

1 **Electric Mars: A large trans-terminator electric**
2 **potential drop on closed magnetic field lines above**
3 **Utopia Planitia**

Glyn Collinson^{1,2}, David Mitchell³, Shaosui Xu³, Alex Gloer¹, Joseph
Grebowsky⁷, Takuya Hara³, Robert Lillis³, Jared Espley¹, Christian
Mazelle^{5,6}, Jean André Sauvaud^{5,6}, Andrei Fedorov^{5,6}, Mike Liemohn⁴, Laila
Andersson⁷, Bruce Jakosky⁷

Author Manuscript

Corresponding author: Glyn Collinson, Heliophysics Science Division, NASA Goddard Space
Flight Center, Greenbelt, MD, USA. (glyn.a.collinson@nasa.gov)

¹NASA Goddard Spaceflight Center,

**This is the author manuscript accepted for publication and has undergone full peer review but
has not been through the copyediting, typesetting, pagination and proofreading process, which
may lead to differences between this version and the Version of Record. Please cite this article**

D R A F T **doi:10.1002/2016JA023389** December 12, 2016, 8:51am

D R A F T

4 Abstract.

5 Parallel electric fields and their associated electric potential structures play
6 a crucial role in ionospheric-magnetospheric interactions at any planet. Al-
7 though there is abundant evidence that parallel electric fields play key roles

Greenbelt, Maryland, USA.

²Institute for Astrophysics and
Computational Sciences, The Catholic
University of America, Washington, District
of Columbia, USA

³Space Sciences Laboratory, University of
California, Berkeley, USA

⁴Space Physics Research Laboratory,
University of Michigan, Ann Arbor,
Michigan, USA

⁵CNRS, Institut de Recherche en
Astrophysique et Planétologie, Toulouse,
France

⁶University Paul Sabatier, Toulouse,
France

⁷Laboratory For Atmospheric and Space
Physics, Boulder, Colorado, USA

8 in Martian ionospheric outflow and auroral electron acceleration, the fields
9 themselves are challenging to directly measure due to their relatively weak
10 nature. Using measurements by the SWEA instrument aboard the NASA
11 *Mars Atmosphere and Volatile Evolution (MAVEN)* Mars Scout, we present
12 the discovery and measurement of a substantial ($\Phi_{MARS} = 7.7 \pm 0.6V$)
13 parallel electric potential drop on closed magnetic field lines spanning the
14 terminator from day to night above the great impact basin of Utopia Plani-
15 tia, a region largely free of crustal magnetic fields. A survey of the previous
16 26 orbits passing over a range of longitudes revealed similar signatures on
17 7 orbits, with a mean potential drop (Φ_{MARS}) of 10.9 ± 0.8 V, suggestive
18 that although trans-terminator electric fields of comparable strength are not
19 ubiquitous, they may be common, at least at these northerly latitudes.

Author Manuscript

1. Introduction

Although magnetic field aligned electric potentials play an important role in planetary ionosphere-magnetosphere interactions, very little is known about electric potential structures at Mars. However, there is abundant evidence for Martian parallel electric fields in charged particle data from the NASA *Mars Global Surveyor* (MGS), the ESA *Mars Express* (MEX), and the NASA *Mars Atmosphere and Volatile Evolution* (MAVEN) Mars Scout class mission [Jakosky *et al.*, 2015]. To date, discussions of Martian parallel electric fields in the literature have focused on two processes: **A.)** ion outflow and **B.)** auroral acceleration.

Ionospheric outflow: Electric fields play a central role in several important ionospheric transport and loss processes at both Mars and Venus. Ion pickup by the penetration of the (perpendicular) motional electric field of the solar wind ($-V \times B$) produces a distinct "plume" of keV ions [see e.g. Curry *et al.*, 2015, and references therein]. Ions are also accelerated by the draping and curving of the Interplanetary Magnetic Field (IMF) around the planet ($J \times B$ forces), which generates energized ion populations readily identified through their close proximity to the current sheet [Barabash *et al.*, 2007]. However, it has long been presumed that processes resulting in parallel electric fields also play a key role in ionospheric transport and escape [Lundin *et al.*, 2006a; Dubinin *et al.*, 2011]. Franklin *et al.* [2010] used fluxes of escaping photoelectrons to make an initial estimate of the escape rates of ionospheric plasmas due to the planetary ambipolar electric field ($E \approx \nabla P_e / en_e$). Later, Collinson *et al.* [2015] used field-aligned superthermal elec-

tron measurements by the Solar Wind Electron Analyzer (SWEA) [Mitchell *et al.*, 2016] aboard *MAVEN* to place an upper limit on the total potential drop of this global “polar wind like” ambipolar electric field of $\leq 2V$. Recently, Collinson *et al.* [2016] used *Venus Express* superthermal electron observations to discover the presence of an extremely powerful ($\sim 10V$) parallel electric potential drop in the ionosphere of Venus, finding it to be stable, persistent, and capable of accelerating oxygen ions directly to escape velocities in an “electric wind”. Thus Venus’ planetary electric potential is important for ion escape and global, whereas at Earth it is weaker ($\leq 2V$ Coates *et al.* [1985]) and confined to the magnetic poles. Given the lower Martian gravity, a $\geq 2V$ electric potential drop would be similarly effective in accelerating Oxygen ions to escape velocity. Thus, whilst parallel electric potentials are thought to play a key role in ion outflow, the direct evidence for their existence at Mars is currently ambiguous.

Auroral acceleration: The evidence for the presence of Martian parallel potential drops is currently ambiguous with regards to auroral electron acceleration above the Martian crustal anomalies. Lundin *et al.* [2006b] reported the discovery by MEX of locally accelerated electrons and ions on the deep nightside of Mars, with ion fluxes sufficient to produce a bright discrete aurora. Brain *et al.* [2006] presented collected MGS and MEX observations of peaked electron energy spectra at the crustal anomalies similar to terrestrial auroral electrons, finding that the most energetic examples occurred during the passage of space weather storm events, similar to the onset of sub-storms at Earth. Lundin *et al.* [2006c] further examined Martian auroral plasma acceleration, finding monoenergetic counterstreaming accelerated ions and electrons consistent with field-aligned electric

64 currents and electric field acceleration. A further study by *Halekas et al.* [2008] of Martian
65 auroral electrons observed by MGS found clear evidence for field-aligned currents, and
66 discrete electron acceleration events reminiscent of the terrestrial auroral zone. *Dubinin*
67 *et al.* [2008] reported observations of electron inverted “V” structures by both MGS and
68 MEX, similar to events in Earth’s auroral acceleration regions, raising the hypothesis
69 that localized aurora on Mars are also generated by electron acceleration through parallel
70 electric fields. A follow up study by [*Dubinin et al.*, 2009] reported the observation of
71 auroral activity on numerous orbits of MEX during a two week period, implying that the
72 auroral acceleration, and therefore the accelerating parallel electric potential drops, are a
73 stable phenomenon, albeit spatially localized.

74
75 In this study we present field-aligned electron measurements by *MAVEN* SWEA to
76 report the new discovery of large ($\sim 10V$) electric potential drops along closed magnetic
77 field lines which span the terminator of Mars, connecting the dayside and nightside iono-
78 spheres. Although the detection of a trans-terminator electric potential is novel, it is
79 important to note that the existence of closed trans-terminator magnetic field lines at
80 Mars is well established in the literature. For example, in their survey of electron data
81 from Mars Express (MEX), *Frahm et al.* [2006] noted that photoelectrons were occasion-
82 ally seen beyond the sunlit ionosphere, suggesting magnetic connectivity between dayside
83 and nightside ionosphere, although without a magnetometer to organize electron spectra
84 by pitch angle it was not possible to confirm this. *Liemohn et al.* [2006] predicted such
85 connectivity with magnetohydrodynamic (MHD) model results, tracing thousands of field
86 lines in a search for day-night connectivity. Furthermore, *Liemohn et al.* [2007] explicitly

87 looked for this trans-terminator magnetic connection in their analysis of MHD simulations
88 for the initial auroral observation conditions, concluding that field lines can connect from
89 many locations across the dayside to the nightside ionosphere.

2. Method

90 In this section, we describe how electric potential drops may be directly measured
91 through examination of outflowing and inflowing planetary electrons, using a technique
92 successfully employed at Earth to measure the large ($\sim 20V$) parallel potential drops
93 that occur above orbiting spacecraft (i.e. between $\sim 3800km$ and the magnetopause)
94 [Winningham and Gurgiolo, 1982; Wilson et al., 1997; Kitamura et al., 2012].

2.1. Measurement topology

95 First let us consider a simplified example. Figure 2 shows a sketch of the overall mea-
96 surement topology. *MAVEN* is above the Martian terminator, on a "closed" magnetic
97 field line (black) which we define as having both ends of the field line intersecting the
98 electron exobase at a "footpoint". In Mars' complex magnetic environment, such a closed
99 topology might result from a fringe field of a crustal magnetic remanent, or from a draped
100 Interplanetary Magnetic Field line that passes below the exobase at both footpoints. For
101 the purposes of this study, we do not differentiate and use the nomenclature "closed" for
102 either case: all that matters is that the ionosphere is magnetically disconnected from the
103 solar wind, and the spacecraft connected to the ionosphere on both ends of the magnetic
104 field line on which it is located. Since electrons (unlike ions) move effectively instanta-
105 neously along field lines and are excellent tracers of magnetic connectivity, such a closed
106 field topology can be established by the presence of ionospheric photoelectrons and the

107 absence of inflowing solar wind electrons which otherwise readily have access to the topside
108 ionosphere [*Brain et al.*, 2007]. In this study, we consider a closed “trans-terminator” field
109 line, with one footpoint on the dayside ionosphere, and the other footpoint terminating
110 below the exobase in the nightside ionosphere.

111
112 The existence of such closed trans-terminator magnetic field lines at Mars is well estab-
113 lished in the literature. For example, in their survey of electron data from Mars Express
114 (MEX), *Frahm et al.* [2006] noted that photoelectrons were occasionally seen beyond the
115 sunlit ionosphere, suggesting magnetic connectivity between dayside and nightside iono-
116 sphere, although without a magnetometer to organize electron spectra by pitch
117 angle it was not possible to confirm this. *Liemohn et al.* [2006] predicted such connectivity
118 with magnetohydrodynamic (MHD) model results, tracing thousands of field lines in a
119 search for day-night connectivity. Furthermore, *Liemohn et al.* [2007] explicitly looked for
120 this trans-terminator magnetic connection in their analysis of MHD simulations for the
121 initial *areosol* observation conditions, concluding that field lines can connect from many
122 locations across the dayside to the nightside ionosphere.

123
124 In this example, there is a 20V electric potential drop distant from the *MAVEN*, occur-
125 ring somewhere between the spacecraft and the nightside ionosphere. In this schematic,
126 two electrons are created in the dayside ionosphere: electron №1 with an energy of 30eV,
127 and electron №2 with an energy of 20eV. Both pass *MAVEN* on their way to the nightside,
128 and both are retarded equally by the electric force. Electron №1 is slowed to an energy
129 of 10eV and impacts the nightside atmosphere. However, electron №2 is at just a low

130 enough energy that it is stopped and then reflected back towards the dayside. This elec-
131 trostatic mirror conserves energy, so any reflected electrons that are observed by *MAVEN*
132 will have regained their original energy. If large numbers of such electrons were created
133 in the dayside ionosphere, *MAVEN* will observe an equal flux of the lower energy ($20eV$)
134 electrons flowing from both the dayside source, and the nightside electrostatic mirror.
135 However, although the lower energy ($20eV$) electrons will be virtually isotropic (by which
136 we mean equal fluxes flowing from the dayside and nightside), the higher energy ($30eV$)
137 electrons will be strongly anisotropic (by which we mean higher fluxes flowing from the
138 dayside than from the nightside).

139
140 Whilst such a potential structure reflects negatively charged electrons back up to the
141 spacecraft, it is important to note that positively charged ions would be accelerated away
142 from the spacecraft, into the nightside. For comparison with electrons, Figure 2C shows
143 how an ion (H^+) would be affected by such a potential drop. The hypothetical ion passes
144 *MAVEN* at $20 eV$, is accelerated to $40 eV$, and lost in the nightside ionosphere. Thus,
145 under this topology, *MAVEN* is remote from the ion acceleration region, and thus will not
146 observe any change in ion energies since this process is happening at a distance from the
147 spacecraft. Additionally, whilst superthermal electrons effectively move instantaneously
148 along such a relatively short magnetic field line, ions move considerably slower due to
149 their substantially greater mass. Electrons are therefore the best particles to use to probe
150 field-aligned electric potential structures.

2.2. Simulated example SWEA electron spectra

151 Having established the basic overall topology we shall now move from hypothetical test
152 particles to simulations of realistic Martian electron energy spectra. We shall first con-
153 sider the control case where no parallel electric fields are present, and then examine how
154 the energy spectra of Martian ionospheric electrons are affected by the presence of our
155 hypothetical $20V$ electric potential drop.

156
157 **Description of Mars-STET model:** The simulated spectra in this study were gen-
158 erated using the Mars SuperThermal Electron Transport (STET) model [*Liemohn et al.*,
159 2003; *Xu and Liemohn*, 2015; *Xu et al.*, 2015], based on the earlier Earth version of the
160 STET code [*Akhanov and Liemohn*, 1995]. STET is a time-dependent, multi-stream
161 model that solves the gyro-averaged Boltzman equation to self-consistently calculate the
162 flux of superthermal electrons at any point along a single magnetic flux tube. STET
163 self-consistently models the generation of photoelectrons from a simulated Martian atmo-
164 sphere from a given solar irradiance, and then simulates their transport to the nightside
165 along the field line above the superthermal electron exobase [e.g. *Xu*, 2015; *Xu et al.*,
166 2016]. STET includes electron-neutral inelastic and elastic collisions, as well as Coulomb
167 collisions with thermal electrons and ions.

168
169 For this particular simulation, the magnetic field line is assumed to be symmetric about
170 the terminator, with a maximum field strength of $20nT$ at the two footpoints (one in
171 the dayside, one on the nightside), and a minimum field strength of $10nT$ at the 400km
172 altitude above the terminator. The dayside footprint (the location where the field line

173 intersects the electron exobase) of the field line is set at 85° solar zenith angle (SZA)
174 and the other in darkness (i.e. no photoionization production). The models of both the
175 neutral atmosphere and thermal plasma environment along the field line are obtained
176 from the Mars Global Ionosphere and Thermosphere Model (M-GITM) [Bougher *et al.*,
177 2015]. Solar irradiance is obtained using the Hinteregger-81 model [Hinteregger *et al.*,
178 1981], with an Earth F10.7 of 130 solar flux units.

179
180 Figures 2A,B show simulated field-aligned and anti-aligned electron spectra for two dif-
181 ferent strengths of electric potential drop: 2A.) an electric potential so weak that it cannot
182 be measured, and 2B.) a hypothetical 20V potential drop, comparable to the large-scale
183 potential drops observed at high altitudes at Earth.

184
185 **Case A.) - No measurable electric potential:** Let us assume *MAVEN* is situated
186 above the terminator of Mars, and is on a closed magnetic field line that has one foot-
187 point in the dayside atmosphere, and the other footpoint somewhere on the nightside (as
188 in Figure 2C). Consistent with the MAVEN observations to be presented later, electrons
189 outflowing from the dayside ionospheric source region are shown in red, and the returning
190 electrons, inflowing back from the nightside are shown in blue. Ionospheric photoelectrons
191 exhibit three key features: **1.)** a cluster of sharp peaks (unresolved) from 23 to 27 eV
192 resulting from the photoionization of CO_2 and O by the intense 30.4-nm He-II line in the
193 solar EUV spectrum; **2.)** an abrupt cut-off at $\approx 75\text{eV}$ due to a sharp drop in the intensity
194 of the solar irradiance at wavelengths shorter than 16nm [Gan *et al.*, 1990; Richards and
195 Peterson, 2008], referred to as the "Aluminium (Al) Edge"; and **3.)** a peak of oxygen

196 auger electrons at $\sim 500\text{eV}$ due to K-shell ionization of CO_2 and O by soft X-rays.

197

198 In the absence of any appreciable electric potential drop (Figure 2A), the only elec-
199 trons returning from the nightside are those which have been magnetically reflected or
200 backscattered from the nightside footpoint (blue spectrum). Although there are sources
201 of electrons on the nightside *Fowler et al.* [2015], these are much weaker than dayside
202 photoproduction, and thus any trans-terminator magnetic field will exhibit a strong pitch
203 angle anisotropy in superthermal electrons at all energies [*Brain et al.*, 2007].

204

205 **Case B.) - A hypothetical +20V electric potential drop:** Now let us impose a
206 hypothetical +20V electric potential drop along the closed magnetic field line, occurring
207 somewhere between the location of *MAVEN* and the nightside footpoint, as in Figure 2C.
208 Figure 2B shows a sketch of how this hypothetical +20V electrostatic mirror will effect the
209 energy spectra of photoelectrons observed by *MAVEN*. Electrons are now isotropic below
210 20eV (by which we mean a quasi-uniform flux at all pitch angles), with all electrons below
211 this energy being reflected back to the spacecraft. However, above this energy, electron
212 distributions transition from isotropic to anisotropic, with greater fluxes coming from the
213 dayside, as electrons above this transition energy have sufficient velocity to overcome the
214 electric potential and impact the nightside atmosphere, and the only electrons returning
215 are those which have been magnetically reflected or backscattered.

216

217 Thus, the total magnitude of electric potential drop may be directly measured through
218 examination of at what energy the outflowing and inflowing electrons transition from

219 isotropy to anisotropy. This transition is very sharp [*Kitamura et al.*, 2012], occurring in
220 less than the energy channel bin width of our instrument and therefore the magnitude of
221 the electric potential drop can be unambiguously determined from the energy at which
222 the ratio between outflowing and inflowing electrons diverges from unity.

223
224 It is also important to note that this (isotropic to anisotropic) “transition” feature is
225 specific to electrostatic mirrors. Whilst a magnetic mirror might also reflect particles
226 back to the dayside, magnetic mirrors also produce loss cones, and are energy indepen-
227 dent and adiabatic. Therefore, magnetic mirroring will not create this specific type of
228 energy-dependent transition from isotropic to anisotropic distributions, and thus we may
229 unambiguously infer the presence of an electrostatic mirror from the presence of such a
230 feature.

231
232 **Description of MAVEN-SWEA:** The MAVEN Solar Wind Electron Analyzer is
233 a top-hat electrostatic plasma analyzer mounted on the end of a 1.5m boom, and is
234 equipped with electrostatic deflector plates which give it a view of 80% of the sky. It
235 measures electrons over the energy range of 3-4600eV with a resolution of 17% ($\Delta E/E$).
236 SWEA operates in two data collection modes: “ionospheric mode” (when MAVEN is
237 below 2000km), and “solar wind” mode (>2000km). Each mode creates two data streams;
238 a “burst mode” and a lower resolution “survey mode” (see table 4 of *Mitchell et al.* [2016]
239 for a summary of SWEA operational modes). Due to the limited data rates available
240 to planetary spacecraft such as MAVEN, only a fraction of burst mode data can be
241 telemetered back. The event presented later in this paper represents our best current

242 example, partly because ionospheric burst mode data was available, providing 64 energy
 243 bins and 8-sec cadence for energy/angle (3D) distributions, and 64 energy bins and 2-sec
 244 cadence for pitch angle distributions (PADs). For full details of SWEA, see *Mitchell et al.*
 245 [2016].

3. Evidence for a trans-terminator electric potential at Mars

246 In this section we present our current best example of a signature consistent with a
 247 trans-terminator electric field, occurring on the 3rd of March 2015. Figure 3 shows a
 248 map of the relevant orbit geometry (№ 821), and Figure 4 shows *MAVEN* observations
 249 for the low-altitude portion of orbit № 821. Periapsis was in the northern hemisphere
 250 near the terminator. The light blue line running parallel to the orbital path in Figure 3
 251 shows the region from which *MAVEN* data are shown in Figure 4. Altitude (4A), and
 252 three-component magnetometer measurements (4B) are shown for reference and context.
 253 Figure 4C shows omni-directional energy spectra measured by SWEA, with time on the
 254 x-axis (with the same range as 4A,B), energy in eV on the y-axis, and colour denoting the
 255 \log_{10} of differential energy flux in $eV cm^{-2} sr^{-1} s^{-1} eV^{-1}$.

256
 257 Shown beneath are three SWEA electron scans which have been ordered by the mag-
 258 netic field vector. In all three scans, the only electron populations visible are planetary
 259 photoelectrons. Since we observe no evidence for any inflowing solar wind electrons [see
 260 e.g. *Collinson et al.*, 2015], we determine that all three scans were taken when *MAVEN*
 261 was on closed magnetic field lines. Each scan (α, β, γ) shows two spectra: one a field-
 262 aligned spectrum (averaged within $\pm 60^\circ$ of $\hat{\mathbf{B}}$) and the other an anti-aligned **spectrum**
 263 (averaged within $\pm 60^\circ$ of $-\hat{\mathbf{B}}$). The two spectra have been coloured so that, consistent

264 with Figure 2, red denotes the distribution of tailward flowing electrons, and blue the
265 sunwards flowing distribution. Below each spectra is a plot showing the ratio of outflow-
266 ing to inflowing electron fluxes. Full 3D pitch angle distributions for each scan are also
267 shown in Figure 5.

268
269 **Scan alpha:** (“ α ” - left panel) represents the case where both footpoints of the closed
270 magnetic field line passing through *MAVEN* are in the dayside photoproduction region,
271 and thus the energy spectra of electrons traveling parallel and anti-parallel to the field
272 are the same.

273
274 **Scan beta:** (“ β ” - center panel) shows the case similar to the simulated electrons in
275 Figure 2A, where *MAVEN* is now on a closed magnetic field line that spans the terminator,
276 but there is no evidence for electrostatic mirroring. Consistent with our simulation, the
277 only electrons inflowing from the nightside (blue) are magnetically reflected or backscat-
278 tered, with a diminished flux compared to those outflowing from the source region (red).

279
280 **Scan gamma:** (“ γ ” - right panel), was measured just over 45 minutes later, when
281 *MAVEN* was at 106° Solar Zenith Angle (SZA), above the Utopia Planitia impact basin,
282 where there are no measurable crustal magnetic sources. *MAVEN* was still on closed
283 magnetic field lines, and would be until $\sim 11:46$, at which point connection to the solar
284 wind is evident in the increase in flux at all energies in the SWEA spectrogram (Figure
285 4C). The outflowing (red) spectra are similar to the outflowing photoelectron spectra in
286 scans α , β , albeit at a lower flux as the source region from which these electrons are

287 outflowing from appears to be closer to the terminator with a lower photo-production
288 rate. Above 6.5eV the electron distributions are anisotropic, consistent with our simu-
289 lations of backscattered photoelectrons on closed trans-terminator magnetic field lines.
290 However unlike scan β , below 6.5eV the electron pitch-angle distribution transitions to
291 isotropic, with equal fluxes of inflowing (blue) and outflowing (red) electrons. This be-
292 haviour can be explained by an electrostatic mirror (Figure 2B), and is consistent with
293 similar (isotropic to anisotropic) transitions presented at Earth by *Kitamura et al.* [2012]
294 associated with the presence of large parallel electric potential structures at high altitudes.

295
296 **Correcting for the electrostatic spacecraft potential:** As electrons arrive at
297 *MAVEN* they fall through an additional electric potential drop due to the electrostatic
298 charging of the spacecraft. If the spacecraft is negatively charged, all electrons are slowed,
299 whereas if the spacecraft is positively charged all electrons are accelerated towards the
300 detector, resulting in an energy shift in the spectra of both inflowing and outflowing elec-
301 trons. Therefore, the 6.5 eV transition from isotropic to anisotropic pitch angle measured
302 in scan γ corresponds to the total potential drop (Φ_{TOTAL}) of $+6.5\text{ V}$, with a $\pm 0.5\text{ V}$ uncer-
303 tainty arising from the width of the energy bin at which the transition occurred. However,
304 Φ_{TOTAL} is the sum of the Martian potential drop (Φ_{Mars}) and the spacecraft potential
305 (Φ_{SC}). Using the *MAVEN* Langmuir Probe and Waves (LPW) experiment, we may in-
306 dependently measure the spacecraft potential during scan γ , finding $\Phi_{SC} = -1.5 \pm 0.5$
307 V. Thus, we may correct for the spacecraft potential to determine the magnitude of the
308 Martian electric potential drop (Φ_{Mars}) using equation 1.

$$\Phi_{Mars} = \Phi_{Total} - \Phi_{SC} \quad (1)$$

309 Applying this correction, we find that the Martian component of the total field-aligned
310 electric potential drop was $\Phi_{Mars} = 7.7 \pm 0.7V$.

311
312 **Expected lack of effect on local ions:** As shown in Figure 1, although such a
313 distant electric potential structure reflects electrons below the "transition" energy (in the
314 case of scan γ , 6.5 eV) from the nightside back to the spacecraft, we would not expect
315 to see any change in the behavior of the local ion population. This is because while the
316 distant $\approx -8V$ electric potential is reflecting negatively charged ($< 8 eV$) electrons back
317 up towards *MAVEN*, it is accelerating positively charged ions away from the spacecraft
318 and down into the distant nightside ionosphere. Therefore, while the potential structure
319 must be accelerating ions, this acceleration is occurring at some distance from the space-
320 craft, and since *MAVEN* is effectively "upstream" of this process we would not expect to
321 see any effect on the local ion population.

322
323 Consistent with this interpretation, no obvious change in the energy of the ions was
324 observed by *MAVEN* STATIC. This gives further weight to the interpretation of a distant
325 (rather than local) electric potential structure, and full details of STATIC observations
326 can be found in the supplemental information section S1, together with continuous SWEA
327 pitch angle distributions, and a comparison between observed magnetic fields and crustal
328 field models.

4. Discussion

4.1. Investigating alternative explanations

329 Before pursuing any further analysis, a series of simulations was first carried out to see
330 whether the parallel and anti-parallel spectra in scan γ could be explained in some way
331 other than an electrostatic mirror. We used the Mars-STET model to investigate the
332 effect of magnetic field configuration, neutral densities, and thermal plasma density on
333 the parallel and anti-parallel electron distributions.

334
335 First, asymmetric field lines are tested, with a high magnetic field strength at the
336 footpoint in sunlight and low field strength at the footpoint in shadow (and vice versa).
337 However, it was found that the results did not affect the electron distributions apprecia-
338 bly. Then, we increased the neutral density in the nightside atmosphere (to increase the
339 collisional interaction with the neutrals) by artificially multiplying the density by factors
340 of 5 and 10. Although this raised the superthermal electron exobase to higher altitudes, it
341 did not affect returning photoelectron fluxes at 310 km (observational altitude). Finally,
342 to investigate whether the equal fluxes of returning superthermal electrons below $7eV$
343 could be explained through coulomb scattering, the nightside ionosphere was significantly
344 enhanced by multiplying its densities by an unrealistic factor of 100 times. We found that
345 although at lower altitudes more uniform fluxes could be observed, at 310km, where scan
346 γ was taken, the simulated return flux was still half of that of the outgoing flux from the
347 source region. However, this is hardly surprising since there is no known mechanism in
348 which backscatter can act like an electrostatic mirror, returning 100% of the incident flux

349 only for a specific energy range.

350

351 We conclude that even large changes in the neutral and plasma densities and variations
352 of the field magnitude at the footpoints cannot account for the observed parallel and
353 anti-parallel energy spectra. Thus, the most plausible explanation for scan γ is electro-
354 static reflection resulting from a potential drop between the spacecraft and the magnetic
355 footpoint at the nightside exobase.

356

4.2. A preliminary search for more events

357 A brief survey was undertaken to see whether any other spectral discontinuities could
358 be observed (and the experiment repeated), or whether the trans-terminator field implied
359 by scan γ on Orbit № 821 represented a rare or potentially unique event. To address
360 this question we examined the 26 orbits preceding 821, finding a total of 10 electric field
361 signatures similar to those in scan γ on 7 of the orbits, suggesting that at least at these
362 northern latitudes, trans-terminator electric fields on closed field lines are a common oc-
363 currence. Table 1 shows a summary of each of these events, showing the orbit number,
364 date and time, the altitude and solar zenith angle of *MAVEN* at the time of the SWEA
365 observation, and the total potential drop (Φ_{TOTAL}) implied from the energy bin of the
366 transition from isotropic to anisotropic pitch angle distributions, the spacecraft potential
367 (Φ_{SC}) as measured by MAVEN-LPW, and the corrected strength of the total Martian
368 electric potential drop (Φ_{MARS}). The uncertainty is taken from the energy resolution
369 ($\Delta E/E = 17\%$) of SWEA and the $\approx \pm 0.5$ V uncertainty in the measurement of the
370 spacecraft potential by LPW. Whilst Orbit № 821 represents our weakest example of a

371 field ($\Phi_{MARS} = 7.7 \pm 0.7$ V), it remains our best because the full “burst” resolution data
 372 was available, whereas on all other orbits only the “survey” mode data was transmitted
 373 to Earth, with reduced energy sampling, reducing our ability to resolve the energy of the
 374 transition, and thus the strength of the potential drop.

375
 376 Although this handful of observations does not represent a statistically significant sam-
 377 ple, we may surmise a few properties of these Martian trans-terminator electric fields along
 378 closed magnetic fields. Firstly, with a mean strength of $\bar{\Phi}_{MARS} = 11$ V, these fields are
 379 strong enough to accelerate ionospheric O^+ and O_2^+ to escape energy, and could thus drive
 380 trans-terminator transport. Secondly, the fact that such electric fields are not observed
 381 on every trans-terminator closed field (e.g. Fig. 4 scan β) indicates that the conditions
 382 responsible for these fields are not always present.

4.3. Speculation as to their origin

383 An electric field in the collisionless region above the ionosphere would be expected to
 384 arise from the electron pressure gradient:

$$E_{\parallel} \approx -\frac{\nabla P_e}{en_e} \quad (2)$$

385 This mechanism is key to the formation of Earth’s polar wind [*Banks and Holzer, 1968*].
 386 At Mars, while there is unquestionably a gradient in electron pressure across the termi-
 387 nator, it is insufficient to explain the 11V trans-terminator potential drops reported here.
 388 Furthermore, if such large potential drops were simply the result of the pressure gradient
 389 across the terminator, then one might expect to observe such potential drops on every
 390 closed trans-terminator field line. However, this is not the case (see scan β , figure 4).

391 Although it is possible that such a strong ambipolar field might briefly form in response
392 to some hypothetical dramatic transient change in the Martian ionosphere, one would
393 not expect such a formation mechanism to create stable potential structures for whole
394 minutes at a time, as was the case for scan γ . Thus, whilst electron pressure gradient
395 driven electric fields must be present, they cannot satisfactorily explain these strong trans-
396 terminator potential drops, and additional unknown physical processes must be at work.

397
398 At present, the cause of these trans-terminator electric fields on closed field lines is
399 unknown, and we cannot propose any particular mechanism, even speculatively. This
400 said however, any future mechanism must explain two observed features: (1.) the electric
401 potential drops generated by these fields are much stronger than that of the background
402 planetary electric field ($< 2V$ [Collinson *et al.*, 2015]), and (2.) the fields are sometimes
403 but not always observed, even when the magnetic topology appears to be the same. This
404 strongly motivates future statistical studies of parallel electric fields at Mars to see if;
405 a.) if they are localized around the crustal magnetic field remanents [Acuña *et al.*, 1998;
406 Connerney *et al.*, 1999]; b.) if they occur at any given longitude, latitude, and altitude;
407 and c.) if they are associated with any particular upstream drivers.

5. Conclusions

408 At 11:45:25 on the 3rd of March 2015, the NASA *MAVEN* Mars Scout was flying
409 above Utopia Planitia, a 3300km wide impact basin, with no measurable crustal mag-
410 netic sources. The spacecraft was ascending through 323km in altitude having just crossed
411 over the terminator from day into night. From electron spectra measured by the Solar
412 Wind Electron Analyzer (SWEA), we know that *MAVEN* was on a closed magnetic field

413 line, with one magnetic footpoint on the dayside and the other on the nightside. This
414 magnetic topology is not unusual, and indeed a similar day-to-night magnetic connection
415 had occurred only 45 minutes previously. However, the parallel and anti-parallel electron
416 distributions show a clear signature for electrostatic reflection, revealing the presence of a
417 $\Phi_{MARS} = 7.7 V$ potential drop from the spacecraft to the nightside footpoint of the mag-
418 netic field line. A brief survey of the previous 26 orbits revealed a total of 10 such similar
419 signatures on 7 orbits, with a mean total electric potential drop of $\Phi_{MARS} = 10.9 \pm 0.8 V$.
420 From this small unstatistical sample, we see that such cross-terminator potentials are
421 sometimes but not always observed, even when the conditions and magnetic topology
422 appear to be similar, at least at these northerly latitudes.

423
424 The fact that such potential drops are not observed on every trans-terminator magnetic
425 field line is highly suggestive that they are the result of either a localized or transient
426 phenomena. Although the generation mechanism of these large parallel electric potential
427 drops is presently unknown, we speculate that they may be associated with the magnetic
428 configuration around localized crustal magnetic sources, although this is pure conjecture.
429 Understanding their locality, origin, and their effects on the ionosphere and thermosphere
430 of Mars are therefore prime objectives for future theory, modeling, and observational
431 investigations.

432 **Acknowledgments.** *MAVEN* data is available from the NASA Planetary Data Sys-
433 tem. This work was partially supported by the CNES for the part based on observations
434 with the SWERA instrument embarked on *MAVEN*.

References

- 435 Acuña, M. H., et al. (1998), Magnetic Field and Plasma Observations at Mars:
436 Initial Results of the Mars Global Surveyor Mission, *Science*, *279*, 1676, doi:
437 10.1126/science.279.5357.1676.
- 438 Banks, P. M., and T. E. Holzer (1968), The polar wind, *J. Geophys. Res.*, *73*, 6846–6854,
439 doi:10.1029/JA073i021p06846.
- 440 Barabash, S., et al. (2007), The loss of ions from Venus through the plasma wake, *Nature*,
441 *450*, 650–653.
- 442 Bougher, S., A. Brecht, R. Schulte, J. Fischer, C. Parkinson, A. Mahieux, V. Wilquet,
443 and A. Vandaele (2015), Upper atmosphere temperature structure at the venusian ter-
444 minators. A comparison of {SOIR} and {VTGCM} results, *Planet. Space. Sci.*, (0), –,
445 doi:http://dx.doi.org/10.1016/j.pss.2015.01.012.
- 446 Brain, D. A., et al. (2006), On the origin of aurorae on Mars, *Geophys. Res. Lett.*, *33*,
447 L01201, doi:10.1029/2005GL024782.
- 448 Brain, D. A., R. J. Lillis, D. L. Mitchell, J. S. Halekas, and R. P. Lin (2007), Electron
449 pitch angle distributions as indicators of magnetic field topology near Mars, *Journal of*
450 *Geophysical Research (Space Physics)*, *112*, A09201, doi:10.1029/2007JA012435.
- 451 Coates, A. J., A. D. Johnstone, J. J. Sojka, and G. L. Wrenn (1985), Ionospheric photo-
452 electrons observed in the magnetosphere at distances up to 7 earth radii, *Planet. Space.*
453 *Sci.*, *33*, 1267–1275, doi:10.1016/0032-0633(85)90005-4.
- 454 Collinson, G., et al. (2015), Electric Mars: The first direct measurement of an upper limit
455 for the Martian "polar wind" electric potential, *Geophys. Res. Lett.*, *42*, 9128–9134,
456 doi:10.1002/2015GL065084.

- 457 Collinson, G. A., et al. (2016), The electric wind of Venus: A global and persistent “polar
458 wind”-like ambipolar electric field sufficient for the direct escape of heavy ionospheric
459 ions, *Geophys. Res. Lett.*, doi:10.1002/2016GL068327.
- 460 Connerney, J. E. P., et al. (1999), Magnetic Lineations in the Ancient Crust of Mars,
461 *Science*, *284*, 794, doi:10.1126/science.284.5415.794.
- 462 Curry, S. M., J. Luhmann, Y. Ma, M. Liemohn, C. Dong, and T. Hara (2015), Com-
463 parative pick-up ion distributions at Mars and Venus: Consequences for atmospheric
464 deposition and escape, *Planet. Space. Sci.*, *115*, 35–47, doi:10.1016/j.pss.2015.03.026.
- 465 Dubinin, E., G. Chanteur, M. Fraenz, and J. Woch (2008), Field-aligned currents and
466 parallel electric field potential drops at Mars. Scaling from the Earths aurora, *Planet.*
467 *Space. Sci.*, *56*, 868–872, doi:10.1016/j.pss.2007.01.019.
- 468 Dubinin, E., M. Fraenz, J. Woch, F. Duru, D. Gurnett, R. Modolo, S. Barabash, and
469 R. Lundin (2009), Ionospheric storms on Mars: Impact of the corotating interaction
470 region, *Geophys. Res. Lett.*, *36*, L01105, doi:10.1029/2008GL036559.
- 471 Dubinin, E., M. Fraenz, A. Fedorov, R. Lundin, N. Edberg, F. Duru, and O. Vaisberg
472 (2011), Ion Energization and Escape on Mars and Venus, *Space Sci. Rev.*, *162*, 173–211,
473 doi:10.1007/s11214-011-9831-7.
- 474 Fowler, C. M., et al. (2015), The first in situ electron temperature and density measure-
475 ments of the Martian nightside ionosphere, *Geophys. Res. Lett.*, *42*, 8854–8861.
- 476 Frahm, R. A., et al. (2006), Carbon dioxide photoelectron energy peaks at Mars, *"Icarus"*,
477 *182*, 371–382, doi:10.1016/j.icarus.2006.01.014.
- 478 Frahm, R. A., et al. (2010), Estimation of the escape of photoelectrons from Mars in 2004
479 liberated by the ionization of carbon dioxide and atomic oxygen, *Icarus*, *206*, 50–63,

480 doi:10.1016/j.icarus.2009.03.024.

481 Gan, L., T. E. Cravens, and M. Horanyi (1990), Electrons in the ionopause boundary
482 layer of Venus, *Journal of Geophysical Research (Space Physics)*, *95*, 19,023–19,035,
483 doi:10.1029/JA095iA11p19023.

484 Halekas, J. S., D. A. Brain, R. P. Lin, J. G. Luhmann, and D. L. Mitchell (2008), Dis-
485 tribution and variability of accelerated electrons at Mars, *Advances in Space Research*,
486 *41*, 1347–1352, doi:10.1016/j.asr.2007.01.034.

487 Hinteregger, H. E., K. Fukui, and B. R. Gilson (1981), Observational, reference and model
488 data on solar EUV, from measurements on AE-E, *Geophys. Res. Lett.*, *8*, 1147–1150,
489 doi:10.1029/GL008i011p01147.

490 Jakosky, B., et al. (2015), The mars atmosphere and volatile evolution (maven) mission,
491 *Space Science Reviews*, pp. 1–46, doi:10.1007/s11214-015-0139-x.

492 Khazanov, G. V., and M. W. Liemohn (1995), Nonsteady state ionosphere-
493 plasmasphere coupling of superthermal electrons, *J. Geophys. Res.*, *100*, 9669–9682,
494 doi:10.1029/95JA00526.

495 Kitamura, N., K. Seki, Y. Nishimura, N. Terada, T. Ono, T. Hori, and R. J.
496 Strangeway (2012), Photoelectron flows in the polar wind during geomagnetically
497 quiet periods, *Journal of Geophysical Research (Space Physics)*, *117*, A07214, doi:
498 10.1029/2011JA017459.

499 Liemohn, M. W., D. L. Mitchell, A. F. Nagy, J. L. Fox, T. W. Reimer, and Y. Ma (2003),
500 Comparisons of electron fluxes measured in the crustal fields at Mars by the MGS
501 magnetometer/electron reflectometer instrument with a B field-dependent transport
502 code, *Journal of Geophysical Research (Planets)*, *108*, 5134, doi:10.1029/2003JE002158.

- 503 Liemohn, M. W., et al. (2006), Mars Global MHD Predictions of Magnetic Connectivity
504 Between the Dayside Ionosphere and the Magnetospheric Flanks, *Space Sci. Rev.*, *126*,
505 63–76, doi:10.1007/s11214-006-9116-8.
- 506 Liemohn, M. W., et al. (2007), *Mars Global MHD Predictions of Magnetic Connectivity*
507 *between the Dayside ionosphere and the Magnetospheric Flanks*, p. 63, doi:10.1007/978-
508 0-387-70943-7-4.
- 509 Lundin, R., et al. (2006a), Ionospheric plasma acceleration at Mars: ASPERA-3 results,
510 *Icarus*, *182*, 308–319, doi:10.1016/j.icarus.2005.10.035.
- 511 Lundin, R., et al. (2006b), Plasma Acceleration Above Martian Magnetic Anomalies,
512 *Science*, *311*, 980–983, doi:10.1126/science.1122071.
- 513 Lundin, R., et al. (2006c), Auroral Plasma Acceleration Above Martian Magnetic Anoma-
514 lies, *Space Sci. Rev.*, *126*, 333–354, doi:10.1007/s11214-006-9086-x.
- 515 Mitchell, D. L., et al. (2016), The MAVEN Solar Wind Electron Analyzer, *Space Sci.*
516 *Rev.*, *200*, 495–528, doi:10.1007/s11214-015-0232-1.
- 517 Richard, P. C., and W. K. Peterson (2008), Measured and modeled backscatter of iono-
518 spheric photoelectron fluxes, *Journal of Geophysical Research (Space Physics)*, *113*,
519 A08321, doi:10.1029/2008JA013092.
- 520 Vignes, D., et al. (2000), The solar wind interaction with Mars: Locations and shapes
521 of the bow shock and the magnetic pile-up boundary from the observations of the
522 MAG/EP Experiment onboard Mars Global Surveyor, *Geophys. Res. Lett.*, *27*, 49–52,
523 doi:10.1029/1999GL010703.
- 524 Wilson, G. F., G. Khazanov, and J. L. Horwitz (1997), Achieving zero current for polar
525 wind outflow on open flux tubes subjected to large photoelectron fluxes, *Geophys. Res.*

- 526 *Lett.*, *24*, 1183–1186, doi:10.1029/97GL00923.
- 527 Winningham, J. D., and C. Gurgiolo (1982), DE-2 photoelectron measurements consistent
528 with a large scale parallel electric field over the polar cap, *Geophys. Res. Lett.*, *9*, 977–
529 979, doi:10.1029/GL009i009p00977.
- 530 Xu, S. (2015), Superthermal electrons at Mars: Photoelectrons, solar wind electrons, and
531 dust storm influences, Ph.D. thesis, University of Michigan.
- 532 Xu, S., and M. W. Liemohn (2015), Superthermal electron transport model for Mars,
533 *Earth and Space Science*, *2*, 47–64, doi:10.1002/2014EA000043.
- 534 Xu, S., M. W. Liemohn, W. K. Peterson, J. Fontenla, and P. Chamberlin (2015), Compar-
535 ison of different solar irradiance models for the superthermal electron transport model
536 for Mars, *Planet. Space. Sci.*, *119*, 62–68, doi:10.1016/j.pss.2015.09.008.
- 537 Xu, S., M. Liemohn, S. Bougher, and D. Mitchell (2016), Martian high-altitude pho-
538 toelectrons independent of solar zenith angle, *Journal of Geophysical Research (Space*
539 *Physics)*, *121*, 3767–3780, doi:10.1002/2015JA022149.

Figure 1. Schematic of a hypothetical trans-terminator magnetic field line, with 20 V electric potential drop on the nightside, and its effect on three particles; electron №1, born with an energy of 30eV and impacting the nightside ionosphere; electron №2, born with an energy of 20 eV and being reflected back to the dayside; and an ion (i^+) born with an energy of 20 eV, and accelerated by the potential drop into the nightside ionosphere.

Figure 2. Sketches showing the predicted spectra of photo-electrons in the tail of Mars for two hypothetical electric potential drops: Panel A.) Electric potential drop $\phi < 3V$ above *MAVEN*; Panel B.) Electric potential drop $\phi = 20V$ above *MAVEN*. The predicted spectra of electrons flowing outwards from the dayside source region are shown in red, and the hypothetical spectra of electrons flowing back from the night side in blue.

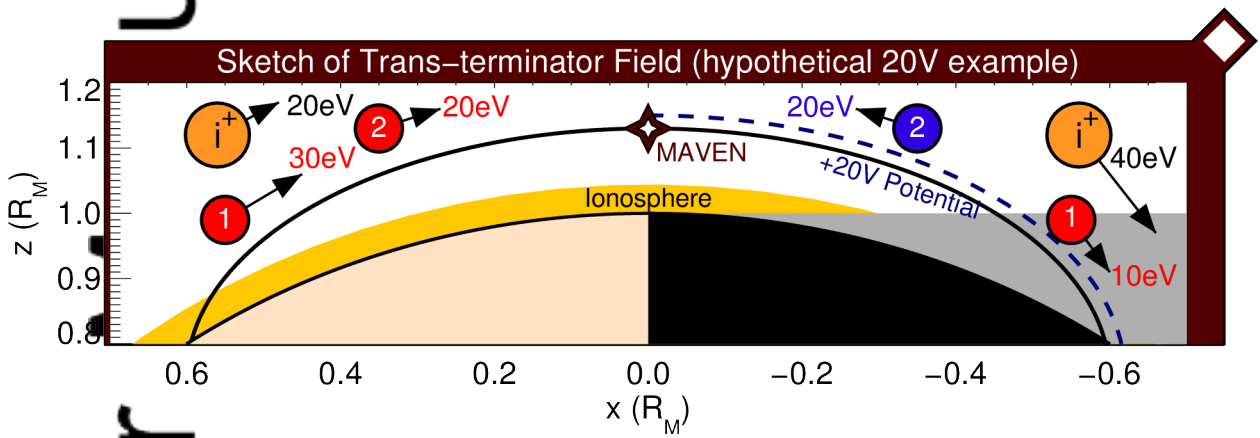
Figure 3. Map of *MAVEN* orbit № 867. The co-ordinate system is Mars Solar Orbital (MSO), where x points towards the sun, y points backwards along the tangent of the planetary orbit, and z completes the right-handed set pointing up out of the plane of the ecliptic out of the northern hemisphere. The orbit of the *MAVEN* Mars Scout is shown in maroon, with a model bow shock and magnetic pileup boundary (ionopause) in black, according to *Vignes et al.* [2000].

Table 1. Preliminary survey of Trans-terminator potential drops

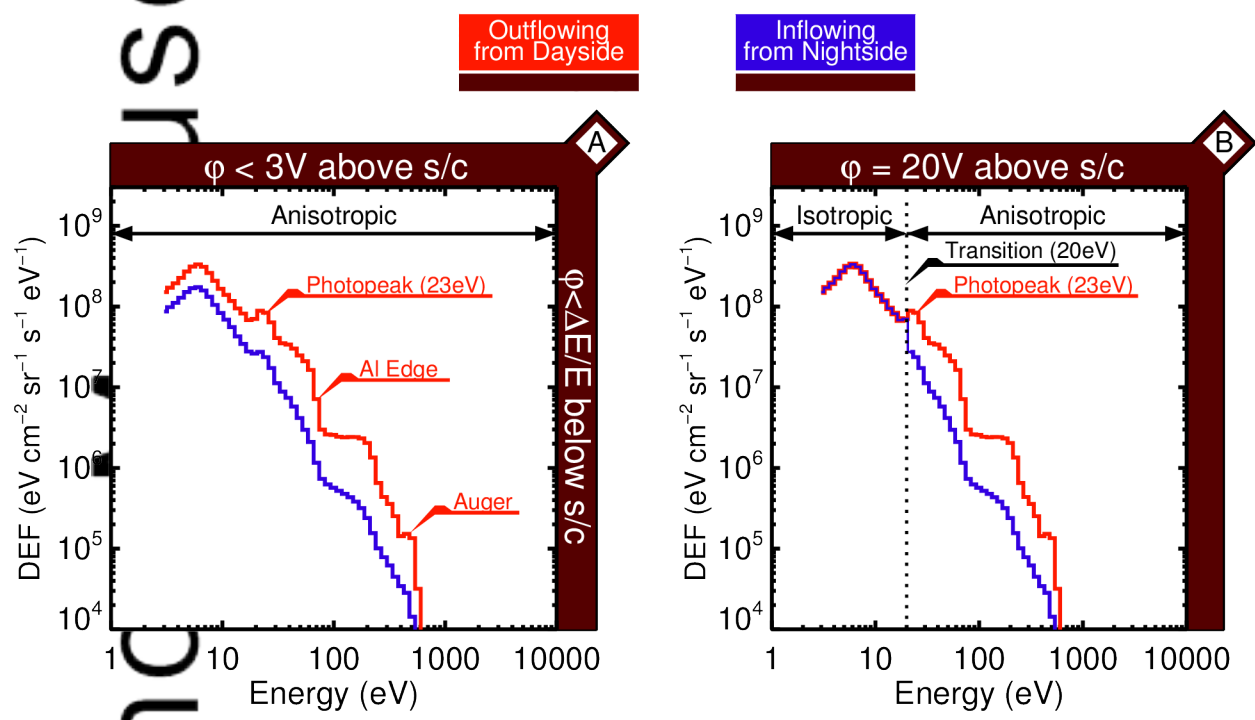
Orbit №	Date/Time (GMT)	Alt.(km)	SZA. (°)	Φ_{TOTAL} (V)	Φ_{SC} (V)	Φ_{Mars} (V)
795	2015-02-26 14:28:03	166	96	11.5±1.0	-1.1±0.2	12.6±1.1
795	2015-02-26 14:29:51	195	102	9.0±0.8	-0.7±0.1	9.7±0.8
798	2015-02-27 04:00:41	188	100	11.5±1.0	-1.9±0.4	13.4±1.1
798	2015-02-27 04:02:17	226	104	11.5±1.0	-3.9±0.8	15.4±1.3
799	2015-02-27 08:31:25	195	101	9.0±0.8	-1.1±0.2	10.1±0.8
799	2015-02-27 08:33:27	248	106	7.2±0.6	-1.9±0.4	9.1±0.7
800	2015-02-27 13:03:47	247	106	9.0±0.8	-1.5±0.3	10.5±0.9
805	2015-02-28 11:37:00	292	108	9.0±0.8	-1.5±0.3	10.5±0.9
818	2015-03-02 22:11:18	242	100	9.0±0.8	-1.3±0.3	10.3±0.8
821	2015-03-03 11:45:25	323	106	6.5±0.5	-1.2±0.2	7.7±0.6
Mean:		232 km	103°	9.3±0.8 V	-1.6±0.6 V	10.9±0.9 V

Figure 4. Magnetic and electron observations from the NASA *MAVEN* Mars Scout on orbit № 867. A - Spacecraft altitude; B - MAVEN Magnetometer observations in the Mars Solar Orbital coordinate system; Panel C - Electron spectrogram from MAVEN SWEA; and three electron scans (α, β, γ), showing energy vs. differential energy flux and energy vs. the ratio between outflowing and inflowing electron fluxes. Scan α - Closed magnetic field line with both footpoints in the dayside ionosphere; Scan β - Closed trans-terminator magnetic field line with no evidence of an electric potential drop; Scan γ - Closed trans-terminator magnetic field line with a 7V electric potential drop.

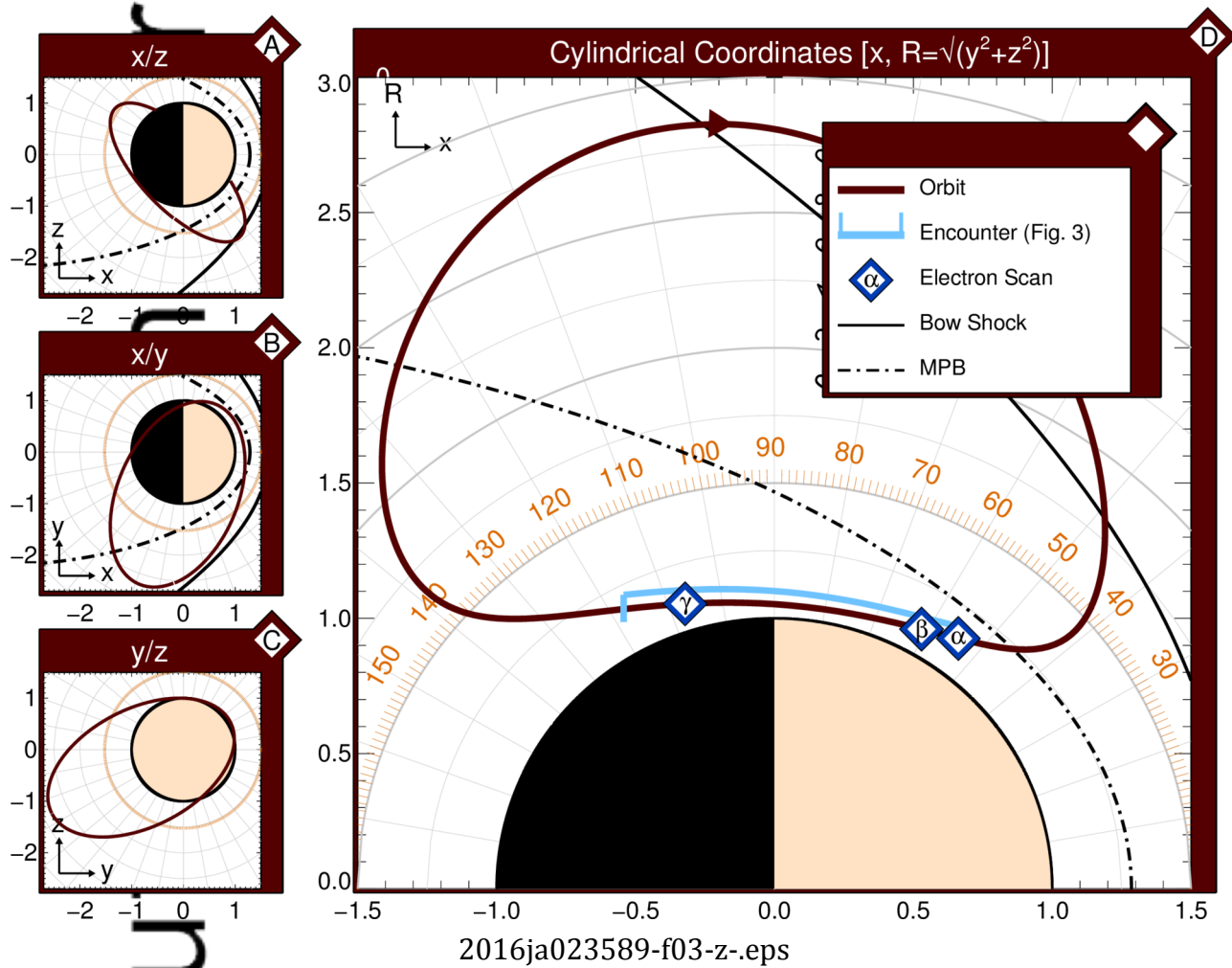
Figure 5. Full pitch-angle distributions for each of the three electron scans: Scan α - Closed magnetic field line with dayside photoelectron sources at both ionospheric footpoints; Scan β - Closed trans-terminator magnetic field line with a single source of photoelectrons on the dayside, and backscattered photoelectrons returning from the nightside resulting in an anisotropic distribution; Scan γ - Closed trans-terminator magnetic field line with a single source on the dayside ionosphere and semi-isotropy below 7eV resulting from electrostatic mirroring.

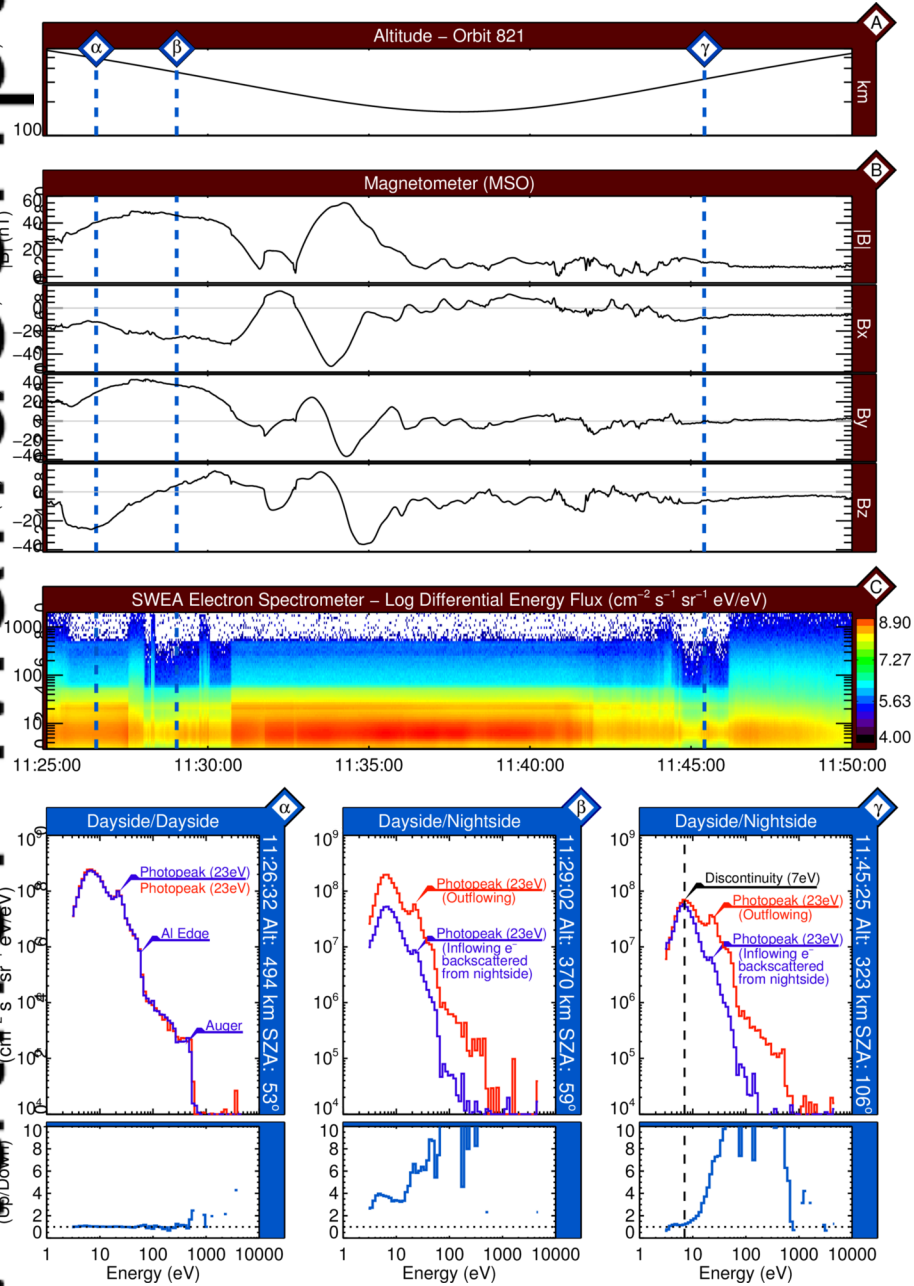


2016ja023589-f01-z-eps

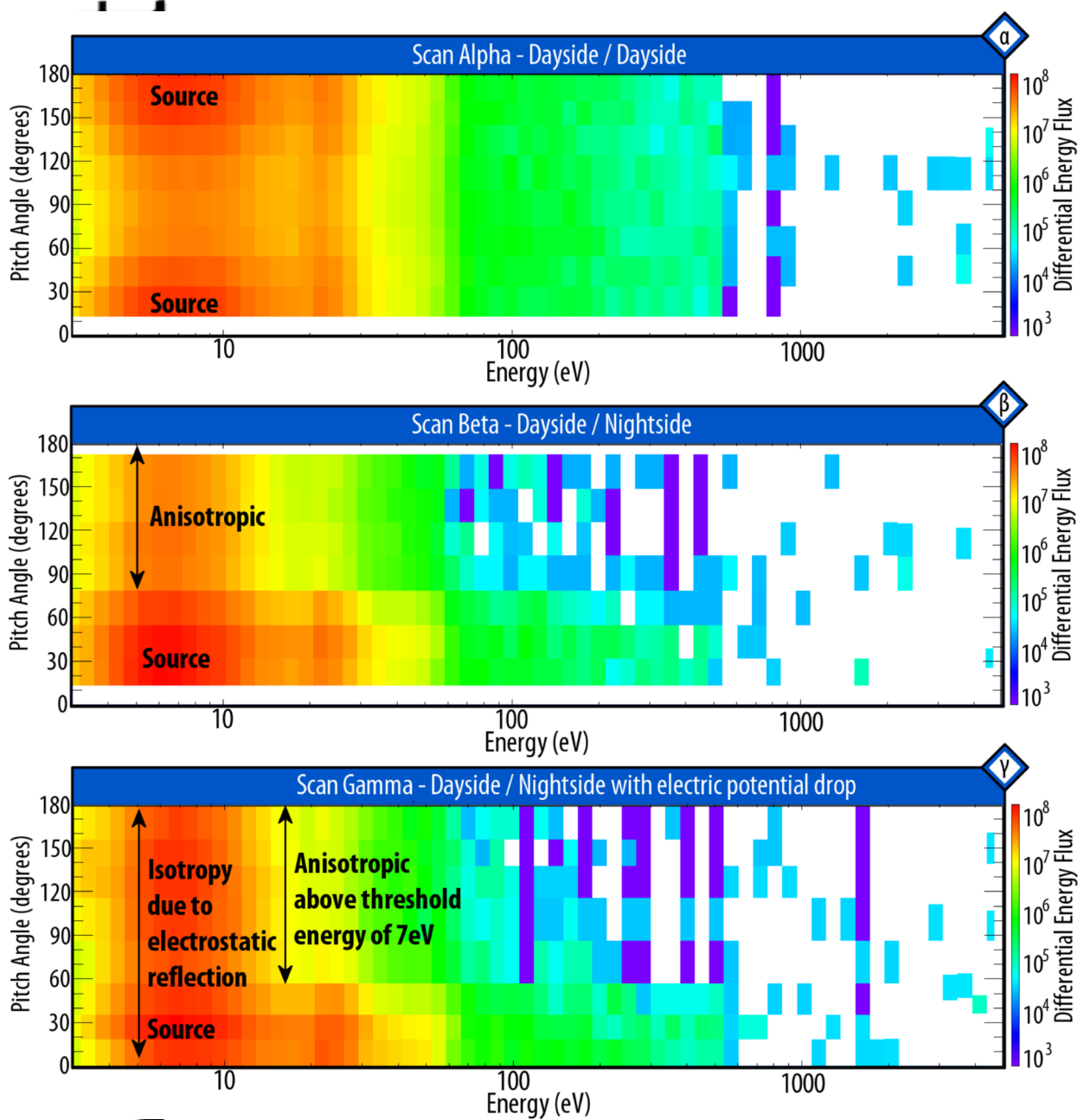


2016ja023589-f02-z-.eps





2016ja023589-f04-z-eps



2016ja023589-f05-z-.eps