

The dependence of H II region properties on global and local surface brightness within galaxy discs

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ABSTRACT

Using B , R and $H\alpha$ images of roughly equal-sized samples of low surface brightness (LSB) and high surface brightness (HSB) galaxies (~ 40 galaxies apiece), we have explored the dependence of H II region properties on local and global disc surface brightness. We have done this by constructing co-added H II region luminosity functions (LFs) according to local and central disc surface brightness and fitting Schechter functions to these LFs. The results show that the shape of the H II region LF within LSB galaxies does not change noticeably as different limiting (i.e. $\mu > \mu_{\text{lim}}$) local surface brightness values are used. However, the LFs for HSB galaxies have larger values of L^* and are less steep at the faint-end than those of LSB galaxies for limiting B -band local surface brightness values as faint as $\mu_{B,\text{lim}} \simeq 23\text{--}24$. Both the LFs and the data for individual H II regions show that luminous ($L > 10^{39}$ erg s $^{-1}$) H II regions are much more common within HSB discs than within LSB discs, implying that the newly formed star clusters are also larger. Taking this into account along with the results of Monte Carlo simulations, the shapes of the LFs imply that the regions within LSB discs and those within the LSB areas of HSB discs are relatively old (~ 5 Myr) while the regions within HSB discs for $\mu_B \lesssim 24$ are significantly younger (< 1 Myr). Since the majority of the LSB galaxies do not have noticeable spiral arms and the majority of the HSB galaxies do, this may indicate a transition within HSB discs from spiral arm-driven star formation to a more locally driven, possibly sporadic form of star formation at $\mu_B \sim 24$, a transition that does not appear to occur within LSB discs.

Key words: galaxies: ISM – galaxies: photometry – galaxies: star clusters.

1 INTRODUCTION

The completeness of the pioneering optical surveys of galaxies suffered from the inevitable bias towards higher surface brightness produced by the intensity of the night sky. Digital detectors and data reduction techniques used in modern times have weakened the effect of this bias. Because of this, the study of so-called low surface brightness (LSB) galaxies has become its own subfield within extragalactic astronomy. Results from this study have since shown LSB galaxies to be larger at the same luminosity as high surface brightness (HSB) galaxies (de Blok, McGaugh & van der Hulst 1996) while still obeying the same Tully–Fisher relationship, but only if the entire baryonic content of the galaxies is taken into account (McGaugh et al. 2000). LSB discs also tend to be less dense

(de Blok & van der Hulst 1998) and have higher gas mass fractions than HSB discs (McGaugh & de Blok 1997). In addition, few LSB galaxies have detectable amounts of molecular gas, despite several searches (O’Neil, Hofner & Schinnerer 2000; Matthews & Gao 2001; O’Neil, Schinnerer & Hofner 2003; O’Neil et al. 2004; Matthews et al. 2005). LSB galaxies also tend to have lower metallicities ($\sim 1/3$ solar; McGaugh 1994) and are more frequently dwarf and/or irregular galaxies without noticeable spiral arms than HSB galaxies (Helmboldt et al. 2004). In addition to this, the H I column densities of LSB galaxies tend to be relatively low, at or below the critical limit for star formation (Kennicutt 1989; de Blok et al. 1996). In light of this, it is not surprising that the star formation rates (SFRs) for LSB galaxies are typically about 10 times lower than those of HSB galaxies (Mihos, Spaans & McGaugh 1999; Helmboldt et al. 2004). Despite this, LSB discs tend to be relatively blue and some have relatively large $H\alpha$ emission line equivalent widths (see, e.g. Helmboldt et al. 2004).

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These properties of LSB galaxies imply that despite conditions that are not conducive to star formation, many LSB galaxies are actively forming stars. Because of this, it is possible that they do not form stars in the same way as typical spiral galaxies. In fact, it has recently been shown that LSB galaxies tend to form fewer high luminosity H II regions and have larger relative amounts of diffuse ionized gas than HSB galaxies (Helmboldt et al. 2005; Oey et al. 2007; O’Neil, Oey & Bothun 2007). While the larger amounts of diffuse ionized gas may reflect differences in the conditions of the interstellar medium within LSB galaxies, the fact that the distribution of H II region luminosities among LSB galaxies is markedly different from that for HSB galaxies points to a possible fundamental difference between how star clusters are formed within these two types of galaxies. It is remarkable that despite the differences in their star formation histories, LSB galaxies still manage to form roughly the same number of stars at a given circular velocity as HSB galaxies.

From recent results, it is not clear if this difference in star formation history results from global or local mechanisms. In other words, does the lack of spiral arms cause star clusters to form differently throughout a typical LSB disc than throughout an HSB disc or do star clusters form in a similar way in the LSB regions of HSB discs as they do throughout LSB discs? To address this question, we have used previously published and newly processed data to compile two roughly equal-sized samples of HSB and LSB galaxies with B , R and $H\alpha$ images. We present here a description of the old and new data (Section 2), the results of a detailed analysis of the H II region properties within these two groups of galaxies and how they depend on both local and global properties (Section 3), and Monte Carlo simulations constructed to help interpret these results (Section 4).

2 GALAXY SELECTION AND OBSERVATIONS

To compare properties of H II regions within HSB and LSB galaxies, we have compiled broad-band B and R images and narrow-band $H\alpha$ images of 89 nearby ($V_r \lesssim 6000 \text{ km s}^{-1}$) disc galaxies, 45 HSB galaxies and 44 LSB galaxies. Based on the so-called Freeman value of $\mu_{0,B} = 21.65 \pm 0.3$ for ‘normal’ disc galaxies (Freeman 1970), we have adopted a definition for HSB galaxies of $\mu_{0,B} < 21.7$. Our working definition of an LSB disc galaxy is any galaxy with a B -band central disc surface brightness $> 22 \text{ mag arcsec}^{-2}$. This is one magnitude brighter than what is typically used (see, e.g. Bothun, Impey & McGaugh 1997; Impey & Bothun 1997), however, it is still more than 1σ fainter than the Freeman value. In addition, as several authors have demonstrated, the distribution of surface brightness among galaxies is continuous and not bimodal (e.g. Jansen 2000; Blanton et al. 2003; Helmboldt et al. 2004), implying this definition is somewhat arbitrary. We therefore chose a limit significantly lower than the Freeman value for normal disc galaxies that also yielded a sample of LSB galaxies similar in number to our sample of HSB galaxies.

2.1 Previously published data

The images were obtained as part of two separate observing runs in 2000 May and October conducted at Kitt Peak National Observatory (KPNO) and Cerro Tololo Interamerican Observatory (CTIO), respectively. The CTIO run consisted of imaging 69 galaxies selected from the H I Parkes All Sky Survey (HiPASS; Barnes et al. 2001) with $V_r < 2500 \text{ km s}^{-1}$. All 45 HSB galaxies were obtained from this sample as well as 20 of the LSB galaxies. The reduction and analysis of the CTIO data was presented in detail by Helmboldt

et al. (2004) and Helmboldt et al. (2005), and we will briefly discuss them here. The images were calibrated using observations of standard stars throughout each night. For the narrow-band observations during the last two nights of the run, this calibration scheme did not work and a relation was derived using the calibration solutions for the other $H\alpha$ and R -band observations given by the following

$$\Delta m_{H\alpha} = \Delta m_R + 2.5 \log \left(C a \frac{\tau_R}{\tau_{H\alpha}} \right) + 0.5, \quad (1)$$

where Δm is the difference between the instrumental magnitude and the actual magnitude, τ is the exposure time in a particular filter, a is the scalefactor each R -band image was multiplied by to be used to subtract the continuum emission from the $H\alpha$ image, and C is a constant empirically determined to be 0.73 ± 0.15 . Here, we define the instrumental magnitude to be $-2.5 \log \text{DN}/\tau + 25$ where Data Numbers (DN) is the flux measured from the image in Analog-to-Digital Units (ADUs).

The $25 \text{ mag arcsec}^{-2}$ isophote was then identified on each calibrated B -band image and an ellipse was fit to it, except for one galaxy for which the $26 \text{ mag arcsec}^{-2}$ isophote was used because it had $\mu_{0,B} \sim 25$. The position, ellipticity and position angle of the fitted ellipse were then used to construct elliptical apertures within which the median B and R -band surface brightnesses were measured to estimate the surface brightness profile of each galaxy. The B and R profiles were then simultaneously fit with an exponential disc model for all radii where the surface brightness was between 0.3 and 2.472 mag above the 1σ limiting isophote. This magnitude range was chosen because for an exponential disc, this corresponds to a span of two disc scalelengths. Each fit was performed with a standard weighted linear least squares routine to obtain the central surface brightness values in each band and a single value for the disc scalelength, h , as well as error estimates for all three quantities. The fits were then used to obtain total extrapolated fluxes by integrating each exponential fit from the radius at which the surface brightness is 2.172 mag above the 1σ limiting isophote to infinity, and then adding this flux to the total flux within that radius measured from the image.

Continuum subtracted $H\alpha$ images were made by scaling the R -band images so that the fluxes of several stars on the R -band and $H\alpha$ images matched, on average, and subtracting the scaled R -band image from the $H\alpha$ image. The $H\alpha$ filter transmission curve and the galaxy redshift were then used to determine the $H\alpha + [\text{N II}]$ flux from each calibrated, continuum subtracted $H\alpha$ image. Using the emission line measurements of Tremonti et al. (2004) from the Sloan Digital Sky Survey (SDSS; York et al. 2000) galaxy spectra, we have derived relationships between $B - R$ colour index and the ratios of the $H\alpha$ flux and equivalent width to the observed $H\alpha + [\text{N II}]$ flux and equivalent width. The Tremonti et al. (2004) measurements included a linear combination of model spectra fit to each galaxy’s spectrum so that the effects of stellar absorption would be fully accounted for. We also corrected both the $H\alpha$ flux and equivalent width measured by Tremonti et al. (2004) for internal dust extinction according to Calzetti (2001) using the ratio of $H\alpha$ to $H\beta$ emission line flux to compute $E(B - V)$ for the ionized gas, assuming $H\alpha/H\beta = 2.87$ (for $T = 10^4 \text{ K}$ and case B recombination; Osterbrock & Ferland 2006). Using emission line ratios, we identified 12 262 nearby ($m_r < 16$) actively star forming galaxies according to Helmboldt, Walterbos & Goto (2008) from the SDSS data of Tremonti et al. (2004) and have measured $B - R$ indices from their spectra. In Fig. 1, we have plotted the ratios of the $H\alpha$ flux and equivalent width to the observed $H\alpha + [\text{N II}]$ flux and equivalent width versus $B - R$ for these galaxies with median values plotted in bins of $B - R$.

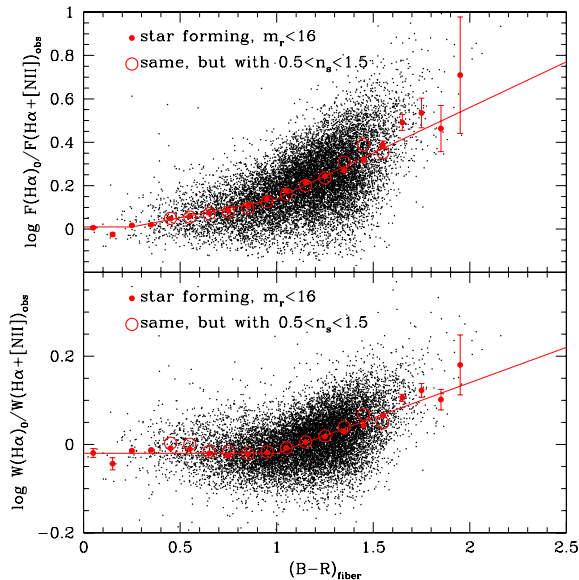


Figure 1. The ratio of the $H\alpha$ flux (upper) and equivalent width (lower) to the observed $H\alpha+[N II]$ flux for 12 262 star forming SDSS galaxies as measured by Tremonti et al. (2004), including corrections for internal dust extinction according to Calzetti (2001), versus $B - R$ colour index as measured from the SDSS spectra (see Section 2.1). Also, plotted are median values within bins of $B - R$ and broken power-law fits to the data [not to the median values; see equation (2)]. Median values for galaxies with Sersic indices from Blanton et al. (2003) between 0.5 and 1.5 (i.e. disc-dominated galaxies) are also plotted as large, open red circles.

Also, plotted are broken power-law functions fitted to the data (not the median values) which have the following form

$$\begin{aligned}
 X &= (0.42 \pm 0.066)(B - R) + (-0.28 \pm 0.085) \\
 &\quad \text{if } (B - R) > 0.992, \\
 &= (0.17 \pm 0.11)(B - R) + (-0.032 \pm 0.11) \\
 &\quad \text{if } 0.992 \geq (B - R) > 0.25, \\
 &= (0.011 \pm 0.11) \text{ if } (B - R) \leq 0.25,
 \end{aligned} \tag{2}$$

$$\begin{aligned}
 Y &= (0.16 \pm 0.066)(B - R) + (-0.18 \pm 0.084) \\
 &\quad \text{if } (B - R) > 1, \\
 &= (-0.020 \pm 0.086) \text{ if } (B - R) \leq 1,
 \end{aligned} \tag{3}$$

where $X = \log F(H\alpha)_0 / F(H\alpha + [N II])_{\text{obs}}$ and $Y = \log W(H\alpha)_0 / W(H\alpha + [N II])_{\text{obs}}$. We must note that the SDSS spectra were obtained using a fibre-fed spectrograph with each fibre having a 3 arcsec aperture. This implies that many of the spectra may be bulge dominated and the above corrections may be more appropriately applied to nuclear regions of galaxies. To check this, we obtained Sersic indices, n_s , for the SDSS galaxies from Blanton et al. (2003) and recomputed the median values plotted in Fig. 1 for galaxies with $0.5 < n_s < 1.5$, that is likely disc-dominated galaxies. These new median values are plotted in Fig. 1 as large, open red circles and are not significantly different from the median values computed for the entire sample of star forming galaxies. We therefore conclude that it is reasonable to apply the corrections given in equations (2) and (3) to galaxy discs based on their $B - R$ colours and not on integrated or nuclear colour indices.

We found no significant correlation between the ratios plotted in Fig. 1 and either absolute magnitude or surface brightness. We also note that different empirically derived corrections were used for these data by Helmboldt et al. (2004) and Helmboldt et al. (2005).

However, those previous relations did not include the effects of stellar absorption, and we have therefore concluded that equations (2) and (3) provide better, more extensive corrections for the $H\alpha+[N II]$ data. These new empirical relations with $B - R$ colour were used to convert the measured $H\alpha+[N II]$ fluxes to extinction corrected $H\alpha$ fluxes using the $B - R$ colour of each galaxy's disc computed using the exponential fits to the B and R surface brightness profiles described above [i.e. $(B - R)_{\text{disc}} = \mu_{0,B} - \mu_{0,R}$]. We used the $(B - R)_{\text{disc}}$ colours because (1) the vast majority of the star formation within these galaxies occurs within their discs and (2) the exponential fits yield $B - R$ colour estimates for the discs that are reasonably free from the effects of any substantial small scale colour fluctuations and bulge or bar contamination.

To measure the properties of the $H II$ regions, the HIIphot algorithm developed by Thilker, Braun & Walterbos (2000) was used. The algorithm is designed to identify $H II$ regions by smoothing the $H\alpha$ images with kernels of different sizes and fitting models to significant peaks in the smoothed images. The boundary of each $H II$ region is then grown until a limiting emission measure gradient is reached (a limit of 1.5 EM pc^{-1} was used) or the boundary reaches the boundary of another region. Following this, surface fits to designated background pixels are used to estimate the amount of diffuse emission along the lines of sight passing through each $H II$ region (see Thilker et al. 2000; Helmboldt et al. 2005, for more detailed descriptions).

2.2 Newly processed data

To supplement the imaging data for LSB galaxies available from the CTIO observing run, a sample of 32 galaxies previously identified as LSB disc galaxies was observed at KPNO. The 32 targets were culled from the literature (Schommer & Bothun 1983; Bothun, Beers & Mould 1985; Kraan-Korteweg 1986; Elmegreen & Elmegreen 1987; Schombert & Bothun 1988; Tully & Fisher 1988; Schneider et al. 1990; O'Neil, Bothun & Cornell 1997) and were chosen to represent the wide variety of known LSB disc galaxies, ranging from dwarf (disc scalelengths ~ 500 pc) to giant (scalelengths ~ 10 kpc) LSB galaxies. However, because of cloudy conditions during the observing run, we did not perform standard star observations and the usual method of calibrating the images could not be used. In the intervening years, the completion of the bulk of the imaging for the SDSS has provided a means to calibrate the B and R band images of 24 of the LSB galaxies observed. We have done this by measuring the fluxes of several ($\sim 6-10$) stars on the KPNO B and R images which have measured g and r magnitudes from SDSS images. The g and r magnitudes were converted to B and R magnitudes using the conversions given by Smith et al. (2002), and the mean difference between the raw KPNO fluxes and these magnitudes was used to calibrate each image. The typical uncertainty in these calibration solutions is about 0.06 and 0.02 mag for the B and R bands, respectively.

There are 12 galaxies from our KPNO sample with B -band magnitudes from the RC3 catalogue. As Fig. 2 demonstrates, our total B -band magnitudes (see below) agree with the RC3 values to within about 1σ on average and that the difference between the two is not systematic and does not depend on the brightness of the galaxy. Note that the 1σ uncertainties in our total B -band magnitudes are combinations of the calibration errors, the photon errors and the errors in the parameters from the fits to the surface brightness profiles used to compute fluxes integrated to radii of ∞ (see below).

For consistency, the surface photometry for these new images was performed in the same way as the CTIO images with one

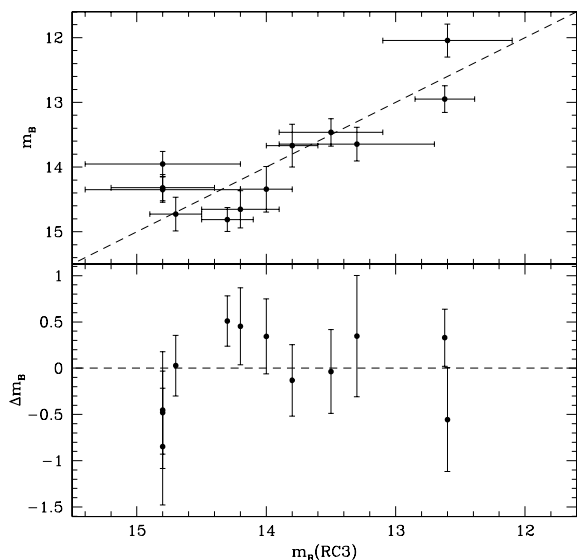


Figure 2. In the upper panel, the total extrapolated B -band magnitude (see Section 2.1) for the 12 LSB galaxies observed at KPNO that are in the RC3 catalogue versus the RC3 B -band magnitudes. The difference between the B -band magnitude measured using our KPNO images and the RC3 magnitude is plotted versus the RC3 magnitude in the lower panel. In each panel, the dashed line represents the case where the magnitudes are the same and is not a fit to the data.

important difference. Since many of these galaxies have relatively low central surface brightnesses, the B -band 25 mag arcsec $^{-2}$ isophote was not ideal to use to characterize the shape of each galaxy if it could be identified at all. Instead, a different isophote was used for each galaxy, determined by eye through trial and error by varying the isophote level in 0.5 mag steps. The isophotes used are listed in Table 1 and range from 26 to 27 mag arcsec $^{-2}$. We have also listed in Table 1 the ellipticity and position angle of this isophote along with the exponential disc scalelength, the B and R central surface brightnesses, and the total extrapolated B and R magnitudes, all determined in the same way as the CTIO data (see above and Helmboldt et al. 2004). The surface brightness profiles for all 24 galaxies are displayed in Fig. 4 along with the exponential disc fits. The $\mu_{0,B}$ distributions plotted in the left-hand panel of Fig. 3 show that the inclusion of the newly processed KPNO data not only doubles the number of LSB galaxies, it also provides a much better sampling of the $23.5 < \mu_{0,B} < 24.5$ region. In the right-hand panel of Fig. 3, we have plotted the distributions of $(B - R)_{\text{disc}}$. These distributions show that the inclusion of the new KPNO data has also introduced a significant number of relatively red LSB discs that supplement the relatively blue LSB discs found among the CTIO data. These redder discs have made the $(B - R)_{\text{disc}}$ distributions for the HSB and LSB galaxies essentially the same.

Without standard star observations or spectra, we have no direct means of calibrating the $H\alpha$ images obtained for these galaxies. However, for all but five of the galaxies observed a KPNO, the same set of 30 Å wide filters, kindly provided by Richard Rand, was employed as was used for most of the CTIO observations. For

Table 1. Properties of LSB Galaxies observed at KPNO.

Name	R.A. (J2000)	Dec. (J2000)	V_r	e	PA	h	$\mu_{0,B}$	$\mu_{0,R}$	m_B	m_R	$\log F(H\alpha)$	$\log W(H\alpha)$	
(1)	(h m s)	($^{\circ}$ ' '')	(km s $^{-1}$)	(5)	($^{\circ}$)	('')	(9)	(10)	(11)	(12)	(erg s $^{-1}$ cm $^{-2}$)	(Å)	
NGC 3447	10 53 26.8	+16 46 44.0	1067	0.32	-73.0	26.0	27.2	22.5	21.6	13.1	12.3	-11.9	1.71
UGC 6151	11 05 56.3	+19 49 31.0	1335	0.18	88.0	27.0	18.2	23.2	22.3	14.8	13.9	-13.0	1.33
U1-4	11 38 25.8	+17 05 03.1	3452	0.38	69.9	26.0	3.4	22.1	21.1	17.3	16.3	-13.7	1.11
NGC 4688	12 47 46.5	+04 20 09.9	986	0.16	39.7	26.5	26.6	22.3	21.4	12.9	12.1	-11.9	1.49
LSBC F574-10	12 50 31.4	+17 28 18.7	844	0.72	-74.5	26.0	6.1	22.9	22.1	16.8	16.1	-14.0	0.94
UGC 8011	12 52 21.1	+21 37 46.2	776	0.08	-12.2	27.0	19.7	23.9	23.0	15.3	14.4	-13.2	1.43
UGC 8061	12 56 44.0	+11 55 55.3	562	0.17	-86.0	27.0	16.6	23.8	22.5	15.5	14.3	-13.5	1.22
UGC 8091	12 58 40.4	+14 13 02.9	214	0.39	-74.1	26.5	16.2	23.0	22.2	14.7	13.9	-12.0	2.16
UGC 8155	13 03 14.7	+07 48 07.9	2925	0.24	-86.2	27.0	30.8	23.5	22.2	13.6	12.4	-12.5	1.07
UGC 8441	13 25 29.1	+57 49 20.0	1519	0.17	1.6	27.0	27.2	23.6	22.6	14.3	13.4	-12.8	1.31
NGC 5409	14 01 46.1	+09 29 25.0	6259	0.01	12.1	26.0	23.8	22.9	21.5	13.7	12.3
UGC 9024	14 06 40.6	+22 04 12.1	2323	0.16	-26.8	27.0	13.5	23.6	22.6	15.6	14.6	-13.3	1.26
UGC 9245	14 25 26.9	+56 19 12.6	1885	0.30	-79.3	27.0	23.0	23.3	22.2	14.3	13.3	-12.7	1.22
UGC 9380	14 34 39.2	+04 15 46.0	1693	0.32	-66.5	27.0	20.2	23.3	22.5	14.7	13.9	-12.8	1.50
NGC 5866B	15 12 07.2	+55 47 06.3	841	0.45	-82.5	26.5	24.3	23.0	21.9	14.0	12.8	-12.7	1.26
LSBC F583-2	15 37 00.5	+20 47 42.0	1719	0.36	38.8	27.0	11.3	23.9	22.9	16.4	15.5	-13.7	1.50
NGC 5964	15 37 36.3	+05 58 24.3	1447	0.09	-27.4	26.0	43.2	22.5	21.4	12.0	11.0	-12.1	1.01
LSBC F583-5	15 45 43.9	+17 18 41.0	3261	0.34	28.4	27.0	4.4	23.3	22.4	18.0	17.1	-13.8	1.69
UGC 10041	15 49 01.4	+05 11 19.0	2171	0.62	89.0	26.5	24.4	22.8	21.8	13.5	12.5	-12.3	1.36
UGC 10061	15 51 15.1	+16 19 46.4	2080	0.34	88.6	27.0	36.1	24.3	23.4	14.3	13.5	-13.0	1.59
LSBC F583-1	15 57 27.5	+20 39 58.0	2264	0.49	86.2	26.5	7.8	22.9	22.1	16.3	15.5	-13.5	1.51
UGC 10281	16 13 20.6	+17 11 34.3	1084	0.04	69.4	27.0	24.9	24.4	25.9	15.1	16.8	-15.6	0.61
LSBC F585-3	16 21 23.8	+20 51 56.0	3100	0.22	42.8	26.5	5.8	22.8	21.8	16.7	15.9	-13.3	1.60
UGC 10398	16 28 10.5	+17 38 23.9	4529	0.19	-16.2	26.5	9.0	23.0	21.9	16.0	15.0	-13.7	1.16

Column (4): radial velocity. Column (5): isophotal ellipticity. Column (6): isophotal position angle (from north through east). Column (7): B -band surface brightness of isophote used for Columns (5) and (6). Column (8): exponential disc scalelength. Column (9): disc central B -band surface brightness. Column (10): disc central R -band surface brightness. Column (11): total extrapolated B -band magnitude (see Section 2), corrected for foreground dust extinction. Column (12): total extrapolated R -band magnitude (see Section 2), corrected for foreground dust extinction. Column (13): log of the total $H\alpha$ flux (see Section 2). Column (14): log of the total $H\alpha$ emission line equivalent width (see Section 2).

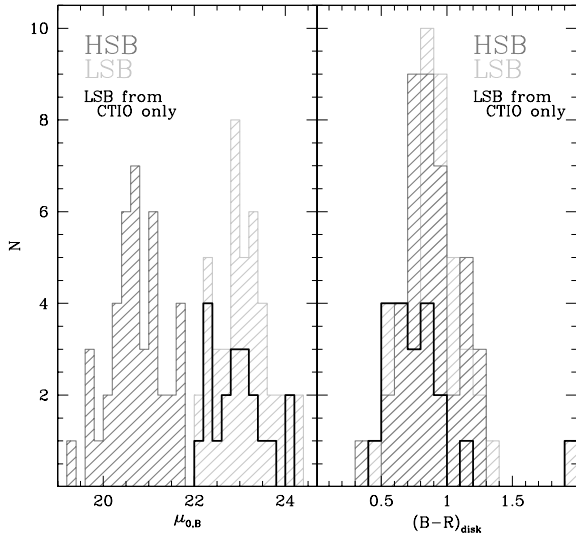


Figure 3. The B -band central surface brightness (left-hand side) and $B - R$ disc colours (right-hand side; see Section 2.1) distributions for the HSB (grey) and LSB (light grey) galaxies. The distributions for the LSB galaxies from the CTIO data only (black histograms) are also included.

the remaining galaxies, two broader $H\alpha$ filters provided at KPNO were used which are similar to those available at CTIO which were employed for 21 of the 69 galaxies observed there. Because the same or similar $H\alpha$ filters were used and since we were again using scaled versions of the R -band images to subtract the continuum emission from the $H\alpha$ images, it was possible to use the derived R -band calibrations with equation (1) to calibrate the $H\alpha$ images. We have done this assuming the mean value of $C = 0.73$ determined using the CTIO observations and have also included the corrections for internal dust extinction, $[N II]$ contamination, and stellar absorption derived from the Tremonti et al. (2004) data given in equations (2) and (3), again using the disc $B - R$ disc colours computed using the exponential fits to the surface brightness profiles (see Section 2.1). We note that the uncertainty in the value of C corresponds to an uncertainty of about 0.09 dex, which is comparable or less than the uncertainties in the corrections given by equations (2) and (3). The resulting calibrated continuum subtracted $H\alpha$ images were then used to measure the total $H\alpha$ flux and equivalent width for each galaxy, which is listed in Table 1. We have also used HIIPhot to identify and perform photometry on $H II$ regions for each of these images using the same settings as were used for the CTIO images. We note that for one of the LSB galaxies observed at KPNO, NGC 5409, no $H\alpha$ image was obtained. The $H II$ region luminosity functions (LFs) of the 23 galaxies with $H\alpha$ images were computed using the HIIPhot results, assuming $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$, and are displayed

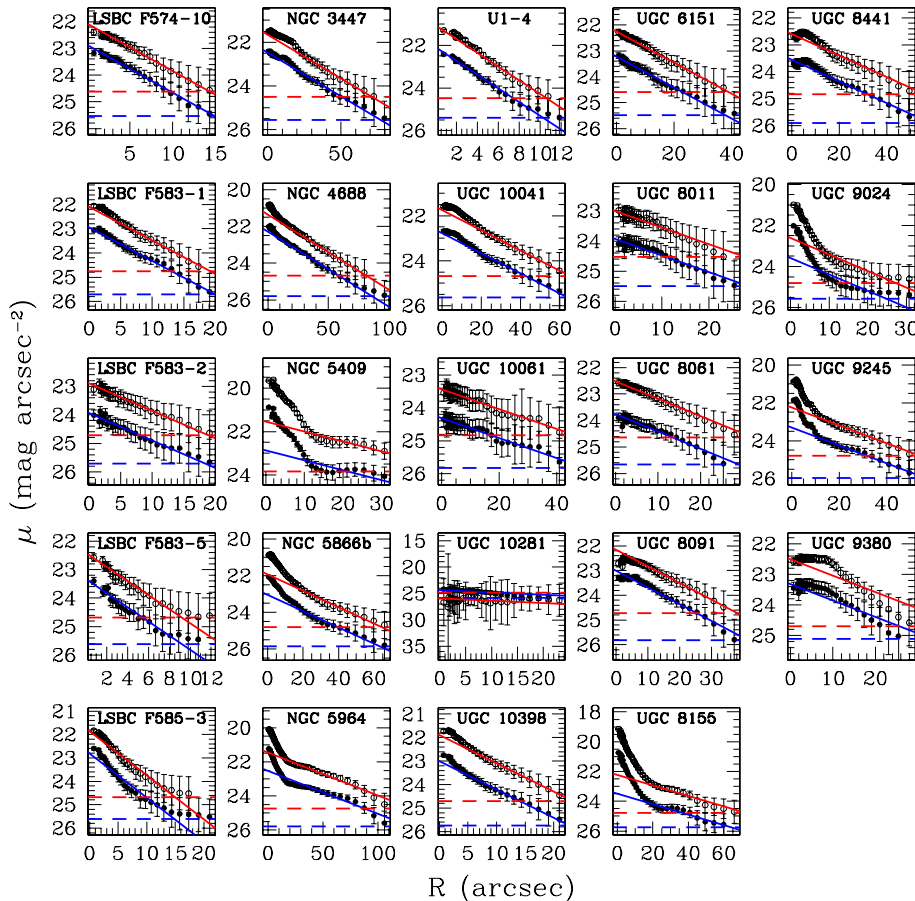


Figure 4. The R (open red points) and B (solid blue points) band surface brightness profiles for the LSB galaxies observed at KPNO; no offsets have been applied to the data. The best-fitting exponential profiles (see Section 2.1) are plotted in red and blue for the R and B band profiles, respectively. The 1σ limiting isophote is plotted as a horizontal dashed line for each band.

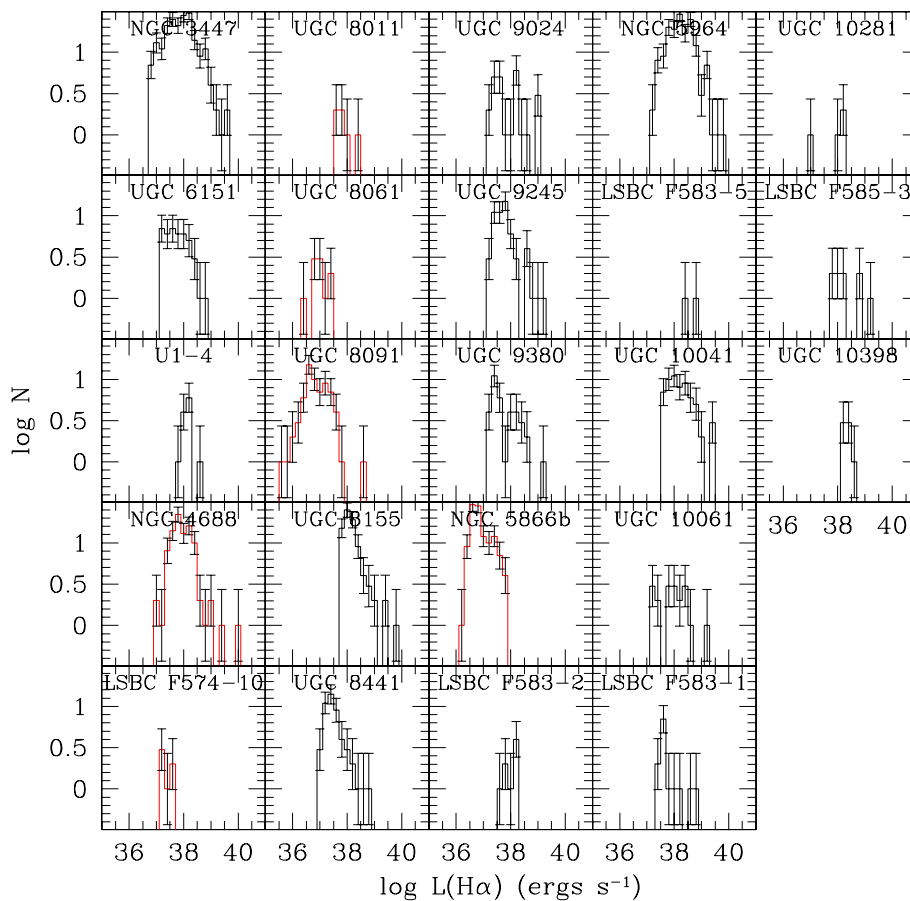


Figure 5. The H II region LFs for the LSB galaxies observed at KPNO. LFs for galaxies with radial velocities $< 1000 \text{ km s}^{-1}$ for which the distance estimates are somewhat unreliable are highlighted in red.

in Fig. 5. LFs for galaxies with radial velocities $< 1000 \text{ km s}^{-1}$ for which the distances may be relatively uncertain are highlighted in red. The H II region properties and LFs of the galaxies observed at CTIO are presented in Helmboldt et al. (2005).

3 H II REGIONS

3.1 Global properties

With B, R, and H α images for 45 HSB and 43 LSB galaxies, we are in a good position to explore the differences in H II region properties between the two classes of galaxies. In Fig. 6, we have plotted the total H α emission line equivalent width as a function of $B - R$ index for the HSB and LSB galaxies as well as similar data from the Nearby Field Galaxy Survey (NFGS; Jansen 2000), which includes a wide range of morphological types. From this plot, it can be seen that the amounts of star formation relative to the stellar continua within both types of galaxies are similar. Both are bluer on average than the NFGS galaxies. For the most part, both the LSB and HSB galaxies have similar H α equivalent widths at the same colour indices as the NFGS galaxies.

However, we start to see distinct differences between HSB and LSB galaxies when we look at H II region properties. In Fig. 7, we have plotted the distributions for two properties (1) the H II region covering factor, the ratio of the total area occupied by H II regions taken from the HIIphot results to the total disc area taken to be $9\pi h^2$, where h is the disc scalelength, and (2) the mean

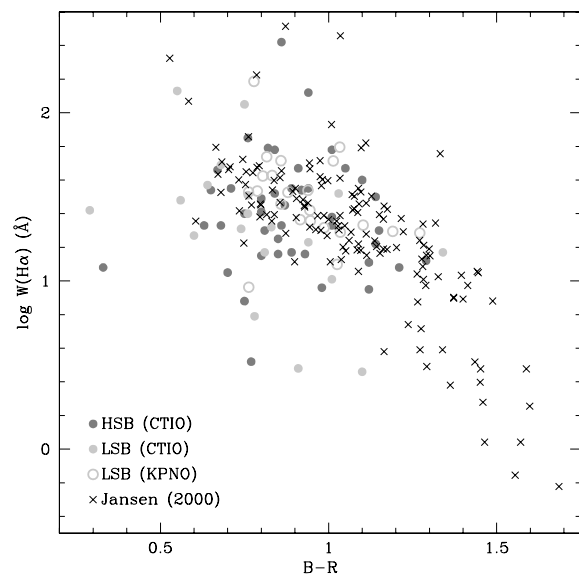


Figure 6. The total H α emission line equivalent width versus the $B - R$ colour index for HSB (grey) and LSB (light grey) galaxies; LSB galaxies observed at KPNO are plotted as open points. For comparison, the galaxies from the NFGS of Jansen (2000) are plotted with \times symbols.

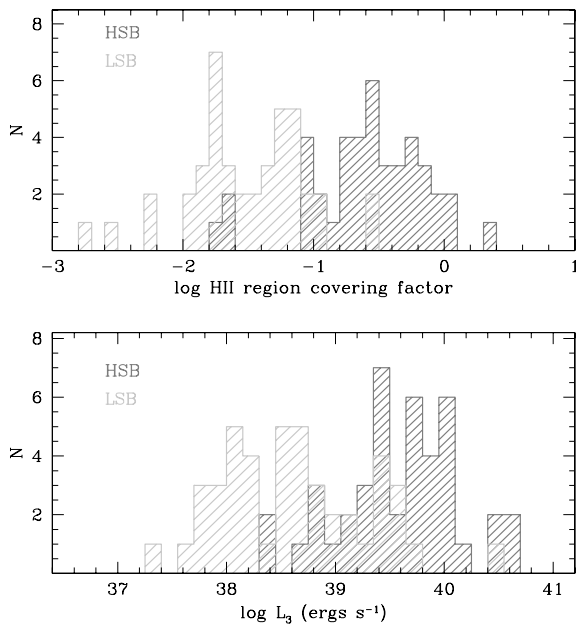


Figure 7. Distributions for the H II region covering factor (see Section 3.1; upper) and the mean H α luminosity of the three brightest H II regions (lower) for the HSB (grey) and LSB (light grey) galaxies.

luminosity of the three brightest H II regions, L_3 . We note that for the surface brightness profiles presented in Helmboldt et al. (2004) and in Fig. 4, the radii are the geometric mean radii of the elliptical apertures used, that is $R = \sqrt{ab}$. Therefore, the assumed total disc area of $9\pi h^2$ is the area of an ellipse with a semi-major axis of $3h(b/a)^{-1/2}$, where b/a is the axis ratio of the elliptical apertures used. No attempt has been made to de-project the estimated total disc area or the area occupied by H II regions. However, we have found no significant correlation between b/a and either the H II region covering factor or the average number of regions per square kiloparsec with Spearman rank correlation coefficients of -0.08 and -0.18 , respectively. Additionally, the mean axis ratios for the LSB and HSB samples are similar ($\langle b/a \rangle = 0.65$ and 0.60 , respectively). It is therefore unlikely that this has had a significant effect on the results presented here.

As Fig. 7 demonstrates, the H II region covering factor is significantly lower on average for LSB galaxies than it is for HSB galaxies, and the same is true for L_3 . These results are generally consistent with what has been found previously (Greenawalt 1998; Helmboldt et al. 2005; Oey et al. 2007; O’Neil et al. 2007).

To compare the typical distributions of H II region luminosities, we have constructed co-added H II region LFs for both the LSB and HSB galaxies. First, we must note that we only have radial velocity distances for our galaxies which implies that those with smaller radial velocities will have larger distance uncertainties. These large errors may artificially alter the shape of any co-added LF, and we therefore must exclude such galaxies from this process. Helmboldt et al. (2005) determined that the optimum radial velocity cut-off for a co-added LF with bins of 0.2 dex in width is $V_r < 840 \text{ km s}^{-1}$. For all co-added LFs discussed in this paper (see Section 3.2 as well), we have excluded galaxies with $V_r < 840 \text{ km s}^{-1}$ and have used the same bin width as Helmboldt et al. (2005). There are 40 of our HSB galaxies and 37 of our LSB galaxies that have $V_r > 840 \text{ km s}^{-1}$.

After applying the radial velocity cut, the next step to co-adding the galaxies’ LFs was to determine a limiting H II region luminosity

for each galaxy by fitting a power law to the H II region luminosity as a function of the signal-to-noise ratio (S/N) of the integrated flux of the region as determined by HIIphot. The limiting luminosity was then taken to be the luminosity given by the power-law fit for $S/N = 5$, plus the rms deviation of the data about the power-law fit. We then constructed an H II region LF for each galaxy above this limiting luminosity. Each LF was then divided by the area of the disc, $9\pi h^2$, to make sure that larger galaxies with more H II regions would not dominate the shape of the co-added LF. As noted above, this estimate of the total disc area has not been de-projected to the face-on disc area. However, as also noted, there is no significant correlation between the isophotal axis ratio, b/a , and the average number of H II regions per square kiloparsec, and the mean values of b/a for the LSB and HSB samples are similar.

For each galaxy group, the average number of regions per square kiloparsec was computed for each luminosity bin by averaging over all galaxies with limiting H II region luminosities larger than the upper boundary of the bin. Uncertainties for each LF were computed using jackknife resampling. Specifically, the galaxies were put in groups of four and the LF was computed 11 different times, each time excluding one group of four, and the standard deviation for each luminosity bin over these 11 computations was taken to be the uncertainty. For instances where the shot noise ($\sqrt{N_{\text{tot}}}$, where N_{tot} is the total number of regions among all co-added galaxies) was larger than the jackknife value, the shot noise was adopted as the uncertainty.

The HSB and LSB galaxy co-added LFs are plotted with these uncertainties in Fig. 8. We can see that while the LSB LF is well approximated by a single power law over more than two decades in H α luminosity, the HSB LF shows a significant amount of curvature. This is similar to what was found by Helmboldt et al. (2005). To quantify this, we have fit to the LFs the following function

$$\frac{dN}{dL} = \phi_* \left(\frac{L}{L_*} \right)^\alpha \exp \left(\frac{-L}{L_*} \right) \quad (4)$$

commonly referred to as a Schechter function where α is the faint-end slope and L_* characterizes the amount of curvature at the bright-end. The Schechter function was chosen because it is basically a

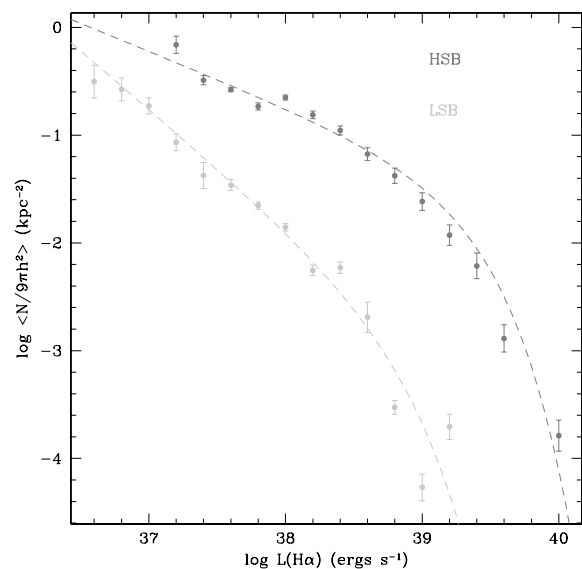


Figure 8. Co-added H II region LFs (see Section 3.1) for the HSB (grey) and LSB (light grey) galaxies with radial velocities $> 840 \text{ km s}^{-1}$ (see Section 3.1). Schechter functions fit to the LFs are also plotted as dashed lines.

truncated power law which is qualitatively what one would expect for an aged population of star clusters with a power-law distribution of initial masses (see the Monte Carlo simulations described in Section 4). The fitted functions are also plotted in Fig. 8. The values of α and L_* from these fits are -2.08 and $5.79 \times 10^{38} \text{ erg s}^{-1}$ for the LSB LF and -1.52 and $1.86 \times 10^{39} \text{ erg s}^{-1}$ for the HSB LF. We note that this difference in LF shape is unlikely to be artificially caused by the application of the corrections given in equation (2) since these corrections are solely functions of $(B - R)_{\text{disc}}$, and Fig. 3 demonstrates that the $(B - R)_{\text{disc}}$ distributions for the LSB and HSB galaxies are essentially the same. We must also acknowledge that the shape of the LF, particularly for HSB galaxies for which the H II region covering factor is larger, may be artificially flattened by blending of H II regions (see Scoville et al. 2001). We will address this issue in more detail in Section 3.3.

3.2 Local properties

While the results presented in Figs 7 and 8 mostly reinforce earlier results, a question still remains regarding whether they arise from differences in the global properties of HSB and LSB galaxies or from local conditions within the LSB and HSB regions of all discs. While results previously presented by Greenawalt (1998) suggest that these properties vary with local surface brightness within HSB discs, we now have the ability to explore whether or not that variation is seen in LSB discs and if it is indeed similar to what has been seen for HSB discs. We have done this by measuring the local B - and R -band surface brightness in the area near each H II region using a circular aperture with a radius of $2\sqrt{n_p}/\pi$ where n_p is the number of pixels occupied by the H II region on the image. We have defined the sizes of the apertures in this way rather than using an aperture with a fixed physical size to ensure that each aperture was small enough to represent the local conditions around each region while being large enough that the region itself did not dominate the local flux. In Fig. 9, we have plotted the H II region luminosity versus local surface brightness. For both LSB and HSB galaxies, there appears to be a trend for H II regions to be more luminous in higher surface brightness areas. The fact that there are few low luminosity H II regions in higher surface brightness areas may be a selection effect as they are harder to detect in such areas. However, the paucity of highly luminous H II regions in more LSB areas cannot be the result of such a bias, implying that while this possible selection effect may influence the observed trend, it cannot be the cause of it. In Fig. 9, for $\mu_B < 23$, the trend for LSB galaxies becomes somewhat steeper than that for HSB galaxies. However, this is heavily influenced by a few LSB galaxies which contain some HSB regions and this difference is not seen in the R -band data. In both the B and R data, the trend appears to become steeper for $\mu_B < 22.5$ and $\mu_R < 21.5$.

To further explore and constrain the results shown in Fig. 9, we have fit a line to $\log L(H\alpha)$ as a function of local surface brightness for each galaxy to obtain a gradient. The distributions for these gradients using both the B and R data are plotted in Fig. 10 separately for LSB and HSB galaxies. On average, the $\log L(H\alpha)$ gradients for LSB galaxies appear to be about $0.1 \text{ dex mag}^{-1} \text{ arcsec}^2$ less steep than those for HSB galaxies with both distributions peaking at values < 0 . This is likely a reflection of the fact that the LSB galaxies have a substantially higher fraction of regions with $\mu_B > 22.5$ and $\mu_R > 21.5$, where the trends seen in Fig. 9 are less steep than HSB galaxies. We note that the gradient distributions for both the B and R bands peak around -0.15 to $-0.05 \text{ dex mag}^{-1} \text{ arcsec}^2$,

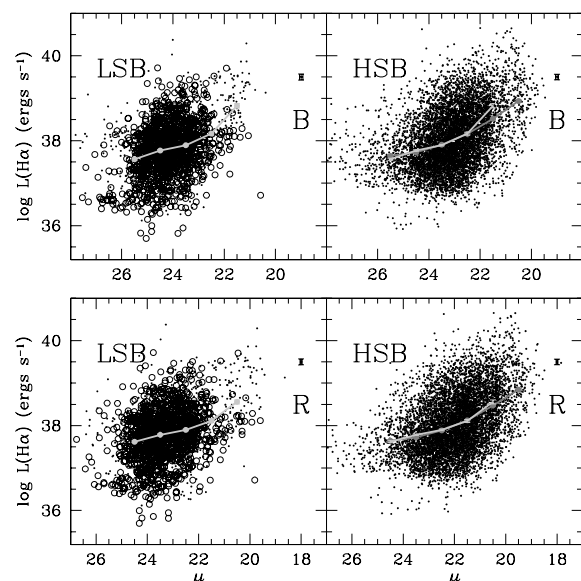


Figure 9. The $H\alpha$ luminosity of H II regions from LSB (left-hand side) and HSB (right-hand side) galaxies as a function of local surface brightness in the B (upper) and R (lower) bands. In the left-hand panels, data for H II regions within the LSB galaxies observed at KPNO are plotted as open circles. Mean values are plotted within surface brightness bins for both LSB (light grey) and HSB (grey) galaxies.

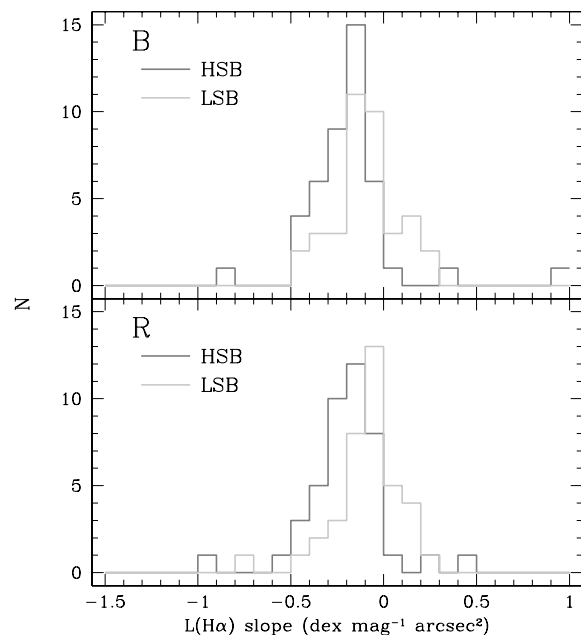


Figure 10. The distributions of the gradients of the log of the H II region $H\alpha$ luminosity (see Section 3.2) as a function of local surface brightness in the B (upper) and R (lower) bands for H II regions within LSB (light grey) and HSB (grey) galaxies.

which is similar to the slopes for the median values plotted in Fig. 9.

The results illustrated by Figs 9 and 10 imply that for a given local surface brightness, the typical H II region luminosity for LSB galaxies is likely similar to that for HSB galaxies. This implies that local surface brightness/density strongly affects star cluster formation within galaxy discs, as suggested by the toy model of O’Neil,

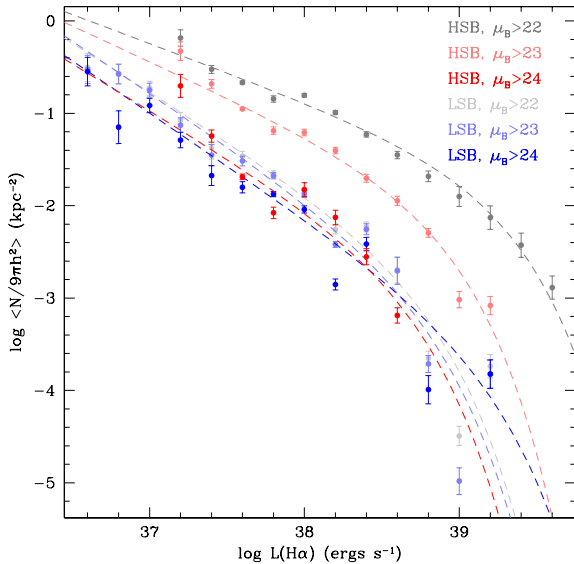


Figure 11. Co-added LFs for H II regions from HSB galaxies (grey to red) and from LSB galaxies (light grey to blue) using different limiting local B -band surface brightness values (see panel).

Bothun & Schombert (1998). However, this does not mean that global disc properties do not have a significant effect on the detailed properties and formation histories of star clusters. For instance, at the same local surface brightness, the shapes of the distributions of H II region luminosities are not necessarily the same for LSB and HSB galaxies, and any difference in shape may point to the influence of global properties on cluster formation. To explore this further, we have constructed co-added LFs for H II regions based on local B -band surface brightness separately for HSB and LSB galaxies. We did this by making three co-added LFs separately for LSB and HSB galaxies using three different limiting local B -band surface brightness values, $\mu_{B,\text{lim}} = 22, 23$ and 24 mag arcsec $^{-2}$. In other words, each co-added LF was remade using only H II regions with local $\mu_B > \mu_{B,\text{lim}}$. The six resulting LFs are plotted in Fig. 11 along with Schechter function fits. From these plots, it can be seen that the HSB LFs have a noticeably larger amount of curvature and a faint-end with a shallower slope than the LSB LFs for $\mu_{B,\text{lim}} < 24$. This is better illustrated in Fig. 12 where we have plotted the parameters of the Schechter functions fitted to the LFs in both Fig. 8 and 11 as functions of $\mu_{B,\text{lim}}$. While the faint-end slope, α , remains roughly constant with $\mu_{B,\text{lim}}$ for LSB galaxies, it appears to become steeper with decreasing limiting surface brightness for HSB galaxies. At $\mu_{B,\text{lim}} = 24$, α appears to be essentially the same for both LSB and HSB galaxies. A similar pattern can be seen for the L_* parameter, except that the values for LSB and HSB galaxies appear to be in good agreement for $\mu_{B,\text{lim}}$ as bright as about 23 mag arcsec $^{-2}$.

3.3 Blending

Since we have used ground-based images of galaxies with radial velocities >840 km s $^{-1}$ to examine H II region properties in disc galaxies, we cannot ignore the influence blending has on the results displayed in Figs 11 and 12. Blending may be particularly important when considering the difference between the shapes of the H II region LFs in HSB and LSB areas within discs. This is evidenced by the fact that in Fig. 11, the HSB galaxy LFs have a significantly

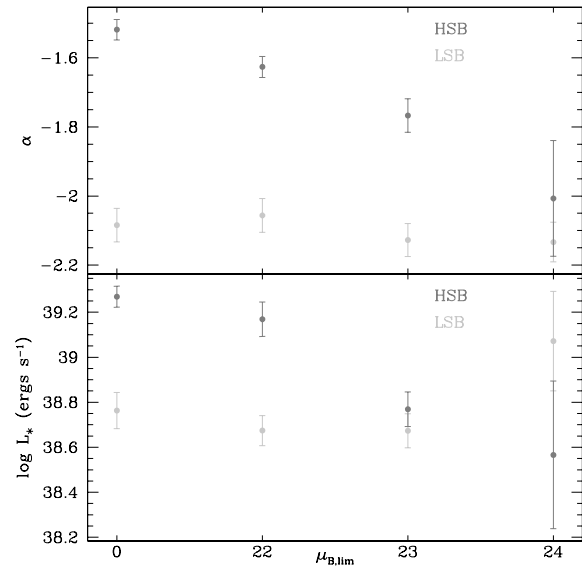


Figure 12. For the LFs displayed in Fig. 8 and 11, the faint-end slope, α (upper), and the value of L_* (lower) for the Schechter functions fit to the LFs as a function of limiting local B -band surface brightness. The parameters from the fits to the LFs using all H II regions are plotted for $\mu_{B,\text{lim}} = 0$. As in previous figures, the results for HSB galaxies are plotted in dark grey and the results for LSB galaxies in light grey.

larger number of regions per square kiloparsec than the LSB galaxy LFs for $\mu_{\text{lim},B} < 23$, implying that blending may be much more important for these LFs. Since the shapes of the LFs for HSB and LSB galaxies are roughly similar for $\mu_{B,\text{lim}} \gtrsim 23$, it is possible that this difference in shape is purely the result of the LFs of regions from HSB areas within discs being artificially biasing towards higher luminosities by blending.

The effect of blending can be demonstrated by smoothing images of more nearby HSB galaxies and re-analysing them with HIIphot. For this, we have selected two HSB galaxies that were imaged during the 2000 October CTIO run, M83 and NGC 1313. These two galaxies were chosen because they are relatively nearby ($V_r = 516$ and 457 km s $^{-1}$, respectively) and have ~ 1000 H II regions identified by HIIphot so that LFs can be constructed for each galaxy for different values of $\mu_{\text{lim},B}$ with significant numbers of regions. In addition to this, M83 and NGC 1313 have two distinct morphologies; M83 is a ‘grand design’ spiral galaxy (type Sb) and NGC 1313 is an SBd galaxy with a somewhat disturbed morphology, especially in H α emission. To simulate the effect of blending at different radial velocities, we smoothed the B , R and H α images of each galaxy with the appropriate Gaussian kernel to simulate the size of the seeing disc measured from stars on the narrow-band image at 10 radial velocities, 560, 600, 700, 800, 900, 1100, 1300, 1500, 2000 and 2500 km s $^{-1}$. The images were then regridded so that the angular size of the pixels relative to that of the seeing disc remained constant. Following this, Gaussian noise was added to each image so that the noise in the smoothed image was similar to that of the original image. HIIphot was then used to identify the H II regions and measure their properties on each of these images.

Following the HIIphot analysis, we constructed H II region LFs for $\mu_{B,\text{lim}} = 0, 22, 23$ and 24 for each galaxy at each radial velocity. The resulting LFs for $V_r = 560, 900, 1500$ and 2500 km s $^{-1}$ are plotted in Fig. 13 with Schechter function fits. In general, as V_r increases, the number of detected regions decreases and the LFs become more

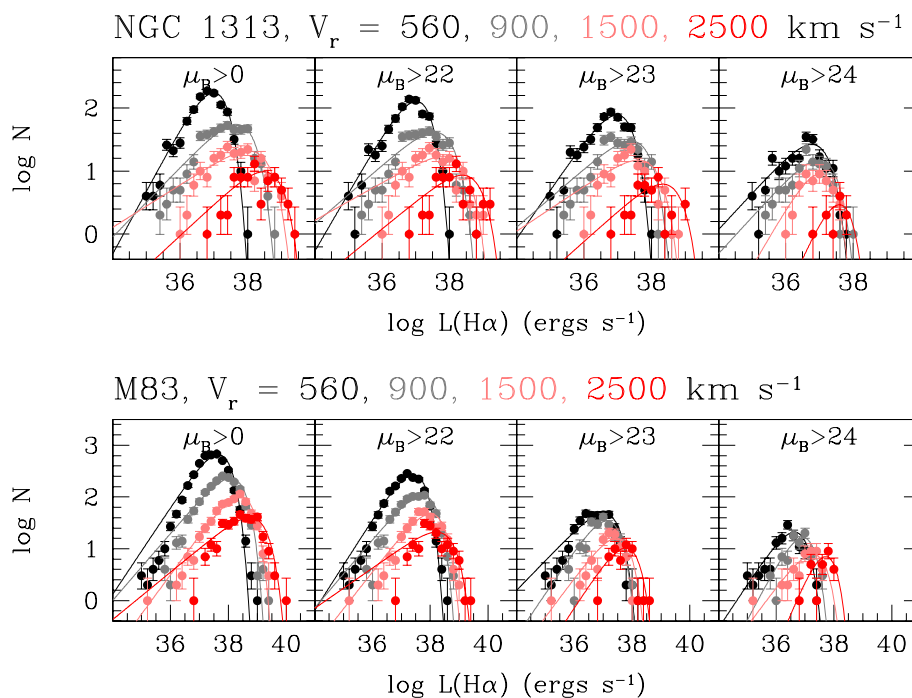


Figure 13. For NGC 1313 (upper panels) and M83 (lower panels), the H II region LFs for four different values of $\mu_{B,\text{lim}}$ (see panels) at four different simulated radial velocities (see Section 3.3). The solid lines are Schechter functions fit to the data.

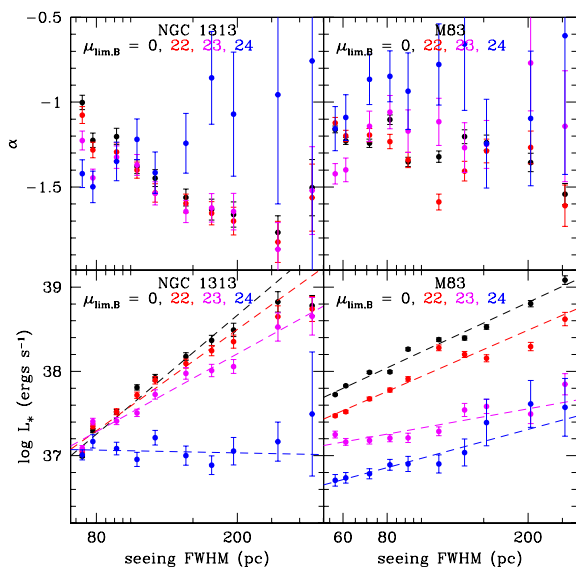


Figure 14. For NGC 1313 and M83, the parameters α and L_* derived from the Schechter functions fit to the LFs measured at 10 simulated radial velocities (560, 600, 700, 800, 900, 1100, 1300, 1500, 2000 and 2500 km s^{-1} ; see Section 3.3) for four different values of $\mu_{B,\text{lim}}$ versus the physical size of the seeing disc. The dashed lines in the bottom panels are weighted, linear least-squares fits to the data plotted to illustrate the convergence (or lack thereof) of the trends between L_* and the spatial resolution among the different values of $\mu_{B,\text{lim}}$.

biased towards higher luminosities. To better demonstrate this, we have plotted the values of α and L_* from the Schechter function fits as a function of the physical size of the seeing disc in Fig. 14. For both galaxies, α is not significantly different for different values of $\mu_{B,\text{lim}}$ for $\mu_{B,\text{lim}} \leq 23$. For $\mu_{B,\text{lim}} = 24$, α is slightly higher. For M83,

α does not appear to change significantly with radial velocity while for NGC 1313, α decreases with increasing V_r for $\mu_{B,\text{lim}} \leq 23$. We therefore find no evidence from these two examples that blending has caused α to be larger in higher surface brightness areas of HSB discs as seen in Fig. 12.

For both NGC 1313 and M83, L_* increases with V_r for $\mu_{B,\text{lim}} \leq 23$ and for all values of $\mu_{B,\text{lim}}$ for M83; L_* remains roughly constant for NGC 1313 for $\mu_{B,\text{lim}} = 24$. For NGC 1313, the value of L_* appears to differ somewhat among the different values of $\mu_{B,\text{lim}}$ with the value for $\mu_{B,\text{lim}} = 24$ being substantially lower at nearly all radial velocities. However, the values of L_* appear to converge at the lowest radial velocities. In contrast, the differences among the values of $\log L_*$ for different limiting surface brightnesses for M83 are roughly constant with V_r . Linear fits to these data (see Fig. 14) suggest that the values of L_* may converge for relatively small (between ~ 2 and 10 pc) spatial resolutions. However, this requires one to extrapolated the trends in the lower right-hand panel of Fig. 14 roughly one order of magnitude lower than the smallest spatial resolution at $V_r = 560 \text{ km s}^{-1}$, implying that there is no evidence that blending alone may be able to reproduce the trend seen in the lower panel of Fig. 12 for M83. From these two examples, it is clear that the values of L_* for the higher surface brightness areas of HSB galaxies are likely inflated by the effects of blending. However, it is also clear that while for some galaxies that are similar to NGC 1313, the trend between L_* and $\mu_{B,\text{lim}}$ may be entirely caused by blending, for other, more typical spiral galaxies like M83, the trend is somewhat unaffected as blending appears to increase L_* by roughly the same factor for all values of $\mu_{B,\text{lim}}$. It is therefore likely that the trend seen in the lower panel of Fig. 12 has a real physical origin and is not entirely the result of blending.

We note that the Schechter fits were not restricted to similar luminosity ranges as those fit to the LFs in Fig. 11 because the radial velocity distances for these two galaxies are highly uncertain. This has possibly led to the inclusion of more faint H II regions and

sharper turnovers at the faint-end. We therefore note that the results presented in Fig. 14 most likely do not reflect the magnitude of the effect of blending on the trend shown in Fig. 12. However, the results are still useful for testing whether or not blending is chiefly responsible for the trends seen for HSB galaxies in Fig. 12. From the results shown in Fig. 14, it appears that blending is likely not the sole cause of these trends.

4 INTERPRETATION AND DISCUSSION

As the results presented in Figs 11 and 12 have shown, the shape of the H II region LF appears to be substantially affected by global processes. In particular, both the faint-end slope and L_* seem to be roughly independent of local surface brightness for LSB galaxies, while both parameters seem to decrease as the limiting local B -band surface brightness is made fainter for HSB galaxies. In Section 3.3, we demonstrated with two examples that the effects of blending have probably inflated the values of L_* , possibly to a larger degree for higher surface brightness areas. However, we also demonstrated that it is likely that blending is not the sole cause of the results shown in the lower panel of Fig. 12. To analyse the physical origins of the difference in LF shapes between LSB and HSB galaxies, we have constructed Monte Carlo simulations of H II region LFs patterned after those developed by Oey & Clarke (1998) which are similar to those used by von Hippel & Bothun (1990). We note that Oey & Clarke (1998) explored different assumptions about the time dependence of the H α luminosity due to stars of different masses and about the formation history of star clusters. However, they found that assuming a constant H α luminosity over the main sequence lifetime of each star which abruptly goes to zero and an instantaneous burst where all clusters form at once was able to reproduce the shapes of typical H II region LFs. We have therefore adopted these two assumptions as well as the same power laws for the distribution of the number of stars per cluster (slope of $\beta = -2$), the relationships between both H α luminosity and main sequence lifetime and stellar mass (slopes of $\delta = 1.5$ and $d = -0.7$, respectively) and the initial mass function (IMF) of Salpeter (1955). We have also assumed the same upper and lower mass limits for the IMF of 17 and $100 M_{\odot}$ and an H α luminosity and main sequence lifetime for $100 M_{\odot}$ stars of 3×10^{38} erg s $^{-1}$ and 2.8 Myr. We note that these power-law relations break down below the adopted lower mass limit for the IMF, which artificially introduces a turnover at $\log L(\text{H}\alpha) \sim 37.7$ and we do not consider the Monte Carlo LFs below this luminosity. Note that since this is a non-physical limit, it was not applied to the fitting of Schechter functions to the observed LFs.

Beyond the assumptions we have made, the two physical factors that determine the shapes of the LFs are the age (from birth) of the clusters and the maximum number of stars per cluster, $N_{*,\text{max}}$. To explore how these two parameters affect the faint-end slope and amount of curvature at the bright-end, we have constructed 121 LFs with 1000 clusters apiece for 11 ages ranging from 0 to 6 Myr and 11 values of $N_{*,\text{max}}$ ranging from 10 to 1000. Each simulated LF used the same set of random numbers to assign a number of stars to each of the 1000 clusters from the assumed power-law distribution, truncated at $N_* = N_{*,\text{max}}$. We note that while ages up to 6 Myr may seem too large for H II regions, the instantaneous burst model of Bruzual & Charlot (2003) estimates that at an age of 6 Myr, the flux of ionizing photons output by a cluster of stars would have dropped by a factor of roughly 30, but would be far from negligible.

An additional factor that must also be considered is the uncertainty in the H α luminosity which varies from one galaxy to another

and can be rather large because of the statistical corrections used for internal dust extinction, [N II] contamination and stellar absorption [see equation (2)]. Also, for galaxies for which the R -band calibration and equation (1) were used to calibrate the H α flux, the uncertainty in the H α luminosity is even larger given the uncertainty in the C parameter in equation (1). This technique was used for a larger fraction of the LSB galaxies and may therefore contribute to any differences between the shapes of the LSB and HSB LFs. To incorporate this into our Monte Carlo-simulated LFs, we first computed an estimate of the full uncertainty in the H α calibration for each galaxy. This computation included the uncertainties in the parameters given in equation (2). It also included the combination of the uncertainty in the C parameter in equation (1) and the uncertainty in the R -band calibration or, in the case of the majority of the CTIO observations, the uncertainty in the calibration solution for the H α filters derived from observations of spectrophotometric standard stars. Finally, a radial velocity uncertainty of 94 km s^{-1} was assumed for each galaxy which is based on the uncertainty in the correction for in-fall into the Virgo cluster (Helmboldt et al. 2005). The resulting uncertainties in $\log L(\text{H}\alpha)$ range from about 0.1–0.2 dex. For each group of galaxies, HSB and LSB, the fraction of the total number of H II regions per square kiloparsec among all galaxies in the group was computed for each galaxy. This was then used to assign the uncertainty of that galaxy to the same fraction of H II regions in the Monte Carlo simulations. This uncertainty was applied to the log of the H α luminosity of each simulated region by adding to it a factor of σg , where g is a random number drawn from a unit normal distribution and σ is the assigned uncertainty. Since all of the H II regions within a particular galaxy have the same H α calibration uncertainty, the same value of g was used for all simulated regions with the same assigned value of σ . Each of the 121 simulated LFs was reconstructed twice, once using the uncertainties for the LSB group, and once using those of the HSB group, resulting in 363 simulated LFs in all.

Schechter functions were fit to all of the simulated LFs, including those with and without the modelled effects of calibration uncertainty. A sampling of the simulated LFs for combinations of six values of star cluster age and six values of $N_{*,\text{max}}$ are displayed in Fig. 15 with their Schechter fits. Those without the modelled effects of H α calibration uncertainty are plotted in black; those with modelled uncertainties based on the HSB and LSB groups are plotted in grey and light grey, respectively. These plots show that for a fixed value of either $N_{*,\text{max}}$ or age, the amount of curvature at the bright-end decreases as the other parameter increases. However, it does appear that $N_{*,\text{max}}$ has a much stronger effect on the amount of curvature than the age. They also show that the introduction of uncertainties somewhat reduces the magnitude of this curvature by generally flattening the LFs. The magnitude of this effect appears to be relatively insignificant for LFs with uncertainties based on the data for either HSB or LSB galaxies.

While other factors such as the assumed IMF, the star formation history and blending also determine the shape of the H II region LF, our Monte Carlo simulations provide some insight into the origin(s) of the difference in shapes among the LFs plotted in Fig. 11 in terms of age and cluster size. From the simulated LFs shown in Fig. 15, it appears that the power-law shape of the LFs seen throughout LSB galaxy discs and in the most LSB regions of HSB discs are plausibly the result of clusters forming with relatively large values of $N_{*,\text{max}}$ and that the curvature at the bright-end seen for the HSB regions of HSB discs requires smaller values of $N_{*,\text{max}}$. However, the paucity of highly luminous H II regions within the LSB regions of discs (see Fig. 9) suggests that it is highly unlikely that $N_{*,\text{max}}$ is smaller for

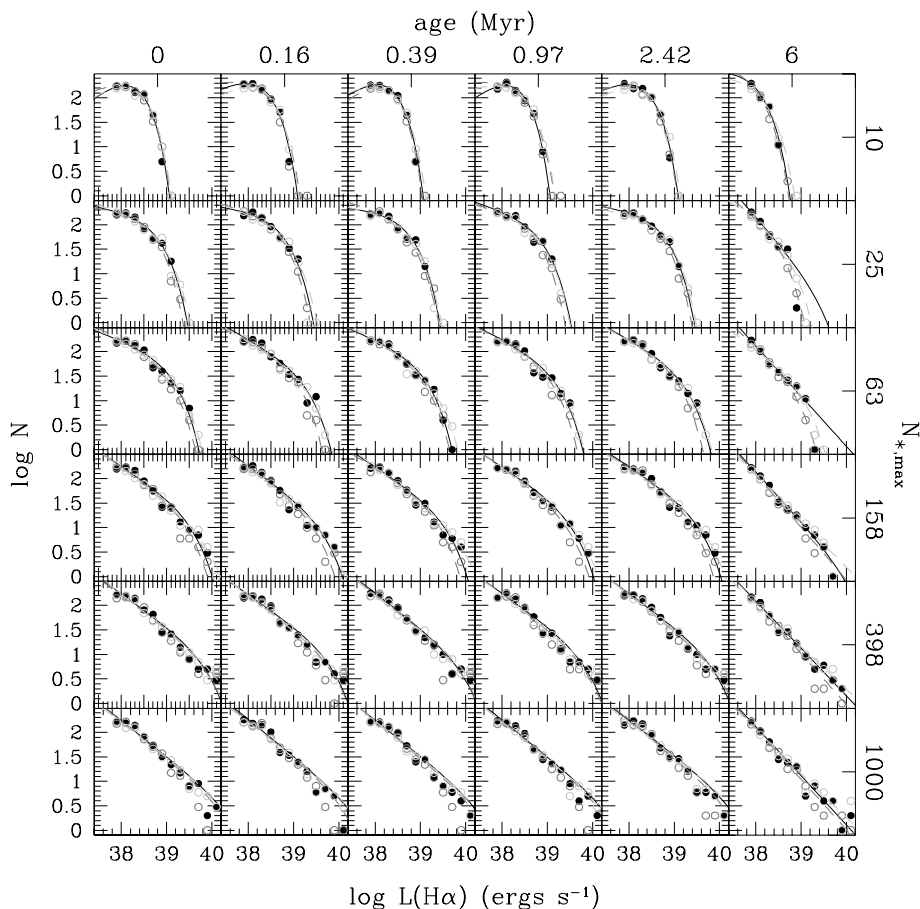


Figure 15. A sampling of the Monte Carlo simulated H II region LFs for six star cluster ages (marked above the panels) and six values of $N_{*,\max}$, the maximum number of stars per cluster (marked to the right of the panels). The simulations used 1000 stellar clusters apiece and were patterned after those of Oey & Clarke (1998, see Section 4.2). The simulated LFs that do not include modelled $L(\text{H}\alpha)$ uncertainties are plotted in black. Those that included modelled uncertainties based on the estimated $\text{H}\alpha$ calibration uncertainties (see Section 4) for LSB and HSB galaxies are plotted in light grey and grey, respectively. The plots shown here also include Schechter function fits [see equation (2)].

clusters within HSB regions than for those within LSB regions. From Fig. 15, it can be seen that if the clusters within LSB regions are relatively old (~ 6 Myr) and the clusters within HSB regions are significantly younger ($\lesssim 1$ Myr), the value of $N_{*,\max}$ can be roughly the same. Since the effect of blending has not been included in the simulated LFs but likely affects the observed LFs, we cannot provide accurate estimates of the typical ages of the clusters within HSB and LSB regions using these simulations. However, it appears that qualitatively, the simulated LFs suggest that the typical star cluster within LSB regions is significantly older than the typical HSB region cluster. We note that a similar explanation has been proposed for the difference in LF shapes between arm and interarm H II regions within spiral galaxies with arm regions being younger on average (von Hippel & Bothun 1990; Oey & Clarke 1998).

As the LFs in Fig. 11 demonstrate, the higher surface brightness regions of HSB discs tend to have more H II regions per square kiloparsec than LSB galaxies have throughout their discs. The Monte Carlo-simulated LFs seem to imply that these more densely packed areas tend to have younger clusters on average. Since the HSB discs have spiral arms more frequently than the LSB discs in our samples, this could point to a fundamental difference between spiral-arm driven star formation and the more sporadic star formation that is likely prominent in LSB discs, particularly for Sm and Im galaxies (see, e.g. Bothun & Caldwell 1984; de Blok, van der Hulst & Bothun

1995; Helmboldt et al. 2005). The results displayed in Figs 11 and 12 would further imply that there may be a transition from spiral-arm driven star formation to a more erratic process within HSB discs that typically occurs at $\mu_B \sim 24$. The galaxy M101 provides some verification of this notion as the luminous H II regions in the outer area of its disc suggest a more erratic process of star formation than is seen within the disc where spiral arms are prevalent (see, e.g. Scowen, Dufour & Hester 1992).

It is also possible that the change in the shape of the H II region LF that appears to occur within HSB galaxies at a local surface brightness of $\mu_B \sim 24$ is linked to a change in the cold gas density. For example, the surface brightness profile we have measured for M83 reaches $\mu_B = 24$ at an apparent radius of about 4.4 arcmin. From the radial profiles presented by Thilker et al. (2005), the gas surface density at this radius is about $25 M_{\odot} \text{pc}^{-2}$, which is roughly six times larger than the typical critical gas density for star formation determined for similar spiral galaxies (see, e.g. Kennicutt 1989). However, at this radius, the profile for the molecular hydrogen drops quickly, and the slope of $\log \Sigma_{\text{H}_2}(R)$ abruptly becomes steeper by a factor of ~ 3 . At this radius, the gas-phase oxygen abundance of M83 has also dropped by factor of about 2 compared to what it is in the inner parts of the disc (Diaz & Tosi 1984). This may be the underlying cause of the lack of both CO gas and the dust grains needed to facilitate the conversion of atomic to molecular

hydrogen. This example demonstrates that it is plausible that the apparent transition to a more sporadic mode of star formation is linked to a change in physical conditions that has caused a sharp decrease in the amount of molecular gas that is readily available in the outer regions of spiral galaxy discs.

Overall, the results we have presented demonstrate that while local conditions appear to have a large effect on star formation within all galaxy discs, it appears that the change in H II region properties from higher to lower density environments is significantly more pronounced within HSB discs than it is within LSB discs. The most plausible explanation appears to be that the star formation within the outer parts of HSB discs and throughout LSB discs is more sporadic than that which occurs within the inner regions of HSB discs. Since LSB discs tend to have spiral arms less frequently, little if any molecular gas, and relatively low metallicities, it may be that this sporadic mode of star formation tends to dominate within regions that have relatively low amounts of dust and molecular gas that are also relatively stable against global instabilities like spiral density waves. This is also consistent with the physical conditions of the outer regions of spiral galaxy discs (see, e.g. Diaz & Tosi 1984; Thilker et al. 2005), which points to a common physical origin for the similarly shaped H II region LFs of both LSB discs and the LSB regions of HSB discs.

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