

The Influence of Chemical Modification on Linker Rotational Dynamics in Metal–Organic Frameworks

Joshua T. Damron⁺, Jialiu Ma⁺, Ricardo Kurz, Kay Saalwächter, Adam J. Matzger,^{*} and Ayyalusamy Ramamoorthy^{*}

Abstract: The robust synthetic flexibility of metal–organic frameworks (MOFs) offers a promising class of tailorable materials, for which the ability to tune specific physicochemical properties is highly desired. This is achievable only through a thorough description of the consequences for chemical manipulations both in structure and dynamics. Magic angle spinning solid-state NMR spectroscopy offers many modalities in this pursuit, particularly for dynamic studies. Herein, we employ a separated-local-field NMR approach to show how specific intraframework chemical modifications to MOF UiO-66 heavily modulate the dynamic evolution of the organic ring moiety over several orders of magnitude.

Metal–organic frameworks (MOFs) are constructed from inorganic clusters bridged by organic linkers and possess ideal architectures for gas storage and mixture separation as well as generating recent interest as nanodevices and molecular machines.^[1] In comparison to traditional inorganic porous materials, MOFs are far more dynamic/flexible due to the incorporation of organic ligands.^[2–4] A key advantage of MOFs is the synthetic ease of chemical modification to the organic moieties without change of framework structure, known as isorecticular synthesis. This allows for many avenues of functional modification including adsorptive selectivity based on favorable ligand–guest interactions,^[5] changes in pore aperture for separations,^[6,7] and shifts in optical behavior.^[8] An additional feature of MOFs, which has captured much attention, is the prevalence of framework dynamics which can lead to dramatic structural shifts such as breathing, swelling, and subnetwork displacements.^[9,10] While deformations in the metal cluster have been implicative of dynamic

events, the nature of the linker and its functionality is a major contributor to framework dynamics,^[10] and play a significant role in determining macroscopic functions such as gas storage and separation, ferroelectricity, spin crossover and luminescence in a variety of MOFs.^[11–17] As such, establishing relationships between local dynamics driven by the organic ligands and their chemical modification is of general interest for materials design in MOFs.^[17,18]

Though the structural picture of MOFs is most informed by analysis with X-ray diffraction, this yields an essentially static picture, failing to capture the dynamic aspects of the framework. However, dynamic measures of crystalline materials are challenging due to the potential of multiple time-scales of coordinated motions in the solid-state. Many dynamic studies have employed solid-state NMR spectroscopy (ssNMR) through ²H-lineshape analysis to characterize the rotational motion of simple *p*-phenylene rings within the MOF framework.^[15,19–24] Separated-local-field (SLF) solid-state NMR spectroscopy is an alternative, attractive means of dynamic characterization as it uses the heteronuclear dipolar coupling as a proxy for molecular dynamics, does not require isotopic labeling, and can be used to characterize motions over a broad dynamic range with atomic resolution.^[25]

Herein, a combination of computation and the DIP-SHIFT^[26,27] (dipolar chemical SHIFT correlation) SLF methodology is employed to characterize ligand dynamics in MOFs; details on the preparation, characterization and calculations are given in the Supporting Information (Figures S1–S6 and Tables S1, S2). For typical heteronuclear dipole–dipole couplings such as ¹³C–¹H and ¹⁵N–¹H, the experiment is sensitive to a dynamic range spanning 0.1 μs–10 ms (see Figure 1).^[28] Though a number of MOFs are suitable for this approach, we focus on UiO-66^[29] and its derivatives (see Figure 2) due to its wide range of potential applications as well as its high thermal and water stability. We show that the dynamics of the ligands in the framework are heavily impacted by ligand functionalization; specifically, that the variable dynamic behavior arises from differences in interactions with the local chemical environment and is highly responsive to temperature.

Terephthalic acid is the most commonly used linker for the construction of MOFs. The phenylene units are symmetric and comprised of four equivalent aromatic ¹³C–¹H bonds making it a suitable candidate for DIPSHIFT measurements. Thus, UiO-66 (Zr₆(μ₃-O)₄(μ₃-OH)₄(terephthalate)₆) (see Figure 2a) was investigated first to understand the dynamic behavior of a simple phenylene unit. Using the ¹³C–¹H bond on the phenylene ring as the local field pair, DIPSHIFT measurements were performed on UiO-66 at 23 °C (Fig-

[*] J. T. Damron,^[†] J. Ma,^[†] Dr. A. J. Matzger, Dr. A. Ramamoorthy
Department of Chemistry, University of Michigan
930 N. University Ave., Ann Arbor, MI 48109-1055 (USA)
E-mail: matzger@umich.edu
ramamoor@umich.edu

R. Kurz, Dr. K. Saalwächter
Institut für Physik-NMR, Halle/Saale (Germany)
Dr. A. J. Matzger
Macromolecular Science and Engineering
University of Michigan
2300 Hayward Avenue, Ann Arbor, MI 48109-1055 (USA)
Dr. A. Ramamoorthy
Biophysics Program, University of Michigan
930 N. University Ave., Ann Arbor, MI 48109-1055 (USA)

[†] These authors contributed equally to this work.

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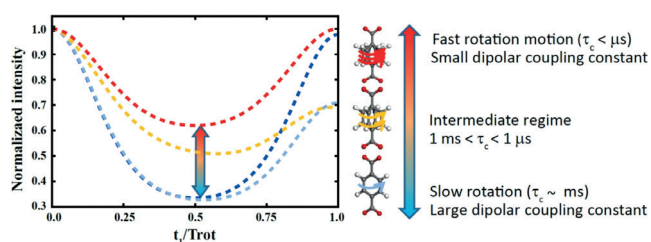


Figure 1. DIPSHIFT curves as a function of dynamic rate on NMR timescale. The lines are model fits assuming diffusive anisotropic motions under the Anderson-Weiss approximation^[28, 30] from UiO-66 data at various temperatures (see the Supporting Information). The curves correspond to estimated correlation times (τ_c) of ms or the rigid limit (dark blue), approximately 350 μs (light blue), approximately 5 μs (yellow), and the fast limit below one microsecond (red). In the present study, DIPSHIFT curves were measured using the pulse sequence shown in Figure S7 and analyzed as explained in the Supporting Information. t_i is the duration of evolution of nuclear spin magnetization under the effect of a combination of ^{13}C - ^1H dipolar coupling and ^{13}C chemical shift (see Figures S7 and S8) and T_{rot} is the sample rotation period.

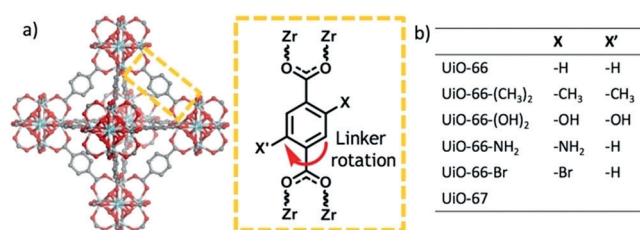


Figure 2. Structure of MOFs. a) UiO-66 structure and Scheme of linker rotation. b) A list of UiO type structures investigated in this study.

ure 3 a) resulting in a heteronuclear dipolar coupling constant (D_{CH}) of $16.4 \pm 0.2 \text{ kHz}$. The exhibited DIPSHIFT curves feature well-depths that are roughly proportional to the strength of the C–H dipolar coupling and can be fit with standard MAS equations to extract the coupling constants (see Figures 3, and S9–S15).^[31, 32] The experimentally determined D_{CH} value is smaller than the theoretical value for a rigid ^{13}C - ^1H bond (ca. 22 kHz), which is due to dynamic averaging resulting from the motional modes, namely ring rotations and librations, available to *p*-phenylene within the framework.^[38] Given the timescales detected by this approach, the primary event detected is associated with the ring flip although smaller, more rigorous librational modes are also expected to contribute to a lesser degree. By adopting an appropriate model, a correlation time for the motional ring flip can be estimated, which is found to be approximately a few milliseconds falling in “slow” regime dynamics for UiO-66 at 23 °C (see the Supporting Information) and is confirmed by higher temperature experiments. This result aligns with reports on MIL-53^[19] and MOF-5^[33] both containing *p*-phenylene units in a sterically unhindered environment. It is notable that the DIPSHIFT data is symmetric around half the rotor period for UiO-66 at 23 °C, which is a model-free confirmation of slower motional time scales assuming the dynamics are not already in the fast limit. Asymmetries arising from intensity decay due to transverse relaxation

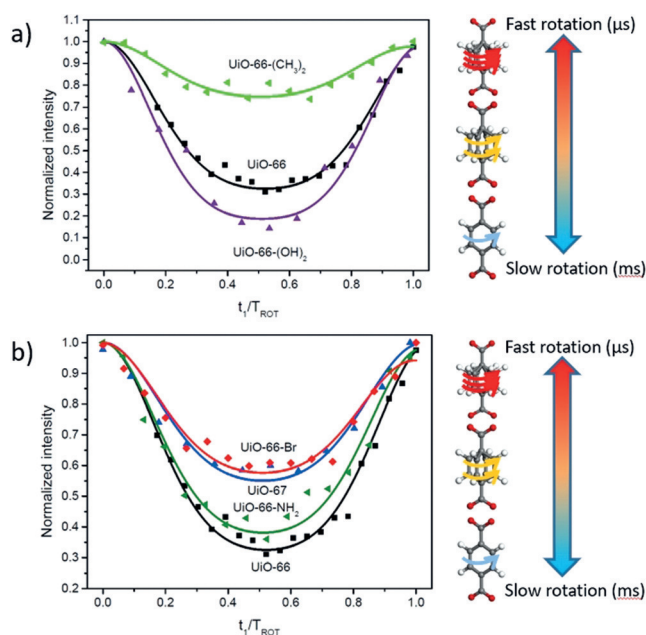


Figure 3. Effect of substituents on the mobility of phenylene group measured from DIPSHIFT experiments at 23 °C. a) UiO-66, UiO-66-(OH)₂ and UiO-66-(CH₃)₂; b) UiO-66, UiO-66-NH₂, UiO-66-Br and UiO-67. The symbols are the experimental data points, while the curves are simulated from the fits of the dipolar coupling constants (see details in Section S3; Figures S9–S15).

effects during the encoding time occur when the molecular motion is accelerated (e.g. temperature induced) to time-scales coinciding with the inverse of the coupling constant, known as the “intermediate motional regime,” as depicted in Figure 1.^[28]

After establishing the effectiveness of DIPSHIFT based characterization in the UiO-66 sample, we next turned our investigation to ligand-functionalized UiO-66 MOFs. Methyl and hydroxy groups are among the most commonly used functionalities and are particularly relevant for applications such as separations and catalysis.^[34, 35] The bulkiness of the methyl group and the hydrogen bonding potential of the hydroxy groups make them a diametric pair to study their effects on rotational dynamics. Indeed, experiments on UiO-66-(OH)₂ (Zr₆(μ_3 -O)₄(μ_3 -OH)₄(2,5-dihydroxyterephthalate)₆) and UiO-66-(CH₃)₂ (Zr₆(μ_3 -O)₄(μ_3 -OH)₄(2,5-dimethylterephthalate)₆) resulted in dramatically different DIPSHIFT curves from one another (Figure 3). UiO-66-(OH)₂ gives a higher D_{CH} value (ca. $19.0 \pm 0.5 \text{ kHz}$) than the parent UiO-66 ($D_{\text{CH}} \approx 16 \text{ kHz}$) indicating reduced mobility; on the other hand, the measured D_{CH} value for UiO-66-(CH₃)₂ is only 10 kHz. This stark reduction of the C–H dipolar coupling constant for UiO-66-(CH₃)₂ evidences much faster rotary motion than UiO-66, placing it well into the fast regime ($\tau_c < 1 \mu\text{s}$). These simple substitutions cause dramatic changes to the ligand framework interaction which induces approximately four orders of magnitude difference in the dynamic rotation. These same trends hold at different temperatures: UiO-66-(OH)₂, the least dynamic of the three MOFs, must be heated to approximately 105 °C (60 °C higher than UiO-66) to achieve rotational times on the order of 100 μs (see the Supporting

Information). In contrast, the 10 kHz value for the methyl substituent is similar to the plateau value observed for UiO-66 in the fast regime ($\tau_c < \mu\text{s}$) at 155°C (ca. 11 kHz). Indeed, when UiO-66-(CH₃)₂ is cooled to 6°C, an asymmetric curve is detected (see Figure S12), suggesting that it is in a comparable dynamic regime to UiO-66 at temperatures approximately 50°C lower.

To complement the experimental results, DFT simulations on a model complex were used as a tool to understand the details of rotational dynamics. In the case of UiO-66, a model of Zr₆O₄(OH)₄(COOH)₁₁(COOC₆H₅) bearing a single phenyl rotor was constructed and its geometry was optimized. A relaxed potential energy surface scan was performed with the phenyl ring rotation through its principle axis. A rotational energy barrier of 43.5 kJ mol⁻¹ was obtained. Despite the simplicity of the model, the obtained activation energy is similar to the one found for MOF-5^[36] and MIL-53.^[19] As with UiO-66, Zr-oxo complexes were constructed and calculations were also performed on models with hydroxy and methyl functional groups. A high activation energy of 75.3 kJ mol⁻¹ was found for the model Zr₆O₄(OH)₄(COOH)₁₁(COOC₆H₅-*o*-OH). This is expected to dramatically slow down the ring dynamics which supports the experimental data. A closer inspection of the model shows that the hydroxy group hydrogen bonds with the carboxylate group, an interaction that must be broken upon rotation (Figure 4). The hydrogen bonding explains the high energy barrier and suggests that it is the dominant factor leading to the experimentally observed high C–H dipolar coupling constant and slow rotational dynamics in UiO-66-(OH)₂. On the other hand, the model for UiO-66-(CH₃)₂ (Zr₆O₄(OH)₄(COOH)₁₁(COOC₆H₅-*o*-CH₃)) shows a lower activation energy of 33.5 kJ mol⁻¹. This reduced energy barrier is attributed to repulsive interactions between the bulky methyl group and the cluster. The simulated energy barrier results are fully consistent with the experimental

observations for rotational frequencies: UiO-66-(*o*-OH)₂ ≪ UiO-66 ≪ UiO-66-(*o*-CH₃)₂. The model above, however, does not account for geometric restrictions and electronic effects imposed on the rotors from Zr₆ clusters at both ends. To simulate these effects, two spatially fixed clusters held at crystallographic distances and bridged by a linker were constructed; the relaxed potential energy scan was performed with a semi-empirical method (Section S2). The energy barrier was then found using DFT single point energy calculations on the highest and lowest energy conformations determined from the relaxed torsional scan. Activation energies with this approach resulted in 65.77 kJ mol⁻¹, 23.39 kJ mol⁻¹ and 91.71 kJ mol⁻¹ for models of UiO-66, UiO-66-(CH₃)₂ and UiO-66-(OH)₂, respectively. These results are consistent with the results of the simpler model and further support the influence of local steric and electronic effects in determining linker dynamics.

UiO-66-Br (Zr₆(μ₃-O)₄(μ₃-OH)₄(2-bromoterephthalate)₆) and UiO-66-NH₂ (Zr₆(μ₃-O)₄(μ₃-OH)₄(2-aminoterephthalate)₆) were also synthesized to test the influence of linker dynamics from another bulky substituent other than a methyl group and a potential hydrogen bond donor. The measured D_{CH} for UiO-66-Br is closer to UiO-66 at 15 kHz but exhibits a slight asymmetry in the curve (Figure 3b).^[39] Despite the higher value than UiO-66-(CH₃)₂, the temperature behavior seen in Figure S13 of UiO-66-Br also shows that it is in the fast regime limit at room temperature. This suggests that the bromine atom also induces a steric effect which lowers the rotational energy barrier similar to the function brought by methyl groups, but is not as extreme given the higher fast limit value. UiO-66-NH₂ shows a D_{CH} value of 15.3 ± 0.3 kHz, value close to UiO-66. However, the temperature behavior is quite constant in a similar temperature range (see Figure S14). This behavior could be interpreted as a combination of effects: a balance between a steric effect and the hydrogen bonding potential of the NH₂ group.

Changing the length of the linker while maintaining the topology, known as isorecticular synthesis, is an important strategy that has been used to increase the surface area of MOFs without altering the topology.^[37] Thus UiO-67 (Zr₆(μ₃-O)₄(μ₃-OH)₄(biphenyl-4,4'-dicarboxylate)₆) was examined for its larger pore window (ca. 9–11 Å) compared to the parent structure UiO-66 (ca. 5–7 Å). The measured D_{CH} for UiO-67 is 12.4 ± 0.3 kHz, which is smaller than that obtained from UiO-66. The lowering of the rotational energy barrier is attributed to the junction of the biphenyl ring, which has a much lower frictional constraint versus the Zr cluster interaction. At higher temperature, the depth of the curve decreases substantially, but no intermediate dynamics are detected (Figure S15). This is possible in the fast limit regime when the reorientational angle of motion increases; it is plausible that the increased length and biphenyl junction could contribute additional bending modes to facilitate such behavior.

In summary, DIPSHIFT solid-state NMR experiments performed on a series of Zr MOFs with various structural modifications showed different ligand dynamics. The presence of functional groups such as methyl or hydroxy groups can greatly speed or slow down the rotational motion of the

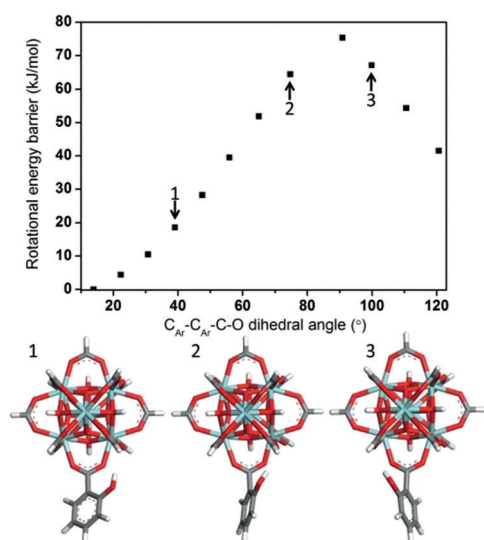


Figure 4. DFT simulation of rotational dynamics in MOFs. Rotational energy barrier calculation results and snap shot of the model complexes incorporating a single cluster with C_{Ar}-C_{Ar}-C-O at 40°, 75° and 100°. Coordinates and energies for all models used in the calculations are given in the Supporting Information.

linker in the Zr MOFs. Insights from DFT simulation reveals that such dynamics changes arise from the local chemical interactions such as hydrogen bonding or steric repulsion. Isoreticular structure expansion and factors such as temperature also influence the ligand dynamics. The insights gained are relevant to applications, where dynamics of the ligands are heavily involved, such as gas separations and kinetically limited processes such as molecular sieving.

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Conflict of interest

The authors declare no conflict of interest.

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- [38] Motions occurring at much faster rates than the coupling constants will result in a “pre-averaged” tensor producing a reduced coupling constant. The motions responsible for the pre-averaging are on distinctly faster time-scales than the motions detected in the experiment.
- [39] This value is larger than UiO-67, despite similar curve depths in the plot because UiO-66-Br measurements were completed at higher spinning frequencies (see the Supporting Information).

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