
9	206.5	529.3	41.60	0.00936	0.280
10	255.0	480.3	41.83	0.01269	0.347
11	298.0	437.3	42.25	0.01617	0.405
12	398.0	337.3	43.03	0.02742	0.541

$$m = 0.0282 \text{cc}^{-1} \quad (\text{From Figure 29})$$

$$I = 0.00 \text{cc}^{-1}$$

$$v_m = \frac{1}{I + m} = 34.46 \text{ cc}$$

$$\text{Total Surface Area} = \frac{(34.46 \text{cc})(6.02 \times 10^{23})(16.2)}{(22,400)(10^{10})^2} = 154.4 \text{ sq. meters}$$

$$\text{Specific Surface} = \frac{154.4 \text{ sq. meters}}{0.3023 \text{ gms}} = 511 \text{ sq. meters/gm.}$$

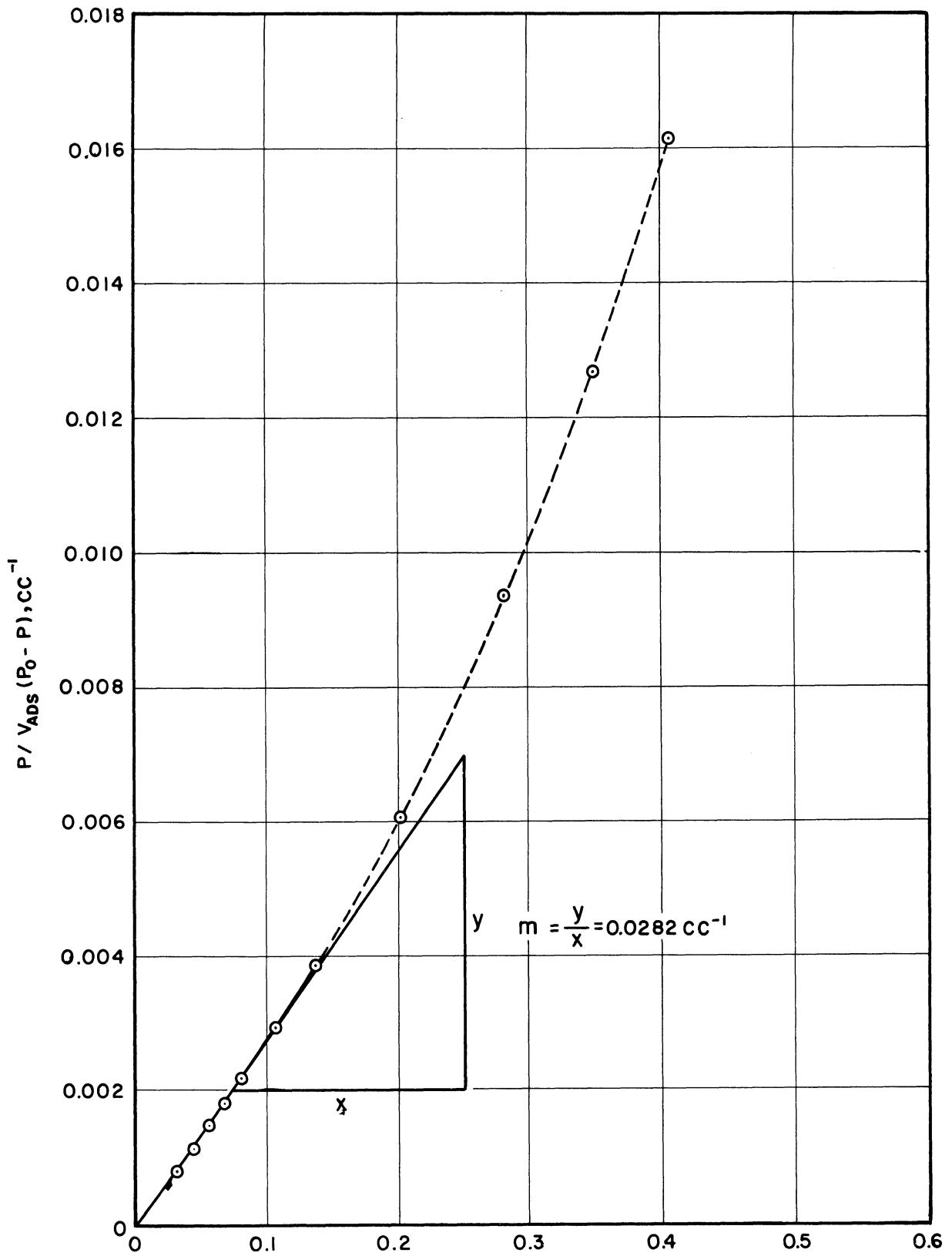


Figure 29. "B E T" Plot for Surface Area Determination 2

Sample - 5A Molecular Sieve

## APPENDIX D

### PRESSURE VESSEL DESIGN CALCULATIONS

#### A. Design Calculation

Using the thick-walled formula from the Unfired Pressure Vessel Code for computing the allowable working pressure (Para. VA-1)

$$P = SE(z-1/z+1)$$

where

$$Z = (R_o/R_o-t)^2$$

$$E = \text{Weld efficiency} = 1.0 \text{ for seamless tube}$$

$$P = \text{allowable pressure}$$

$$R_o = \text{outside radius (in.)} = 1.9 \text{ in}$$

$$S = \text{allowable stress (psi)} = 18,750 \text{ psi}$$

$$t = \text{thickness (in.)} = 0.4 \text{ in.}$$

Then

$$Z = (0.95/0.95-0.4)^2 = 2.99$$

$$P = 18,750 \times 1.0(2.99 - 1/2.99 + 1)$$

$$P = 9,350 \text{ psi.}$$

For the end design the Code (Para. UG - 34) specifies that for flat heads the thickness shall be determined using:

$$t = d \sqrt{CP/S}$$

where

$$t = \text{minimum thickness}$$

$$d = \text{the inside diameter}$$

$$c = \text{appropriate constant} = 0.25$$

as the criteria.

Then 
$$t = 1.1 \sqrt{0.25 \times 9,350/18,750}$$
$$= 0.125 \text{ inches minimum}$$

B. Effect of Temperature and Pressure on Free Volume

The effect on the calibrated volume of the absorption cell by both high pressures and low temperatures may be considerable.

Bartlett<sup>(5)</sup> found that:

$$\Delta V = \frac{V_0}{E(r_2^2 - r_1^2)} [(1-2u)(P_2 r_2^2 - P_1 r_1^2) + 2(1+u)(P_2 - P_1)r_2^2]$$

where

$$V_0 = 124 \text{ cc. at 1 atm.}$$
$$E = 29.3 \times 10^6 \text{ lbs./in.}^2$$
$$u = 0.28$$
$$r_1 = 1.1''$$
$$r_2 = 1.9''$$
$$P_2 = \text{operating pressure}$$
$$P_1 = \text{test pressure}$$

represents the effect of pressure on the volume of a vessel at constant temperature. Calculations indicate that the change of volume resulting from a maximum operating pressure of 2000 psi, is 0.050 cc, or 0.04%.

The effect of a change of temperature on the volume of the adsorption cell was determined using the standard volumetric expansion equation:

$$V_{T_2} = V_{T_1} [1 + 3\alpha(T_2 - T_1)]$$

where the linear coefficient of expansion

$$\alpha = 16 \times 10^{-6} \text{ } ^\circ\text{C}^{-1} \quad \text{for 316 stainless steel.}$$

The minimum operating temperature of 123°K resulted in a change of 0.9 per cent in the volume.

As the total change resulting from the maximum change in the operating conditions is less than one per cent, the approximate error in the calibration, no correction was made for temperature and pressure.



APPENDIX E

SAMPLE CALCULATION OF ADSORPTION EQUILIBRIUM

Sample Calculation

Run No. 76

A. Calculation of Gas Density by B-W-R Equation at Initial Feed Conditions

Experimental Data

$$P = 328.2 \text{ psia}$$

$$T = 300.4^\circ\text{K}$$

$$Y_{\text{CH}_4} = 0.7257$$

Calculation

$$P = \frac{328.2}{14.696} = 22.33 \text{ atm.}$$

Mixed Coefficients

$$\begin{aligned} B_{\text{om}}(\text{BM})^* &= \sum_i B_{oi} y_i \\ &= 0.0426000 \times 0.7257 + 0.0484824 \times 0.2743 \\ &= 0.0442135 \end{aligned}$$

$$\begin{aligned} A_{\text{om}}(\text{AM}) &= \left( \sum_i A_{oi} \frac{1}{2} y_i \right)^2 - 0.1000 (y_{\text{CH}_4} \times y_{\text{N}_2}) \\ &= \left( (1.8550)^{\frac{1}{2}} \times 0.7257 + (1.27389)^{\frac{1}{2}} \times 0.2743 \right)^2 \\ &\quad - 0.1 (0.7257 \times 0.2743) \\ &= 1.6648 \end{aligned}$$

$C_{\text{om}}(\text{CM})$ ,  $D_{\text{om}}(\text{DM})$  and  $\gamma_m(\text{IM})$  follow squared combinatorial rule, without the added correction. Thus:

$$\begin{aligned} C_{\text{om}}(\text{CM}) &= \left( (22570.0)^{\frac{1}{2}} \times 0.7257 + (4273.0)^{\frac{1}{2}} \times 0.2743 \right)^2 \\ &= 16,648.6 \end{aligned}$$

\* The letters in parentheses are used in the computer program.

$$\begin{aligned} D_{om}(DM) &= (7.61781 \times 10^6) \times (0.2743)^2 \\ &= 5.73168 \times 10^5 \end{aligned}$$

$$\begin{aligned} \gamma_m(IM) &= ((0.006000)^{\frac{1}{2}} \times 0.7257 + (0.006500)^{\frac{1}{2}} \times 0.2743)^2 \\ &= 0.006135 \end{aligned}$$

$$\begin{aligned} b_m(EM) &= \left( \sum_i b_i^{1/3} y_i \right)^3 \\ &= ((0.00338004)^{1/3} \times 0.7257 + (0.00232373)^{1/3} \\ &\quad \times 0.2743)^3 \\ &= 0.00306385 \end{aligned}$$

$a_m(FM)$ ,  $c_m(GM)$ ,  $\delta_m(HM)$  and  $\alpha_m(JM)$  follow the cubic combinatorial rule also, therefore:

$$\begin{aligned} a_m(FM) &= ((0.0494000)^{1/3} \times 0.7257 + (0.0178444)^{1/3} \times 0.2743)^3 \\ &= 0.0385993 \end{aligned}$$

$$\begin{aligned} c_m(GM) &= ((2,545.0)^{1/3} \times 0.7257 + (475.00)^{1/3} \times 0.2743)^3 \\ &= 1,748.9 \end{aligned}$$

$$\begin{aligned} \delta_m(HM) &= 0.08320 \times 10^6 \times (0.2743)^3 \\ &= 0.1717 \times 10^5 \end{aligned}$$

$$\begin{aligned} \alpha_m(JM) &= ((0.000124359)^{1/3} \times 0.7257 + (0.00015300)^{1/3} \\ &\quad \times 0.2743)^3 \\ &= 0.000131824 \end{aligned}$$

The second and fourth virial coefficients may now be evaluated independent of the density.

$$\begin{aligned} C_1 &= B_{om}^{RT} - A_{om} - C_{om}/T^2 - D_{om}/T^4 \\ &= 0.0442135 \times 0.08206 \times 300.4 - 1.6648 - \frac{16,648.6}{(300.4)^2} \\ &\quad - \frac{5.73168 \times 10^5}{(300.4)^4} \\ &= -0.753636 \end{aligned}$$

$$\begin{aligned}
 C_3 &= a_m \alpha_m \\
 &= 0.385993 \times 0.000131824 \\
 &= 0.508830 \times 10^{-5}
 \end{aligned}$$

The third virial coefficient

$$C_2 = bRT - a + (C/T^2 - \delta/T^4)(1 + \gamma\rho^2)/e^{\gamma\rho^2}$$

can only be evaluated after a density has been estimated, and if

$$\rho = 0.93058 \text{ gm. moles/liter}$$

then

$$\begin{aligned}
 C_2 &= 0.00306385 \times 0.8206 \times 300.4 - 0.0385993 \\
 &+ \left( \frac{1,748.9}{(300.4)^2} + \frac{1.717 \times 10^4}{(300.4)^4} \right) \times \left( \frac{1 + (0.93058)^2 \times 0.006135}{\rho(0.93058)^2 \times 0.006135} \right) \\
 &= 0.0553097
 \end{aligned}$$

The Pressure (PI) may now be evaluated from the equation of state

$$\begin{aligned}
 PI &= RT + C_1\rho^2 + C_2\rho^3 + C_3\rho^6 \\
 &= 0.08206 \times 300.4 - 0.75363 \times (0.93058)^2 + 0.0563097 \\
 &\times (0.93058)^3 + 5.0883 \times 10^4 \times (0.93058)^6 \\
 &= 22.332 \text{ atm.}
 \end{aligned}$$

Since

$$PI = P$$

the density

$$\rho = 0.930585 \text{ gm. moles/liter}$$

is the correct density.

## B. Calculation of Amount Loaded

### 1. Initial Amount in Feed Reservoir

$$\begin{aligned}
 \text{Methane: GMCI} &= V_{\text{res}} \times \rho_{\text{INT}} \times y_{\text{CH}_4} = VI \times \text{DAI} \times \text{MFCHI} \\
 &= 961.1 \text{ cc.} \times 0.000930585 \text{ gm. moles/liter} \times 0.7527 \\
 &= 0.64906 \text{ gm. moles.}
 \end{aligned}$$

$$\begin{aligned}\text{Nitrogen: GMNI} &= 961.1 \times 0.000930585 \times 0.2473 \\ &= 0.24533 \text{ gm.moles.}\end{aligned}$$

2. Final Amount in Feed Reservoir

$$\text{Density: DAF} = 0.43538 \text{ gm.moles/liter at } 155.7 \text{ psia and } 300.4^\circ\text{K}$$

$$\begin{aligned}\text{Methane: GMCF} &= 972.2 \times 0.00043538 \times 0.7527 \\ &= 0.30718\end{aligned}$$

$$\begin{aligned}\text{Nitrogen: GMNF} &= 972.2 \times 0.00043538 \times 0.2473 \\ &= 0.11611\end{aligned}$$

3. Net Amount Loaded Into Cell

$$\begin{aligned}\text{Methane: LOADC} &= \text{GMCI} - \text{GMCF} \\ &= 0.64906 - 0.30718 \\ &= 0.34188 \text{ gm. moles}\end{aligned}$$

$$\text{Nitrogen: LOADN} = 0.12922$$

C. Calculation of Equilibrium Amount Adsorbed

1. Experimental Data

$$\text{Pressure: PC} = 111.7 \text{ psia}$$

$$\text{Temperature: TC} = 174.87^\circ\text{K}$$

$$\text{Free Space: VC} = 89.72 \text{ cc.}$$

$$\text{Gas Phase Comp, CH}_4, \text{ MFCH} = 0.4908$$

$$\text{N}_2, \text{ MFN} = 0.5092$$

2. Computation

$$\text{Pressure: PC} = 111.7/14.696 = 7.60 \text{ atm.}$$

$$\text{Density: DAC} = 0.0005558 \text{ gm, moles/liter}$$

Amount in Gas Phase

$$\begin{aligned}\text{Methane: } \text{GMGC} &= \text{VC} \times \text{DAC} \times \text{MFCH} \\ &= 89.72 \times 0.0005558 \times 0.4908 \\ &= 0.024476 \\ \text{GMGN} &= 89.72 \times 0.0005558 \times 0.5092 \\ &= 0.02539\end{aligned}$$

Amount in Adsorbed Phase

$$\begin{aligned}\text{GAC} &= \text{LOADC} - \text{GMGC} \\ &= 0.34188 - 0.02448 \\ &= 0.3174 \\ \text{GAN} &= 0.12922 - 0.025394 \\ &= 0.103828 \\ \text{XC} &= \text{GAC}/\text{GAN} + \text{GAC} \\ \text{XC} &= 0.7535\end{aligned}$$

Relative Volatility

$$\begin{aligned}\alpha &= (\text{MFN}/\text{XN})/(\text{MFCH}/\text{XC}) \\ &= 3.28\end{aligned}$$

D. Desorption or Amount Lost in Sampling

Sample calculation for first desorption

1. Experimental Data

Pressure: PS = 12.447 psia

Volume: VS = 11,780 cc.

Temperature: TS = 301.5

Composition: MFCH = 0.6949, MFN = 0.3051

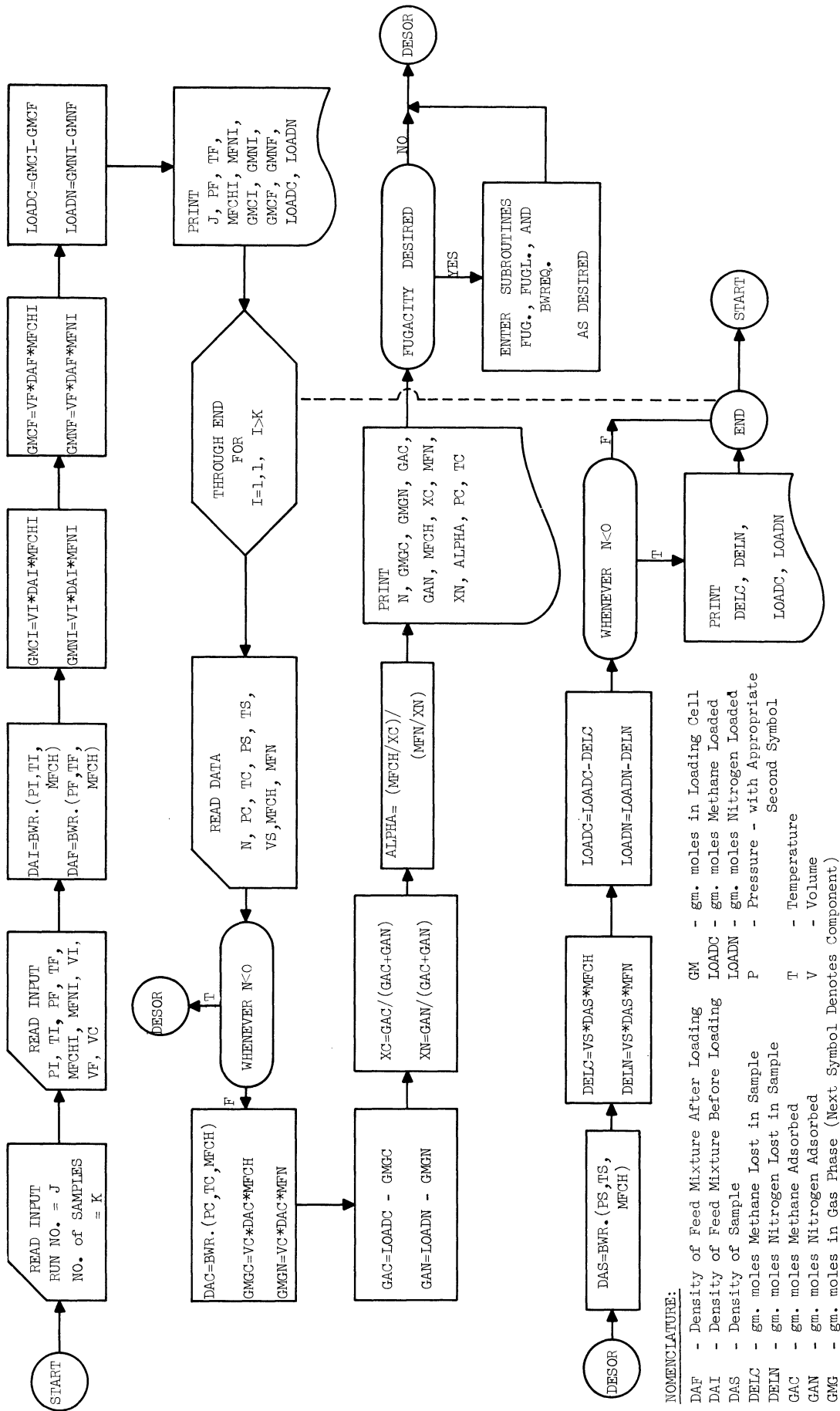
2. Calculation

$$\text{DAS} = 0.000034 \text{ gm. moles/cc.}$$

$$\begin{aligned}\text{Desorbed DELC} &= \text{VS} \times \text{DAS} \times \text{MFCH} \\ &= 11,780 \times 3.4 \times 10^{-5} \times 0.6949 \\ &= 0.2805 \text{ gm. moles} \\ \text{DELN} &= 11,780 \times 3.4 \times 10^{-6} \times 0.3051 \\ &= 0.1231 \text{ gm. moles}\end{aligned}$$

Net Remaining in System

$$\begin{aligned}\text{Methane LOADC}_{i-1} &= \text{LOADC}_i - \text{DELC} \\ &= 0.34188 - 0.2805 \\ &= 0.0613 \text{ gm. moles}\end{aligned}$$



**NOMENCLATURE:**

- DAF - Density of Feed Mixture After Loading
- DAI - Density of Feed Mixture Before Loading
- DAS - Density of Sample
- DELC - gm. moles Methane Lost in Sample
- DELN - gm. moles Nitrogen Lost in Sample
- GAC - gm. moles Methane Adsorbed
- GAN - gm. moles Nitrogen Adsorbed
- GMG - gm. moles in Gas Phase (Next Symbol Denotes Component)
- GM - gm. moles in Loading Cell
- LOADC - gm. moles Methane Loaded
- LOADN - gm. moles Nitrogen Loaded
- P - Pressure - with Appropriate Second Symbol
- T - Temperature
- V - Volume

Figure 30. Computer Flow Diagram for Equilibrium Adsorption Loading Calculation.

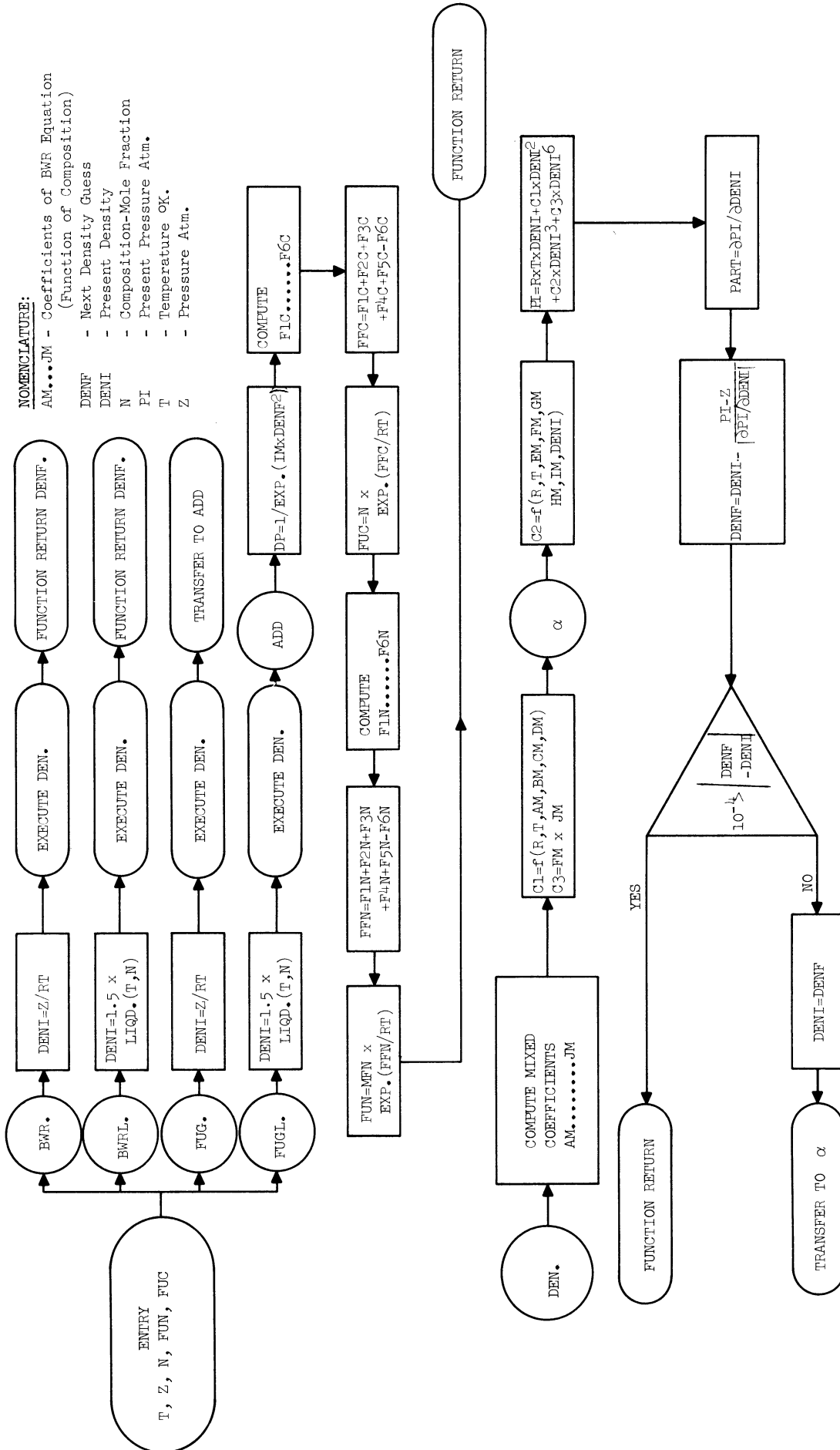


Figure 31. Computer Flow Diagram for B-W-R Equation Solution for Density and Fugacity



TABLE XXVII

SAMPLE PRINT OUT FROM COMPUTER CALCULATION

ADSORPTIVE EQUILIBRIUM CH4-N2

RUN NO. 76

INITIAL CONDITIONS

INT.PRESSE= 10.59 ATM. INT. TEMP = 300.4 KELVIN

M.FRAC.CH4=0.7257 M.FRAC.N2 = 0.2743

G.MOL1CH4= 0.64906 G.MOL1.N2 = 0.24933 G.MOL2.CH4= 0.30718 G.MOL2.N2 = 0.11611

LOADED G.MOL.CH4= 0.34188 G.LOL.N2 = 0.12922

SAMP NO	G.MOLES CH4	G.MOLES N2	MOLE FRAC. CH4	MOLE FRAC. N2	ALPHA
1	0.024476	0.317400	0.073626	0.926374	0.315
	ADSORB	ADSORB	ADSORB	ADSORB	
	0.024476	0.317400	0.073626	0.926374	0.315
	GAS	GAS	GAS	GAS	
	0.024476	0.317400	0.073626	0.926374	0.315
	1 FUGACITY-GAS PHASE CH4 =	3.47222	N2 =	3.77847	7.60
	VAR. DENSITY=	0.55584	GM/L		174.9
				TOTAL	7.25561

SAT. VALUES

CH4-TEMP= 143.66409 DENSITY= 23.81894 GM/L

N2-TEMP= 99.87026 DENSITY= 24.84852 GM/L

MIXTURE DENSITY= 24.06463 GM/L

SAT. VAR. PRESSE. CH4= 28.14911 N2= 156.57950 MIXTURE= 93.54937 ATM

SAT. DENSITY= 14.76709GM/L

FSCCH4= 14.35203 FSCN2= 32.61521 FSCMIX= 43.13835

N\*VCPD= 0.18769737 N\*VCH4= 0.15065651 N\*VCH2= 0.04724088 CC/GRAM

LNCFS/FD/VE= 0.04583569 GMOL/CC T\*LNCFS/FD/VE= 8.01528692 K GMOL/CC

LNCSUMPS/SUMPD/VE= 0.04785440

-2	GMOL.CH4 DESORBED=	0.280508	GMOL. N2 DESORBED=	0.123158
	GMOL. IN CELL-CH4=	0.060857	N2=	0.005533

-3	GMOL.CH4 DESORBED=	0.058192	GMOL. N2 DESORBED=	0.000000
	GMOL. IN CELL-CH4=	0.002675	N2=	0.005533

-4	GMOL.CH4 DESORBED=	0.004074	GMOL. N2 DESORBED=	0.000000
	GMOL. IN CELL-CH4=	-0.001359	N2=	0.005533

## APPENDIX F

## ADSORPTION OF HELIUM

During the early exploratory work, a number of runs was made to determine the adsorption of helium at higher pressures. Several runs were made at room temperature as well as one at  $-68^{\circ}\text{C}$  and one at  $-197^{\circ}\text{C}$ . The results are listed in Table XXVII below.

TABLE XXVIII  
ADSORPTION OF HELIUM

Temp. $^{\circ}\text{C}$	Pressure ATM	Amount Adsorbed milligram moles/gm.
21	71.45	0.0
21	41.17	0.0
-68	44.2	2.0
-197	27.2	2.9

The experimental procedure was identical to that used in determining methane and nitrogen equilibrium.

The calculation of the equilibrium amount adsorbed was similar to that outlined in Appendix E which shows a sample calculation for methane-nitrogen adsorption. The P-V-T behavior of helium at room temperature was represented by the Beattie-Bridgman equation with appropriate constants. At the lower temperatures,  $-68^{\circ}\text{C}$  and  $-197^{\circ}\text{C}$ , the reduced state correlation of Hamrin and Thodos<sup>(23)</sup> was used to obtain the gas phase density of helium.

One determination of the adsorption equilibria of a mixture of nitrogen and helium was made at room temperature. The data from this run, though inconclusive, indicate that the helium acts as a diluent at this temperature. The amount adsorbed was, however, about fifteen per cent

greater than would be expected if one assumes an ideal mixture between the two components.

The results obtained in this exploratory study are included to serve as a guide for future work as the helium-nitrogen system is of particular interest today from a conservation standpoint.

## BIBLIOGRAPHY

1. Antropff, von A. Kolloid Z., 129, 11 (1952).
2. "Boiler and Pressure Vessel Code." ASME, Sect. VIII, 1956ed., 7-12, 109-113.
3. Baly, E. C. C. Proc. Roy. Soc., A160, 465 (1937).
4. Barrer, R. M. J. Soc. Chem. Ind., 64, 130 (1945).
5. Bartlett, E. P. J. Am. Chem. Soc., 49, 687 (1927).
6. Benedict, M., Webb, G. B. and Rubin, L. C. J. Chem. Phys., 8, 334 (1940).
7. Ibid, 10, 747 (1942).
8. Berényi, L. L. Z. Physik. Chem., 94, 628 (1920).
9. Berényi, L. L. Z. Angw. Chem., 35, 237 (1922).
10. Bloomer, O. T., Parent, J. D. I.G.T. Bull. No. 17, 1952.
11. Bloomer, O. T., Roa, K. N. I.G.T. Bull. No. 18, 1952.
12. Breck, D. W., Eversole, W. G., Milton, R. M., Reed, T. B., and Thomas, T. L. J. Am. Chem. Soc., 78, 5963 (1956).
13. Brunauer, S. The Adsorption of Gases and Vapors, Vol. I, Princeton Univ., Princeton, 1945.
14. Brunauer, S., Demming, L. S., Demming, W. E. and Teller, E. J. Am. Chem. Soc., 62, 1723-32 (1940).
15. Brunauer, S., Emmet, P. and Teller, E. J. Am. Chem. Soc. 60, 309 (1938).
16. Coolidge, A. S. J. Am. Chem. Soc., 46, 596 (1924).
17. Darby, H. S. Pressure - Volume - Temperature Relationship of Methane - Nitrogen Mixtures at Low Temperatures, Unpublished Master's Thesis, I.G.T., 1951.
18. Dubinin, M. M. Chem. Rev., 60, 235 (1960).
19. Ellington, R. T., Bloomer, O. T., Eakin, B. E. and Gami, D. C. "Thermodynamic and Transport Properties of Gases, Liquids and Solids," ASME, New York, 102 (1959).

20. Fowler, R. H. and Guggenheim, E. A. Statistical Thermodynamics, Macmillan, New York, 1939.
21. Freundlich, H. Colloid and Capillary Chemistry, Methuen and Co., London, 1926.
22. Goldman, F. and Polanyi, M. Z. Physik. Chem., A132, 321 (1928).
23. Hamrin, C. E. Jr. and Thodos, G. AIChE Journal, 4, 480 (1958).
24. Hill, T. L. J. Chem. Phys., 14, 263 (1946).
25. Ibid., 268-7.
26. Hill, T. J. Am. Chem. Soc., 72, 5347 (1950).
27. Hiza, M. J. Chem. Eng. Progr., 56, No. 10, 68 (1960).
28. Honig, J. M. Anal. N. Y. Acad. Sci., 58, 741 (1958).
29. Johnson, V. J. Adv. in Cryogenic Eng., 3, Paper A-2, Plenum Press, New York, 1960.
30. Josefowitz, S. and Othmer, D. F. Ind. Eng. Chem., 40, 739 (1948).
31. Joyner, L. G., Weinberger, E. B. and Montgomery, C. W. J. Am. Chem. Soc., 67, 2182 (1945).
32. Keyes, F. G. and Burks, H. G. J. Am. Chem. Soc., 50, 1100 (1928).
33. Langmuir, I. Phys. Rev., 8, 149 (1916).
34. Langmuir, I. J. Am. Chem. Soc., 40, 1361 (1918).
35. Lee, V. J. Unpublished Paper, University of Michigan, 1959.
36. Lewis, W. K., Gilliland, E. R., Chertow, B. and Codogan, W. P. Ind. Eng. Chem., 42, 1319 (1950).
37. Ibid., 1326 (1950).
38. McBain, J. W. and Bakr, A. M. Jr. J. Am. Chem. Soc., 48, 690 (1926).
39. McBain, J. W. and Britton, G. T. J. Am. Chem. Soc., 52, 2198 (1930).
40. Markham, E. D. and Benton, A. F. J. Am. Chem. Soc., 53, 497 (1931).
41. Milne, W. E. Numerical Calculus, Princeton University, Princeton, 1949.
42. Morris, H. E. and Mass, O. Can. J. Res., 9, 240 (1923).

43. Nikolaev, K. M. and Dubinin, M. M. Izvest. Akad. Nauk SSSR, Otdel. Kheim. Nauk, 1165 (1958).
44. Pace, J. C. Compressibility Factors of Nitrogen, Methane, and Three of Their Mixtures. Unpublished Master's Thesis, I. G. T., 1951.
45. Pease, R. N. J. Am. Chem. Soc., 45, 1196 (1923).
46. Peierls, R. Proc. Camb. Phil. Soc., 32, 471 (1936).
47. Polanyi, M. Ber. Deut. Physik. Ges., 16, 1012 (1914).
48. Polanyi, M. Ver. Deut. Physik. Ges., 18, 55 (1916).
49. Polanyi, M. Z. Electrochemie, 26, 370 (1920).
50. Polanyi, M. Z. Physik. Chem., A138, 459 (1928).
51. Rabe, A. and Thomas, T. L. Private Communication.
52. Redlich, O. and Peterson, D. L. J. Phys. Chem., 63, 1024 (1959).
53. Reed, T. B. and Breck, D. W. J. Am. Chem. Soc., 78, 5972 (1956).
54. Rideal, E. K. Surface Chemistry, Cambridge University, Cambridge, 1930.
55. Sault, R. A. Compressibility of Certain Methane - Nitrogen Mixtures. Unpublished Master's Thesis, I.G.T., 1952.
56. Schay, G. J. Chem. Phys., (Hungary), 53, 691 (1956).
57. Schay, G. Fejes, P. and Szethmary, J., Acta Chim. Aca. Sci., Hungary, 12, 299 (1957).
58. Sips, R. J. Chem. Phys., 16, 490 (1948).
59. Ibid., 18, 1024 (1950).
60. Tarpnell, B. M. W. Chemisorption. Butterworth, London, 1955.
61. Wang, J. Proc. Roy. Soc., A161, 127 (1938).
62. Williams, A. M. Proc. Roy. Soc., Edinburg, 38, 23 (1917-18).
63. Zeldowitsh, J. Acta Physiochim. USSR, 1, 961 (1934).