# The influence of IMF clock angle on dayside flux transfer events at Mercury

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### Key Points:

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- Large statistical study of dayside FTEs at Mercury
- FTEs at Mercury more prevalent during periods of near-southward IMF

• FTEs form preferentially in the pre-noon sector of Mercury's dayside magnetopause

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#### 10 Abstract

Analysis of MESSENGER data has shown for the first time that the orientation of the Inter-

planetary Magnetic Field (IMF) in the magnetosheath of Mercury plays a crucial role in the

formation of flux transfer events (FTEs) at the dayside magnetopause. During the first 4 Her-

<sup>14</sup> mean years of MESSENGER's orbit around Mercury, we have identified 805 FTEs using mag-

netometer data. Under conditions of near-southward IMF, at least one FTE was detected on

nearly 70% of passes through the magnetopause but the observation rate during northward IMF

17 was less than 20%. FTEs were also observed preferentially in the pre-noon sector.

# 18 **1 Introduction**

Mercury was first discovered to have an intrinsic global magnetic field by Mariner 10 19 [Ness et al., 1974, 1975], and details of the nature of its magnetosphere were refined through 20 measurements made by the MErcury Surface, Space ENvironment, GEochemistry, and Rang-21 ing (MESSENGER) spacecraft when it became the first satellite to orbit Mercury [Anderson 22 et al., 2011, 2012]. Mercury's close proximity to the Sun exposes it to the extreme solar wind 23 conditions present at an orbital distance of 0.31-0.47 AU, including an interplanetary magnetic 24 field (IMF) strength of 20-40 nT [Blomberg et al., 2007], ~5 times that measured at Earth, and 25 solar wind number density of 30-70  $cm^{-3}$  [Blomberg et al., 2007], an order of magnitude greater 26 at Mercury [Baumjohann et al., 2006]. Furthermore, the planetary dipole moment at Mercury 27 is about 3 orders of magnitude lower than that at Earth [Johnson et al., 2012; Johnson and Hauck, 28 2016], with a value of 195 nT  $R_M^3$  [Anderson et al., 2011] (where  $R_M = 2440$  km is the ra-29 dius of Mercury). The combination of this weak planetary field and the solar wind conditions 30 means the Hermean magnetosphere is extremely small and strongly driven by variable con-31 ditions in the solar wind [Slavin et al., 2009]. The mean distance to the magnetopause at the 32 subsolar point is only  $1.45R_M$  [Winslow et al., 2013], but during extreme solar wind condi-33 tions the magnetopause can be compressed or eroded sufficiently to barely hold the solar wind 34 off the surface, with observations as low as  $1.03R_M$  [Slavin et al., 2014]. 35

Magnetic reconnection is an important factor in the interaction between the solar wind 36 and the magnetosphere, eroding the dayside magnetosphere [Slavin et al., 2010a; Heyner et al., 37 2016] and driving the Dungey cycle of magnetic flux circulation [Dungey, 1961; Imber and 38 Slavin, 2017], thus allowing entry of solar wind plasma into the magnetosphere [Raines et al., 39 2015]. At Earth, reconnection on the dayside magnetopause occurs at low latitude primarily 40 when the magnetic shear angle between the planetary field and the IMF in the magnetosheath 41 is high [e.g. Dungey, 1961; Fairfield and Cahill, 1966; Perreault and Akasofu, 1978; Sonnerup 42 et al., 1981]. Antiparallel reconnection at a single X-line connects magnetospheric field lines 43 to draped IMF in the magnetosheath. The newly open field lines are dragged away from the 44 reconnection site by the magnetosheath flow. Helical bundles of open magnetic flux, known 45 as flux transfer events (FTEs) [Russell and Elphic, 1978], are commonly observed at the mag-46 netopause of Earth, often with a large azimuthal extent [Fear et al., 2008]. Following the first 47 observation of FTEs at Earth by Russell and Elphic [1978], Lee and Fu [1985] suggested that 48 the observed bipolar signature in the magnetic field component normal to the magnetopause 49 that is attributed to FTEs could be explained by reconnection occurring at multiple parallel 50 X-lines. This produces a flux rope with its long axis aligned with the X-line, and connected 51 magnetically to both the IMF and the planetary magnetic field. 52

FTEs have been observed at Earth at all locations on the magnetopause under a wide 53 range of solar wind conditions by single spacecraft such as International Sun-Earth Explorer 54 1 (ISEE-1) [Kawano and Russell, 1996, 1997] and Interball-1 [e.g. Sibeck et al., 2005; Koro-55 tova et al., 2012], in addition to many multi-spacecraft missions, including Cluster [e.g. Fear 56 et al., 2008], Time History of Events and Macroscale Interactions during Substorms (THEMIS) 57 [e.g. Korotova et al., 2011; Trenchi et al., 2016], and most recently Magnetospheric Multiscale 58 (MMS) [e.g. Eastwood et al., 2012; Farrugia et al., 2016; Hasegawa et al., 2016], allowing for 59 accurate determination of the orientation and scale size of the FTEs. Such detailed measure-60

ments are not possible with the single MESSENGER spacecraft, however observations have nonetheless not only confirmed the presence of FTEs at Mercury, but also shown them to be ubiquitous in nature [*Slavin et al.*, 2009, 2010b,a, 2012; *Imber et al.*, 2014]. Indeed, studies by *Slavin et al.* [2012] and *Imber et al.* [2014] have demonstrated that FTEs at Mercury occur more frequently than those seen at Earth, and are considerably larger with respect to the size of the magnetosphere. This is attributed to the reconnection-driven formation of FTEs being greatly enhanced due to the stronger interaction between the IMF and the Hermean magnetic field.

69 One way of quantifying the reconnection rate is to calculate the ratio of inflow velocity at a reconnection site to the Alfvén velocity of the outflow [Sonnerup, 1974]. This dimen-70 sionless reconnection rate can also be expressed as a ratio of the component of the magnetic 71 field normal to the boundary to the total field just inside the magnetopause [Sonnerup et al., 72 1981]. At Earth, reported values vary considerably, ranging from as little as 0.01 [Fuselier et al., 73 2005] to  $\sim 0.1$  [Sonnerup et al., 1981]. Many of these values were obtained from case stud-74 ies of individual magnetopause crossings, however, and in the largest statistical study to date 75 Mozer and Retinò [2007] analysed 22 events and determined an average reconnection rate of 76 0.046. At Mercury, only one study has investigated this quantity at the dayside magnetopause. 77 DiBraccio et al. [2013] used measurements of the magnetic field for 43 magnetopause cross-78 ings, and calculated a mean dimensionless reconnection rate of 0.15, validating the theory of 79 stronger interactions between the planetary field and the IMF at Mercury [Slavin and Holzer, 80 1979]. However, DiBraccio et al. [2013] found that the dimensionless reconnection rate dis-81 played very little dependence on the magnetic shear angle between the two regimes, contrary 82 to similar investigations at Earth [e.g. Sonnerup, 1974]. They attributed this to a low Alfvén 83 Mach number,  $M_A$ , and low plasma  $\beta$  (the ratio of thermal pressure to magnetic pressure) in the Hermean magnetosheath. Under these conditions, a large plasma depletion layer forms due 85 to the pile-up of magnetic flux in the magnetosheath [Gershman et al., 2013], leading to en-86 hanced reconnection rates and enabling reconnection over a wider range of shear angles than 87 observed at the Earth. 88

In this paper we present a large statistical study of FTEs observed near the dayside magnetopause using data obtained by MESSENGER's Magnetometer [*Anderson et al.*, 2007] during the first four Hermean years after orbital insertion. Our analysis suggests that the formation of FTEs at Mercury exhibits a strong dependence on the orientation of the IMF, with a considerably enhanced production rate for magnetopause crossings during which the magnetic shear angle was large.

#### 95 **2** Observations

On 18 March 2011, MESSENGER orbital insertion placed the spacecraft into an eccen-96 tric, high-inclination orbit about Mercury with a period of 12 h. The orbital plane was fixed 97 in inertial space such that the periapsis precessed completely around the planet once every Her-98 mean year (88 days). In this study, we have used data obtained by the Magnetometer onboard 99 MESSENGER, which at full resolution provided 20 samples/s [Anderson et al., 2007], dur-100 ing the interval spanning orbital insertion until 9 March 2012. By including exactly 4 Hermean 101 years, we have ensured approximately even coverage of all magnetic local time (MLT) sec-102 tors over the duration of this study, with the exception of 19 orbits between 24 May and 2 June 103 2011, when the Magnetometer collected no data near the dayside magnetopause traversals. These 104 orbits are symmetric about 12 h MLT and confined to a small MLT range, however, so no dawn-105 dusk bias is introduced by the lack of data in this period. Furthermore, the number of miss-106 ing passes is small compared to the total number of passes in the affected MLT sectors, so no 107 biases have been introduced. Data are presented in the Mercury solar magnetospheric (MSM) 108 coordinate system, in which the X axis points towards the Sun, the origin is centered on the 109 internal dipole of Mercury and the Z axis is aligned with magnetic north. This coordinate sys-110 tem is then rotated to account for Mercury's changing orbital motion with respect to an av-111

erage solar wind velocity of 400  $km s^{-1}$ , producing the resultant aberrated MSM coordinate system (MSM').

Our focus in this study is the dayside magnetosphere, therefore the magnetic field data 114 have been examined for every encounter of MESSENGER with the magnetopause sunward 115 of X' = -0.5  $R_M$ . An example of a MESSENGER orbit is shown in Figures 1(e-f), with model 116 locations for the bow shock and magnetopause, as given by Winslow et al. [2013], and the com-117 ponents of the magnetic field measured by the MESSENGER magnetometer are shown in pan-118 els (a-d). Panels (g-l) show a subsection of these data, spanning the inbound crossings of the 119 120 bow shock and magnetopause on this orbit. Several large amplitude FTEs are present in the data, as indicated by the arrows in Figure 1j. 121

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#### 2.1 Identifying magnetopause crossings and flux transfer events

Every spacecraft pass through the dayside magnetopause during the time interval con-123 sidered was visually inspected for individual magnetopause crossings and FTE signatures in 124 the magnetic field data. A pass here refers to a traversal of the magnetopause region, during 125 which multiple individual magnetopause crossings may be observed. The magnetopause cross-126 ings were identified by a sudden large change in the magnetic field strength or, for cases when 127 the magnitude varied only slightly, by a rotation in the magnetic field vector. In both scenar-128 ios, the identification of crossings was aided by a significant reduction in the amplitude and 129 frequency of fluctuations in the magnetic field on the magnetospheric side of the magnetopause. 130 Flux transfer events were initially identified on the basis of a clear increase in the total field 131 strength compared to the background level, accompanied by a bipolar signature in one or more 132 field components. Throughout the period considered here, in 727 passes during which the mag-133 netopause was traversed sunward of X' = -0.5  $R_M$ , we identified a total of 1717 individual 134 magnetopause crossings and 805 FTEs for which the above conditions were satisfied. In the 135 306 passes on which these FTEs were observed, 818 individual magnetopause crossings were 136 identified, yielding an average observation rate of 0.98 FTEs per magnetopause crossing on 137 passes containing FTEs. 138

#### **3** Analysis

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#### 3.1 Magnetopause and FTE locations

The location of each of the 1717 magnetopause crossings identified in this work is projected into the MSM X' - Y' and X' - Z' planes in Figures 2a and 2b. Due to the highly elliptical polar orbit of the MESSENGER spacecraft, the inbound portion of the orbit through the dayside magnetosphere often passes through the northern magnetic cusp. The spacecraft therefore regularly skims the magnetopause at high northern latitudes, resulting in multiple detectable magnetopause crossings on a single orbit. Additionally, ongoing reconnection or variable solar wind conditions can result in a magnetopause that repeatedly moves back and forth over the spacecraft, again leading to the observation of multiple crossings on a single pass.

Figure 2a shows that crossings were observed approximately equally in all MLT sectors 149 in the dayside magnetosphere, and that on average the magnetopause crossings occurred near 150 to the location given by the Winslow et al. [2013] model for the majority of orbits considered 151 here. There appears to be a substantial spread in the distance of the observed crossings from 152 the model location, which is likely due to crossings occurring during a range of Hermean sea-153 sons, resulting in significant changes to the compression of the magnetosphere by the solar 154 wind between aphelion and perihelion [Zhong et al., 2015]. The location of the FTEs iden-155 tified in this study are presented in panels (c) and (d) as red circles, with the magnetopause 156 crossings indicated in grey for context. It can be seen that the majority of FTEs were observed 157 near local noon, and the approximately equal data coverage in MLT means this is manifested 158 as a greater percentage observation of FTEs within 3 h MLT of local noon. 159

Figures 2b and 2d show two distinct latitudinal groupings of both magnetopause crossings and FTEs, which can be attributed to orbital bias. The group near the subsolar point have been observed during MESSENGER's "hot season" orbits, when periapsis was on the dayside and the spacecraft passed outwards through the dayside magnetopause at low latitude. Half a Hermean year later, the orbital trajectory of MESSENGER carries it into the magnetopause thigh latitude, close to the northern cusp, producing the higher latitude group of magnetopause crossings and FTEs.

In a previous study of a smaller number of events over a different time period, Imber 167 et al. [2014] observed a larger number of FTEs in the dawn sector than the dusk, a bias that 168 is also present in these data. This is more apparent in Figure 3a, which shows that the largest 169 number of FTEs are seen at a magnetic local time of 10 h, with 288 FTEs observed between 170 9-11 h MLT compared to 238 between 13-15 h MLT. This asymmetry may be due to the un-171 usual conditions observed in the IMF during the period examined [James et al., 2017; Lock-172 wood et al., 2017], whereby in the majority of passes IMF  $B_X$  is positive, leading to a sim-173 ilar bias towards  $-B_Y$  due to the Parker spiral [Parker, 1958]. This in turn leads to increased 174 probability of near-antiparallel fields in the pre-noon sector of the portion of the magnetosphere 175 sampled by MESSENGER. 176

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#### 3.2 Influence of IMF clock angle on FTE formation

Many studies have investigated the parameters influencing dayside reconnection rates at 178 Earth [e.g. Akasofu, 1981; Mozer and Retinò, 2007; Milan et al., 2007, 2012; Newell et al., 2007], 179 however there has only been one such study at Mercury. DiBraccio et al. [2013] analysed the 180 magnetic field data from 43 magnetopause crossings to determine a dimensionless reconnec-181 tion rate, and concluded that for their dataset there was no significant variation with magnetic 182 shear angle. The FTEs observed in this study were formed by reconnection on the dayside mag-183 netopause, and have an average duration of 3.27s, calculated by recording the start and end 184 time of the bipolar signature of each event. This is similar to the  $\sim$ 2-3 s durations observed 185 by Imber et al. [2014] and Slavin et al. [2012]. Given the high velocities of these structures 186 observed at Earth, and the small spatial scale of the Hermean magnetosphere, it is reasonable 187 to assume that the IMF direction had not changed significantly from the time of formation of 188 the FTEs to their observation. In this study, we analyse the dependence of FTE observation 189 on IMF orientation. 190

The orientation of the magnetosheath field was recorded over 1 minute just outside the 191 outermost magnetopause crossing on each orbit to give a measurement of the clock angle in 192 the magnetosheath, where  $0^{\circ}$  is directed northwards and  $+90^{\circ}$  is directed towards  $+B_{Y'}$  and 193 the total number of FTEs in each  $30^{\circ}$  bin has been plotted in Figure 3b. In agreement with 194 studies at equivalent locations in the Earth's magnetosphere [Kawano and Russell, 1997; Sibeck 195 et al., 2005, e.g.], this shows a clear general trend towards greater FTE occurrence during in-196 tervals of near-southward IMF, and therefore nearly anti-parallel fields, although any poten-197 tial statistical bias introduced by multiple FTEs in a single pass or an uneven distribution of 198 observed IMF orientations needs to be accounted for. 199

A histogram of the occurrence frequency of the magnetosheath clock angle for every pass 200 on which at least 1 FTE was observed is presented in Figure 4a. Multiple FTEs observed on 201 a single crossing are therefore grouped into a single event, resulting in a similar distribution 202 to that presented in Figure 3 with some asymmetries removed. FTEs were observed on 306 203 of the 727 total passes inspected, during which 818 magnetopause crossings were detected, 204 and Figure 4b shows the distribution of clock angles observed across all magnetopause encoun-205 ters. The approximately equal coverage of all clock angle orientations indicates that variations 206 in observation rates cannot be attributed to sampling bias. By dividing the values in Figure 207 4a by those in Figure 4b we obtain the percentage occurrence of at least 1 FTE for each clock 208 angle, as indicated in Figure 4c. For clock angles close to zero, indicating a magnetosheath 209 magnetic field pointing approximately along the positive  $B_{Z'}$  axis, FTEs have been detected 210

on fewer than 20% of passes, whereas for near-southward IMF the observation rate increases
to nearly 70%. During periods of northward IMF, the reconnection X-line is expected to exist tailward of the cusp regions, therefore we would not expect to observe any FTEs generated at low latitudes near the dayside magnetopause. However, MESSENGER's orbit samples
significant portions of the high latitude magnetosphere, so we would still expect to observe
FTEs that have formed under northward IMF if reconnection is taking place in these locations.

Out of a total of 727 passes, events exhibiting the required magnetic field signature were 217 observed on 306, although many crossings contained multiple events. Considering how ubiq-218 uitous FTEs have been found to be at Mercury in previous studies [Slavin et al., 2012; Imber et al., 2014], this ratio is perhaps lower than expected. However, the formation of FTEs at the 220 dayside magnetopause has been shown for the first time to be significantly less likely during 221 northward IMF, and these orientations contribute a substantial portion of the data examined 222 here. Therefore, the higher ratios seen in previous studies could be explained by an IMF ori-223 entation during those periods that is more favourable for FTE formation. Furthermore, in re-224 quiring a clear increase in the core field component, we have restricted our sample to those 225 events for which MESSENGER entered the flux rope directly. As a result, many events exhibiting similar features have not been included, such as the travelling compression regions 227 identified by Slavin et al. [2012]. 228

In addition to the effect of the IMF clock angle on the observation rate of FTEs in the Hermean magnetosphere, the events in this study were also found to exhibit a small dependence on the strength of the magnetosheath field. In general, a stronger magnetosheath field resulted in the observation of more FTEs per pass, reaching a maximum at  $\sim$ 140 nT, above which there were too few occurrences for results to be statistically significant. However, this increase is only small, resulting in a trend that is considerably less significant than the clock angle effects presented here.

There are several reasons why the results presented here contrast so strongly with those 236 observed by DiBraccio et al. [2013]. First of all, although the formation of FTEs requires re-237 connection, the reconnection rate itself is not measured here, so it is difficult to directly com-238 pare the results. Secondly, the sample size used by DiBraccio et al. [2013] was considerably 239 smaller than that utilised here. The large dataset investigated over a long time interval in this 240 study is likely to have averaged out the effects of other parameters, thereby producing a more 241 accurate reflection of how the IMF orientation alone influences the observation rate of FTEs 242 at Mercury. Furthermore, the analysis performed by DiBraccio et al. [2013] utilised only cross-243 ings with a well defined normal direction to the magnetopause, as determined from minimum 244 variance analysis of the magnetic field data. The presence of FTEs during a crossing may re-245 sult in a poorly defined magnetopause normal, therefore crossings containing FTEs may have 246 been excluded from their analysis, possibly leading to a calculation of the reconnection rate 247 only under conditions less favourable to FTE formation. 248

#### 249 4 Conclusions

<sup>250</sup> 727 passes of magnetic field data taken by the MESSENGER spacecraft were visually <sup>251</sup> inspected for flux transfer event signatures near the dayside magnetopause encounters. Obser-<sup>252</sup> vation of FTEs is shown to be strongly dependent on the orientation of the IMF in the mag-<sup>253</sup> netosheath. FTEs with clear signatures were identified in 306 of the 727 passes through the <sup>254</sup> magnetopause sunward of MSM  $X' = -0.5 R_M$ , with a total of 805 FTEs observed. During <sup>255</sup> periods of near-southward IMF at least 1 FTE was observed on nearly 70% of passes, whereas <sup>256</sup> during northward IMF the observation rate is less than 20%.

The spatial distribution of the identified FTEs peaks at a magnetic local time of 10 h, and more FTEs were observed throughout the pre-noon sector than post-noon, corroborating the results of *Imber et al.* [2014]. Additionally, the identified magnetopause crossings agree well with the *Winslow et al.* [2013] model for large parts of the dayside magnetosphere. Some crossings on the dawn and dusk flanks are seen closer to Mercury than predicted, but these
 occurred during perihelion, when stronger solar wind forcing produced a more compressed magnetosphere.

The upcoming BepiColombo mission will provide the opportunity to expand further on the analysis performed herein, due to improved instruments including a magnetometer with even greater temporal resolution than the MESSENGER magnetometer [*Glassmeier et al.*, 2010] and additional plasma measurements [*Saito et al.*, 2010]. Additionally, the orbital paths will provide considerably greater magnetopause coverage, allowing for the observation of FTEs across a much larger range of latitudes, including for the first time significant coverage of the dayside magnetopause in the southern hemisphere.

Figure 1. Magnetic field data in MSM' coordinates for a complete MESSENGER orbit. Panels (a)-(d)

- show  $B_{X'}$ ,  $B_{Y'}$ ,  $B_{Z'}$  and |B| respectively. The spacecraft trajectory during the course of this orbit is projected onto the (e)  $Y' \cdot X'$  and (f)  $Z' \cdot X'$  planes. Model locations of the bow shock (blue) and magnetopause
- (green), as given by the *Winslow et al.* [2013] models, are also shown. Panels (g)-(l) show the same as (a)-(f)
   above, but for a shorter interval spanning the inbound bow shock and magnetopause crossings with some FTE
   signatures visible, as indicated by the arrows.
- Figure 2. Locations of the magnetopause crossings in this study, projected onto the (a) MSM X' Y' and (b) MSM X' - Z' planes. The locations of the identified FTEs are shown in the same projections in panels (c) and (d), with the magnetopause crossings also indicated in grey for comparison. The model magnetopause location predicted by *Winslow et al.* [2013] is indicated by the dashed line.
- **Figure 3.** Histograms showing (a) the locations of the observed FTEs in MLT and (b) how the total number of FTEs observed varies with the clock angle of the IMF in the magnetosheath. The total number of FTEs, n, is also indicated.
- Figure 4. Histograms showing (a) the number of passes during each IMF orientation for which at least 1
   FTE was observed, (b) the occurrence of each clock angle, and (c) percentage of magnetopause crossings
   under each IMF orientation during which at least 1 FTE was observed. The number of passes with at least 1
- <sup>287</sup> FTE, n, is also indicated.

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The event list is available on request.

# 295 References

Akasofu, S.-I. (1981), Energy coupling between the solar wind and the magnetosphere,
 *Space Sci. Rev.*, 28(2), 121–190, doi:10.1007/BF00218810.

298 299 300	Anderson, B. J., M. H. Acuña, D. A. Lohr, J. Scheifele, A. Raval, H. Korth, and J. A. Slavin (2007), The Magnetometer Instrument on MESSENGER, <i>Space Science Reviews</i> , 131(1), 417–450. doi:10.1007/s11214-007-9246-7.
301	Anderson B I C L Johnson H Korth M E Purucker R M Winslow I A Slavin
302	S. C. Solomon, R. L. McNutt, J. M. Raines, and T. H. Zurbuchen (2011). The global
303	magnetic field of Mercury from MESSENGER orbital observations., Science (New York,
304	<i>N.Y.</i> ), <i>333</i> (6051), 1859–62, doi:10.1126/science.1211001.
305	Anderson, B. J., C. L. Johnson, H. Korth, R. M. Winslow, J. E. Borovsky, M. E. Purucker,
306	J. A. Slavin, S. C. Solomon, M. T. Zuber, and R. L. McNutt (2012), Low-degree struc-
307	ture in Mercury's planetary magnetic field. <i>Journal of Geophysical Research: Planets</i> .
308	<i>117</i> (E12), doi:10.1029/2012JE004159.
309	Baumjohann, W., A. Matsuoka, K. Glassmeier, C. Russell, T. Nagai, M. Hoshino,
310	T. Nakagawa, A. Balogh, J. Slavin, R. Nakamura, and W. Magnes (2006),
311	The magnetosphere of mercury and its solar wind environment: Open issues
312	and scientific questions, Advances in Space Research, 38(4), 604 – 609, doi:
313	http://dx.doi.org/10.1016/j.asr.2005.05.117, mercury, Mars and Saturn.
314	Blomberg, L. G., J. A. Cumnock, KH. Glassmeier, and R. A. Treumann (2007), Plasma
315	Waves in the Hermean Magnetosphere, Space Sci. Rev., 132(2-4), 575-591, doi:
316	10.1007/s11214-007-9282-3.
317	DiBraccio, G. A., J. A. Slavin, S. A. Boardsen, B. J. Anderson, H. Korth, T. H. Zur-
318	buchen, J. M. Raines, D. N. Baker, R. L. McNutt, and S. C. Solomon (2013), MES-
319	SENGER observations of magnetopause structure and dynamics at Mercury, Journal of
320	Geophysical Research: Space Physics, 118(3), 997-1008, doi:10.1002/jgra.50123.
321	Dungey, J. W. (1961), Interplanetary Magnetic Field and the Auroral Zones, Phys. Rev.
322	Lett., 6, 47-48, doi:10.1103/PhysRevLett.6.47.
323	Eastwood, J. P., T. D. Phan, R. C. Fear, D. G. Sibeck, V. Angelopoulos, M. Ieroset, and
324	M. A. Shay (2012), Survival of flux transfer event (FTE) flux ropes far along the tail
325	magnetopause, J. Geophys. Res. Sp. Phys., 117(8), doi:10.1029/2012JA017722.
326	Fairfield, D. H., and L. J. Cahill (1966), Transition region magnetic field and polar mag-
327	netic disturbances, J. Geophys. Res., 71(1), 155–169, doi:10.1029/JZ071i001p00155.
328	Farrugia, C. J., B. Lavraud, R. B. Torbert, M. Argall, I. Kacem, W. Yu, L. Alm, J. Burch,
329	C. T. Russell, J. Shuster, J. Dorelli, J. P. Eastwood, R. E. Ergun, S. Fuselier, D. Ger-
330	shman, B. L. Giles, Y. V. Khotyaintsev, P. A. Lindqvist, H. Matsui, G. T. Marklund,
331	T. D. Phan, K. Paulson, C. Pollock, and R. J. Strangeway (2016), Magnetospheric
332	Multiscale Mission observations and non-force free modeling of a flux transfer event
333	immersed in a super-Alfvénic flow, Geophys. Res. Lett., 43(12), 6070–6077, doi:
334	10.1002/2016GL068758.
335	Fear, R. C., S. E. Milan, A. N. Fazakerley, E. A. Lucek, S. W. H. Cowley, and I. Dan-
336	douras (2008), The azimuthal extent of three flux transfer events, Annales Geophysicae,
337	26(8), 2353–2369, doi:10.5194/angeo-26-2353-2008.
338	Fuselier, S. A., K. J. Trattner, S. M. Petrinec, C. J. Owen, and H. Réme (2005), Comput-
339	ing the reconnection rate at the Earth's magnetopause using two spacecraft observations,
340	J. Geophys. Res. Sp. Phys., 110(A6), A06,212, doi:10.1029/2004JA010805.
341	Gersnman, D. J., J. A. Slavin, J. M. Kaines, I. H. Zurbuchen, B. J. Anderson, H. Kortn, D. N. Dahar and S. C. Salarman (2012). Magnetic flux gilaxy and plasma deplation
342	D. N. Baker, and S. C. Solomon (2013), Magnetic flux plieup and plasma depiction
343	doi:10.1002/20131A.010244
344	Classemaier K H H II Auster D Havner K Okrafka C Carr G Berghofer
345	B I Anderson A Balogh W Baumiohann D Cargill U Christenson M Dalva
346	M Dougherty K H Fornacon T S Horbury F A Lucek W Magnes M Man
349	dea A Matsuoka M Matsushima II Motschmann R Nakamura V Narita
349	H. O'Brien, I. Richter, K. Schwingenschuh H. Shibuya I. A. Slavin C. Sotin B. Stoll
350	H. Tsunakawa, S. Vennerstrom, J. Vogt, and T. Zhang (2010). The fluxgate magne-
351	tometer of the BepiColombo Mercury Planetary Orbiter, Planet. Space Sci., 58(1-2),

352	287–299, doi:10.1016/j.pss.2008.06.018.
353	Hasegawa, H., N. Kitamura, Y. Saito, T. Nagai, I. Shinohara, S. Yokota, C. J. Pollock,
354	B. L. Giles, J. C. Dorelli, D. J. Gershman, L. A. Avanov, S. Kreisler, W. R. Paterson,
355	M. O. Chandler, V. Coffey, J. L. Burch, R. B. Torbert, T. E. Moore, C. T. Russell, R. J.
356	Strangeway, G. Le, M. Oka, T. D. Phan, B. Lavraud, S. Zenitani, and M. Hesse (2016),
357	Decay of mesoscale flux transfer events during quasi-continuous spatially-extended
358	reconnection at the magnetopause, Geophys. Res. Lett., doi:10.1002/2016GL069225.
359	Heyner, D., C. Nabert, E. Liebert, and KH. Glassmeier (2016), Concerning reconnection-
360	induction balance at the magnetopause of Mercury, J. Geophys. Res. Sp. Phys., 121(4),
361	2935–2961, doi:10.1002/2015JA021484.
362	Imber, S. M., and J. A. Slavin (2017), MESSENGER Observations of Magnetotail Load-
363	ing and Unloading: Implications for Substorms at Mercury. J. Geophys. Res. Sp. Phys.
364	doi:10.1002/2017JA024332.
365	Imber, S. M., J. A. Slavin, S. A. Boardsen, B. J. Anderson, H. Korth, R. L. McNutt, and
366	S. C. Solomon (2014). MESSENGER observations of large dayside flux transfer events:
367	Do they drive Mercury's substorm cycle? <i>Journal of Geophysical Research: Space</i>
368	<i>Physics</i> , 119(7), 5613–5623, doi:10.1002/2014JA019884, 2014JA019884,
369	James, M. K., S. M. Imber, E. J. Bunce, T. K. Yeoman, M. Lockwood, M. J. Owens,
370	and J. A. Slavin (2017). Interplanetary magnetic field properties and variability near
371	Mercury's orbit, J. Geophys. Res. Sp. Phys., doi:10.1002/2017JA024435.
372	Johnson C. L. and S. A. Hauck (2016). A whole new Mercury: MESSENGER reveals a
373	dynamic planet at the last frontier of the inner solar system. J. Geophys. Res. Planets.
374	doi:10.1002/2016JE005150.
375	Johnson, C. L., M. F. Purucker, H. Korth, B. J. Anderson, R. M. Winslow, M. M. H.
376	Al Asad, J. A. Slavin, I. I. Alexeev, R. J. Phillips, M. T. Zuber, and S. C. Solomon
377	(2012), MESSENGER observations of Mercury's magnetic field structure. <i>Journal of</i>
378	Geophysical Research: Planets, 117(E12), doi:10.1029/2012JE004217.
379	Kawano, H., and C. T. Russell (1996), Survey of flux transfer events observed with the
380	ISEE 1 spacecraft: Rotational polarity and the source region, J. Geophys. Res. Sp. Phys.,
381	<i>101</i> (A12), 27,299–27,308, doi:10.1029/96JA02703.
382	Kawano, H., and C. T. Russell (1997), Survey of flux transfer events observed with the
383	ISEE 1 spacecraft: Dependence on the interplanetary magnetic field, J. Geophys. Res.
384	Sp. Phys., 102(A6), 11,307–11,313, doi:10.1029/97JA00481.
385	Korotova, G. I., D. G. Sibeck, A. Weatherwax, V. Angelopoulos, and V. Styazhkin (2011),
386	THEMIS observations of a transient event at the magnetopause, J. Geophys. Res. Sp.
387	Phys., 116(7), 1–13, doi:10.1029/2011JA016606.
388	Korotova, G. I., D. G. Sibeck, and V. I. Petrov (2012), Interball-1 observations of flux
389	transfer events, Ann. Geophys., 30(10), 1451-1462, doi:10.5194/angeo-30-1451-2012.
390	Lee, L. C., and Z. F. Fu (1985), A theory of magnetic flux transfer at the Earth's magne-
391	topause, Geophys. Res. Lett., 12(2), 105-108, doi:10.1029/GL012i002p00105.
392	Lockwood, M., M. J. Owens, S. M. Imber, M. K. James, E. J. Bunce, and T. K. Yeo-
393	man (2017), Coronal and heliospheric magnetic flux circulation and its relation
394	to open solar flux evolution, J. Geophys. Res. Sp. Phys., 122(6), 5870-5894, doi:
395	10.1002/2016JA023644.
396	Milan, S. E., G. Provan, and B. Hubert (2007), Magnetic flux transport in the Dungey
397	cycle: A survey of dayside and nightside reconnection rates, J. Geophys. Res. Sp. Phys.,
398	112(1), doi:10.1029/2006JA011642.
399	Milan, S. E., J. S. Gosling, and B. Hubert (2012), Relationship between interplanetary
400	parameters and the magnetopause reconnection rate quantified from observations of the
401	expanding polar cap, J. Geophys. Res. Sp. Phys., 117(3), doi:10.1029/2011JA017082.
402	Mozer, F. S., and A. Retinò (2007), Quantitative estimates of magnetic field reconnection
403	properties from electric and magnetic field measurements, J. Geophys. Res., 112(A10),
404	A10,206, doi:10.1029/2007JA012406.

405	Ness, N. F., K. W. Behannon, R. P. Lepping, Y. C. Whang, and K. H. Schatten (1974),
406	(New York NY) 185(4146) 151–160 doi:10.1126/science 185.4146.151
407	Ness, N. F., K. W. Behannon, R. P. Lepping, and Y. C. Whang (1975). The mag-
409	netic field of Mercury, 1, Journal of Geophysical Research, 80(19), 2708–2716, doi:
410	10.1029/JA082i019p02828.
411	Newell, P. T., T. Sotirelis, K. Liou, CI. Meng, and F. J. Rich (2007), A nearly universal
412	solar wind-magnetosphere coupling function inferred from 10 magnetospheric state
413	variables, J. Geophys. Res. Sp. Phys., 112(A1), doi:10.1029/2006JA012015.
414 415	Parker, E. N. (1958), Dynamics of the Interplanetary Gas and Magnetic Fields., <i>Astrophys. J.</i> , 128, 664, doi:10.1086/146579.
416	Perreault, P., and SI. Akasofu (1978), A study of geomagnetic storms, <i>Geophys. J. R.</i>
417	Astron. Soc., 54(3), 547–573, doi:10.1111/j.1365-246X.1978.tb05494.x.
418	Raines, J. M., G. A. DiBraccio, T. A. Cassidy, D. C. Delcourt, M. Fujimoto, X. Jia,
419 420	V. Mangano, A. Milillo, M. Sarantos, J. A. Slavin, and P. Wurz (2015), Plasma Sources in Planetary Magnetospheres: Mercury, doi:10.1007/s11214-015-0193-4.
421	Russell, C. T., and R. C. Elphic (1978), Initial ISEE magnetometer results - Magnetopause
422	observations, Space Science Reviews, 22, 681–715, doi:10.1007/BF00212619.
423	Saito, Y., J. Sauvaud, M. Hirahara, S. Barabash, D. Delcourt, T. Takashima, and
424	K. Asamura (2010), Scientific objectives and instrumentation of Mercury Plasma
425	Particle Experiment (MPPE) onboard MMO, <i>Planet. Space Sci.</i> , 58(1-2), 182–200,
426	doi:10.1016/j.pss.2008.06.003.
427	Sibeck, D. G., G. I. Korotova, V. Petrov, V. Styazhkin, and T. J. Rosenberg (2005), Flux
428	has represented by the magnetopause. The robust values, Ann. Geo- $has 23(11)$ 3549–3559 doi:10.5194/angeo-23-3549-2005
430	Slavin, J. A., and R. E. Holzer (1979). The effect of erosion on the solar wind stand-off
431	distance at Mercury, J. Geophys. Res., 84(A5), 2076, doi:10.1029/JA084iA05p02076.
432	Slavin, J. A., M. H. Acuña, B. J. Anderson, D. N. Baker, M. Benna, S. A. Boardsen,
433	G. Gloeckler, R. E. Gold, G. C. Ho, H. Korth, S. M. Krimigis, R. L. McNutt, J. M.
434	Raines, M. Sarantos, D. Schriver, S. C. Solomon, P. Trávníček, and T. H. Zurbuchen
435	(2009), MESSENGER Observations of Magnetic Reconnection in Mercury's Magneto-
436	sphere, <i>Science</i> , 324(5927), 606–610, doi:10.1126/science.1172011.
437	Slavin, J. A., B. J. Anderson, D. N. Baker, M. Benna, S. A. Boardsen, G. Gloeckler,
438	R. E. Gold, G. C. Ho, H. Korth, S. M. Krimigis, R. L. McNutt, L. R. Nittler, J. M. Paines, M. Sarantos, D. Schriver, S. C. Solomon, P. D. Starr, P. M. Trávnícek, and
439	T H Zurbuchen (2010a) MESSENGER observations of extreme loading and unloading
441	of Mercury's magnetic tail. <i>Science</i> , 329(5992), 665–8, doi:10.1126/science.1188067.
442	Slavin, J. A., R. P. Lepping, CC. Wu, B. J. Anderson, D. N. Baker, M. Benna, S. A.
443	Boardsen, R. M. Killen, H. Korth, S. M. Krimigis, W. E. McClintock, R. L. McNutt,
444	M. Sarantos, D. Schriver, S. C. Solomon, P. Trávnícek, and T. H. Zurbuchen (2010b),
445	MESSENGER observations of large flux transfer events at Mercury, Geophysical Re-
446	search Letters, 37(2), doi:10.1029/2009GL041485, 102105.
447	Slavin, J. A., S. M. Imber, S. A. Boardsen, G. A. DiBraccio, T. Sundberg, M. Saran-
448	tos, T. Nieves-Chinchilla, A. Szabo, B. J. Anderson, H. Korth, T. H. Zurbuchen, J. M.
449	Kaines, C. L. Jonnson, K. M. Winslow, K. M. Killen, R. L. McNutt, and S. C. Solomon (2012) MESSENCER absorbations of a flux transfer event absorber of Mersen at Merse
450	(2012), MESSENGER OBSERVATIONS OF a HUX-transfer-event snower at Mercury, Jour- nal of Geophysical Research: Space Physics 117(A12), doi:10.1020/20121A017026
452	a00M06.
453	Slavin, J. A., G. A. DiBraccio, D. J. Gershman, S. M. Imber, G. K. Poh, J. M. Raines
454	T. H. Zurbuchen, X. Jia, D. N. Baker, KH. Glassmeier, S. A. Livi, S. A. Boardsen.
455	T. A. Cassidy, M. Sarantos, T. Sundberg, A. Masters, C. L. Johnson, R. M. Winslow,
456	B. J. Anderson, H. Korth, R. L. McNutt, and S. C. Solomon (2014), MESSENGER
457	observations of Mercury's dayside magnetosphere under extreme solar wind con-
458	ditions, Journal of Geophysical Research: Space Physics, 119(10), 8087-8116, doi:

- 459 10.1002/2014JA020319, 2014JA020319.
- Sonnerup, B. U. Ö. (1974), Magnetopause reconnection rate, *J. Geophys. Res.*, 79(1),
   1546–1549, doi:10.1029/JA079i010p01546.
- Sonnerup, B. U. Ö., G. Paschmann, I. Papamastorakis, N. Sckopke, G. Haerendel, S. J.
  Bame, J. R. Asbridge, J. T. Gosling, and C. T. Russell (1981), Evidence for magnetic
  field reconnection at the Earth's magnetopause, *J. Geophys. Res. Sp. Phys.*, 86(A12),
- <sup>465</sup> 10,049–10,067, doi:10.1029/JA086iA12p10049.
- Trenchi, L., R. C. Fear, K. J. Trattner, B. Mihaljcic, and A. N. Fazakerley (2016), A se quence of flux transfer events potentially generated by different generation mechanisms,
   *J. Geophys. Res. A Sp. Phys.*, *121*(9), 8624–8639, doi:10.1002/2016JA022847.
- Winslow, R. M., B. J. Anderson, C. L. Johnson, J. A. Slavin, H. Korth, M. E. Purucker,
- 470 D. N. Baker, and S. C. Solomon (2013), Mercury's magnetopause and bow shock from
- 471 MESSENGER Magnetometer observations, Journal of Geophysical Research: Space
- 472 Physics, 118(5), 2213–2227, doi:10.1002/jgra.50237.
- Zhong, J., W. X. Wan, Y. Wei, J. A. Slavin, J. M. Raines, Z. J. Rong, L. H. Chai, and
   X. H. Han (2015), Compressibility of Mercury's dayside magnetosphere, *Geophys. Res. Lett.*, 42(23), 10,135–10,139, doi:10.1002/2015GL067063.

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