

Received Date : 25-Feb-2016

Revised Date : 14-Jul-2016

Accepted Date : 16-Jul-2016

Article type : Article

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## **Structure-property relationships in aligned electrospun barium titanate nanofibers**

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**Abstract:** Barium titanate nanofibers were uniaxially aligned by electrospinning onto a rotating copper wire drum and alignment was maintained during calcination of the fibers. Two methods for maintaining alignment during calcination were tested, by either using carbon tape or a peeling off method to remove the aligned fibers from the mandrel followed by calcination. The carbon tape removal method led to the formation of shorter aligned nanowires while the peeling off method resulted in longer nanofibers. Additionally, the effects of calcination temperature and time on crystal structure were also examined. The degree of tetragonality in the barium titanate nanofibers increased at higher calcination temperatures and times. Piezoelectricity was confirmed in the nanofibers calcined using piezoresponse force microscopy (PFM), yielding a  $d_{33}$  value of 15.5 pm/V. Using the methods presented here, large quantities of aligned piezoelectric barium titanate and other ceramic fibers or wires can be produced to fulfill their demand in novel microelectronics.

This is the author manuscript accepted for publication and has undergone full peer review but has not been through the copyediting, typesetting, pagination and proofreading process, which may lead to differences between this version and the [Version of Record](#). Please cite this article as [doi: 10.1111/jace.14455](https://doi.org/10.1111/jace.14455)

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## 1. Introduction

Electrospinning is a versatile method for producing nano- to micro-meter sized fibers that are suitable for both rapid prototyping and large scale production.<sup>1</sup> Electrospinning has been used to synthesize a wide array of nanoscale materials including polymeric,<sup>2-4</sup> ceramic,<sup>5-8</sup> and even composite bi- and multi-phasic fibers and particles.<sup>9-12</sup> Fibers formed via electrospinning are promising candidates for use in many applications including filtration, catalysis, wound healing, and drug transport and release.<sup>1</sup> Traditionally, electrospinning results in the formation of randomly oriented fibers. However, fiber alignment can be achieved by tailoring the geometry of the counter electrode.<sup>4, 13, 14</sup> In some applications, electrospinning a large mat of randomly oriented fibers is sufficient, whereas for others alignment or a desired orientation in the electrospun fibers may be desired.<sup>4, 13, 14</sup> This paper demonstrates a route to produce large quantities of uniaxially aligned barium titanate nanofibers in a manner that can also be applied to a wider range of ceramic and metallic electrospun nanofibers.

The piezoelectric and ferroelectric properties and high dielectric constant of barium titanate have enabled its incorporation into a wide array of electrical applications including actuators,<sup>15</sup> medical imaging devices,<sup>16</sup> ultrasound transducers,<sup>17</sup> and high dielectric capacitors.<sup>18</sup> Nanoscale barium titanate has been considered a promising material for high density ferroelectric random access memories.<sup>14, 19-22</sup> Additionally, electrospun barium titanate nanofibers have been used to construct high performance humidity sensors in place of barium titanate thin films due to their high surface area and large aspect ratio.<sup>21, 22</sup> A key figure of merit for barium titanate nanowires is their piezoelectric coefficient,  $d$ , which is a measure of the magnitude of polarization that is generated in response to an applied load. In particular the  $d_{33}$  piezoelectric coefficient is often reported, which measures the piezoelectric response along the axis of the applied electric field.

Given the push towards incorporating barium titanate and similar perovskite nanofibers and nanowires into electrical devices, it is important to develop methods to synthesize large quantities of these nanofibers.<sup>14, 19, 20</sup> Control of their orientation is desirable since randomly oriented fibers may not be best suited for use in microelectronics. While well aligned barium titanate/polymer nanofibers have been previously synthesized with the use of two parallel electrodes separated by an insulating air gap,<sup>14, 20</sup> the area of fibers which can be produced with this method is limited by the distance between the electrodes and their length. The method is further limited by a loss in alignment as subsequent fiber layers are deposited.<sup>4</sup> This has been overcome in polymer systems by parallel electrodes spaced around a rotating mandrel.<sup>4</sup> However, for ceramic nanofibers compared to polymer fibers there is an additional challenge of maintaining alignment during calcination. The as-spun fibers are amorphous thus a calcination step is required to burn off the polymer binder and to form the final crystalline product.

This paper demonstrates the first example of alignment in ceramic nanofibers that persists during calcination by electrospinning onto a rotating mandrel. In this paper, two methods of maintaining fiber alignment throughout the calcination process are explored. The first method involves the removal of the as-spun fibers using carbon tape and the other involves the removal of sections of fibers by peeling them off the wire drum as sheets. In this work we also examine the effects of other processing variables, including calcining time and temperature on the structure and properties of the resultant nanofibers.

## **2. Experimental Process**

### **2.1. Materials**

Barium acetate, poly(vinyl pyrrolidone) (PVP, MW=1,300,000), acetic acid, and titanium isopropoxide ( $\geq 97.0\%$ ) were obtained from Sigma Aldrich.

### **2.2. Barium Titanate Sol-Gel Preparation**

The barium titanate precursor solution was prepared in a 1:1 molar ratio of barium to titanium by first dissolving 2.55 g barium acetate in 8 ml of acetic acid followed after 2 hours by the dropwise addition of 2.95 ml titanium isopropoxide.<sup>11, 19, 23</sup> A polymer solution was simultaneously prepared, consisting of 1.1 g poly(vinyl pyrrolidone) (PVP, MW=1,300,000) in 10 ml ethanol and was also stirred for 2 hours. Then the barium titanate solution was added to the polymer solution and allowed to stir for 10 minutes before electrospinning.

### **2.3. Electrospinning Aligned Barium Titanate Fibers**

To electrospin aligned nanofibers a counter electrode consisting of parallel copper wires was constructed on a rotating mandrel (Figure 1). The copper wires had a separation of 1 cm and the drum was rotated by a DC electric motor at a rate of 21 rpm, determined by the rotational speed of the DC motor and the pulley diameters.

For electrospinning, the barium titanate precursor solution was fed through a stainless steel needle (20 gauge) with a flow rate of 2 mL/h. The needle was connected to a high voltage power supply, with an applied voltage of 13kV, and positioned normal to the surface of the electrically grounded rotating copper wire drum, with a separation distance of 16 cm. After electrospinning, two methods were tested for the removal of as-spun fibers from the copper wire mandrel for calcination: removal by peeling the fibers off as sheets or by adhering the aligned fibers to a piece of carbon tape. The as-spun fibers were calcined at 750°C, 875°C, and 1000°C for 2, 4, or 6 hours.

### **2.4. Characterization**

The viscosity of five barium titanate precursor solutions were measured using a Brookfield DV-II+ Pro rotational viscometer with a spindle speed of 20 rpm. The average viscosity of the as-prepared precursor solutions was 137 cP with a standard deviation of 14 cP.

The nanofiber samples were imaged with a FEI Phillips XL40 FEG scanning electron microscope (SEM). Fiber diameters were measured using ImageJ, where diameter measurements were taken on at least 150 fibers per sample. The measured fiber diameters were then fit to multiple normal distributions using normalmixEM<sup>24</sup> in Rstudio. A fit for multiple log normal distributions was found by first taking the log transform the diameter measurements, fitting the data with normal mixEM, and converting the normal fits of the log transformed data back into log normal distributions. The crystal structure of the nanofibers was determined via X-ray diffraction (PANalytical X'Pert Powder) with subsequent Rietveld refinement in PANalytical Highscore Plus. Additional crystallographic structure information was obtained using Raman spectroscopy; Raman spectra were obtained using a Renishaw Invia Raman microscope with a 633 nm laser and a 20× objective lens.

Piezoelectric force microscopy (PFM) was performed using a Park Systems XE 70 atomic force microscope modified with a function generator (Agilent, 33210A). PFM samples were prepared by electrospinning directly onto platinum coated silicon wafers with a layer of titanium dioxide between the platinum and silicon to promote adhesion of the thin film, after which they were calcined for 2h at 750°C. The  $d_{33}$  value of a barium titanate nanofiber calcined at 750°C for 2h was obtained via linear regression on the positive voltage sweep using 20 cycles.

### 3. Results and Discussion

After electrospinning onto the copper wire drum, scanning electron microscope (SEM) images were obtained to verify alignment and to observe the morphology of the as-spun fibers (Figure 2a). After the alignment of the as-spun fibers was verified, the fibers were removed from the copper wire mandrel using two different methods. In the first method fibers were removed by adhering them to carbon tape. The carbon tape was then placed on a ceramic dish, and was calcined in air at 750°C for 2 hours, resulting in the oxidation and removal of the carbon tape from the fibers (Figure 2b). The adhesive on the carbon tape restricted the shrinking of the fiber as the PVP was removed, leading to the fibers being pulled apart lengthwise, breaking to form aligned nanowires instead of longer nanofibers (Figure 2b-c).

The second method of maintaining alignment during calcination involved removing the fibers by peeling off sections of sheets from the copper wire drum. Using this method, the fibers were free to shrink along their length during calcination. Figures 2c-d show that this sheet removal method maintained alignment and resulted in uniaxially aligned nanofibers, after calcination at 750°C and 1000°C, respectively.

Figures 3a and 3b show the fiber diameters distributions for both calcined and uncalcined fibers, respectively, revealing a bimodal distribution of fiber diameters. This multimodal distribution has been previously observed in both polymer<sup>25, 26</sup> and sol gel electrospinning<sup>27</sup> and has been attributed to branching, the formation of secondary polymer/sol gel jets from the primary jet, during the

electrospinning process. This branching phenomena is due to an instability of the circular cross section of the fiber at high potentials. This instability causes the formation of cusps from which secondary jets can be emitted.<sup>25</sup> Though previously such multimodal fiber distributions have been modeled using multiple normal distributions,<sup>26</sup> we found that the diameters of the electrospun fibers were better fit by multiple log normal distributions. The origin of this log normal fiber diameter distribution can be related to solution inhomogeneities that are present during electrospinning. In electrospinning a polymer or sol-gel solution is placed in a syringe that is connected to a high voltage power supply, which leads to the formation of an elongated droplet at the syringe tip, referred to as the Taylor Cone. The radius of the Taylor cone has been shown to be a determining factor of the droplet size in electrospraying,<sup>28</sup> and it has also been predicted that it is related to the final jet diameter for electrospinning.<sup>29</sup> Additionally, the radius of the Taylor cone is proportional to the viscosity of the fluid,<sup>30</sup> where the viscosity of the polymer is related exponentially to its concentration, as described by the Martin equation, which varies randomly throughout the solution.<sup>31</sup> These random variations then give rise to a viscosity that is log normally distributed, resulting in a Taylor cone radius and fiber diameters that also exhibit a log normal distribution (Figure 3).

The fiber diameter distributions from both normal and log normal models were graphically compared to the experimental cumulative distribution (Figure 3c-d). The fit of each model distribution was also quantitatively evaluated using the Kolmogorov-Smirnov (K-S) test to determine if the model distribution is a reasonable fit for the experimental results.<sup>32</sup> Based on the K-S test, the multiple normal distributions could be rejected for both the uncalcined ( $p = 1.96 \times 10^{-4}$ ) and calcined ( $p = 2.60 \times 10^{-14}$ ) fibers. The multiple log normal distributions could not be rejected with the K-S test with a 95% confidence level, thus the multimodal log normal distribution model was accepted. The multiple log normal distribution for the uncalcined fibers can be defined as follows: the first distribution had a weight of 0.58, a mean diameter of 232 nm +/- 117 nm, while the second had a weight of 0.420, a mean of diameter of 463 nm +/- 58 nm. For the calcined multiple log normal distribution the first distribution had a weight of 0.447, a mean diameter of 180 nm +/- 98 nm, while the second had a weight of 0.553, a mean diameter of 302 nm +/- 47 nm.

The crystallographic properties of the calcined nanofibers were analyzed using X-ray diffraction (XRD). The XRD spectra are shown in Figure 4, revealing peaks characteristic of perovskite barium titanate.<sup>14, 19, 33</sup> Rietveld refinement indicated the presence of the tetragonal perovskite structure for all calcination conditions. However, Figure 4 shows a lack of observable peak splitting in the (200) peak at 45°, which can be attributed to peak broadening in the nanocrystalline samples.<sup>33-35</sup> The Rietveld refinement also indicated the presence of a cubic and hexagonal phase of barium titanate, in addition to the presence of barium carbonate (witherite). The metastable hexagonal phase has been previously shown to coexist with the tetragonal phase of barium titanate in nanoparticles. The phase transformation between

the hexagonal and tetragonal phases is a reconstructive phase transformation, and thus characterized by sluggish kinetics.<sup>36</sup> Therefore, as anticipated the amount of hexagonal phase present in the sample decreases with increasing calcination temperature, along with a corresponding increase in the amount of tetragonal phase present, when calcination time is kept constant (Figure 5). It is important to note, that all samples had a barium carbonate impurity phase, which varied from 10-14 wt. % depending on the calcination temperature. Across all calcination times temperature was negatively ( $R=-0.621$ ) and significantly ( $p=0.0005$ ) correlated with the wt.% of the intermediate phase as more of the recombination transformation occurred. Furthermore the percent of the tetragonal phase was positively ( $R=0.522$ ) and significantly ( $p=0.005$ ) correlated with increasing calcination temperature and attributed to the more complete transformation of the metastable hexagonal phase at higher temperatures, where more energy was available to promote the kinetically sluggish transformation. Similarly, longer calcination times also corresponded with an increasing wt.% of the tetragonal phase, and a decrease in the amount of hexagonal phase present. However, we observed that calcination time has a lesser impact on the phases present compared with temperature for the temperatures and times tested. Lastly, both cubic barium titanate and barium carbonate decreased in weight percent with increasing calcination times and temperatures. Rietveld refinement results are given in Table 1, where three samples were synthesized and analyzed under each calcination condition.  $R_{wp}$  is a measurement of the distance between the observed and fitted spectra,  $R_{exp}$  is essentially a measure of the data quality and essentially the expected  $R_{wp}$  for an ideal fit, and the reduced chi square statistic  $\chi_v^2 = \frac{R_{wt}^2}{R_{exp}^2}$  gives an indication of the quality of the Rietveld refinement on the spectra. For example a  $\chi_v^2 < 1$  indicates the spectra may be over fitted,  $\chi_v^2 \gg 1$  indicates a poor fit, and  $\chi_v^2 \approx 1$  signifies a good fit.

Raman spectroscopy was also employed to further confirm the presence and transformation between the metastable hexagonal and tetragonal phases of barium titanate. Figures 6a and 6b show the evolution of normalized Raman spectra as a function of calcination temperature and time, respectively. Peaks that are attributed to tetragonal (T),<sup>36-39</sup> cubic (C),<sup>40-42</sup> and hexagonal (H)<sup>36, 37, 43</sup> barium titanate as well as barium carbonate (witherite, W)<sup>43</sup> are labeled within the spectra.. Figure 6a reveals that the tetragonal (T) peak intensities increase with increasing temperature, alongside a corresponding decrease in the intensities of the metastable hexagonal (H) phase. However, Figure 6b shows the evolution of Raman spectra at calcination times of 2h, 4h, and 6h, for a constant calcination temperature of 875°C. It is apparent here, that the calcination temperature has a greater effect on the transformation from the hexagonal to the tetragonal phase, as also observed in the XRD data.

It is anticipated that barium titanate nanofibers with the highest fraction of the tetragonal phase will result in the best ferroelectric and piezoelectric properties. Therefore, the effects of calcination on the

grain size of the tetragonal barium titanate phase were also determined by relating peak broadening to the average grain size. Since strain can also contribute to peak broadening strain was also considered in the Rietveld refinement, however no significant strain contributions to peak broadening were found. Figure 7 shows a plot of the natural logarithm of grain size of the barium titanate nanofibers as a function of inverse temperature. From Figure 7 the activation energy of grain growth of tetragonal barium titanate was calculated to be  $2.641 \pm .377$  eV/(atom K). This fits in to a trend previously observed in barium titanate nanofibers by Peng et al. where the activation energy of grain growth was found to decrease with decreasing fiber diameter.<sup>45</sup>

Lastly, the piezoelectric properties of the electrospun barium titanate nanofibers were determined using piezoresponse force microscopy (PFM). Figure 8a shows the butterfly loop, with corresponding phase plot figure 8b demonstrating a  $180^\circ$  phase shift, obtained from PFM analysis on a barium titanate nanofiber calcined at  $750^\circ\text{C}$  for 2h, confirming piezoelectric behavior. The  $d_{33}$  value, extracted via linear regression on figure 8a was found to be 15.5 pm/V with a standard deviation of 0.1 pm/V, in a method established by Zhou et al.<sup>46</sup>. The standard deviation reported here refers to that of the fitting parameter based on the 20 repeated scans used to obtain the butterfly loop that is shown in figure 8a. The value reported here is also in good agreement, with a 20 pm/V value previously reported for piezoelectric barium titanate nanofibers.<sup>48</sup> However, it is important to note that there are several biases present in PFM measurements, including the effect of surface charges which act upon the body of the PFM cantilever, in addition to the forces that are acting solely on the tip.<sup>47</sup> For cases where domains are more randomly oriented and relatively small in comparison to the cantilever dimensions this bias is reduced, if not eliminated.<sup>47</sup> Figure 9 shows the phase angle scan of a barium titanate nanowire, revealing the as-calcined nanofibers exhibit a random distribution of domains, that are much larger than the cantilever tip. It is anticipated that all calcination conditions would result in the formation of a random domain structure, requiring subsequent poling steps. Although 1-dimensional piezoelectrics cannot easily be poled using conventional techniques, methods have been developed where arrays of piezoelectric nanowires are assembled across electrodes, whereby electrical contacts can be made across both ends of the nanowire so as to apply a voltage along the length of the wire.,<sup>49</sup> Additionally, a corona discharge poling have been demonstrated as a means to pole piezoelectric barium titanate nanofibers/nanowires.<sup>50, 51</sup>

#### 4. Conclusion

In this work, we have demonstrated a method for obtaining aligned barium titanate nanofibers using a rotating copper wire drum collector, where the alignment was maintained during calcination by one of two removal methods. We found and quantitatively tested a distribution model that appropriately fit the measured diameters of the barium titanate nanofibers. Additionally, we presented links between processing conditions and phase evolution in electrospun barium titanate nanofibers. By increasing the

calcination time and temperature the formation of a tetragonal barium titanate phase was promoted, leading to the formation of barium titanate nanofibers with enhanced piezoelectric properties. The use of a rotating drum collector, alongside optimized calcination conditions hold potential for the large-scale synthesis of electrospun ceramic nanofibers with enhanced properties, greatly increasing their ability to be incorporated into novel devices and applications.

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### Figure Captions:

**Figure 1.** Photograph of the rotating mandrel set-up for the electrospinning of aligned nanofibers showing the as-spun amorphous fibers aligned across the parallel copper wires on the rotating mandrel.

**Figure 2.** Scanning electron micrographs of (a) uncalcined, aligned nanofibers (b) aligned barium titanate nanofibers, calcined at 750°C for 2h after removal via the carbon tape method, and aligned nanofibers calcined at (c) 750°C for 2h and (d) 1000°C for 2h after removal by the peeling off method.

**Figure 3.** Fiber diameter distributions for (a) uncalcined barium titanate fibers and (b) calcined barium titanate fibers with the log normal probability density function given by the black line. Experimental cumulative distribution of the (c) uncalcined and (d) calcined barium titanate fibers and along with fitted bimodal normal and log normal distributions.

**Figure 4.** a) X-ray diffraction spectra for barium titanate nanofibers calcined for 6h at 750°C, 875°C, and 1000°C. (b) X-ray diffraction spectra for barium titanate nanofibers calcined at 875°C for 2h, 4h, and 6h.

**Figure 5.** The weight percent of the (a) tetragonal, (b) cubic, and (c) hexagonal phase of barium titanate as a function of calcining temperature, for a constant calcining time of 6 h.

**Figure 6.** (a) Raman spectra for barium titanate nanofibers calcined for 6h at 750C, 875C, and 1000C and (b) Raman spectra for barium titanate nanofibers calcined at 875C for 2h, 4h, and 6h. Peaks are labeled as being attributed to tetragonal barium titanate (T), cubic barium titanate (C), hexagonal barium titanate (H), or barium carbonate (witherite, W).

**Figure 7.** The natural logarithm of average grain size plotted versus reciprocal temperature for the tetragonal phase. The slope from these plots was used to calculate the activation energy for grain growth in the barium titanate nanofibers.

**Figure 8.** (a) Butterfly loop and (b) phase plot of barium titanate nanofiber calcined at 750°C for 2h.

**Figure 9.** Phase angle scan of a barium titanate nanofiber calcined at 750°C for 2h, revealing the domain structure of the fiber, the width of the scan shown is 1 micron.

**Tables:**

**Table 1: Agreement indices for the Rietveld refinements on the barium titanate fibers**

<i>Temp</i> (°C)	<i>Time</i> (h)	<i>Trial</i>	$R_{exp}$ (%)	$R_{wt}$ (%)	$\chi^2_v$
750	2	1	5.03	5.99	1.42
		2	6.12	6.50	1.13
		3	6.49	7.08	1.19
750	4	1	6.42	7.95	1.54
		2	5.76	6.13	1.13
		3	6.87	7.45	1.18
750	6	1	4.89	5.05	1.07
		2	6.22	7.52	1.46
		3	6.64	7.56	1.30
875	2	1	5.09	6.00	1.39
		2	5.88	6.22	1.12
		3	6.60	7.16	1.18
875	4	1	4.95	6.19	1.57
		2	6.14	7.08	1.33
		3	6.75	7.44	1.21
875	6	1	5.13	5.88	1.31
		2	6.29	7.36	1.37

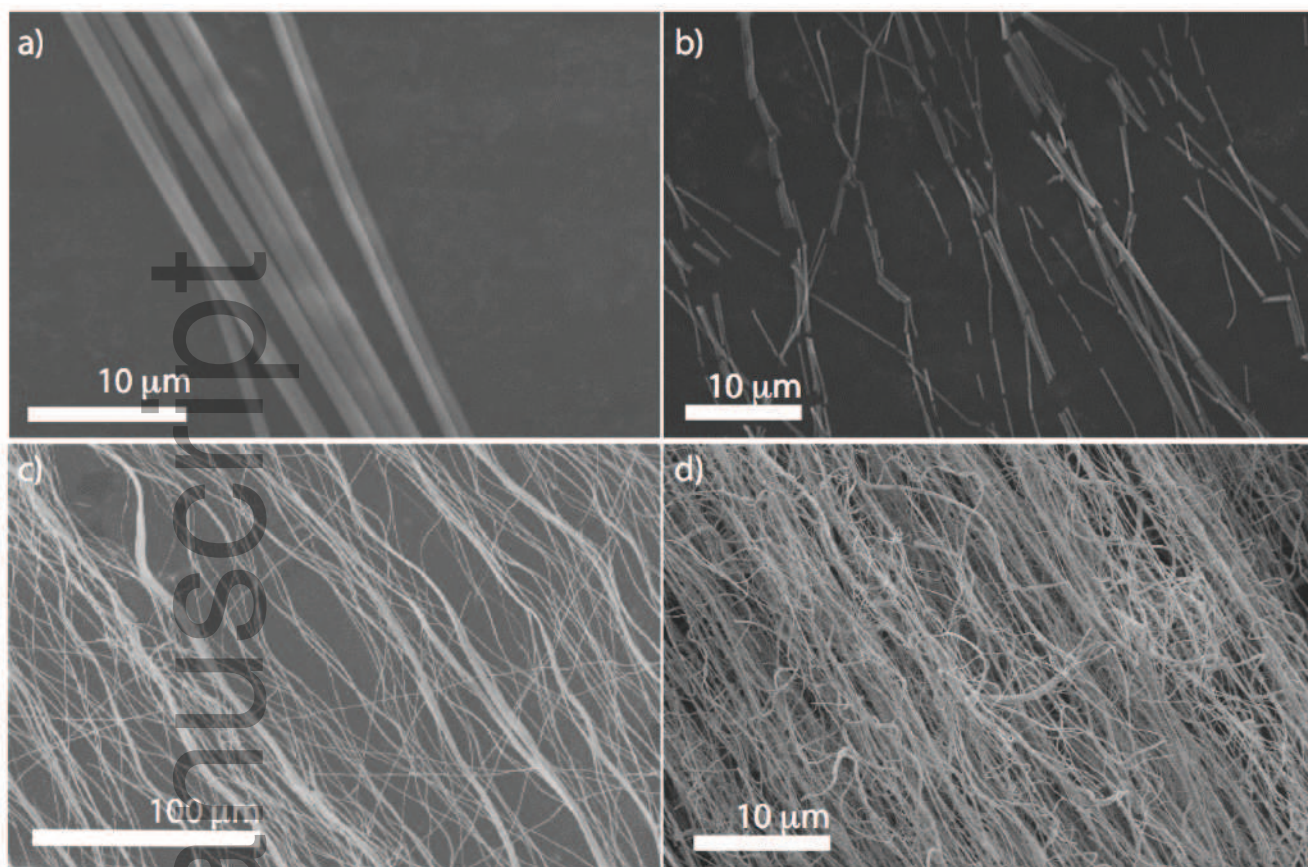
		3	6.61	7.03	1.13
1000	2	1	6.52	8.47	1.69
		2	6.65	8.34	1.57
		3	6.86	7.68	1.25
1000	4	1	4.92	6.34	1.66
		2	6.31	7.28	1.33
		3	6.69	7.69	1.32
1000	6	1	6.42	7.95	1.54
		2	6.54	7.94	1.48
		3	6.82	8.97	1.73

Agreement indices for the Rietveld refinements performed.  $R_{wt}$  and  $R_{exp}$  are the weighted and expected R-factors respectively and  $\chi^2_{\nu}$  is the reduced chi-square statistic.

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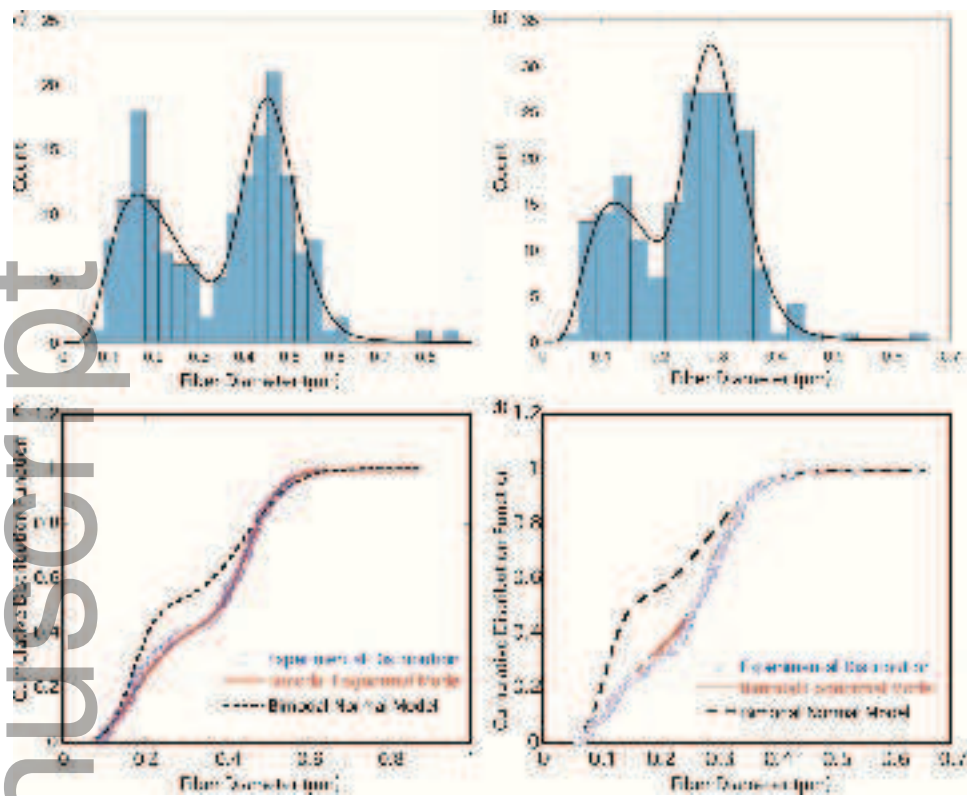


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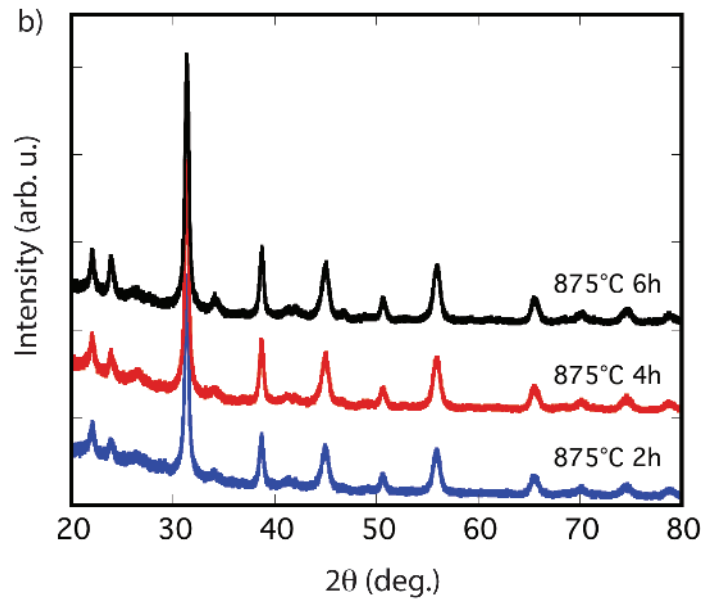
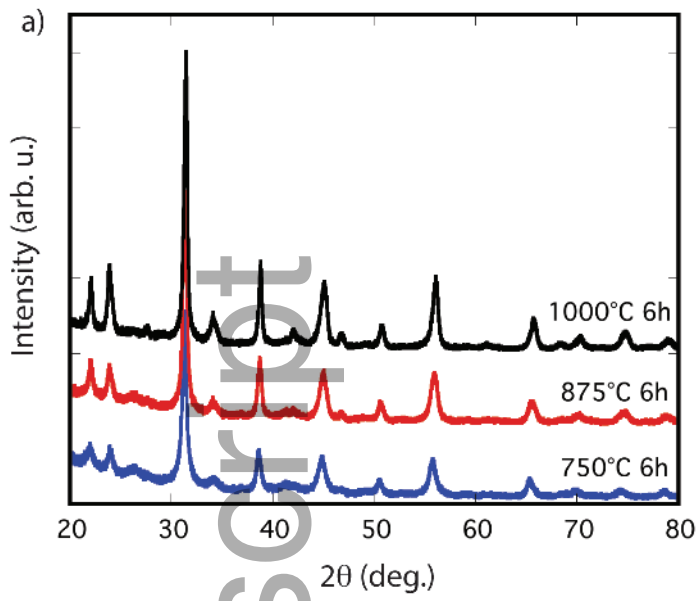


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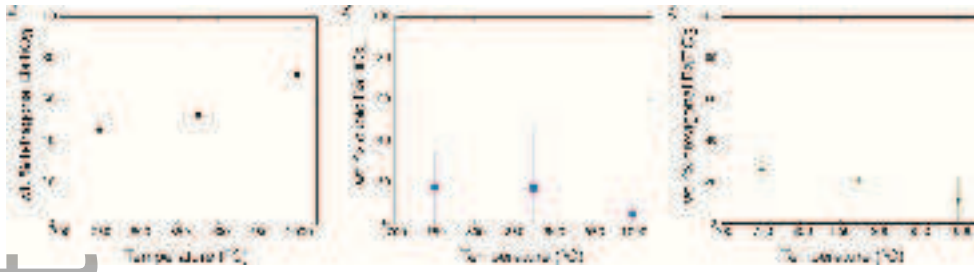


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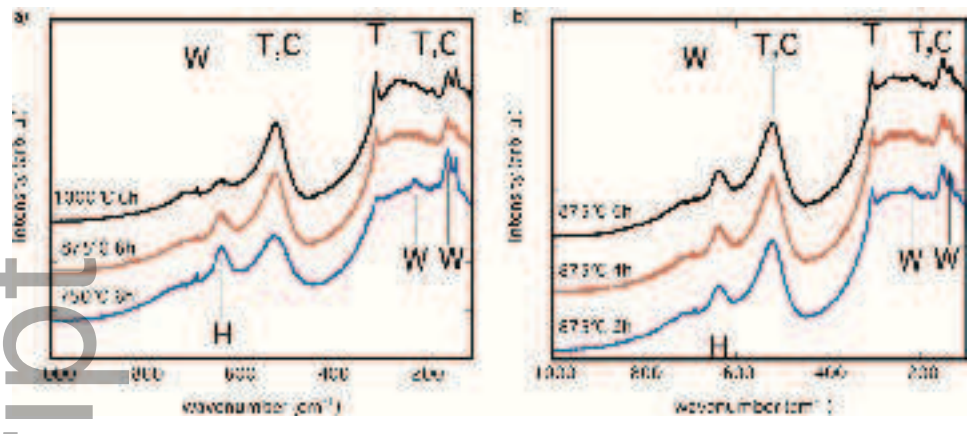


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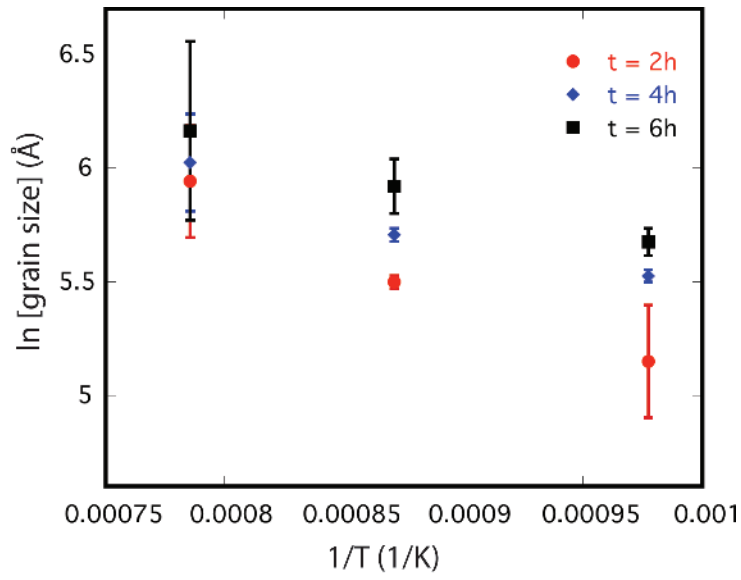
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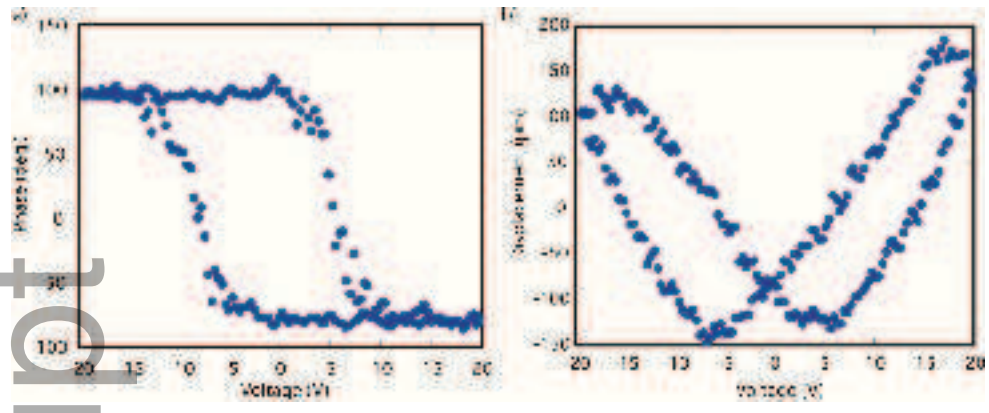
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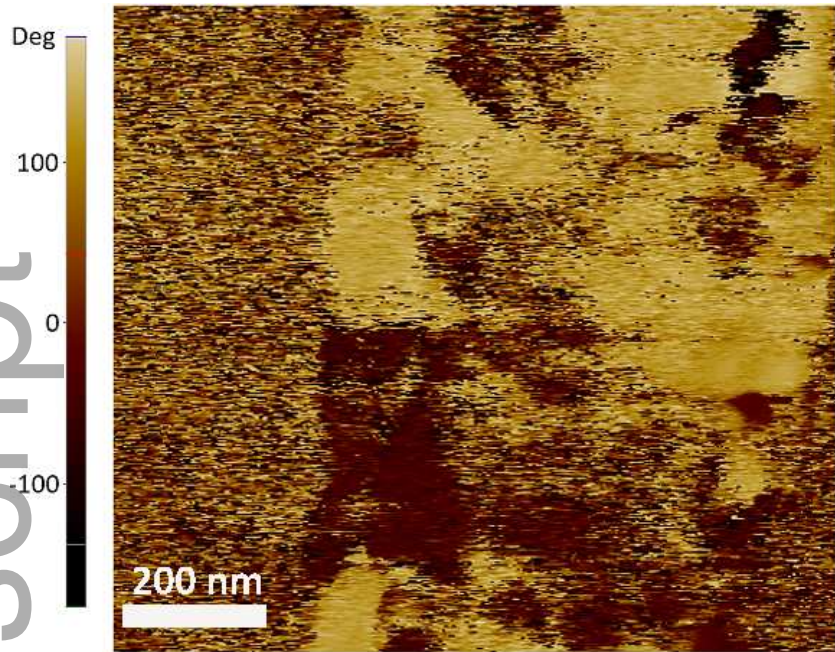
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