

where $u_s(\mathbf{r}_1)$ is the gas-solid energy and $u(r_{ij})$ the mutual interaction energy of the gas molecules with one another. The volume elements are $d\mathbf{r}_1$ and $d\mathbf{r}_2$ if molecules 1 and 2 are located at \mathbf{r}_1 and \mathbf{r}_2 respectively.

Gregg and Sing² have discussed the shape of the adsorption isotherm by means of an analysis of the point of inflection in the BET equation.³ By applying the same ideas to Eq. (1) it is easy to deduce that the pressure P_i corresponding to the point of inflection of the adsorption isotherm is given by

$$P_i = -kT C_{AAS} / 3 D_{AAAS} \quad (5)$$

while the corresponding number of adsorbed molecules, $N_a(i)$ is, in a first approximation, given by

$$N_a(i) = -C_{AAS} B_{AS} / 3 D_{AAAS} \quad (6)$$

These equations allow some considerations to be made in relation to physisorption isotherms of types II and III according to the classical classification.⁴ Since type III isotherms are those that do not show a point of inflection, according to Eq. (5) they would correspond to the case in which P_i is negative, which is true only if C_{AAS} and D_{AAAS} have the same sign. This has been confirmed in our laboratory for some systems whose experimental data we have published previously.^{5,6} Physically this argument is related to the existence of characteristic temperatures observed in adsorption systems.⁷

It can then be said that the existence of a type III isotherm implies the same form of energy contribution between structure AAS and AAAS, while that of type II implies the opposite.

Since P_i is not a function of B_{AS} , it can be seen that the shape of the isotherm is independent of the AS configuration. However, Eq. (6) shows that, when the pressure is P_i , the number of adsorbed molecules is a function of the three virial coefficients. On the other hand, even though $N_a(i)$ does not belong to the monolayer, it may correspond, at least in the majority of type II isotherms, to a first approximation of the value.²

¹W. A. Steele, *The Interaction of Gases with Solid Surfaces* (Pergamon, Oxford, England, 1974).

²S. J. Gregg and K. S. Sing, *Adsorption Surface Area and Porosity* (Academic, New York, 1967).

³S. Brunauer, P. H. Emmett, and E. Teller, *J. Am. Chem. Soc.* **60**, 309 (1938).

⁴S. Brunauer, *The Adsorption of Gases and Vapours* (Clarendon, Oxford, 1945).

⁵L. Alzamora, S. Contreras, and J. Cortés, *J. Colloid Interface Sci.* **50**, 503 (1975).

⁶A. Tornquist, E. Valencia, L. Alzamora, and J. Cortés, *J. Colloid Interface Sci.* **66**, 415 (1978).

⁷J. R. Sams, Jr., G. Constabaris, and G. D. Halsey, Jr., *J. Chem. Phys.* **36**, 1334 (1962).

ERRATA

Erratum: Representations of molecular force fields. V. On the equilibrium structure of methane [J. Chem. Phys. 68, 1213 (1978)]

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Expressions for the stretch-bend-bend cubic constants should be amended to read:

$$\tilde{f}_{122} = r_0(-F' + F + F_3)/6$$

and

$$\tilde{f}_{144} = r_0(-5F' - 3F + F_3)/6.$$

Associated numerical values of the methane cubic constants are changed from those reported by less than 0.2 mdyn/Å². The corrections ($\langle r \rangle - \langle r_p \rangle$) of Table II should be increased by 0.0004 to 0.0005 Å for the different isotopic species and different parameterizations (KB or MUB-2).