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4 **On-Demand Reconfiguration of Nanomaterials—When Electronics Meets Ionics**

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2 switching, resistive switching

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4 Rapid advances in the semiconductor industry, driven largely by device scaling, are now approaching
5 fundamental physical limits and face severe power, performance, and cost constraints.

6 Multifunctional materials and devices may lead to a paradigm shift towards new, intelligent, and
7 efficient computing systems, and are being extensively studied. This Review attempts to examine

8 how, by controlling the internal ion distribution in a solid-state film, a material's chemical

9 composition and physical properties can be reversibly reconfigured using an applied electric field, at

10 room temperature and after device fabrication. Reconfigurability is observed in a wide range of

11 materials, including commonly used dielectric films, and has led to the development of new device

12 concepts such as resistive random-access memory (RRAM). Physical reconfigurability further allows

13 memory and logic operations to be merged in the same device for efficient in-memory computing

14 and neuromorphic computing systems. By directly changing the chemical composition of the

15 material, coupled electrical, optical, and magnetic effects can also be obtained. A survey of recent

16 fundamental material and device studies that reveal the dynamic ionic processes is included, along

17 with discussions on systematic modeling efforts, device and material challenges, and future research

18 directions.

19

20 **1. Introduction**

1 The rapid growth of modern electronics has had unparalleled impact on the ways humans work,
2 live and socialize in the past a few decades. This exponential growth has mostly been driven by
3 device scaling, *i.e.* constant shrinking of the size of transistors that form the basic building block of
4 integrated circuits, following Moore's law.^[1] However, continued geometrical scaling will likely come
5 to a halt within the next decade as the device size approaches fundamental physical limits at ca. ~
6 5nm. New materials that offer new functionalities and provide new state variables that can be used
7 for computing and data storage are of great interest to maintain the growth of the semiconductor
8 industry.

9 A close examination of today's electronics systems reveals that they are mostly based on
10 "functional reconfigurable" materials and devices. That is, the device's electrical and optical
11 properties can be modulated upon the application of a stimulation, normally an electric field, but the
12 change is temporary and the device will recover its original form after the stimulation removal. For
13 example, light-emitting diodes (LEDs) can emit light when a voltage is applied but the material itself
14 is not modified during the process. Light emission will stop when the voltage is removed. Similarly, a
15 transistor can be turned "on" by a gate voltage but will return to the "off" state when the gate
16 voltage is removed. On the other hand, if a material is not only functional reconfigurable, but also
17 "physically reconfigurable" at the atomic scale, its electrical, optical, mechanical and magnetic
18 properties can be modified *on the fly* in a *non-volatile yet reversible* fashion, thus can potentially
19 facilitate a paradigm shift in building highly efficient computing systems. For example, compute-in-
20 memory and compute-with-memory concepts have received tremendous interest recently as an
21 alternative architecture to perform "big data" tasks compared to conventional computing
22 architectures. However, the performance of these memory-first approaches is often limited by the

1 properties of the memory, and by the ability to integrate large amounts of memory directly with the
2 logic components. Emerging memories based on physical reconfiguration effects, such as resistive
3 random-access memory (RRAM), offer advantages of nonvolatile storage, high-density, and fast and
4 random access, and can dramatically simplify the memory hierarchy and improve the system
5 performance.^[1-4] In an RRAM, the different stored data are represented by the different resistance
6 values of the storage material, through physical ion migrations and associated chemical redox
7 processes within the film.^[2,5,6] By directly changing the properties of the storage material itself, the
8 device also directly modulates the signals propagating through it, thus allowing logic functions to be
9 performed at the same physical locations as the memory element, leading to significant speed and
10 energy benefits.^[4]

11 For any device, physical reconfiguration involves moving matter (*i.e.* atoms) within the device. In
12 this review, we focus on devices where the reconfiguration of the material and device is achieved
13 through electric field driven redistribution of ions over large distances (compared to the lattice
14 constant/atomic spacing), allowing the devices and circuits to be reconfigured conveniently *in-situ*,
15 through simple changes in bias voltages^[2,6] One representative example of such reconfigurable
16 devices is RRAM, which exhibits resistive switching effects as a result of the reversible ion migration.
17 In this Review, we aim to first provide a summary of the coupled electronic and ionic processes that
18 govern resistive switching effects, by analyzing recent experimental observations and theoretical
19 simulations. Afterwards, applications of these reconfigurable materials and devices that offer
20 tunable electrical, optical, and magnetic properties will be introduced, with a focus on memory,
21 logic, and neuromorphic computing systems. Finally, discussions on new research directions that can
22 lead to promising multifunctional systems will be provided.

1

2 **2. Coupled electronics and ionics: mechanisms**

3 In conventional electronic devices such as transistors, the information is processed by controlling
4 the mobile carriers, *e.g.*, electrons, schematically shown in **Figure 1**. Ions such as dopants are
5 necessary to build the device during fabrication, but are not actively modulated/programmed during
6 device operation. Thus, no physical changes to the devices and materials are involved during device
7 operation. On the contrary, physical reconfiguration of a material requires moving masses (atoms or
8 ions) during the device operation. Although devices based on ionic processes have been extensively
9 studied for different types of applications (*e.g.* batteries), these devices are normally based on
10 materials that are incompatible with semiconductor electronics and have not been seriously
11 considered as electronic devices. For example, fast ion transport and storage are desirable for
12 conventional ionic devices such as Li-ion batteries and supercapacitor devices,^[7,8] while electronic
13 effects are considered secondary. As a result, traditionally, electronic devices and ionic devices
14 utilize distinct classes of materials (the former aiming for materials with high electronic mobility and
15 essentially zero ion mobility, while the latter aiming for materials with high ionic mobility) and target
16 dramatically different applications.^[9-11]

17 Recent advances in thin-film and fabrication technologies, combined with developments in
18 characterization and measurement techniques, have now made it possible to control the migration
19 of individual atoms (ions) in even conventional insulators and semiconductors.^[6] These devices can
20 be programmed at high speed with other operation parameters also compatible with conventional
21 electronic circuits, and thus enable the merge of ionics with electronics. For example, ionic-based

1 resistive switching devices that operate by controlled physical reconfiguration of the switching
2 material, schematically shown in Figure 1, have been extensively studied with excellent performance
3 metrics.^[2,6] Specifically, these devices possess the simplicity of a resistor – with the whole device
4 structure consisting of only two electrodes sandwiching a “switching” medium. Unlike a simple
5 resistor, however, the device offers the capability to change its resistance on demand, and more
6 importantly, the ability to store (memorize) the new resistance state as needed, thus the name
7 memristor (memory + resistor)^[12-14] or RRAM in memory applications. The resistance change and
8 storage (termed resistive switching, RS) effects are achieved by the redistribution of ions (atoms) in
9 the switching layer, which in turn modulate the local resistivity and the overall electrical resistance
10 of the device.^[2,6] In this case, the functionality is realized through ionic processes but the information
11 is sensed and processed through electrical signals, achieving efficient data storage and computing
12 functions in an extremely compact structure.

13 Nanionic devices such as resistive switching devices can be classified in several different ways
14 from materials, switching mechanisms, to switching characteristics and applications.^[1,2,5,6] Based on
15 the switching mechanism, commonly studied resistive switching devices can be generally divided
16 into two groups: devices based on the migration and redistribution of metal cations such as Ag or Cu
17 ions; and devices based on the migration of anions such as oxygen ions, leading to redistribution of
18 oxygen vacancies (V_{Os}). In both types of devices, the RS process is generally triggered by ionic
19 migration leading to the formation of conductive paths for mobile carriers (*e.g.* electrons), while
20 each type of device may exhibit its own distinct characteristics due to the different material
21 properties and internal dynamics involved in the RS process. In this section, characteristics of each
22 device group will be reviewed based on the microscopic RS mechanism.

1

2 **2.1 Electrochemical metallization (ECM) devices**

3 RS devices based on cation migration are usually termed electrochemical metallization (ECM)
4 cells, sometimes also called conductive bridge random access memory (CBRAM) in the
5 literature.^[2,5,15] The basic structure of an ECM cell consists of an electrochemically active electrode
6 such as Ag or Cu, an “insulator” layer serving as the switching layer, and an electrochemically inert
7 counter electrode.^[2,5,6,15,16] The insulator layer in ECM cells plays the role of a solid electrolyte for Ag
8 or Cu cation migration, and can be based on either conventional electrolyte materials such as
9 chalcogenides^[2,17] or even common dielectric materials such as SiO₂^[18], Al₂O₃^[19,20] or amorphous Si
10 (a-Si)^[21,22].

11 During RS, an essentially new material is created in the switching material in an ECM cell, through
12 electrochemical processes and ion migration processes. For example, a Ag filament can be created in
13 an a-Si switching layer under the applied bias, converting the device from the insulating “off” state
14 to the conducting “on” state through this physical reconfiguration process.^[21,22] Three steps are
15 normally involved during this material reconfiguration process: oxidation, migration and reduction of
16 the cations to form the filaments^[2,6,15,17], schematically shown in **Figure 2a**. First, the active anode
17 material (e.g., Ag or Cu) becomes ionized into Ag or Cu cations under an external voltage bias. The
18 ionization rate can be exponentially speeded up under the applied bias, as described by the Butler-
19 Volmer equation.^[23,24] Second, facilitated by the applied electric field, the Ag or Cu cations migrate
20 towards the inert electrode through the solid electrolyte. Again, under high field this ion-transport

1 process can be exponentially speeded up with the barrier lowering effect, allowing fast ion transport
 2 in the switching layer at room temperature, as shown in Figure 2c and Equation (1).^[15,25]

$$3 \quad v = a \cdot f \cdot \exp\left(-\frac{E_a}{k_b T}\right) \cdot \sinh\left(\frac{qaE}{2k_b T}\right) \quad (1)$$

4 Here v is the drift velocity of the cation, a is the distance between two sites that can store the ion
 5 (hopping distance), f is an attempt frequency, E_a is the energy barrier for ion migration, E is the
 6 applied electric field, k_b is Boltzmann's constant, q is the ion's charge and T is absolute temperature.
 7 Note the high field condition can be readily met in nanoscale thin films even at moderate voltages
 8 (e.g. 1V).^[6] The long-distance migration of the ions is a key feature in RRAM and differentiates these
 9 devices from other non-volatile memories that are also based on ion/atom re-arrangements such as
 10 phase-change memory (where local-ordering of the atoms are changed) and ferro-electric memory
 11 (where the ion displacement is within a unit cell).

12 Finally, the displaced Ag or Cu cations become reduced to elemental Ag or Cu atoms after
 13 capturing electrons, leading to the nucleation and growth of metal clusters within the switching
 14 layer through nucleation and deposition processes, eventually resulting in the creation of nanoscale
 15 metal filaments (Figure 2a).^[2]

16 The completion of a metal filament creates an electrically conductive path inside the dielectric
 17 layer that leads to a low-resistance state (LRS), with an abrupt conductance increase in the I - V curve
 18 as shown in Figure 2d.^[2,6,15,17,26] The reverse process, where a segment of the filament is broken by
 19 physically and chemically removing a portion of the filament material, leads to the reset process that
 20 switches the device back to the high-resistance state (HRS).

1 The dynamics of the physical and chemical ionic processes during RS in ECM devices have been
2 studied through cyclic voltammetry (CV) measurements using very slow sweep rates. (Ionic) Current
3 peaks corresponding to characteristic oxidation and reduction processes of Cu ions were clearly
4 observed in a Cu/SiO₂/Pt device^[27,28], as shown in **Figure 3a**. These measurements help identify the
5 partial electrochemical redox reactions of the cations during the RS process. For example, during the
6 subsequent negative sweep after a positive sweep, Cu ions are reduced and thus the concentration
7 of Cu ions near the Cu/SiO₂ interface is significantly decreased. As a result, the current is limited by
8 the diffusion of Cu ions injected into SiO₂ through the first oxidation process, leading to the negative
9 current peaks marked as $J_{p,red}$ in Figure 3a. In addition, the peak current values are proportional to
10 the square root of the voltage sweep rate, as shown in Figure 3b, as expected from the Randles-
11 Sevcik equation for electrochemical processes:

$$12 \quad j_p = 2.99 \times 10^5 \times z^{3/2} \times c_{ox} \times \sqrt{\alpha D_{ox} \nu} \quad (2)$$

13 Here j_p , z , c_{ox} , α , D_{ox} , and ν are the peak current density, the number of electrons transferred
14 during the redox reaction, the concentration of the ions, the charge transfer coefficient, the
15 diffusion coefficient of the ions and the sweep rate, respectively. These quantitative analyses rule
16 out parasitic effects such as the capacitance effect as possible causes of the observed current peaks.
17 Additionally, it has been found that kinetic parameters such as the ion diffusion coefficient and the
18 ion mobility depend on the ion concentration, which is in turn affected by the voltage sweep rate, as
19 shown in Figure 3c and d.^[28,29] Specifically, the ion diffusion coefficient and the ion mobility decrease
20 as ion concentration increases (Figure 3d), indicating that the electrolyte near the reactive electrode
21 behaves as a concentrated solution instead of a diluted solution in CV measurements.^[28,29] In other
22 words, ion-ion interactions are not negligible and thus transport of ions is more retarded under

1 higher ion concentration, leading to smaller diffusion coefficient and ion mobility. In particular, a
2 high ionic diffusion coefficient is expected at high sweep rates, supporting the feasibility of
3 electrochemically assisted ion migration even in conventional dielectric materials and at low
4 temperatures (*i.e.* room temperature).^[28] Clear redox peaks were also observed from CV
5 measurements in other ECM systems, such as Ag/SiO₂/M and Ag or Cu/Ta₂O₅/M (where M = Pt, Ir,
6 Ru, etc.).^[29,30]

7 A key characteristic of the ECM cell is the use of an active electrode which provides the active
8 ionic species (metal cations), while the switching layer acts mainly as a physical storage medium (to
9 first order) and does not participate in the chemical redox processes.^[2,6,15,17,26] For example, it has
10 been found that Ag or Cu does not form chemical compounds with common insulator-based host
11 switching material such as SiO₂, and instead exists as elemental metal inside the dielectric.^[31,32] As a
12 result, dramatic changes to the material properties can be obtained as the new material (filament) is
13 injected and removed from the switching material. Additionally, during reset it is possible to fully
14 remove the filament material to create a gap region that maximally recovers the excellent insulating
15 property of the host material (schematically shown in Figure 2a), thus the reset can lead to a very
16 high off-state resistance in ECM devices and a high on/off resistance ratio, as seen in the *I-V*
17 characteristics during RS of a typical ECM device (Figure 2d).

19 **2.2 Valence change memory (VCM) devices**

20 Physical and chemical ionic processes can also occur in oxide-based solid electrolytes, where the
21 active ionic species are oxygen ions or oxygen vacancies. Resistive switching devices based on such

1 devices are sometimes termed valence change memory (VCM) or simply referred to as oxide RRAM
2 in the literature.^[2] In VCM devices, inert electrodes are used at both the anode and the cathode, and
3 the active species (*e.g.* V_{O_s}) are present internally inside the switching material stack (*e.g.*, in a non-
4 stoichiometric sub-oxide^[33-35] or at an interface formed through an oxygen gettering material such as
5 a Ti electrode layer^[36,37]), as shown in Figure 2b.

6 RS in VCM devices is driven by the internal redistribution of oxygen vacancies between the V_{O-}
7 rich and the V_{O-} -poor layers, again involving physical ion migrations and chemical redox processes.
8 Specifically, under a high applied field, oxygen vacancies migrate from a V_{O-} -rich reservoir layer into
9 the near-stoichiometric switching layer that initially has high resistance. As the V_{O-} concentration in
10 the switching layer increases, percolated paths for electron conduction can form at high V_{O-}
11 concentrations, and switch (set) the device from a high-resistance “off” state to a low-resistance
12 “on” state.^[4,38] The reverse process corresponds to the reset step that switches the device from “on”
13 to “off”

14 One key feature of VCM devices is that the switching layer will also actively participate in the
15 redox chemical processes, *e.g.*, during set that forms a conducting sub-oxide channel (filament) the
16 metal cation in the oxide is reduced to a lower oxidation state^[2,39,40], and recovers back to the near-
17 stoichiometric oxide composition during reset, schematically shown in Figure 2b. Because the active
18 species, V_{O_s} , are native defects in these oxides that also chemically react with the switching material,
19 it is difficult to completely remove all V_{O_s} from the switching layer and fully restore the near-
20 stoichiometric composition during reset. As a result, the off-state resistance of VCM devices is
21 generally leakier than that of ECM devices, leading to smaller on/off ratios^[4,41], as shown in Figure
22 2e. On the other hand, VCM devices generally exhibit more reliable switching characteristics with

1 better write/erase endurance due to the use of native ionic species instead of foreign ionic species
2 as in the case of ECM devices, as the repeated insertion and removal of foreign cations in a solid
3 dielectric film may lead to plastic deformation of the film and irrecoverable damage if programming
4 conditions are not optimally set.^[34,42]

6 3. Fundamental experimental studies and modeling

7 3.1 *In-situ* TEM observations

8 To experimentally reveal the microscopic origin of resistive switching processes, material
9 characterization techniques such as transmission electron microscopy (TEM)^[31,32,40,43-47] energy
10 dispersive X-ray spectroscopy (EDS)^[31,48,49] X-ray absorption spectroscopy (XAS)^[50-52], electron energy
11 loss spectroscopy (EELS)^[44,45,47], photoemission electron microscopy (PEEM)^[53-55], and conductive
12 atomic force microscopy (C-AFM)^[26,56] have been employed. As the most direct method, physical
13 reconfigurations of the switching layer have been analyzed *via* microscopic imaging techniques.
14 Specifically, several TEM-based studies using high-resolution imaging and spectroscopic analysis
15 tools combined with electrical measurements have provided direct evidence of (long distance) ion
16 migration during RS processes for both ECM and VCM types of devices.^[31,32,40,43-45,47,50,57-59]

18 3.1.1 ECM devices

19 **Figure 4a and b** highlight the formation/rupture of nanoscale Ag filaments after set/reset in an
20 Ag/SiO₂/Pt-based ECM device (fabricated in a planar structure on a SiN_x membrane, inset).^[31] The as-

1 fabricated device was turned into the low-resistance state with an abrupt increase in current after
2 application of a constant positive voltage bias, as shown in Figure 4c. The electrical resistance
3 change was found to accompany changes in the electrode and the switching material (Figure 4a and
4 c), where the active electrode material, *i.e.*, Ag, was found to be injected into the insulating SiO₂
5 layer (acting as the switching layer) under the applied electric field. The device is switched “on”
6 when the two electrodes are connected by the Ag filament, abruptly changing the device’s electrical
7 resistance. Upon voltage bias of opposite polarity, the device switched back to the high-resistance
8 state (Figure 4d), corresponding to the rupture of the filament (Figure 4b) due to similar redox
9 reactions and ionic migration in the reverse direction. These observations unambiguously reveal the
10 physical nature of the RS process, where the physically displaced Ag atoms lead to dramatic changes
11 to the device’s electrical properties.

12 Besides allowing direct observation of the metal filaments during RS, *in-situ* TEM analyses provide
13 important information concerning the microscopic dynamic ionic processes during filament growth.
14 Figure 4e shows real-time structural evolutions of an Ag/a-Si/W device under a constant voltage bias
15 from an *in-situ* TEM study.^[31] The filament growth process can be directly correlated with the
16 electrical measurements, allowing one to identify the different steps during filament growth in real-
17 time. Subsequent systematic *in-situ* TEM analyses provided a complete picture of the different
18 dynamic processes that can occur in an ECM device, *e.g.*, filament growth direction, geometry, and
19 morphology.^[32] Based on these experimental studies, a model that incorporates key kinetics factors
20 such as the ion mobility and the oxidation/reduction rate has been developed that can reconcile
21 previously reported, seemingly contradicting experimental findings, depending on which ionic
22 process becomes rate-limiting.^[6,32,60]

1 Specifically, the filament growth/dissolution dynamics in ECM devices can be classified into four
2 categories based on the ion mobility and the redox reaction rate (Figure 4f). Note both parameters
3 depend on the material properties but can also be affected by programming conditions such as the
4 applied electric field and temperature:^[32] (1) In cases with high ion mobility and high redox rates, the
5 active cations injected into the electrolyte during set will be reduced at the inert (cathodic) electrode
6 due to the high ion mobility, leading to filament growth initiated from the inert electrode. A large
7 amount of ion supply due to the high oxidation rate results in the formation of cone-shaped
8 filaments with a large base at the inert electrode interface.^[60] Such inverted cone-shaped filaments
9 have been observed in ECM devices using classical solid electrolytes such as GeTe^[43] that are known
10 to be good ionic conductors. (2) In cases with high ion mobility and low redox rate, filaments will
11 grow from the cathode because of the high ion mobility, similar to the previous case. The filament
12 shape and filament dissolution behaviors, however, will be different because the oxidation process is
13 slow and rate-limiting, so that the supply of cations is limited. The limited cations will be more likely
14 to get reduced at the tip of existing filaments where the field is enhanced, leading to the formation
15 of branched filaments such as the case observed in porous sputtered-SiO₂ (Figure 4a and b).^[31,57,60]
16 (3) and (4) correspond to systems with low cation mobility, which will exhibit a different filament
17 growth direction compared to the high cation mobility cases because ionic transport is now the rate-
18 limiting process. In particular, the nucleation of metal clusters could occur inside the electrolyte
19 before the cations reach the inert electrode interface due to the slow movement of ions.^[31,60] The
20 precipitated metal clusters inside the electrolyte will extend the filament growth from the active
21 electrode toward the inert electrode during set.^[6,32] Indeed, forward growth of filaments was
22 experimentally observed in ECM devices using non-conventional electrolyte such as amorphous Si^[31]

1 and oxides such as ZrO_2 ^[58] and PECVD-grown SiO_2 ^[59] that are known to have lower ion mobility
2 compared to conventional electrolytes. The differences in redox rates can further lead to different
3 filament shapes, with continuous extension of the anode towards the cathode when the redox rate
4 is high, case (3),^[32] and discrete nanoclusters and overall forward cone-shaped filaments when the
5 redox rate is low, case (4), as shown in Figure 4e^[31].

6 The exponential field-acceleration, a key feature of the ionic-based RS devices, has also been
7 verified by systematically examining the wait time for the abrupt current change as a function of the
8 applied field. Several studies have demonstrated that the wait time depends exponentially on the
9 electric field, and shows an Arrhenius relationship with the temperature,^[31] as shown in Figure 4g
10 and h respectively, consistent with predictions from Equation (1) in cases where ion migration is the
11 rate-limiting process.

13 3.1.2 VCM devices

14 Similar to ECM devices, there have been extensive studies aimed to observe conductive filaments
15 and to identify the origin of RS behaviors in oxide-based devices such as TiO_2 ^[40,50], Ta_2O_5 ^[45,47], and
16 ZnO ^[44] systems, using TEM techniques in conjunction with spectroscopic analysis,.

17 **Figure 5** shows high-angle annular dark-field (HAADF) scanning TEM (STEM) images of the same
18 location in a $\text{Pt}/\text{SiO}_2/\text{Ta}_2\text{O}_{5-x}/\text{TaO}_y/\text{Pt}$ device, after applying negative (set) and positive (reset)
19 voltages, respectively, using an *in-situ* experimental setup. Clear contrast differences in the HAADF-
20 STEM images can be observed, indicating changes in the composition of the film at different
21 resistance states. Regions corresponding to localized, nanoscale conductive channels with a different

1 composition compared to the host material are indicated by the yellow arrows in Figure 5a. EELS
2 line-scan analysis confirmed that the contrast changes in the STEM images are due to differences in
3 oxygen concentration in the switching regions. Specifically, the local oxygen concentration after the
4 set operation is reduced compared to the off-state, shown in Figure 5d, verifying the role of V_O
5 migration during the operation of VCM devices.

6 Like metal ions in ECM devices, the migration of V_O s and the associated redox processes cause
7 direct changes to the chemical composition and physical properties of the switching layer. In this
8 case, the redistribution of oxygen vacancies modulates the stoichiometry of the oxide-based
9 switching layer and in turn the film's electrical conductivity, although the changes in material
10 properties may be less dramatic compared with ECM devices where complete new materials are
11 created in the host. Similar to ECM devices, the modulation of the chemical composition and the
12 electrical properties of the switching material in VCM devices can be reversed by applying a voltage
13 bias in the opposite polarity. In this case, typically longer write/erase endurance cycles can be
14 obtained in VCM devices due to the involvement of only native ionic species during the chemical and
15 physical processes.^[33,34]

16 17 **3.2 Other experimental techniques**

18 It is generally more challenging to directly observe filament evolution in VCM devices^[61] since the
19 filaments consist of native defects, *i.e.* oxygen vacancies rather than foreign metallic species. Careful
20 spectroscopic analysis is typically required to confirm the composition changes. A host of state-of-
21 the-art techniques such as X-ray absorption spectromicroscopy^[51,52], photoemission electron

1 microscopy (PEEM)^[53-55], and electron holography^[62] have recently been employed to
2 comprehensively characterize the changes in film microstructure, composition and chemical states
3 accompanying the resistive switching effects in VCM devices. However, continued improvements in
4 spatial and temporal resolution are still needed to clearly reveal in real-time the nanoscale filament
5 growth dynamics.

7 3.2.1 XAS analysis

8 Spatially-resolved XAS using scanning transmission X-ray microscope (STXM) was performed to
9 non-destructively investigate the chemical and structural changes during RS.^[50] After electroforming
10 and set/reset cycling, a localized region showing a different absorption spectrum from that of the
11 surrounding matrix was clearly observed in STXM image of a TiO₂-based device, as shown in **Figure**
12 **6a.**^[50] The altered absorption spectrum after the electrical cycling matches that of reduced titanium
13 oxide, in which the valence state of the Ti ions is reduced from +4 to +3, while the absorption
14 spectrum of the titanium oxide layer in virgin-state with no prior electrical treatment is consistent
15 with stoichiometric titanium dioxide (Figure 6b). Electron diffraction measurement revealed that the
16 localized region consists of the Ti₄O₇ *Magneli* phase with metallic conductivity, supporting the
17 formation of localized conducting channels due to oxygen vacancy migration during resistive
18 switching. Localized conducting channels with reduced oxygen concentration after electrical
19 forming/set were also found in a Ta₂O₅-based device from X-ray measurements, consistent with
20 results obtained from cross-sectional TEM and EELS studies.^[63,64]

1 Thermally-driven radial migration of oxygen ions/vacancies and the associated local composition
2 changes were also verified by X-ray photoemission electron spectromicroscopy upon application of
3 multiple high-voltage pulses.^[52] Figure 6c shows the O K-edge X-ray transmission intensity map of an
4 electrically-cycled Ta₂O₅-based device, where a ring-like feature with a bright inner core and a dark
5 perimeter is observed. Besides field-driven drift of oxygen ions/vacancies in the vertical direction
6 (parallel to the field), temperature gradients due to Joule heating near the localized conduction
7 paths produce thermophoretic forces that cause lateral, inward-migration of V_Os (*i.e.* Soret
8 Diffusion)^[65], which is in turn balanced by Fick diffusion driven by the V_O concentration gradient. As a
9 result, lateral segregation of oxygen-deficient and oxygen-rich regions occurs, corresponding to the
10 bright and dark areas in Figure 6a, respectively. It was found that electrical cycling with high power
11 amplified the lateral V_O migration, eventually inducing irreversible changes to the material, including
12 strong clustering of defects and accelerated device failure. The change of the Ta oxidation state (due
13 to V_O migration) rather than the Ta concentration (due to Ta migration) in the oxygen-deficient
14 region was also verified by Ta L-edge X-ray transmission analysis, consistent with the now-generally
15 accepted theory of VCM devices. On the other hand, it has been recently observed that cation
16 migration is also possible in VCM devices under high bias,^[66] which should be generally avoided and
17 may affect the reliability of the device operations.

18 In addition, XAS has been employed to investigate the detailed information of local structural
19 changes and chemical reactions during the operation of RRAM devices. Specifically, the X-ray fine
20 structure (XAFS) technique can provide short-range structure information (*e.g.*, local atomic
21 arrangement, bond distances, and coordination number) from extended X-ray absorption fine
22 structure (EXAFS) analysis, as well as local chemical environment (*e.g.*, oxidation state and local

1 coordination) through X-ray absorption near-edge structure (XANES) analysis.^[67] For example, the
2 short-range structural changes in the TiO_{2-x} switching layer^[68] and the intermediate nanoscale oxide
3 layers that spontaneously form at the Me/MeO_x interface (where Me = Ti, Hf and Ta, and MO_x =
4 Ta₂O₅, HfO₂, and SiO₂)^[69] were observed by comparing electron momentum-weighted oscillation
5 spectra, $k^2\chi(k)$ in EXAFS, as shown in Figure 6e. The observed dependence of the electrode oxidation
6 rate on the oxide's thermodynamic stability and the metal's oxygen affinity in Ref. 69 was found to be
7 consistent with the RS behaviors of Ta₂O₅-based devices with different electrodes, as well as
8 predictions from first-principle calculations.^[70] More importantly, this sensitive technique helps point
9 out the role of the interfacial oxide layer as an oxygen reservoir in mediating reliable ionic migration,
10 and the mechanism of device failure due to the transition of short-range ordered interfacial oxide
11 toward disordered main oxide matrix.^[69] In another study, chemical dissolution of Ag into a GeS_x
12 electrolyte accompanied by Ag ionization was revealed through XANES analysis^[71] that is known to
13 be sensitive to local chemical bonding and thus the metal's oxidation state^[67]. Specifically, higher-
14 energy shift was observed in the XANES derivative spectra peaks from a Ag/GeS_x film compared to a
15 Ag foil, whereas data from a Ag/SiO₂ film showed no change compared with the Ag foil.^[71] This
16 observation indicates that the valence number of Ag in SiO₂ is null (*e.g.*, Ag exists as elemental Ag⁰)
17 while Ag in GeS is positively charged, implying that not only metallic Ag atoms but also the formation
18 of the conductive Ag-Ge-S ternary compound may account for the RS behavior in chalcogenide-based
19 ECM devices.

21 3.2.2 XPEEM analysis

1 Besides binary oxides, complex perovskite oxides such as SrTiO₃ are also known to exhibit
2 resistive switching behaviors based on the redistribution of oxygen vacancies.^[72,73] Nanoscale redox
3 reactions in SrTiO₃ were verified by X-ray PEEM analysis in a device with a graphene top electrode
4 that provides both electrical conductivity and photoelectron transparency to allow electrical
5 programming and spectromicroscopy measurements simultaneously *in situ*, in a non-destructive
6 fashion.^[54] When the device was switched between the LRS and the HRS, noticeable differences
7 were observed in PEEM images in localized areas (**Figure 7a and c**), which can be attributed to the
8 formation/dissolution of a conductive filament. The decrease of the O concentration in the LRS was
9 observed from the O K-edge spectra (Figure 7b and d), revealing the correlation between the
10 changes in electric properties and the chemical composition modulation, controlled by *in-situ*
11 electric stimuli. In addition to O migration, redox reactions involving changes of the cation valence
12 state (*e.g.*, Ti⁴⁺ ↔ Ti³⁺) were verified, by comparing the relative spectra intensity at different photon
13 energies that are sensitive to specific Ti oxidation states. Specifically, the Ti³⁺ concentration
14 decreases after the reset process and increases again after the set process, as shown in Figure 7e,
15 indicating the valence state change of the Ti ions induced by the O ion migration during resistive
16 switching. Similar results are also consistently observed in other studies.^[50,53,55] On the other hand,
17 even the HRS does not show the same spectra as in the reference area, as shown in Figure 7d,
18 indicating the incomplete restoration during the reset process in oxide-based devices. Due to the
19 considerable amount of oxygen vacancies left in the switching layer that cannot be completely
20 removed, the resistance of the off-state in oxide-based devices is much lower than in the prior-
21 forming, virgin state, leading to a small memory window as discussed earlier.

22

1 3.2.3 C-AFM analysis

2 Besides direct observations of changes in the switching material's composition (*e.g.*, Ag/Cu
3 injection in electrolytes or cation/oxygen ratio in oxides) and chemical states (*e.g.*, oxidation state of
4 cations in oxides), changes in material properties such as the local electrical conductivity can be
5 measured by the 3D conductive atomic force microscope (C-AFM) tomography technique. 3D
6 conductivity tomography is based on scraping the material of interest layer-by-layer through a
7 conductive diamond tip in a C-AFM setup, allowing the local conductivity of the film to be revealed
8 even for regions normally buried under an electrode, as schematically shown in **Figure 8a**. Figure 8b
9 and c show 3D conductivity tomography results obtained in a Cu/Al₂O₃ device in the LRS and the
10 HRS, respectively.^[26] Due to the high spatial resolution of the scanning probe microscopy and the
11 subnanometer precision in the material removal process, the shape and size of the conductive
12 filament can be accurately determined with nanometric resolutions in all three dimensions.
13 Specifically, conductive filaments observed in the Cu/Al₂O₃ device exhibit a conical shape with the
14 narrow constriction formed close to the inert-electrode (Figure 8d), consistent with observations
15 through earlier TEM studies and predictions from the electrochemical theoretical framework
16 discussed in Section 3.1.1.^[31,32] The ruptured region of the filament in the HRS (Figure 8c) was found
17 to be near the inter-electrode interface, as expected from the geometry the filament since rupture is
18 expected to occur at the narrowest region of the filament. These results identify the actual physical
19 location of the switching events, and is critical for continued device optimizations. The scraping C-
20 AFM technique was also successfully employed for oxide-based VCM devices (Figure 8e).^[56]
21 Specifically, direct measurements from the 3D conductivity tomogram verified that the conductive
22 filament size can be controlled to be less than $7 \times 7 \text{ nm}^2$ in a Hf/HfO₂ device, by controlling

1 programming conditions such as the current compliance as shown in Figure 8e, thus supporting the
2 excellent scalability offered by RRAM devices below the 10 nm regime as demonstrated in previous
3 studies.^[74,75]

4 5 **3.3 Theoretical simulations and modeling**

6 Accompanying the experimental findings, modeling efforts with different levels of physical details
7 have been performed to examine the underlying physical processes in these coupled ionic/electronic
8 systems. These models in turn help predict the device switching characteristics and provide the
9 needed insights towards continued device and material design and optimization.

10 Recent atomistic simulations^[76,77] based on molecular dynamics provide a feasible physical
11 picture for the ECM devices during RS processes. The simulation implemented reactive interatomic
12 potentials, including the effects of electrostatic potential on partial atomic charges, to describe the
13 complex chemical reactions involved in the RS process beyond conventional molecular dynamics
14 approaches.^[78,79] Upon the application of a set/reset voltage, the formation/rupture of a conductive
15 bridge connecting the two electrodes were successfully observed^[76,77], as shown in **Figure 9a and b**.
16 More importantly, the growth of the filament by the progressive reduction of dissolved Cu ions was
17 observed during the forming process (Figure 9c-f), which is consistent with experimental
18 observations in systems with fast ionic transport.^[31] During reset the filament became ruptured in
19 the proximity of the inactive electrode as a result of the dissolution of Cu atoms connected to the
20 positively-biased electrode (Figure 9g-i), again supporting earlier experimental results on the
21 location of the filament rupture.^[31] The concentration of Cu ions in the SiO₂ electrolyte increases in

1 the gap region after the initial forming and reset cycle, which facilitates the re-construction of a new
2 filament at a lower voltage during subsequent set processes, again consistent with widely reported
3 experimental findings.^[2,26,31,60] Additionally, simulations on the atomic trajectories during the atomic
4 reconfiguration provide insight into critical kinetic parameters that allow stable switching at
5 nanoscale. For example, the aggregation of Cu atoms with more than three metal-metal bonds lead
6 to metallic clusters and subsequently stable filaments by the percolation of the clusters, while single-
7 atom-chain bridges are found to be not thermodynamically stable in this SiO₂/Cu system, probably
8 due to the low activation energy for Cu migration.

9 Atomistic simulations based on molecular dynamics and first-principles calculations using density
10 functional theory (DFT) have also been performed to investigate the physical processes during VCM
11 device switching. The changes in electronic structure induced by the modulation of oxygen vacancy
12 concentration in Ta₂O₅ were confirmed by density-of-state calculations using amorphous cells having
13 different stoichiometry.^[80,81] Additionally, a V_O-rich region providing continuous conduction path for
14 electrons was observed in a cell with low O concentration.^[81] Attractive cohesive interaction
15 between oxygen vacancies was found in TiO₂, supporting experimental observations of a conductive
16 V_O-rich filament with *Magneli* phase and ordered vacancies.^[40,82] In addition, first-principles
17 calculations have been performed to calculate important thermodynamic and kinetic factors such as
18 the V_O formation energy and the migration barrier more accurately, without using empirical
19 parameters as *a priori*. The calculated formation energy for oxygen vacancies and the migration
20 barrier in various oxides show good agreement with experiments^[83-86], providing not only clear
21 theoretical basis for the roles played by oxygen vacancies but also guidelines for device design and
22 optimization. For example, the predicted effects of the filament growth conditions (*e.g.*, the oxygen

1 chemical potential and thus the vacancy formation energy) from theoretical calculations were
2 successfully implemented to experimentally engineer devices with improved I - V characteristics by
3 selecting proper materials as the reactive electrode.^[70,87]

4 Although the atomistic simulations are based on accurate underlying physical processes, the
5 spatial and temporal scales analyzed in the simulations are normally experimentally unattainable.
6 For example, the dynamical evolution of the atomic filament configuration was simulated within a
7 very short time, \sim ns, in Ref. 76. Additionally, comprehensive, atomic-level descriptions of VCM
8 devices that can capture all relevant dynamic processes and can be closely linked to experimental
9 observations still remain to be developed. On the other hand, macroscopic-level, phenomenological
10 models that can quantitatively describe resistive switching behaviors in oxide-based devices have
11 been extensively studied. By solving a set of partial differential equations (*i.e.*, ionic drift/diffusion
12 and continuity equation for ion transport, electronic continuity equation for current conduction, and
13 Fourier equation for Joule heating), the dynamic switching processes including the evolution of the
14 atomic configurations and the associated conductance changes can be reliably modeled.^[88,89] The
15 obtained I - V characteristics from a simulated cell that mimics the actual device structure (**Figure**
16 **10a**) exhibit remarkable agreement with the experimentally measured data, as shown in Figure 10b.
17 Dynamic switching processes including the depletion of oxygen vacancies close to the top electrode
18 during the reset process, and the refilling of the gap during the set process, were accurately
19 captured as shown in the V_{O} concentration profile (Figure 10c and d). In addition to binary resistive
20 switching, rich phenomena based on controlled oxygen vacancy redistribution, such as multiple
21 conductance states for multilevel storage, analog conductance modulation that is useful for

1 neuromorphic computing, and complementary resistive switching (CRS)^[90,91] effects, can be
2 accurately captured in the model and used to support experimental findings.^[92]

3 The rich and controlled dynamic processes observed in coupled ionic/electronic systems make
4 these devices promising candidates as a reconfigurable platform for multifunctional devices and
5 circuits. A broad range of material and device parameters, *e.g.*, the ion migration activation energy,
6 ion hopping distance, electrode material, and oxide stoichiometric/thickness, that can affect the RS
7 characteristics have been examined in simulations and experiments, through techniques such as
8 controlled doping^[93], interface/stack engineering^[94], and control of the electrical forming process^[95].
9 In addition, the role of Joule heating was extensively analyzed during the operation of VCM devices.
10 The heat generation inside the filament due to Joule heating, as shown in Figure 10d, can greatly
11 facilitate thermally-activated processes including ionic diffusion and drift, and can thus be further
12 utilized to improve the RS characteristics such as the dynamic range of the device^[96]. Additionally,
13 utilizing the internal ionic dynamics may allow realistic emulation of synaptic behaviors in these
14 solid-state devices, including activity-dependent synaptic plasticity where the local temperature acts
15 as a short-term state variable (discussed in more detail in Section 4.2.4).^[97]

16 17 **4. Applications**

18 **4.1. Resistive switching devices – memory**

19 The most obvious application of the conductance modulation effect is in the form of RRAM for
20 non-volatile, high density and fast data storage. RRAM devices an attractive candidate in
21 applications such as storage-class memory (SCM), by potentially combining the desirable

1 characteristics of DRAM and Flash memory and bridging the performance gap within the memory
2 hierarchy.^[1,3,4,98] Specifically, in RRAM the resistance change is associated with material-
3 reconfiguration driven by redox reactions and ion transport that can be exponentially accelerated by
4 electric field and/or temperature during programming.^[88,89] This extremely non-linear response to
5 stimulation ensures that the new atomic-configuration state will remain unchanged when the
6 voltage bias is removed or under low-bias read conditions, thus enabling the devices' ability of both
7 fast programming (within ~ns) and long data retention (~years). Indeed, sub-ns switching speed^[99,100]
8 and decades-long retention time^[34,101] along with other excellent performance metrics, *e.g.*, large
9 on/off ratio^[25], low energy consumption^[18,19] and high endurance^[33,34] have been demonstrated.
10 Extensive research efforts have been made to further improve the performance of RRAM devices for
11 very large scale, practical applications in the last decade, and continued device optimizations are still
12 currently underway.

14 4.1.1 Device variability

15 One of the most challenging issues in device-level research is to improve the device uniformity,
16 for both ECM and VCM devices. Specifically, RRAM devices are inherently stochastic during the
17 resistive switching process, originating from the stochastic, thermally activated ion migration and
18 redox processes. These stochastic effects lead to temporal variations^[102], in which the switching time
19 or switching voltage is broadly distributed from cycle to cycle even for the same device.^[103,104] The
20 stochastic growth of the filament, combined with the disordered nature of the amorphous switching
21 material, also results in large device to device variations. To minimize bit errors during programming,

1 feedback schemes that verify the device state after each programming operation, or excessive
2 programming voltage/pulse width conditions can be used at the expense of system speed or device
3 reliability, respectively.^[105] In addition, device variations can arise from non-ideal fabrication
4 processes such as inhomogeneous film thickness, film stoichiometry, and line edge roughness, which
5 can be mitigated through tighter process control.^[104]

6 Engineering of the switching materials and the device structure has been extensively carried out
7 to improve RS uniformity. Doping of the oxide layer in VCM devices, such as Al^[106], Gd^[107], and Ge^[108]
8 doping in HfO₂, has been reported to result in narrower distributions of the HRS/LRS resistance and
9 the forming, set, and reset voltages. Dopants can lower the oxygen vacancy formation energy and
10 thus oxygen vacancies are likely to be formed near the dopants instead of randomly, thus help
11 confine the V_o filament formation to the same location among cycles and improve cycle-to-cycle
12 uniformity. Improved uniformity in the operating voltages and resistance states was also achieved by
13 introducing nanodots^[109-111] and nanocone-shaped electrodes^[112]. The concentrated electric-field
14 resulting from these engineered nanostructures can help guide ion transport and redox reactions,
15 thus suppressing random and uncontrolled filament growth.^[109] The positive effect of filament
16 confinement on device uniformity was also verified in a graphene-inserted RRAM device, where the
17 ion transport is allowed only through an engineered nanopore in the graphene film working as an
18 ion-blocking layer, leading to a confined/controlled filament with improved device variations.^[49,113]

19

20 *4.1.2 Memory window and set/reset current in VCM devices*

1 Other crucial parameters that need to be further optimized for RRAMs, particularly VCM devices,
2 include the set/reset current and the memory window.^[114] Relatively high operation current (\sim mA)
3 and small on/off ratio ($< \sim 10^2$) are typically observed in VCM devices. A lower operation current is
4 desirable not only to achieve low power consumption but also to mitigate the series resistance issue
5 from the interconnect wires in high-density arrays, and to relax the current driving capability
6 requirement of the programming circuitry and the selector element,^[19] while a larger memory
7 window with high on/off ratio relaxes device variation requirements and allows larger arrays to be
8 build.

9 Multi-layered device structure may be used to reduce the operating current, in which additional
10 semiconductor or dielectric layers are inserted in the device stack to control the supply of oxygen
11 vacancies or to act as a barrier layer. For example, an asymmetric and gradual profile of V_O chemical
12 potential was obtained by employing multiple oxide layers with different O-scavenging effects,
13 leading to the confined filament formation with lower current levels.^[115,116] A reduced operation
14 current was also observed when the supply of oxygen vacancies is limited only through an
15 engineered nanopore in a graphene layer, in which the operation current was found to be directly
16 proportional to the diameter of the nanopore.^[49] Simply scaling the device size may also lead to
17 lower operation current of VCM devices. Along with increased HRS resistance due to suppressed
18 leakage current through the bulk film, higher LRS resistance values ($> 0.1\text{M}\Omega$) were obtained in
19 aggressively scaled devices^[74,75,117], especially when the device is scaled to the sub-10nm dimension,
20 *i.e.* comparable to the nanoscale filament size of $\sim 10\text{nm}$ ^[56]. Following a similar concept, low-current
21 operation of a HfO_2 -based VCM device was demonstrated in a C-AFM based structure, where the C-
22 AFM tip serves as the top electrode having a sub-10nm diameter.^[118] It is important to note that in

1 these cases the on/off ratio is maintained with significantly increased HRS resistance, contrary to
2 cases when the low programming current is obtained in large size devices through current
3 compliance which typically leads to the degradation of the on/off ratio. Scaling the filament size by
4 limiting the oxygen vacancy supply may further increase the HRS resistance since only a small
5 number of V_{O} s need to be removed to completely restore the excellent insulating properties of the
6 host material during reset (*i.e.* achieving a deep off-state)^[118] thus resulting in a high on/off ratio that
7 has been highly desirable but difficult to achieve in VCM devices.

8

9 4.1.3 Switching reliability in ECM devices

10 Even though the operations of both VCM and ECM devices originate from similar physical
11 mechanisms, distinct RS characteristics are typically observed in these devices. For example, large
12 memory windows and low operation current levels can be readily obtained in ECM devices, while
13 such properties remain elusive for VCM devices, as discussed in the last section. Even with an
14 ultralow programming current ($\sim 1\text{ nA}$), a $\text{Cu}/\text{Al}_2\text{O}_3$ -based ECM device still exhibits excellent on/off
15 ratio with good retention characteristics.^[19,119] The feasibility of very low-current operation in ECM
16 devices can be attributed to the microscopic origin of the RS process, that is, the injected Cu and Ag
17 cations do not chemically react with the host material in certain insulator-based switching layers and
18 can thus be completely removed from the switching layer, leading to efficient recovery of the
19 excellent insulator property at HRS. Indeed, although compounds may be formed in some solid
20 electrolytes such as chalcogenides,^[17,71] in insulator-based switching layers such as SiO_2 it has been
21 found that the injected Cu and Ag cations exist as pure metal element instead of chemical

1 compound. For example, selective area electron diffraction (SAED) patterns of filaments in a Ag/a-Si
2 device, and fast Fourier transformation (FFT) patterns of a HRTEM image on nanoclusters in a
3 Ag/SiO₂ device are indexed as elemental phase of Ag.^[31,32] As a result, the deep off-state provides
4 sufficient on/off margin even in cases with a very low on-state current, in devices based on a
5 barrier/tunneling layer or a partially formed filament using an external current compliance^[19,120].

6 On the other hand, the repeated insertion of foreign materials into the switching layer in ECM
7 devices can potentially cause inevitable mechanical stress and result in deteriorated device
8 reliability.^[42] Indeed, the reported write/erase endurance of ECM devices is typically limited to the
9 order of 10⁶, comparably lower than that of VCM devices (~ 10¹²).^[33,34] Improved endurance has
10 been reported by introducing an additional layer such as TiW^[121], Al₂O₃^[120], and graphene^[122] in the
11 device stack that prevents excessive cation injection into the switching layer. However, resolving the
12 endurance and reliability issue still remains a critical challenge in ECM device development,
13 particularly for logic-type applications where long endurance is needed.

14 15 4.1.4 Crossbar arrays

16 A main feature of RRAM is the high storage-density the device offers. Specifically, RRAM's simple,
17 two-terminal structure enables device integration in a crossbar form with high-density and high-
18 connectivity, having one device formed at each crosspoint in the crossbar, as shown in **Figure 11a**. In
19 the crossbar structure, each device can be randomly accessed and the cell size is the minimum $4F^2$ (F
20 representing the smallest feature size), allowing highest possible 2D density. Additionally, RRAM
21 fabrication typically involves only a few additional steps, using low-temperature processes and

1 materials that are mostly compatible with complementary-metal-oxide-semiconductor (CMOS)
2 processes. This makes it possible to directly integrate RRAM arrays on top of CMOS circuitry in the
3 same chip, as well as 3D stacked memory structures with even higher density. Combined with the
4 sub-10nm scalability discussed earlier, RRAM is generally viewed as a promising candidate for a
5 broad range of applications from high-density data storage, hybrid CMOS logic circuits to
6 portable/wearable electronics.^[22,123]

7 Figure 11b shows an RRAM crossbar array directly integrated on top of a CMOS chip. Each Ag/a-Si
8 switching element in the 40×40 array can be successfully programmed into two distinct resistance
9 states with a read margin of at least $20\times$ in the worst cases. Different 1600-pixel bitmap images were
10 programmed in the 40×40 array through CMOS decoders underneath, and successfully retrieved, as
11 shown in Figure 11c.

12 The successful operation of large-scale crossbar arrays requires the development of non-linear
13 “selector” devices to suppress leakage currents *via* sneak paths in the passive crossbar, along with
14 optimized programming and read bias schemes.^[124,125] Various selector devices based on
15 mechanisms such as Schottky barriers,^[126,127] crested tunneling barriers,^[128,129] Mixed-Ionic-
16 Electronic-Conduction (MIEC),^[130] threshold switching^[131] and metal-insulator-transition^[132-134], have
17 been proposed. These devices either show a rectifying behavior that suppresses the reverse current
18 conduction or a very non-linear I - V characteristic that suppress leakage current in the low-bias
19 regime. It should be noted that even though rectifying selectors are helpful during the read
20 operation, they do not reduce power dissipation through the half-selected devices during the
21 programming stage.^[125] As a result, selectors with high I - V nonlinearity is preferred. Recently, large
22 nonlinearity (10^{10}) has been reported in a Ag/HfO₂-based threshold switching device, with the

1 formation of the Ag filament at high bias and spontaneous rupture at low bias providing the desired
2 nonlinearity.^[135,136] In addition, RRAM devices with intrinsic nonlinear characteristics, *e.g.* by adding
3 a barrier layer in the device stack,^[119,137,138] can simplify the fabrication process by not having to build
4 a middle electrode^[127,135].

6 4.1.5 Three-dimensional (3D) integration

7 To further improve the storage density, three-dimensional integration of memory structures is
8 necessary. 3D RRAM architectures can be achieved since only low-temperature processes are
9 needed and the device does not rely on a single crystalline substrate. There are generally two types
10 of 3D architectures.^[139] The first involves stacking multiple 2D crossbar arrays on top of each other,
11 as shown in Figure 11d and e, in which the storage density increases as the number of stacked layers
12 increases. The fabrication and functionality of 3D stacked structures have been demonstrated using
13 Ag-based ECM devices^[140], polymer(PI:PCBM)-based organic resistive switching devices^[141], and
14 TaON-based VCM devices^[142]. To resolve the sneak current issue, a current rectifier, *e.g.*, a TaO_x
15 diode shown in Figure 11f, was inserted in series with each RRAM cell in the cross-point 3D structure
16 in Ref. 142. 3D stacking offers ultra-high scalability with minimal changes to the fabrication
17 processes. However, the number of lithographic steps and masks used in the process increases
18 linearly with the number of stacks.^[139,143]

19 As a cost-effective 3D architecture, the vertical RRAM structure was proposed.^[144] In the vertical
20 RRAM structure the MIM cells are formed in parallel at the sidewall of a vertical electrode (Figure
21 11g-i). Several studies have demonstrated the feasibility of the proposed vertical RRAM structure,

1 mostly based on transition metal oxide-based systems *e.g.*, TaO_x^[145], HfO_x^[146], TaO_x/TiO₂^[147], WO_x^[148],
2 and AlO₃/Ta₂O_{5-x}/TaO_y^[149]. The integrated RRAM devices in the vertical structure exhibit typical
3 resistive switching behaviors, *e.g.*, >1000× on/off ratio, >10¹⁰ endurance, 10⁴ s retention at 125°C.^[149]
4 More importantly, good uniformity between devices at different layers was observed, supporting
5 the feasibility of vertical 3D integration.^[146,149] Since the electrode size in the vertical RRAM structure
6 is determined by the width of the vertical electrode and the thickness of the horizontal electrode,
7 potentially better scaling can be achieved since film thickness control during deposition is generally
8 easier than pattern size control during lithography. For example, a graphene layer with atomic
9 thickness has been employed as a horizontal electrode in a vertical RRAM structure, with improved
10 power consumption due to the reduced device size.^[150] Approaches to overcome the sneak current
11 problem in vertical RRAM structures have also been carried out, using devices having intrinsic
12 nonlinearity or self-rectifying behaviors^[147], or through the implementation of a vertical transistor^[146]
13 or a metal-insulator-transition (MIT)-based selector^[134] device.

14

15 **4.2. Resistive switching devices – synaptic devices**

16 *4.2.1 Neuromorphic computing approaches*

17 Beyond memory applications, an emerging, attractive field for resistance switching devices is bio-
18 inspired neuromorphic computing.^[151-153] Non-conventional computing architectures are now being
19 seriously considered since the current digital computing architecture, in which the central processing
20 unit (CPU) is physically separated from the memory, increasingly suffers from the von Neumann
21 bottleneck problem resulting in the loss of energy efficiency and throughput due to intensive data

1 movements. Furthermore, rapid advances in social networks, mobile devices and sensors require
2 efficient, real-time storage and analysis of large amounts of data, while current hardware systems
3 are not optimized for these data-intensive tasks. Neuromorphic computing systems, taking
4 inspiration from the human brain, can potentially offer enormous computing capability through
5 massive parallelism at extremely low power consumption, and have attracted strong interest as a
6 promising option in the search of new computing paradigms.^[154-159] In this section, we will review
7 recent progress in the development of neuromorphic computing hardware based on resistive
8 switching devices. In particular, key requirements for practical emulation of biological elements and
9 learning rules will be discussed, followed by discussions on recent network-level implementations of
10 functional neuromorphic systems.

11 In these discussions we will follow the *memristive systems* theoretical framework, due to both
12 fundamental arguments that this framework elegantly incorporates the native, internal dynamic
13 processes of the device, as well as practical reasons that the analog switching characteristics, desired
14 by neuromorphic computing systems, can be readily predicted and modeled using this approach.

16 4.2.2 Memristive systems theoretical framework

17 The concept of memristors and memristive systems was first proposed by Leon Chua and further
18 generalized a few years later by Chua and Kang.^[12,13] In the generalized *memristive systems*
19 framework, the device dynamics are described mathematically as:

$$20 \quad v = R(w, i)i \quad (3)$$

$$1 \quad \frac{dw}{dt} = f(w, i) \quad (4)$$

2 where v , i , and w are the voltage, current and an internal state variable of the device, respectively.

3 From Equations (3) and (4), it becomes clear that in a memristive system, the resistance cannot be
4 determined directly by the instantaneous inputs (voltage and current). Rather, only the change of
5 the device state can be predicted through the current state and the instantaneous inputs, as shown
6 in Equation (4). As a result, the device's state (represented by variable w) is naturally history-
7 dependent. This simple but powerful mathematical description fundamentally explains the history-
8 dependent hysteresis behaviors observed in RS devices, and implies that the inherent memory effect
9 is a distinct feature of a memristive system.

10 The memristive framework started to draw significant attention after the work by HP labs in 2008
11 that directly linked the experimentally observed RS effects in a TiO_x based device with the memristor
12 concept.^[14] Specifically, the device was modeled as two resistive elements connected in series,
13 corresponding to the V_o -doped region and undoped region, respectively, with an adjustable
14 boundary. In this configuration, the thickness of the doped region considered as the internal state
15 variable, which is in turn modulated by the electric field-induced V_o migration. The dynamic, time
16 evolution of the state variable in response to the applied voltage is mapped to the dynamic Equation
17 (4) in the memristive system framework, and the device exhibits characteristic pinched hysteresis
18 loops with a clear frequency-dependence, as expected from the memristor model.

19 Beyond theoretical importance, the memristive device framework is a very flexible and powerful
20 tool that can help guide the understanding and optimization of devices, as different I - V and dynamic
21 equations, often complex and non-linear, can be used in Equations (3) and (4) to accurately predict

1 the device operation, along with properly identified state variable(s). For example, a model with two
2 resistive elements in parallel was proposed to successfully describe the incremental, analog
3 resistance changes often observed in oxide-based RRAM devices,^[160] in which the conductive
4 channel area was identified as the internal state variable that reflects the nature of localized
5 conductive paths. Additionally, exponential ionic drift, instead of the linear drift model as originally
6 used, was employed to more accurately capture the V_0 migration dynamics at high fields, and a
7 decay term can also be included in the rate equation (4) to account for spontaneous, lateral diffusion
8 of ions.^[160,161]

9 Based on these analyses, it has now been generally accepted that all resistive switching devices
10 can be modeled within the memristive framework by mapping the current Equation (3) and the
11 dynamic Equation (4) to the corresponding physical processes during the device operation, with
12 properly identified internal state variable(s).^[162,163] It should be noted that the dynamic Equation (4)
13 also often implies physical reconfiguration of the material, *e.g.*, through the change of the
14 conductive region length or area due to ionic transport, so that key features of memristive systems,
15 such as incremental conductance changes and other history-dependent changes of the material's
16 properties, directly reflect the reconfigurability of the system. Therefore, proper understanding and
17 modeling of these dynamic processes are essential to effectively utilize the reconfigurable elements
18 towards successful system-level implementations such as building efficient neuromorphic computing
19 hardware.

21 4.2.3 Synaptic plasticity implementations

1 A biological network consists of interconnected neurons through reconfigurable connections –
2 synapses, as illustrated in **Figure 12a**. The connection strength, the synaptic weight, can be
3 strengthened (*via* potentiation) or weakened (*via* depression) depending on the neurons' firing
4 patterns, *e.g.*, spiking timing and rates. The synapse can retain the updated weight^[164] from tens of
5 milliseconds to a few minutes in the form of short-term plasticity^[165], and from minutes (and hours)
6 to days (and months) in the form of long-term plasticity^[166-168]. These effects, termed synaptic
7 plasticity, are believed to play a critical role in the processes of learning and memory.^[169] The ability
8 of memristors to modulate its conductance and evolve into different states, dependent on the input
9 history, makes these devices ideal building blocks as artificial synapses for hardware implementation
10 of neuromorphic systems.^[4,163,170-172] Additionally, analog-type switching, in which the device
11 conductance changes incrementally in response to a series of input signals (spikes), allows simple
12 implementation of online learning.^[4,173] Such incremental conductance modulation effects, as shown
13 in Figure 12b, were experimentally implemented in different types of memristive devices including
14 WO_x , Ag/a-Si, Ta_2O_5 and $\text{Al}_2\text{O}_3/\text{TiO}_2$.^[93,151,160,174]

15 Different synaptic learning rules including timing- and rate-dependent plasticity that enable the
16 efficient operation of biological systems have been successfully demonstrated in memristors. For
17 example, spike timing dependent plasticity (STDP) is a critical biological phenomenon during
18 learning, and states that the synaptic potentiation or depression depends on the relative timing of
19 the pre- and post-synaptic neuron firings. Specifically, if the pre-synaptic neuron fires before the
20 post-synaptic neuron, causality may exist between the two events and the synaptic connection
21 between the two neurons should be strengthened. Additionally, the potentiation should be stronger
22 if the firing events are closer (representing stronger causality). This timing-based synaptic weight

1 change has been emulated in several different memristive systems,^[175,176] where carefully designed
2 pulses are employed to convert the relative timing information from the spiking events to a
3 parameter that can directly program the device, *i.e.* pulse height or duration, depending on the
4 different overlap patterns of the pre- and post-neuron pulses.

6 4.2.4 Biorealistic emulation of synaptic effects

7 Beyond using the gradual conductance changes to emulate desired synaptic weight updates
8 through engineering of the programming pulses, an intriguing alternative is to use the internal ionic
9 dynamic processes in memristors to directly emulate the underlying physical and chemical processes
10 in biological synapses, naturally at different time scales. This approach not only allows more bio-
11 realistic emulation, but also offers significant power advantages and enables direct processing of
12 temporal, spiking inputs.^[177,178]

13 For example, in biological synapses the timing information is naturally encoded internally, and
14 spikes do not overlap with each other. In fact, spiking events are sparse (thus overlapping is very
15 rare) and *sparsity* is a key factor that makes such systems so power efficient. Recent experimental
16 and modeling efforts have shown that the internal dynamic processes within the synapse, *e.g.* the
17 rise and decay of the Ca^{2+} concentration in short-term, essentially provides an internal timing
18 mechanism that enables the different timing- and rate-based synaptic functions observed
19 experimentally.^[179-181] Closer examination of the memristor model shows that similar internal
20 dynamics can also be implemented in semiconductor memristive devices, where the (short-term)
21 dynamics of one state variable can *natively* encode the relative timing information of the external

1 stimuli, and eventually lead to the (long-term) change of another state variable that controls the
2 device conductance, *i.e.* synaptic weight. These devices, with multiple state variables that evolve at
3 different time scales, are termed “second-order” memristors and have been recently demonstrated
4 experimentally.^[97,182] In such a second-order memristor as illustrated in **Figure 13a**, the dynamics of
5 the first state variable, such as the conductive channel area that directly determines the device
6 conductance, is affected by a second state variable that evolves at a much shorter time scale, such
7 as the internal temperature^[97] and the oxygen vacancy mobility^[182]. Similar to the Ca^{2+} dynamics in
8 biological synapses, the rise and spontaneous decay of the short-term state variable provide an
9 internal timing mechanism and allows the device to naturally emulate the rate- and timing-
10 dependent synaptic effects, in a biorealistic fashion.

11 Experimentally, different synaptic functions at both short-term and long-term scales such as pair-
12 pulse facilitation (PPF) and STDP have been directly implemented in a second-order memristor,
13 without manually designed complex waveforms or overlapping pulses, as shown in Figure 13b and
14 c.^[182] The same concept was recently applied to a system based on two memristors, one with short-
15 term dynamics (so-called diffusive memristor) which in turn modulates the second device with long-
16 term conductance changes. The short-term dynamics in the diffusive memristor were directly
17 observed from *in-situ* TEM studies (Figure 13d), in which Ag nanoclusters were found to
18 spontaneously diffuse and disperse due to interfacial energy minimization effects.^[136]

19 The studies on second-order memristors and diffusive memristors reveal that rich internal
20 dynamics can exist in a memristor device, despite its simple appearance. Different physical processes
21 involving ion drift, diffusion, chemical reactions, and Joule heating can be included in the dynamic
22 Equation (4), and allow the device to exhibit diverse dynamic behaviors beyond simple conductance

1 changes. The ability to directly and biorealistically emulate the underlying synaptic processes will
2 likely in turn open the possibility to significantly extend the function and efficiency of hardware
3 systems built with such dynamic devices.

4 On the other hand, the electrochemical processes during the memristor device operation can
5 lead to chemical potential gradients arising from the Nernst potential due to the difference of metal
6 (M) chemical potential from one interface to the other, the diffusion potential because of
7 inhomogeneous distribution of M^{z+} and OH^- ions and the Gibbs-Thomson potential due to the
8 differences in surface free energy between the electrode and the nanoscale sized filament.^[183] These
9 chemical potential gradients, formed in a non-equilibrium state during device operation, in turn
10 introduce electromotive forces and result in the experimentally observed nanobattery effect.^[183]
11 Specifically, it can be argued that the electrochemical origin of the nanobattery effect resembles the
12 electrochemical processes observed in neurobiological systems, and allows the solid-state ionic-
13 based memristive devices to faithfully emulate the fundamental underlying physical and chemical
14 processes in biological structures such as neurons and synapses, beyond what conventional
15 electronic devices can afford.

16 17 *4.2.5 Memristive crossbar as neuromorphic network hardware*

18 An important feature of the cortical system is its large connectivity, enabling massively parallel
19 processing. For example, the human brain has roughly 100 billion neurons with on average 7,000
20 synaptic connections per neuron^[184]. The two-terminal structure of memristive devices can
21 potentially offer high connectivity, high density, and random access that are required to implement

1 large-scale neuromorphic systems, in the form of memristor crossbar arrays. Indeed, a neural
2 network can be readily mapped onto a crossbar structure as shown in Figure 12c and d. In this case,
3 each input neuron (*e.g.* connected to a horizontal electrode) is connected to every output neuron
4 (*e.g.* connected to a vertical electrode) with a memristor device acting as a synapse. Equally
5 importantly, since current directly flows through the memristor between the input and output
6 neurons, the “resistive” feature of the device allows it to directly modulate the information (in the
7 form of current) flow through Ohm’s law. In this sense, the device allows memory and computing
8 functions to be performed simultaneously at the same physical location, without having to move
9 data between different components.

10 Specifically, in a crossbar array, when input voltage pulses are applied to the rows of the crossbar,
11 the output currents at the columns are determined by the product of the input voltage and the
12 stored conductance values in the crossbar matrix. In this way, the memristor crossbar naturally
13 performs the (input) vector- (stored synaptic weight) matrix multiplication operation^[163,170,174,185,186],
14 a key operation in machine learning, analog computing, and other data intensive computing tasks.
15 Note here the vector-matrix operation is achieved directly “for free” through physical phenomena
16 (*e.g.* Ohm’s law and Kirchhoff’s law, where the output is obtained *via* a *single* read of the output
17 current), without having to perform computation and move data between separate processor
18 weight storage units. The ability to change the device conductance *in situ* in turn allows the network
19 to adapt to the input patterns and learn useful features from the input data using online learning
20 algorithms,^[163,170,171,174] thus making the crossbar-based system a natural fit for mapping machining-
21 learning and bio-inspired computing algorithms.

1 Experimental implementations of functional memristor-based neuromorphic networks have been
2 recently demonstrated. In one example, a single layer perceptron was mapped to a 12×12 crossbar
3 array,^[174] using carefully optimized $\text{Al}_2\text{O}_3/\text{TiO}_2$ stacks with nonlinearity and low forming voltages as
4 synaptic devices. Classification of 3×3 black-and-white images into three classes, *e.g.*, z, v, and n,
5 was successfully demonstrated, as shown in **Figure 14a-d**. The neuromorphic network was trained
6 online using a supervised training algorithm, using 30 input images including three ideal letters and
7 nine noisy versions of each letter. After applying each training pattern, the corresponding memristor
8 conductance was updated using the Manhattan update rule^[187], a binary-quantized version of the
9 batch-mode delta rule, as shown in Figure 14c. The weight update was performed in parallel for each
10 column of the crossbar, validating the intrinsic parallelism of the crossbar-based neuromorphic
11 system during learning as well during inference. As shown in Figure 14d, all patterns were correctly
12 classified after on average 23 training iterations, from six training runs with different initial states.

13 In another example, a Ta_2O_5 -based memristor network was used to perform principal component
14 analysis (PCA)^[188-189] by extracting features (*e.g.*, the principal components) from sensory data and
15 performing data clustering using unsupervised, online learning.^[190] 9-dimensional unlabeled data,
16 representing breast cell mass properties measured in 9 categories from the University of Wisconsin
17 Hospital, were mapped to a 2-dimensional output space using a 9×2 crossbar array (inset of Figure
18 14e) to separate benign cells from cancer cells. The two principal components were represented by
19 memristor weights associated with the first two output neurons, and obtained through online
20 learning based on Sanger's rule^[191]. After the unsupervised learning process, the memristor network
21 successfully clustered the 583 test data points into two groups. With the help of a decision boundary
22 obtained using logistic regression, the memristor-based hardware system achieved a classification

1 accuracy of 97.1%, as shown in Figure 14e. These results, obtained experimentally from a small
2 memristor network, are already comparable to results (97.6% accuracy) obtained from direct
3 calculation of the eigenvectors in conventional PCA analysis. The high accuracy that can be obtained
4 experimentally even in the presence of device variations further suggests that memristor-based
5 neuromorphic hardware can indeed be reliably used for efficient feature extraction, using
6 unsupervised, online learning algorithms.

7 A particularly attractive property of biological neural systems is their excellent energy-efficiency
8 in processing massive, complex data. Sparse coding, in which massive sensory inputs such as vision,
9 auditory, and touch are represented by a relatively small set of strongly activated neurons, has been
10 believed as an essential underlying mechanism by which biological systems can efficiently handle the
11 large amount of high-dimensional data.^[192-194] Experimental implementation of sparse coding
12 algorithms was recently demonstrated in a memristive crossbar neural network based on WO_x
13 devices (Figure 14f), in which the critical operations required by sparse coding, including pattern
14 matching and lateral neuron inhibition, were effectively mapped on to the crossbar structure.^[195]
15 Specifically, input images were encoded by (sparse) output neuron activities, so that the input can
16 be efficiently reconstructed using only a few active neurons and the associated dictionary elements
17 stored in the memristor crossbar. Besides feedforward connections, lateral inhibition among all
18 output neurons is a critical feature of sparse coding, as inhibition allows an optimal representation
19 to be obtained, out of many possible representations.^[196] In Ref. 195, lateral neuron inhibition was
20 achieved by iterative forward and backward passes through the same crossbar, where the
21 reconstructed input (obtained from the backward pass) is subtracted from the original input, thus
22 essentially allowing active neurons to suppress other neurons that share similar dictionary elements.

1 Careful analysis of the experimentally-obtained network dynamics verified that the network can
2 indeed find an efficient and optimal solution out of several possible solutions, while adjusting the
3 sparsity parameter can balance the sparsity requirement and reconstruction accuracy. Finally, the
4 system was used to successfully process natural images (120×120 pixels) as shown in Figure 14g,
5 following predictions of the algorithm. These successful demonstrations clearly highlight the
6 potential of memristive systems to process complex data with desired throughput and energy-
7 efficiency in demanding environments, and serve to stimulate continued device and system
8 developments based on these nanoscale, physically reconfigurable components.

10 *4.2.6 Continued device optimizations for neuromorphic computing*

11 Successful, large-scale implementation of memristor-based neuromorphic networks critically
12 depends on the device properties, thus continued optimizations are still required beyond the initial
13 demonstrations discussed above. Specifically, criteria for efficient logic operations such as
14 neuromorphic computing are different from those for memory applications where the key is to be
15 able to distinguish two stored states. For example, binary devices with large on/off are ideally suited
16 for memory applications, while analog-type switching devices are more desirable for neuromorphic
17 applications. Such analog switching devices, unfortunately typically exhibit smaller dynamic ranges
18 (usually < 10) which can cause systematic errors.^[197] On the other hand, linear conductance
19 modulation, where the conductance change can be predicted by the applied pulse width or
20 amplitude, without prior knowledge of the current conductance value, is desirable for online
21 learning implementations while practical devices typically show non-linear, asymmetric conductance

1 modulation characteristics. In a typical device, the conductance can be increased quickly from a
2 state close to the HRS, and the rate of change will slow dramatically as the device approaches the
3 LRS.^[94,96] Combined with the device variation issue, the limited dynamic range and the non-linear
4 weight updates can significantly degrade the performance of memristor-based networks for
5 neuromorphic applications,^[173,197] with asymmetric nonlinearity of weight update having the most
6 significant effect on the performance degradation while symmetric nonlinearity is less significant.^[198]
7 Additionally, it has been shown that read noise due to thermal, 1/f, and random telegraph effects, as
8 well as write noise due to temporal variations can also degrade the network performance.^[198] The
9 high switching current (typically > 100 μ A) in typical analog-switching devices will in turn affect the
10 system's power efficiency and limit the network size due to series-resistance issues. Finally, forming-
11 free memristor devices or devices with low enough forming voltages are required for reliable
12 operation of the passive crossbar array. Otherwise, during the forming process half-selected cells
13 that are already formed may be exposed to a high voltage and can become damaged (*e.g.*, stuck-at-
14 1), even with carefully designed protective voltage schemes.^[174,190]

15 These device optimizations will likely require synergistic efforts from materials scientists,
16 physicists, chemists and device and computer engineers, using systematic investigations starting
17 from first-principle calculations and atomic-level characterizations that reveal the fundamental ionic
18 processes and structural changes, to careful device engineering and characterizations, and clever
19 architecture design and optimizations that can maximally utilize the properties of the devices while
20 sufficiently mitigating the device non-ideality effects.

21

1 4.3. Resistive switching devices – logic applications

2 As the transistor size is approaching fundamental physical limits with diminishing returns on
3 performance, alternative computing architectures including reconfigurable computing systems,
4 analog computing, and neuromorphic computing, have been extensively studied to extend
5 computing functionality and increase the logic performance.^[4,172,199] Below we discuss a few
6 examples where RRAM devices have been used in these alternative logic computing applications.

8 4.3.1. Reconfigurable logic circuits

9 It has been demonstrated that RRAM crossbar arrays on top of CMOS circuits can serve as
10 reconfigurable interconnects that wire up CMOS components so that different logic gates can be
11 connected to form custom configurations, *i.e.* in the form of Field-Programmable Gated Array (FPGA)
12 systems.^[200] In this RRAM/CMOS hybrid FPGA system, RRAM devices can work as the configuration
13 bits and switches in a data routing network. The CMOS-compatible fabrication processes, small
14 device/switch size, and nonvolatile functionality that helps eliminate power consumption for refresh
15 make this approach attractive. Indeed, TiO₂-based RRAM/CMOS integrated circuits with FPGA
16 functionality were experimentally demonstrated.^[200] Successful logic operations such as NOT, AND,
17 NAND, NOR and D flip-flop, and more importantly, reconfiguration of the CMOS gate connections
18 were achieved by programming the states of the RRAM device overlaying the CMOS circuitry.
19 Additionally, the defect-tolerant capability of the hybrid integrated circuits based on the redundant
20 data paths in crossbar-array structure was demonstrated as pointed by previous numerical
21 simulations.^[199,200]

1

2 **4.3.2. Material implication logic**

3 It is also shown that RRAM devices can be used for logic operations in which the resistance of the
4 memristive switches was used to directly drive subsequent logic gates in the form of material
5 implication (IMP) logic. IMP is a fundamental Boolean logic operation equivalent to $(\text{NOT } p) \text{OR } q$ (q
6 and q are logic variables), and was successfully executed with two memristive switches in 1×17
7 crossbar array.^[201] In this implementation, logic values (*i.e.*, “0” and “1”) are represented by
8 resistance of the memristive devices and the IMP logic operation was achieved by applying voltage
9 pulses resulting in different modulation of device resistance depending on the initial states of the
10 two devices. Here, the conditional toggling properties for IMP were achieved through the voltage
11 divider effect between the memristive devices and a carefully selected series resistor, with $R_{\text{on}} <$
12 $R_{\text{series}} < R_{\text{off}}$, so that whether the state (resistance) of output device q can be modified depends on the
13 state (resistance) of the input device p . Subsequently, the universal NAND operation was
14 demonstrated using the IMP operation and an erasing operation in three interconnected memristive
15 junctions.^[201]

16

17 **4.4. Beyond resistance change**

18 The ability to directly change the materials’ chemical composition through ionic processes allows
19 a host of physical properties to be reconfigured, beyond electrical resistance changes. In principle,
20 optical, magnetic, and mechanical properties of the material can also be potentially modified on-
21 demand, using simple electrical signals, thus offering tremendous opportunities in building

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1 multifunctional and reconfigurable systems. In this section, we discuss other functions that can
2 benefit from material reconfiguration, involving changes in composition, structure, geometry, and
3 chemical states induced by electrochemical reactions and ionic transport, and highlight the potential
4 of these ionic systems as an important building block in multifunctional-embedded applications.^{[202-}

5 ^{206]}

7 *4.4.1. Electrically reconfigurable optical properties*

8 In an ECM device a new material, in the form of metal filaments, is created inside the host
9 dielectric material during the RS process. As a result, the process can significantly alter the film's
10 optical properties as well as electrical properties.

12 *Plasmonic memristor:*

13 In plasmonics, the signal is encoded as surface plasmon polaritons (SPPs) that propagate along a
14 metal–dielectric interface with strong coupling with free electrons at the metal surface.^[207] The SPP
15 modes can be confined below the diffraction limit, offering potential for the implementation of
16 nanoscale optical devices and circuits.^[208,209] The built-in metal-dielectric interface in the MIM
17 structure of RRAM devices makes RRAM devices naturally compatible with plasmonic circuits and
18 can lead to devices such as plasmonic switches and modulators, that provide coupled electro-optical
19 memristive effects. These “plasmonic memristors” can open new opportunities for electro-optical
20 system design and implementation, by offering low-power, electrically-controlled, non-volatile yet

1 reversible optical modulations, beyond conventional electro-optical devices based on thermal^[210],
2 free carrier dispersion^[211], and Pockels effects^[209].

3 In a recent study, a Ag-based ECM device was used as a plasmonic component that is directly
4 integrated on a silicon-on-insulator (SOI) waveguide, shown in **Figure 15a and b.**^[202] The propagated
5 light through the waveguide is coupled to the fundamental plasmonic mode supported by the
6 memristive device having a Ag metal electrode, and thus different levels of optical transmission can
7 be obtained depending on whether a metallic filament exists inside the a-Si layer, schematically
8 shown in Figure 15c. Reversible electrical modulation of the optical signal with a clear optical
9 hysteresis was demonstrated (Figure 15d). Here, the optical bistable behavior showing different
10 optical transmission is attributed to the changes of the absorption and scattering loss of the
11 fundamental plasmonic mode due to the formation/annihilation of the lossy metal filament. In
12 particular, by applying a write voltage, the ECM device is switched ON in which a Ag filament is
13 formed inside the a-Si waveguide that increases scattering of the propagating light, leading to lower
14 light transmission (Figure 15c left). Note since the plasmonic field is highly confined in the a-Si
15 insulating layer of the MIM stack^[207], even minuate changes of the insulating layer due to the
16 formation of a nanoscale Ag filament can lead to measurable changes in transmission. After the
17 application of an erase voltage, the transmitted optical signal recovers back to its initial value, due to
18 erase of the filament and reduced scattering in the a-Si waveguide (Figure 15c right). In a
19 subsequent study, a high extinction ratio of 12 dB was demonstrated using a similar concept in a
20 coupled plasmonic/memristive device, as shown in Figure 15e, demonstrating the functionality of
21 plasmonic memristor (*i.e.*, the electrical-control of optical properties).^[203]

22

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1 *Memory with optical readout functionality:*

2 Plasmonic memristor with optical readout functionality can be utilized to couple information
3 storage with optical communication, enabling optical readout (\sim ps) that is much faster (\sim 2 order of
4 magnitude) than electrical readout and reducing the number of optical-to-electrical signal
5 conversions in optical communication systems.^[212] Recent experimental demonstrations of optical
6 readout of an RRAM device integrated with an SOI waveguide using a 1550 nm-wavelength light
7 system that is widely used in telecommunications support the feasibility of this application.^[202-204]

8
9 *Latching optical switch:*

10 Another interesting characteristic of the plasmonic memristor is the latching behavior that is
11 highly desirable for low power optical switch applications.^{[202][203]} Specifically, since the optical
12 modulation is based on the direct, atomic reconfiguration of the waveguide material that is both
13 non-volatile and reversible, latching optical switches, where the optical state can be maintained
14 without constant energy consumption, can be achieved at the nanometer scale as shown in Figure
15 15d and e. The power consumption can thus be significantly reduced compared to conventional
16 optical switches based on volatile switches that require a constant voltage supply.^[207] Plasmonic
17 memristors can also offer very low energy consumption during the switching process (*e.g.* 12.5
18 nW)^[204] compared with devices based on thermo-plasmonic effects (*e.g.* 13.1 mW).^[213] Optical
19 modulators working in the visible-range that can achieve reflectance changes up to 78%, operated
20 within an ultralow voltage of 100 mV have also been recently demonstrated in a Ag/Al₂O₃/ITO device
21 structure.^[214]

1

2 *Atomic-scale optical modulator:*

3 Historically, although electronic devices have been aggressively scaled down to nanometer and
4 even atomic dimensions in the form of single atom transistors and atomic switches^[215,216], photonic
5 devices generally need to be much larger to allow sufficient light coupling and readout.^[217-219]

6 Recently, prototype photonic systems that can operate at the atomic scale have been developed by
7 placing a single atom in a photonic cavity, which can enhance the weak interaction between optical
8 photons and matter on the atomic scale.^[219,220] In these demonstrations, atoms were delivered to
9 optical resonators through slow laser cooling and magneto-optic trapping processes, requiring high
10 vacuum and very low temperature. The ability to reliably move individual atoms at high enough
11 speed, at room temperature in solid-state devices, will significantly help the development of
12 nanoscale integrated photonic circuits.

13 An atomic-scale optical switch was demonstrated by employing a resistive switching device showing
14 conductance quantization.^[204] The integrated Ag/a-Si/Pt memristive device in a planar structure on
15 top of a silicon waveguide enables the creation of a plasmonic cavity by moving a single or at most a
16 few atoms at the metallic tip between two planar electrodes, as shown in Figure 16a. The relocation
17 of metallic atoms short-circuits the two pads and modifies the electric field in the junction, causing
18 both resistive switching with quantized conductance and plasmonic switching (Figure 16b). Indeed,
19 abrupt switching characteristics between two distinct plasmonic resonance states were observed in
20 the device, acting as an atomic scale optical switch, as shown in Figure 16c and d. Additionally, a
21 blue-shift of the plasmonic resonance was observed in the LRS (Figure 16e) as a result of the

1 reduction of the local field due to the present of the filament, further confirming that the distinct
2 plasmonic resonance states are a result of the atomic reconfiguration that accompanies electrical
3 switching effects. Additionally, electro-optical switching in the visible spectra range with a very small
4 active volume of $(5\text{nm})^3$, comparable to the sizes of modern electronic devices and much smaller
5 than common optical devices, was demonstrated based on atomic reconfiguration in a Au/Al₂O₃/Ag
6 device, as shown in Figure 16f and g, confirming the potential of nanoscale integrated photonic
7 circuits through material reconfiguration based on ionic effects.^[221]

8 9 4.4.2. Electrically reconfigurable magnetic properties

10 In spintronics, modulations of the materials' magnetic properties have been traditionally
11 achieved using either external magnetic fields or induced fields from electric currents, neither is very
12 scalable. More efficient control of magnetization has been recently achieved through the spin-
13 transfer torque (STT) effect in ferromagnetic nanostructures such as those in a magnetic tunnel
14 junction.^[222-224] However, continued decrease of the current that drives the STT effect without
15 sacrificing switching speed and data retention is still on on-going task.^[223,225] As a result, achieving
16 magnetism control by means of an electric field (instead of current) is currently the focus of
17 extensive investigation that can lead to higher speed, lower power consumption, and improved
18 reliability.^[225-228] Electric field control of magnetism has been suggested through a number of
19 magnetoelectric coupling mechanisms,^[225] including carrier density modulation,^[229] electron orbital
20 occupation modulation,^[226] and exchange coupling in multiferroics.^[228,230] In addition to these
21 magnetoelectric effects, the ability of direct, *in-situ* and on-demand modulation of the

1 ferromagnetic materials through field-driven ionic processes (*i.e.*, electric-field control of magnetism
2 based on atomic reconfiguration) can offer potential for strong modulation, low-power, and non-
3 volatility and can significantly broaden the potential of spintronic and magnetic materials and
4 devices.

5

6 *Magneto-ionic effect:*

7 In a recent study on a Co/GdO_x bilayer structure sandwiched by electrodes (**Figure 17a**), bipolar,
8 voltage-controlled annihilation and restoration of magnetic hysteresis were observed, as shown in
9 Figure 17b.^[205] The pronounced, non-volatile changes in magnetic anisotropy energy (MAE) were
10 attributed to the modulation of the metal-oxide interface by an electric field. Specifically, the oxygen
11 concentration in the Co layer was reversibly modulated through field-driven oxygen ion migration
12 during the voltage cycling, as confirmed through STEM-EELS analysis (Figure 17c). Changes in the
13 oxygen concentration modulates Co's oxidation state and in turn the film's magnetic properties,
14 leading to the observed modulation of the magnetic hysteresis. The magneto-ionic effect can lead to
15 changes in the interfacial magnetic anisotropy energy by >0.75 erg·cm⁻² at 2 V bias, outperforming
16 efficiency achieved by conventional magneto-electric coupling mechanisms.^[205] The ionic transport
17 nature was examined through temperature, voltage, and oxide thickness dependence, where faster
18 switching speed was observed at higher temperature, under larger voltage, and in thinner oxides,
19 consistent with the proposed mechanism. The localized, non-destructive magneto-ionic coupling
20 effect allows reversible tuning of the magnetic properties with excellent spatial control, unlike
21 previous approaches that rely on irreversible ion-beam irradiation processes^[231,232]. For example,

1 localized laser heating facilitates ionic transport only in the illuminated area, enabling selective MAE
2 programming as shown in Figure 17d. These results suggest a path towards electrical tuning of the
3 interfacial chemistry in reconfigurable systems, allowing a broad range of physical and chemical
4 properties to be modulated.

5

6 *Control of magnetism beyond interface/surface:*

7 One distinct feature of the magneto-ionic effect is the possibility to control magnetism in the bulk
8 material, beyond the interface/surface in heterostructures. The aforementioned magnetoelectric
9 effects such as charge carrier modulation and orbital reconstruction are mostly limited to small
10 regions of the material^[233,234], *i.e.* only a few monolayers in close proximity to the interfaces or
11 surfaces, due to the short screening length in ferromagnetic metals^[235] and semiconductors^[236] and
12 the underlying mechanism based on interfacial chemical bonding, respectively.^[225,237,238] On the
13 other hand, it has been demonstrated that electrically-induced ion migration and electrochemical
14 processes modulate the materials' magnetic properties in the bulk form beyond the
15 interfaces.^[237,239,240] Reversible modulation of magnetic moments shown in **Figure 18a** was achieved
16 by electrochemically-driven lithium intercalation and de-intercalation in a ferromagnetic iron oxide
17 spinel, $\gamma\text{-Fe}_2\text{O}_3$.^[237] As Li ions are inserted into the iron oxide, donated electrons from the Li ions lead
18 to chemical reduction of the octahedral-site Fe ions, from Fe^{3+} ($5\mu_B$) to Fe^{2+} ($4\mu_B$) shown in Figure
19 18b, as confirmed by an analysis of the Mössbauer data exhibiting an increase of the Fe^{2+} fraction
20 from 3% to 33% after lithiation.^[237] The valence change of the octahedral-site Fe ions and the
21 associated magnetic state change in turn directly lead to the modulation of the overall bulk

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1 magnetism. In this approach, the amount of the ion-intercalation should be carefully controlled to
2 avoid irreversible structural destruction. The critical limit is ~ 1 Li per formula unit,^[237] and 1 Bohr
3 magnetron (μ_B) can be reduced by one intercalated Li-ion. Based on this understanding and
4 intuition, ferrite systems that have precisely 1 μ_B net magnetic moment per formula unit such as
5 CuFeO_4 was further investigated to obtain large and reversible magnetic property modulations
6 towards on-and-off magnetism.^[239] Upon the electrically-controlled Li-ion exchange, a large change
7 in magnetization was achieved in CuFeO_4 system through the reduction of Cu^{2+} to Cu^{1+} ions, although
8 fully on-and-off switching was still not obtained due to the increased net magnetic moment from the
9 exchange of octahedral-site Cu^{2+} and tetrahedral-site Fe^{3+} ions in the as-prepared nanoparticles.^[241]
10 Additionally, shift of the Neél temperature was observed in the Li-ion exchange processes from
11 which $\sim 70\%$ magnetic response was obtained through magnetic phase transition in ZnFe_2O_4 , as
12 shown in Figure 18c.^[239]

13 The magneto-ionic motion based on oxygen migration can also extend far beyond the interface
14 limit. In a $\text{Si/Pd/AlO}_x/\text{GdO}_x/\text{Co/Pd}$ stack, polarized neutron reflectometry (PNR) depth profile analysis
15 verified significant field-driven oxygen migration into the Co layer throughout the entire thickness
16 (15 nm), as shown in Figure 18d.^[240] This O migration accompanied the reduction in the oscillation
17 amplitude of spin asymmetry, indicating a decrease of magnetization saturation and/or structural
18 changes. The oxidation of the Co layer was also verified from XAS measurements in which additional
19 peaks were observed in electrically programmed samples whereas the peaks were not observed in
20 the as-grown sample. Accompanying these material modulations, significant changes in coercivity
21 (68%) and remanent (55%) magnetization were observed after electric and thermal treatment
22 compared with the as-grown sample.

1

2 *Magnetization switching coupled with resistance change:*

3 Nanoscale electrical control of magnetism through ion migration that is directly coupled with
4 resistance change was also observed^[206], paving ways to tailor the magnetic properties of materials
5 at the nanoscale. In a device based on a lithium ferrite (LiFe_5O_8) structure^[242], magnetization
6 switching was observed accompanying resistive switching through the application of an external
7 electric field, as shown in **Figure 19a-c**. Magnetic and electrical measurements verified that electric
8 field-induced atomic reconfiguration can induce non-volatile and reversible modulation of the local
9 magnetization at room temperature. Specifically, de-intercalation/intercalation of Li ions driven by
10 the electric field alters the chemical state of the iron ions, leading to modulations of the magnetic
11 moment.^[243] At the same time, the Li-ion migration modulates the local concentration of electron
12 hopping centers - Li vacancies, leading to the resistive switching effects. Scanning magnetic-force
13 microscopy (MFM) analysis shows that the areas with significant magnetization changes exactly
14 overlap with the high conductivity areas in a C-AFM current map, confirming the magnetization
15 modulation originates from the same underlying atomic processes with resistive switching effects.
16 The dynamic evolution of magnetic domains under voltage bias was also observed, showing an initial
17 switching at a localized region followed by lateral expansion, again consistent with the
18 characteristics of filament formation/growth processes observed in resistive switching devices. By
19 utilizing the coupled magnetization and resistance changes, a multi-mode storage device in which
20 two bits, represented respectively by the direction and the intensity of the film's magnetization, can
21 be stored in one physical cell. The magnetization direction and the intensity were independently
22 switched by an external magnetic field and an electric pulse, respectively. Subsequently, the stored

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1 data with four different states per cell can be properly read out through MFM measurements (Figure
2 19d).

3 4 **5. Challenges, looking into the future and conclusions**

5 In this article, we discussed atomic-level reconfiguration of materials through field-driven ionic
6 processes, which we believe may lead to a new class of fundamental building blocks for future
7 electronic, optical, and energy devices and multifunctional systems. Recent studies have clearly
8 shown that physical reconfiguration of the material can occur through *in-situ*, on-demand
9 manipulation of masses (atoms) by electrical stimuli, leading to new functionalities beyond
10 conventional mobile charge-based effects. Experimental studies such as *in-situ* TEM analysis
11 combined with spectroscopy characterizations have provided direct evidence showing that coupled
12 ionic and electronic effects, such as chemical redox reactions and long-distance ion migration, can
13 occur in solid-state semiconducting and dielectric films. Theoretical studies based on molecular
14 dynamics and first-principles calculations have provided insight into the fundamental physical and
15 chemical processes underlying the atomic processes involved in the material reconfigurations. A
16 broad range of applications based on the reconfigurable materials and systems, ranging from non-
17 volatile data storage, neuromorphic computing, to optical switches and modulators, have been
18 proposed and demonstrated, verifying the potential of such reconfigurable materials.

19 Looking into the future, a number of challenges at the materials/device level still need to be
20 addressed to pave the way for practical, large-scale applications. For example, both cation- and
21 anion-based devices suffer from variability issues that affect their applications in large-sized arrays.

1 Cation-based devices are often found to show moderate endurance while anion-based devices
2 typically exhibit low on/off ratio and high operating current. Advanced characterization techniques
3 that can provide high spatial, temporal, and spectral resolution, thus enabling real-time, non-
4 destructive analysis of the reconfiguration dynamics at the atomic scale will dramatically speed up
5 the material and device developments. With the continued progress in materials, devices and
6 characterization techniques, we expect precise control of matter at the individual atom level can be
7 eventually achieved, which will enable precise design and engineering of these reconfigurable
8 materials. This level of control not only improves device scaling but also help address issues such as
9 device variability, endurance and energy. Indeed, the principle of atomic-scale reconfiguration of
10 materials has been demonstrated in studies of atomic switches in which quantum conductance
11 effects were observed due to the formation of quantum point contacts (QPCs) having a single atom
12 or few atoms in the filament.^[216,244,245] Highly reproducible and controllable QPC formations have
13 been demonstrated experimentally in both ECM and VCM devices, and the detailed atomic
14 configurations of QPCs consisting of only a few atoms were evaluated by first-principles
15 calculations.^[246-249] The formation of Ag clusters with an atomic contact bridging an STM tip and a
16 substrate was also directly observed, supporting the feasibility of atomic control of mass movement
17 through electrochemical processes^[250]. Continued research will still be needed to improve the device
18 reliability and to overcome the instability near the point-contact regime^[76,251] and consequent
19 retention loss^[252]. Additionally, we note that recent studies on manipulation of individual atoms (or
20 vacancies) with atomic precision using a STM tip method^[253-255] have already led to the
21 demonstration of a kilobyte rewritable atomic memory, in which data are encoded by the positions
22 of individual chlorine vacancies, as shown in **Figure 20**. With further developments of the coupled

1 ionic/electronic systems as discussed in this Review that may allow *in situ* control of single atoms in
2 solid-state materials/systems, an ultimate goal may be to replace the slow and low-temperature
3 STM-tip based process with simple field-driven processes that can be easily implemented at room
4 temperature, after the circuits have been completed. This kind of capability will no doubt lead to
5 paradigm shifts in the design and fabrication of electronic, optical and energy systems, and will
6 require truly multidisciplinary efforts from physicists, chemists, materials scientists, engineers and
7 computer scientists to take full advantage of such reconfigurable materials.

8

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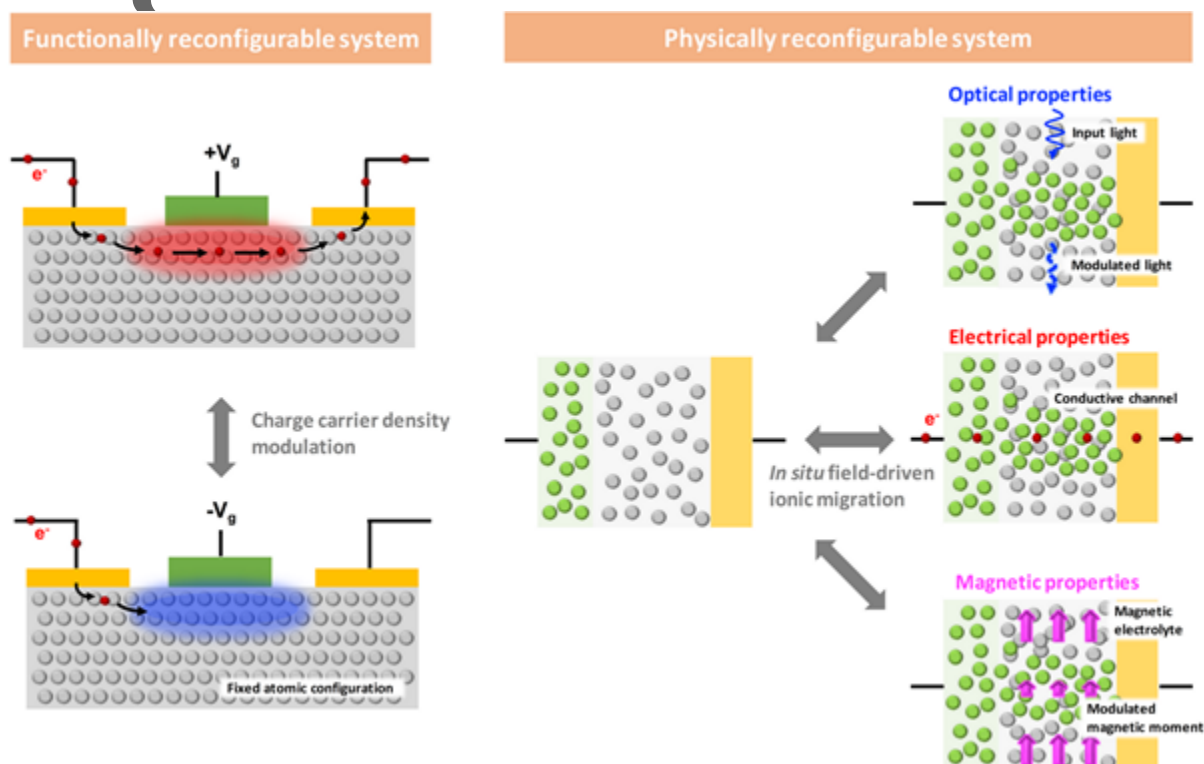
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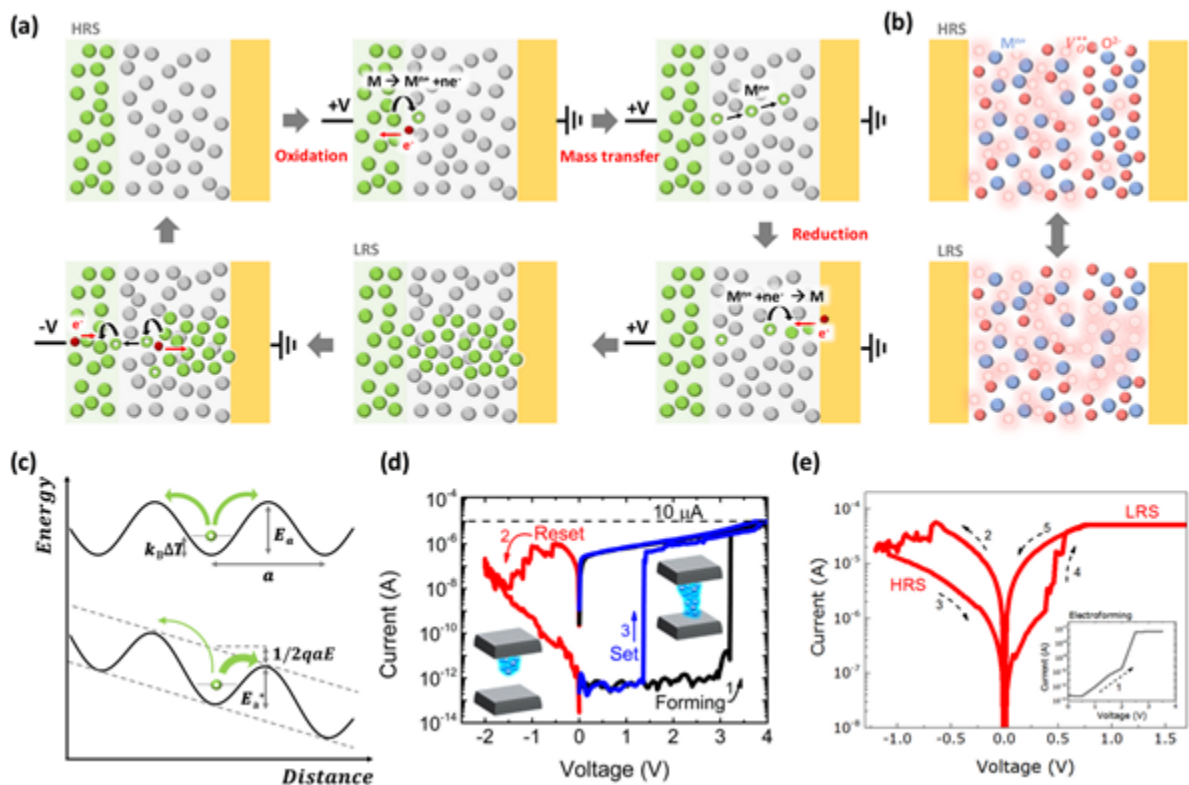
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5
 6 **Figure 1.** Schematic illustration showing how functionally (left) and physically (right) reconfigurable
 7 systems are differently operated. The operation of conventional electronic devices such as
 8 transistors is based on the modulation of mobile charge carriers, using a fixed atomic configuration;
 9 while the operation of coupled ionic/electronic devices involves physical changes based on ionic
 10 migration and chemical reactions, allowing in situ, on demand control of electronic, optical, and
 11 magnetic properties of materials and devices. For example, transmission through a dielectric-based

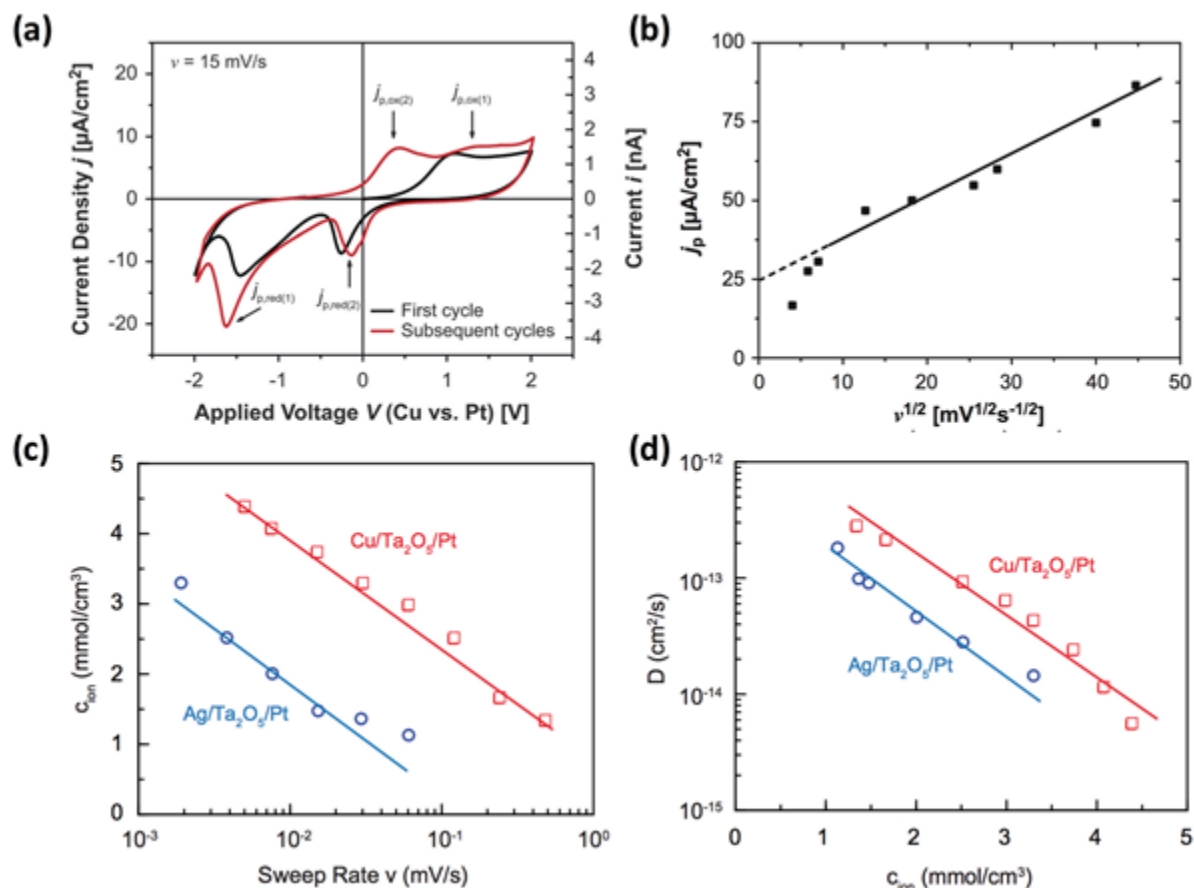
1 waveguide can be modulated by the formation of an electrically-formed metallic filament (upper
 2 panel). Resistive switching behaviors have been commonly observed in anion- and cation-based
 3 memory devices (middle panel), and magnetization of a ferromagnetic material may be controlled
 4 by ion intercalation/de-intercalation, followed by redox reactions, leading to changes of the
 5 magnetic moment (bottom panel).



6
 7 **Figure 2.** a,b) Schematic illustration of the resistive switching process in ECM (a) and VCM (b)
 8 devices. c) Field-driven acceleration of ionic transport, where the effective energy barrier is lowered
 9 by the applied field. d) $I-V$ characteristics of a typical ECM device showing large on/off ratio with
 10 high off-state resistance. Reproduced with permission.^[26] Copyright 2014, American Chemical
 11 Society. e) $I-V$ characteristics of a typical VCM device. Reproduced with permission.^[56] Copyright

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2

3 **Figure 3.** a) Cyclic voltammetry of a Cu/SiO₂/Pt cell, showing current density peaks corresponding to
 4 Cu-ion oxidation and reduction processes. Reproduced with permission.^[28] Copyright 2012, Royal

5 Society of Chemistry. b) Current density peak value vs voltage sweep rate, showing a linear relation

6 between j_p and $v^{1/2}$. Reproduced with permission.^[27] Copyright 2011, AIP Publishing LLC. c) ion

7 concentration vs voltage sweep rate, showing a decrease of the ion concentration as the sweep rate

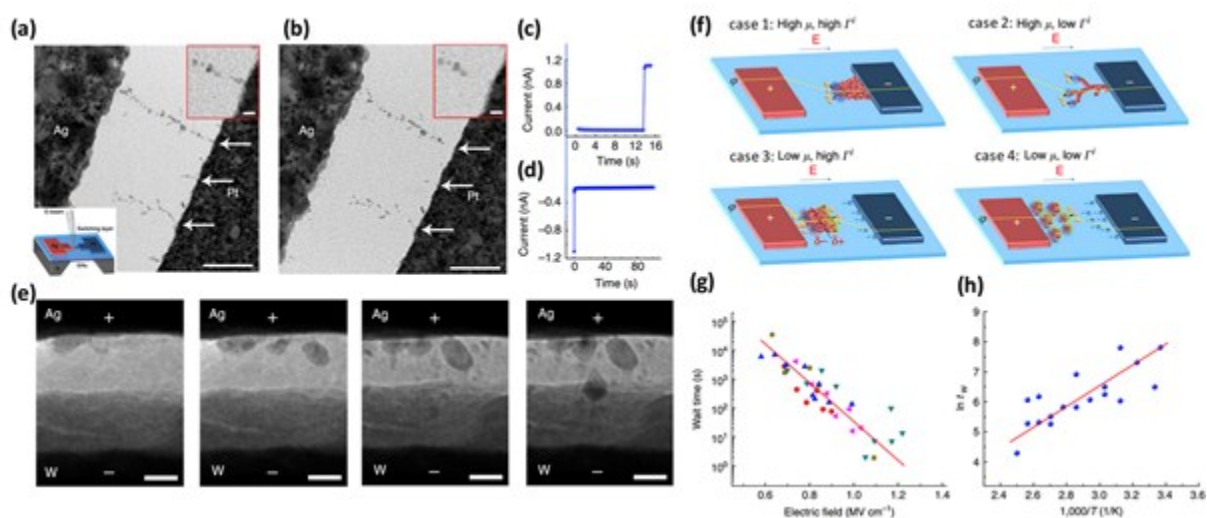
8 increases.^[28] d) Ion diffusion coefficient vs ion concentration, showing a decrease of the ion diffusion

9 coefficient as ion concentration increases, and implying that the electrolyte behaves as a

10 concentrated solution with non-negligible ion–ion interactions. Reproduced with permission.^[29]

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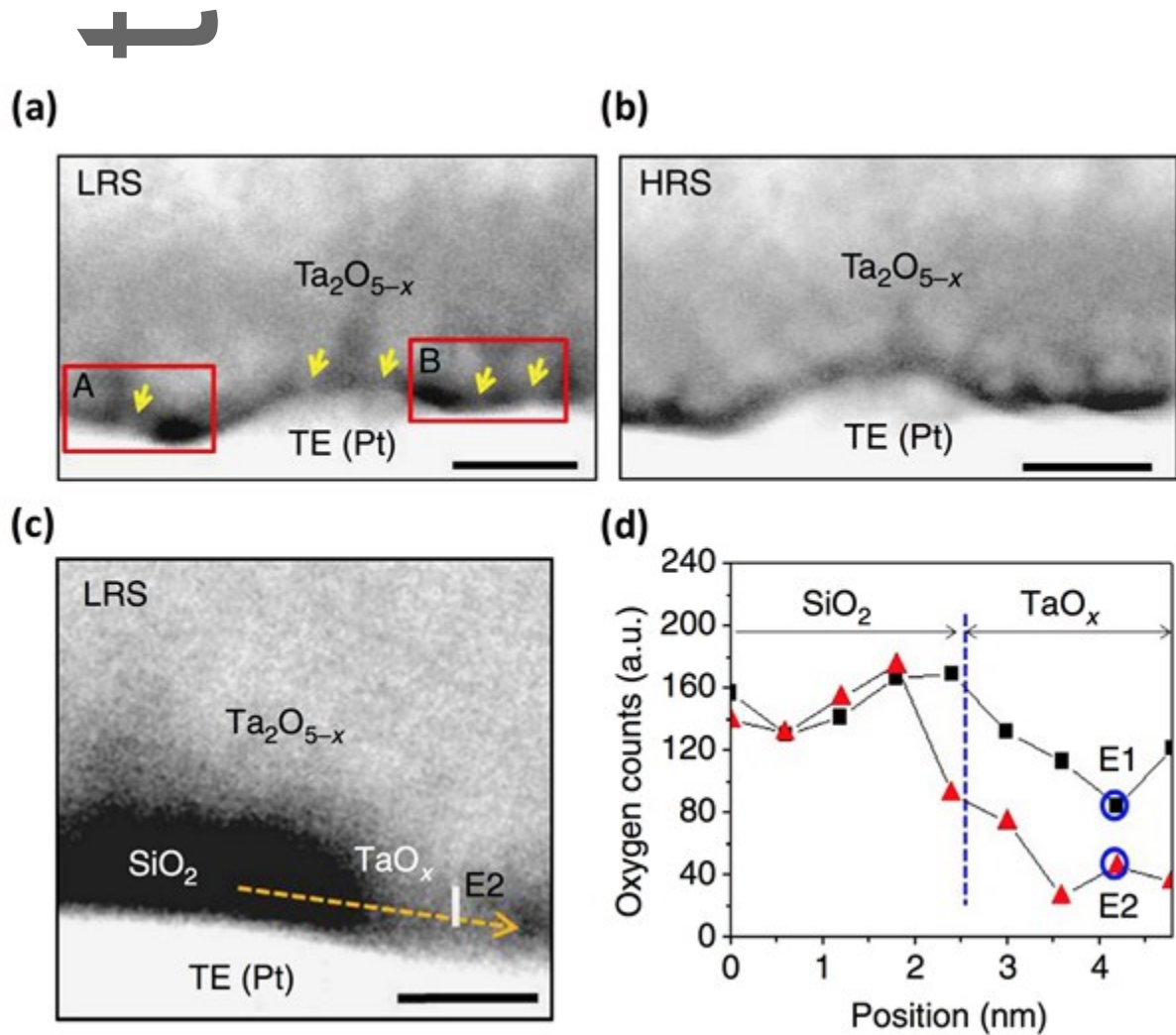
2

3 **Figure 4.** a) TEM image of an Ag/SiO₂/Pt-based device after the forming process, showing the
 4 formation of conductive filaments. Inset: schematic of the planar-structured device for TEM
 5 observation. b) TEM image of the same device in (a) after reset showing the dissolution of the
 6 filaments at the inert electrode interface. c) $I-t$ curve during the forming process under constant
 7 bias. The abrupt current increase corresponds to the formation of filaments shown in (a). d) $I-t$ curve
 8 during the reset process under constant bias. The abrupt current drop corresponds to the rupture of
 9 filaments shown in (b). e) Real-time structural evolutions of an Ag/a-Si/W-based device obtained
 10 through in-situ TEM observation, showing dynamic filament growth process that initiates from the
 11 reactive electrode. Reproduced with permission.^[31] Copyright 2012, Nature Publishing Group. f)
 12 Schematic illustration of filament growth dynamics at different cases, depending on kinetic
 13 parameters such as the redox rates, Γ , and ion mobility, μ . Reproduced with permission.^[32]
 14 Copyright 2014, Nature Publishing Group. g) Wait time as a function of the applied electric field
 15 during the forming process, showing exponential dependence consistent with Equation (1). h)

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1 Arrhenius plot of the wait time vs temperature, consistent with Equation (1).^[31]

2

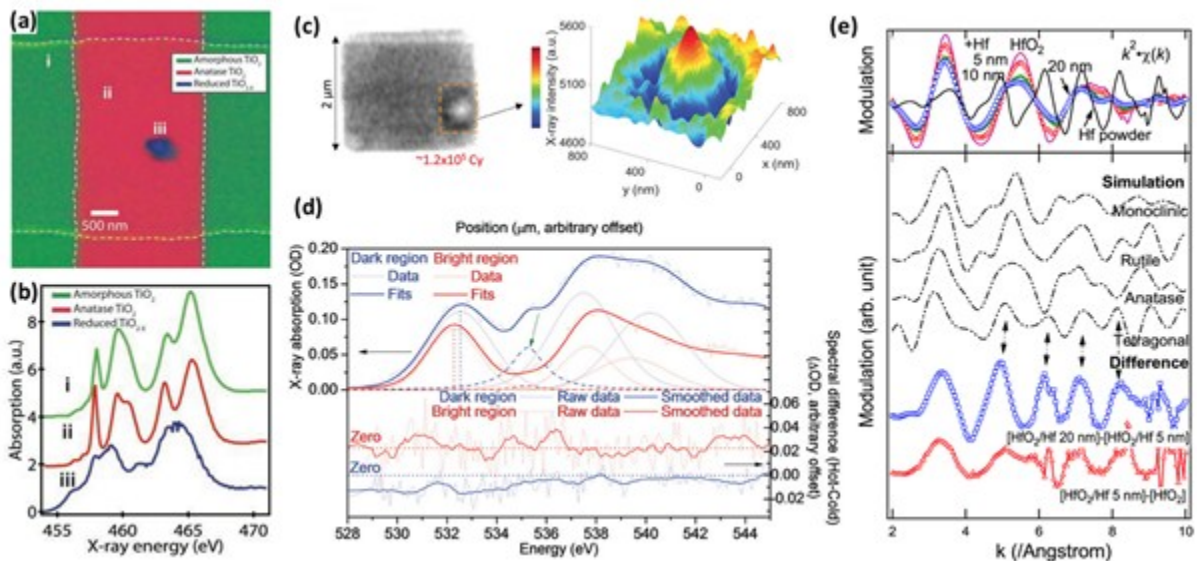


3

4 **Figure 5.** a,b) In-situ HAADF-SETM image of a Pt/SiO₂/Ta₂O_{5-x}/TaO_y/Pt device in LRS (a) and HRS (b),
 5 showing the formation/annihilation of nanoscale channels in a thin SiO₂ layer, as indicated by yellow
 6 arrows in (a). Scale bars: 5 nm. c) Magnified image of a channel region in LRS. Scale bar: 2 nm. d)
 7 Oxygen EELS line scan profiles obtained from LRS (red triangle) and HRS (black square) along the
 8 horizontal direction as indicated by the yellow arrow in (c). Significant changes of oxygen
 9 concentration between the LRS and the HRS are observed in the channel region. Reproduced with

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2

3 **Figure 6.** a) Scanning Transmission X-ray Microscopy (STXM) image of a TiO_2 -based device after
4 electrical cycling, showing structural changes and the formation of a localized channel.^[50] b)

5 Corresponding Ti L-edge X-ray absorption spectra from the three regions in (a); stoichiometric
6 amorphous TiO_2 (green), stoichiometric anatase TiO_2 (pink), and reduced TiO_{2-x} (blue). Reproduced
7 with permission.^[50] Copyright 2010, John Wiley and Sons. c) O K-edge X-ray transmission

8 map (left) and 3D intensity plot (right) of a Ta_2O_5 -based device after 120 000 electrical cycles. A ring-
9 like feature with a bright core and a dark perimeter, corresponding to low and high O

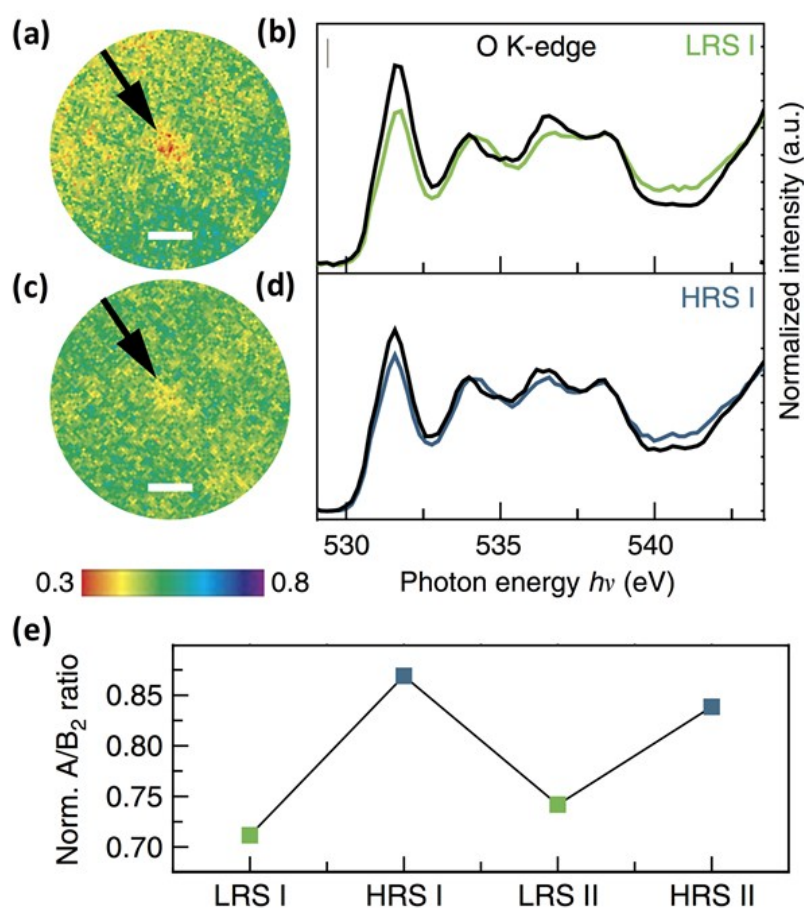
10 concentrations, respectively, is observed.^[52] d) Upper panel: O K-edge absorption spectra from the
11 bright (red) and dark (blue) regions, showing the downshift of the lowest conduction band position
12 (indicated by vertical dotted lines) in the bright region leading to a higher electric conductivity.

13 Lower panel: O K-edge spectra difference between the hot and cold states (corresponding to with
14 and without current flow, respectively) from the bright and dark regions. More recognizable changes

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1 are observed in the bright region, indicating that current mainly flows through the bright region.

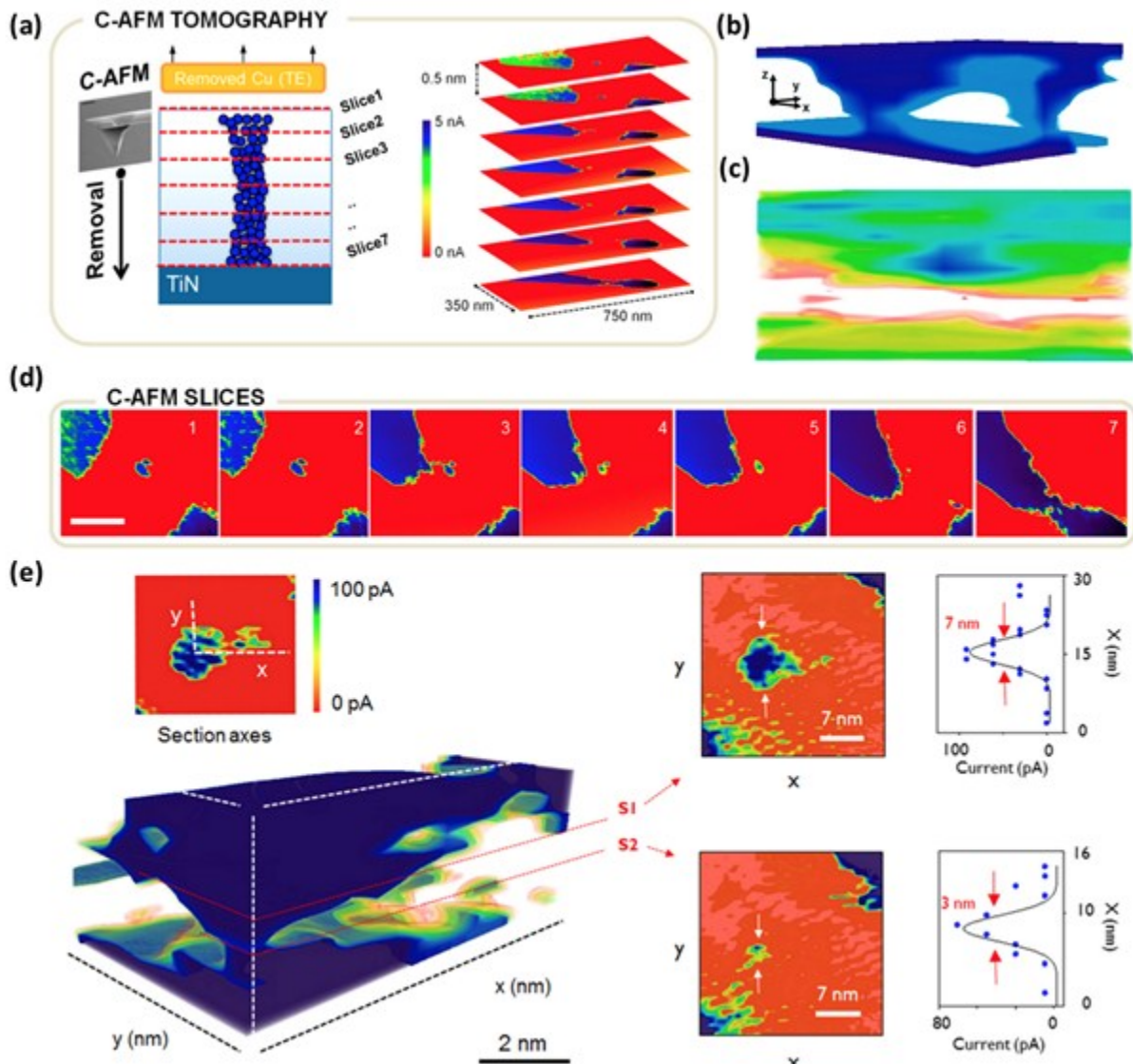
2 Reproduced with permission.^[52] Copyright 2016, John Wiley and Sons. e) Upper panel: $k^2\chi(k)$ spectra
 3 of HfO_2 (30 nm)/Hf (5,10, and 20 nm) samples with a reference of Hf powder. Lower panel: Simulated
 4 $k^2\chi(k)$ spectra for different crystalline structures including monoclinic, rutile, anatase, and tetragonal,
 5 revealing that the local structure of the HfO_2 main matrix is monoclinic-like. The $k^2\chi(k)$ difference
 6 (e.g., [HfO_2/Hf 20 nm] - [HfO_2/Hf 5 nm] and [HfO_2/Hf 5 nm] - [HfO_2]) is most similar to that of
 7 tetragonal structure instead of monoclinic, suggesting that the interfacial oxide has a different local
 8 structure from the matrix oxide. Reproduced with permission.^[69] Copyright 2017, American Chemical
 9 Society.



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1 **Figure 7.** a,c) XPEEM images of a SrTiO₃-based device in LRS (a) and HRS (c), showing a localized
2 filament with changes of the O concentration. Scale bars: 1 μm. b,d) O K-edge spectra obtained from
3 the filament region (green curve for LRS in (b) and blue curve for HRS in (d)) and the surrounding
4 region (black curves). e) Normalized A/B₂ ratio, which is inversely proportional to the Ti³⁺
5 concentration, in different device states, indicating the valence change of Ti ions induced by O
6 migration during the resistive switching process. Reproduced with permission.^[54] Copyright 2016,
7 Nature Publishing Group.

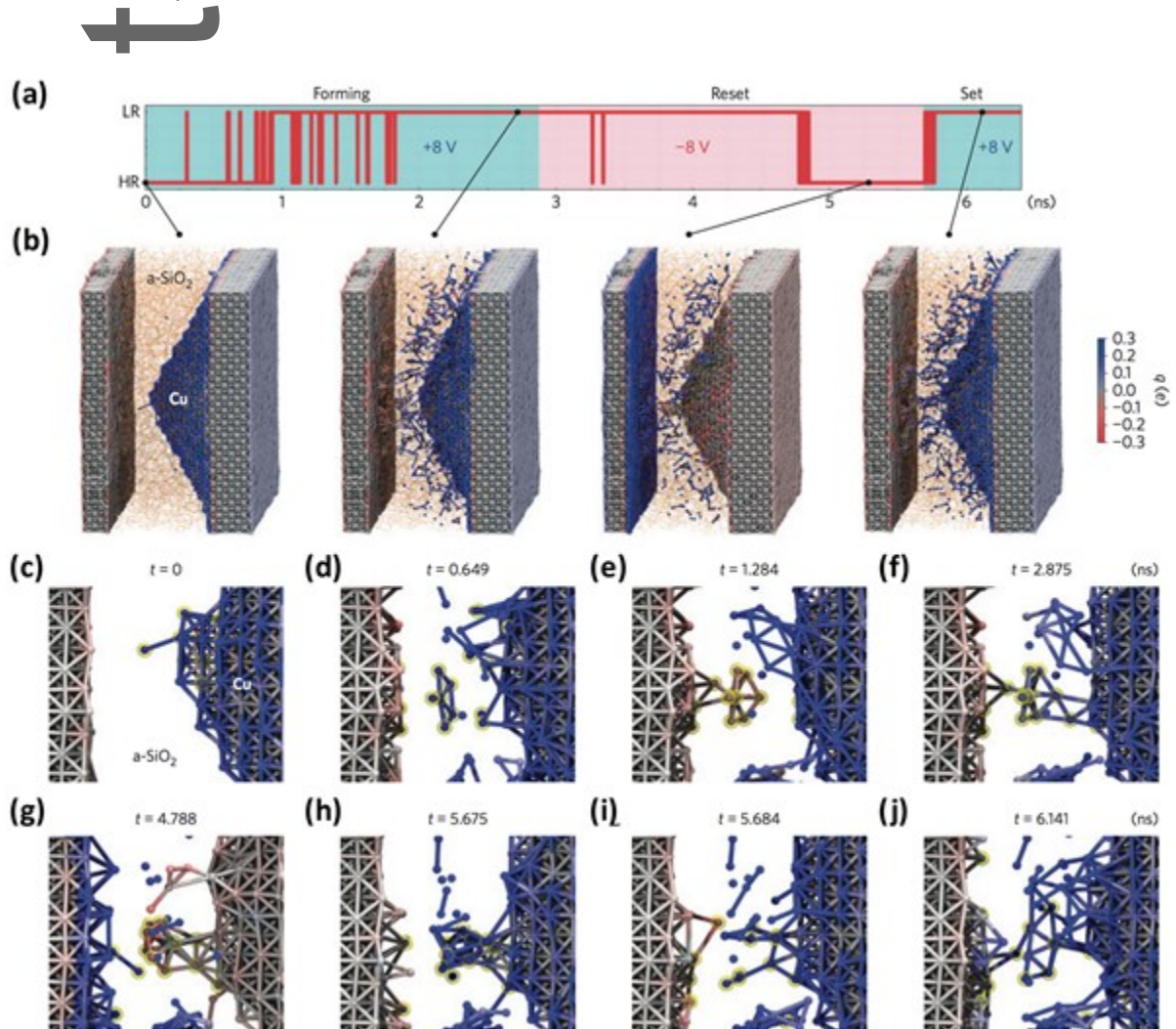


1

2 **Figure 8.** a) Schematic illustration of the 3D C-AFM tomography technique, showing the layer-by-
 3 layer scraping and the 2D C-AFM image slices collected at different heights. b,c) Reconstructed 3D C-
 4 AFM tomogram of a Cu/Al₂O₃-based ECM device, at LRS (b) and HRS (c). d) A set of 2D C-AFM slices
 5 used for the 3D reconstruction, indicating a CF having a conical shape with the narrow constriction
 6 close to the inert-electrode. Scale bar: 80 nm. Reproduced with permission.^[26] Copyright 2014,
 7 American Chemical Society. e) 3D C-AFM tomogram of a HfO₂-based VCM device, and C-AFM spectra

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1 indicating the sub-10 nm size of the CF. Reproduced with permission.^[56] Copyright 2015, American
 2 Chemical Society.



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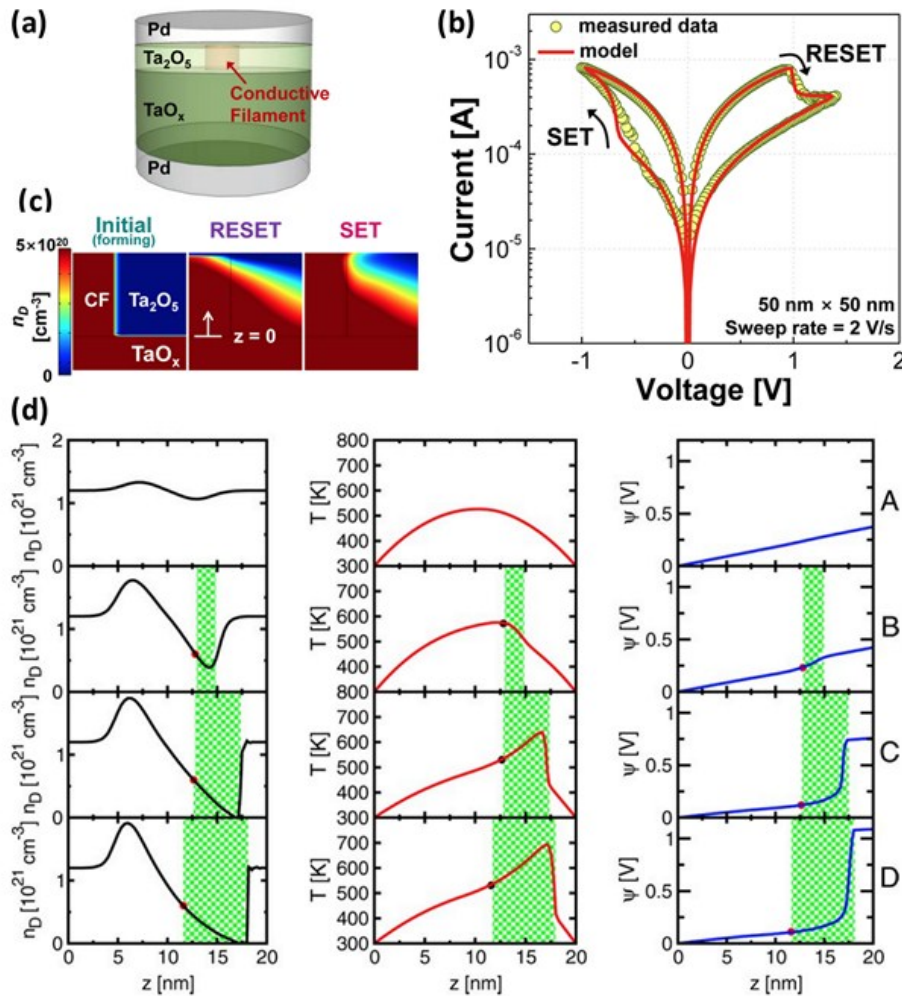
4 **Figure 9.** a) Evolution of the cell state for a Cu/SiO₂-based simulation cell between low resistivity (LR)
 5 and high resistivity (HR), under forming/set (cyan color) and reset (pink color) bias conditions. b)
 6 Snapshots of the atomic configuration at different moments in (a). c–j) Zoomed-in snapshots during
 7 the forming (c–f)/reset (g)/set (h–j) processes, showing the connection/disconnection of the two
 8 electrodes by the formation/rupture of a Cu-based CF upon voltage biases with opposite polarity,

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1 supporting experimental observations. Reproduced with permission.^[76] Copyright 2015, Nature

2 Publishing Group.

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4

5 **Figure 10.** a) Schematic of a simulation cell having a Ta₂O₅/TaO_x bilayer structure. b) *I*–*V*

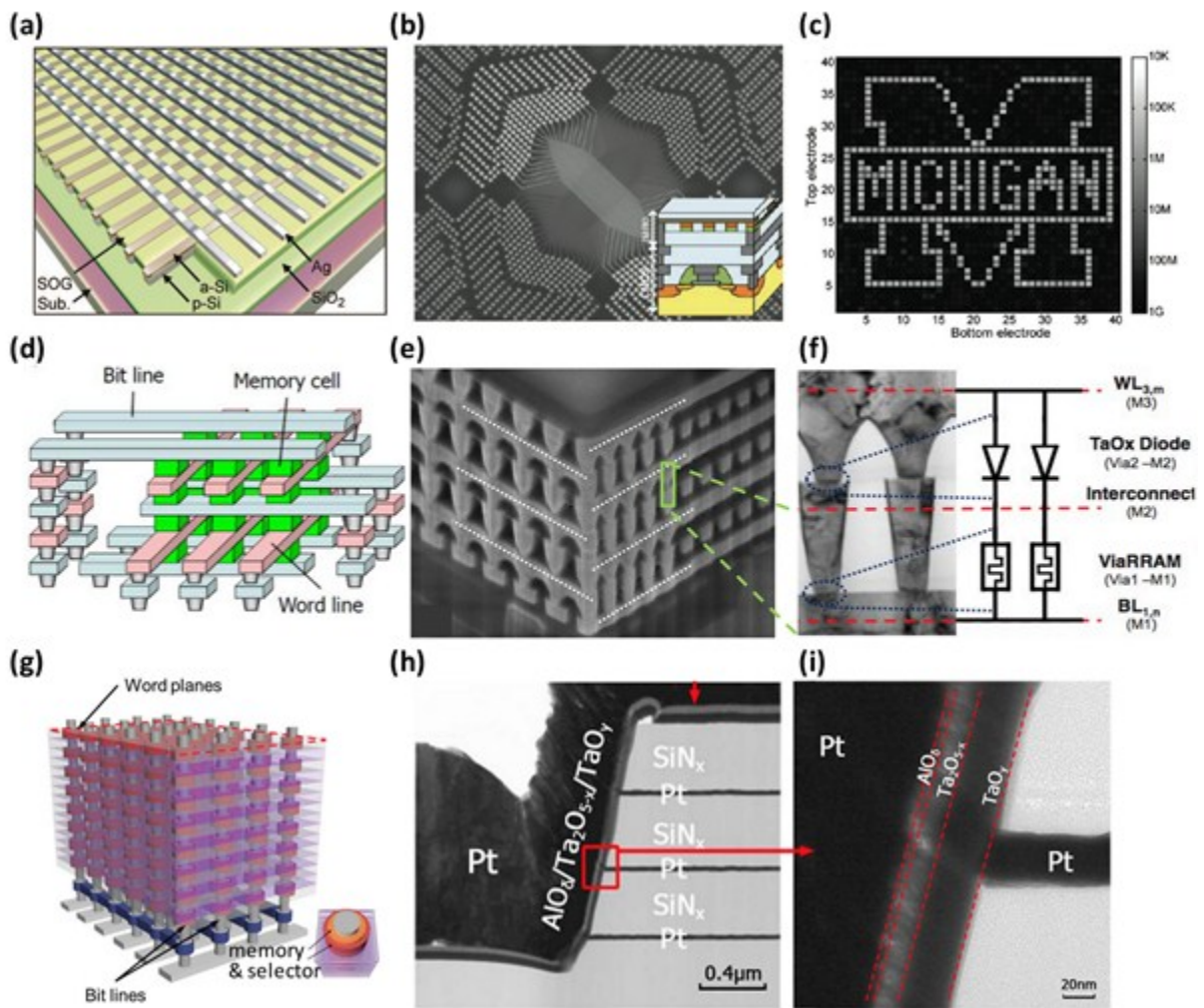
6 characteristic during RS, showing good agreement of the model (solid lines) with experimental

7 measurements (circles). c) Calculated 2D *V*₀-concentration profiles in the initial state, HRS and LRS.

8 Reproduced with permission.^[89] Copyright 2014, American Chemical Society. d) Calculated 1D

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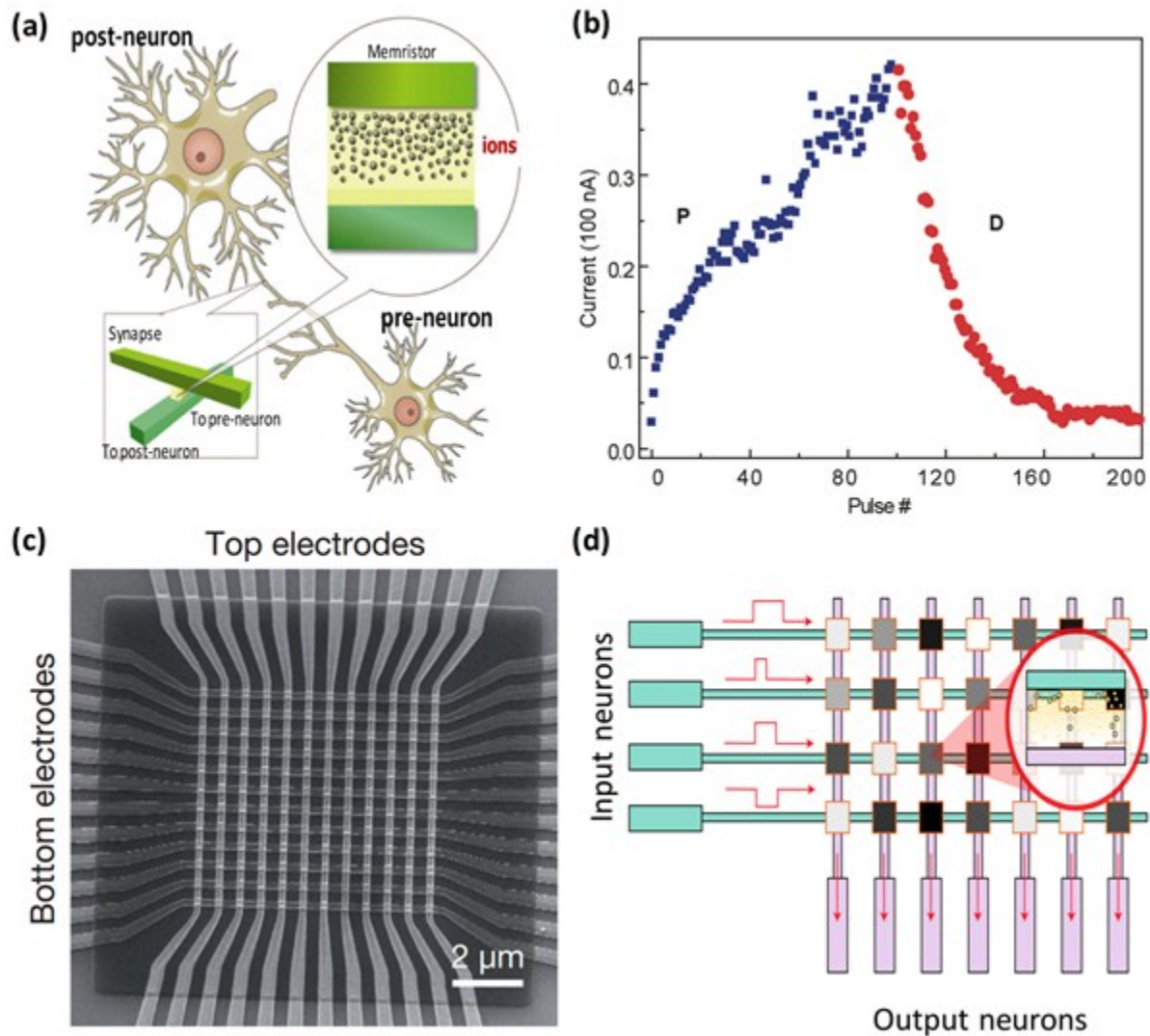
1 profiles of the V_O -concentration, temperature, and electrostatic potential at different states of
 2 another model VCM cell. Evolution of the depletion gap, defined as the region where the V_O -
 3 concentration is lower than $6 \times 10^{20} \text{ cm}^{-3}$, is indicated by the shaded area. Reproduced with
 4 permission.^[88] Copyright 2012, IEEE.



6
 7 **Figure 11.** a) Schematic illustration of a crossbar array in which a Ag/a-Si/Si RRAM cell is formed at
 8 each crosspoint. Reproduced with permission.^[21] Copyright 2009, American Chemical Society. b) SEM

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1 image showing a high-density crossbar array fabricated on top of a CMOS chip. Inset: schematic of
2 the vertical integration of the RRAM crossbar array with the CMOS circuitry. c) A 1600-pixel bitmap
3 image obtained by programming and retrieving data in the 40×40 crossbar array in (b). Reproduced
4 with permission.^[123] Copyright 2012, American Chemical Society. d) Schematic illustration of a 3D,
5 stacked cross-point structure. Reproduced with permission.^[145] Copyright 2011, IEEE. e) SEM image
6 and f) cross-sectional TEM image of a 3D stacked cross-point RRAM array based on the ViaRRAM
7 concept. A TaO_x diode is inserted in series with the RRAM element to reduce the sneak current.
8 Reproduced with permission.^[142] Copyright 2013, IEEE. g) Schematic illustration of the vertical RRAM
9 structure. Reproduced with permission.^[98] Copyright 2015, John Wiley and Sons. h) Cross-sectional
10 TEM image showing two RRAM cells formed along the sidewall of a Pt bitline. i) Zoomed-in TEM
11 image of a $\text{AlO}_\delta/\text{Ta}_2\text{O}_{5-x}/\text{TaO}_y$ sidewall device. Reproduced with permission.^[149] Copyright 2014,
12 Nature Publishing Group.

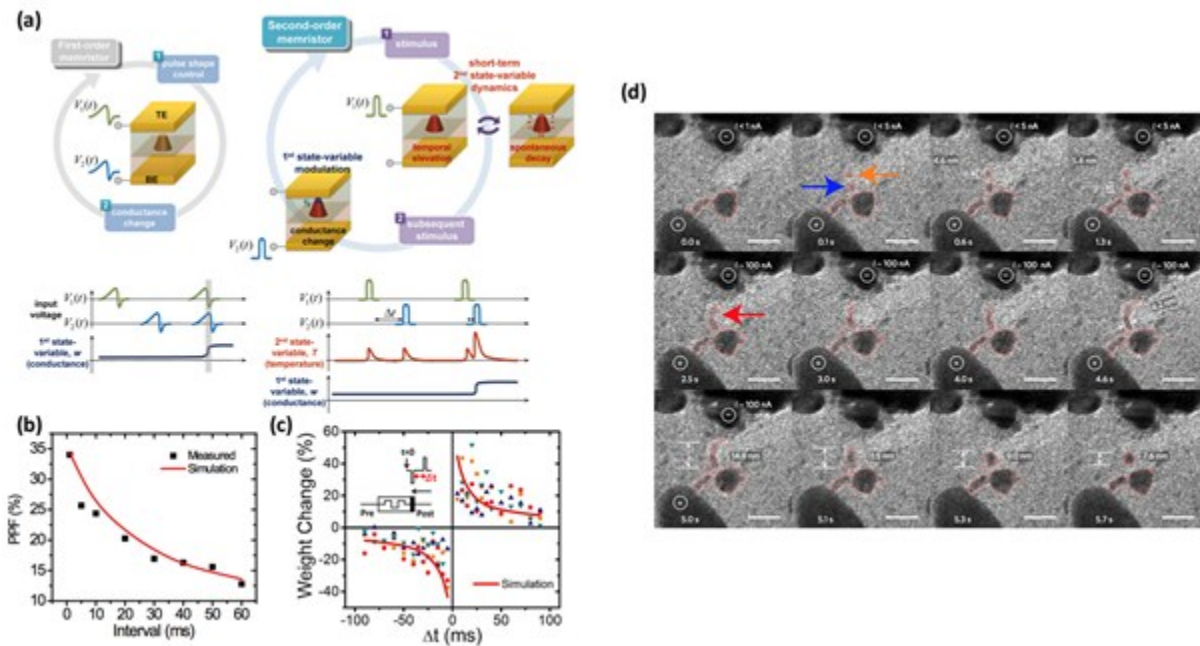


1

2 **Figure 12.** a) Schematic illustration of a memristor acting as a synapse bridging a pair of neurons,
 3 with the ability to gradually modulate the synaptic weight (conductance) by controlling the internal
 4 ionic configuration. b) Incremental conductance modulation of a memristor, in response to identical,
 5 consecutive potentiating (blue square) and depressing (red circle) pulses. Reproduced with
 6 permission.^[151] Copyright 2010, American Chemical Society. c) A fabricated 12×12 memristive
 7 crossbar array used for classification of input patterns. Reproduced with permission.^[174] Copyright

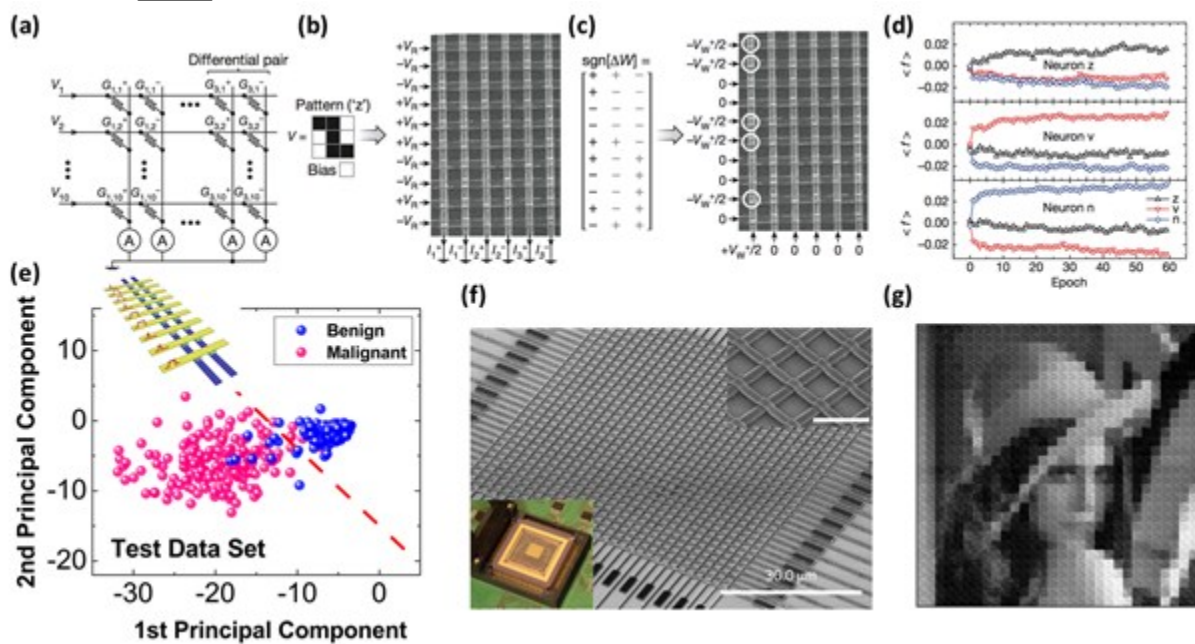
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1 2015, Nature Publishing Group. d) Schematic illustration of an artificial neural network based on a
 2 crossbar structure, in which memristors acting as synaptic devices at each crosspoint regulate signals
 3 propagating through them. Reproduced with permission.^[195] Copyright 2017, Nature Publishing
 4 Group.



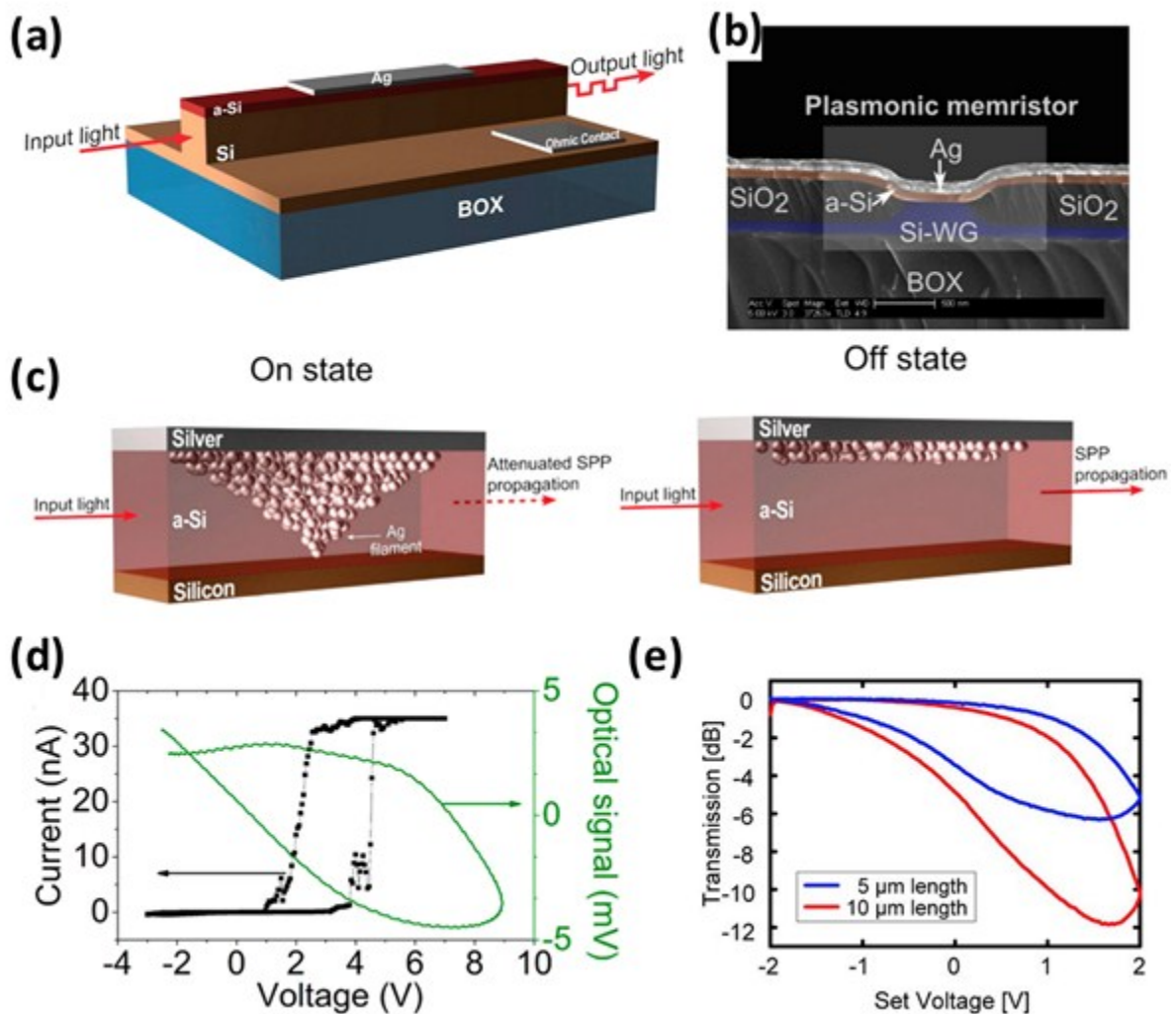
5
 6 **Figure 13.** a) Conceptual illustration of the operation of a second-order memristor. (left) Operation
 7 of a first-order memristor. The conductance is directly controlled by input voltage pulses. (right)
 8 Operation of a second-order memristor. The input pulses drive the short-term evolutions of a
 9 second state variable, which in turn facilitates the evolution of the first state variable and modifies
 10 the device conductance. The internal short-term dynamics allows the device to naturally implement
 11 rate- and timing-dependent synaptic effects with simple, nonoverlapping spikes. Reproduced with
 12 permission.^[97] Copyright 2015, American Chemical Society. b) A PPF results obtained from a WO_x -
 13 based memristor, showing interval- and frequency-dependent behavior analogue to biological

1 synapses. c) Weight change of the memristor device as a function of relative timing between the
 2 presynaptic and postsynaptic pulses, emulating STDP behavior with non-overlapping pulses.
 3 Reproduced with permission.^[182] Copyright 2015, John Wiley and Sons. d) In situ TEM images
 4 showing the field-driven filament formation ($t < 5.0$ s) and the diffusion-driven, spontaneous
 5 relaxation of the filament to nanoclusters ($t > 5.0$ s) in a diffusive memristor. Scale bars: 20 nm.
 6 Reproduced with permission.^[136] Copyright 2016, Nature Publishing Group.



7
 8 **Figure 14.** a–c) Illustrations of a neuromorphic network based on a 10×6 crossbar array, acting as a
 9 perceptron for pattern classification applications. d) Evolution of the output signals (f)
 10 experimentally obtained from the three output neurons in the perceptron, for different classes of
 11 inputs during training. Reproduced with permission.^[174] Copyright 2015, Nature Publishing Group. e)
 12 PCA analysis results obtained in a memristor array after unsupervised training, along with a decision
 13 boundary (red dotted line) obtained using logistic regression. Inset: Schematic of memristor-based

1 network operation. The input voltage signals with different pulse widths are applied to each row,
 2 and the output current signals are collected from each column. Reproduced with permission.^[190]
 3 Copyright 2017, American Chemical Society. f) SEM image of a fabricated 32×32 WO_x -based
 4 memristor crossbar array. Inset: magnified SEM image of the crossbar (upper right) and a wire-
 5 bonded memristor chip on the test board (bottom left).^[195] g) An experimentally reconstructed
 6 image based on a sparse-coding algorithm implemented using the crossbar in (f). Reproduced with
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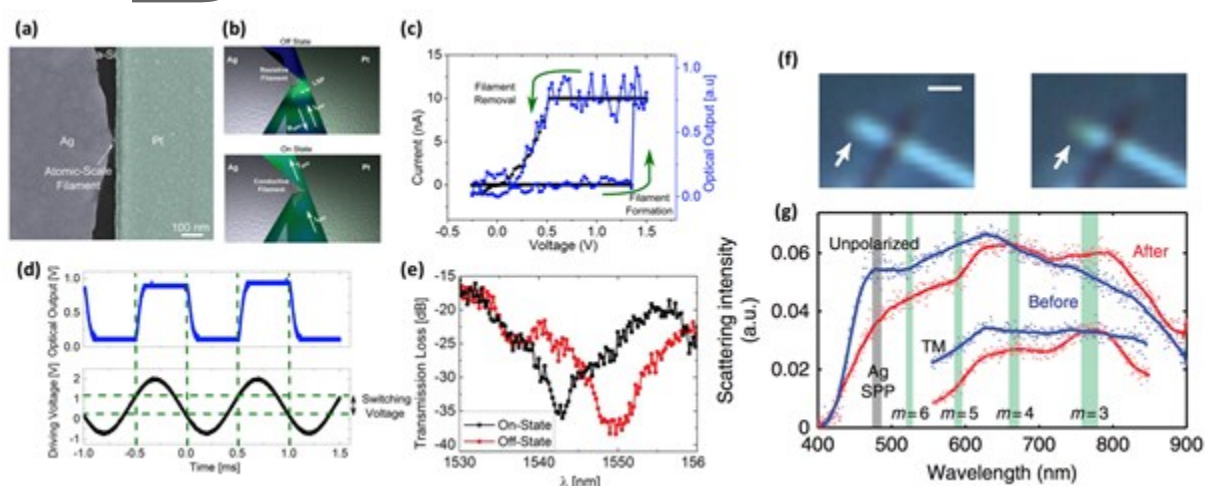


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A

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1 **Figure 15.** a) Schematic illustration of the plasmonic memristor structure with optical readout
 2 functionality. b) Cross-sectional SEM image showing a Ag/a-Si-based memristor integrated with a SOI
 3 waveguides. c) Schematic illustration of the plasmonic memristor. The formation (left)/annihilation
 4 (right) of a Ag-based conductive filament inside the a-Si electrolyte modulates the coupling of light
 5 with the plasmonic mode of the electrode, and thus light transmission through the waveguide. d)
 6 Correlation of the resistive switching I - V characteristics with the observed optical bistable behavior.
 7 Reproduced with permission.^[202] Copyright 2013, American Chemical Society. e) Latching optical
 8 switching behavior with an extinction ratio of 12 dB (6 dB) in another plasmonic memristor.
 9 Reproduced with permission.^[203] Copyright 2014, OSA publishing.

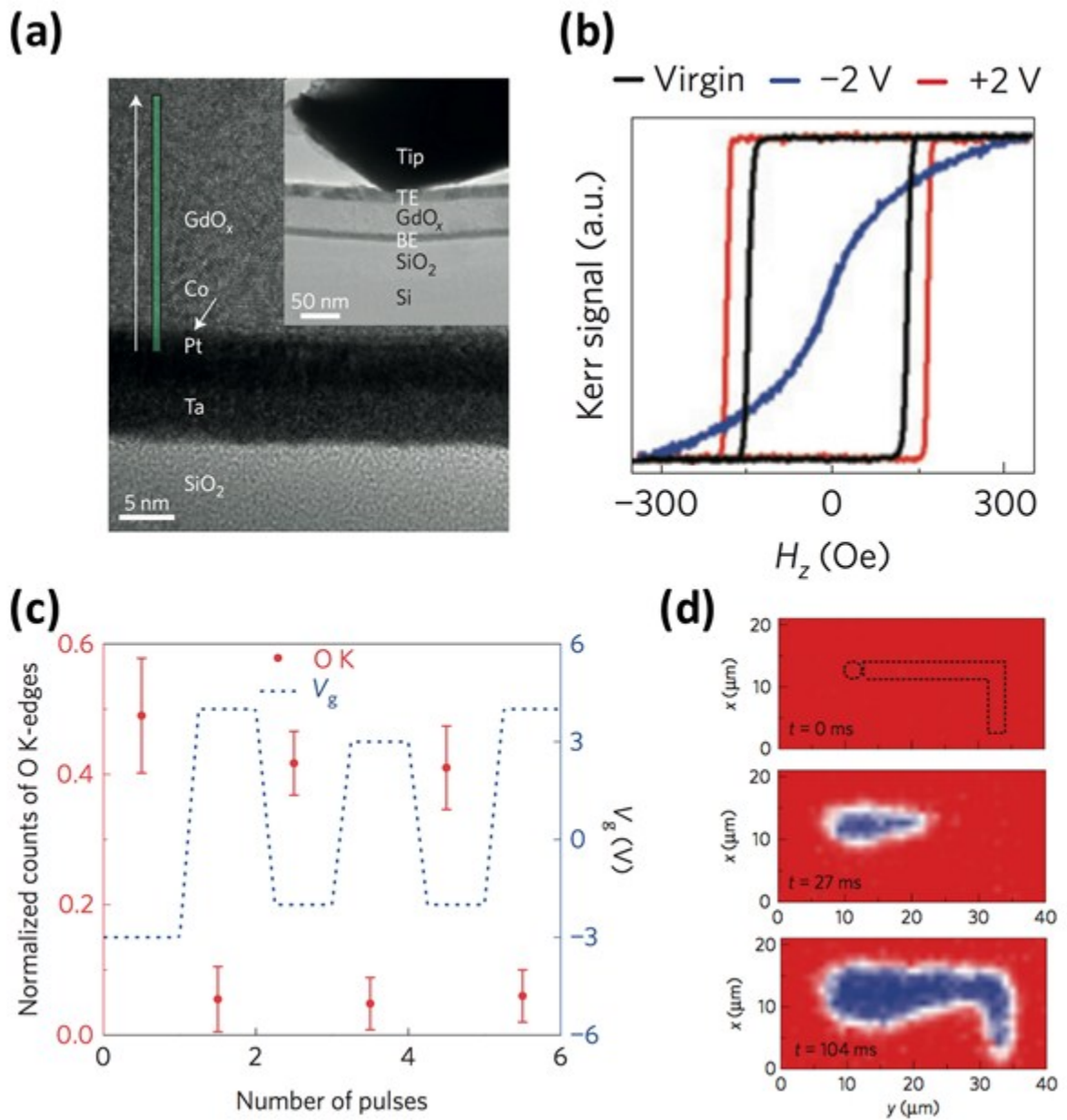


10

11 **Figure 16.** a) Top-view SEM image of a fabricated Ag/a-Si-based atomic scale plasmonic switch with
 12 the electrically formed Ag-filament. b) Schematic illustration of the operation of an atomic-scale
 13 plasmonic switch. The electrically controlled migration of a single or only a few atoms into a
 14 plasmonic cavity can short-circuit the two metallic pads and alter the plasmonic cavity resonance
 15 state. c) Resistive switching I - V characteristics (black circle) and the corresponding optical output

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1 transmission (blue square) that are simultaneously measured, indicating that the abrupt optical
2 transition is correlated with resistive switching based on atomic reconfiguration. d) Optical response
3 to a sinusoidal electrical input signal showing reversible, digital optical switching. e) Optical
4 transmission as a function of wavelength showing two different resonance states, with a blue-shift
5 of the resonance in the on state. Reproduced with permission.^[204] Copyright 2016, American
6 Chemical Society. f) Bright-field optical microscopy images of a Au/Al₂O₃/Ag-based memristive
7 optical antenna before (left) and after (right) electrical switching. A change in color from blue to
8 green is observed, suggesting that light scattering behaviors of the antenna are sensitive to even
9 small changes in the junction as light is effectively concentrated through the excitation of surface
10 plasmons. g) White-light-scattering spectra confocally obtained from the device before and after
11 electrical switching. Noticeable changes in the resonant scattering responses are observed in both
12 unpolarized and transverse magnetic (TM)-polarized spectra. Reproduced with permission.^[221]
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1

2 **Figure 17.** a) Cross-sectional TEM image of a device stack consisting of SiO₂/Ta/Pt/Co/GdO_x layers.

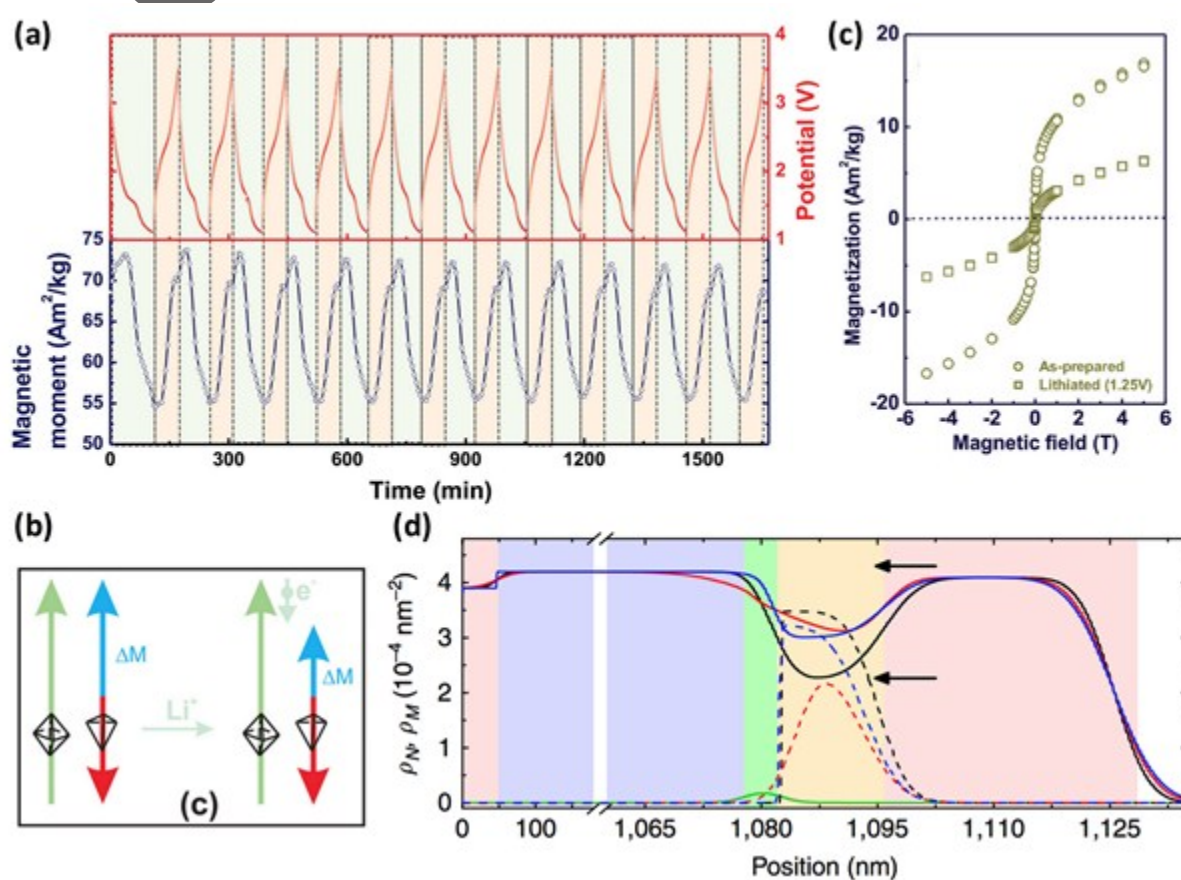
3 Inset: The in-situ TEM measurement setup where a Pt/Ir probe is contacted to a device with a Ta/Au

4 TE and Ta/Pt/Co BE. b) Polar magneto-optic Kerr effect (MOKE) hysteresis loops measured from the

5 virgin state and after the application of the negative and positive voltages at room temperature,

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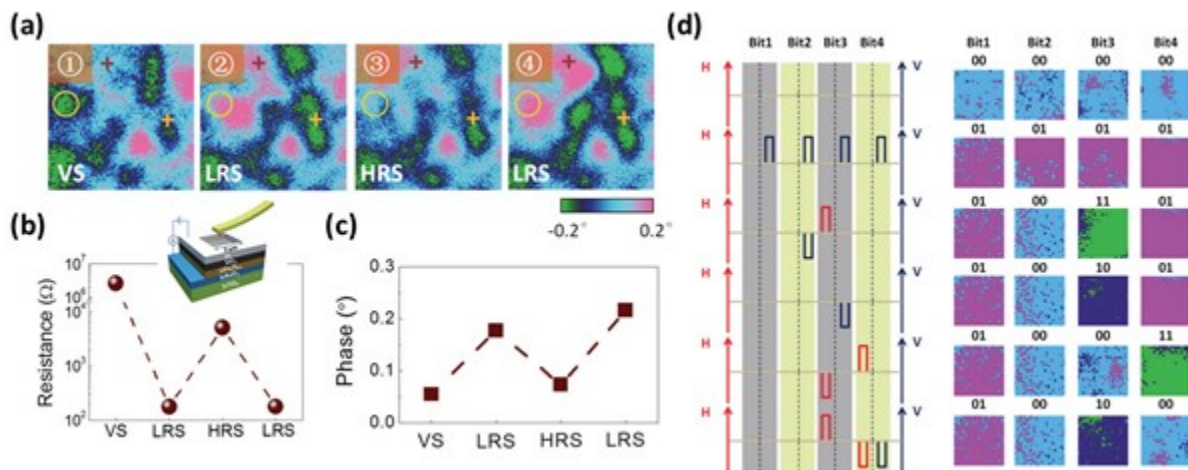
1 showing the complete, reversible removal and restoration of the hysteresis induced by electrical
 2 stimuli. c) Normalized O K-edge EELS intensity measured in the middle of the Co layer (red circles)
 3 under the application of positive and negative voltage cycles (blue dot line) showing the reversible
 4 modulation of oxygen concentration. d) Time-resolved polar MOKE maps showing the evolution of
 5 domain expansion following a laser-defined pattern, due to the selectively enhanced ionic transport
 6 by laser heating. Reproduced with permission.^[205] Copyright 2014, Nature Publishing Group.



7
 8
 9 **Figure 18.** a) Reversible modulation of magnetic moment in response to electrical stimuli, due to
 10 electrochemically driven delithiation/lithiation in spinel iron oxide ($\gamma\text{-Fe}_2\text{O}_3$) nanoparticles. b)

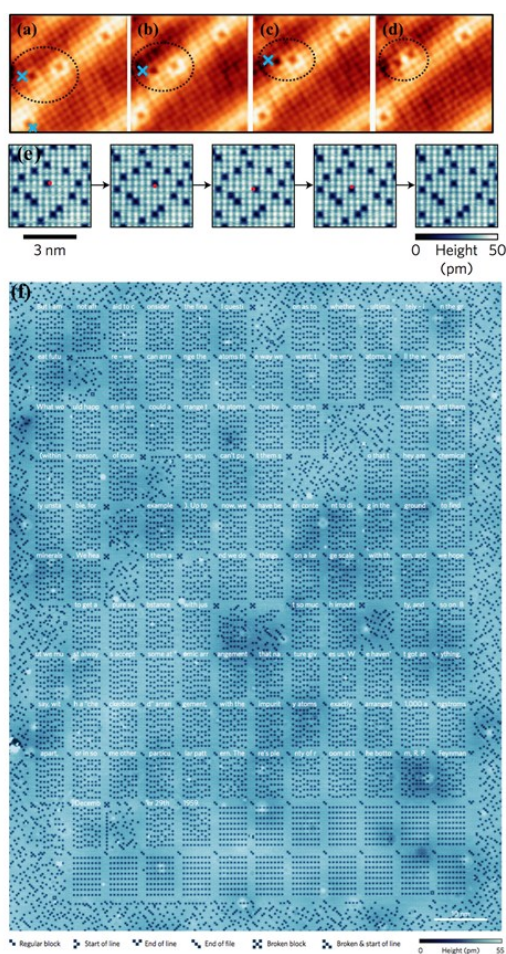
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1 Schematic illustration of the magnetization modulation induced by lithium intercalation that reduces
 2 the octahedral-site Fe ions from +3 to +2 state, thus lowering its magnetic moment (as represented
 3 the green arrow) and decreasing the total magnetization ΔM (blue arrows). Reproduced with
 4 permission.^[237] Copyright 2014, John Wiley and Sons. c) Changes in magnetization of as-prepared
 5 and lithiated ZnFe_2O_4 samples, showing 70% loss of the saturation magnetic moment. Reproduced
 6 with permission.^[239] Copyright 2016, John Wiley and Sons. d) Polarized neutron reflectometry (PNR)
 7 depth profiles for as-grown (black), +40 V conditioned (red), and ± 40 V conditioned (blue) samples.
 8 The solid and dashed lines represent nuclear scattering length density (SLD), ρ_N , and magnetic SLD,
 9 ρ_M , respectively, and background colors represent Pd (red), AlO_x (blue), GdO_x (green) and Co (yellow)
 10 layers, respectively. The significant increase of the ρ_N curve inside the entire Co layer indicates that
 11 the applied electric field drives oxygen ions deep into the Co film, leading to decrease of the
 12 magnetization of the Co film. Reproduced with permission.^[240] Copyright 2016, Nature Publishing
 13 Group.



14
 15 **Figure 19.** a) MFM images of a LiFe_5O_8 -based device at different states during resistive switching

1 cycles shown in (b). The areas with significant changes of magnetization are highlighted. c) MFM
2 phase data collected during electrical cycling in (b). The data are measured at the location marked by
3 the brown cross in the MFM maps in (a), showing reversible modulation of magnetization that is
4 correlated with the resistive switching behavior. d) Demonstration of multilevel data storage based
5 on the coupled magnetization and resistance switching effects. (left) Magnetic field (red) and electric
6 pulse (blue) programming sequences for each domain, in which two bits are stored using the
7 direction and the intensity of the domain's magnetization, respectively. (right) Read-out of the
8 stored data from MFM measurements, showing four distinct states. Reproduced with permission.^[206]
9 Copyright 2016, John Wiley and Sons.



1

2 **Figure 20.** a–d) STM topography images of Cl vacancies in a NaCl/Au surface demonstrating atomic-
 3 level manipulation of the material configuration. Individual Cl vacancies can be laterally moved by
 4 the application of voltage with a STM tip. Reproduced with permission.^[254] Copyright 2015, American
 5 Chemical Society. e) STM topography images of Cl vacancies (represented by dark blue squares) in a
 6 Cl/Cu surface showing that a Cl vacancy can be repositioned in all four directions by a STM tip, where
 7 the position of the STM tip is marked by a red dot. f) STM image of a kilobyte atomic memory
 8 through atom-by-atom manipulation of Cl vacancies, offering a high areal density of 502 terabits per
 9 square inch. Reproduced with permission.^[255] Copyright 2016, Nature Publishing Group.

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1

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5 nanoionics for memory and neuromorphic computing applications.

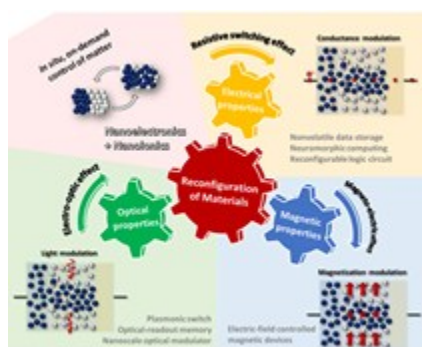


6

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9 research fellow at Harvard University, Cambridge, MA. In 2005, he joined the faculty of the Electrical
10 Engineering and Computer Science Department at the University of Michigan and is currently a
11 Professor. He is the Director of the Lurie Nanofabrication Facility and Associate Editor for Nanoscale.
12 His research interest includes resistive-random access memory, memristor-based logic circuits,

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1 neuromorphic computing systems, aggressively scaled transistor devices, and electrical transport in
 2 low-dimensional systems.



4
 5 **By controlling the internal ion distribution in a solid-state film**, the material's chemical composition
 6 and physical (i.e. electrical, optical and magnetic) properties can be reversibly reconfigured, in-situ,
 7 using an applied electric field. The reconfigurability has been achieved in a wide range of materials,
 8 and can lead to the development of new memory, logic, and multifunctional devices and systems.

10 **Reconfigurable Materials**

12 J. Lee, W. D. Lu*

14 **On-Demand Reconfiguration of Nanomaterials—When Electronics Meets Ionics**

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