

## Multispacecraft Observations and Modeling of the June 22/23, 2015 Geomagnetic Storm

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### Key Points:

- MHD models can reproduce well the dipolarizations seen at MMS and VAP. Space weather forecasting can predict Kp variations within 0.5 step.
- Beams of O<sup>+</sup> flowing downstream appear to cross the separatrix and become a second energized population of the tail plasma sheet.
- MHD models successfully reproduced the polar cap convection patterns and cross-polar cap potential drops for a range of IMF conditions.

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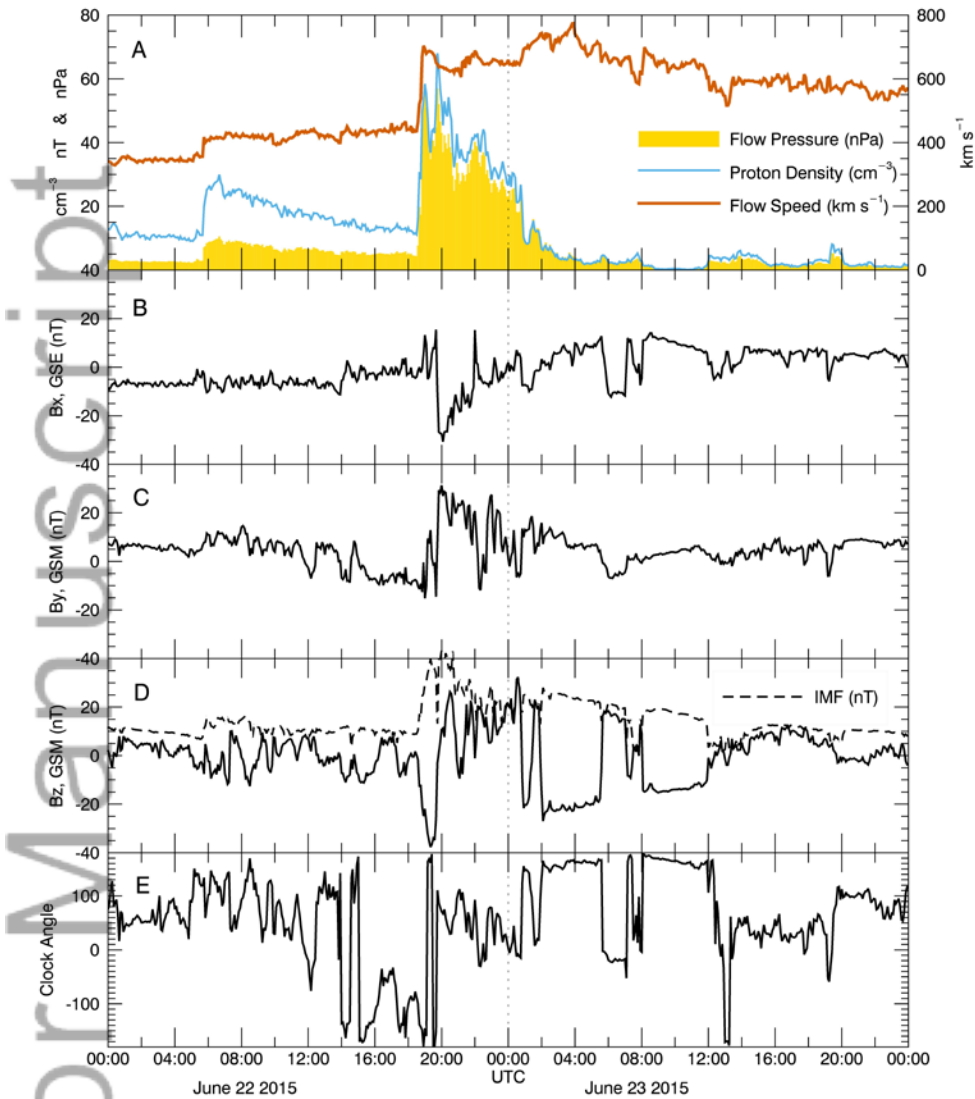
## Abstract

The magnetic storm of June 22-23, 2015 was one of the largest in the current solar cycle. We present in situ observations from the Magnetospheric Multiscale Mission (MMS) and the Van Allen Probes (VAP) in the magnetotail, field-aligned currents from AMPERE, and ionospheric flow data from DMSP. Our real-time space weather alert system sent out a “red alert”, correctly predicting Kp indices greater than 8. We show strong outflow of ionospheric Oxygen, dipolarizations in the MMS magnetometer data, and dropouts in the particle fluxes seen by the MMS FPI instrument suite. At ionospheric altitudes, the AMPERE data show highly variable currents exceeding 20 MA. We present numerical simulations with the BATS-R-US global magnetohydrodynamic (MHD) model linked with the Rice Convection Model (RCM). The model predicted the magnitude of the dipolarizations, and varying polar cap convection patterns, which were confirmed by DMSP measurements.

## 1 Introduction

On June 22, one moderate and one giant coronal mass ejection (CME) passed the ACE spacecraft at 04:51 and 17:58 UT respectively. The larger shock was observed by the MAG instrument [Smith *et al.*, 1998] as a jump in the IMF from about 10 to over 40 nT and by the SWEPAM instrument [McComas *et al.*, 1998] as a jump in solar wind density from 20 to over 45 particles/cm<sup>3</sup>, with a corresponding increase in pressure to over 50 nPa. When propagated to the bowshock, it was forecast to impact at about 18:36 UT after a smaller shock at about 05:40 UT (Figure 1A). Coupled with a strong southward IMF (Figure 1D), the Boyle Index reached nearly 500 kV, prompting our forecast system [<http://mms.rice.edu/realtime/forecast.html>] to send out a “yellow alert” at 06:04 and a “red alert” at 18:34, even before the CME impacted the bow shock. The y-component of the IMF was very strong, with the IMF clock angle (Figure 1E), rotating anticlockwise then clockwise nearly 360 degrees. The neural network forecaster [Bala and Reiff, 2012] predicted Kp of over 8 and the Kp forecast status went to “red” at 19:02.

The magnetospheric flotilla of spacecraft included MMS, VAP, Themis, and Cluster in the magnetosphere, plus AMPERE, DMSP, and ISS at low altitudes. This paper will concentrate on the large-scale features of the activity. It will include selected MMS, VAP, AMPERE and DMSP results and compare to the BATS-R-US model simulations. Other papers in this issue (e.g. Nakamura *et al.*; Baker *et al.*) will focus on other aspects of this event.



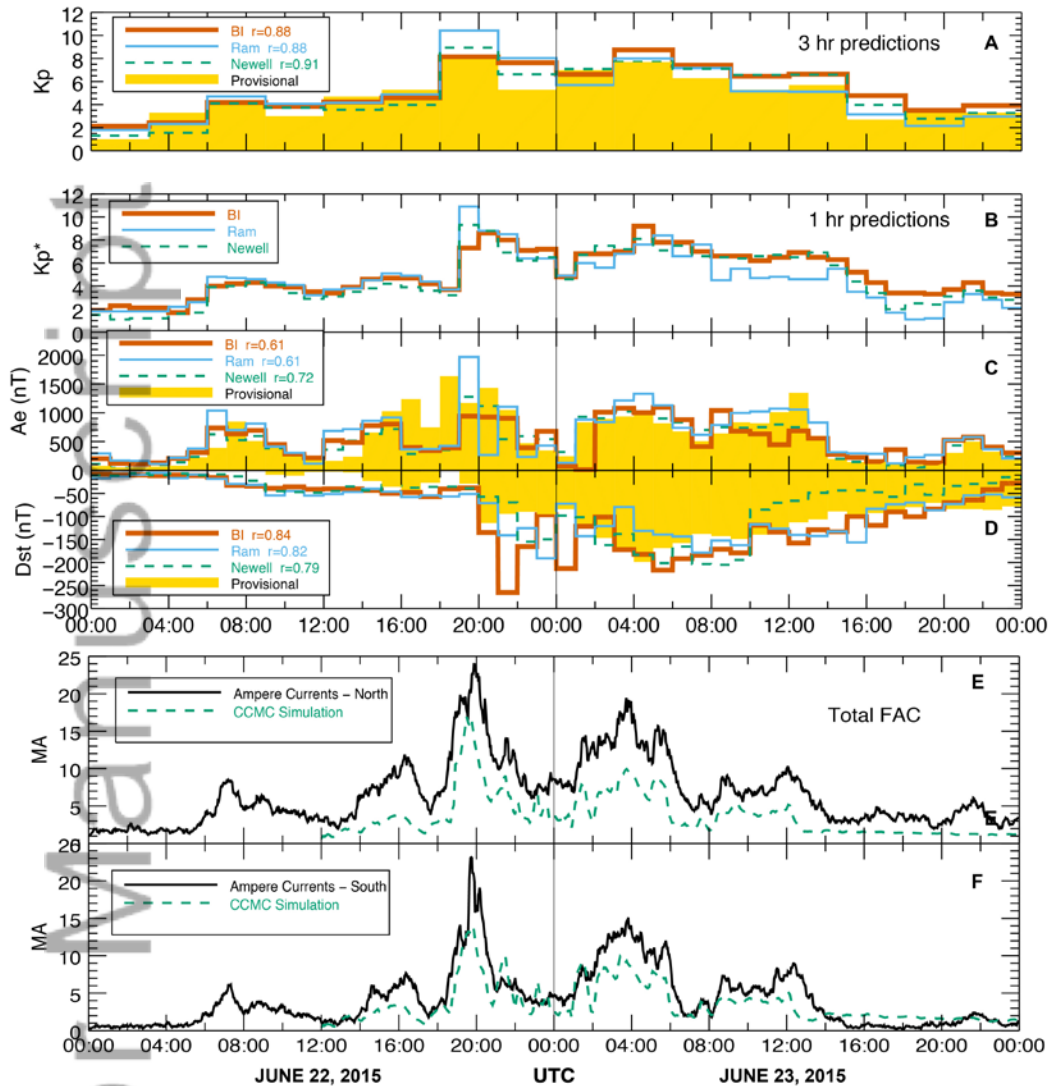
**Figure 1.** Interplanetary conditions during the event. A: solar wind density, velocity, and flow pressure; B, C, D: IMF Bx, By, Bz (and |B| shown dashed) components (GSM); E: IMF clock angle (0 = +Bz; 90 = +By direction). The values are plotted from OMNI data, propagated to the bow shock.

## 2 Space Weather Forecasting

The Rice University space forecast system [<http://mms.rice.edu/realtime/forecast.html>] predicts the Kp, AE, and Dst indices for one and three hours ahead of real time, at a cadence of 15 m. The

prediction is based on a neural network forecast that uses one or more base functions and look-back data of up to 24 hours. Yellow alerts are sent out if the forecast Kp is greater than 4, and red alerts if the forecast Kp is 6 or greater [Bala and Reiff, 2012]. Statistically, the accuracy of the predicted Kp is approximately one unit. Three base formulae are used in the forecast: the Boyle Index (marked as BI), the Boyle Index with a ram pressure term (RAM), and the Newell function (Newell) [Bala and Reiff, 2014]. All have similar prediction efficiencies, and the three predictions generally span the actual measured quantities. For this event, all three formulae gave similar predictions, with the AE index least well predicted ( $r = 0.61$  to  $0.72$ ), whereas Kp and Dst were successful with  $r = 0.79$  to  $0.91$  (Figure 2).

A comment on timing of “predictions” is in order. The forecast algorithm was trained on one-hour averages of AE, Dst, and Kp\*, with the Kp\* values being an overselection of Kp (because it is intrinsically a 3-hour average). The algorithm predicts the one-hour AE, Dst and Kp\* values for the hour following, based on the solar wind data of the previous hour. The calculation is done at the top of the hour after the realtime solar wind data are received, and a forecast typically posted at 6 minutes past the hour for the hour of the forecast (e.g. 18:06 for the hour 18-19). In this case, because the shock hit at the very end of hour 17, the AE predictions for hour 18-19 were not as high as the provisional AE turned out to be (Figure 2C), because the forecast only included two minutes of post-shock data, and because the high solar wind speed meant that the shock arrived at earth well before the end of hour 18. The forecast 3-hour Kp, based on three-hour averages of the one-hour predictions (Figure 2A), does show a very good fit ( $r = .88$  to  $.91$ , implying an accuracy of prediction from 0.36 to 0.47 step in Kp).



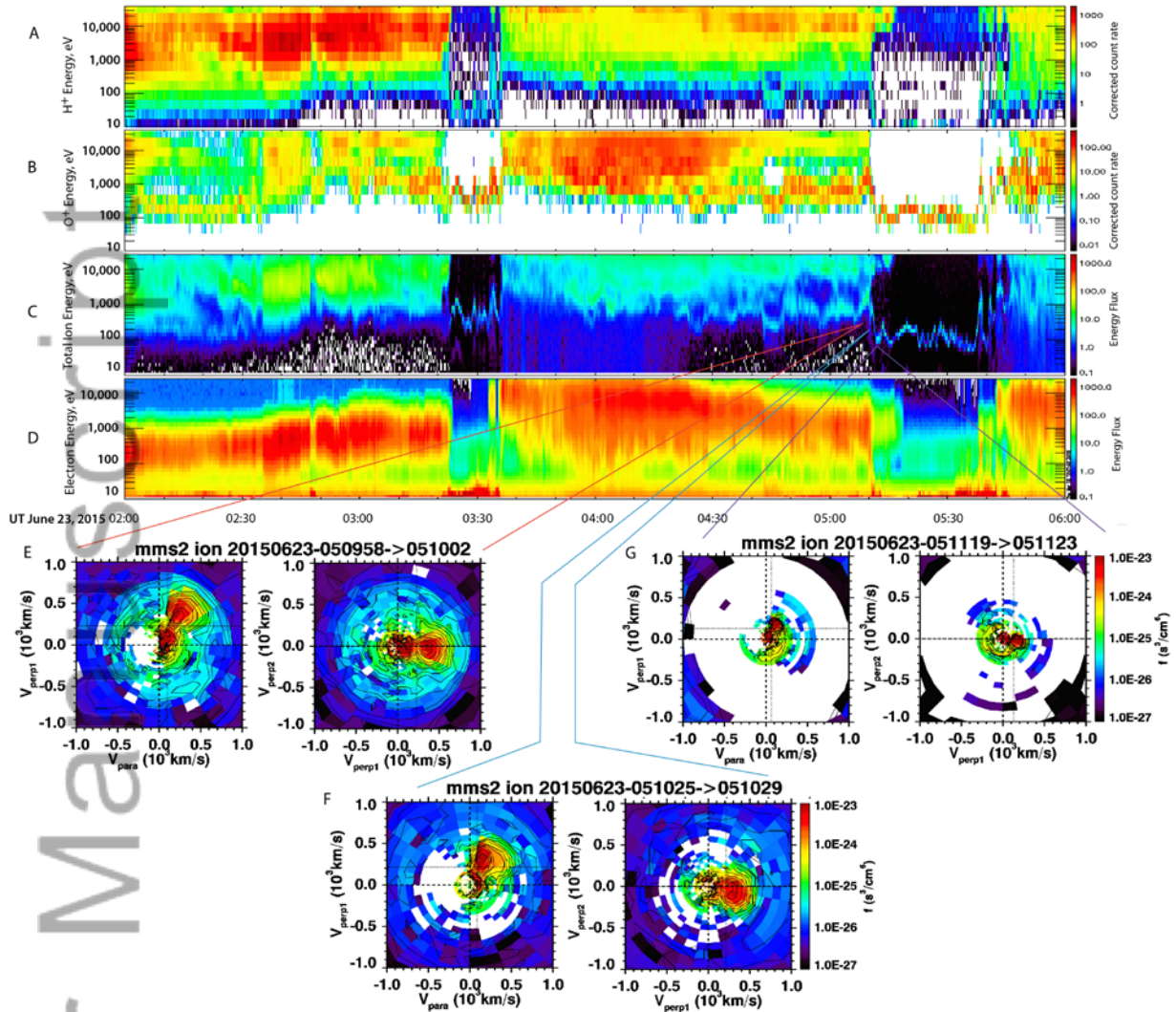
**Figure 2.** Comparison of predicted and modeled values compared to data. A shows the predicted 3-hour Kp index for June 22-23, 2015, compared to the provisional Kp, with a 90% correlation coefficient. Predicted (histograms) and observed provisional (yellow) Kp\* (panel B), AE (panel C), and Dst (panel D) indices. All three base functions: BI (thick line), Ram (thin line) and Newell (dashed) performed well, some slightly better than others. Modeled integrated field aligned currents (dashed) versus those calculated from AMPERE measurements for the Northern (panel E) and Southern (panel F) hemispheres.

### 3 Magnetotail Observations

Since the MMS spacecraft suite [Burch *et al.*, 2015] was in the “commissioning” phase of its mission, not all of the instruments were fully operational. All of the magnetometers [Russell *et*

*al.*, 2015] and many of the particle detectors were making measurements and saw the dynamic response of the magnetosphere. The Fast Plasma Instrument (FPI) onboard MMS2 [Pollock *et al.*, 2015] saw a number of excursions between the plasma sheet and the lobe, shown in Figure 3 as particle flux dropouts at about 3:20-3:30 and 5:11-05:40 UT (Panels C and D). Clearly visible in the lobe are antisunward flowing ions, at a few hundred eV (Panels B and C). The Hot Plasma Composition Analyzer (HPCA) onboard MMS1 [Young *et al.*, 2015] identified these beams as  $O^+$  (Panel B) presumably from auroral outflow [Lu *et al.*, 1992] or possibly from the dayside cusp [Liao *et al.*, 2010] though cusp ions will more likely reach the neutral sheet farther downtail [Liao *et al.*, 2012]. The HPCA instrument, because of a new AC sweep field, can reduce background  $H^+$  fluxes by nearly two orders of magnitude, making the determination of the heavy ion species much more reliable. This cool lobe ion beam appears to be accelerated and heated as it crosses into the plasmashet. In Figure 3, Panels E- G, we show sample total-ion particle distributions from FPI MMS2 as the spacecraft exits the plasmashet into the lobe around 05:10:30. In the lobe around 05:11:19 (panel G), we see two separated cold ion beams: one with a parallel velocity of around 50 km/s and another with a parallel velocity of  $\sim 200$  km/s. The right image of that pair shows that those beams are also convecting at about the same  $\mathbf{ExB}$  speed ( $v_{\text{perp1}}$ ) as the parallel speeds, with no drift in the other direction ( $v_{\text{perp2}}$ ) perpendicular to  $\mathbf{B}$ . From the HPCA data (Panels A and B), it is clear that the less energetic ions are  $H^+$  and the ones with apparently higher velocity are just  $O^+$  ions with the same parallel and perpendicular velocity as the  $H^+$  ions. This is confirmed by the HPCA particle distribution plots (not shown). As the spacecraft exits the plasmashet boundary layer at around 05:10:30 (Panel F), we see the same two beams, but now with a  $\mathbf{ExB}$  velocity double that in the lobe. This middle distribution does show some  $v_{\text{perp2}}$ , but it is not clear whether this is a time aliasing as the fields change direction during the four-second measurement of the distribution function. The distribution function measured deepest in the plasmashet (Panel E, at 05:09:58), shows the highest parallel and  $\mathbf{ExB}$  drift speeds, and no significant  $v_{\text{perp2}}$ . For additional information on this event see Nakamura *et al.* (this volume).

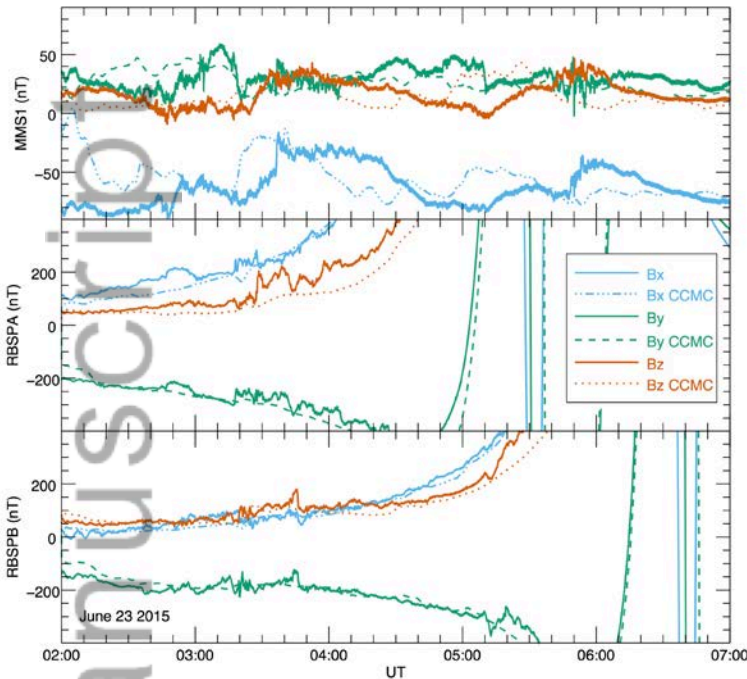
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**Figure 3.** HPCA (MMS1) and FPI (MMS2) measurements in the magnetotail from 0200 to 0600 of June 23. Energy spectrograms of the HPCA  $H^+$  (Panel A) and  $O^+$  (Panel B) corrected counts are shown in the top two panels with Panel B showing an  $O^+$  beam from the ionosphere merging into the plasma sheet at each lobe/plasmasheet transition. FPI ion and electron spectrograms of differential energy flux are shown respectively in Panels C and D. The bottom three panels (E, F and G) show three pairs of FPI distribution functions as the spacecraft exited the plasmasheet to the lobe. The left of each pair shows  $v_{\text{parallel}}$  and  $v_{\text{perp1}}$  (along  $\mathbf{ExB}$ ) components of the particle distribution functions, and the right of each pair shows the two perpendicular velocities,  $v_{\text{perp1}}$  and  $v_{\text{perp2}}$ .

At the same time, the magnetic fields observed at MMS in the tail and by VAP closer to the Earth showed dramatic dipolarizations as the magnetotail responded to the Northward turnings of the IMF (Figure 4). At MMS-1 (top), the measured  $B_x$  went from -90 to -20nT, and  $B_z$

increased from near zero to 30nT. The transitions from MMS from the plasmashet to the lobe resulted primarily because of the thinning and expansion of the plasmashet and partially because of the flapping of the magnetotail current sheet up and down.



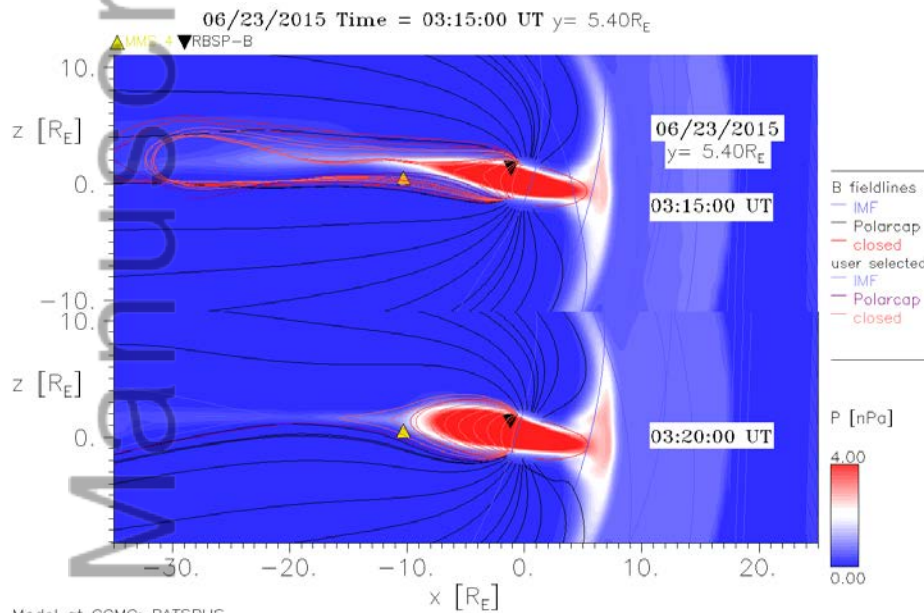
**Figure 4.** Tail magnetic fields observed by MMS1 (top) and Van Allen Probes A (middle) and B (bottom) (solid lines) and modeled (dashed lines) using a high-resolution run of BATS-R-US from the CCMC.

## 4 Modeling Results

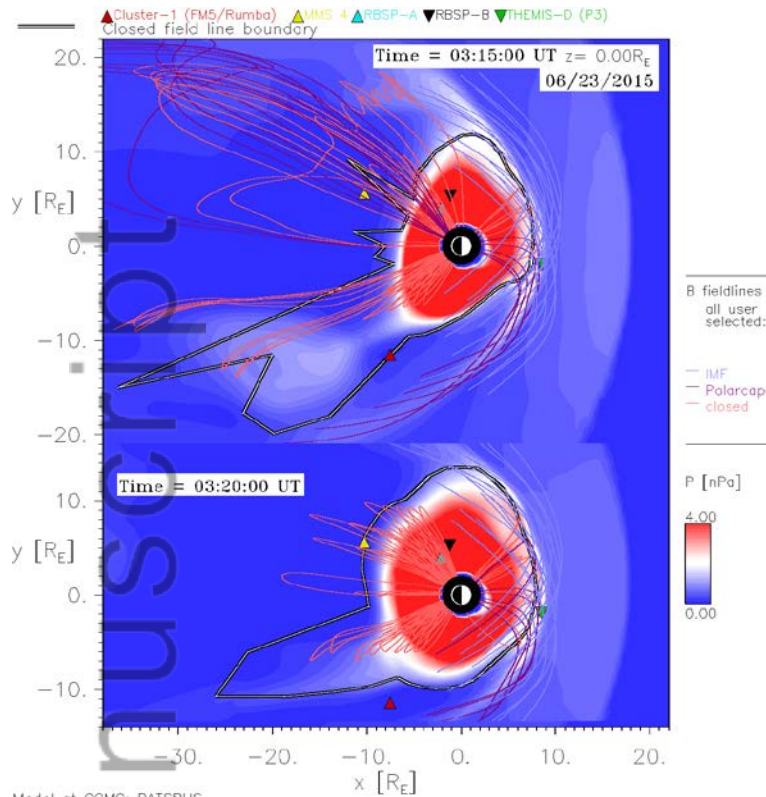
To put these observations into context we ran the BATS-R-US model [Powell *et al.*, 1999; Tóth *et al.*, 2005, 2012] with RCM in highest resolution available from the Coordinated Community Modeling Center (CCMC) (5 minute cadence and 1/8 RE at the inner boundary), using the measured propagated solar wind and IMF from OMNIdata. The model showed many tail reconnection events in the time frame of 00:00 UT on June 22 to 24:00 UT on June 23. Movies showing various cuts of the model for this event, including a cut that dynamically followed the MMS trajectory, can be found in the Supplemental material and at <http://mms.rice.edu/June22>. The model did especially well at capturing the several stretching and rapid dipolarizations observed in the 2-6 UT June 23 time frame (Figure 4, dashed). The first dipolarization in the model was between 3:15 and 3:20 (Figure 5A) and was observed at MMS at 3:16 [Nakamura *et al.*, this volume]. The second dipolarization in the model occurred between 04:55 and 05:00, seen at MMS at ~5:10 (See the movie *62223MMScut.mp4* in the supporting information). The



magnitude of the field changes predicted along the MMS path were quantitatively accurate, but with some variances in timing, reaching peak values somewhat earlier than measured. The model suggests that at the times of the dipolarizations, an x-line is just downstream of MMS, which is very near the separatrix (Figure 5A). Since the tail was so stretched, a modest flapping of the tail affects the location of MMS relative to the model. The dipolarization signatures observed at VAP RBSP-A (middle panel of Figure 4) were accurate in location but low in amplitude. The dipolarization signatures observed at RBSP-B (lower panel of Figure 4) showed good agreement both in magnitude and location. Thus we argue that the  $O^+$  ions seen flowing downstream in the lobe near the edge of the plasmashet may be captured by the reconnection process and become the energetic  $O^+$  seen in the plasmashet.



**Figure 5A.** Two frames from a movie of the CCMC run, showing MMS 1 spacecraft located very near the separatrix of an x-line in the magnetotail. The dipolarization in the model is well observed between the stretched configuration of 03:15 and the more dipolar configuration of 03:20. The full movie including similar frames for the 05:00 dipolarization is in the supplementary materials and is available at <http://mms.rice.edu/June22/>.



Model at CCMC: BATSRUS

**Figure 5B.** Similar to 5a, but showing the equatorial plane locations of the various magnetospheric spacecraft (MMS, VAP-A and B, Themis, Cluster) at the time of the dipolarization. The field lines in both times were started at the same locations on the equatorial plane. Note how less stretched each field line is after the dipolarization. A movie of this view is provided in the supplementary materials and on our website <http://mms.rice.edu/June22>

When mapped to the ionosphere, the model showed extremely large field-aligned currents (FAC). As compared to the FAC's inferred from the AMPERE data [e.g. *Coxon et al., 2014*], which reached an integrated value of nearly 25 MA, the predicted magnitudes were 20-50% smaller than that observed, although the timing structure was quite accurate (Figure 2E-F).

## 5 Polar Cap Convection Patterns

Because of the large magnetic field strengths (~40 nT) and the large and variable IMF  $y$ -component, the model-predicted ionospheric convection varied from normal two-cell convection patterns during times of large negative  $B_z$  (Figure 6A), to a four-cell during high positive IMF  $B_z$  intervals (Figure 6B) to a single convection cell rotating in opposite directions (clockwise in the North, counter-clockwise in the South) in the two hemispheres after a long interval of strong  $+B_y$  (Figure 6C). A true single cell convection pattern is very rare: in this case the single cell was only predicted (and observed) in the North. A movie of the modeled convection patterns

62223\_ionosphere.mp4 is part of the supplemental information and posted at <http://mms.rice.edu/June22>.

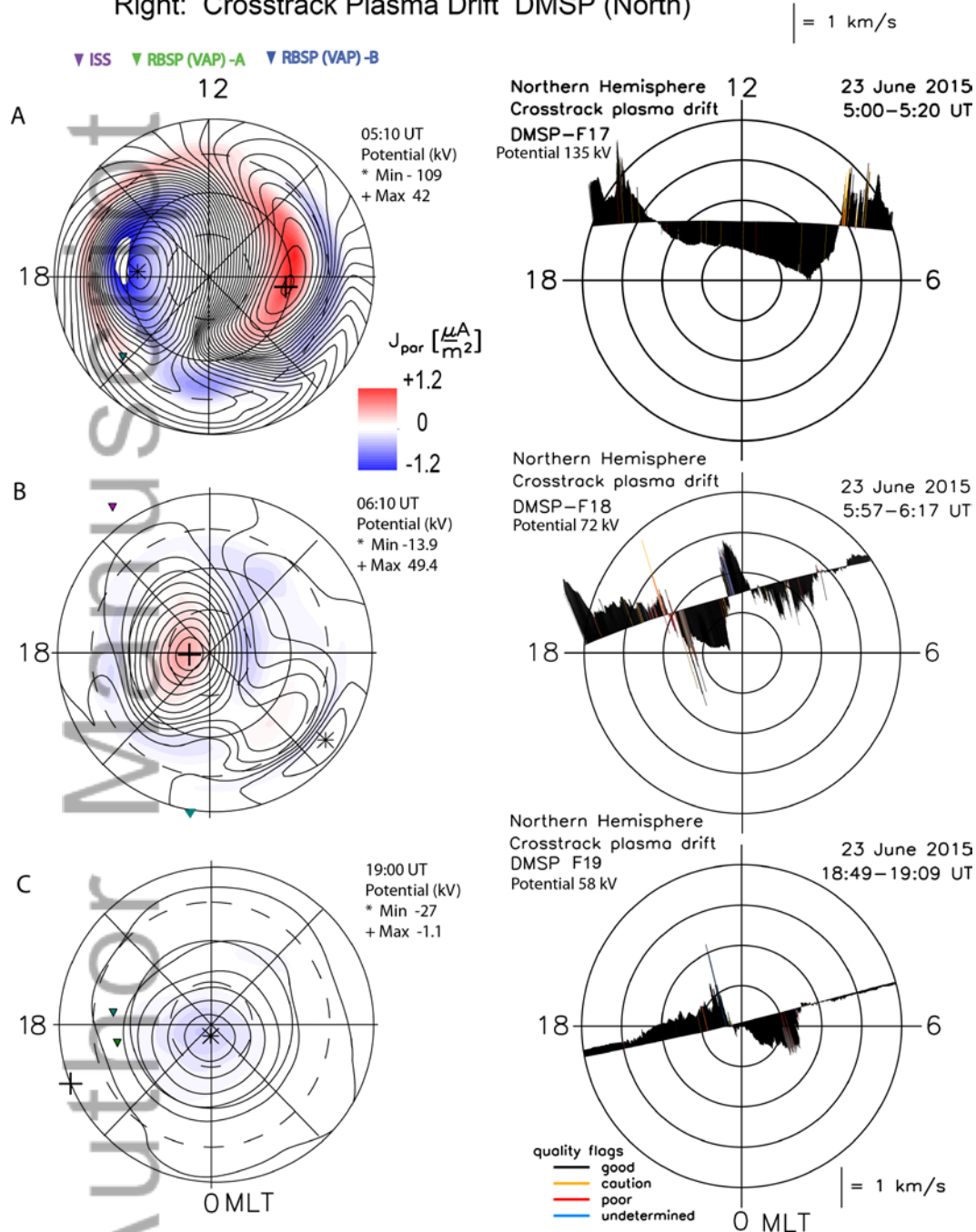
A comparison with DMSP plasma flow data from the same intervals confirms the basic features of the convection, but interesting differences are observed. For example, in the southward IMF case (6A), the polar cap is larger in the data than in the model by about 5 degrees. The sunward flow channel in the data during northward IMF (6B) is much narrower than in the model, which predicts a very large area of sunward flow in the polar cap. The model also missed the strong low-latitude sunward flows on the dusk side. The DMSP flow data confirm presence of a single-cell convection pattern in the North (Figure 6C, bottom right), with just a hint of a viscous cell at 7MLT, 75 ILAT. The cross-polar cap potential drops estimated from the model and from the flow data are comparable. Both the model and the data show the center of the cell just duskward of the magnetic pole. The lobe cell convection is more common in times of high sunward dipole tilt [Crooker, 1992], and this event which occurred on June 22- 23, had maximum sunward tilt in the North.

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### Ionospheric Convection Patterns June 23, 2015

Left: CCMC model (BATS-R-US)

Right: Crosstrack Plasma Drift DMSP (North)



**Figure 6.** Northern polar cap convection equipotentials predicted from the BATS-R-US run (left image in each row). Colored areas are the field-aligned current densities, and equipotential lines are 4kV apart

in each plot. The movie is available at <http://mms.rice.edu/June22>. Right side: measured crosstrack plasma drifts from DMSP at the same time as the predicted potentials. Panel A is a normal two-cell convection pattern during southward IMF; Panel B is a four-cell pattern with reversed flow in the central polar cap during strong northward IMF; Panel C is a single-cell clockwise convection observed during strong positive Y-component of the IMF.

## 6 Conclusions

This event represents the first major storm of the new Heliospheric Flotilla era. With well-instrumented spacecraft strategically placed in the magnetosphere, and new computational models, our understanding of magnetospheric dynamics, especially its response during dramatic events such as the one presented in this paper, is taking a leap forward. Despite this being an anomalously intense event with large magnetic fields, the BATS-R-US model did an admirable job of reproducing the amount of field change during the dipolarizations, estimating the polar cap convection and currents, and the approximate location of MMS near the separatrix during a very dynamic magnetotail sequence.

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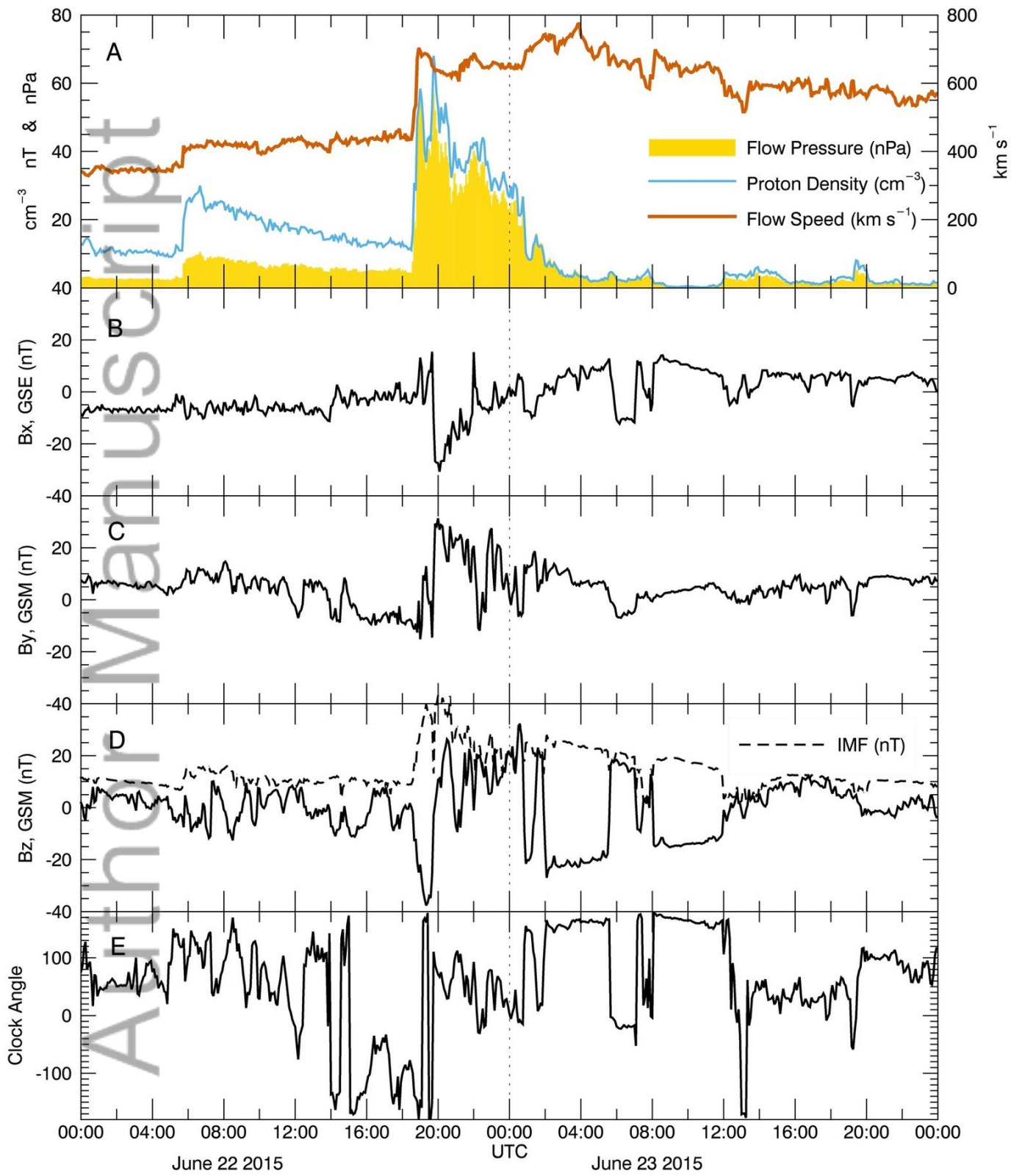
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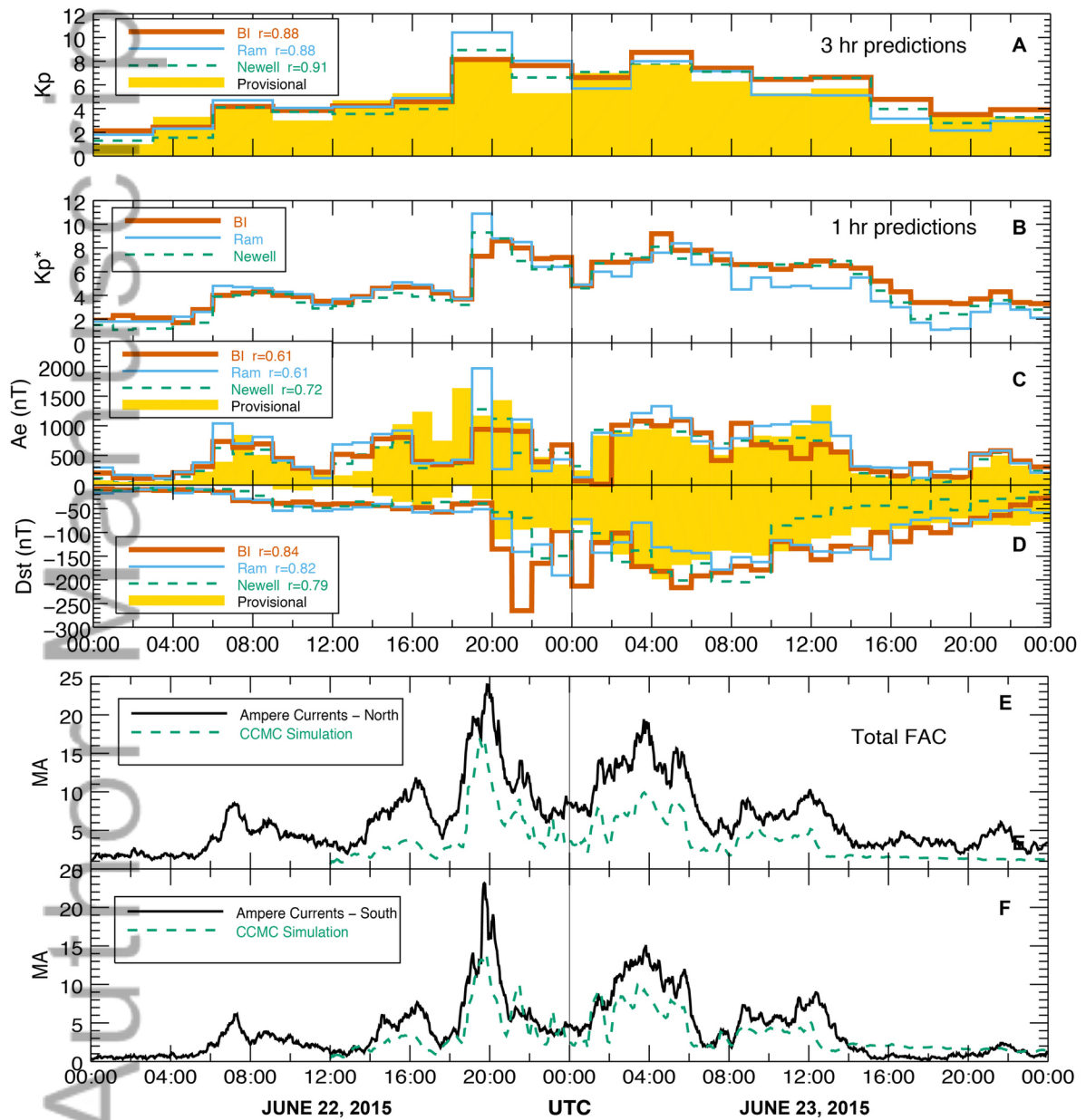
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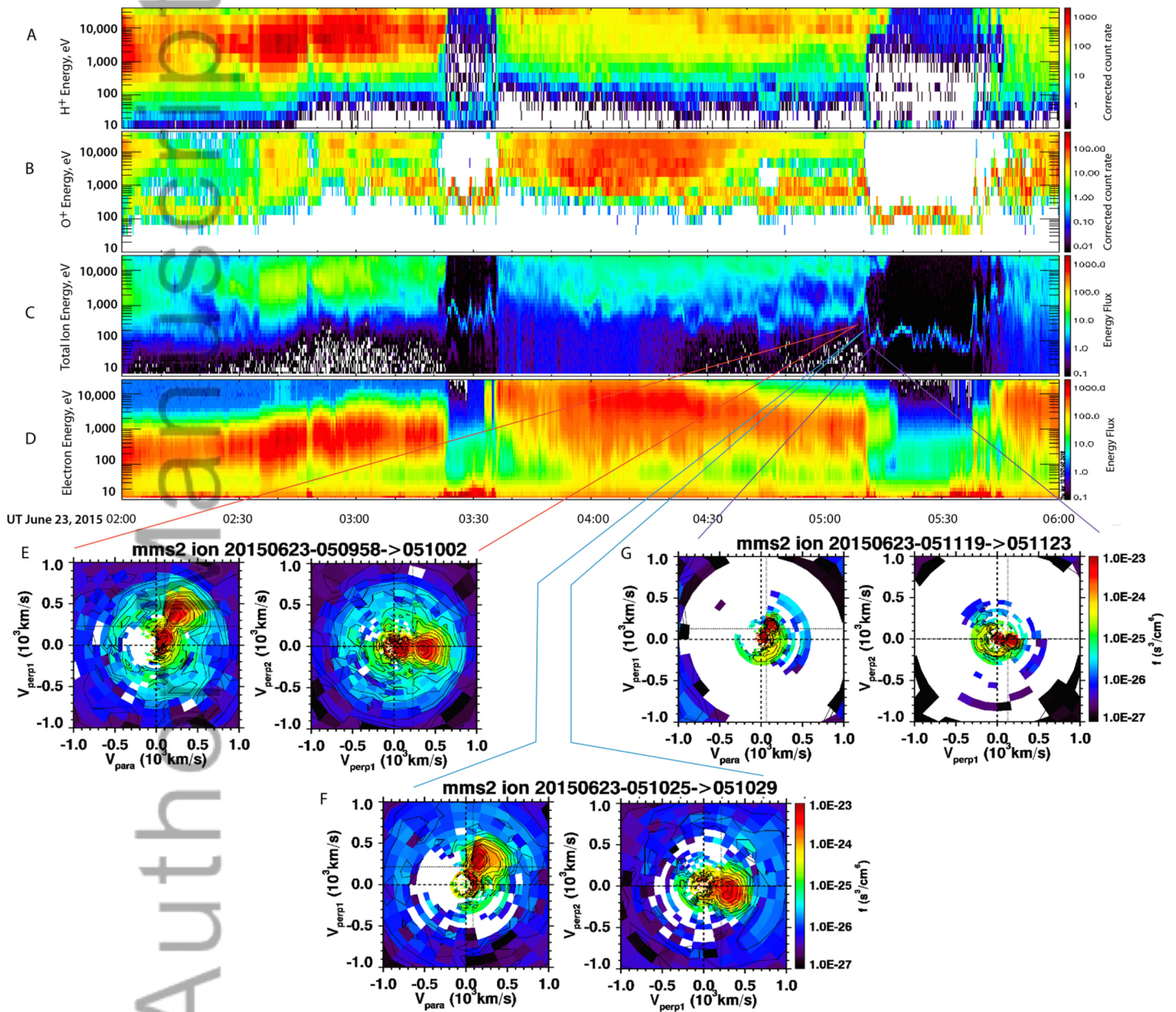


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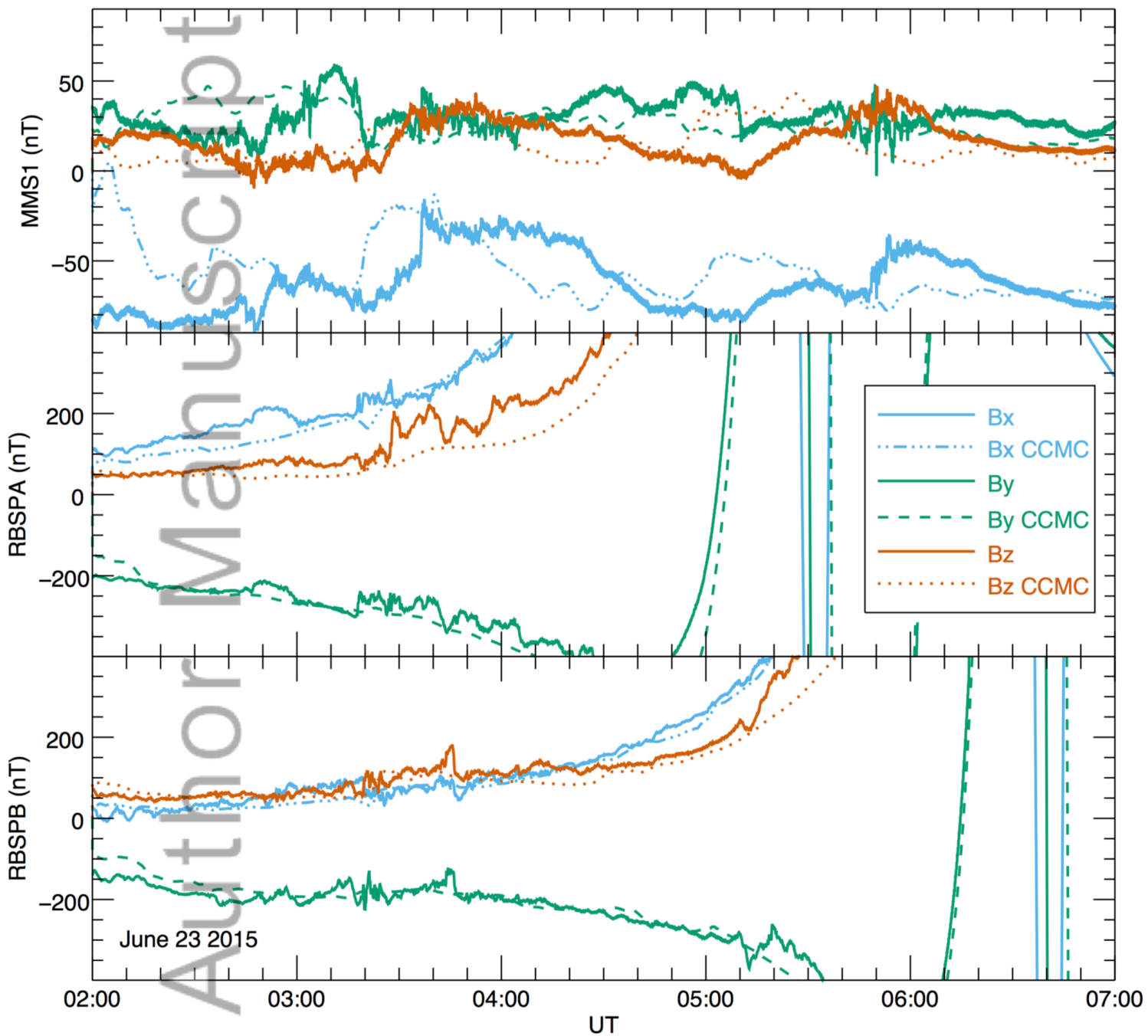


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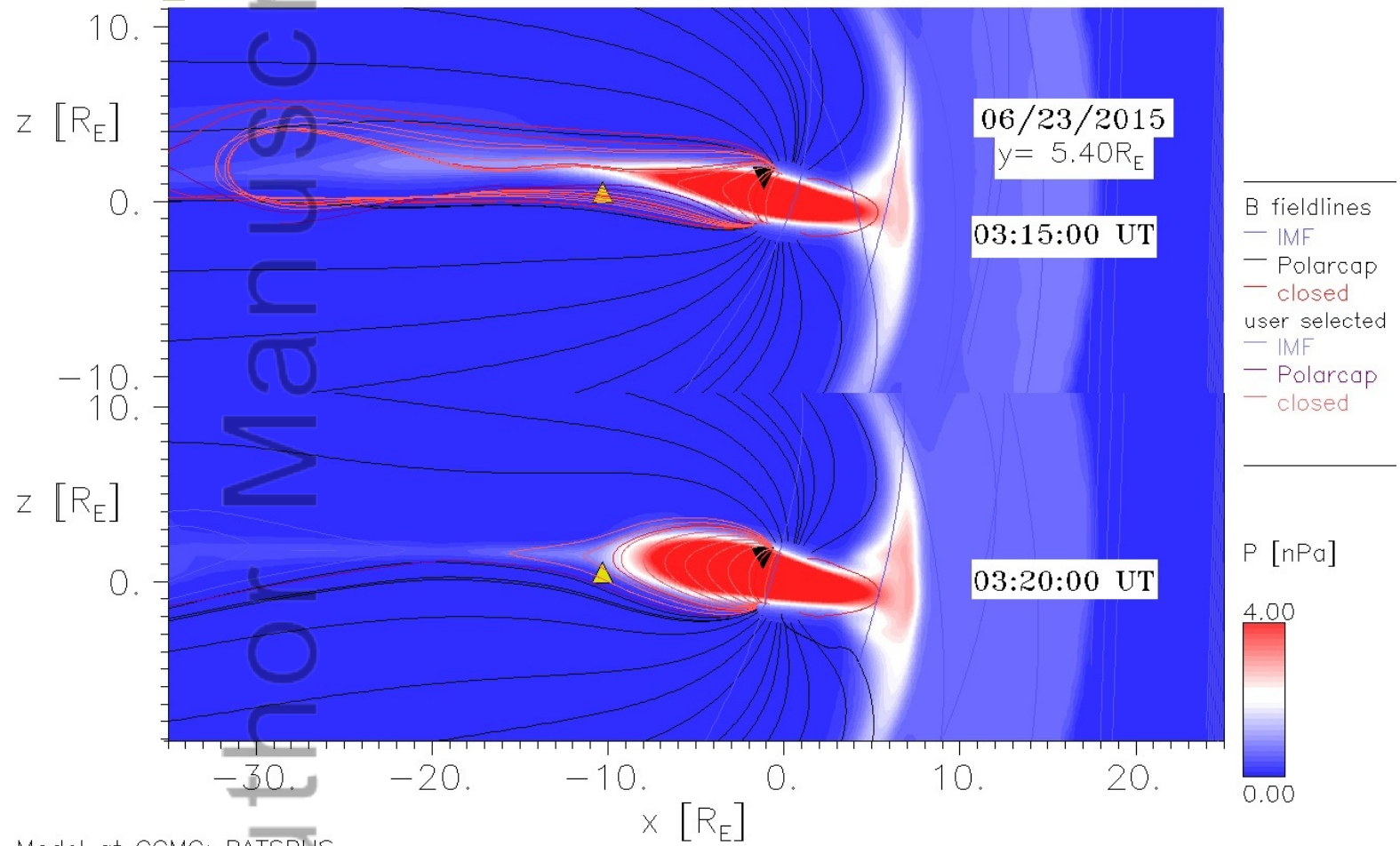


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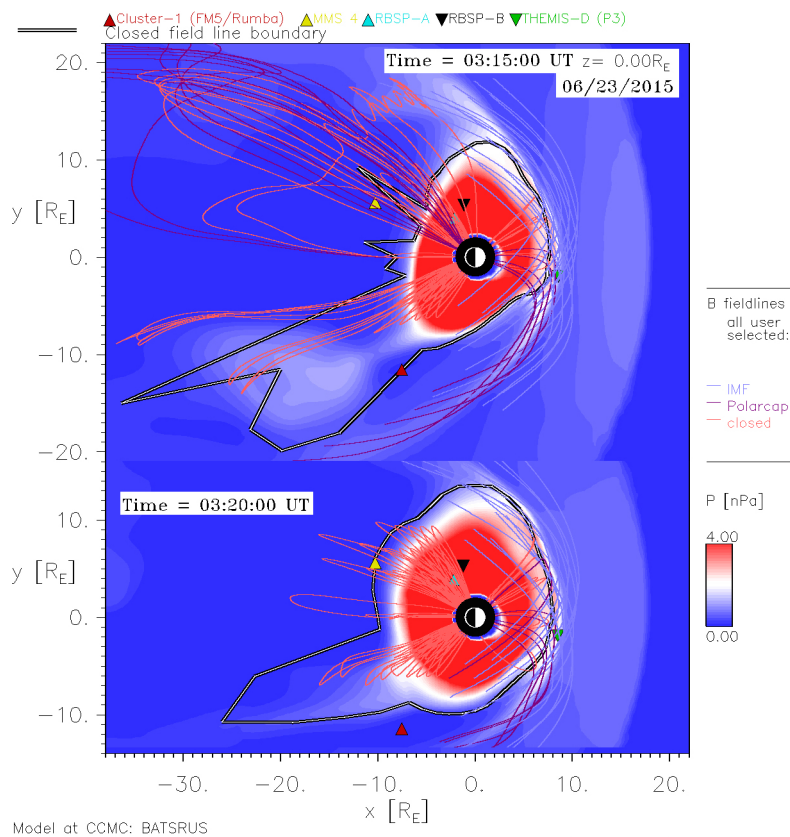
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Model at CCMC: BATSRUS

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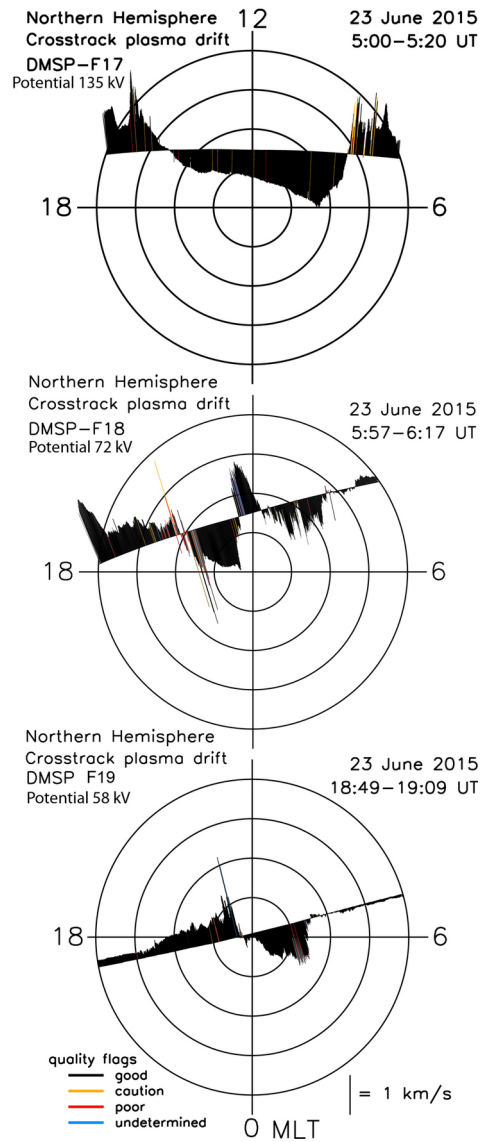
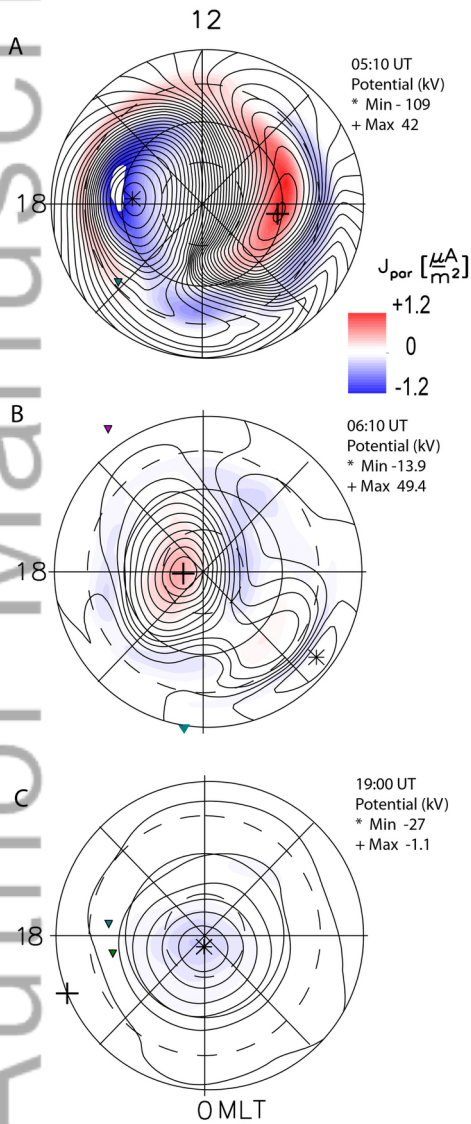
**Ionospheric Convection Patterns June 23, 2015**

Left: CCMC model (BATS-R-US)

Right: Crosstrack Plasma Drift DMSP (North)

| = 1 km/s

▼ ISS ▼ RBSP (VAP)-A ▼ RBSP (VAP)-B



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