

RADIATIVE IMPACT OF CRYOSPHERE ON THE CLIMATE OF EARTH AND MARS

by

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Dedication

To my parents

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ABSTRACT

Radiative Impact of Cryosphere on the Climate of Earth and Mars

by

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Chair: Mark Flanner

Snow- and ice-covered surfaces are the most reflective regions on Earth and Mars, and their extent can change substantially with small changes in climate. The presence of Earth's cryosphere greatly alters the planet's albedo and changes in cryospheric extent and reflectivity therefore partially determine the sensitivity of climate to anthropogenic and external forcings. Carbon dioxide ice is abundant on the Martian surface, and plays an important role in the planet's energy budget.

Firstly, we quantify the shortwave Cryosphere Radiative Effect (CrRE) on Earth. Relatively high resolution ($0.05^\circ \times 0.05^\circ$) MODIS (Moderate Resolution Imaging Spectroradiometer) data along with radiative kernel datasets are used to estimate the global land shortwave CrRE. We perform multiple analyses to determine the sensitivity of our estimates to the use of different thresholds for snow cover determinations, different climatologies for missing data, and radiative kernels generated with different distributions of clouds produced with various versions of the Community Atmosphere Model. We estimate a global land-based CrRE of about -2.6 Wm^{-2} during 2001-2013, with about 59% of the effect originating from Antarctica.

Secondly, we adapt the terrestrial Snow, Ice, and Aerosol Radiation (SNICAR) model to simulate CO_2 snow albedo across the solar spectrum ($0.2\text{-}5.0 \mu\text{m}$). We apply recent laboratory derived refractive indices of CO_2 ice, which produce higher broadband

CO₂ snow albedo (0.93–0.98) than previously estimated. We perform multiple analyses to determine the sensitivity of cryosphere spectral albedo to the amount and type of dust, co-presence of CO₂ and H₂O ices, ice grain size, snow layer thickness, and solar zenith angle. In addition, we also compare our simulations with observed Mars surface albedo, and achieved a reasonable fit between the two.

Finally, SNICAR is implemented with the Laboratoire de Météorologie Dynamique Mars GCM to prognostically determine ice cap (both H₂O and CO₂) albedos interactively in the model. We then explore the impact of dust on surface cryosphere albedo and its impact on Mars' shortwave energy flux. After integrating SNICAR into the Mars GCM, we find that the impact caused by dust is about 1.5 times the impact caused by the presence of snow itself on shortwave flux at the surface.

Chapter 1

Introduction

1.1 Albedo

Albedo or reflection coefficient (from the Latin *albus* for "white") is defined as the ratio of reflected radiation from the surface to incident radiation upon it. Albedo depends on the frequency of the radiation and the properties of the surface. In general, the albedo also depends on the directional distribution of incident radiation, except for Lambertian surfaces, which scatter radiation uniformly in all directions and have an albedo that is independent of the incident distribution.

Black-sky albedo (directional hemispherical reflectance) is defined as albedo in the presence of unidirectional incident light and is a function of solar zenith angle. It is the integration of the bi-directional reflectance over the viewing hemisphere. Bidirectional Reflectance Distribution Function (BRDF) gives the reflectance of a target as a function of illumination geometry (solar zenith angle for planetary albedo), and viewing geometry (e.g., point of view of a satellite). The BRDF depends on wavelength and is determined by the structural and optical properties of the surface, including effects such as shadow-casting, multiple scattering, mutual shadowing, transmission, reflection, absorption, and emission by surface elements, facet orientation distribution, and facet density. White-sky albedo (bi-hemispherical reflectance) is defined as albedo in the absence of a direct

component when the diffuse component is isotropic. It is the integration of the directional albedo over the illumination hemisphere.

Black-sky albedo (α_{bs}) and white-sky albedo (α_{ws}) mark the extreme cases of completely direct and completely diffuse illumination. The actual albedo is a value which is interpolated between these two, and can be written as:

$$\alpha = f_{dif} \cdot \alpha_{ws} + (1 - f_{dif}) \cdot \alpha_{bs} \quad (1.1)$$

where f_{dif} is the fraction of incident skylight that is diffuse, which is itself a function of the atmospheric optical depth at the frequency of interest [e.g., *He et al.*, 2012]. These terms are important because they allow the albedo to be calculated for any given illumination conditions from a knowledge of the intrinsic properties of the surface [*Roman et al.*, 2010]. Two common albedos used in astronomy are discussed below:

1.1.1 Geometric Albedo

Geometric albedo is the ratio of the global brightness of a planet, viewed from the direction of the Sun, to that of a hypothetical, white, and fully reflecting Lambertian sphere of the same size at the same distance [*Morrison and Owen*, 1996]. From *Chesley et al.*, [2002], the relationship between geometric albedo (p_v) with absolute magnitude (H), and diameter (D) of a celestial object is:

$$D = 1329 \text{ km} \times 10^{-H/5} p_v^{-1/2} \quad (1.2)$$

Absolute magnitude (dimensionless) is the measure of intrinsic brightness of a celestial object. Absolute magnitude is defined to be the apparent magnitude an object would have if it were located at a distance of 10 parsecs (32.6 light years) from Earth. The apparent magnitude of a celestial object is a number that is a measure of its brightness as

seen by an observer on Earth. The brighter an object appears, the lower its magnitude value. Some objects can be more reflective than the hypothetical idealized reference sphere, therefore geometric albedos higher than 1.0 are possible.

1.1.2 Bond Albedo

Bond albedo (named after US astronomer George P. Bond), is defined as the fraction of the total incident solar radiation reflected by a planet back into space. Bond albedo is a measure of the planet's solar energy balance and dependent on the spectrum of the incident radiation defined over the entire range of wavelengths. Since Bond albedo (A) tells us how much solar energy is absorbed, we can use it to calculate an effective temperature (T_e) of a planet using Stefan-Boltzmann law:

$$T_e = \left[\frac{(1 - A)I_0}{4\sigma} \right]^{1/4} \quad (1.3)$$

where I_0 is the incident solar flux-density (insolation) on the planet, and σ is the Stefan-Boltzmann constant ($5.67 \times 10^{-8} \text{ Wm}^{-2}\text{K}^{-4}$).

Bond albedo can never be greater than 1.0. Enceladus, a moon of Saturn, has a geometric albedo of 1.38 [Verbiscer *et al.*, 2007] and Bond albedo of 0.9 [Buratti and Veverka, 1984; Verbischer and Veverka, 1994]. For rest of the introduction, Bond albedo will simply be referred to as albedo.

1.2 Cryosphere

The portion of Earth's surface covered by frozen water (e.g. ice caps, glaciers, etc.) is known as cryosphere. In this thesis, I will use the term *cryosphere* to refer to ice or snow present on a planet's surface due to deposition or freezing of any chemical compound or gases, e.g., water, carbon dioxide, methane, and nitrogen. Due to a very high albedo (H_2O

~0.8, CO₂ ~0.97) of pure snow, the cryosphere can play a significant role in driving global climate of any planet. Albedo feedback is also one of the most important positive feedbacks within the Earth's climate system (Chapter 2), which reinforces the importance of cryosphere in a climate system.

1.2.1 Earth

Earth's cryosphere consists of frozen water in various forms, including sea ice, lake ice, river ice, snow cover, glaciers, ice sheets, and frozen ground (which includes permafrost). Earth's permanent ice sheets are presently located on Greenland and Antarctica, with smaller glaciers present elsewhere at high elevations and latitudes. The cryosphere reflects a significant part of incoming solar energy as compared to open water (albedo ~0.07). The cryosphere exchanges heat with warmer areas of the planet, and plays a significant role in the global climate in response to global changes. Being a strong feedback mechanism, albedo feedback can significantly delay or amplify a global seasonal change like melting of glaciers.

1.2.2 Mars

Carbon dioxide ice is abundant on the Martian surface [*Leighton and Murray, 1966; Herr and Pimentel, 1969; Larson and Fink, 1972; Forget, 1998; Bibring et al., 2005*], and plays an important role in the planet's energy budget due to its high reflectivity (pure CO₂ snow albedo of ~0.97) and seasonal variation. H₂O ice is also present at the surface of the perennial ice caps, though in lesser abundance than CO₂ ice, and with small seasonal deposition in other parts of the planet [e.g., *Kieffer et al., 2000; Bibring et al., 2004; Brown et al., 2014*]. Due to the lack of any hydrogen bonding, CO₂ does not exist in liquid phase,

so on Mars the transition occurs between two phases of CO₂, viz. solid and gas, resulting in a large amount of energy exchange during the phase changes.

1.2.3 Other celestial bodies in our solar system

Many other celestial bodies in our solar system have also shown evidence of cryospheric presence on their surface. Moons of Jupiter including Europa, Ganymede, and Callisto; Moons of Saturn including Titan, Hyperion, and Enceladus; and Moons of Uranus including Titania, Umbriel, and Oberon have all shown evidence of water ice (along with ices from other chemical compounds) present on their surface [Fink *et al.*, 1973; Spencer, 1987; Cruikshank, 1980; Buratti and Veverka, 1984; Verbischer and Veverka, 1994]. Titan is known to have liquid hydrocarbon lakes, making it the only other body in our solar system with presence of stable liquid at the surface [Mitri *et al.*, 2007; Stofan *et al.*, 2007]. Enceladus has been found to be one of the most geologically dynamic objects in the solar system [Kargel, 2006]. Recently, evidence for presence of H₂O, CH₄ and N₂ ice has been found at the surface of Pluto [Stern *et al.*, 2015; Cruikshank *et al.*, 2015; Trafton, 2015].

Enceladus, Europa and Titan also show evidence of cryovolcanism [Kargel, 2006; Spencer *et al.*, 2009; Greeley, 1998; Quick *et al.*, 2013; Sotin *et al.*, 2005; Mitri *et al.*, 2007; Stofan *et al.*, 2007]. A cryovolcano is a volcano that erupts volatiles such as water, ammonia or methane, instead of molten rock. The presence of cryosphere and cryovolcanism are reasons for relatively higher surface albedo of these bodies. Although Venus also has a high albedo of ~0.75, it does not have a cryosphere present on its surface. On the contrary, Venus is one of the hottest planets (mean surface temperature of 735 K) in our solar system, caused by the high amount of greenhouse gases in its atmosphere. Venus's high albedo is caused by thick, highly reflective clouds present in the planet's

atmosphere. These clouds are made up of sulfur dioxide, and concentrated sulfuric acid [Pierrehumbert, 2010]. Interestingly, the night side of Venus' atmosphere above 100 km is so cold (temperature as low as 100 K) that it was named the 'cryosphere' [Bertaux *et al.*, 2007].

1.3 Climate of Mars

The Martian atmosphere consists primarily of carbon dioxide with small percentages of nitrogen (N₂), argon (Ar), and trace amounts of other gases like oxygen, water vapor, carbon monoxide, and hydrogen. The mass of Mars is about 11% that of Earth, and it has gravitational acceleration about 38% that of Earth's. Therefore Mars has a very thin atmosphere with mean surface pressure of about 6 mbar (about 0.6% of Earth's mean sea level pressure). Mars' thin atmosphere and its greater distance from the sun mean that Mars is much colder than Earth. The average temperature is about -55 °C (218 K), although it can vary from -153 °C (120 K) near the poles during the winter to as much as 20 °C at midday near the equator [Hess *et al.*, 1977]. Actual temperature measurements at the Viking landers' site range from -143 °C (130 K) to 17 °C (290 K) [Kieffer *et al.*, 1977].

The axial tilt of Mars' (25°) is nearly the same as that of Earth (23.5°). Mars therefore experiences similar seasonal amplitudes, except the seasons are roughly twice as long as on Earth. In each hemisphere a polar ice cap forms each fall and winter under an obscuring layer of clouds, reaching its maximum extent at the start of spring. It then recedes with the seasonal warming, shrinking to its residual cap by the end of summer. The orbital eccentricity of Mars' orbit is 0.1, much greater than Earth's present eccentricity of about 0.02. Due to the large eccentricity, winters in the southern hemisphere are long and cold while those in the North are short and relatively warm. The planet receives a lesser amount

of net solar energy than the Earth (Mean Solar constant = 590 W/m^2) due to larger distance from Sun.

Another unique feature of Martian climate is dust storms. Although Mars' atmosphere is thin, it supports weather, clouds, and winds. Giant dust devils routinely lift the oxidized iron dust that covers Martian surface. Dust storm activity can contribute to regional surface albedo changes and the largest of these storms can generate a cloud of dust that can encircle both hemispheres of the planet for more than a Martian season [*Leovy et al.*, 1973; *Cantor*, 2007]. Most of the dust storms' generation can be explained in terms of the classical atmospheric tidal theory developed for the earth's atmosphere [*Zurek and Leovy*, 1981; *Ryan et al.*, 1981]. The radiative heat of sunlight reaching the surface of the planet is fundamentally what drives these dust storms. As sunlight is absorbed by the ground, it warms the air closest to the surface, while the upper air remains cooler. Thus, the Martian atmosphere becomes unstable, with warm air rising and lifting dust with it. Once dust is in the atmosphere, it provides a potent source for heating which have an impact on for vertical and horizontal winds. Rising plumes of warm air create weather systems ranging from small dust devils, similar to those that form in deserts on Earth, to larger continent-sized storms. These larger storms sometimes combine into global storms which cover the entire planet in atmospheric dust [*Leovy et al.*, 1973; *Ryan et al.*, 1981]. Due to higher eccentricity of orbit, these larger storms typically only occur during summer in Mars' southern hemisphere, when radiative heat forces are strongest. Once started, bigger storms can last weeks to months.

1.4 Mars General Circulation Model (GCM)

Mars GCMs have been developed to study the nature of the general circulation of the atmosphere of Mars, how that circulation is driven, and how it affects the climate of Mars in the long term. I use a Mars GCM to determine the surface albedo, and its impact on Martian climate (Chapter 4).

The GCM of the Martian atmosphere developed by the Laboratoire de Météorologie Dynamique (LMD) of the CNRS in Paris in collaboration with the Atmospheric and Oceanic Planetary Physics sub-department in Oxford [Hourdin *et al.*, 1993, 1995; Forget *et al.*, 1998, 1999] is one of the most widely used 3D models of Martian climate by the scientific community. The current model includes the *Non-local thermodynamic equilibrium* (NLTE) radiative transfer code [López-Puertas and López-Valverde, 1995] valid at up to 120 km altitude, tracer transport, the water cycle with water vapor and ice, the "double mode" dust transport model (Madeleine *et al.*, 2011), and optional photochemistry and extension into the thermosphere up to 250 km altitude. The model uses a sigma- P (sigma-pressure) hybrid coordinate which is equivalent to sigma coordinates near the surface, and gradually shifting to purely pressure coordinates with increasing altitude. The distribution of the vertical layers is irregular, to enable greater precision at ground level. The resolution of the model is 64×48 grid cells horizontally, corresponding to 3.758 degrees latitude by 5.6258 degrees longitude. 25 levels are used to represent the atmosphere to a height of 80 km, 32 levels for simulations up to 120 km, or 50 levels to include the thermosphere. The first layer describes the first few meters above the ground, whereas the upper layers span several kilometers.

The radiative transfer equation accounting for the absorption and scattering of solar radiation by dust and other atmospheric constituents is based on the numerical scheme of *Fouquart and Bonel* [1980]. The upward and downward fluxes are obtained from the reflectances and transmittances of the layers, computed using the Delta-Eddington approximation. The available studies of Martian dust optical properties have suggested that the single-scattering properties are characterized by an abrupt change around wavelengths of 0.5 μm , near the peak of the incident solar flux [*Ockert-Bell et al.*, 1997]. The LMD Mars GCM therefore uses two broad bands (0.1-0.5 μm and 0.5-5 μm) to resolve spectral dust properties and represent radiative transfer.

The deposition and sublimation of carbon dioxide on the ground is controlled by relatively simple physical processes. When the surface temperature falls below the condensation temperature, CO_2 condenses, releasing the latent heat required to keep the solid-gas interface at the condensation temperature. Conversely, when CO_2 ice is heated, it partially sublimates to keep its temperature at the frost point temperature. Two kinds of mechanisms which lead to atmospheric condensation are radiative cooling and adiabatic cooling [*Forget et al.*, 1998].

1.5 Organization of the Thesis

In Chapter 2 [*Singh et al.*, 2015], we discuss the quantification of the shortwave Cryosphere Radiative Effect (CrRE) on Earth. This study uses relatively high resolution ($0.05^\circ \times 0.05^\circ$) MODIS (Moderate Resolution Imaging Spectroradiometer) data along with radiative kernel datasets as applied by *Perket et al.*, [2014], to derive an estimate of the global land shortwave CrRE. We also perform multiple analyses to determine the sensitivity of our estimates to the use of different thresholds for snow cover determinations,

different climatologies for missing data, and radiative kernels generated with different distributions of clouds produced with various versions of the Community Atmosphere Model (CAM).

In Chapter 3 (*manuscript under review*), we adapt the terrestrial Snow, Ice, and Aerosol Radiation (SNICAR) model to simulate CO₂ snow albedo across the solar spectrum (0.2-5.0 μm). In this study, we apply recent laboratory derived refractive indices of CO₂ ice, which produce higher broadband CO₂ snow albedo (0.93–0.98) than previously estimated. We perform multiple analyses to determine the sensitivity of cryosphere spectral albedo to the amount and type of dust, co-presence of CO₂ and H₂O ices, ice grain size, snow layer thickness, and solar zenith angle. In addition, we also compare our simulations with observed Mars surface albedo.

Finally, in Chapter 4 we describe the integration of SNICAR into the LMD Mars GCM to compute real time albedo in the model. We then explore the impact of dust on surface cryosphere albedo and its impact on Mars' energy budget. Chapter 5 summarizes the key results, and also discusses the relevant research directions for future work.

Chapter 2

The Global Land Shortwave Cryosphere Radiative Effect during the MODIS era

As appears in:

Singh, D., M. G. Flanner, and J. Perket, (2015), *The global land shortwave cryosphere radiative effect during the MODIS era*, *The Cryosphere*, 9, 2057-2070, doi:10.5194/tc-9-2057-2015.

2.1 Introduction

Snow- and ice-covered surfaces are the most reflective regions on Earth, and their extent can change substantially with small changes in climate. The presence of Earth's cryosphere greatly alters the planet's albedo and changes in cryospheric extent and reflectivity therefore partially determine the sensitivity of climate to anthropogenic and external forcings. After water-vapor and cloud feedback, the albedo feedback is the third most powerful positive feedback mechanisms operating within the current climate system [e.g., *Bony et al.*, 2006; *Winton*, 2006; *Randall et al.*, 2007; *Soden et al.*, 2008; *Shell et al.*, 2008; *Flato et al.*, 2013].

Earth's Cryosphere has shown compelling indications of climate change during recent decades, including mass loss from ice sheets and glaciers [e.g., *Rignot et al.*, 2011, *Gardner et al.*, 2013], rapid ablation of autumn Arctic sea-ice [e.g., *Serreze et al.*, 2007; *Stroeve et*

al., 2012], and reduced seasonal snow coverage [e.g., *Dery and Brown*, 2007; *Brown and Robinson*, 2011]. The Arctic is one of the most sensitive regions on Earth to global climate change [*Manabe et al.*, 1992; *Manabe and Stouffer* 1994; *Miller and Russell* 2000; *Meehl and Washington* 1990]. Several recent studies [e.g., *Chapman and Walsh*, 2007, *Monaghan et al.*, 2008, *Steig et al.*, 2009] have shown that climate is also warming over west Antarctica, and is related to Pacific Ocean warming [*Ding et al.*, 2011] and circumpolar winds.

The shortwave Cryosphere Radiative Effect (CrRE) is the instantaneous influence of snow- and ice-cover on Earth's top of atmosphere (TOA) solar energy budget [*Flanner et al.*, 2011; *Hudson*, 2011; *Perket et al.*, 2014]. CrRE depends not only on snow and sea-ice coverage, but also on local insolation, cloud cover, and properties of the snow, ice and their underlying surface that determine reflectance. These features determine the impacts of cryospheric presence on net TOA solar flux [e.g., *Winton*, 2006; *Qu and Hall*, 2005]. Changes in the extent of seasonal snow cover and sea-ice can drive large changes in CrRE on sub-decadal timescales, whereas the areal coverage of ice sheets and glaciers tend to evolve on much longer timescales. The presence of the cryosphere also perturbs Earth's longwave energy budget, e.g., through changes in emissivity and surface temperature resulting from the insulating effect of snow and the change in surface elevation induced by ice sheets. This study, however, concentrates exclusively on the shortwave component of CrRE (hereafter referred to simply as CrRE).

Our work focuses on developing a global, gridded, time-resolved dataset of the land-based CrRE, using modern remote sensing observation of surface albedo and snow presence; combined with radiative kernels that provide TOA radiative impacts. *Flanner et*

al [2011] derived a 30-year record of the northern hemisphere CrRE from coarse-resolution determinations of snow cover extent. Here we apply higher-resolution, higher quality remote sensing data from the MODerate-resolution Imaging Spectroradiometer (MODIS) to derive global land-based CrRE over 2001-2013, helping inform on the utility of the longer-term record derived by *Flanner et al* [2011], and broadening the scope of these estimates to include the southern hemisphere. We generate CrRE for both all-sky and clear-sky conditions to help assess the masking effect of clouds and atmospheric aerosols. We provide statistics including global, hemispheric, glaciated and non-glaciated land CrRE averages. We also perform multiple analyses to determine the sensitivity of our estimates to the use of different thresholds for snow cover determinations, different climatologies for missing data, and radiative kernels generated with different distributions of clouds. These sensitivity analyses help us identify the sources of uncertainty that have relatively high impact on CrRE. In this paper, we focus only on land based CrRE, and refer readers to other recent estimates of CrRE from Arctic sea-ice [*Pistone et al.*, 2014; *Cao et al.*, 2015].

2.2 Methods

2.2.1 Satellite Dataset Used

MODIS MCD43C3 collection 5 surface albedo data [e.g., *Schaaf et al.*, 2002] and accompanying snow coverage statistic (both provided at spatial resolution $0.05^\circ \times 0.05^\circ$) are the primary input datasets we use to generate land based CrRE (LCrRE) at 16-day resolution. The snow presence parameter (0-100%) is a measure of the fraction of native measurements within each 16-day period and each 0.05° pixel in which the presence of snow is detected. We assume any pixel with snow coverage greater than zero has a surface albedo that was affected by the presence of snow. We apply data with quality flag 4 and

better to maximize the spatial and temporal coverage of albedo measurements. Missing data in this collection arise from cloud cover and absence of sunlight at high latitudes during winter. To derive a spatially and temporally continuous LCrRE record, we also apply snow-cover information from the Near-real-time Ice and Snow Extent (NISE) dataset [Nolin *et al.*, 1998], as described in the next subsection. NISE provides daily binary snow-cover at 25 km Equal-Area Scalable Earth Grid (EASE-grid) spatial resolution, projected on a polar stereographic grid. Because it is determined from microwave remote sensing observations, it offers estimates of snow presence under all conditions, including environments with clouds and low illumination. Since these datasets have different spatial resolutions, the NISE dataset has been remapped to the higher resolution MODIS grid prior to analysis.

We use global gridded estimates of snow-free albedo derived from MODIS [Moody *et al.*, 2008], also provided at 16-day and 0.05 degree resolution, as a baseline for determining the surface albedo contrast induced by snow. These data are climatological 16-day averages derived from 2000-2004 MODIS measurements.

2.2.2 Global Albedo Climatology for filling missing data

A 16-day global surface-albedo climatology with spatial resolution of $0.05^\circ \times 0.05^\circ$ is generated to fill any missing MODIS data-points, using the following steps:

Step 1: For each pixel and each 16-day period of the year, we take the average albedo over all the years (2001-2013) of MODIS data during periods with valid measurements and non-zero snow cover. Since this is a seasonally-varying gridded

climatology, it is primarily used to replace albedo of missing MODIS data in situations deemed to be snow-covered in the NISE dataset.

Step 2: We take the annual mean of the albedo values generated from the previous step for each snow-covered unfilled pixel in step 1. We only apply this average at locations and times when the climatology from step 1 does not provide valid data (e.g., at pixels and 16-day intervals that had substantial cloud cover during each of the 13 years of MODIS observations).

Step 3: Albedo values generated from the previous two steps are averaged spatially over all pixels within each land classification type defined in the MODIS MCD12C1 product using IGBP (Type 1) land cover classification. This procedure produces an annual-mean snow-covered albedo climatology by land classification (listed in Table A.1), and is used to fill any remaining missing pixels unfilled by steps 1 and 2.

MCD12C1 provides the global dominant land cover types at $0.05^{\circ} \times 0.05^{\circ}$ spatial resolution. It is continuous and therefore completely eliminates the chance of having any missing pixels after applying step 3.

Although the snow-free surface albedo dataset is continuous, it is undefined in regions with large solar zenith angle or near-permanent snow cover. To define snow-free albedo in these regions, we apply annual averages of the snow-free albedo values for each pixel (similar to method applied in step 2).

The ice-free surface albedo for permanent glaciated areas (e.g. Greenland, Antarctica) is assumed to be 0.26, an average value of snow-free albedo over barren land

(*Flanner et al.*, 2011). This assumption enables a rough estimate of the TOA impact associated with presence of the full ice sheets. While our estimates of absolute LCrRE in these regions are therefore subject to ambiguities (such as the type of vegetation that would thrive without the ice sheet), seasonal and inter-annual changes in glacier surface albedo, e.g., as caused by altered insolation, melt extent, and snow metamorphic state, drive changes in LCrRE that are unaffected by this assumption, since the ice-free albedo is assumed to be static.

Considering the current bed topography removal of the West Antarctic Ice Sheet would expose open-ocean (albedo 0.07). Although, we still use barren land albedo because of several other uncertain effects that would occur in the event of total ice sheet ablation, e.g., isostatic rebound of land, sea-level rise, encroachment of vegetation over open land. Our objective is to quantify the instantaneous LCrRE and quantifying all the aforementioned uncertainties in this study are beyond the scope of this paper. We also note that the MODIS land mask applied in our study excludes ice shelves.

2.2.3 Methodology

Our interest lies in the change in solar energy reflected because of snow, and we therefore assume that measured albedo increase in the presence of snow, relative to the snow free state, is caused entirely by snow. *Flanner et al* [2011] employed a definition of CrRE that utilizes snow cover fraction, in order to facilitate the use of snow extent data without coincident albedo measurements from several decades ago. Here we utilize a simpler definition that omits snow cover fraction, since we have direct measurements of surface albedo from MODIS. Using this approach, the mathematical framework to

describe CrRE at time t within a region R that is composed of N partially snow- or ice-covered grid-cells i is:

$$CrRE(t, R) = \frac{1}{A(R)} \sum_{i=1}^N \text{Max}[\alpha(t, i) - \alpha_{\text{snow-free}}(t, i), 0] \frac{\partial F}{\partial \alpha}(t, i) A(i)$$

(2.1)

where A is area, $\alpha - \alpha_{\text{snow-free}}$ is the albedo contrast ($\Delta\alpha$), and $\partial F/\partial\alpha$ is the change in TOA net solar energy flux with changing surface albedo. The ‘Max’ function is used to avoid any negative albedo contrast values. We determine $\partial F/\partial\alpha$ using radiative kernels that provide the instantaneous effect on TOA energy budget associated with small perturbations in surface albedo [e.g., *Shell et al.*, 2008]. This equation is only applied to gridcells where snow presence has been detected, and LCrRE is otherwise assumed to be zero.

We created kernel datasets using the general framework of *Perket et al.*, [2014], using the Community Atmosphere Model versions 4 and 5 (CAM4 and CAM5). TOA energy fluxes were calculated with and without surface albedo perturbations every model time-step for one year of simulation, and flux differences were then averaged into monthly resolved kernels. We also apply radiative kernels generated previously with the CAM3 model [*Shell et al.*, 2008] and Geophysical Fluid Dynamics Laboratory Atmosphere Model (AM2) [*Soden et al.*, 2008].

A graphical representation of the basic algorithm used in our analysis is shown in Figure 2.1, and can be summarized as follows:

- a) We first check for the presence of snow in a particular pixel using MODIS data. If snow was present at any time during the 16-day retrieval (snow flag >0) then the difference between actual (16-day) surface albedo and snow-free albedo is taken as the albedo contrast.
- b) If MODIS data were missing then presence of snow is checked using NISE data. If snow is present during some or all of the 16-day period, then albedo contrast is determined by taking the difference between climatological snow-covered albedo for that gridcell (and time of year, as described in the previous subsection) and snow-free albedo. Since the NISE dataset is continuous and daily-resolved, we are able to determine those situations when a particular pixel was covered with snow for only a portion of the 16-day period of MODIS measurements. In those cases, if a pixel is covered with snow for D days out of 16, then the albedo contrast is multiplied with a scaling factor of $D/16$.
- c) If MODIS determines a pixel is not snow-covered (snow flag = 0) or in the absence of MODIS data if NISE does not indicate any snow ($D = 0$), then albedo contrast and LCrRE are set to zero for that pixel.
- d) After determining the albedo contrast, this term is multiplied with the various radiative kernels to derive different estimates of all-sky and clear-sky LCrRE. Except as noted in sensitivity studies described later, subsequent LCrRE results are derived from the CAM4 radiative kernel,

which simulates an intermediate level of cloud masking compared with the other kernels.

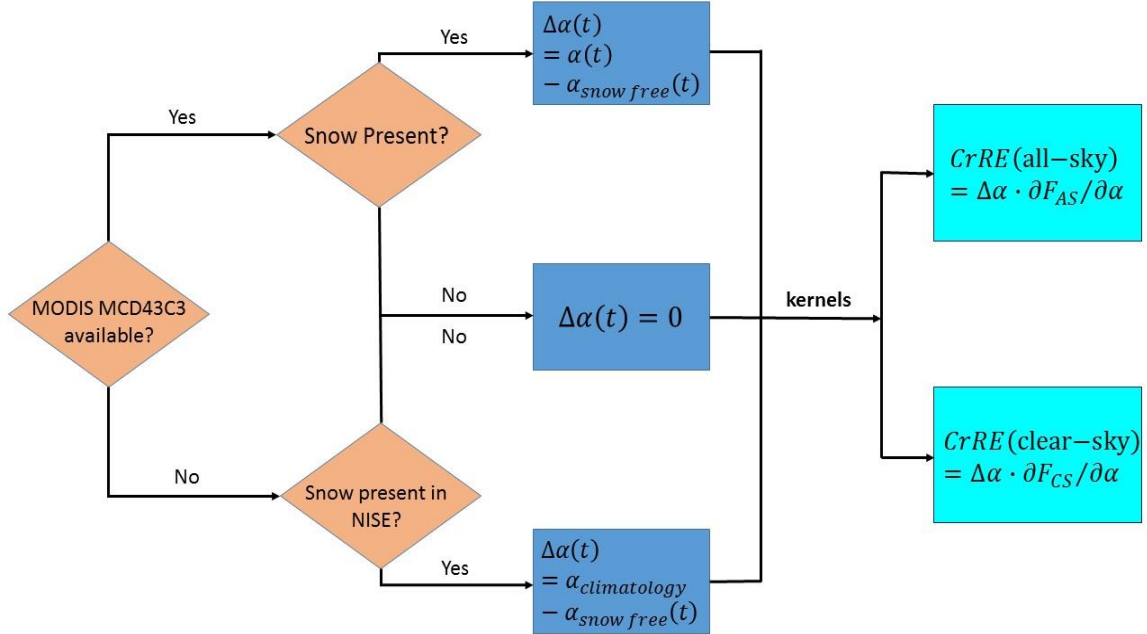


Figure 2.1: Algorithm for calculating CrRE

2.3 Results and Discussion

2.3.1 Spatial and Seasonal Variability of the Mean Climate State

All-sky global annual-mean LCrRE during 2001-2013 calculated using the CAM4 kernel is -2.58 Wm^{-2} , and ranges from -2.16 to -2.96 Wm^{-2} with application of different radiative kernels. Table 2.1a shows LCrRE averages over different domains and calculated with different radiative kernels. Permanent glaciated regions (e.g., Greenland and Antarctica) contribute about two-thirds of the net global LCrRE, due to their persistently high albedo around the year. For this discussion, land classified as snow or ice in the MODIS MCD12C1 land type dataset (Appendix A) is considered as permanently glaciated.

Due to the size of the Antarctic Ice Sheet, the Southern Hemisphere contributes about 60% of the global LCrRE. On the other hand, non-glaciated regions in Northern Hemisphere contribute about 32% of the global LCrRE. Because the areal extent of seasonal snow has little "memory" beyond a year, non-glaciated component of the LCrRE can respond rapidly to climate change and drive albedo feedback on sub-decadal timescales. Impact of non-glaciated region in Southern Hemisphere is negligible because of less land presence at mid and high latitudes (Figure 2.2), contributing only about 0.08% of the global LCrRE. In the Northern Hemisphere, non-glaciated regions contribute about 3.7 times more to LCrRE than glaciated areas. LCrRE associated with permanent glacier cover in mountain regions (e.g., the Himalaya) is also clearly visible, even though they are situated at much lower latitudes (Figure 2.2). Table 2.1 (b) and 2.1 (c) shows percentage contribution of different domains and separate land masses to the global LCrRE respectively.

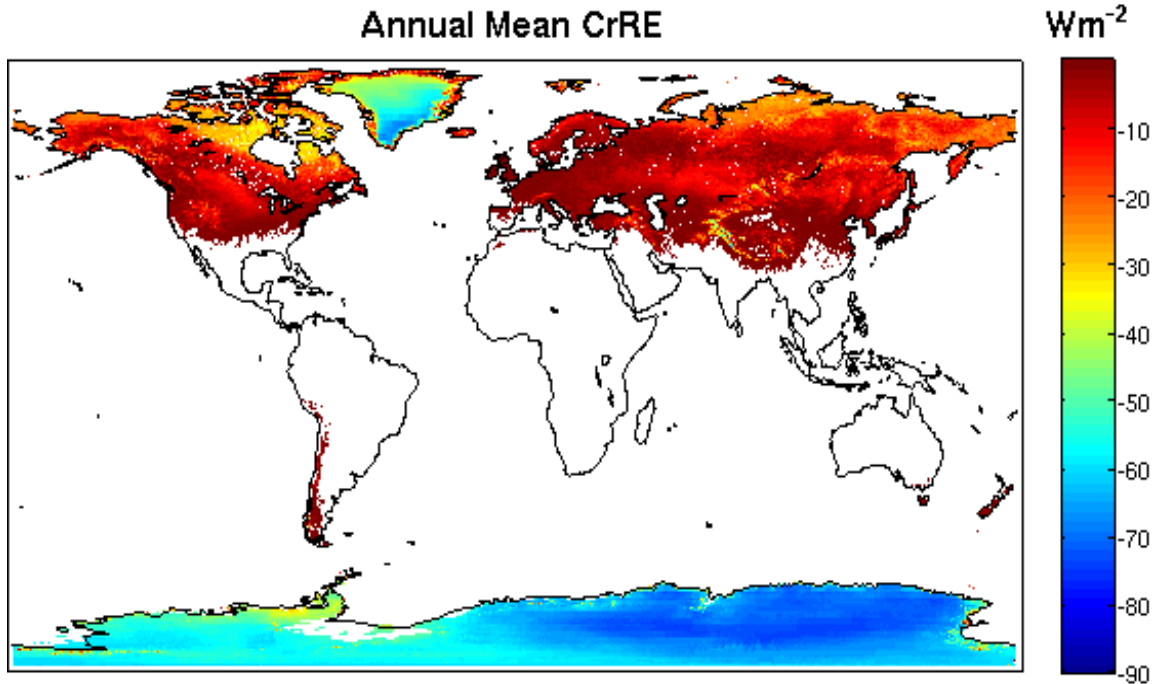


Figure 2.2: Map showing annual-mean all-sky land CrRE, averaged over 2001-2013, derived with the CAM4 kernel.

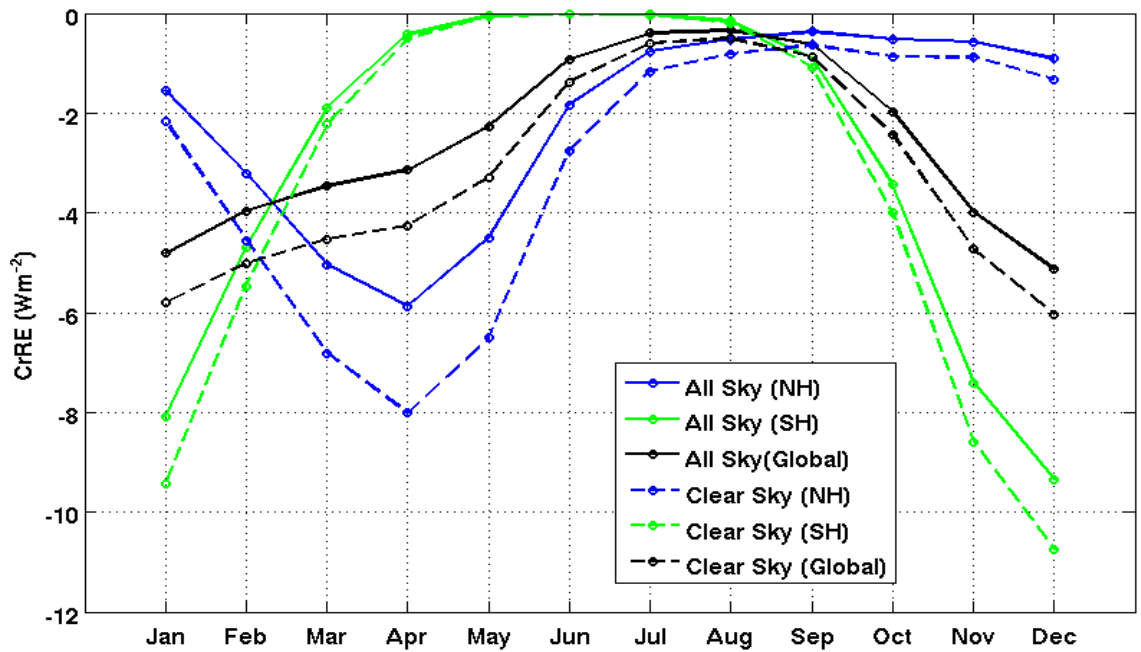


Figure 2.3: Annual cycle of hemispheric and global LCrRE for both all-sky and clear-sky conditions, derived from the CAM4 radiative kernels.

Table 2.1: (a) All-sky (clear-sky) cryosphere radiative effect averaged over different domains and derived with various radiative kernels (Wm^{-2}). (b) Percentage contribution of different domains to the global LCrRE using CAM4 kernels. (c) Percentage contribution of different land masses to the global LCrRE using CAM4 kernels.

(a)		CAM4 kernel	CAM5 kernel	CAM3 kernel	AM2 kernel
Northern Hemisphere	Glaciated	-0.45 (-0.62)	-0.56 (-0.63)	-0.36 (-0.61)	-0.45 (-0.60)
	Non-Glaci-ated	-1.67 (-2.4)	-1.95 (-2.48)	-1.16 (-2.42)	-1.40 (-2.24)
Southern Hemisphere	Glaciated	-3.08 (-3.58)	-3.41 (-3.61)	-2.79 (-3.58)	-3.16 (-3.59)
	Non-Glaci-ated	-0.004 (-0.007)	-0.005 (-0.008)	-0.003 (-0.007)	-0.004 (-0.007)
Global	Glaciated	-1.77 (-2.1)	-1.99 (-2.12)	-1.58 (-2.09)	-1.81 (-2.09)
	Non-Glaci-ated	-0.84 (-1.2)	-0.98 (-1.25)	-0.58 (-1.21)	-0.70 (-1.12)
Global	Global	-2.58 (-3.28)	-2.96 (-3.37)	-2.16 (-3.31)	-2.51 (-3.22)
(b)		% Contribution			
Northern Hemisphere	Glaciated	8.7 (9.5)			
	Non-Glaci-ated	32.4 (36.6)			
Southern Hemisphere	Glaciated	59.7 (54.6)			
	Non-Glaci-ated	0.08 (0.1)			
Global	Glaciated	68 (64)			
	Non-Glaci-ated	32 (36)			
Global	Global	100 (100)			
(c) Land Mass		LCrRE	% Contribution		
Antarctica	-1.51 (-1.76)	58.6 (53.6)			
Europe+Asia	-0.55 (-0.79)	21.1 (23.94)			
North America	-0.34 (-0.49)	13.02 (14.83)			
Greenland	-0.19 (-0.25)	7.2 (7.5)			
South America	-0.0024 (-0.0043)	0.09 (0.13)			
Australia	-2.28E-04 (-4.37E-04)	~0			
Africa	-5.59E-05 (-6.2E-05)	~0			
Global	-2.58 (-3.28)	100 (100)			

Clear-sky LCrRE values are expectedly higher than all-sky LCrRE values because of the absence of cloud scattering. Similarly, all-sky LCrRE derived from the CAM5 kernel is higher than that derived from the CAM4 kernel because cloud masking is substantially less in CAM5, due to the prevalence of thinner clouds (*Kay et al., 2012, Perket et al., 2014*). AM2 and CAM4 values are similar, indicating a similar degree of cloud masking in these two kernels, whereas the CAM3 kernel provides substantially greater attenuation of surface albedo anomalies at the TOA (*Flanner et al., 2011*). *Cao et al., (2015)* determined that the CAM3 and AM2 radiative kernels likely mask too much of the Arctic sea-ice radiative effect, but it is unclear if these kernel biases exist over land and have persisted in the more modern CAM4 and CAM5 models.

Figure 2.3 shows global and hemispheric monthly variations of LCrRE for both all-sky and clear-sky conditions. Seasonal variations are apparent in both hemispheres. LCrRE values peak (become most negative) during April in the Northern Hemisphere, about 2 months before the peak boreal solstice insolation. On the other hand, Southern Hemisphere LCrRE peaks in December, during maximum austral insolation. These differences occur because the Northern Hemisphere LCrRE is dominated by seasonal snow, while Southern Hemisphere LCrRE is dominated by permanent glaciated Antarctica, and thus the seasonal cycle of LCrRE is determined more directly by insolation in the Southern Hemisphere. Spatial distributions of tri-monthly seasonal averages of LCrRE are shown in Figure 2.4, indicating the same insolation-modulated tendencies over glaciated terrain described earlier. The timing of peak LCrRE in regions with seasonal snow, however, depends on the timing of melt onset, which almost always precedes the peak insolation period.

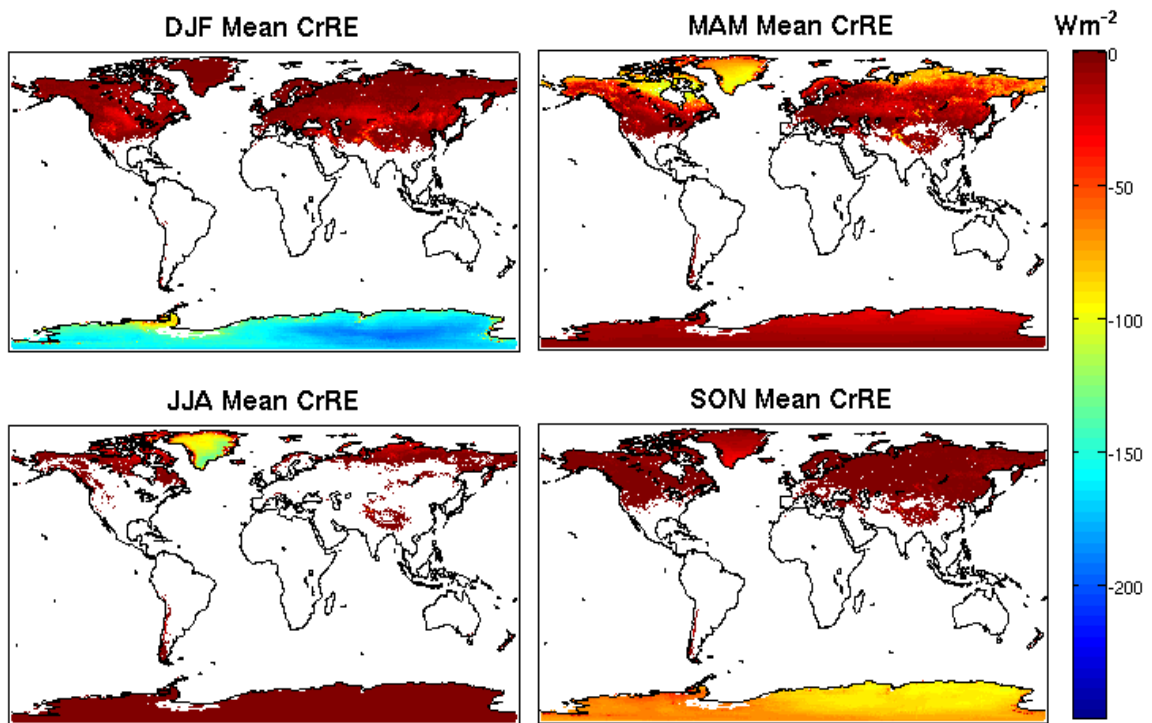


Figure 2.4: Tri-monthly seasonally averaged LCrRE derived with the CAM4 all-sky kernel. (DJF – December, January, February; MAM – March, April, May; JJA – June, July, August; SON – September, October, November).

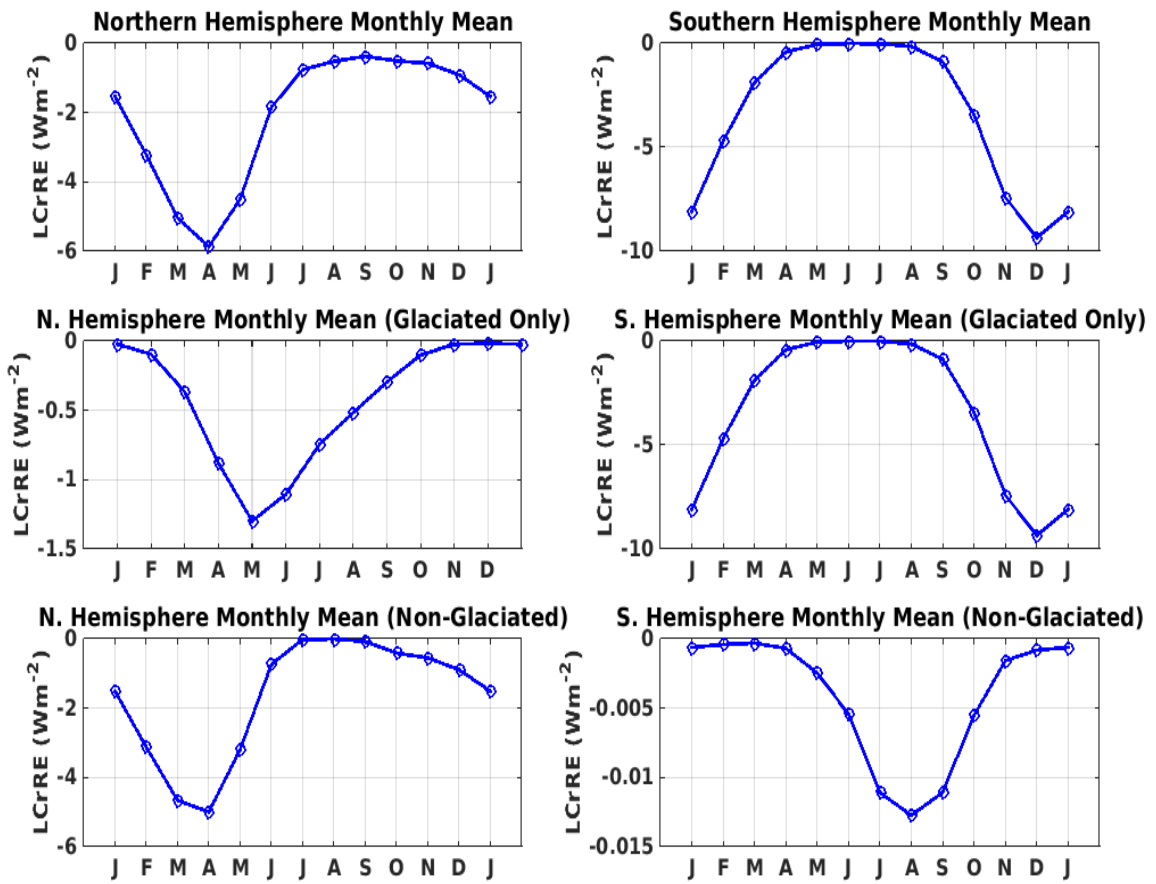


Figure 2.5: All-sky LCrRE averaged over the Northern (left) and Southern (right) Hemisphere, shown as contributions from all land within the hemisphere (top), permanently glaciated areas only (middle), and non-glaciated areas only (bottom). Data were derived with the CAM4 radiative kernel.

Peak LCrRE over glaciated regions of the Northern Hemisphere occurs in May, lagging the peak over non-glaciated regions by about one month (Figure 2.5). Glaciated LCrRE in the Northern Hemisphere peaks before the summer solstice because Greenland surface albedo decreases between May and the end of June as summer melt commences. This tendency is not apparent over Antarctica because a smaller portion of the Antarctic Ice Sheet experiences surface melt during summer. Also the LCrRE change is more gradual for glaciated than non-glaciated regions during both melting and freezing seasons. As the temperature rises, temporary snow over non-glaciated land ablates relatively quickly. This causes a rapid decrease in surface albedo and hence lowering of the albedo contrast. Similar observations cannot be made for Southern Hemisphere due to negligible LCrRE contribution from non-glaciated region. Global LCrRE values show less seasonal variation than the hemispheric averages because they are basically averages of two hemispheric seasonal cycles that are somewhat out of phase with each other. Nonetheless, global LCrRE does exhibit a minimum during July and August, when there is little seasonal snow cover in the Northern Hemisphere and Antarctica receives little insolation.

We observe a contribution of about -0.037 Wm^{-2} (1.43% of the global total) from high mountainous regions (Himalayas, Andes, etc.) towards global LCrRE. The Himalayan region alone contributes about 93.8% of the total mountainous region LCrRE. (Here, we define the high mountainous contribution to LCrRE as that coming from permanent glaciated areas between latitudes of 60°S and 60°N .) The contribution from these high-altitude areas is relatively smaller but consistent throughout the year. Our use of 0.05 degree resolution data allows us to determine LCrRE over many regions with patchy snow extent, though mountain snow cover varies substantially on even smaller spatial scales.

Figure 2.6 shows LCrRE contribution of the Andes and the Himalayas averaged over the MODIS era.

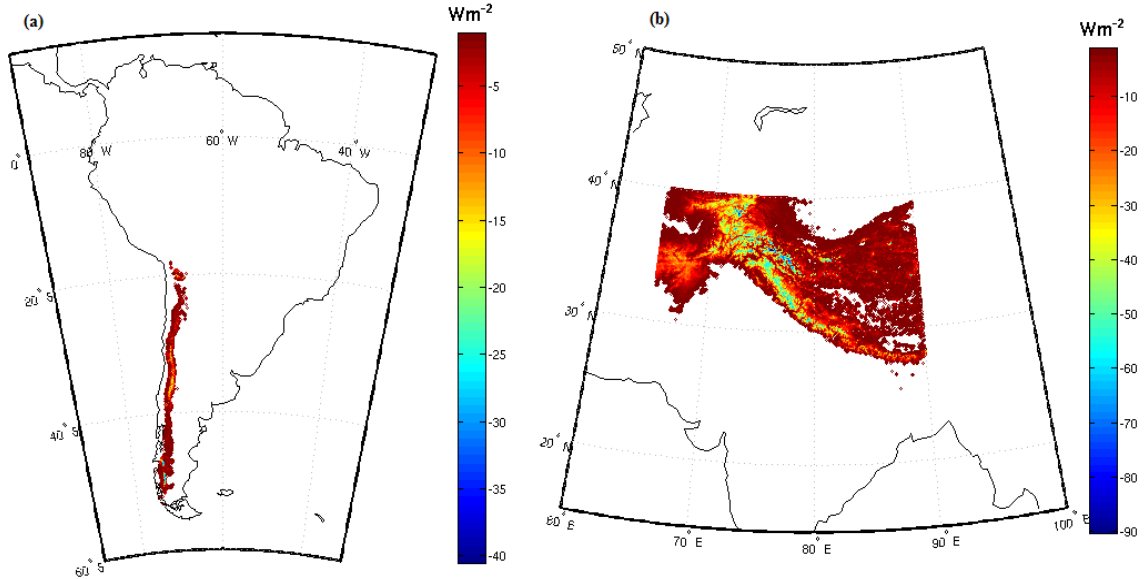


Figure 2.6: Maps showing all-sky LCrRE, averaged over 2001-2013 over (a) Andes and (b) Himalayan mountain ranges.

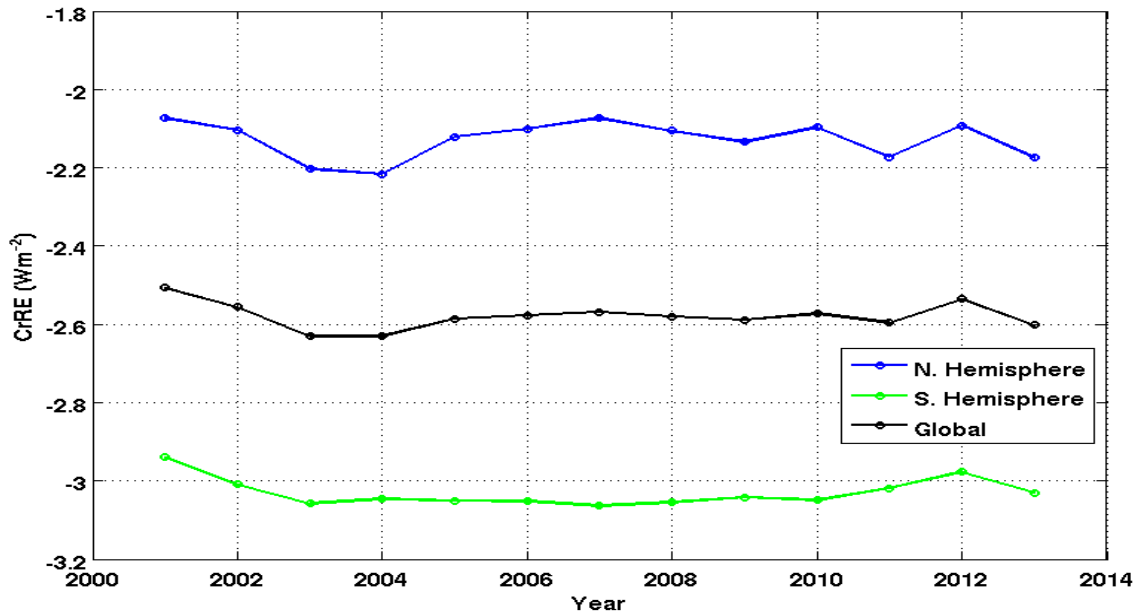


Figure 2.7: Annual mean land LCrRE time series during 2001-2013 for the globe and each hemisphere.

2.3.2 *Sensitivity Analysis*

A sensitivity analysis of LCrRE has been done to estimate its dependence on various parameters included in the algorithm (Table 2.2). In this part of the study, CAM4 kernels are used and other parameters are varied. The first column represents the LCrRE values from the default case using the algorithm discussed in section 2.3 (Figure 2.1). The second column lists LCrRE values generated using only the NISE snow flag (i.e., snow presence is determined exclusively with NISE measurements, rather than MODIS). Differences between these estimates are due to the different remote sensing techniques applied to determine snow presence and different spatial and temporal resolutions of the NISE and MODIS datasets.

The third and fourth columns are similar to first and second columns, respectively, but apply the land class climatology of albedo contrast instead of spatially- and temporally-varying albedo contrast (section 2.2.2). The derivation of these land class climatology values is described in step 3 of section 2.2.2. A list of mean albedo values and standard deviations for the different land classes is provided in Appendix A. In these cases land class climatology albedo values have been used irrespective of data availability in MODIS, in combination with snow presence determined from MODIS+NISE (column 3) and exclusively from NISE (column 4). This sensitivity study is designed to assess the utility of using much simpler estimates of snow-covered albedo than our more involved space- and time-dependent estimates.

LCrRE values in the second, third and fourth columns (Table 2.2) are very similar to each other. In all three cases the estimated LCrRE is lower than the original analysis

(column 1). This indicates the dependence of LCrRE on MODIS snow flag and global albedo climatology, and changing either of those drops the LCrRE estimates by about 8% globally. Also when land class climatological albedos are used, LCrRE is very similar for different choice of snow flags. This may be due to the fact that the land class climatology only depends on land class type and does not vary with time.

MODIS MCD43C3 albedo data are accompanied by quality flags, indicating the fraction of input measurements to each 16-day data point that were made under cloud-free conditions with sufficiently small solar zenith angle [e.g., *Schaaf et al*, 2002]. Table 2.3 shows a comparative study of LCrRE determined without quality flag filtering (i.e., quality flag 4 and better), and determined using only quality flag 2 or better albedo data. Quality flag 2 is mixed, with 75% or less of the underlying data derived from inverting reflectance for the BRDF retrievals, and 25% or less of the underlying data filled. Using better albedo data (lower quality flag) does not make a significant difference in our determination of global LCrRE. This indicates that the fill values applied in the MCD43C3 retrieval algorithm are similar to our developed climatology (section 2.2.2).

2.3.3 *Inter-annual trends*

Annual global LCrRE averages do not show significant inter-annual trends during the MODIS era (Figure 2.7), perhaps partially due to the relatively short duration of this period. Slopes for global, northern hemisphere and southern hemisphere LCrRE are $-2 \times 10^{-3} \text{ Wm}^{-2}\text{yr}^{-1}$, $-1.3 \times 10^{-3} \text{ Wm}^{-2}\text{yr}^{-1}$ and $7.3 \times 10^{-4} \text{ Wm}^{-2}\text{yr}^{-1}$, respectively. These trends are not significant at the 95% confidence level.

To better understand the LCrRE trend, pixel-by-pixel trend maps have been generated using the Mann-Kendall regression technique with 95% significance level (Figure 2.8). A strong positive trend (indicating less negative LCrRE) is observed over some regions of the Northern Hemisphere, especially over Northern Asia during the spring season (Figure 2.8b). Positive trends over seasonally snow covered areas are most likely due to a decrease in snow coverage [e.g., *Qu and Hall, 2007; Vaughan et al., 2013*], and the trend in LCrRE that we find over Northern Asia is consistent with studies showing recent declines in snow cover over this region [e.g., *Dery and Brown, 2007; Brown and Robinson, 2011; Derksen et al., 2014*]. Possible reasons for the positive trends seen over Greenland include: 1) increased snow metamorphism and black carbon deposition [e.g., *Box et al, 2012; Keegan et al., 2014*], 2) transport and deposition of dust and other light-absorbing impurities over the ice-sheets due to increased dust source areas associated with increased snow-free area [*Dumont et al., 2014*], 3) higher melt extent across the Greenland ice sheet expose more fresh water at the surface (*Tedesco et al., 2014*) and 4) MODIS sensor degradation on the Terra satellite during recent years [*Sun et al., 2014; Lyapustin et al., 2014*], which would indicate a spurious decline in albedo.

We also observe slightly negative LCrRE trends over Antarctica, which may be due to increases in snowfall that have decreased the duration that surface snow has to “age”, thereby increasing surface albedo [*Picard et al., 2012*]. Interestingly, LCrRE has also become more negative over some regions of Central Asia, presumably due to increased temporal coverage of snow.

Table 2.2: All-sky (Clear-Sky) CrRE (Wm-2) derived with different algorithms. All cases apply the CAM4 radiative kernels.

		Both flags with original climatology	Only NISE flag with original climatology	Both flags with land class climatology	Only NISE flag with land class climatology
Northern Hemisphere	Glaciated	-0.45 (-0.62)	-0.37 (-0.51)	-0.39 (-0.54)	-0.38 (-0.51)
	Non-Glaci- ated	-1.67 (-2.4)	-1.4 (-2.0)	-1.41 (-2.05)	-1.47 (-2.1)
Southern Hemisphere	Glaciated	-3.08 (-3.58)	-3.03 (-3.51)	-2.91 (-3.38)	-2.91 (-3.38)
	Non-Glaci- ated	-0.004 (-0.007)	-0.004 (-0.006)	-0.005 (-0.008)	-0.008 (-0.012)
Global	Glaciated	-1.77 (-2.1)	-1.7 (-2.0)	-1.65 (-1.96)	-1.65 (-1.95)
	Non-Glaci- ated	-0.84 (-1.2)	-0.7 (-1.0)	-0.71 (-1.03)	-0.74 (-1.05)
Global	Global	-2.58 (-3.28)	-2.4 (-3.01)	-2.36 (-2.99)	-2.38 (-3.0)

Table 2.3: All-sky (Clear-Sky) CrRE (Wm-2) with different quality flag filters, derived using CAM4 kernels in all cases.

		Quality flag 4 or better	Quality flag 2 or better
Northern Hemisphere	Glaciated	-0.45 (-0.62)	-0.43 (-0.59)
	Non-Glaci- ated	-1.67 (-2.4)	-1.6 (-2.29)
Southern Hemisphere	Glaciated	-3.08 (-3.58)	-3.07 (-3.57)
	Non-Glaci- ated	-0.004 (-0.007)	-0.004 (-0.006)
Global	Glaciated	-1.77 (-2.1)	-1.75 (-2.08)
	Non-Glaci- ated	-0.84 (-1.2)	-0.8 (-1.15)
Global	Global	-2.58 (-3.28)	-2.55 (-3.23)

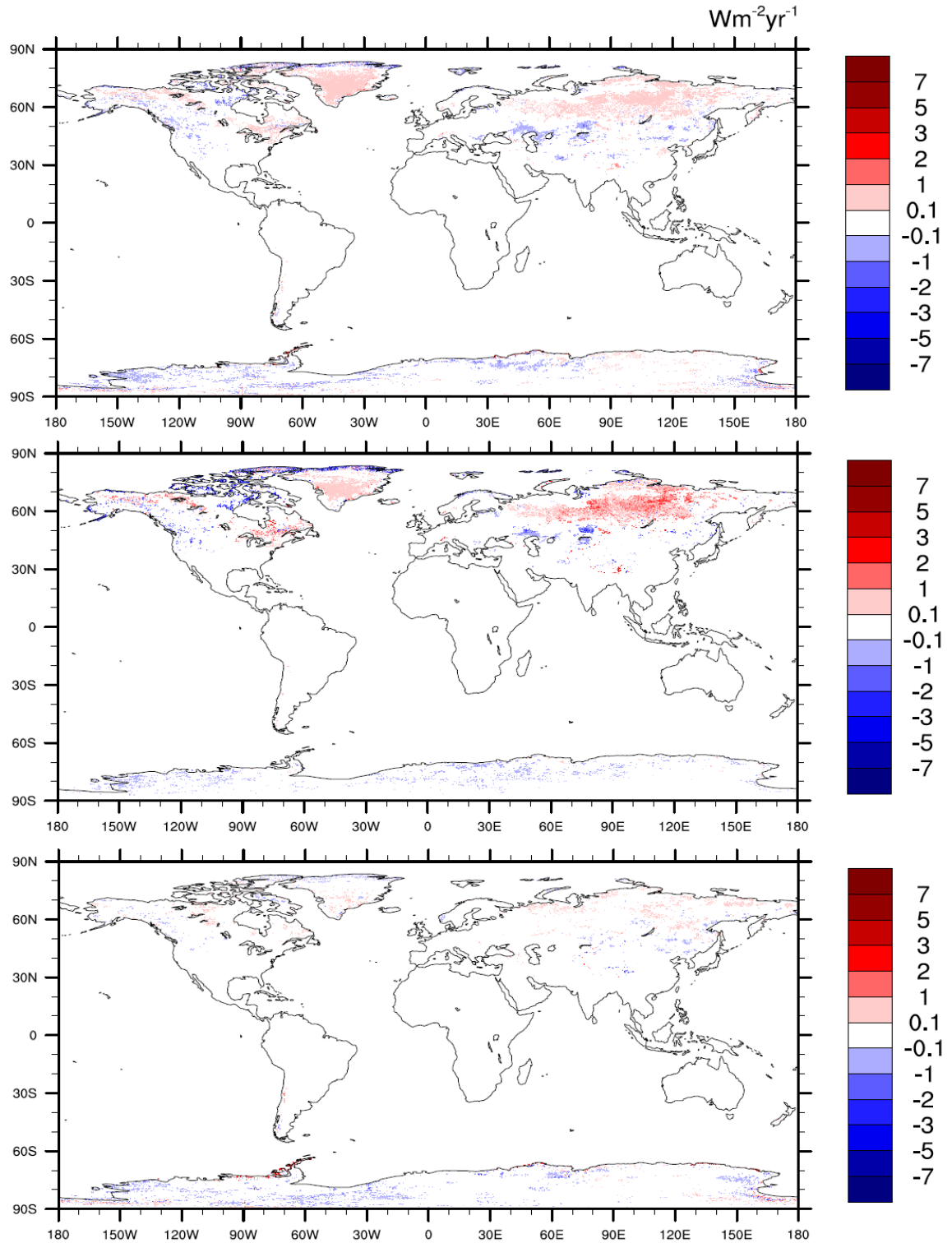


Figure 2.8: LCrRE trend maps for annual (Top), MAM (Middle) and SON (Bottom) seasons (MAM – March, April, May; SON – September, October, November). Trends were derived using Mann-Kendall regression technique. Only regions with significant trends ($p=0.05$) are shown in the map.

2.3.4 Comparison with previous study

The current study and that of *Flanner et al.*, [2011] have a common overlap period of 2001-2008. To compare the derivations of LCrRE between these two studies, we re-derived Northern Hemisphere LCrRE during 2001-2008 using the CAM3 radiative kernel and an assumption of ice-free albedo over Greenland of 0.316, as in *Flanner et al.* [2011]. For this overlap period, and using these assumptions, the current study gives a LCrRE value of about -1.41 Wm^{-2} over the northern hemisphere as compared to -1.72 Wm^{-2} from *Flanner et al.*, [2011], who determined snow presence during 1979-2008 from coarse-resolution AVHRR data [e.g., *Robinson et al.*, 2000]. (The 1979-2008 mean northern hemisphere LCrRE found by *Flanner et al.*, [2011] was -2.0 Wm^{-2} , indicating less boreal cryospheric influence during 2001-2008 compared with 1979-2001). To understand this difference, an LCrRE difference map between the current study and *Flanner et al.*, (2011) has been generated (Figure 2.9). This difference was taken between annual mean LCrRE values over the common overlap period. We find that the use of coarse-resolution snow cover data [*Flanner et al.*, 2011] likely leads to overestimation of LCrRE along the continental edges and in mountainous regions like the Himalayas. Some of the differences along the continental edges (especially of Greenland), depicted in Figure 2.9 at 1 degree resolution, are likely due to poor resolution of land cover fraction in the AVHRR snow cover data used by *Flanner et al.* [2011], and possibly also the influence of land-fast sea-ice. Snow in mountainous regions is difficult to resolve in any model because of irregular topography, and is uncertain in both datasets. The annual mean LCrRE time-series (Figure 2.10) of both datasets show similar patterns and exhibit a correlation coefficient of 0.78, lending confidence in the application of longer-term AVHRR data to derive estimates of large-

scale LCrRE trends. We expect that the current analysis will be useful in determining correction algorithms for biases in the Himalayas and along the coastal margins of Greenland that may result from use of coarse-resolution AVHRR data.

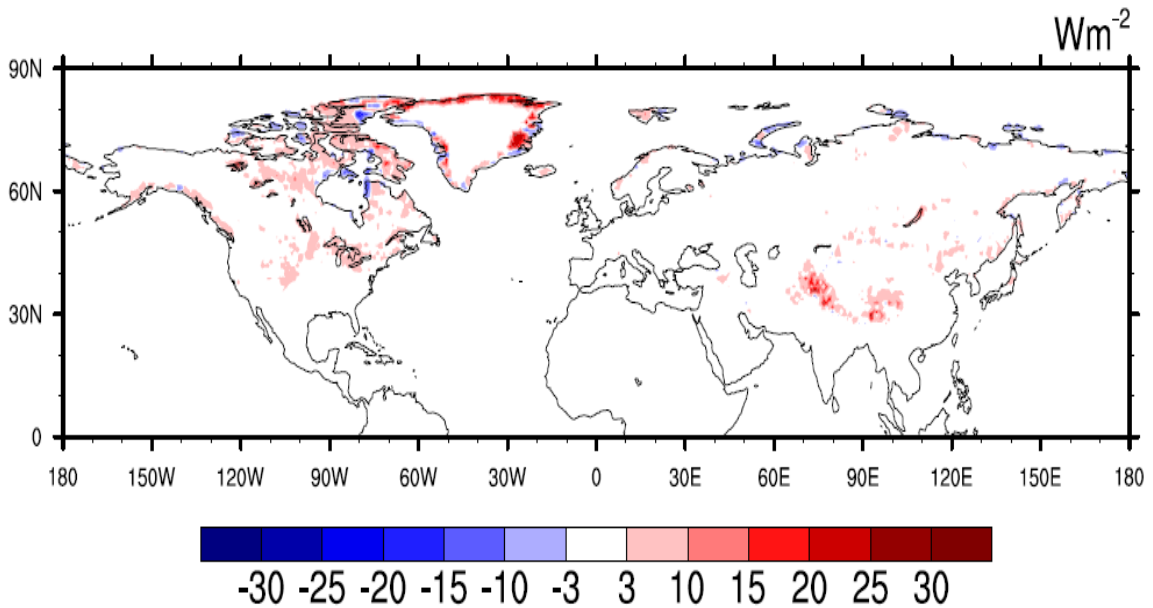


Figure 2.9: Difference of 2001-2008 mean LCrRE between the current study and Flanner et al., (2011) using CAM3 kernels. Areas of red indicate a stronger (more negative) LCrRE determined by Flanner et al [2011] than determined here.

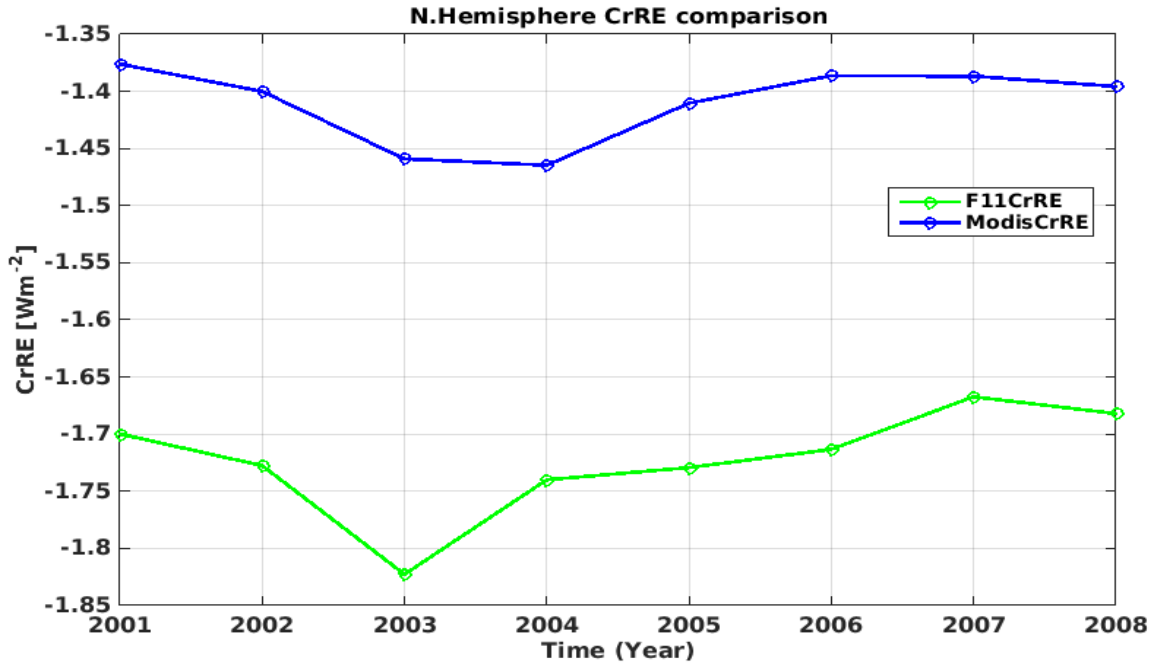


Figure 2.10: Annual mean Northern Hemisphere LCrRE timseries derived for the current study (ModisCrRE) and by Flanner et al., [2011] (F11CrRE) using CAM3 kernels.

2.4 Conclusions

We have estimated a global land-based CrRE (LCrRE) of about -2.6 Wm^{-2} during 2001-2013, with about 59% of the effect originating from Antarctica. For a particular region, LCrRE can vary significantly (e.g., from 0 to -250 Wm^{-2} over Antarctica; Fig. 5) depending on the season. Due to a large contribution from seasonal snow, northern hemisphere LCrRE peaks during April, whereas southern hemisphere LCrRE peaks during the December solstice. About 78% of Northern Hemisphere LCrRE originates from seasonal or non-glaciated snow while about 99.9% of Southern Hemisphere LCrRE comes from glaciated areas. Mountainous region provide a small (-0.037 Wm^{-2}) yet consistent contribution towards LCrRE throughout the year. Higher resolution data has improved the LCrRE estimates relative to previous studies by better resolving the snow impact in mountainous regions and in coastal areas, leading to a 18% smaller estimate of northern

hemisphere LCrRE during 2001-2008 in this study than in *Flanner et al.* (2011). Inter-annual variations between these two studies are well correlated ($r=0.78$), indicating that these data are useful in determining longer term trends in LCrRE.

Snow cover, however, is patchy on substantially smaller spatial scales than the 0.05 degree resolution analysis applied here. No significant trend over time has been observed in global annual LCrRE values between 2001 and 2013, though significant reductions in cryospheric influence are evident over large areas of northern Eurasia, especially during spring. Slightly negative LCrRE trends (indicating increased albedo) are evident over Antarctica, possibly due to increased snowfall frequency. LCrRE primarily depends on albedo contrast induced by snow and the propagation of surface albedo anomalies to top-of-atmosphere energy fluxes. A slight variation in either of the two can cause a major impact on LCrRE estimates. Sensitivity analysis shows a secondary dependency of LCrRE on MODIS snow flag and global climatology derived using MODIS surface albedo product, and changing either of those drops the LCrRE estimates by about 8% globally. Using different MODIS data quality flags (Q2 & Q4) does not have a significant effect on our LCrRE estimates. Using different radiative kernels can cause a variation of about 3-16% in global LCrRE as compared to LCrRE with the CAM4 kernel, depending on the atmospheric attenuation present in each particular model due to cloud cover and aerosols.

For the consistency, we tried not use any partial year data available at the time of analysis. Also, adding another year of the data (year 2014) would not have made any significant impact on the results. However, it would be beneficial to look into long term trends (about 20-30 years) and comparison with previous research, with more MODIS data available in future.

Appendix A: Snow-covered albedo climatology by land classification

Definitions for different land classes used in the MODIS MCD12C1 product are provided by the International Geosphere-Biosphere Programme (IGBP). Table A.1 provide mean (μ) and standard deviation (σ) of surface albedo for each land class type using section 2.2 algorithm.

Table A.1: Mean (μ) and standard deviation (σ) of snow-covered surface albedo for different land class types. The averaging filter used to determine snow presence was based on both MODIS and NISE datasets of snow, and thus these averages are specific to the current study. MODIS provides surface albedo only over land, so albedo for water cannot be determined.

Land Class Type	Mean (μ)	Std Dev (σ)
Water	N/A	N/A
Evergreen Needleleaf forest	0.30	0.008
Evergreen Broadleaf forest	0.21	0.005
Deciduous Needleleaf forest	0.33	0.024
Deciduous Broadleaf forest	0.33	0.018
Mixed forest	0.29	0.013
Closed shrublands	0.42	0.023
Open shrublands	0.54	0.046
Woody savannas	0.40	0.025
Savannas	0.46	0.031
Grasslands	0.48	0.038
Permanent wetlands	0.41	0.027
Croplands	0.49	0.036
Urban and built-up	0.37	0.019
Cropland/Natural vegetation mosaic	0.41	0.026
Snow and ice	0.72	0.033
Barren or sparsely vegetated	0.37	0.023

Chapter 3

An improved carbon dioxide snow spectral albedo model:

Application to Martian conditions

3.1 Introduction

The Martian atmosphere consists primarily of carbon dioxide, and Martian polar caps are covered primarily with carbon dioxide ice [*Leighton and Murray, 1966; Herr and Pimentel, 1969; Larson and Fink, 1972; Forget, 1998; Bibring et al., 2005*]. H₂O ice is also present at the surface of the perennial ice caps, with small amounts of seasonal deposition in other parts of the planet [e.g., *Kieffer et al., 2000; Bibring et al., 2004; Brown et al., 2014*]. Significant portions of atmospheric CO₂ (25-30%) deposit seasonally in each hemisphere, as indicated by model simulations and surface pressure measurements [*Tillman et al., 1993; Forget et al., 1998; Kieffer and Titus, 2001*]. To understand the impact of these ices on the planet's cryosphere albedo, it is important to accurately determine the spectral dependence of CO₂ snow albedo and influences of properties such as dust content, ice grain size, and snow thickness, as well as the albedo effects of mixing and layering of CO₂ and H₂O snow.

Our work determines the albedo of CO₂ snow by extending the Earth-based Snow, Ice, and Aerosol Radiation (SNICAR) model [*Flanner et al., 2007; Flanner et al., 2009*], originally designed for H₂O snow. SNICAR utilizes the multiple scattering, multi-layer

two-stream radiative approximation described by *Toon et al.* [1989], with the delta-hemispheric mean approximation. We extend the current version of SNICAR from 470 bands (over the wavelength range 0.3-5.0 μm) to 480 bands spanning 0.2-5.0 μm at 10 nm spectral resolution. We include these extra 10 bands in the ultraviolet (UV) spectrum because of the lack of ozone in the Martian atmosphere compared to Earth [*Montmessin and Lefèvre, 2013*], meaning more UV radiation reaches the Martian surface and interacts with snow. A single-layer implementation of SNICAR can be operated interactively on the web at: <http://snow.engin.umich.edu>.

Very few studies have focused on modeling of Martian CO₂ snow albedo across the UV, visible and near-IR spectrum [*Warren et al., 1990; Hansen, 1999; Bonev et al., 2008*]. *Langevin et al., [2007]* and *Appéré et al., [2011]* present modeled near-IR albedo of Martian cryospheric surfaces, as discussed in Section 4. With the availability of more accurate and spectrally-resolved laboratory measurements of CO₂ ice complex refractive indices across the solar spectrum [*Hansen, 1997; Hansen 2005*], we provide improved and updated spectral albedos of carbon dioxide snow with applicability to Martian conditions. The presence of light-absorbing impurities generally lowers snow albedo. We simulate the impact of Martian dust [*Wolff et al., 2006; Wolff et al., 2009; Wolff et al., 2010*] and palagonite [*Clark et al., 1990; Clancy et al., 1995*] on surface cryosphere albedo. Palagonite is a volcanic rock and serves as a terrestrial analog for Martian dust [*Banin et al., 1997*]. We perform multiple analyses to determine the sensitivity of cryosphere spectral albedo to the amount and type of dust, presence of both ices, ice grain size, snow layer thickness, and solar zenith angle. We also compare the results of our simulations with observed Mars surface albedo derived from Compact Reconnaissance Imaging Spectrometer (CRISM)

measurements, and the OMEGA instrument [Appéré *et al.*, 2011] (section 4). Apparent albedo, derived by *Brown et al* [2014] using CRISM measurements, were provided by Adrian Brown (SETI Institute, personal communication), and observations and modeled albedo from the OMEGA instrument were provided by Thomas Appéré (IPAG, Grenoble, personal communication).

3.2 Data and methodology

Hansen [1997, 2005] made extensive laboratory measurements of the complex refractive indices of solid CO₂ ice in the spectral range from 0.174 μm to 333 μm . We apply these data along with Mie calculations to derive optical properties of different lognormal size distributions of CO₂ ice particles, reported with effective radius (r_{eff}), or the surface area-weighted mean radius of the size distribution. Our simulations of H₂O snow albedo utilize refractive indices of H₂O ice provided by *Warren and Brandt* [2008]. We use “central hematite” dust mineral fractions from *Balkanski et al.* [2007] as a proxy of typical Earth dust. These mineral abundances are representative of aeolian dust from the Saharan desert. Refractive indices for this dust mixture are derived using the Maxwell-Garnett mixing approximation, following e.g., *Sokolik and Toon* [1999], along with measurements of mineral refractive indices from various sources. Refractive indices of Martian dust (hereafter Mars dust) were provided by Mike Wolff (Space Science Institute, personal communication), and were derived using data from instruments onboard the Mars Reconnaissance Orbiter (MRO), as described by *Wolff et al.* [2009], and *Wolff et al.* [2010]. We combine refractive indices of palagonite from *Clark et al.* [1990] over the 0.2 μm to 0.6 μm spectral range and measurements from 0.6 μm to 5 μm made by *Clancy et al.* [1995]

to derive data over the solar spectrum. In this study, we divide our broadband (0.2-5.0 μm) into two sub-regions: Visible (0.2-0.7 μm) and near-IR (0.7-5.0 μm).

We determine the spectrally-varying single scatter albedo (ω_0), scattering asymmetry parameter (g), and mass extinction cross-section for all three dust types using Mie Theory with an assumed gamma size distribution [*Hansen and Travis, 1974*] with $r_{\text{eff}}=1.5\mu\text{m}$ and effective variance (v_{eff})=0.3 [*Wolff et al., 2006*]. Figure 3.1 shows the optical properties for these dust types. Mass absorption cross-section (Figure 3.1d) is the product of single scatter co-albedo (Figure 3.1a) and mass extinction cross-section (Figure 3.1b). We assume the same dust density of 2000 kg/m^3 for both palagonite and Mars dust [*Stroncik and Schmincke, 2002*], and a dust density of 2570 kg/m^3 for Earth dust, based on the volume fraction-weighted densities of the constituent minerals [*Balkanski et al., 2007*].

The bulk optical properties (extinction optical depth (τ), ω_0 , and g) for each snow layer are calculated from the abundances of each constituent [e.g., *Flanner et al, 2007*], with τ calculated as the sum of that associated with each constituent, ω_0 as a τ -weighted average of each constituent, and g as a scattering optical depth (product of τ and ω_0) weighted average of the asymmetry parameter of each constituent. These bulk layer properties are then applied in the two-stream multiple scattering approximation adopted by SNICAR. We utilize this extended framework of SNICAR to first simulate the albedo of pure CO_2 snow across the visible and near-IR spectrum. We then explore and compare the impacts of Earth dust, Mars dust and palagonite on CO_2 and H_2O snow albedo. Finally, we adopt a two layer model (with the bottom layer being semi-infinite) to calculate the impacts of H_2O snow presence on top of CO_2 snow and vice-versa. In all other cases we assume only a single snow layer, either with semi-infinite thickness or varying shallow thicknesses.

Here, a layer with thickness of 100 m is referred to as semi-infinite. Figure 3.2 shows simplified diagrams of the various model configurations applied in this study.

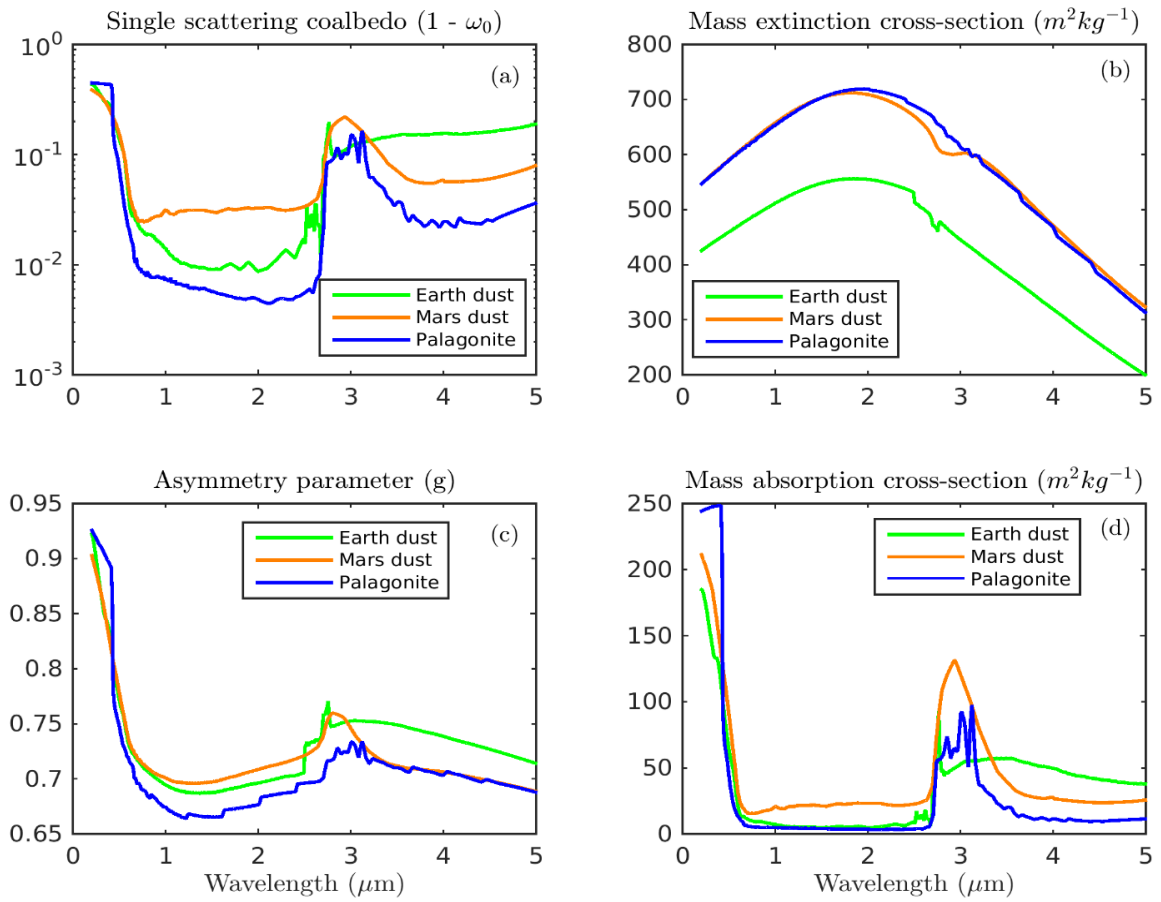


Figure 3.1: Spectral optical properties of the Earth dust, Martian dust, and palagonite applied in this study.

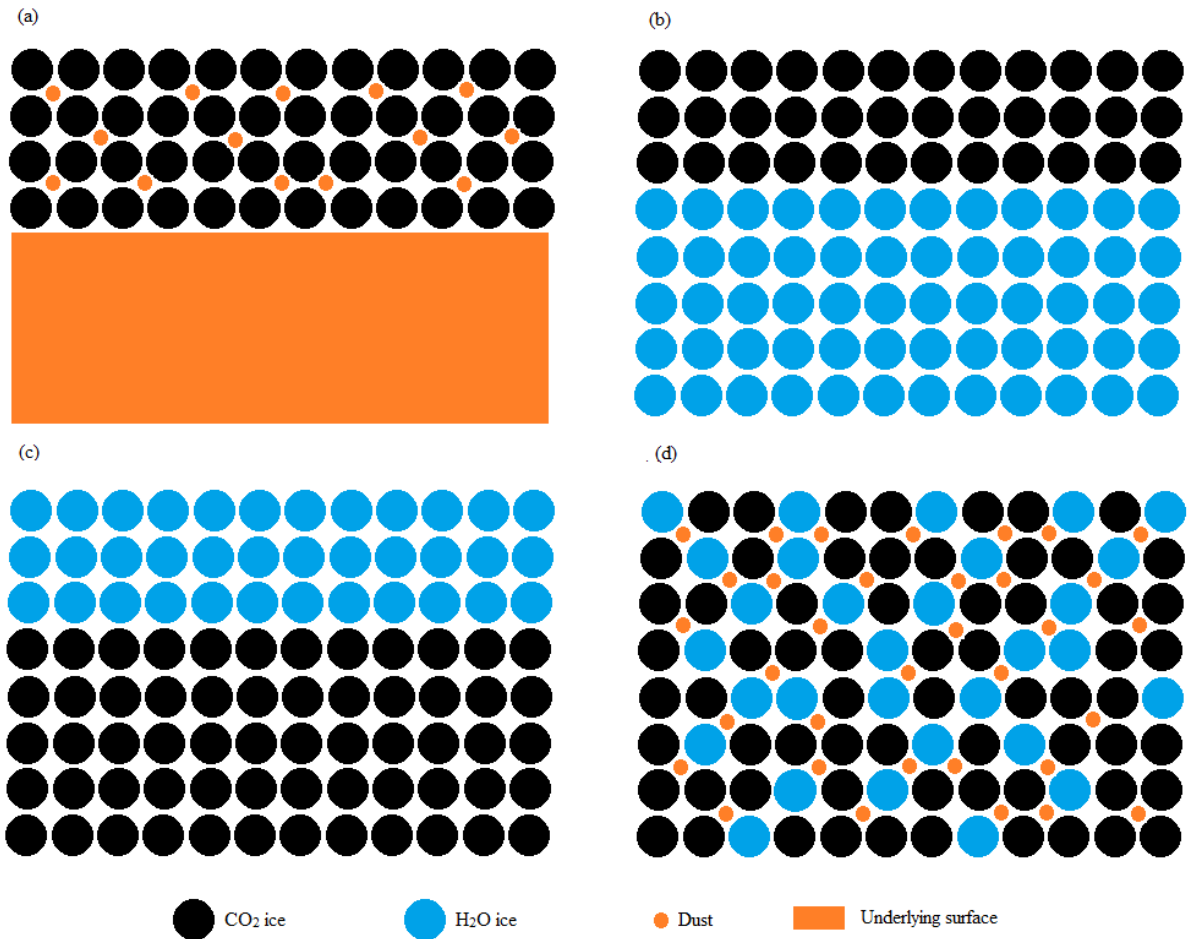


Figure 3.2: Schematic of the various model configurations applied in this study: (a) a single CO_2 or H_2O snow layer (with dust) on top of solid underlying surface; (b) CO_2 snow with variable thickness on top of a semi-infinite H_2O snow layer; (c) H_2O snow with variable thickness on top of a semi-infinite CO_2 snow layer; (d) mixed snow layers with dust.

3.3 Results

Figure 3.3 shows the spectral hemispheric albedo of pure, semi-infinite CO₂ and H₂O snow with spherical grain effective radius of 100 μm and solar zenith angle of 60°. It is evident that the CO₂ is more reflective than H₂O snow, especially in the near-IR spectrum. Table 3.1 compares solar broadband, visible and near-IR albedo for both types of snow, where band-averaged values are weighted with solar spectral irradiance measurements from *Labs and Neckel* [1968]. H₂O snow albedo is only about two-thirds of the CO₂ snow albedo in the near-IR spectrum, although this ratio is grain size dependent (section 3.3.2). With the presence of both CO₂ and H₂O ices on Martian polar caps, the contrasting reflectivity of these types of snow could significantly impact planetary shortwave energy fluxes, both at the surface and top of atmosphere.

Snow albedo depends on many physical quantities, including grain size, solar zenith angle, layer thickness, and type and amount of impurities. We performed a sensitivity analysis to understand the effect of these parameters on carbon dioxide snow albedo, and describe this analysis below. We assume an effective grain size of 100 μm and solar zenith angle of 60° for all analyses, unless stated otherwise. For comparing CO₂ and H₂O albedo, we also assume identical pore volume fractions for each type of ice. A typical bulk density of settled snow on Earth is 200 kg/m³ (e.g., *EN 1991-1-3*). Because the bulk densities of H₂O and CO₂ ice are assumed to be 917 kg/m³ and 1500 kg/m³, respectively, we therefore assume H₂O and CO₂ snow densities of 200 kg/m³ and 327.15 (= 200×1500/917) kg/m³, respectively.

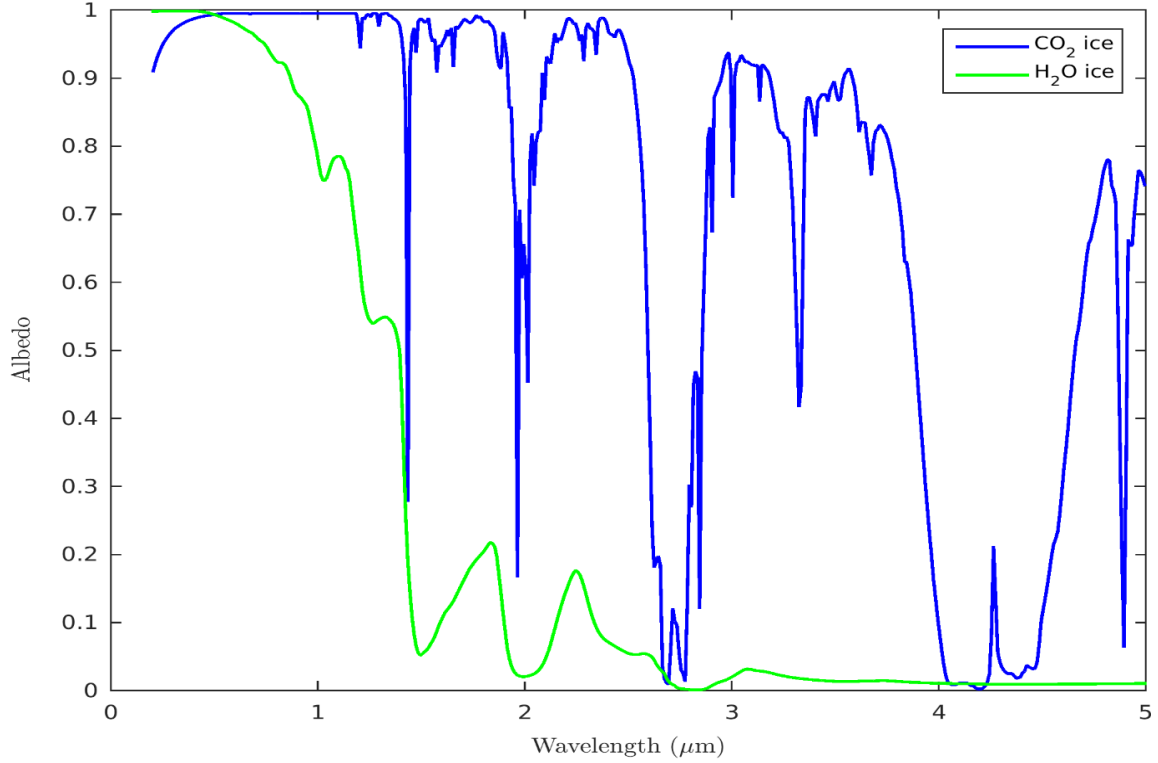


Figure 3.3: Comparison of CO_2 and H_2O snow spectral albedo simulated with the SNICAR model (Grain size = $100 \mu\text{m}$ and solar zenith angle = 60°)

Table 3.1: Albedo values for pure, semi-infinite CO_2 and H_2O snow in different spectral bands (Grain size = $100 \mu\text{m}$ and solar zenith angle = 60°)

	pure CO_2	pure H_2O
Visible (0.2-0.7 μm)	0.991	0.988
Near-IR (0.7-5.0 μm)	0.952	0.609
Broadband (0.2-5.0 μm)	0.970	0.787

3.3.1 Solar zenith angle

We see little effect of solar zenith angle on CO₂ snow albedo (Figure 3.4). The broadband albedo only changes from 0.98 to 0.96 over the 80° change in zenith angle. This compares with H₂O snow albedo, which varies from 0.75 to 0.82 for the same set of parameters. Variability in albedo with solar zenith angle is greater in more absorptive parts of the spectrum (i.e., the near-IR), but for CO₂ snow the variability is small in the parts of the spectrum containing most of the solar energy.

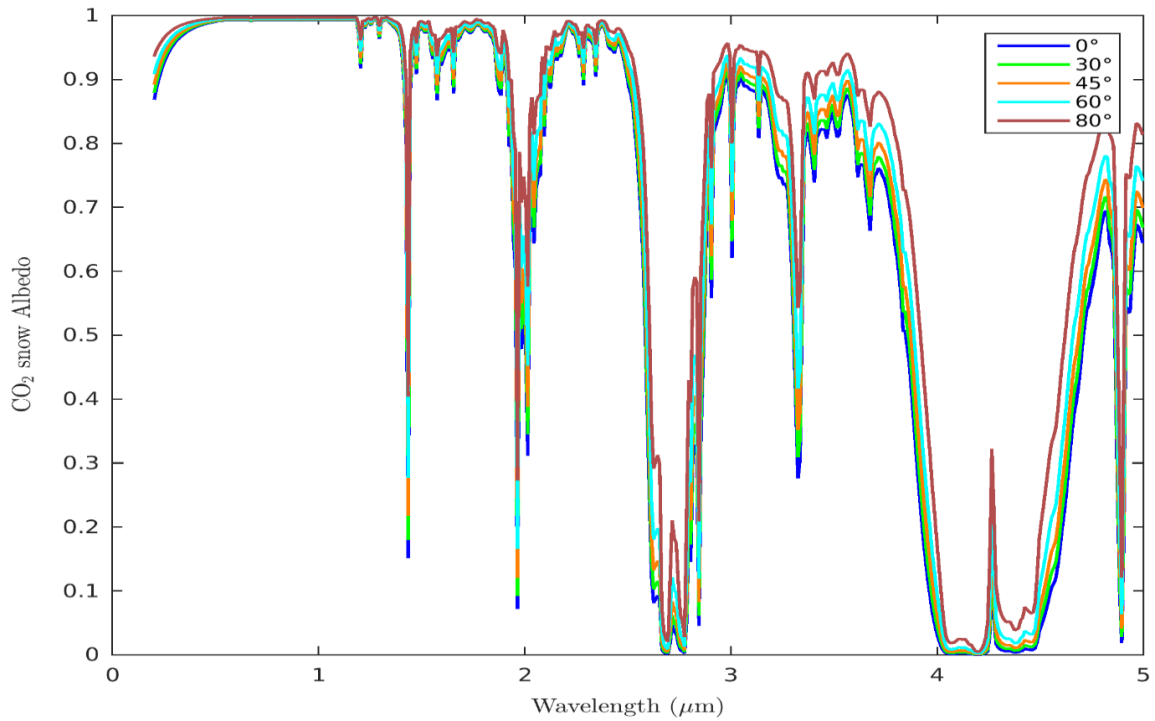


Figure 3.4: Variation of CO₂ snow albedo with solar zenith angle

3.3.2 Grain Size

The results of our simulation indicate a monotonic decrease in snow albedo with increasing effective grain size (Figure 3.5). With larger grain size, ice volume per surface area increases, resulting in greater absorption of light, as also occurs with H₂O snow [e.g., *Wiscombe and Warren, 1980*]. The broadband (near-IR) albedo of CO₂ snow drops from 0.98 to 0.93 (0.96 to 0.89) for grain size increasing from 50 μm to 1500 μm . This is also smaller variability than exhibited by H₂O snow, whose broadband (near-IR) albedo changes from 0.82 to 0.65 (0.66 to 0.37) over the same range in effective grain size. CO₂ snow albedo varies less with grain size because CO₂ ice is inherently less absorptive in the near-IR than H₂O ice. Consequently the incremental increase in absorption associated with increasing photon path length within larger ice grains is smaller for CO₂ snow than H₂O snow.

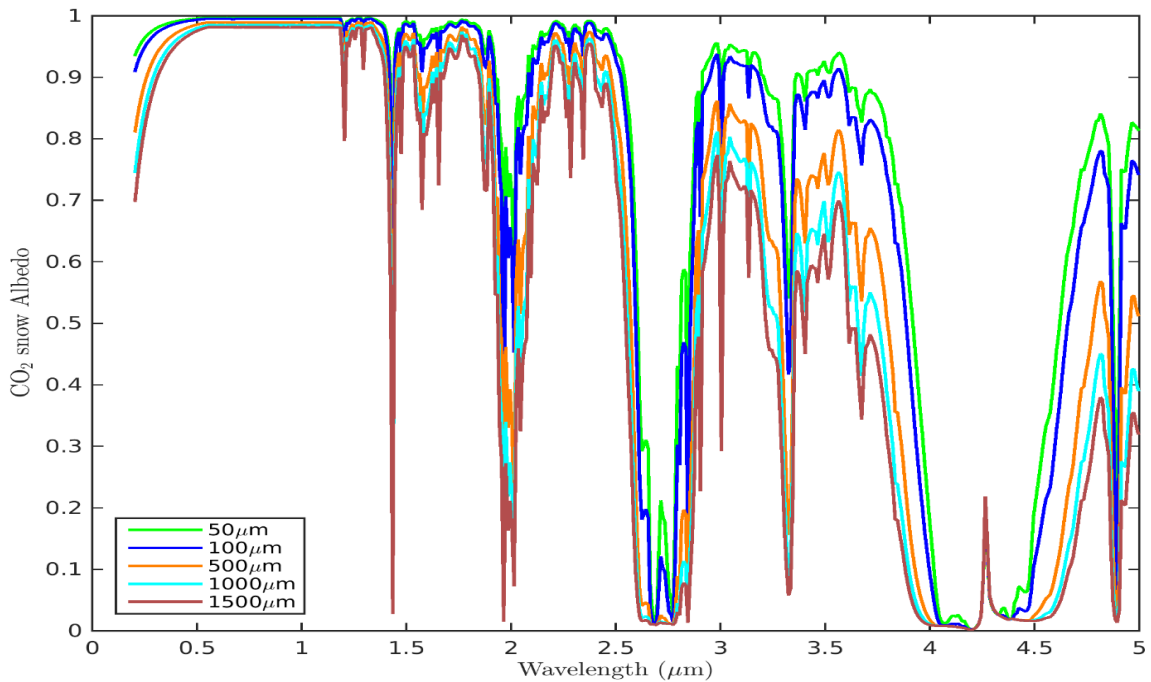


Figure 3.5: Variation of CO₂ snow albedo with effective grain radius

3.3.3 *Snow thickness*

Figure 3.6 shows the combined impacts of snow thickness and grain size on CO₂ and H₂O broadband snow albedo, for snow overlying a surface with spectrally constant albedo of 0.2 (model configuration shown in Figure 3.2a). For a given grain size, CO₂ snow albedo has a larger variation over the entire thickness range compared to H₂O snow albedo. CO₂ snow albedo is almost independent of grain size when the snow layer is thick, and consequently all curves saturate at a similarly high albedo ranging from 0.93-0.98. However, H₂O snow albedo has greater dependence on grain size, as each curve saturates at a different albedo ranging from 0.65-0.82.

We define “saturation thickness” as the snow thickness needed for the broadband albedo to differ by less than 0.01 from its semi-infinite value. Once the snow layer thickness exceeds the saturation thickness, the impact of the underlying surface becomes negligible as an insignificant amount of light penetrates through the snow to interact with the underlying surface. We note, however, that the penetration depth of radiation in snow depends strongly on wavelength, with multiple scattering leading to much deeper penetration and influence of underlying substrate at wavelengths where the ice absorbs weakly [e.g., *Wiscombe and Warren, 1980*]. The saturation thickness also depends on snow density, with higher density producing lower saturation thickness. Table 3.2 presents the saturation thickness required for different grain sizes for both types of snow. The saturation thickness is grain size-dependent, increasing from about 6 to 100 cm, and from about 5 to 83 cm as grain size increases from 50 to 1500 μm for CO₂ and H₂O snow, respectively. With other factors equal, CO₂ snow has higher saturation thickness than H₂O snow because CO₂ ice is less absorptive than H₂O ice, especially in the near-IR region, enabling multiple-

scattered photons to penetrate deeper in the snow. Also the difference in saturation thicknesses increases with larger grain size because of less dependency of CO₂ snow albedo on grain size compared to H₂O snow (section 3.3.2). Finally, we note that studies on optical properties of ice particles have found that the scattering asymmetry parameter is generally smaller for non-spherical ice particles than equal volume/area ice spheres [e.g., *Fu*, 2007; *Libois et al.*, 2013; *Räisänen et al.*, 2015]. An implication of this is that saturation thickness will be smaller for non-spherical ice particles than values reported here, as photons are scattered less preferentially into the forward direction.

Table 3.2: Saturation snow layer thickness with various grain sizes for each type of snow. CO₂ and H₂O snow densities are 327.15 and 200 kg/m³, respectively.

Grain Size	CO₂ snow saturation thickness	H₂O snow saturation thickness
50 μm	6.5 cm	5.0 cm
100 μm	10.9 cm	8.9 cm
250 μm	24.0 cm	18.8 cm
500 μm	43.1 cm	33.7 cm
1000 μm	73.8 cm	59.4 cm
1500 μm	100.2 cm	83.2 cm

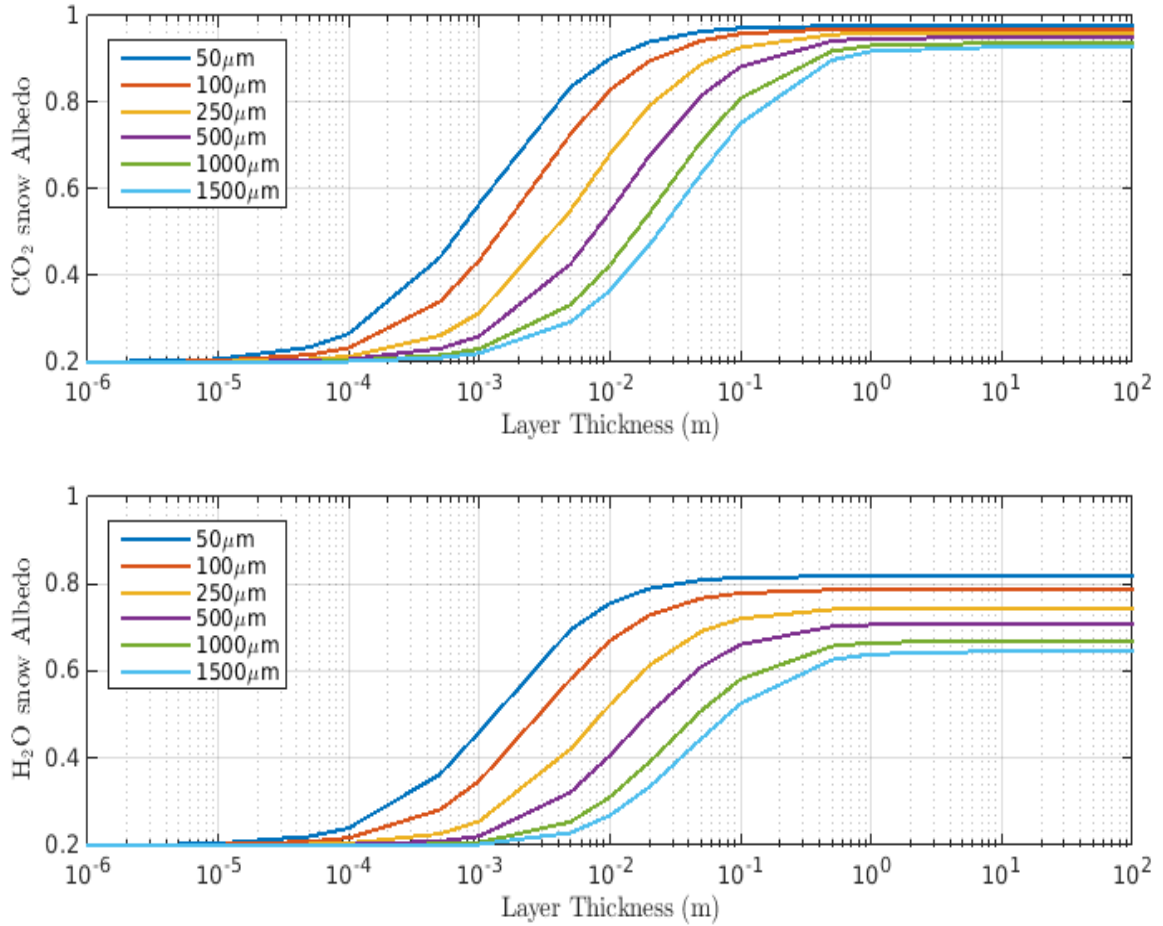


Figure 3.6: Broadband CO_2 (top) and H_2O (bottom) snow albedo dependence on snow layer thickness for various effective grain sizes. The underlying surface is assumed to have a constant albedo of 0.2. The snow density is 327.15 kg/m^3 and 200 kg/m^3 for CO_2 and H_2O snow, respectively.

3.3.4 *Dust type*

The snow albedo impacts of different types of dust depend on their optical properties, and in particular on their mass absorption cross-section. Here we compare the impacts of Earth dust, Mars dust and palagonite on semi-infinite snow albedo (Figure 3.7). Again, identical size distributions are assumed for each type of dust, and only the refractive indices are varied as input to Mie calculations. Mars dust has the greatest albedo impact of the three between wavelengths of 0.5 and 2.5 μm , which contains most of the incident solar energy, while palagonite has the least. Mars dust is relatively dark due to the presence of higher amount of iron oxides [*Bell III* et al., 1990; *Bell III*, 1996; *Christensen* et al., 2000; *Christensen* et al., 2001a; *Christensen* et al., 2001b], combined with its large mass extinction cross-section. Although average Earth dust has a lower near-IR single-scatter albedo than the Mars dust, it also has a lower mass extinction cross-section across the spectrum, due to its larger density, leading to lower absorption per unit mass of dust, expressed via the mass absorption cross-section (Figure 3.1d).

We consider “mass mixing ratio” as the mass of impurity divided by the mass of ice in which the impurity is mixed. Table 3.3 presents the effective CO_2 snow albedo in the presence of 0.01% mass mixing ratio of dust (kg of dust / kg of ice). Dust causes a larger albedo drop in the visible region (~ 0.6) compared to the near-IR region (~ 0.3) because all three types of dust have very low single scatter albedo in the blue and green spectra while ice grains scatter very efficiently at these wavelengths (Figure 3.1). Furthermore, the drop in albedo in the blue spectrum (~ 0.8) is much higher than the drop in the red spectrum (~ 0.3), helping explain the planet’s red appearance even in its cryospheric regions.

Although palagonite is less absorptive than Mars dust, the spectral variations in albedo impacts of the two species are similar.

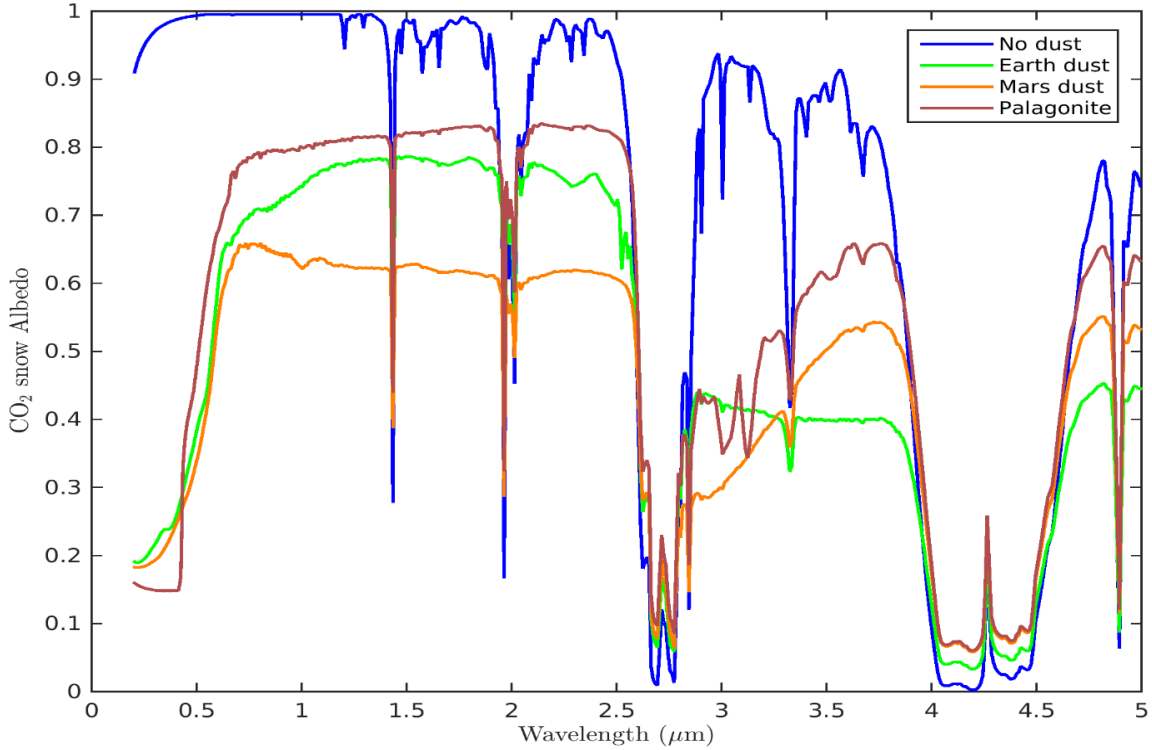


Figure 3.7: Impact of different types of dust on CO_2 snow albedo, with dust mass mixing ratio of 0.01%. The assumed snow grain size is $100 \mu\text{m}$.

Table 3.3: Impact of different types of dust on CO_2 broadband snow albedo (grain size of $100 \mu\text{m}$), with dust mass mixing ratio of 0.01%.

	pure CO_2	Earth dust	Mars Dust	Palagonite
Visible (0.2-0.7 μm)	0.991	0.429	0.395	0.477
Near-IR (0.7-5.0 μm)	0.952	0.718	0.616	0.779
Broadband (0.2-5.0 μm)	0.970	0.583	0.513	0.637

Table 3.4: Impact of different types of dust on H₂O broadband snow albedo (grain size of 100 μm), with dust mass mixing ratio of 0.01%

	pure H₂O	Earth dust	Mars Dust	Palagonite
Visible (0.2-0.7 μm)	0.988	0.450	0.415	0.494
Near-IR (0.7-5.0 μm)	0.609	0.551	0.503	0.599
Broadband (0.2-5.0 μm)	0.787	0.504	0.461	0.550

Table 3.4 presents the effective H₂O snow albedo in the presence of 0.01% mass mixing ratio of dust. Pure H₂O snow is relatively absorptive in the near-IR spectral region (Table 3.1 & Table 3.4), so dust has little impact on near-IR albedo, or in the most absorptive portions of the spectrum, it even increases albedo (Figure 3.8). Fundamentally, this occurs because the single-scatter albedo of dust exceeds that of H₂O ice grains at wavelengths longer than about 1.5 μm. Similar to CO₂ snow, the drop in albedo is large in the visible region for H₂O snow. However in the near-IR region the change is less than 0.1, leading to smaller broadband albedo impacts of dust on H₂O snow albedo than on CO₂ snow albedo. These differences indicate that dust can have greater impact on Martian cryosphere albedo than Earth cryosphere albedo due to the higher abundance of CO₂ ice compared to H₂O ice.

3.3.5 Dust Concentration

Figure 3.9 shows the effect of varying amounts of Mars dust on CO₂ snow albedo (Figure 3.2a). It is obvious that albedo decreases with increasing amount of dust, except in the strong absorption bands of CO₂ ice (e.g., near 1.4 μm, 1.9 μm, and 2.7 μm). Again we see maximum impact in the visible part of the spectrum. Saturation of the albedo from this type of dust is evident with dust mixing ratios exceeding 0.01%, with almost no difference

in albedo between scenarios with 0.1% and 1% dust. Also, the strong absorption features of CO₂ snow at various wavelengths disappear with high dust concentrations.

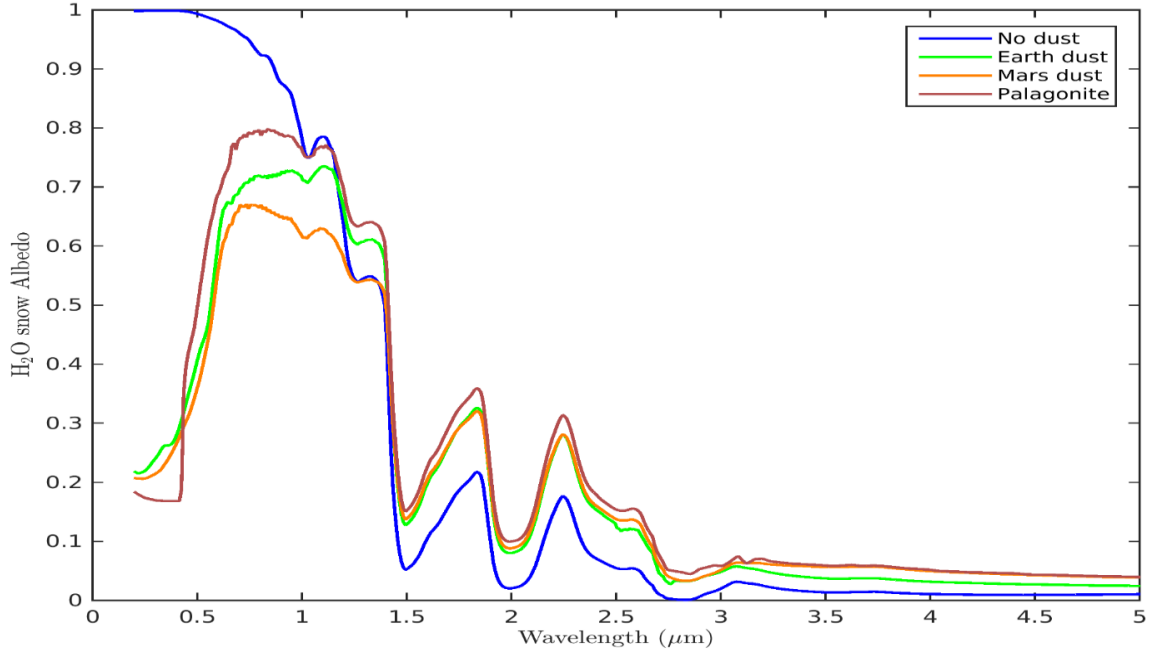


Figure 3.8: Impact of different types of dust on H₂O snow albedo, with dust mass mixing ratio of 0.01%. The assumed snow grain size is 100 μm .

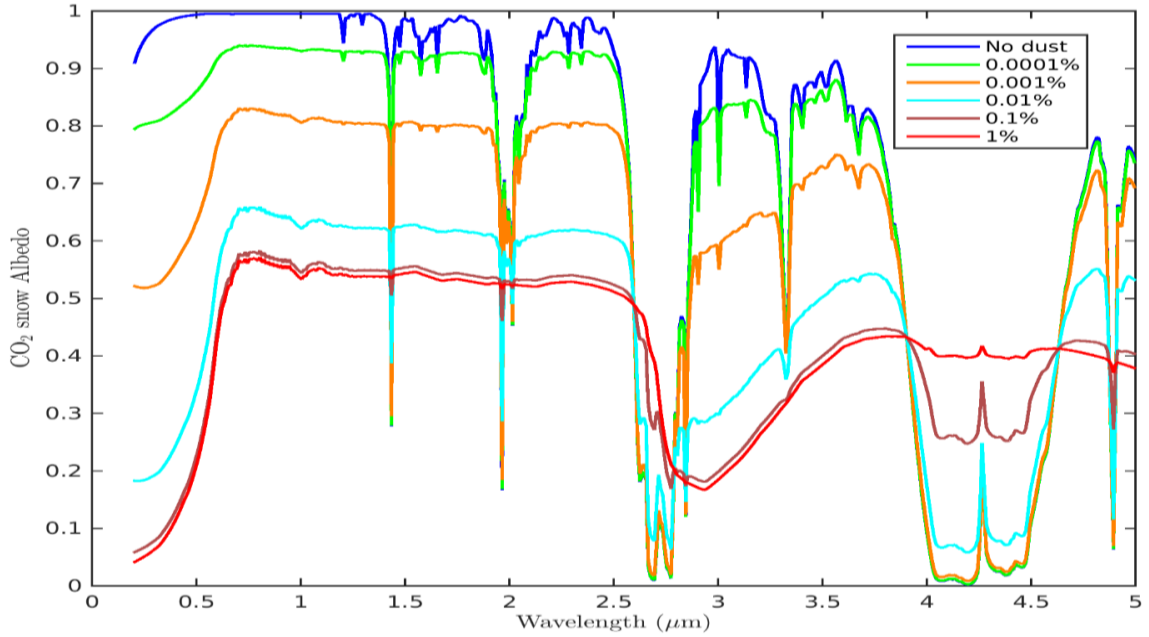


Figure 3.9: Impact of various amounts of Martian dust on CO₂ snow albedo. Snow effective grain radius is 100 μm .

3.3.6 Presence of both H₂O and CO₂ ice

3.3.6.1 Separate layers

H₂O and CO₂ ice interact quite differently with radiation in the visible and near-IR portions of the spectrum and also exhibit different sensitivities to different physical parameters. These ice types are also known to co-exist on the surface of Mars [e.g. *Byrne et al., 2008; Brown et al., 2014*]. Here, we explore the albedo effects associated with thin slabs of one type of ice overlying the other, building on sensitivity studies conducted by *Warren et al., [1990]*. Figure 3.10 shows spectral albedo of surfaces with H₂O snow on top of CO₂ snow (scenario shown in Figure 3.2c), and Figure 3.11 shows albedo with the reverse situation (scenario shown in Figure 3.2b).

Albedo decreases markedly with the presence of thin slabs of H₂O snow on top of CO₂ snow. This decrease occurs more rapidly with increasing thickness of H₂O ice than the rate of albedo increase that occurs with increasing thickness of CO₂ snow overlying H₂O snow. For a grain size of 100 μm, about 2.5 cm of H₂O snow is required to completely mask out the albedo effect of CO₂ snow (Table 3.5), whereas about 5.5 cm of CO₂ snow is required to completely mask out the albedo effect of underlying H₂O snow (Table 3.6). Again, differences are more pronounced in the near-IR spectrum (e.g., Figure 3.3, Table 1), where the contrast in absorptivity between CO₂ and H₂O ice is greater (Table 3.5 & Table 3.6). Rows marked with an asterisk (*) correspond to saturation thickness. In summary, very little H₂O snow is needed to mask the presence of underlying CO₂ snow, whereas a larger thickness of CO₂ snow is needed to prevent underlying, more absorptive H₂O snow from effecting surface albedo.

Warren et al., [1990] determined that the presence of H₂O snow on top of CO₂ snow will increase the net surface albedo, contrary to the analysis presented here. The CO₂ ice refractive indices applied in that study, however, were measured using unpurified commercial dry ice [*Egan and Spagnolo*, 1969], and data in the 1.0-2.5 μm spectral region were extrapolated from the 0.3-1.0 μm spectral region. The higher quality measurements on pure CO₂ ice provided by *Hansen* [1997, 2005], and applied here, indicate that CO₂ ice is much less absorptive than previously assumed, leading to our opposite conclusion.

Figure 3.12 shows the broadband net surface albedo dependence on layer thickness for different snow types. The black lines indicate the saturation thickness (section 3.3.3) for the two scenarios discussed above. To further explore the behavior of albedo with layer thickness, we divide each curve into two regions: with thickness lesser and higher than saturation thickness. Table 3.7 lists the slopes of albedo per layer thickness for these two regions. For top layer thickness less than saturation thickness (first region), higher slope (in absolute values) of data shown in Figure 3.12a than Figure 3.12b substantiates the conclusion that H₂O snow is harder to mask out compared to CO₂ snow. However, for the second region with thickness higher than saturation thickness, the scenario with overlying CO₂ snow (Figure 3.12b) demonstrates higher absolute slope than the reverse situation. This indicates that H₂O snow is less sensitive to thickness variation (once a minimum threshold is reached), owing to its more absorptive nature in the near-IR spectrum compared to CO₂ snow.

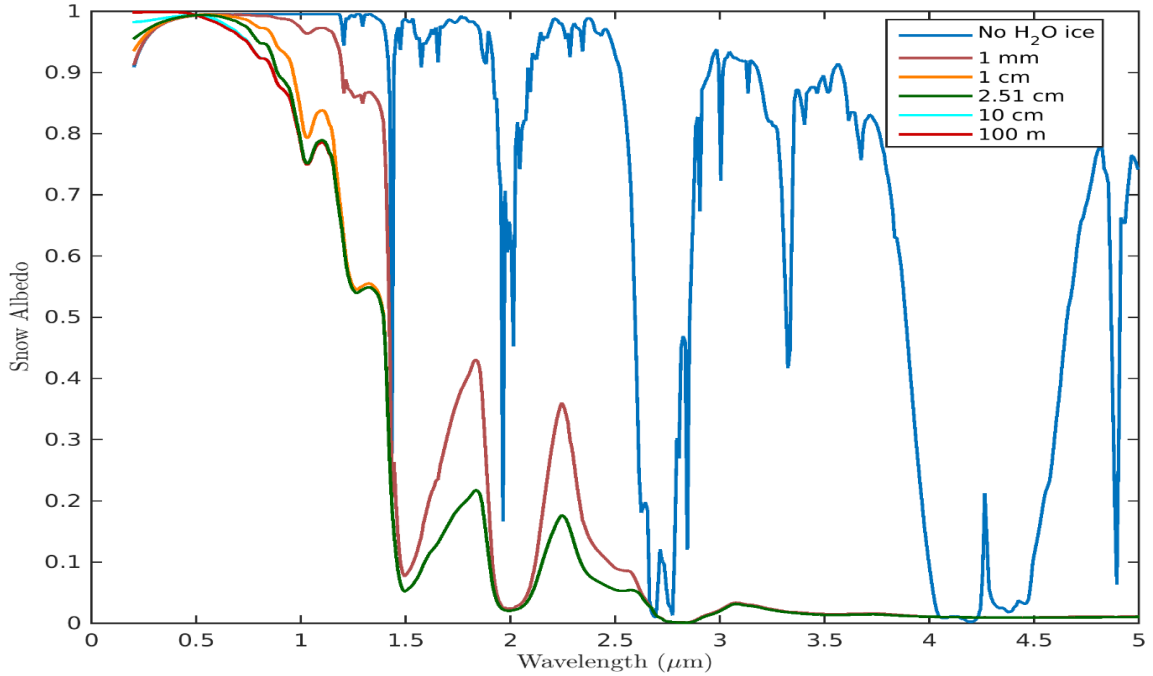


Figure 3.10: Spectral albedo of surfaces with H₂O snow layers of varying thickness on top of CO₂ snow. The green curve thickness corresponds to saturation thickness.

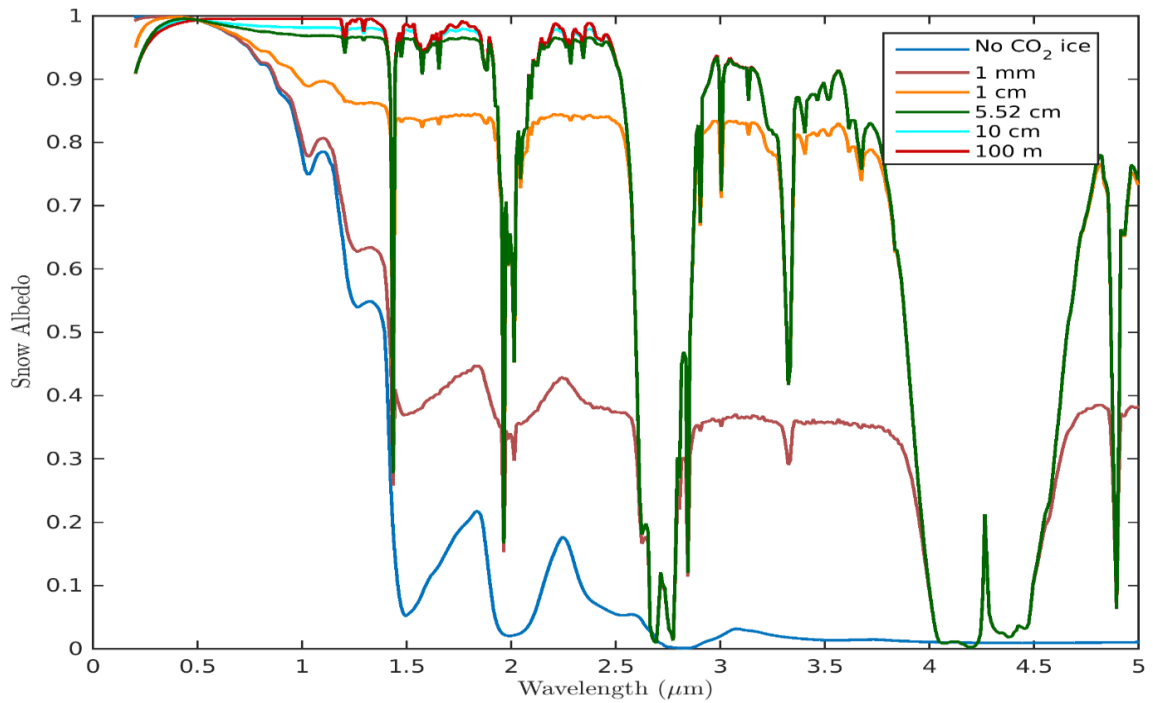


Figure 3.11: Spectral albedo of surfaces with CO₂ snow layers of varying thickness on top of H₂O snow. The green curve thickness corresponds to saturation thickness.

Table 3.5: Broadband albedos of surfaces with H₂O snow layers on top of CO₂ snow

H₂O layer thickness	Vis	Near-IR	Broadband
0	0.991	0.952	0.970
1 mm	0.991	0.739	0.857
1 cm	0.991	0.639	0.804
2.51 cm*	0.987	0.618	0.793
10 cm	0.990	0.610	0.787
100 m	0.988	0.609	0.787

*Saturation thickness (section 3.3.3)

Table 3.6: Broadband albedos of surfaces with CO₂ snow layers on top of H₂O snow

CO₂ layer thickness	Vis	Near-IR	Broadband
0	0.988	0.609	0.787
1 mm	0.988	0.703	0.837
1 cm	0.989	0.869	0.925
5.52 cm*	0.989	0.934	0.960
10 cm	0.990	0.943	0.965
100 m	0.991	0.952	0.970

*Saturation thickness (section 3.3.3)

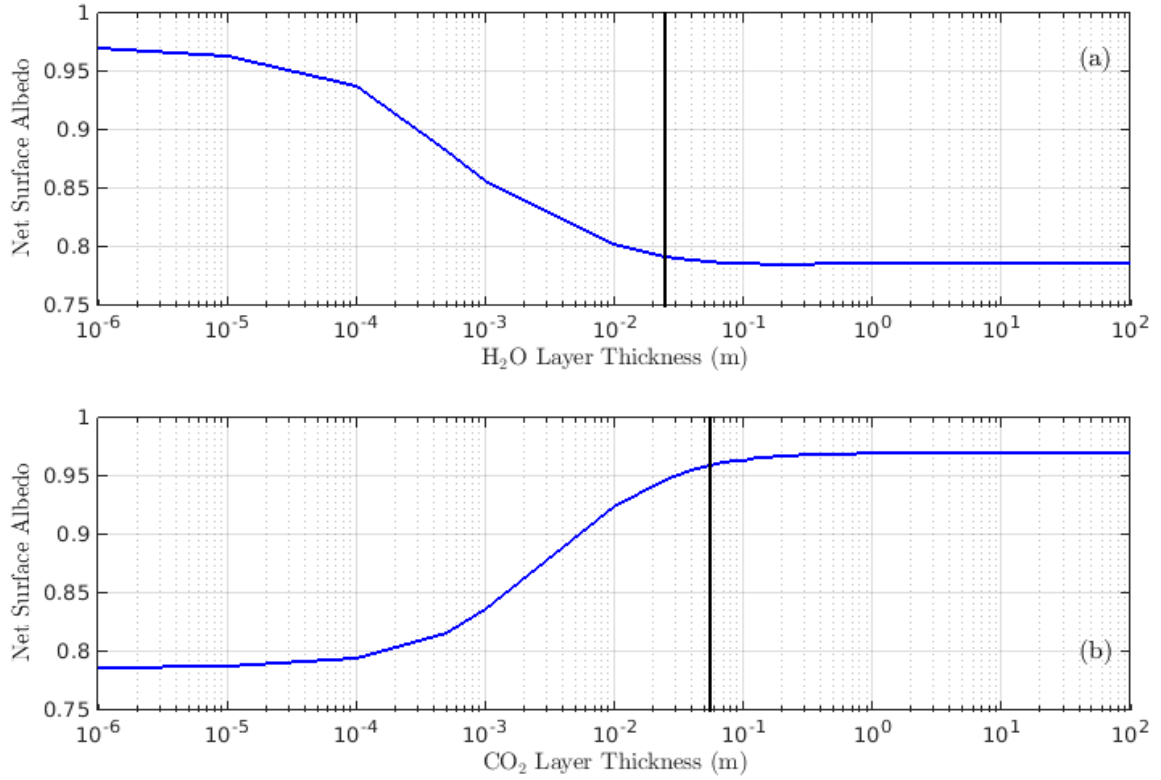


Figure 3.12: Dependence of net surface broadband snow albedo on layer thickness for (a) an H₂O snow layer on top of a semi-infinite CO₂ snow layer, and (b) a CO₂ snow layer on top of a semi-infinite H₂O snow layer. The black line indicates the saturation thickness for each case. Snow grain size is assumed to be 100 μm in both cases.

Table 3.7: Slopes of albedo per change in layer thickness for the two scenarios shown in figure 12

Thickness regime	H ₂ O snow layer on top of CO ₂ snow layer (Fig 11a)	CO ₂ snow layer on top of H ₂ O snow layer (Fig 11b)
Less than saturation thickness	-7.10 m ⁻¹	+3.14 m ⁻¹
Higher than saturation thickness	-5.54×10 ⁻⁵ m ⁻¹	+1.05×10 ⁻⁴ m ⁻¹

3.3.6.2 *Mixed layers*

In addition to the possibility of different snow types being present as separate layers, both types of snow can also become mixed together on the surface of Mars [e.g. *Brown et al., 2014*]. In this section, we explore the net surface albedo changes caused by one snow type becoming mixed with the other snow type. In practice, this is simulated by treating the less abundant ice type as an externally-mixed “impurity” with specified mass mixing ratio (scenario shown in Figure 3.2d), analogous to the treatment of dust described earlier. Figure 3.13 shows the net surface albedo when H₂O snow is present as an impurity within CO₂ snow, and Figure 3.14 shows the net surface albedo for the reverse situation. We assume a semi-infinite snow layer, and effective grain size of 100 μm for both snowpacks.

Broadband (near-IR) albedo drops by ~0.18 (~0.34) for an associated increase of H₂O snow from 0% to 10% within CO₂ snow. However, broadband (near-IR) albedo increases by ~0.10 (~0.18) in the reverse situation. This is consistent with our earlier findings that H₂O snow is relatively darker (especially in the near-IR region) compared to CO₂ snow, hence causing larger impact on net surface albedo. 0.01% of H₂O snow is sufficient to reduce the net broadband albedo of the mixture by 0.03, while 1% of CO₂ snow is required to increase the net surface broadband albedo of the mixture by same amount.

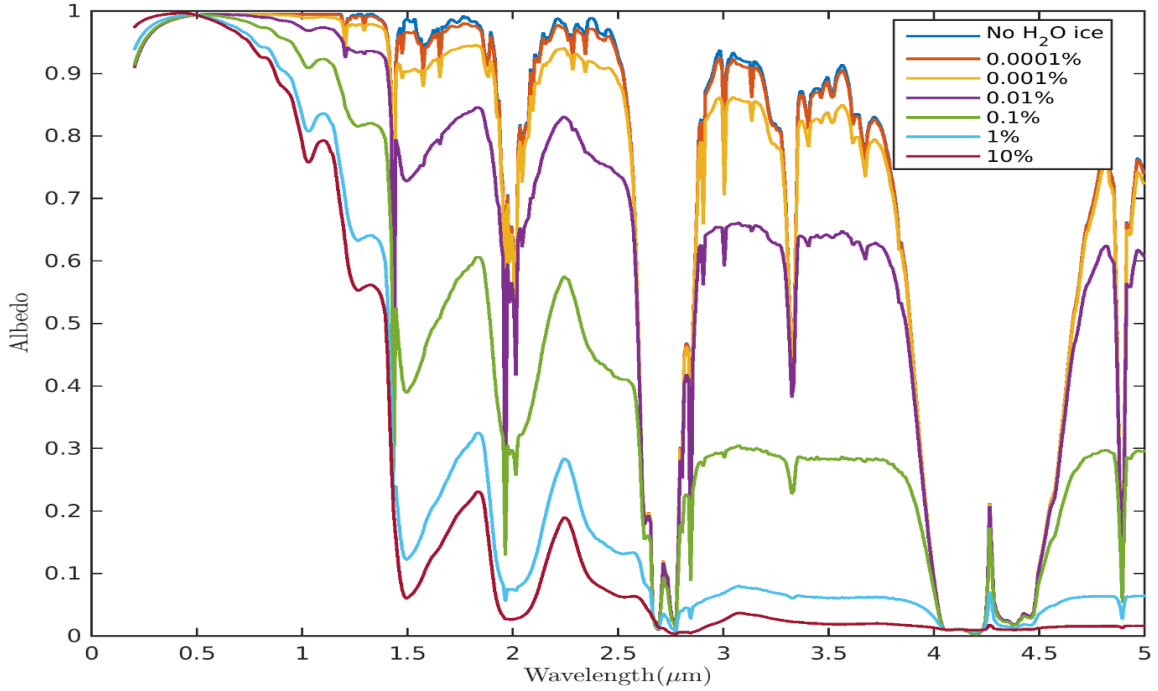


Figure 3.13: Spectral albedo of surfaces with H_2O ice present as an impurity within CO_2 snow. Effective grain size for both snow types is $100 \mu m$ and the layer is semi-infinite.

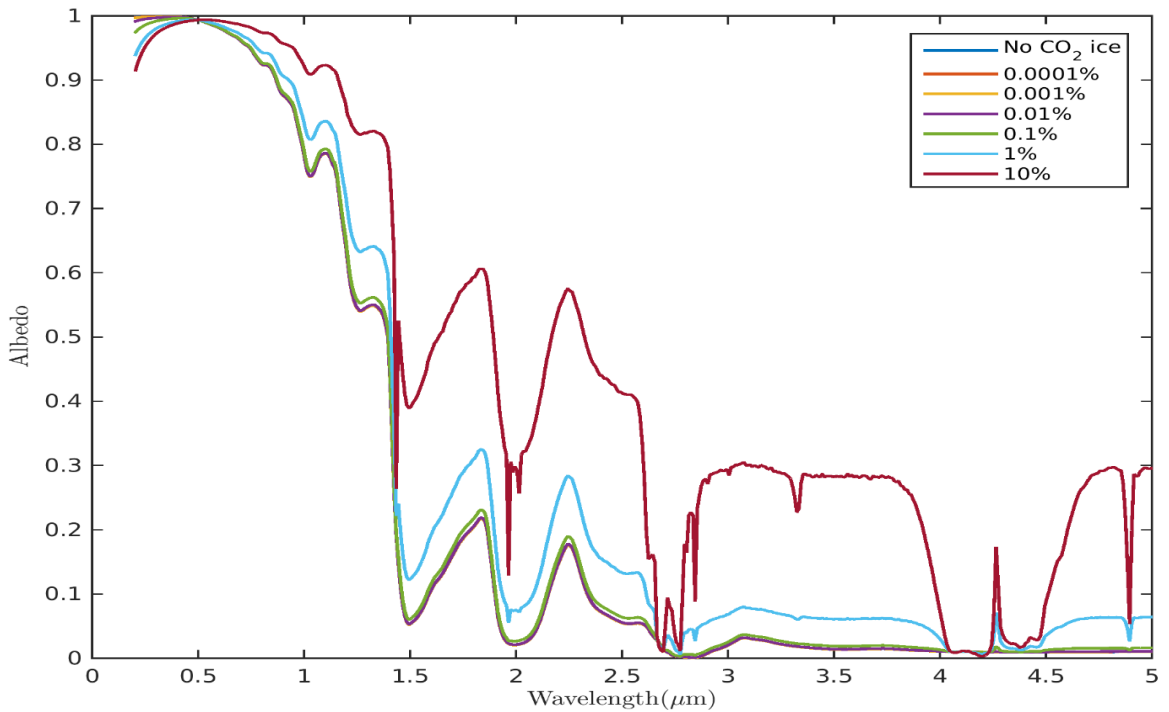


Figure 3.14: Spectral albedo of surfaces with CO_2 ice present as an impurity within H_2O snow. Effective grain size for both snow types is $100 \mu m$ and the layer is semi-infinite.

3.4 Comparison with observed Martian albedo

In this section, we compare the results of our simulations with some observations of surface albedo in the Martian cryosphere and try to identify parameter combinations that produce reasonable agreement between the modeled and observed data.

3.4.1 Comparison with CRISM measurements

We obtained spectral albedo data for one location in the Southern Hemisphere [point A, Fig. 3 of *Brown et al.*, 2014] identified to have a majority of CO₂ ice (hereafter location S), and one location from the Northern Hemisphere [point B, Fig 4c of *Brown et al.*, 2012] identified to have mostly H₂O ice (hereafter location N). The observed albedos were obtained using radiance measurements from CRISM in the 1-4 μm spectral range. Therefore we only compare our simulations in this spectral range.

Since the albedo data display a time dependency, we first take the average of all available spectra to generate a mean spectral curve for each location. Then we determine the best fit for each location by minimizing the weighted RMSE between the mean spectral curve and our simulations, which span a wide range of the parameter space described earlier. Spectral weighting for the RMSE calculation is done with the same solar spectral irradiance measurements [*Labs and Neckel*, 1968] applied throughout this study.

Table 3.8 lists the best fit values of various parameters along with spectrally-weighted RMSE values for both locations. The presence of any CO₂ ice tends to increase the RMSE at location N. This happens because observations from location N occurred during early spring to mid-summer in that hemisphere, when there is likely no CO₂ ice present on the surface [e.g. *Brown et al.*, 2014]. Location S needs about one-third the amount of dust as compared to location N for optimal fitting. One reason for this is that

Location S is deemed to have some coincident H₂O ice, which functions as a competing impurity to dust because H₂O snow substantially darkens albedo in the near-IR spectral region. For location S, changing CO₂ ice grain size does not have any significant impact on RMSE, while the H₂O ice grain size has significant impact on RMSE for location N. This is consistent with our results in section 3.3.2.

Table 3.8: Best fit parameter combinations for each Mars location. The H₂O ice at Location S is simulated as a mixture within the CO₂ snow, as opposed to a distinct layer.

Location S		Location N	
CO₂ ice grain size	1000 μm	H₂O ice grain size	250 μm
H₂O ice grain size	200 μm	CO₂ ice grain size	n/a
Amount of H₂O ice	0.07%	Amount of CO₂ ice	0%
Amount of dust	0.007%	Amount of dust	0.02%
RMSE	0.0676	RMSE	0.0542

Figure 3.15 and Figure 3.16 show the observed spectral albedo curves along with simulated curves using parameter values shown in Table 8. We observe an outlier point at 2.7 μm wavelength in the observed data for both locations and at all times. Albedo values at this wavelength are uncharacteristically higher than at nearby wavelengths, whereas both CO₂ and H₂O snows have very high absorptivity at this wavelength (Figure 3.3). This leads us to conclude that the high observed albedos at this wavelength are spurious, and may be due either to an artefact of the measuring spectrometer (CRISM), atmospheric anomalies associated with dust or clouds [Brown *et al.*, 2010], and/or a calibration issue in the retrieval of surface reflectance at that wavelength.

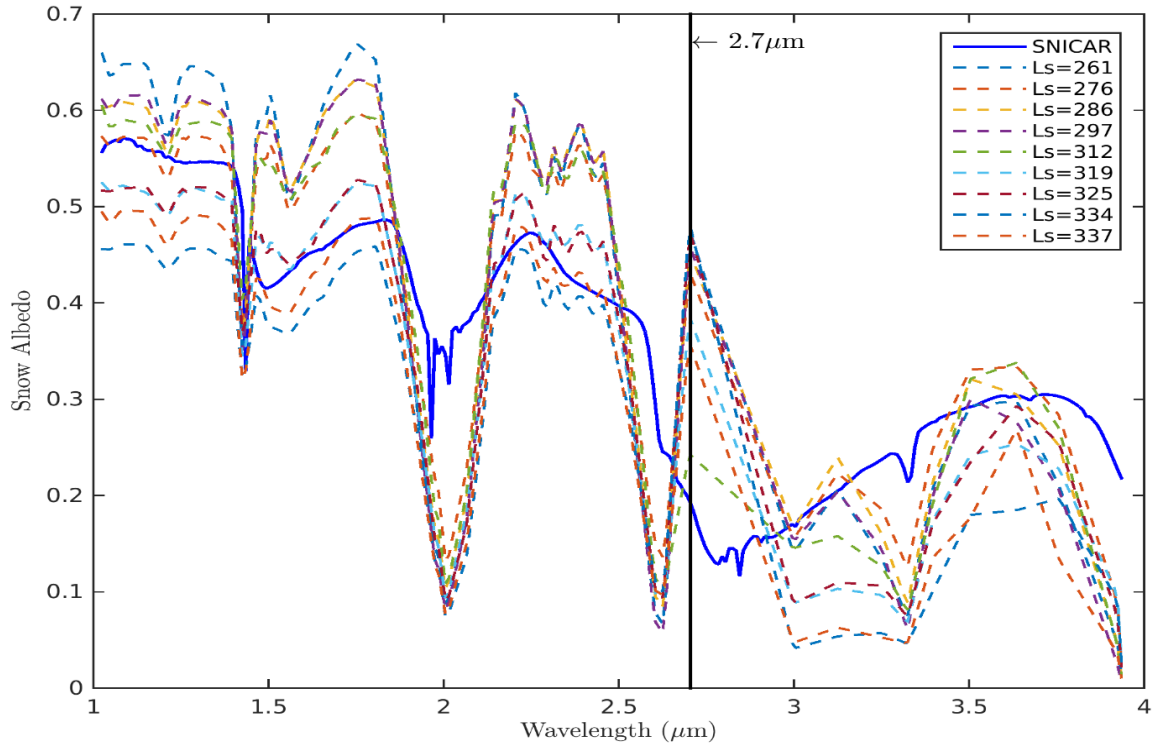


Figure 3.15: Observed albedo for location S (dotted curves) along with modeled albedo using best-fit parameters (solid blue curve). The black vertical line at $2.7\ \mu\text{m}$ indicates the outlier described in the text.

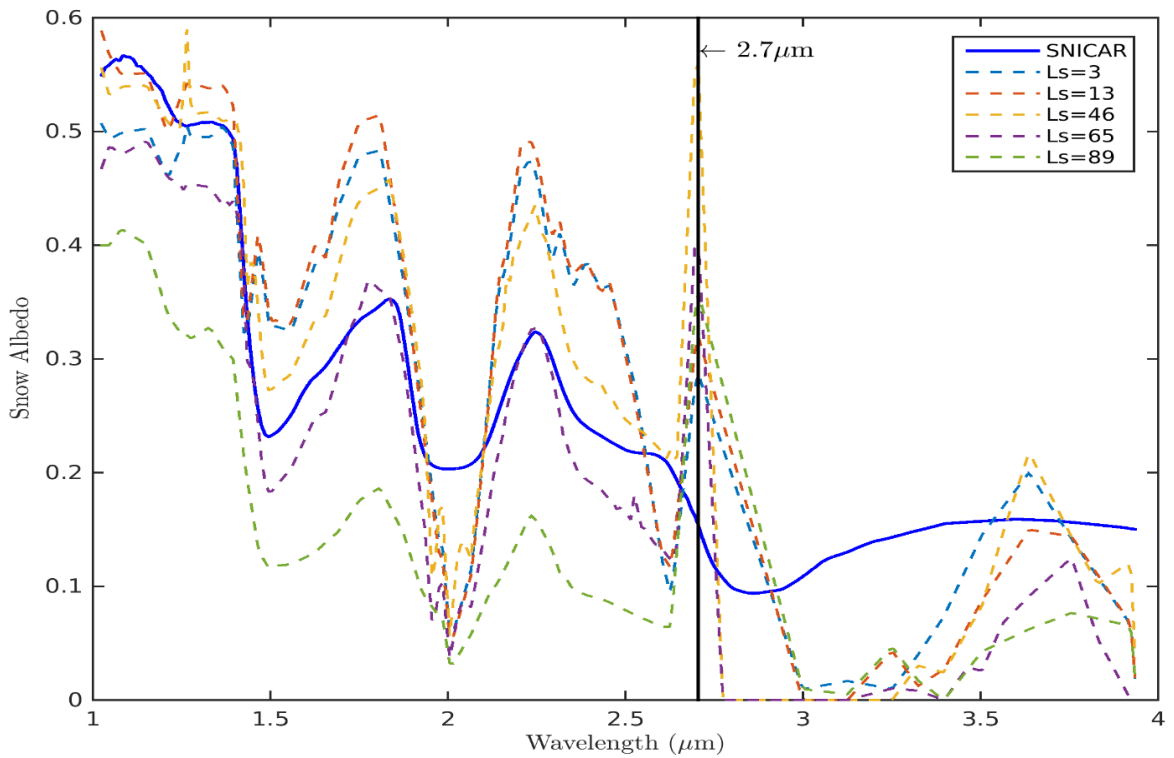


Figure 3.16: Same as fig. 15 but for location N

3.4.2 Comparison with OMEGA measurements

We also obtained spectral albedo for Observation 2621_1 (extracted at 65.3°N/341°E acquired at $L_s = 2.4^\circ$) from the OMEGA instrument [Fig. 18 from *Appéré et al.*, 2011; shown later here], and modeled spectral albedo of CO₂ ice [Fig. 1 from *Appéré et al.*, 2011] simulated with a different radiative transfer model developed by *Douté and Schmitt* [1988]. Figure 3.17 shows the comparison of CO₂ snow albedo spectra from both models. The grain size applied in both models is 1 mm. The model used by *Appéré et al.*, [2011] shows weaker absorption in the key absorption bands (e.g., near 1.4 μm , 1.9 μm , and 3.3 μm), consistent with the use of optical constants in that study from *Quirico and Schmitt* [1997]. Updated CO₂ ice optical constants applied here were measured from thick samples [*Hansen*, 2005], which have significant advantages (described in detail by *Hansen*, 2005) over the measurements of optical constants conducted by *Quirico and Schmitt* [1997] using thin films. The differences in CO₂ snow spectral albedo simulated by these two models translates into a net spectrally-weighted RMSE of 0.090 for CO₂ snow. These differences, while not that large, could affect the determination of best-fit model parameters derived from observations.

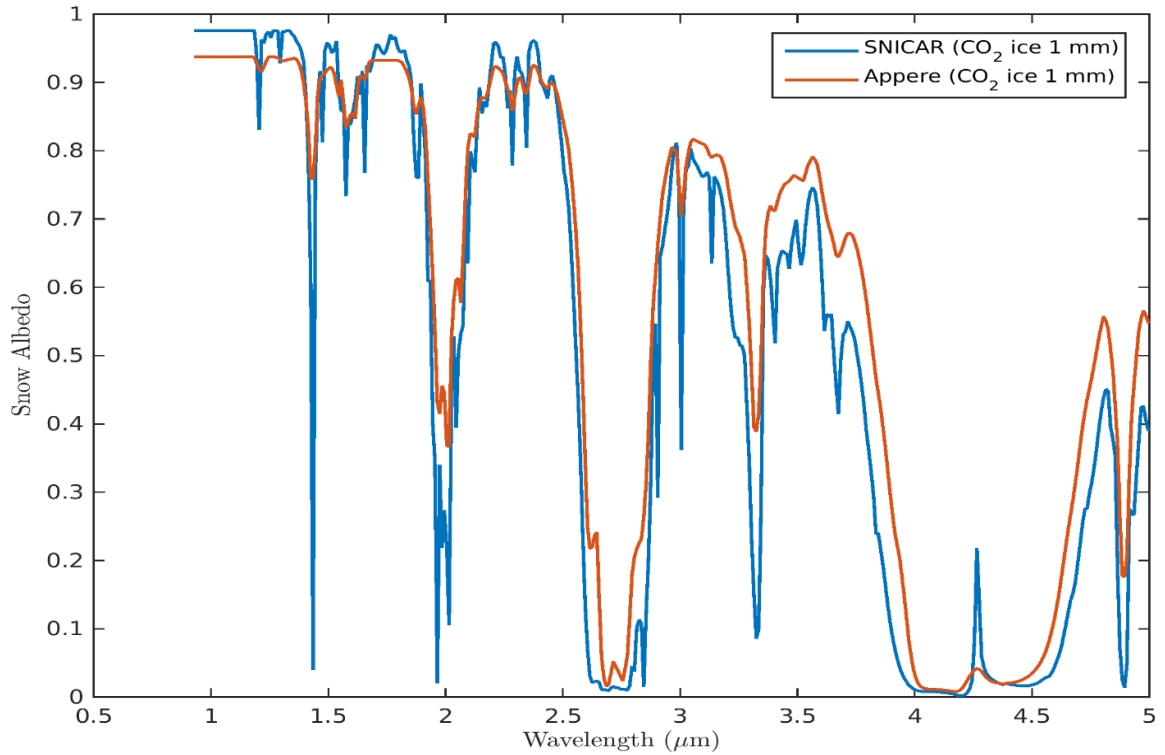


Figure 3.17: Comparison of CO₂ snow albedo simulated with SNICAR and shown in figure 1 of Appéré et al. [2011]. The snow grain size assumed in both models is 1mm.

Table 3.9: Comparison of best-fit parameters between SNICAR and Appéré et al., [2011]

	SNICAR	Appéré et al., [2011]
CO₂ ice grain size	3000 μm	7 cm
Amount of CO₂ ice	99.83%	99.75%
H₂O ice grain size	100 μm	200 μm
Amount of H₂O ice	0.16%	0.19%
Dust grain size	1.5 μm	13 μm
Amount of dust	0.0094%	0.06%
RMSE	0.027	0.016

As shown in Table 3.9 and Figure 3.18, *Appéré et al.*, [2011] results are in good agreement with simulated near-IR snow albedo and OMEGA observations of a CO₂-rich deposit. This result was obtained by assuming a mixed CO₂/H₂O snow configuration with a CO₂ grain size of 7 cm, which interestingly is much larger than what is deemed to be realistic under most conditions [*Barr and Milkovich*, 2008]. We also obtained a reasonable agreement (spectrally-weighted RMSE of 0.027) with the OMEGA observations presented by *Appéré et al.*, [2011] using a mixed CO₂ and H₂O configuration of SNICAR applying a smaller CO₂ grain radius and slightly different mass fractions of dust and H₂O (Table 3.9; Figure 3.18). However, our results do not agree as well as that of *Appéré et al.* [2011]. Interestingly, the quality of our fit ceases to change much with larger grain sizes of CO₂ ice.

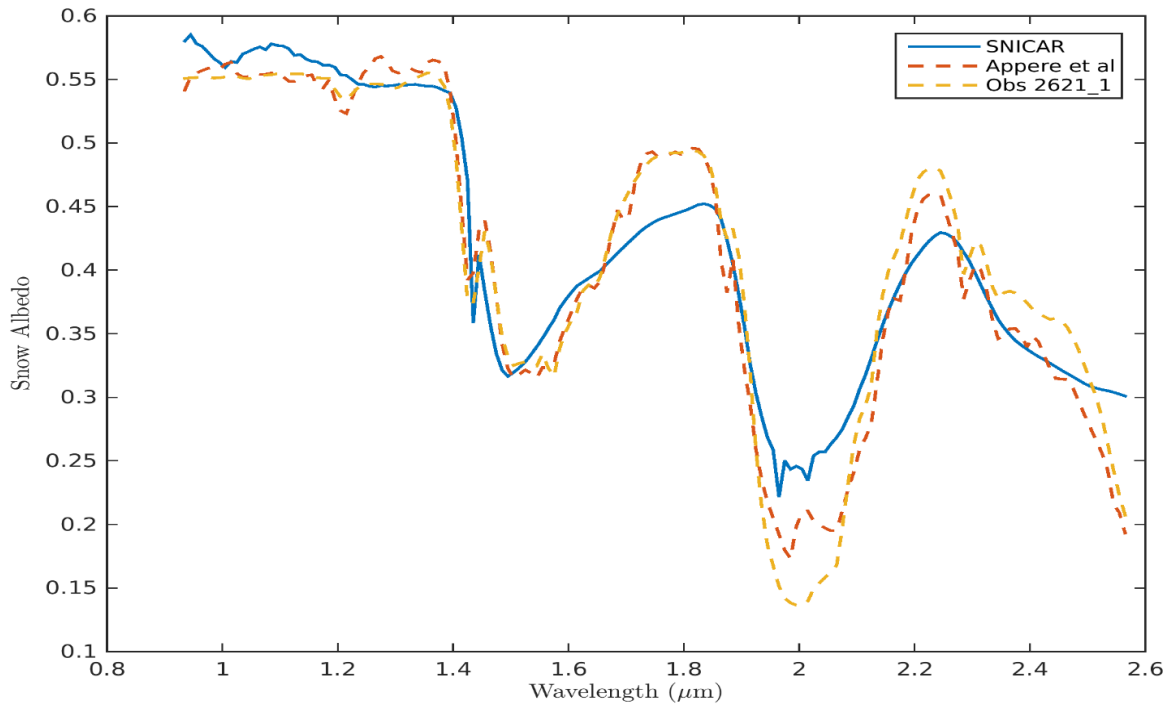


Figure 3.18: Comparison of best fit albedo using SNICAR and the model applied by *Appéré et al.*, [2011] against data from OMEGA observation 2621_1.

3.5 Conclusions

We have simulated the spectral albedo of CO₂ snow using an enhanced 480-band version of SNICAR, which was previously developed for terrestrial snow. We explored CO₂ snow albedo across the entire solar spectrum, including UV, visible and near-IR wavelengths. Our analysis shows significant differences between H₂O and CO₂ snow albedos. H₂O snow is about 8 times more absorptive than CO₂ snow in the near-IR region, and 7 times more absorptive averaged over the entire solar spectrum. CO₂ snow albedo has very little dependence on solar zenith angle and a weaker dependence on grain size than H₂O snow. The broadband albedo of CO₂ snow decreases by only 0.064 as effective grain size increases from 50 to 1500 μm .

The presence of thin snow layers exposes underlying surfaces to incoming radiation, hence impacting the surface albedo. The saturation thicknesses for CO₂ snow and H₂O snow range from approximately 6.5 – 100 cm and 5 – 83 cm, respectively, for effective grain sizes ranging from 50 to 1500 μm , though we caution again that these thicknesses may be smaller with non-spherical ice particles. Non-spherical grains scatter less strongly in the forward direction, thereby decreasing the penetration depth of solar radiation. Thicker CO₂ snow is required to negate the impact of underlying surface because CO₂ ice grains scatter more strongly than H₂O ice grains, especially in the near-IR spectrum. The presence of 0.01% dust reduces the broadband albedo of CO₂ snow by about 50%, with Martian dust being the darkest type of dust explored here, followed by typical Earth dust and palagonite. The spectral shape of albedo changes caused by palagonite, which is often used as a Martian dust analog, closely follow those of Mars dust, but palagonite is not as

absorptive as Mars dust. The impact of dust on CO₂ snow albedo saturates after its mass mixing ratio exceeds roughly 0.1%.

Because of the contrasting properties of H₂O and CO₂ ice in the near-IR spectrum, different layering and mixing configurations of these types of snow can have substantial impact on net surface albedo. With an effective grain size of 100 μm, only about 2.5 cm of H₂O snow is needed to mask the influence of underlying CO₂ snow on net surface albedo (Table 3.5), while more than double this amount of CO₂ snow is needed to mask the influence of H₂O snow (Table 3.6). Such effects are relevant for the perennial H₂O ice caps of Mars, and where water vapor from the atmosphere condenses on top of CO₂ ice in other areas of the planet. When both snow types are present as a mixture rather than separate layers, 0.01% of H₂O snow reduces the broadband albedo of CO₂ snow by 0.03, while 1% of CO₂ snow is required to increase the broadband albedo of H₂O snow by the same amount.

With the identification of optimal snow parameter combinations, our results show a decent agreement between modeled and observed spectral albedo of the Martian polar ice caps in the 1–4 μm spectral range. The CRISM data exhibit anomalously high albedo at 2.7 μm, which cannot be explained with presence of either CO₂ or H₂O ice, as both media are highly absorptive at this wavelength. SNICAR also provides realistic best-fit parameters for matching OMEGA near-IR observations. Simulations presented here could potentially be used in combination with observed data to refine the calibration of surface albedo retrieval algorithms. This study also indicates that spectral albedo observations of the Martian cryosphere could be used in combination with inverse modeling to loosely infer parameters such as the amount of dust, presence of each ice type, and H₂O effective

grain size. Model results can also be used to interpolate measurements to higher spectral resolution. The new spectrally resolved albedos for CO₂ snow presented here have potentially wide applicability to any planet or system with CO₂ ice, especially Mars.

Chapter 4

Integration of SNICAR with LMD Mars GCM

4.1 Introduction

As discussed in the previous chapters, surface albedo plays an important role in any planet's energy budget and driving its climate system. Water-vapor, cloud and, albedo feedbacks are the three most powerful positive feedback mechanisms operating within the current climate system on Earth [e.g., *Bony et al.*, 2006; *Winton*, 2006; *Randall et al.*, 2007; *Soden et al.*, 2008; *Shell et al.*, 2008; *Flato et al.*, 2013]. Due to the small amounts of water-vapor on Mars (as compared to Earth), cloud (carbon dioxide) and albedo feedbacks are primary feedback mechanisms of climate change. Dust is a critical component of Martian climate, and plays an important role on the Martian climate [e.g., *Madeleine et al.*, 2011]. It can either amplify or weaken the other feedback mechanisms significantly.

Dust is tightly coupled with the other components of the climate system. Global dust storms can cause changes in meteorological phenomena (i.e., dust storms, dust devils, clouds, recession of the polar caps, and surface temperatures) that can persist for periods ranging from few weeks to more than one Mars year [*Cantor*, 2007]. The dust and CO₂ cycles have long been known to greatly affect the present-day climate of Mars. Although

dust cycles are highly variable on the Mars, CO₂ condensation cycle is highly repeatable on multi-annual timescales [Hess *et al.*, 1980]. The dust and CO₂ cycles on Mars show no correlation between them [e.g., Kahre and Haberle, 2010]. Viking landers measurements (VL-1 and VL-2) show a highly repeatable CO₂ cycle on multi-annual timescales [Hess *et al.*, 1980], while demonstrating a highly variable dust cycle (including large dust storms) over the years [e.g., Zurek and Martin, 1993]. Snow albedo is largely affected by the amount of impurity present in it.

Current version of LMD Mars GCM (hereafter MGCM) uses annually repeating albedo values from the Thermal Emission Spectrometer (TES) observations. These observations also show significant seasonal variation of snow (both H₂O and CO₂) albedo [Kieffer *et al.*, 2000; Kieffer and Titus, 2001]. Small changes in albedo would have compelling impact on the climate dynamics, especially for very low pressure (<1% of Earth's atmosphere) atmosphere of Mars. Therefore, it is important to prognostically determine ice cap (both H₂O and CO₂) albedos interactively in real-time inside the model.

4.2 Methodology

We apply the extended SNICAR model (chapter 3) to the current version of MGCM to calculate real-time interactive surface broadband snow albedo. Extended version of SNICAR utilizes 480 bands spanning 0.2-5.0 μm at 10 nm spectral resolution. We divide our broadband (0.2-5.0 μm) into two sub-regions (0.2-0.5 μm and 0.5-5.0 μm) which are used for atmospheric radiative transfer calculations in original MGCM. These band-averaged values are weighted with solar spectral irradiance measurements from *Labs and Neckel* [1968]. We determine the optical properties of Martian dust [Wolff *et al.*, 2006; Wolff *et al.*, 2009; Wolff *et al.*, 2010] using Mie Theory with an assumed gamma size

distribution [Hansen and Travis, 1974] with $r_{eff}=1.5\mu\text{m}$ and effective variance (v_{eff})=0.3 [Wolff et al., 2006] (chapter 3, section 3.2). SNICAR replaces the original scheme which used a simple prescribed surface albedo. SNICAR coupled with MGCM calculates surface albedo interactively using prognostic surface dust deposition flux (already present in the model) as an input. As pure snow albedos are much higher than observed albedos, first step in this process was to determine a base amount of dust to bound simulated albedos within realistic limits. First the original model was run for a long-term with dust accumulation throughout. The dust amount at the time-step for which minimum difference between TES observed albedo and modeled albedo is achieved is used as a “baseline dust” for model initialization. This albedo difference was minimized only for snow-covered regions.

Next step in the process was to have realistic amount of dust in snow for long duration model runs. Dust exchange on Martian surface is caused by the winds around the entire planet. We used an exponential decay mechanism to achieve consistent dust amount in ice caps. For a given time ‘ t ’, net concentration of dust in snow $d(t)$ at the surface is given by:

$$d(t) = [d(t-1) + d_0] e^{-t_0/\tau} \quad (4.1)$$

where ‘ d_0 ’ is the dust deposition flux calculated by original MGCM for that time step, ‘ t_0 ’ is the physical time-step (1/48 sol; 1 sol = 88775 seconds) of model run, and ‘ τ ’ is the decay constant. For $t = 0$, we use baseline dust to represent $d(t-1)$. To calculate the value of ‘ τ ’, first we run the model for one Martian Year (MY) without any decay mechanism. Next, we replace $d(t)$ with average baseline dust determined earlier, $[(d(t-1) + d_0)]$ with average model accumulated dust at the end of run, and ‘ t_0 ’ with the

total run-time (in seconds) of the model. We found the value of ' τ ' to be equal to 1.36×10^8 seconds (~1528 sols) using this technique.

Finally, the new version of SNICAR was implemented in MGCM to provide higher resolution, and more precise albedo simulations for snow/ice covered surface. We also determine the impact of cryosphere and dust on shortwave energy flux of the Mars. We assume zero change in all physical quantities at non-snow covered regions for all scenarios in this chapter. Therefore changes over non-snow covered regions are not plotted in any maps. All global average parameters are area-weighted means for that particular quantity.

4.3 Results and Discussion

Figure 4.1 shows the global annual mean surface albedos for various scenarios on Mars. Snow-free albedo is bare ground albedo (prescribed in the model) without the presence of any snow. Dust-free albedo is pure snow albedo, and original MGCM albedo is simulated using original version of MGCM (without SNICAR). SNICAR albedo (global annual mean) is the albedo computed after the integration of SNICAR into the MGCM. This results in an albedo that has a correlation of 0.999 with original MGCM albedo. Dust-free and snow-free albedos are derived from diagnostic dust-free and snow-free albedo calculations respectively within MGCM. These diagnostic calculations do not affect the simulated climate state. The global annual mean values are not far apart for both original MGCM and SNICAR scenarios (Table 4.1). Table 4.1 also presents mean albedo values calculated over snow-covered surface only. Dust-free albedos are much higher than the SNICAR albedo due to the presence of high amount of dust within snow (chapter 3, section 3.3.5). Also CO₂ snow is more susceptible to dust darkening as compared to H₂O snow, especially in the near-IR spectral region (chapter 3, section 3.3.4).

Table 4.1: 7-year global annual mean albedo for various scenarios on Mars. First row correspond to values plotted in Figure 4.1. All means are area-weighted.

Scenario	SNICAR	Original MGCM	Dust-free	Snow-free
Mean Albedo (Global)	0.24	0.23	0.3	0.2
Mean Albedo (Snow-covered)	0.34	0.33	0.63	0.21

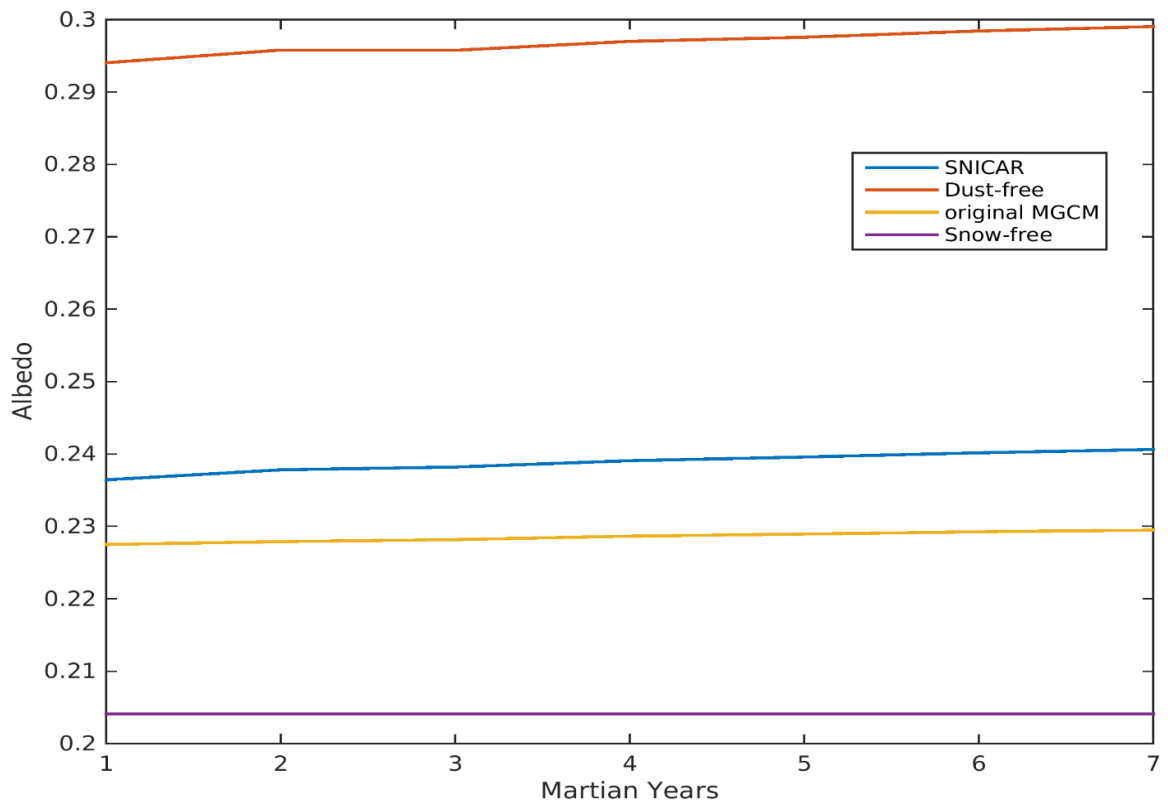


Figure 4.1: Annual global mean surface albedo of Mars calculated for various scenarios.

4.3.1 Comparison between SNICAR and MGCM simulations

Figure 4.2 shows the global map of difference in albedo simulated using SNICAR and original MGCM. Difference is taken by subtracting MGCM annual mean values from SNICAR annual mean values. In general, SNICAR albedos are higher than MGCM albedos near poles while smaller near mid-latitudes. MGCM simulates much higher dust deposition near mid-latitudes as compared to poles (Figure 4.3), resulting in lower snow albedo simulated by SNICAR. This effect was not present in original version of MGCM because the albedo was insensitive to the dust variability on surface. Dust deposition flux in Figure 4.3 is consistent with the original MGCM generated dust flux, although the distribution has been amplified due to baseline dust used with SNICAR implementation. Figure 4.4 shows the variation of albedo difference, dust deposition flux and net surface wind speed. We estimate a high correlation coefficient of 0.69 between wind speed and dust deposition flux. This indicates that the wind pattern is primary cause for specific dust distribution features. Albedo changes also follow the similar distribution as dust deposition flux (positive with lower dust concentration and negative with higher dust concentration), especially in Northern Hemisphere. We observe a relatively higher change in albedo in Southern Hemisphere due to higher abundance of CO₂ ice [e.g., *Kieffer et al.*, 2000], which is brighter than H₂O ice (chapter 3). These changes in albedo significantly impact the shortwave flux, especially at the surface (Figure 4.5). After SNICAR integration, we estimate global annual mean change in albedo and shortwave flux (at surface) of -0.0053 and 0.47 W/m² respectively from original MGCM.

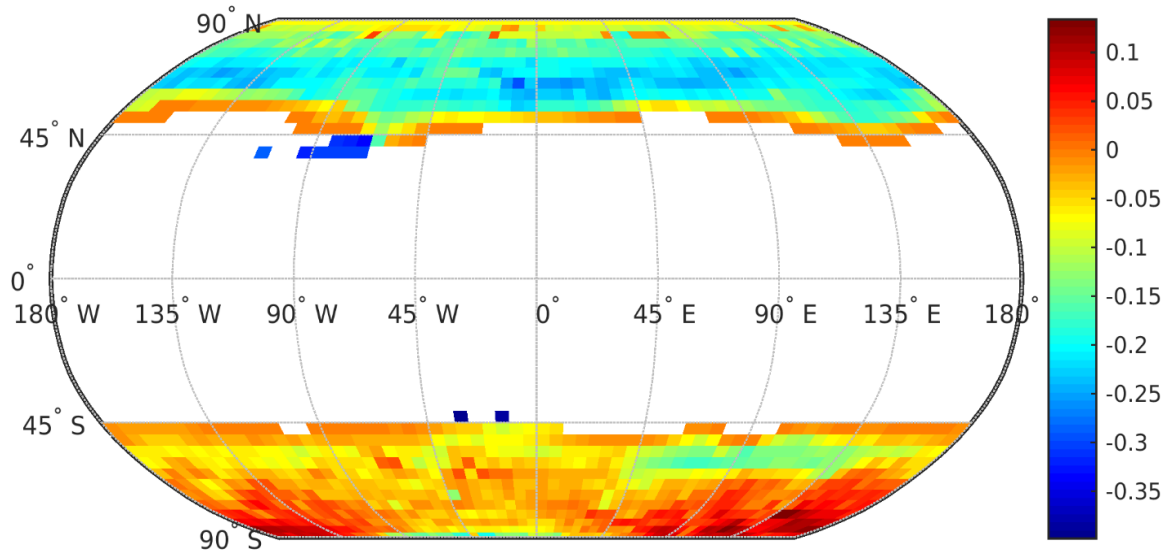


Figure 4.2: Global map of mean albedo difference between SNICAR and MGCM (mean SNICAR albedo – mean MGCM albedo).

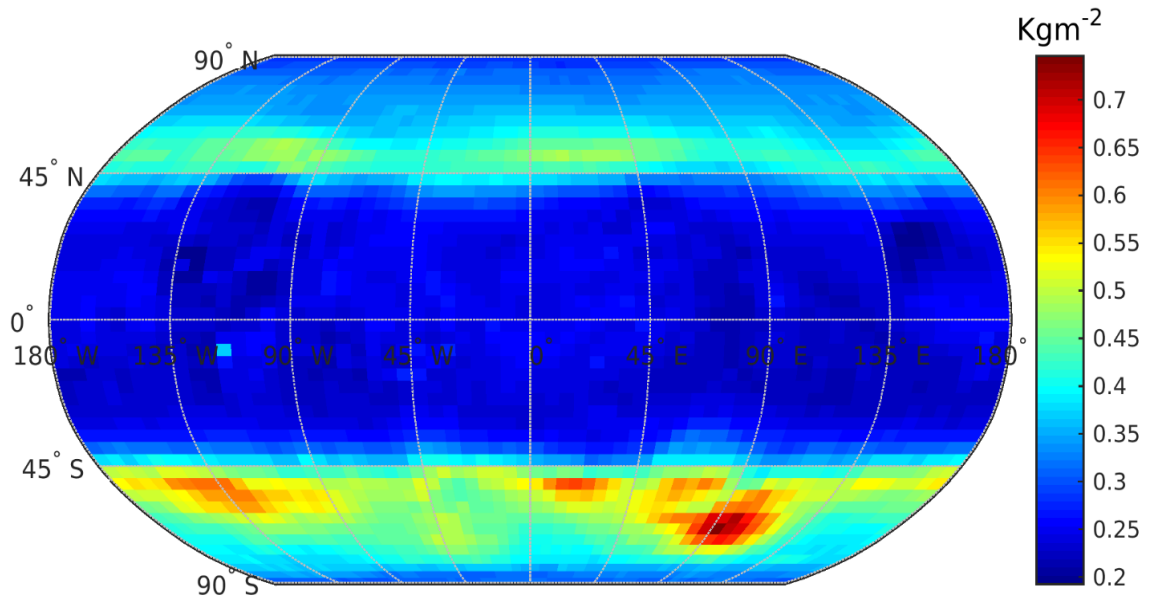


Figure 4.3: Global map of 7-years annual mean dust deposition on Mars surface.

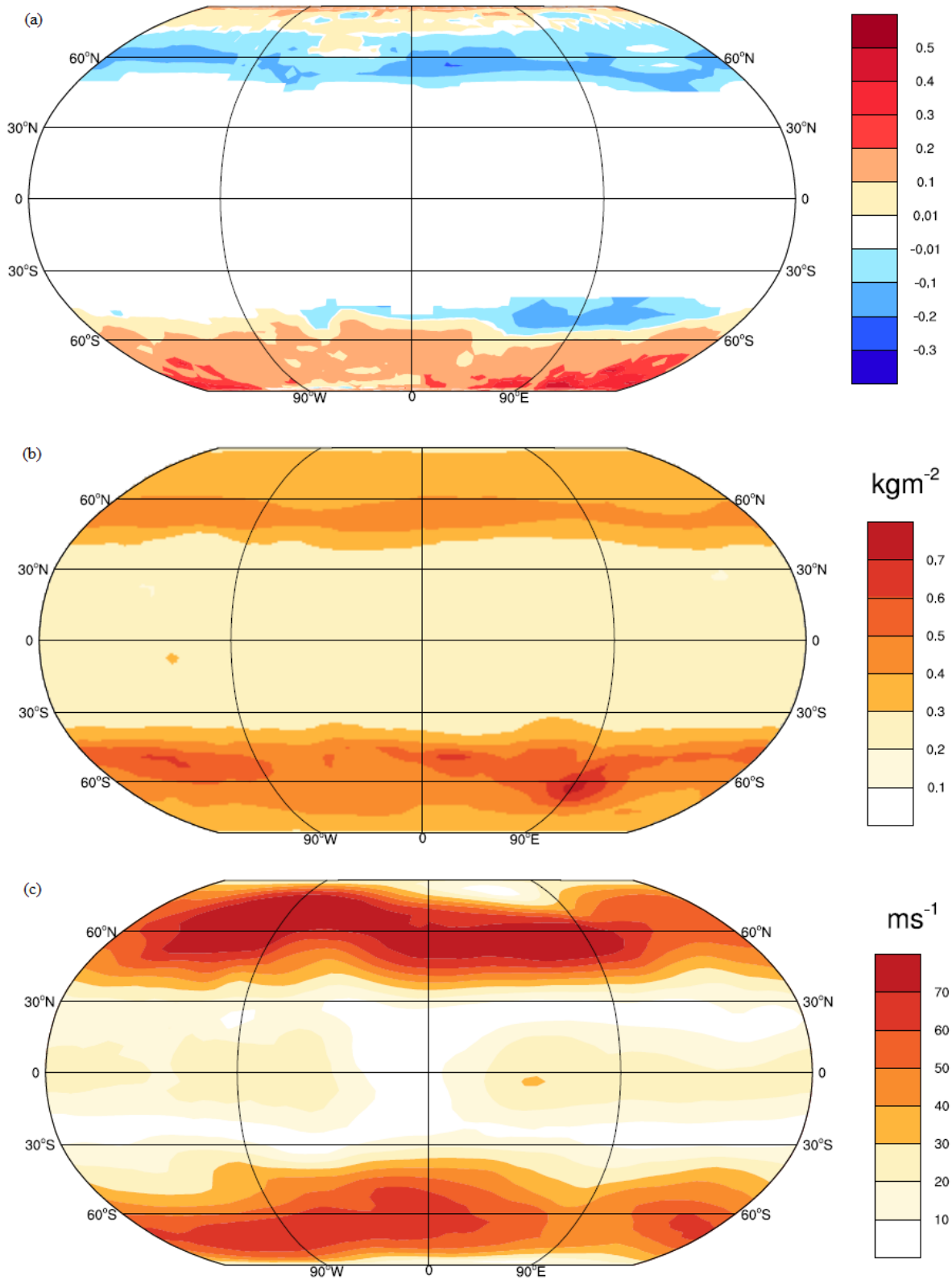


Figure 4.4: Annual mean maps highlighting (a) albedo difference (re-plot of Fig. 4.2); (b) net dust deposition flux (re-plot of Fig. 4.3); (c) absolute surface wind speed.

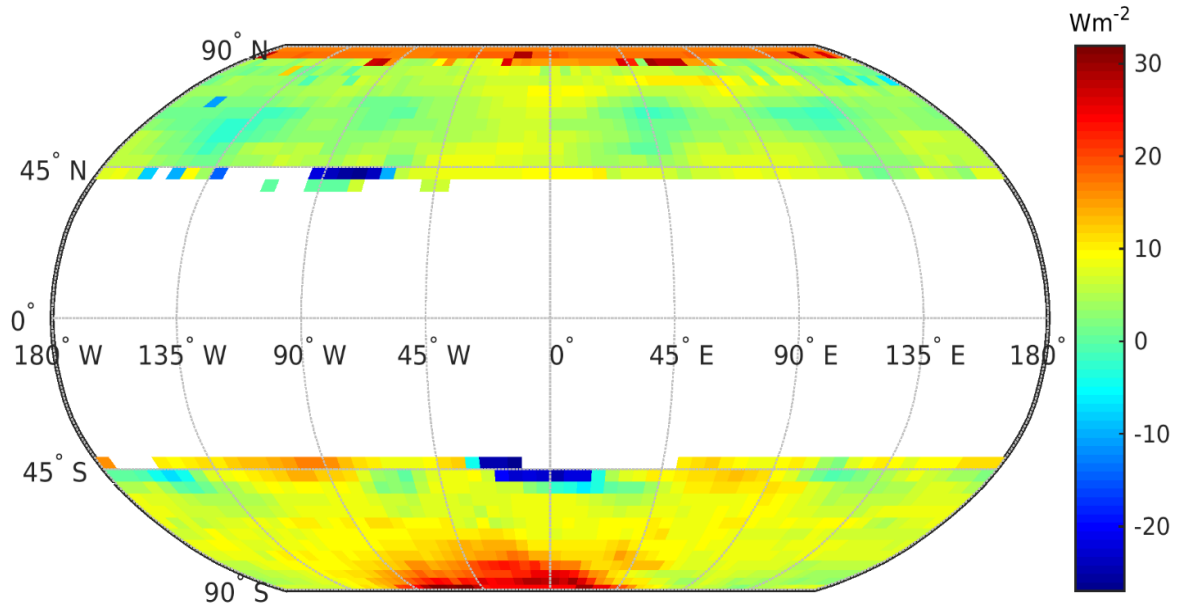


Figure 4.5: Global map of mean shortwave flux difference at the surface between SNICAR and MGCM (mean SNICAR albedo – mean MGCM albedo).

4.3.2 Comparison between snow-covered and snow-free scenarios

After integrating SNICAR into the MGCM, we estimate the impact of cryosphere on Martian albedo and shortwave energy flux. Both snow-free and snow-covered simulations are done with baseline dust and model generated dust deposition flux present within the system. Net concentration of dust in snow is determined using Equation 4.1. Figure 4.6 shows the global map of difference of bare ground albedos from snow-covered albedos. Few grid-points near mid-latitudes show a slightly higher albedo for bare ground as compared to snow-covered albedo. In these regions snowfall frequency and amount are both lower as compared to Polar Regions. A high concentration of dust within small amount of snow can decrease the surface albedo significantly, and bring it lower than the actual bare ground albedo. Also, the strong absorption features of CO_2 snow at various wavelengths disappear with high dust concentrations (chapter 3, section 3.3.5). Figure 4.6 shows the changes in shortwave energy flux at surface between snow-covered and snow-

free scenarios. On Mars, cryosphere causes a global annual mean change of 0.015 and -0.27 W/m² in albedo and shortwave flux (at surface) respectively.

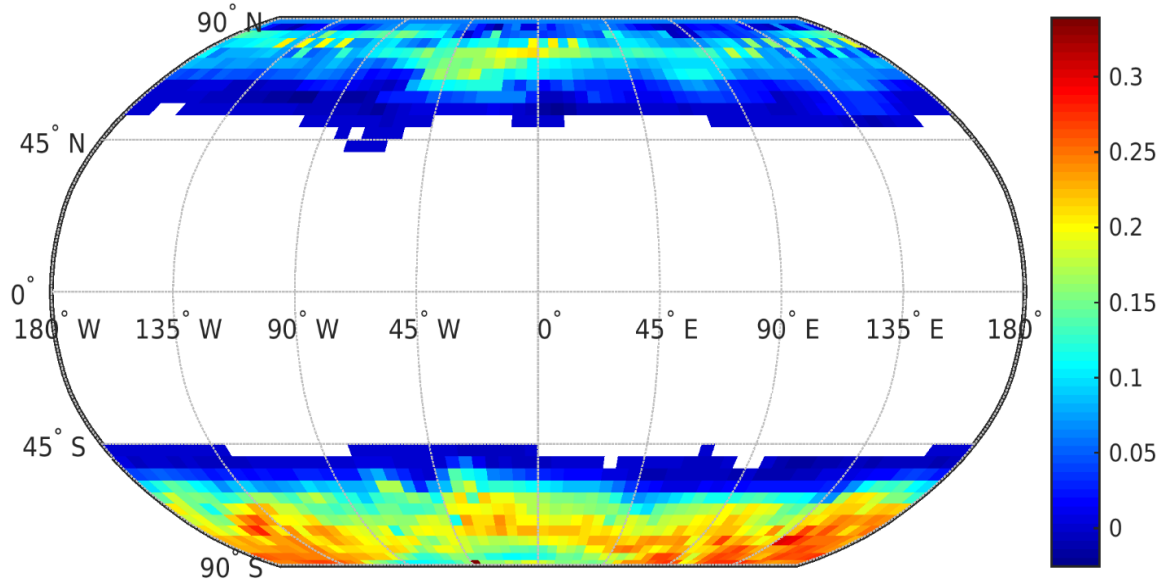


Figure 4.6: Global map of mean albedo difference between snow-covered and snow-free albedo

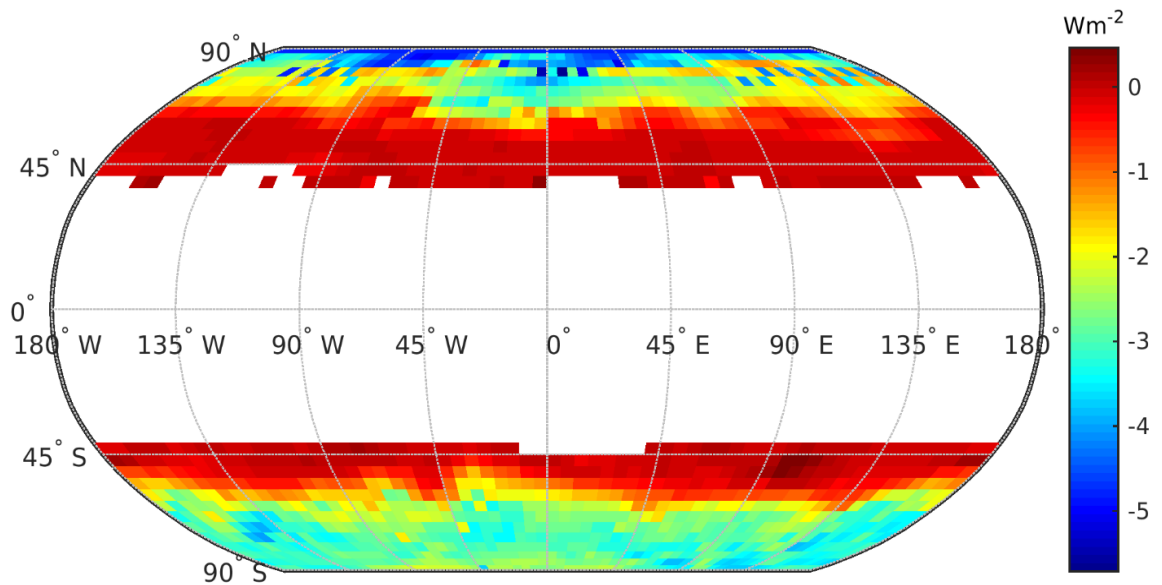


Figure 4.7: Global map of mean shortwave flux difference at the surface between snow-covered and snow-free scenarios using SNICAR.

4.3.3 Impact of dust

Next we calculate the impact of dust darkening on Martian snow albedo and shortwave energy flux. Net global annual mean dust-free albedo is higher by about 0.06 than the actual albedo calculated using SNICAR (Figure 4.8). Differences are negligible near the cap edge due to scarcity of snow. Except for those edges, albedo change is higher near mid-latitudes as compared to poles. This happens due to decrement in dust amount from mid-latitudes to poles (Figure 4.4b). Dust causes a global annual mean change of 0.41 W/m^2 in shortwave flux at the surface (Figure 4.9). Change caused by dust in snow is higher (absolute values) than the snow+dust cryosphere itself. This indicates the importance of dust in the Martian climate system.

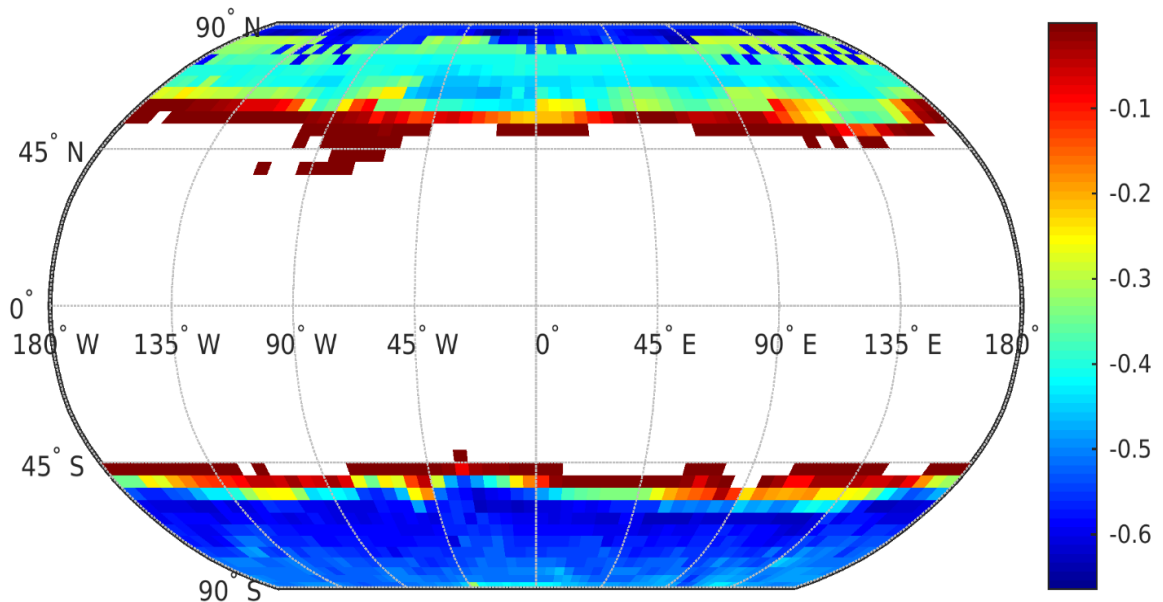


Figure 4.8: Global map of mean albedo difference between SNICAR albedo and dust-free albedo.

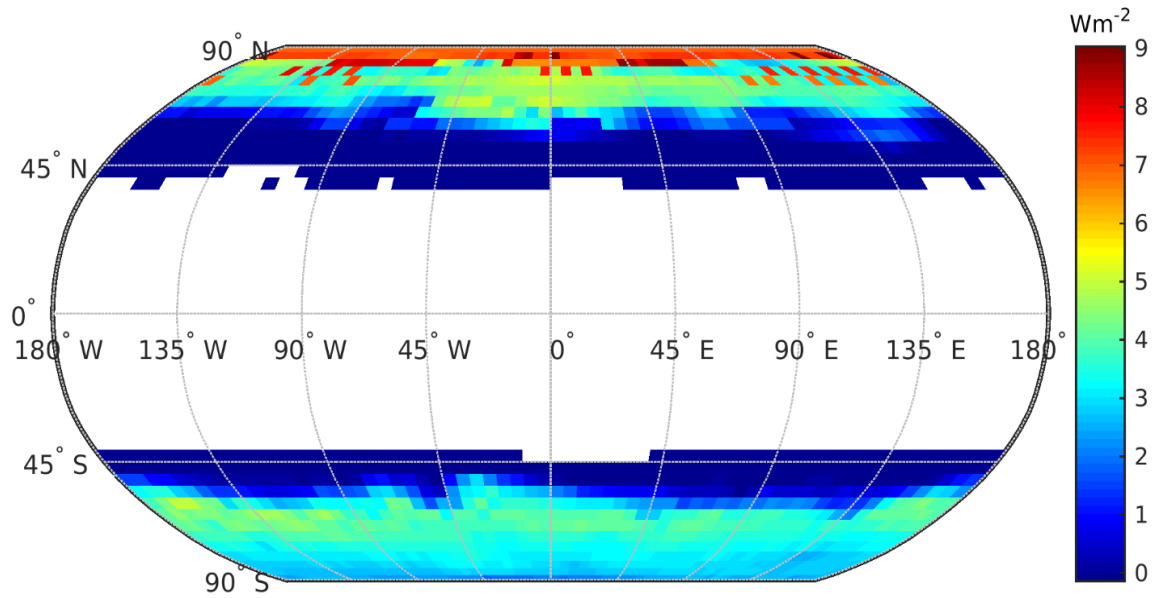


Figure 4.9: Global map of mean shortwave flux difference at the surface between SNICAR albedo and dust-free scenarios using SNICAR.

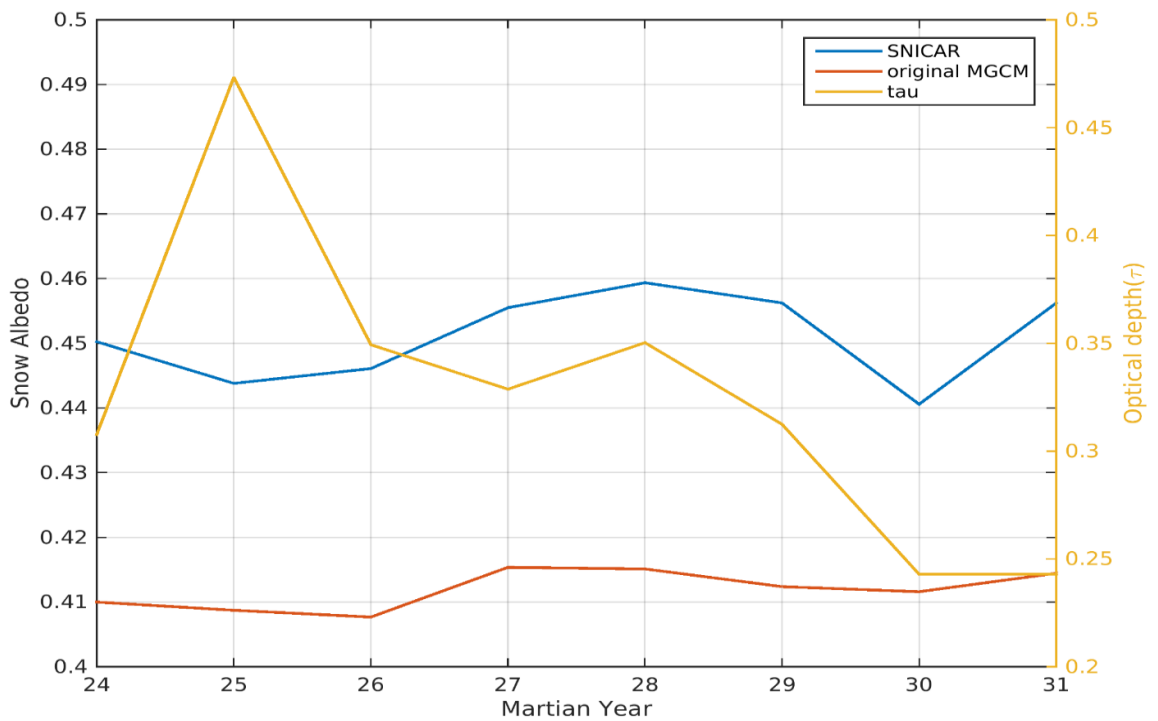


Figure 4.10: Global annual mean of snow albedos using SNICAR and MGCM, and optical depth of dust for various Martian years.

Finally, we run the MGCM for various Martian Years (MY24 – MY31) separately using different initial conditions specific to each year shown in Figure 4.10. We initially observe that the original MGCM albedos are not very sensitive to dust optical depth over the years. We observe a higher sensitivity of albedo with dust with SNICAR integrated into the model. We achieved a correlation of 0.64 with SNICAR as compared to 0.5 with original MGCM (both values are absolute). MY30 has some issues (yet to be determined) with SNICAR integrated run, and therefore not included in the correlation calculations for either cases.

4.4 Conclusions

We successfully integrated SNICAR into the current version of the MGCM. We use an initial dust content to constrain the surface albedo within realistic boundaries for Mars. We also use an exponential decay method to release dust on the surface over time. Although the new SNICAR integrated model produces slightly higher albedo than the original MGCM, it generates albedos with higher sensitivity to surface dust content as compared to original MGCM. In some cases SNICAR albedo is smaller because of higher dust in that region. Changes in albedo and surface dust content also impact the shortwave energy flux at the surface. SNICAR integrated model simulates a change of 0.47 W/m^2 in shortwave flux. We observe a high dependency (correlation coefficient = 0.69) of dust deposition at surface with surface wind speed. Changes in albedo also show the similar distribution as dust concentration over the globe.

Real snow albedos are much lower than the dust-free albedos due to the high amount of dust deposition on Mars, and higher susceptibility of CO_2 snow albedo to dust darkening as compared to H_2O snow. Dust is one of the key elements in determining the climate state

of Mars. The impact caused by dust is about 1.5 times (absolute values) of the impact caused by the presence of snow itself on shortwave flux at the surface. Runs with different initial conditions for various years show higher correlation between dust optical depth and surface albedo with the SNICAR integrated model. There are some minor technical issues which are yet to be resolved. We are hopeful that these issues will be resolved soon, and we will be able to include SNICAR in the new version of MGCM.

Chapter 5

Conclusions and Future Work

5.1 Summary of Results

Due to the strong feedback effect of ice-albedo on planetary climate systems, it is necessary to accurately account for this effect for producing realistic climate simulations on Earth and Mars. Being a strong feedback mechanism, albedo feedback can significantly delay or amplify a global seasonal change like melting of glaciers on Earth, CO₂ snow cover on Mars etc.

In Chapter 2, we apply measurements from the MODIS, combined with microwave retrievals of snow presence and radiative kernels produced from 4 different models, to derive CrRE over global land during 2001–2013. We estimate global annual mean land CrRE during this period of -2.6 W/m^2 , with variations from -2.2 to -3.0 W/m^2 resulting from use of different kernels, and variations of -2.4 to -2.6 W/m^2 resulting from different algorithmic determinations of snow presence and surface albedo. Slightly more than half of the global land CrRE originates from perennial snow on Antarctica, whereas the majority of the northern hemisphere effect originates from seasonal snow. Mountain glaciers resolved in 0.05 degree MODIS data contribute about -0.037 W/m^2 (1.4%) of the global effect, with the majority (94%) of this contribution originating from the Himalayas. During a common overlap period of 2001–2008, our MODIS estimates of the northern hemisphere land CrRE are about 18% smaller (less negative) than previous estimates

derived from coarse-resolution AVHRR data, though inter-annual variations are well correlated ($r=0.78$), indicating that these data are useful in determining longer term trends in land CrRE.

In Chapter 3, we adapt the terrestrial SNICAR model to simulate CO₂ snow albedo across the ultraviolet, visible, and near-IR spectrum (0.2-5.0 μm). We apply recent laboratory derived refractive indices of CO₂ ice, which produce higher broadband CO₂ snow albedo (0.93–0.98) than previously estimated. Compared with H₂O snow, we find that CO₂ snow albedo is much higher in the near-IR spectrum, less dependent on ice grain size, less dependent on solar zenith angle, and more susceptible to darkening from dust. The presence of small amounts of H₂O snow on top of CO₂ snow can substantially decrease the surface albedo. Whereas 2.5 cm of H₂O snow (with 100 μm grain-size) can completely mask the impact of underlying CO₂ ice or the surface, roughly twice as much overlying CO₂ snow is required to mask underlying H₂O snow. Similarly, a 10% mixing ratio of H₂O ice embedded in CO₂ snow decreases broadband albedo by 0.18 (relative to pure CO₂ snow), while 10% CO₂ ice elevates H₂O snow broadband albedo by 0.10. We also achieve a reasonable agreement between modeled and observed Martian spectral surface albedo in the near-IR.

In Chapter 4, we successfully integrate the extended SNICAR model into the current version of MGCM to calculate real-time interactive surface broadband snow albedo. We use an initial dust content to constrain the surface albedo within realistic boundaries for Mars, and an exponential decay method to release surface dust over time. The SNICAR integrated model simulates a change of 0.47 W/m² in shortwave flux as compared to original MGCM. Real snow albedos are much lower than the dust-free albedos due to high

amount of dust deposition included with SNICAR implementation, and higher susceptibility of CO₂ snow to darkening as compared to H₂O snow. We estimate that the impact caused by dust is about 1.5 times (absolute values) of the impact caused by the presence of snow in shortwave flux at surface. Runs with different initial conditions for various years show higher correlation between albedo variation and dust opacity with SNICAR integrated model.

5.2 Future Work

This thesis opens multiple research topics which I plan to explore in future. In Chapter 4, we observe some technical issues occurring with SNICAR integrated MGCM for some specific initial conditions for long term runs of the model. These issues are most likely caused by erroneous dust deposition simulations on the surface within the model. I plan to resolve these issues as soon as possible so that we can release SNICAR integrated MGCM to the public.

I would also be interested in improving other aspects of MGCM that are dependent on snow/ice coverage of Mars. Dust in Mars's climate system significantly impacts (both darkening and brightening) the planetary albedo. Large global dust storms put enough dust in the air to completely cover the planet and block out the sun, hence largely affecting the planetary energy exchange with Sun and outer space. Although MGCM captures several aspects of the current Mars climate, they require further improvements to better understand surface albedo and emissivity feedbacks and thus to accurately determine the energy balance of the planet at any given time. Therefore, I would like to investigate the coupling of the dust, CO₂, and water cycles through surface albedo and emissivity feedbacks using variability of albedo feedback, radiative cooling, and adiabatic cooling caused by presence

of dust. Additionally, I would also like to explore the impact of these processes on present-day climate system of Mars. Improving planetary climate model is the key to better understand the evolution of, and to predict the future events of any planet.

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