

Mechanical equilibrium of hot, large-scale magnetic loops on T Tauri stars

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ABSTRACT

The most extended, closed magnetic loops inferred on T Tauri stars confine hot, X-ray-emitting plasma at distances from the stellar surface beyond the X-ray-bright corona and closed large-scale field, distances comparable to the corotation radius. Mechanical equilibrium models have shown that dense condensations, or ‘slingshot prominences’, can rise to great heights due to their density and temperatures cooler than their environs. On T Tauri stars, however, we detect plasma at temperatures hotter than the ambient coronal temperature. By previous model results, these loops should not reach the inferred heights of tens of stellar radii where they likely no longer have the support of the external field against magnetic tension. In this work, we consider the effects of a stellar wind and show that indeed hot loops that are negatively buoyant can attain a mechanical equilibrium at heights above the typical extent of the closed corona and the corotation radius.

Key words: stars: activity – stars: coronae – stars: flare – stars: magnetic field – stars: pre-main-sequence.

1 INTRODUCTION

Observations of T Tauri stars (TTs) have revealed the presence of magnetic structures much larger than those seen on our Sun. Indeed, analysis of these structures has invoked solar-flare-based models: to infer the length-scales of confining loops, Reale et al. (1997) developed the uniform cooling loop (UCL) model, linking solar-flare X-ray decay slopes to spatially resolved lengths of the loops confining the emitting gas. The UCL model was applied to data from the Chandra Orion Ultradeep Project (COUP, Getman et al. 2005), and it was discovered that evidently the post-flare loops of TTs can reach up to multiple stellar radii in extent (Favata et al. 2005, results summarized in Fig. 1). This result is surprising: T Tauri coronae are typically believed to be compact, as evidenced by the common detection of rotationally modulated coronal X-ray emission (Flaccomio et al. 2005). X-ray-bright regions are generally thought to be confined within the corona, with loops generally not reaching more than a stellar radius in height.

In the mechanical equilibrium model of Jardine & van Ballegoijen (2005), prominences form after reconnection occurs in the open field, and it is a straightforward matter to show that cool, dense loops can reach heights well beyond corotation (the radius at which the orbital velocity equals the stellar rotational velocity) and the closed stellar coronal field. Hot loops such as those in-

ferred by the COUP, however, appear to need the support of the stellar field, and thus should be limited to heights within the closed coronal magnetic field: a fraction of the heights inferred by UCL models. It was initially suggested that perhaps these hot, large-scale loops are anchored to circumstellar disc material for stability. Testing this idea, follow-up studies of the COUP objects that searched for discs were unable to find conclusive evidence for dusty disc material within the reach of the loops (e.g. Getman et al. 2008; Aarnio, Stassun & Matt 2010). Furthermore, the apparent anticorrelation of large-scale loops and close-in disc material suggests it is the absence of inner disc material that allows the loops to reach greater extents than they would were close-in disc material present (as suggested by Jardine et al. 2006, close-in circumstellar discs could strip away these stars’ outer coronae). Potentially, this indicates some missing physics in earlier models, and it raises the question of how such large loops remain stable for the duration of the X-ray flare decay phase (in some cases, multiple rotation periods).

The potential ramifications of highly energetic events on stellar and circumstellar evolution are far-reaching: large-scale magnetic structures, in their formation and eventual disruption, could regulate stellar angular momentum evolution as well as circumstellar disc evolution, and thus planet formation. The combination of the high frequency with which large-scale magnetic structures have been inferred and the great heights at which they form above the stellar surface indicates circumstances are ideal for the generation of substantial torques opposing stellar rotation. Large-scale magnetic reconnection, in addition to shedding mass and angular

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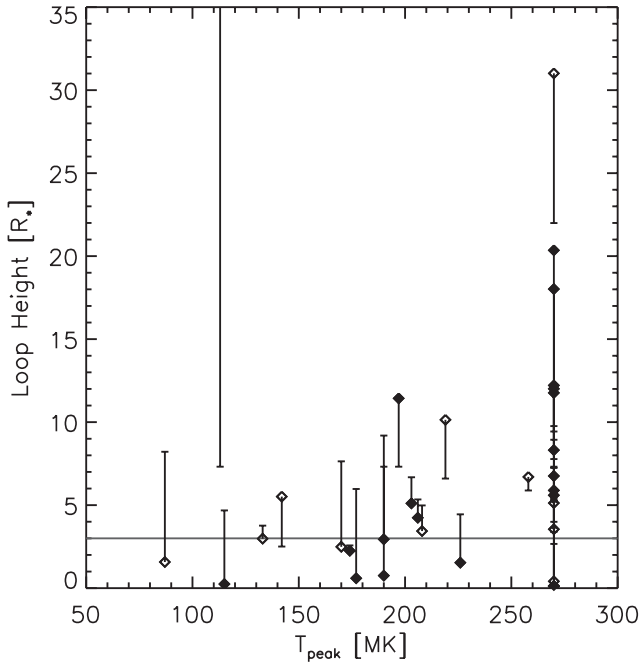


Figure 1. A summary of the results of Favata et al. (2005): inferred loop height as a function of measured loop temperature. The vertical lines connect the loop height to the star’s corotation radius. We have scaled the figure for ease of viewing; the tallest loop (its vertical corotation line shown) is $\sim 63R_*$ in height. The filled points denote cases in which any parameters in calculation of the corotation radius were assumed (i.e. any values of R_* , M_* or P_{rot} not drawn from surveys of the ONC; fiducial values were used as discussed in Section 3.1). The grey horizontal line at $3R_*$ indicates the location of the source surface (the extent of the closed coronal magnetic field) in our models. There is a pileup of points at 270 MK because, to be conservative, the authors adopted a maximum best-fitting temperature and applied it to all cases above a given threshold (beyond this temperature, the fits to their X-ray data are statistically equivalent, and a higher temperature simply yields a shorter loop).

momentum, also bombards circumstellar material with high-energy photons and particles. This could potentially play a pivotal role in planet formation by stimulating grain growth (cf. flash-heated chondrule formation, Miura & Nakamoto 2007). High-energy radiation and particle fluxes can influence disc chemistry (cf. solar nebula isotopic abundances, Feigelson, Garmire & Pravdo 2002) and structure: Glassgold, Najita & Igea (1997) showed that hard stellar X-rays could potentially ionize the disc sufficiently for accretion to occur via magnetohydrodynamic instability. The large heights inferred for these loops would put X-ray-emitting plasma closer to the disc, reducing the column depth through which the hard X-rays would have to travel before ionizing disc material.

Given these potential impacts of large-scale magnetic loops (and their destabilization) on early star and disc evolution, we aim to ascertain whether the existence of hot, large, magnetic loops, as inferred from the analysis of the brightest, most energetic events in X-ray surveys of TTSs, is physically plausible. To do this, we address an evident buoyancy deficit in the previous models by including a stellar wind in the model physics. We do not distinguish between the single-loop UCL model or the possibility of multiple loop events; we only seek to demonstrate that it is possible to find hot loops beyond the typical ‘quiescent’ X-ray-emitting corona, and in some cases, beyond corotation.

2 MECHANICAL EQUILIBRIUM MODEL

Cool, extended prominences have been observed in abundance on the young rapid rotator AB Dor, ranging in height from 2 to $5R_*$, well above the corotation radius of $1.7R_*$ (Collier Cameron & Robinson 1989), and slingshot prominences much like those seen on the Sun have been observed on TTSs (e.g. Massi et al. 2008; Skelly et al. 2008). From a theoretical perspective, other groups (Jardine & Collier Cameron 1991; Jardine & van Ballegoijen 2005, hereafter JCC91 and JvB05, respectively) have assessed the stability of cool loops, determining where in parameter space mechanical equilibria can be found. Loops found beyond the closed coronal field and possibly even beyond corotation are likely embedded in the stellar wind (these loops could form via reconnection of open field lines in the wind; see JvB05, fig. 2).

This work builds upon the prescription of JCC91 and JvB05, save our addition of a stellar wind (Section 2.2). We list below the four major forces our model accounts for (and the parameters upon which each depends) that dictate the height and shape of a magnetic loop:

- (i) gravity (M_* , R_* , ω);
- (ii) buoyancy (ρ , g , T , $\partial p/\partial y$);
- (iii) magnetic pressure gradient (assumed field structure);
- (iv) magnetic tension (footpoint separation, external field strength and structure).

A magnetic loop is in equilibrium when the pressure gradients internal and external to the loop, magnetic tension, and gravity are evenly balanced: the loop is stable, neither expanding nor contracting. We describe the general mechanical equilibrium model, and the addition of a stellar wind term. In Section 3, the case of a simple TTS is demonstrated.

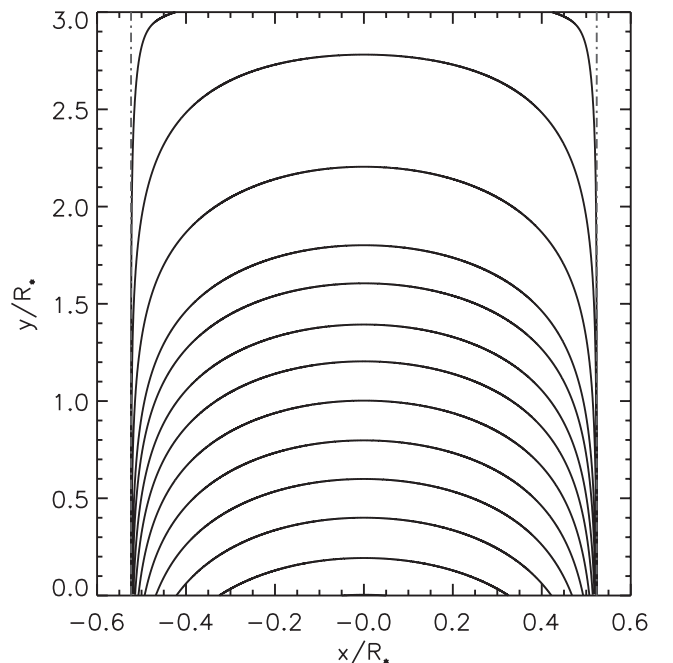


Figure 2. A graphical representation of the form of the external field (see equation 2). Our chosen source surface height, where the magnitude of the field is zero, is $3.0R_*$; beyond this height, the field is open. The vertical lines indicate the width of the arcade ($|x| = \pi/k$, where, here, we choose $k = 3$ so the loop covers 60° in longitude about the stellar equator).

2.1 Equilibrium conditions

We incorporate a number of fairly standard, simplifying assumptions in our 2D model. First, the arcade in which the loop is embedded is not twisted [i.e. no shear; Browning & Priest (1984) found, in the solar case, that including twist in the flux tube only serves to reduce the maximum loop footpoint separation]. Secondly, the external field is potential (i.e. current free) and closed up to a height y_s ; beyond this point, the source surface, the external field is open. The source surface sets, then, the maximum height quiescent coronal loops can reach. The magnetic loop itself is assumed to be narrow, its width less than the length-scale of the external field, and it is isothermal. Finally, the loop is treated as non-disruptive of its surroundings. The equation of motion for plasma in the corona is

$$0 = -\nabla \left(p + \frac{B^2}{8\pi} \right) + \frac{1}{4\pi} (\mathbf{B} \cdot \nabla) \mathbf{B} + \rho \mathbf{g}, \quad (1)$$

where the first term expresses the plasma and magnetic pressure gradients, the second term represents the magnetic tension, and the final term is the balance of gravitational and centrifugal forces. In all the following derivations, we refer to quantities inside the loop as ‘internal’ and coronal properties as ‘external’ (variable subscripts i and e , respectively).

We choose a simple external field which is potential and two dimensional:

$$B_{ex} + iB_{ey} = B_0(e^{2ikz} + e^{-2ky_s})^{1/2}, \quad (2)$$

where $z = x + iy$ (see Fig. 2, illustration of field) and B_0 is the field strength at the stellar surface. Rewriting the Lorentz force in terms of the current density, \mathbf{j} , we have

$$\nabla \left(\frac{B^2}{8\pi} \right) + \frac{1}{4\pi} (\mathbf{B} \cdot \nabla) \mathbf{B} = \mathbf{j} \times \mathbf{B}. \quad (3)$$

Since the Lorentz force is zero in the direction of the field (hence $\mathbf{j} \times \mathbf{B} = \mathbf{0}$), we can combine equations (1) and (3) to obtain an expression for pressure balance along the loop:

$$\frac{dp}{ds} = \rho g_s, \quad (4)$$

where $g_s = \mathbf{g} \cdot \mathbf{B}/|\mathbf{B}|$ is the component of gravity along the loop. In a Cartesian coordinate system where $\mathbf{g} = (0, g)$, this reduces to

$$\frac{dp}{dy} = \rho g(y). \quad (5)$$

We assume the plasma along a given loop is in hydrostatic equilibrium, so we solve for it as

$$p = p_0 \exp \left(\frac{\mu m_H}{k_B T} \int_0^y g dy \right), \quad (6)$$

where p_0 is the loop base pressure (i.e. where $y = 0$).

We note that pressure balance across the loop requires

$$B_i^2 = B_e^2 + 8\pi(p_e - p_i). \quad (7)$$

Following the derivation and boundary conditions of JvB05, we consider gradients normal to the field by taking the scalar product of $\hat{\mathbf{n}}$ and both sides of equation (3), arriving at an expression that governs the shape of the loop:

$$\left(y' \frac{\partial}{\partial x} - \frac{\partial}{\partial y} \right) \frac{B_i^2}{2} = -\frac{y''}{1 + (y')^2} B_i^2. \quad (8)$$

The loop shape [given by $y(x)$] is therefore determined by the requirement that the Lorentz force is zero and hence that the magnetic pressure gradients are balanced by the magnetic tension. At

the loop summit, magnetic tension acts directly downwards, so we must have $\partial B^2/\partial y < 0$. The behaviour of B_i^2 , however, is determined through pressure balance (equation 7) by the nature of the external magnetic field and the variation of the internal and external plasma pressures. The requirement that $\partial B_i^2/\partial y < 0$ reduces to

$$\frac{1}{8\pi} \frac{\partial B_e^2}{\partial y} < (\rho_i - \rho_e)g. \quad (9)$$

From equation (9), we begin to see some necessary physical conditions for stable loops. The maximum loop height is constrained by how buoyancy [i.e. the gradient of the pressure difference, $\partial(p_e - p_i)/\partial y$] depends on height. In the absence of a wind, the gas pressure external to the loop falls with increasing height until a turning point is reached at corotation: at this radius, centrifugal forces dominate over gravitational, and the buoyancy changes sign. If equation (9) is satisfied, the loop can be supported by the external field; if not satisfied, the loop will lose equilibrium.

Meeting these conditions becomes a problem in the T Tauri case, as densities measured in hot loops are greater than ambient coronal densities, the internal temperatures are likely greater than the external temperatures, yet the inferred heights put them beyond corotation. By our previous understanding of the physics, this combination of parameters would not result in a mechanically stable loop within a hydrostatic corona.

For cool prominences, equilibria existed beyond corotation only for negatively buoyant loops; in addition to the loop density being greater than its surroundings, the plasma β (ratio of gas to magnetic pressure) is sufficiently large, but less than unity. We can adjust the plasma β in the previous models and find solutions for tall, hot loops, but the value of β required for solutions beyond corotation is unphysical; for β values consistent with measured parameters on TTSs, hot loops are positively buoyant to large distances. To provide additional upward force to support the hot loop, we add a stellar wind external to the loop. In this case, hot loops can reach a greater range of heights because the wind term extends the range over which buoyancy is negative.

2.2 Including a stellar wind

In order to maintain stability of very extended, hot loops, the external plasma pressure must decrease approaching and passing corotation in order to allow the loop to maintain negative buoyancy at large heights. Far enough from the star, the external magnetic pressure becomes negligible, and so the internal field strength (and thus the loop shape) is determined by the plasma pressure difference. Now, we include an additional term to include a lowering of the pressure due to the flow of a stellar wind along the open external field in which the magnetic loop is embedded; this will provide more upward force on the loop, potentially allowing greater heights to be obtained by loops with a wider range of parameters (consistent with inferred parameters of TTS magnetic loops).

We have added to the basic formulation of JvB05 a simple, Parker wind (Parker 1964). Assuming an inviscid parcel of gas is corotating with a star, the total force on that parcel is

$$\mathbf{F} = -\nabla p + \rho \mathbf{g} + \mathbf{j} \times \mathbf{B}. \quad (10)$$

As this is an expression of Newton’s second law, it can be rewritten as follows:

$$\mathbf{F} = \rho \frac{D\mathbf{v}}{Dt}, \quad (11)$$

where $\frac{D}{Dt}$ is the substantive derivative,

$$\frac{D}{Dt} = \frac{\partial}{\partial t} + (\mathbf{v} \cdot \nabla). \quad (12)$$

We will assume here both a potential field and a steady flow, so the momentum equation along the direction of the loop (i.e. parallel to the direction of the field: $\hat{s} = \mathbf{B}/|\mathbf{B}|$) is

$$\rho \hat{s} \cdot (\mathbf{v} \cdot \nabla) \mathbf{v} = -\hat{s} \cdot \nabla p + \rho \mathbf{g} \cdot \hat{s}. \quad (13)$$

Under the assumption of an isothermal gas, density and pressure are related by the local sound speed of the plasma, $\rho = \frac{p}{c_s^2}$. Conservation of mass requires that the flow across a given cross-sectional area remains constant (i.e. $\rho v A = \text{constant}$), and similarly the conservation of magnetic flux dictates no sinks or sources of magnetic field ($BA = \text{constant}$).

Bearing these relationships in mind, the momentum equation along the loop can be rewritten as

$$\frac{\partial}{\partial s} \left(\frac{v^2}{2} \right) = \frac{-c_s^2}{p} \frac{\partial p}{\partial s} + \mathbf{g} \cdot \hat{s}. \quad (14)$$

Integrating equation (14) from the loop base upwards, we obtain the following expression relating the external field with pressure:

$$0 = \ln(p_e) + \frac{1}{2} M^2 \left(\frac{B_e}{p_e} \right)^2 - \frac{1}{2} M^2 - \frac{1}{c_s^2} \int g dy. \quad (15)$$

where $M = \frac{v_0}{c_s}$ is the initial sonic Mach number of the flow. The external pressure and field, p_e and B_e , are scaled by their base values (p_0 and B_0 , respectively).

If the initial Mach number is 0, the momentum equation reduces to the hydrostatic case (equation 6). The only physical solution of equation (15) is the transonic solution, which gives a slow, subsonic speed close to the star, passes through a sonic point (where the flow velocity equals the local sound speed), and is supersonic at greater distances from the star (see fig. 1 of Parker 1965). To illustrate how including a stellar wind can impact the maximum possible loop height, we consider the specific case of a TTS below.

3 SPECIFIC CASE: A T TAURI STAR

We have applied the mechanical equilibrium model plus wind to a generic, simple TTS. The basic outline of the model and analyses over robust parameter spaces are described fully in previous work (JCC91; Unruh & Jardine 1997, JvB05); here we give full attention to the new addition of a stellar wind term and its effects on the physics of the model. First, we discuss possible equilibria for hot loops embedded within the closed corona; secondly, we assess maximum heights for loops extending beyond the source surface.

3.1 Fiducial T Tauri parameters

As input parameters to the model, we require stellar mass, radius and rotation period; we have adopted representative mean values from the Orion Nebula Cluster (ONC) from the study of Hillenbrand (1997). The mean mass of members (with 50 per cent membership probabilities or greater) is $0.55 M_\odot$, while the mean radius is $1.7 R_\odot$. Our fiducial TTS rotation period of 5.3 d is derived from the assembled rotation period data of Getman et al. (2005). We adopt a wind temperature of 2 MK, consistent with T Tauri wind observations (Dupree et al. 2005). Internal loop temperatures have been taken from the range of values measured for the COUP ‘superflares’ inferred by Favata et al. (2005); these vary from ~ 50 to 200 MK (see Fig. 1). We set the initial Mach number to 0.05. For the extent

of the closed coronal field, we show our results for $y_s = 3R_*$, which is within the corotation radius ($5.2R_*$). It is unclear for TTSs how far from the stellar surface the closed corona may extend, but as we note in Section 4, our results do not depend strongly on the choice of source surface height (see Section 3.3 for further discussion of y_s).

3.2 Pressure and buoyancy

For the hydrostatic case, the functional form of the external plasma pressure is exponentially decreasing close to the star, with the argument of the exponential function dependent upon gravitational and centrifugal forces. In the case of a TTS specifically, given the same mass and rotation period as a main-sequence counterpart but twice the radius, the effective gravity is approximately four times weaker, thereby extending the distances to physically important heights. The corotation radius for the TTS is a factor of 2 closer, and the difference in pressure (i.e. $p_e - p_i$) remains below the base pressure difference to a distance approximately three times closer than the main-sequence star. The corotation radius is important because it is where pressure difference and buoyancy change sign.

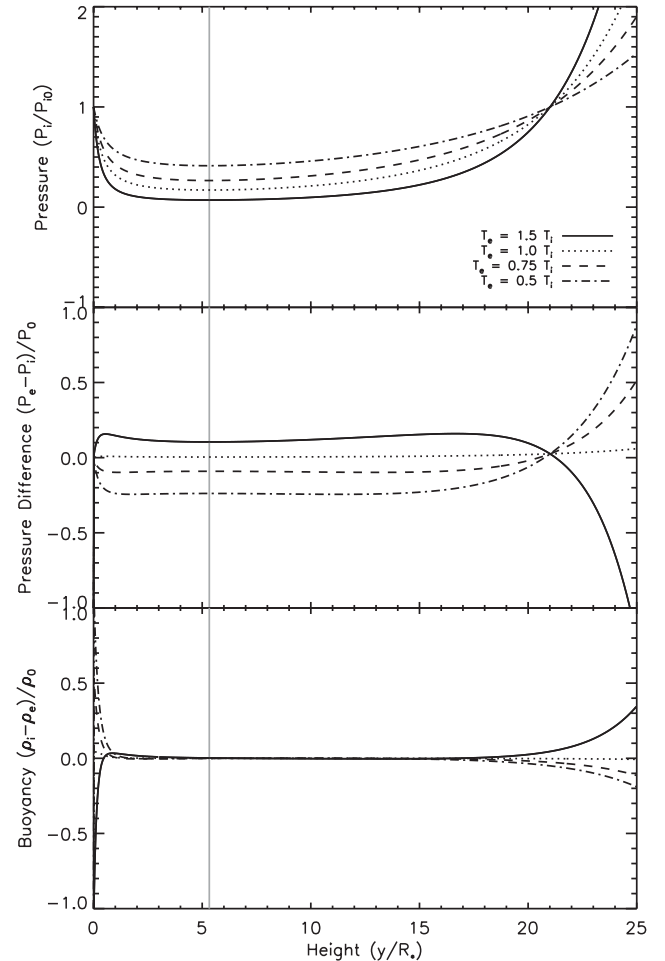


Figure 3. Pressure, pressure difference and buoyancy as a function of height (upper, middle and lower panels, respectively). The external plasma temperature is 2 MK, and temperature ratios (T_e/T_i) shown are 1.5 (solid line), 1.0 (dotted line), 0.75 (dashed line) and 0.5 (dot-dashed line). The corotation radius is located at $5.2R_*$ (grey vertical lines); at $20.2R_*$ the internal pressure equals its base value, and the pressure difference changes sign.

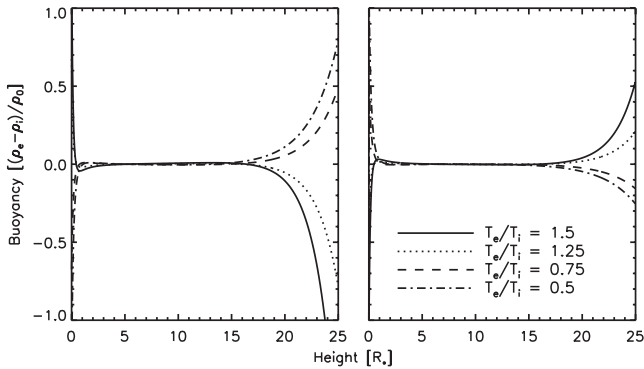


Figure 4. A comparison of loop buoyancy using the physics of JvB05 (left-hand panel) and our present model including a stellar wind (right-hand panel). For our fiducial T Tauri parameters, under the previous model physics, loops much hotter than their surroundings were negatively buoyant only fairly close to the stellar surface, within the source surface.

We choose a base pressure and temperature ratio, and with these we solve equation (15) for the equilibrium shape of the loop. In Fig. 3, we show, for the parameters outlined above, the internal pressure, pressure difference ($p_e - p_i$) and buoyancy $[\partial(p_e - p_i)/\partial y]$ as a function of height for various temperature ratios T_e/T_i . It is important to note the heights at which the loop’s buoyancy is negative when including the stellar wind; in Fig. 4, we illustrate how our addition of the stellar wind affects loop buoyancy. Where cool prominences had negative buoyancy out to great heights previously (Fig. 4, left-hand panel, solid/dotted lines), we now see that including the wind means hot loops have negative buoyancy at large distances from the stellar surface (Fig. 4, right-hand panel, dashed/dot-dashed lines).

3.3 Loop shapes and maximum heights

The critical parameter we are altering in this analysis is the external pressure (and thus the pressure balance of the loop with its environs). We have improved on previous models that assumed a hydrostatic external plasma pressure to include the effect of a wind. The height of the source surface (the height beyond which all field lines are open) is not a well-constrained parameter, as interpretations of multiwavelength observations imply conflicting values. Solar observations show a white light corona which extends to greater heights than the X-ray-bright corona. In the stellar case, densities derived from X-ray emission measures suggest more compact coronae, while the lack of apparent rotational modulation in some cases, as well as the formation of large-scale prominences, indicates the X-ray coronae are extended. Our choice reflects a field more extended than main-sequence, solar-type stars, but less so than a purely dipolar field (based on analysis of the relationship between field complexity and the observed coronal X-ray emission measure, Jardine et al. 2006).

In Fig. 5, we show the shape of hot loops when the source surface is placed $3R_*$ from the stellar surface. Within the closed corona, the pressure gradient of the external field is directed upwards, providing support for the loop against the downward force of magnetic tension. This set of solutions are effectively identical to those found for hot loops by JvB05; we too find that equilibrium solutions can be found for both negatively and positively buoyant loops out to the source surface. The coolest of the loops plotted (30 and 40 MK) are able to reach heights above the source surface and above corotation due to the extra support of the wind.

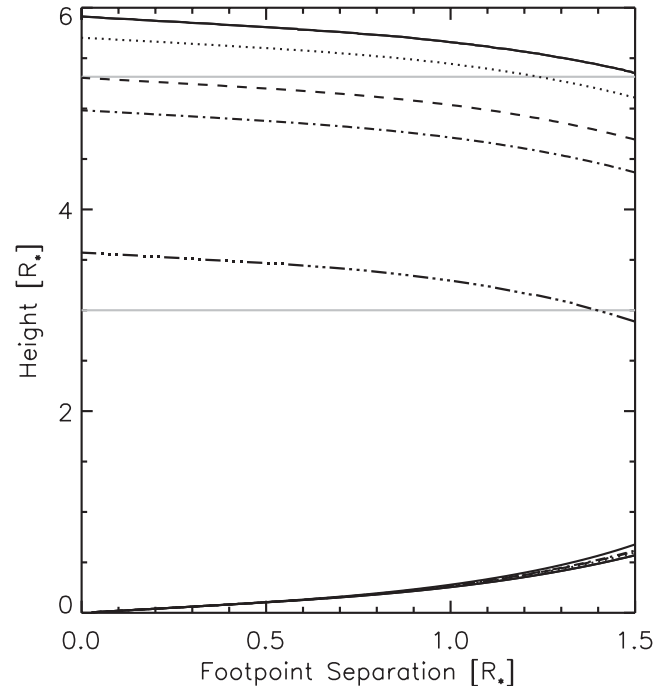


Figure 5. Loop heights as a function of footpoint separation. Our loop temperatures probe the parameter space illustrated in Fig. 1: for loop temperatures from 30 to 200 MK, we find loops with heights above the closed coronal field (consistent with the inferences of Favata et al. 2005). The loop temperatures are: 30 MK (solid line), 40 MK (dotted line), 50 MK (dashed line), 60 MK (dot-dashed line) and 200 MK (triple-dot-dashed line). The grey, horizontal lines denote the source surface (extent of closed coronal field) and corotation radii at $3R_*$ and $5.2R_*$, respectively. T_e is 2 MK, the plasma β is 0.05, and the stellar parameters used are as described in Section 3. For all loops shown here, p_i is ~ 10 – 20 per cent lower than p_e . Footpoint separations beyond $1.05R_*$, though mathematically possible, are non-physical solutions, as we have defined the arcade width to be $x = \pi/k$, where $k = 3$. For a fairly narrow range of loop temperatures, we find a wide range of possible maximum loop heights, the coolest loops reaching just beyond corotation, and the hottest loops confined within the closed corona.

4 CONCLUSIONS

We have addressed the issue of mechanical equilibrium for large, hot loops. Since cases have been found in which hot loops several stellar radii in height seem to not be disc supported, we propose that these loops reach mechanical equilibria beyond the closed coronal magnetic field due to the stellar wind. We have shown that, for a generic TTS, our equilibrium loop heights are consistent with the inferred loop lengths for the most powerful flares observed by the COUP. We also recover solutions for hot loops within the range of heights below the source surface anticipated by previous models, suggesting that large and small stable post-flare loops can occur under similar conditions. In the physics of previous models, which were hydrostatic and did not consider the role of stellar winds, loops hotter than the quiescent corona were confined to heights within the closed coronal field, needing its support. Our present work indicates that, given a stellar wind, hot loops can be found at heights well beyond the source surface, and even beyond corotation.

Young stars are known to be very magnetically active in the T Tauri phase, with frequent flaring and high X-ray luminosities. The apparent abundance of large-scale loops, favourable conditions for their formation in TTS coronae and winds, and potential for consequence on pre-main-sequence stellar (and circumstellar) evolution

make understanding these large-scale loops very important. Having verified that such large-scale magnetic structures can indeed exist, interesting subsequent work will be to determine how the loops' disruption can impact circumstellar disc material as well as the star itself. Large-scale disruptions of this nature could effect stellar rotational evolution and activity as well as circumstellar disc evolution and planet formation.

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