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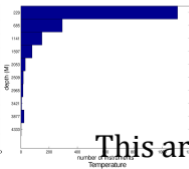
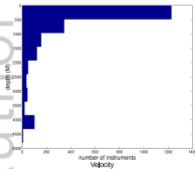
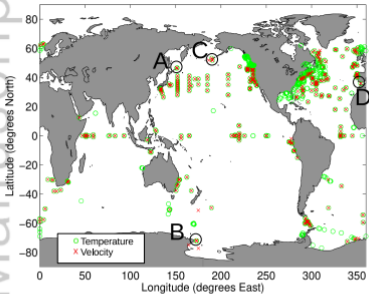
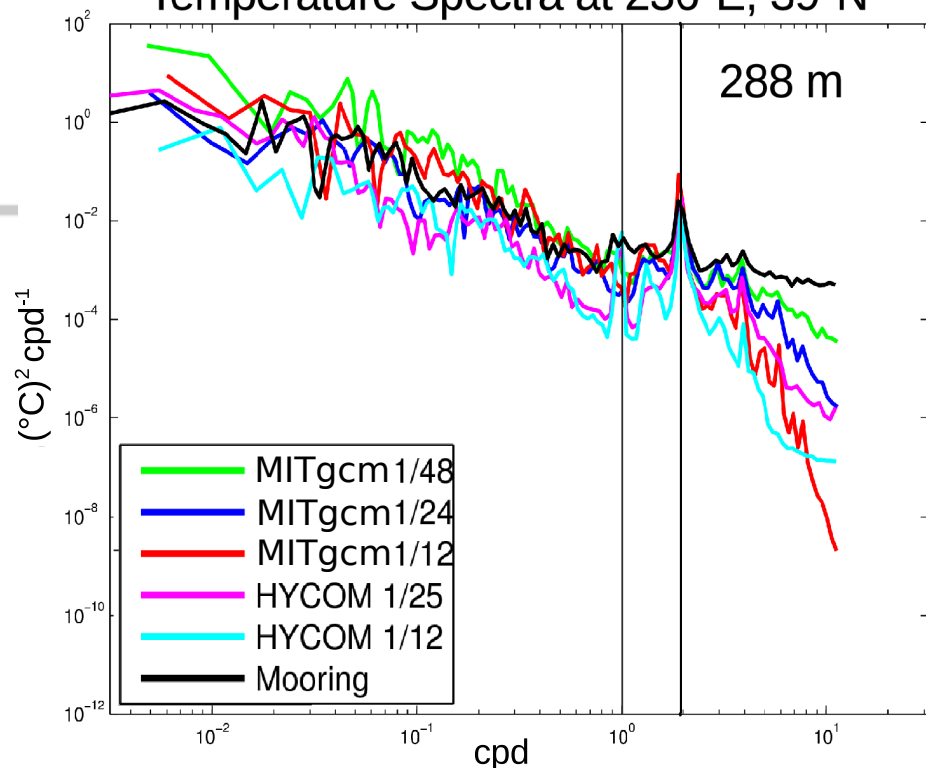


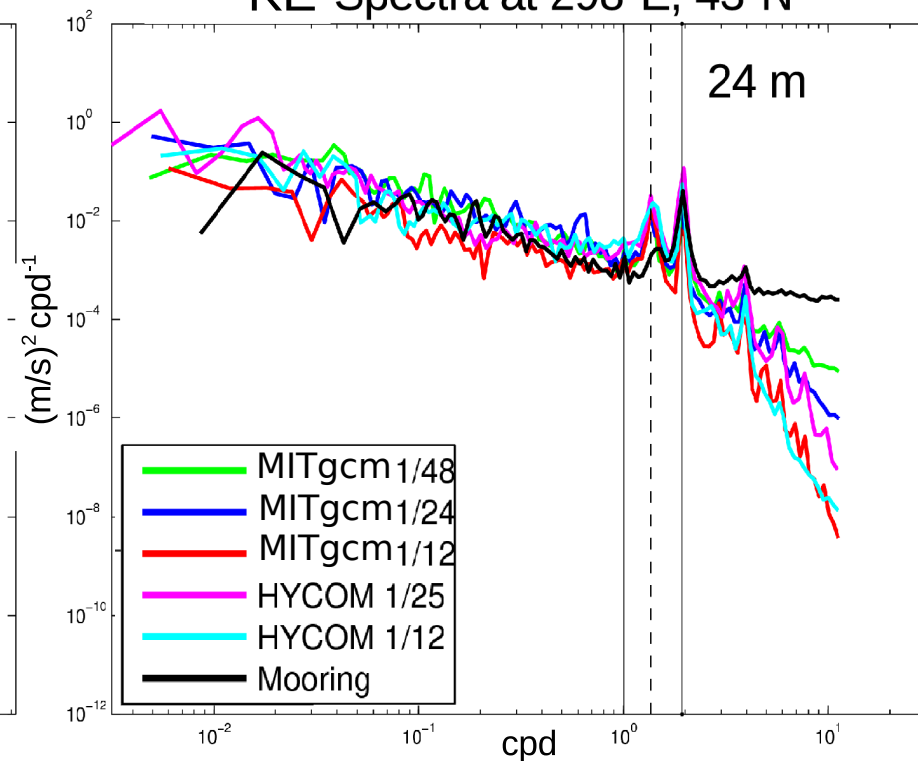
Figure 2.

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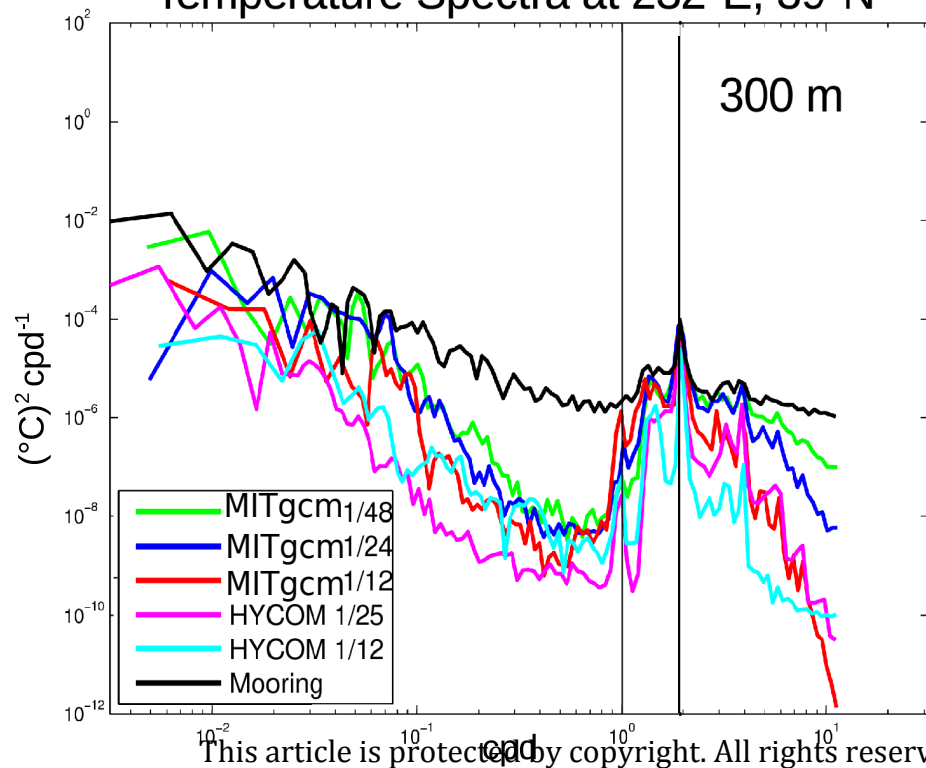
Temperature Spectra at 236°E, 39°N



KE Spectra at 298°E, 43°N



Temperature Spectra at 232°E, 39°N



KE Spectra at 288°E, 39°N

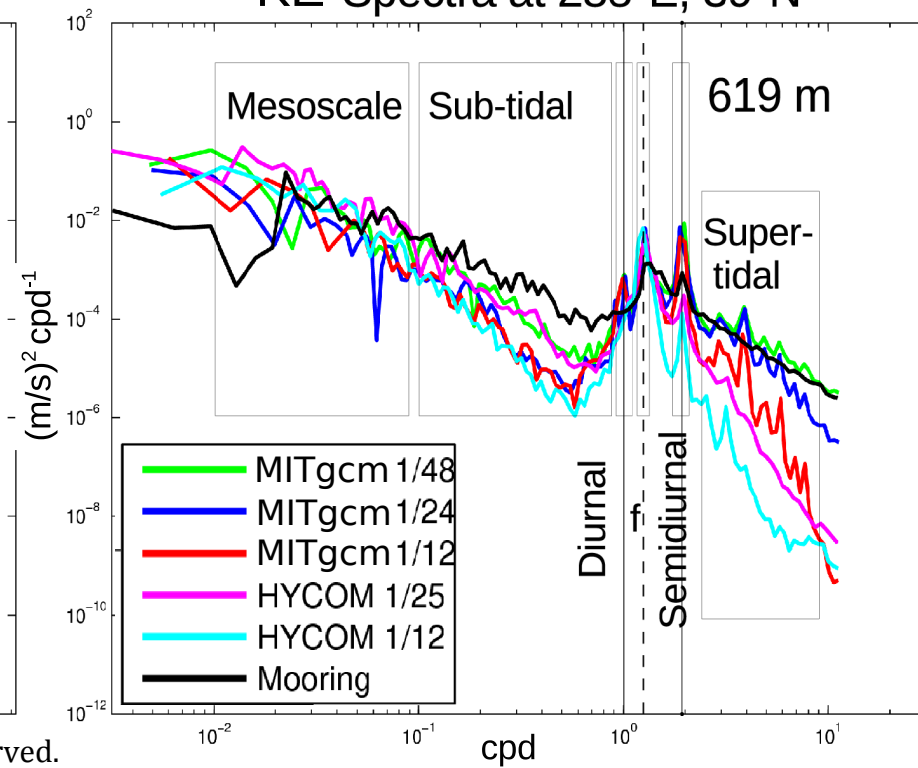
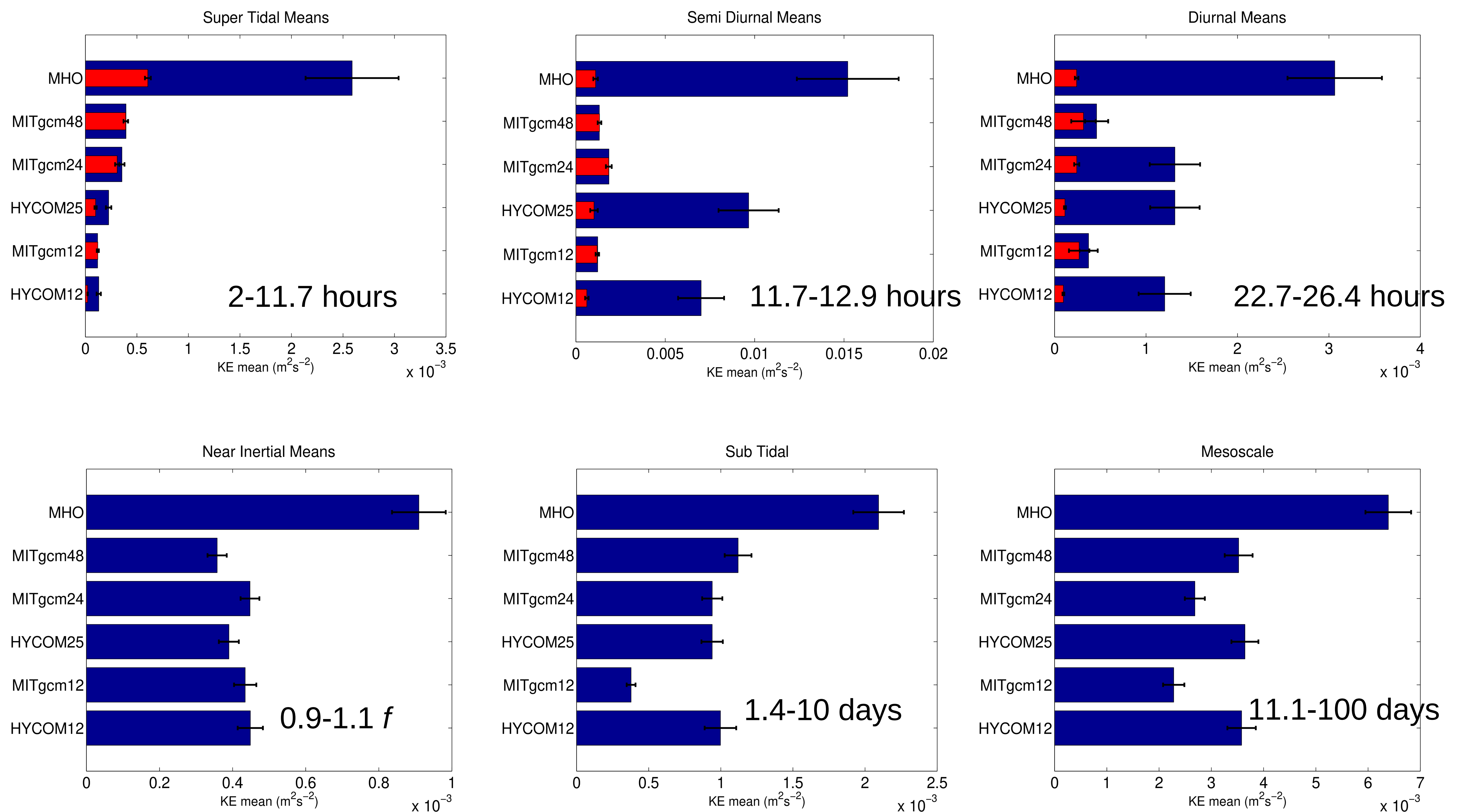


Figure 3.

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Mean KE by Frequency Band



Mean Temperature Variance by Frequency Band

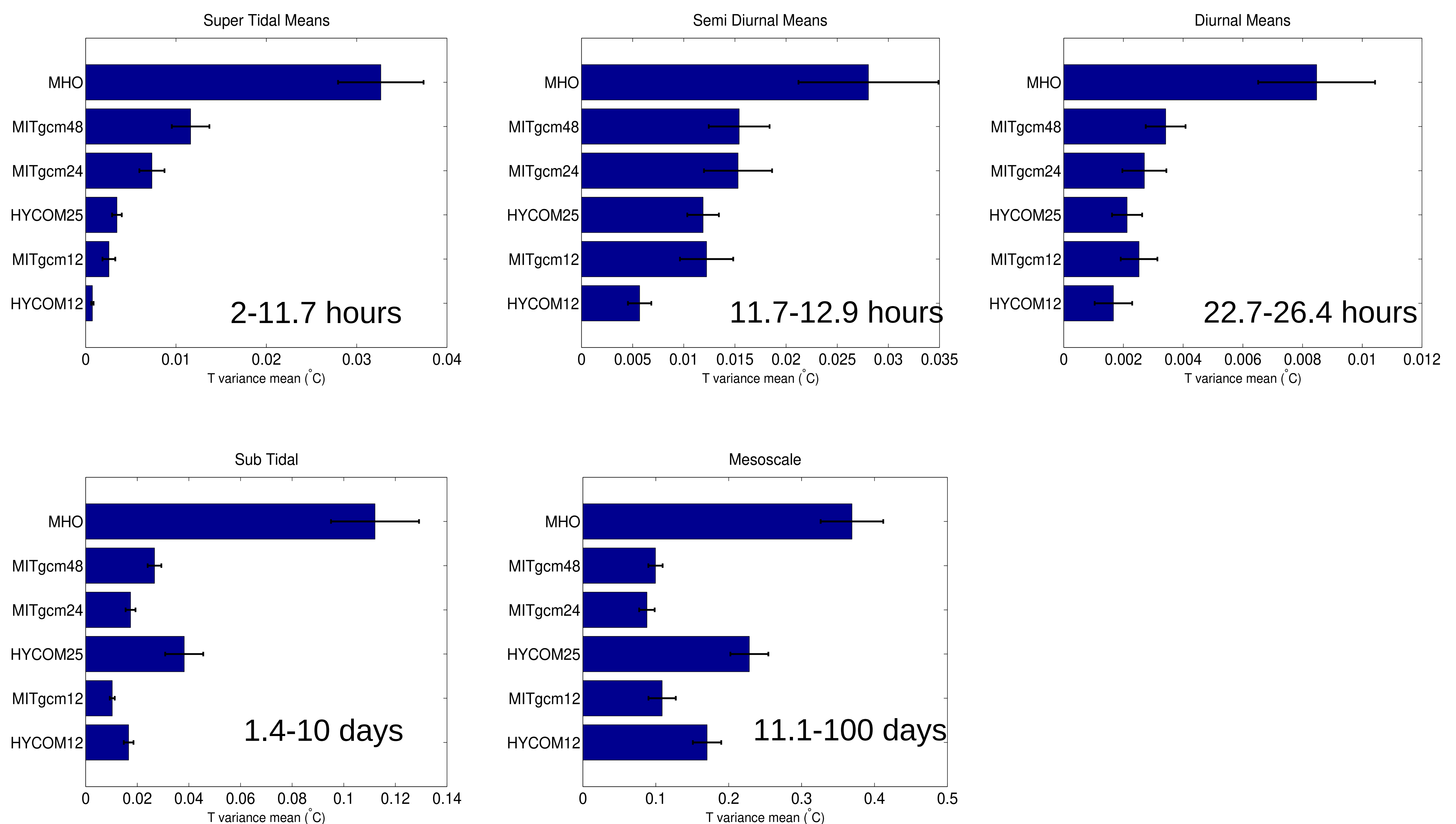
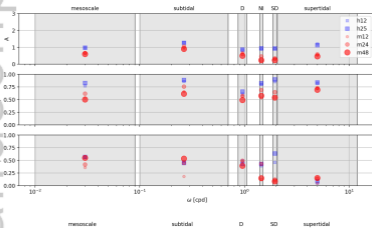


Figure 4.

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KE comparison metrics



temperature variance comparison metrics

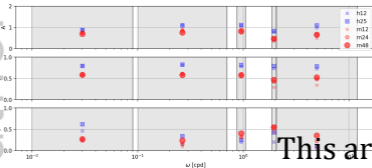
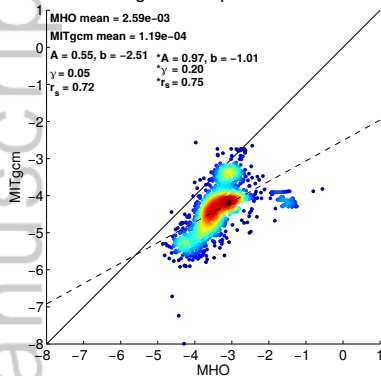


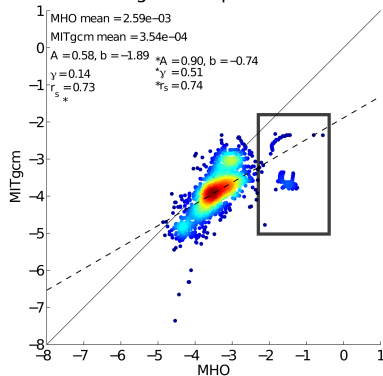
Figure 5.

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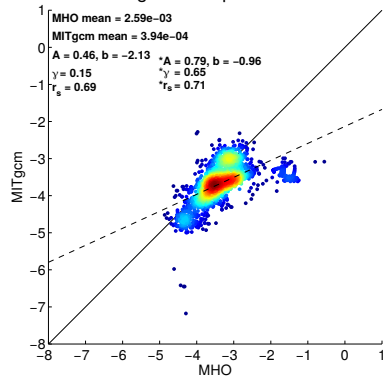
MITgcm12 Super Tidal KE



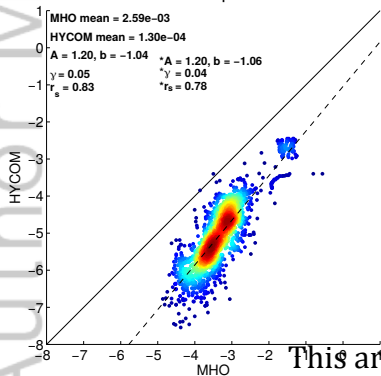
MITgcm24 Super Tidal KE



MITgcm48 Super Tidal KE



HYCOM12 Super Tidal KE



HYCOM25 Super Tidal KE

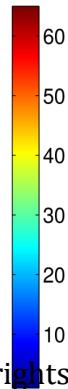
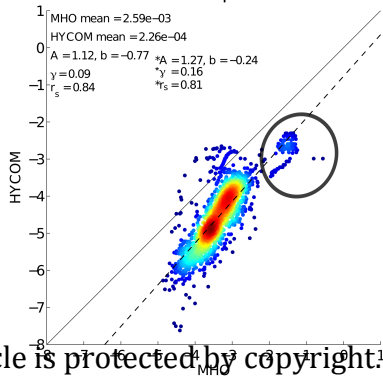
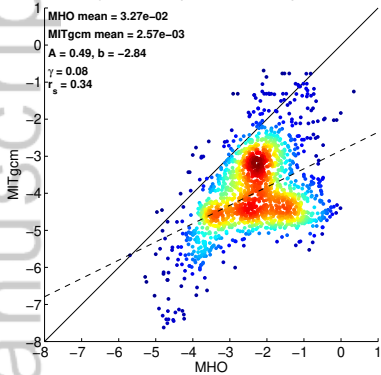


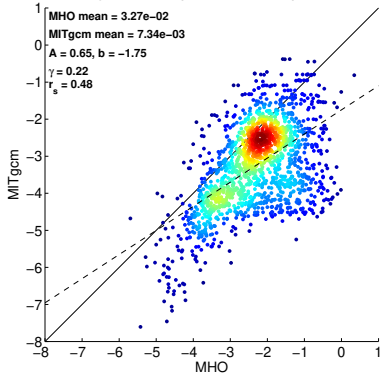
Figure 6.

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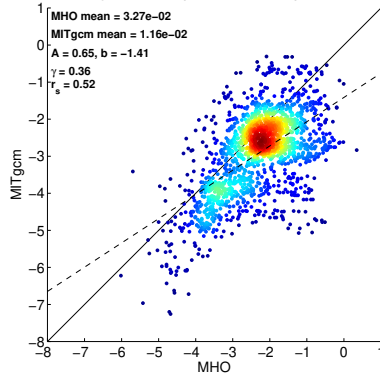
MITgcm12 Super Tidal temperature



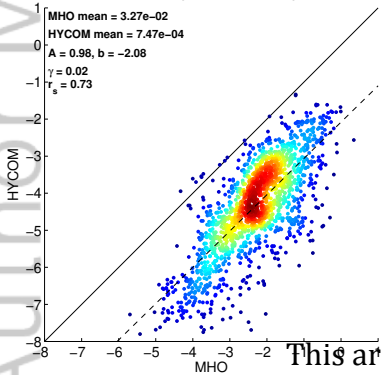
MITgcm24 Super Tidal temperature



MITgcm48 Super Tidal temperature



HYCOM12 Super Tidal temperature



HYCOM25 Super Tidal temperature

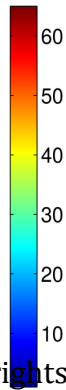
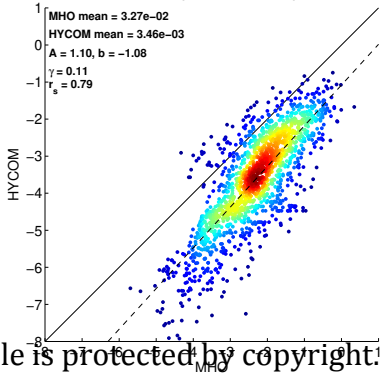
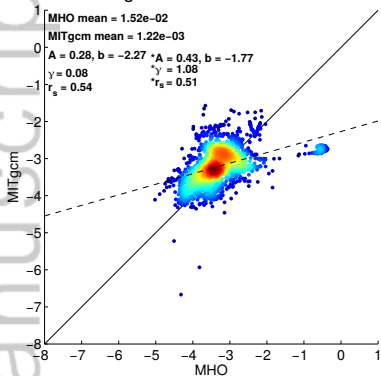


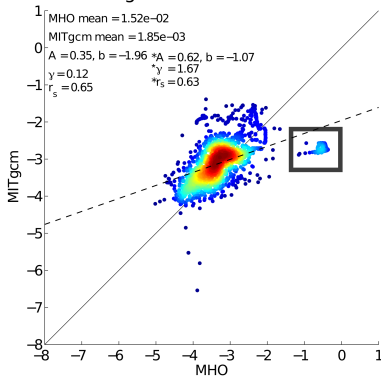
Figure 7.

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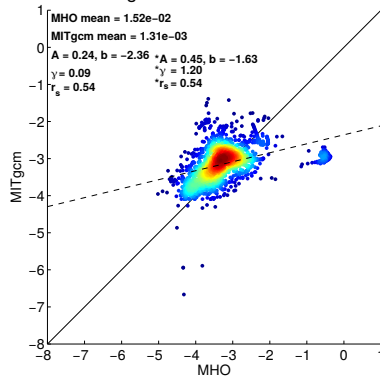
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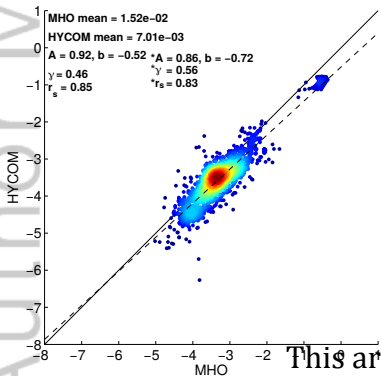
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MITgcm48 Semi-Diurnal KE



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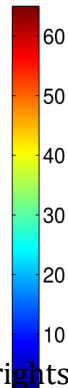
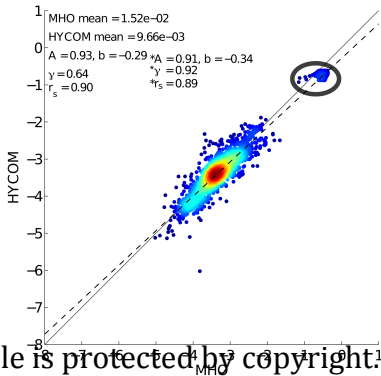
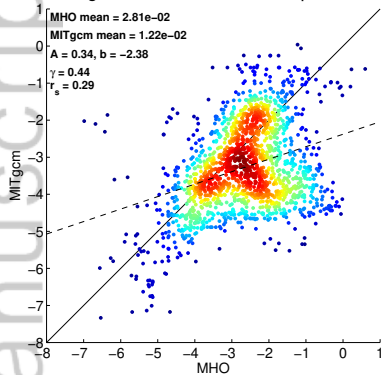


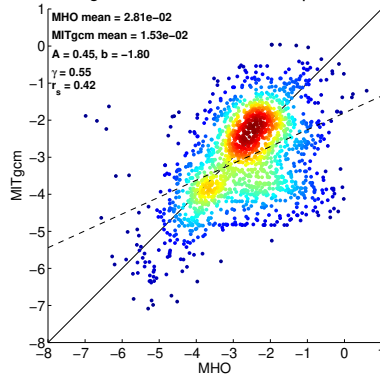
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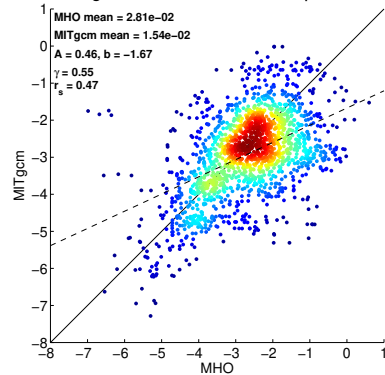
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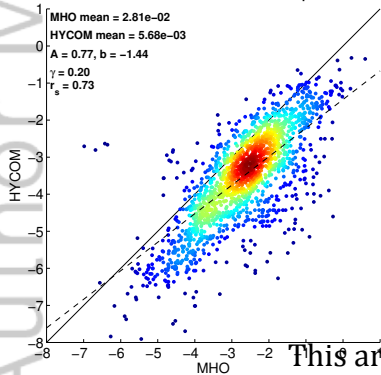
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HYCOM12 Semi-Diurnal temperature



HYCOM25 Semi-Diurnal temperature

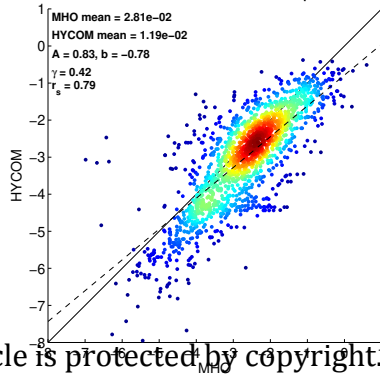
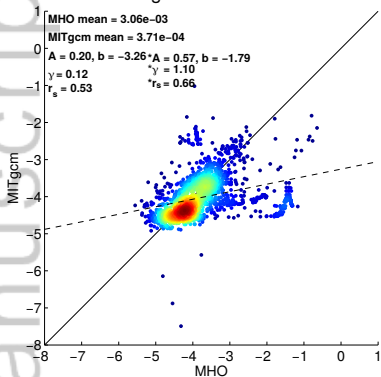


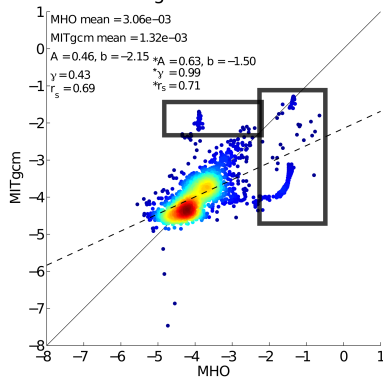
Figure 9.

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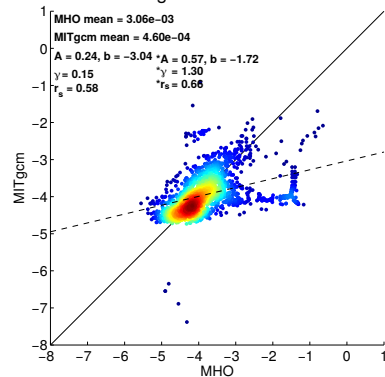
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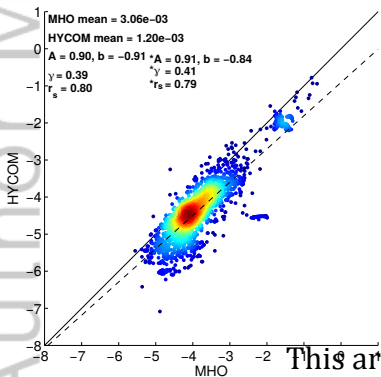
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MITgcm48 Diurnal KE



HYCOM12 Diurnal KE



HYCOM25 Diurnal KE

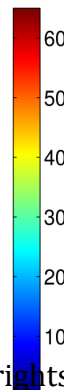
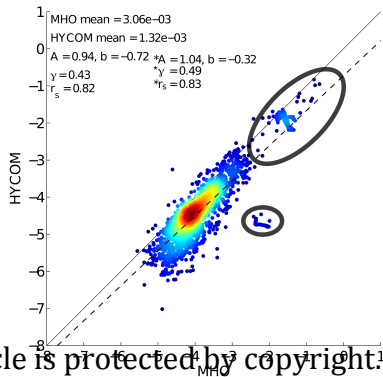
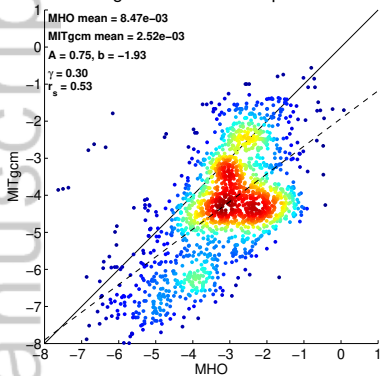


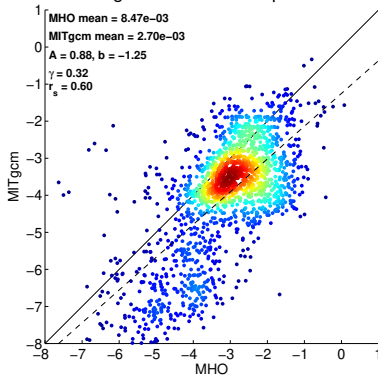
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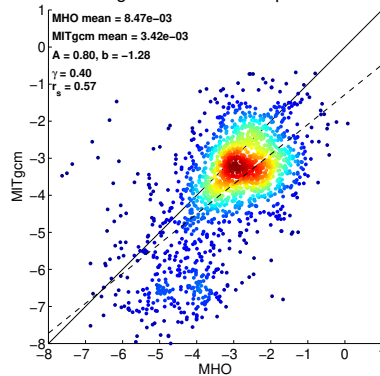
MITgcm12 Diurnal temperature



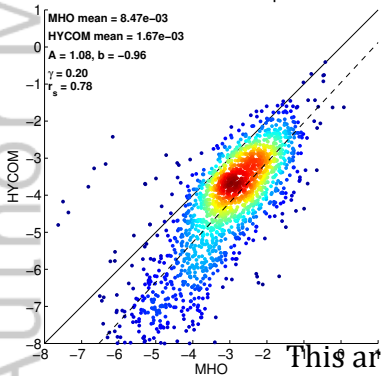
MITgcm24 Diurnal temperature



MITgcm48 Diurnal temperature



HYCOM12 Diurnal temperature



HYCOM25 Diurnal temperature

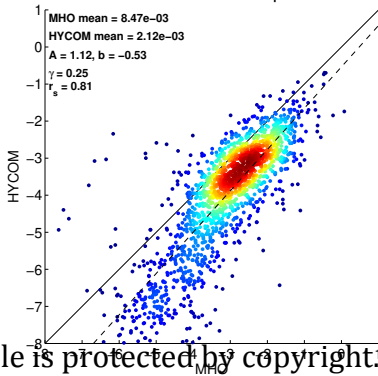
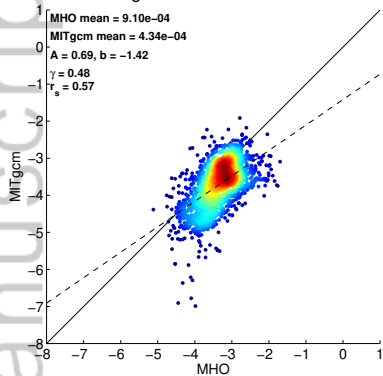


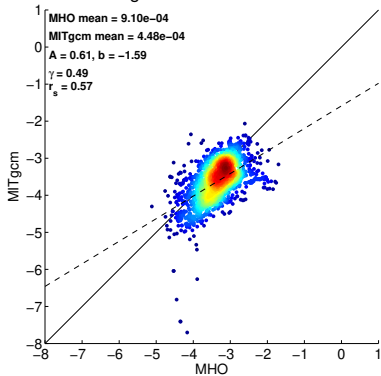
Figure 11.

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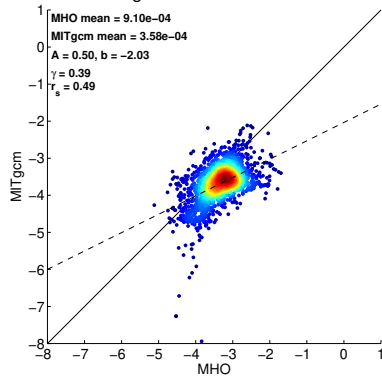
MITgcm12 Near Inertial KE



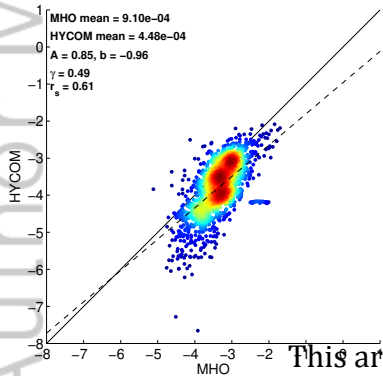
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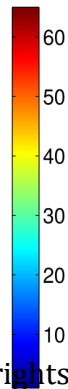
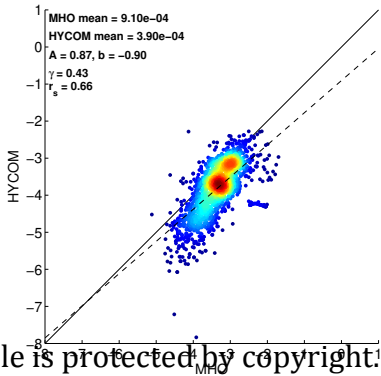
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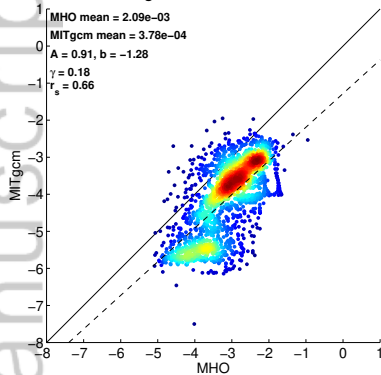
HYCOM12 Near Inertial KE



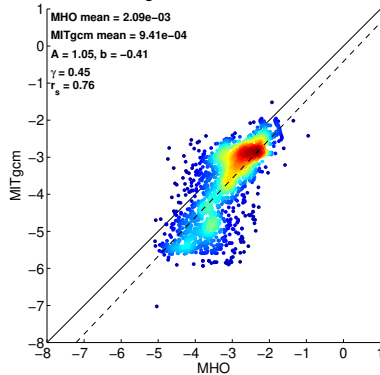
HYCOM25 Near Inertial KE



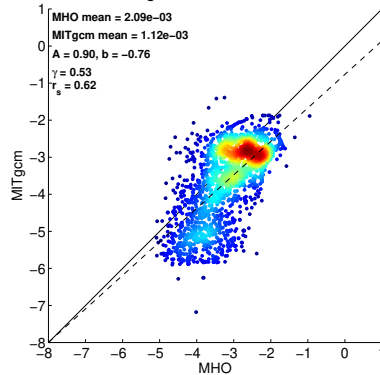
MITgcm12 Sub Tidal KE



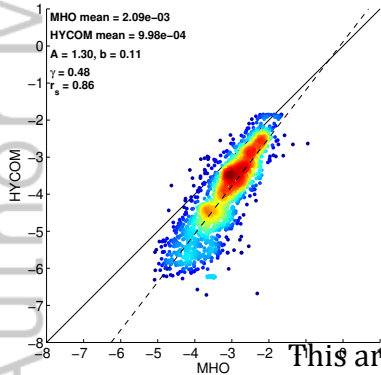
MITgcm24 Sub Tidal KE



MITgcm48 Sub Tidal KE



HYCOM12 Sub Tidal KE



HYCOM25 Sub Tidal KE

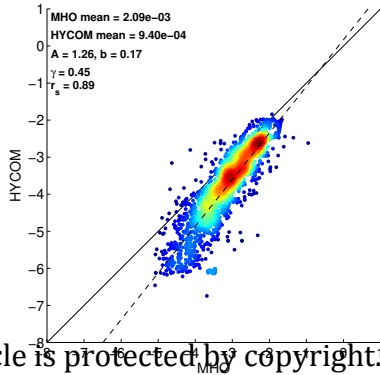
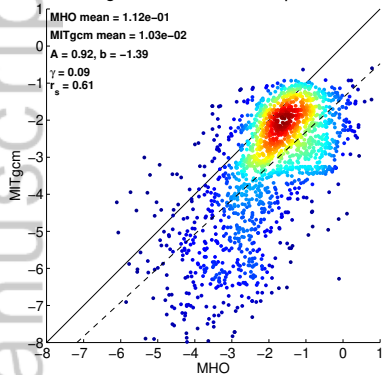


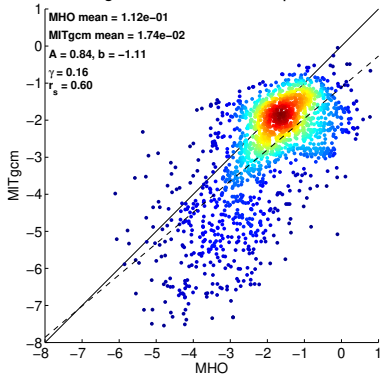
Figure 13.

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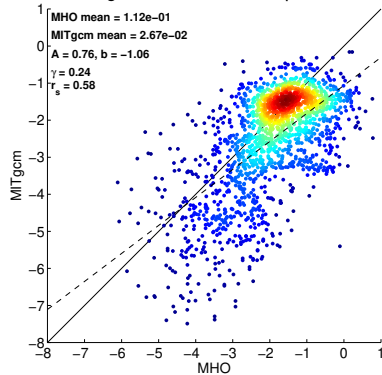
MITgcm12 Sub Tidal temperature



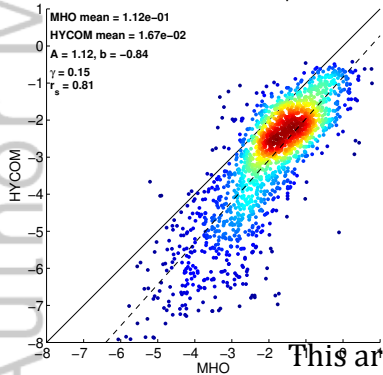
MITgcm24 Sub Tidal temperature



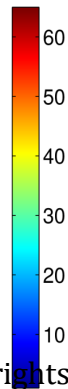
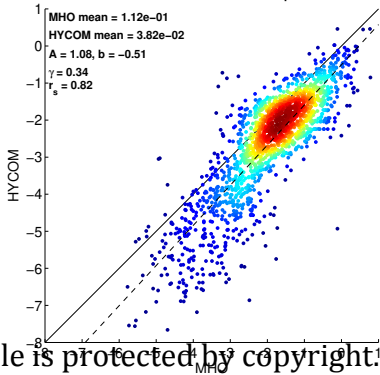
MITgcm48 Sub Tidal temperature



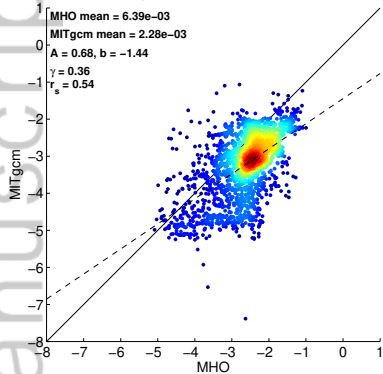
HYCOM12 Sub Tidal temperature



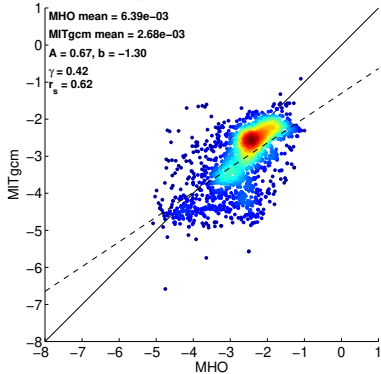
HYCOM25 Sub Tidal temperature



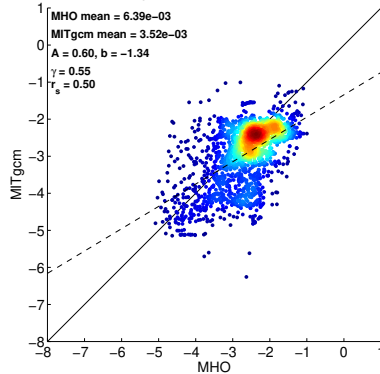
MITgcm12 Mesoscale KE



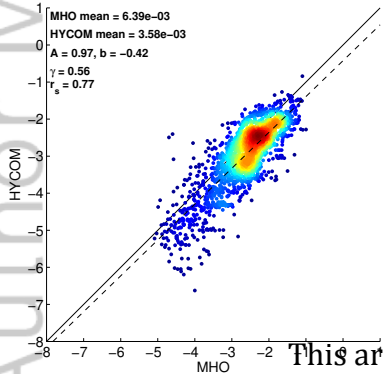
MITgcm24 Mesoscale KE



MITgcm48 Mesoscale KE



HYCOM12 Mesoscale KE



HYCOM25 Mesoscale KE

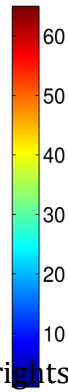
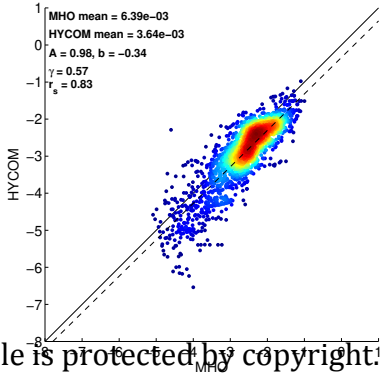


Figure 15.

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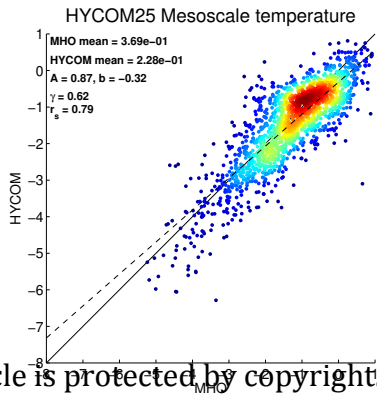
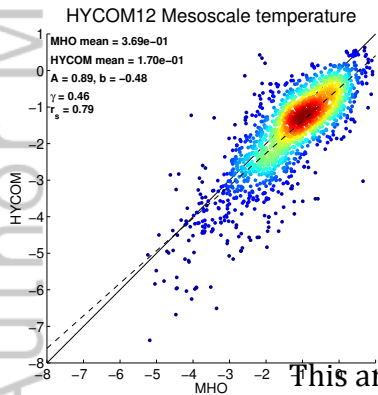
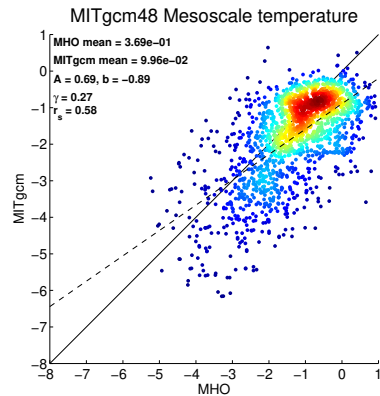
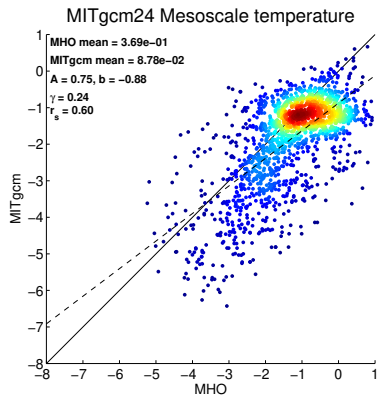
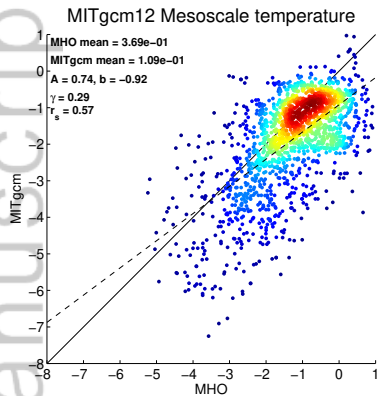


Figure 16.

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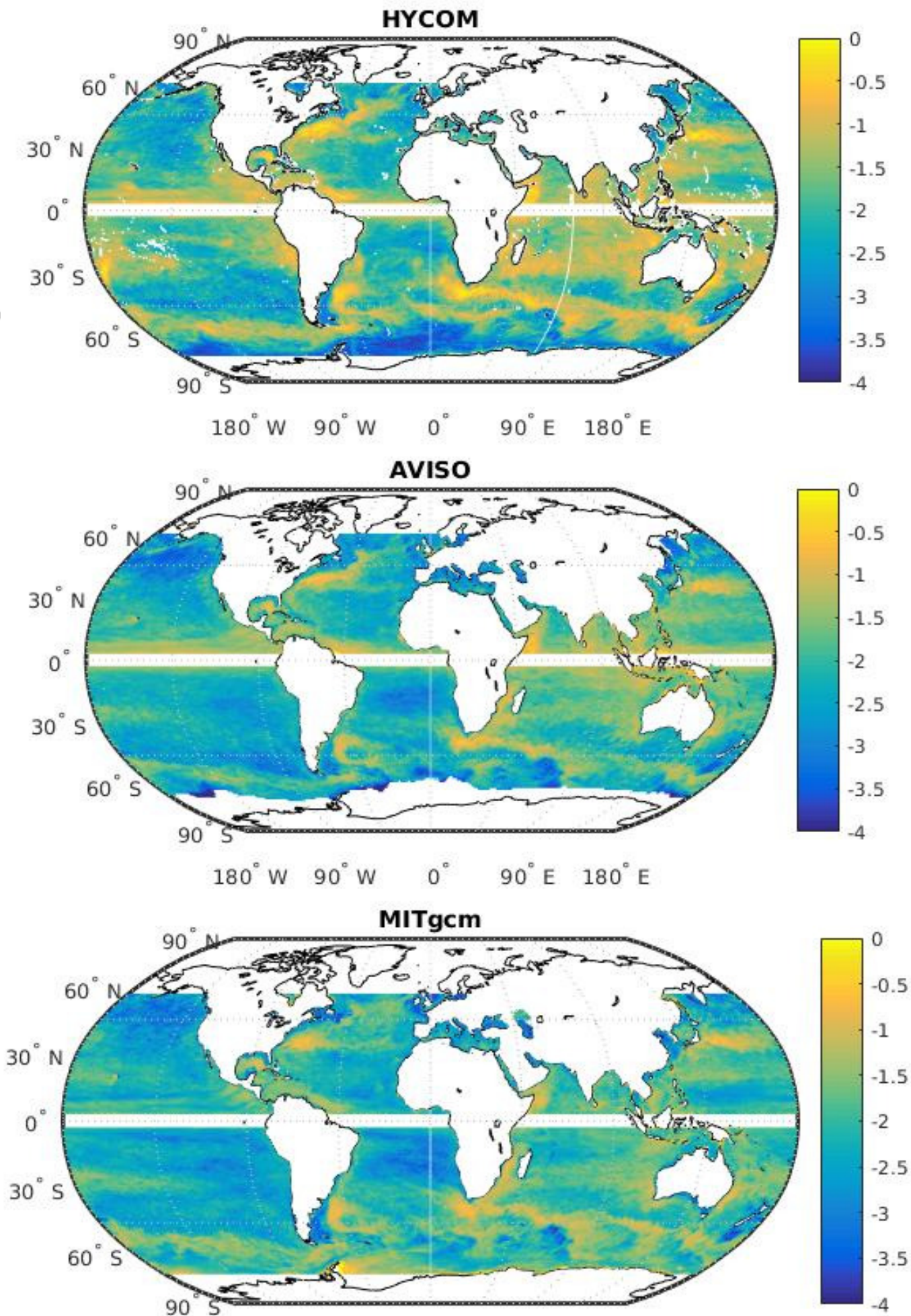
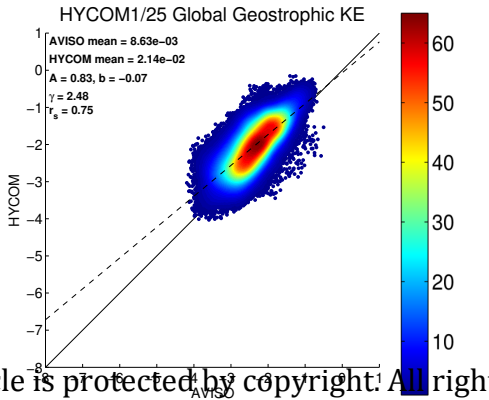
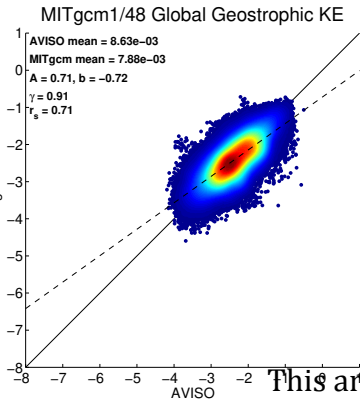
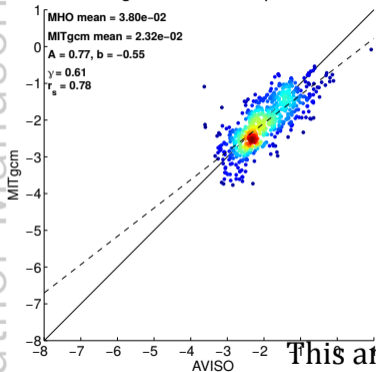


Figure 17.

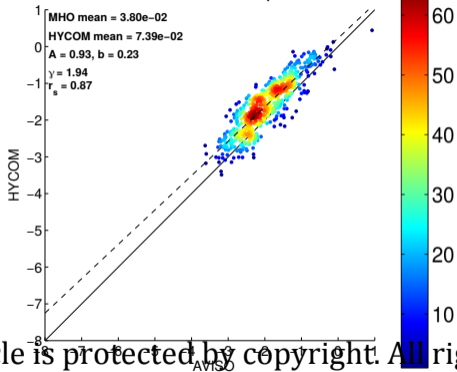
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MITgcm1/48 Geostrophic KE



HYCOM1/25 Geostrophic KE



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Statistical Comparisons of Temperature Variance and Kinetic Energy in Global Ocean Models and Observations: Results from Mesoscale to Internal Wave Frequencies

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Key Points:

- Kinetic energy and temperature variance in global HYCOM and MITgcm simulations are compared with moored observations
- Model resolution has large impact on the supertidal internal gravity wave continuum band
- Spatial correlations with observations are higher in HYCOM than in MITgcm

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Abstract

Temperature variance and kinetic energy (KE) from two global simulations of the HYbrid Coordinate Ocean Model (HYCOM; 1/12, 1/25 degree) and three global simulations of the Massachusetts Institute of Technology general circulation model (MITgcm; 1/12, 1/24, and 1/48 degree), all of which are forced by atmospheric fields and the astronomical tidal potential, are compared with temperature variance and KE from a database of about 2,000 moored historical observations (MHO). The variances are computed across frequencies ranging from supertidal, dominated by the internal gravity wave continuum, to subtidal, dominated by currents and mesoscale eddies. The most important qualitative difference between HYCOM and MITgcm, and between simulations of different resolutions, are in the supertidal band, where the 1/48 degree MITgcm lies closest to observations. Across all frequency bands examined, the HYCOM simulations display higher spatial correlation with the MHO than do the MITgcm simulations. The supertidal, semidiurnal, and diurnal velocities in the HYCOM simulations also compare more closely with observations than do the MITgcm simulations in a number of specific continental margin/marginal sea regions. To complement the model-MHO comparisons, this paper also compares the surface ocean geostrophic eddy KE in HYCOM, MITgcm, and a gridded satellite altimeter product. Consistent with the model-MHO comparisons, the HYCOM simulations have a higher spatial correlation with the altimeter product than the MITgcm simulations do. On the other hand, the surface ocean geostrophic eddy KE is too large, relative to the altimeter product, in the HYCOM simulations.

1 Introduction

In this paper, we examine the ability of high-resolution global ocean models to recreate an accurate temperature variance and kinetic energy spectrum ranging from mesoscale eddies (periods of about 30-200 days) through the high-frequency (supertidal) internal gravity wave (IGW) continuum. This ability is quantified through comparison of models with both moored historical observations and satellite altimetry.

In recent years, global ocean models have been run at finer grid spacings and timesteps. Increased computer power has led to high-resolution, three-dimensional ocean models that are able to simulate mesoscale eddies on a global scale [e.g., *Hecht and Hasumi, 2008; Maltrud and McClean, 2005; Chassignet et al., 2009*]. The HYbrid Coordinate Ocean

54 Model (HYCOM) and Massachusetts Institute of Technology general circulation model
55 (MITgcm) simulations examined here include both astronomical tidal forcing and atmo-
56 spheric forcing. High-frequency atmospheric forcing generates near-inertial waves [e.g.,
57 *Silverthorne and Toole, 2009; Simmons and Alford, 2012*], and barotropic tidal flow over
58 topographic features creates internal tides [e.g., *Garrett and Kunze, 2007*]. In global mod-
59 els with sufficiently high horizontal and vertical resolution, nonlinear wave-wave interac-
60 tions and simultaneous tidal and atmospheric forcing begin to fill out the oceanic internal
61 gravity wave continuum, as first shown in *Müller et al. [2015]*.

62 Tidal and atmospheric forcing were first employed simultaneously in high-resolution
63 general circulation simulations in HYCOM [*Arbic et al., 2010, 2012*], and the HYCOM
64 tidal simulations were shown to have an internal gravity wave continuum spectrum [*Müller*
65 *et al., 2015; Savage et al., 2017a,b*]. Later, MITgcm simulations were run with higher ver-
66 tical and horizontal resolutions, and were also shown to have an IGW spectrum [*Rocha*
67 *et al., 2016a,b; Savage et al., 2017b; Torres et al., 2018; Wang et al., 2019*]. An overview
68 of global internal tide and wave modeling in HYCOM and MITgcm is given in *Arbic et al.*
69 [2018], which also briefly describes two other community global high-resolution simula-
70 tions with simultaneous atmospheric and tidal forcing.

71 Resolving a spectrum of internal gravity waves (hereafter, often referred to simply
72 as “internal waves”, or IGWs) represents a new paradigm for global ocean models. In-
73 ternal waves transport substantial amounts of energy over long distances, and breaking
74 IGWs drive most of the mixing in the open-ocean beneath the mixed layer. The three-
75 dimensional geography of mixing is important for large-scale ocean dynamics [*Munk and*
76 *Wunsch, 1998; Melet et al., 2016; MacKinnon et al., 2017; Kunze, 2017*]. As a first step
77 towards potentially using the simulations reported on here to better understand mixing
78 geography, we must assess the realism of the modeled internal tide and IGW continuum
79 fields. Model-data comparisons of the latest, high-resolution global general circulation
80 models such as the simulations presented here have several benefits to the oceanographic
81 community. Thorough comparisons of these simulations with observations helps to expose
82 the faults and successes of individual changes in model configuration. Additionally, high
83 resolution global simulations have been increasingly used to inform and supplement obser-
84 vational experiments. As a result, there is broad interest in the observational community
85 in the realism of particular model solutions both geographically and with respect to differ-
86 ent physical regimes such as the ones presented in this paper.

87 High-resolution global ocean models have been compared with observations in a
88 number of studies. For instance, low-frequency sea surface height (SSH) variance in high-
89 resolution global models has been compared with altimeter results in works such as *Mal-*
90 *trud and McClean* [2005] and *Chassignet and Xu* [2017], amongst others. Comparisons
91 of low-frequency kinetic energy in high-resolution global models with mooring, surface
92 drifter, and altimeter observations have been done in several studies [e.g., *Maltrud and Mc-*
93 *Clean*, 2005; *Penduff et al.*, 2006; *Scott et al.*, 2010; *Thoppil et al.*, 2011]. Low-frequency
94 eddy available potential energy in global HYCOM simulations without tidal forcing was
95 compared to mooring and Argo float observations in *Luecke et al.* [2017].

96 A number of model-data comparisons have been performed with the global HYCOM
97 and (to a lesser extent) MITgcm tidal simulations. However these comparisons often fo-
98 cus on a specific frequency band, a particular geographic location, or both. For example,
99 barotropic and internal tide sea surface height signatures in HYCOM have been compared
100 with altimetry [*Shriver et al.*, 2012; *Stammer et al.*, 2014; *Ansong et al.*, 2015; *Ngodock*
101 *et al.*, 2016]. HYCOM tidal kinetic energy and internal tide energy fluxes have been com-
102 pared with moored in-situ observations in *Timko et al.* [2012, 2013] and *Ansong et al.*
103 [2017] respectively. In *Buijsman et al.* [2016] the semidiurnal internal tide dissipation
104 rates are compared with Argo inferred dissipation rates. *Savage et al.* [2017b] compares
105 the dynamic height variance in both HYCOM and MITgcm to the variance computed from
106 moored McLane profiler data, across a wide range of frequencies. The McLane profilers
107 [*Doherty et al.*, 1999] used in *Savage et al.* [2017b] were set up to study internal grav-
108 ity waves, and have high vertical and temporal resolution. However *Savage et al.* [2017b]
109 found only 9 such records that were useful for comparing with models in the open ocean.
110 In *Qiu et al.* [2018], geostrophic surface eddy kinetic energy (EKE) in MITgcm is com-
111 pared with AVISO satellite data. *Rocha et al.* [2016a] compared kinetic energy wavenum-
112 ber spectra in the MITgcm with along-track Acoustic Doppler Current Profiler (ADCP)
113 data in one Southern Ocean location. *Wang et al.* [2018] compares MITgcm and ADCP
114 velocity for the purpose of mooring design for ground-truthing future satellite missions.

115 The various observational datasets that we compare our HYCOM and MITgcm solu-
116 tions to all have their advantages and disadvantages. For instance, altimetry is global, but
117 only measures the sea surface heights. The mooring archive that we use here is as exten-
118 sive as we could find, in terms of global coverage, though it still does not offer the dense
119 spatial coverage offered by, for instance, altimetry and surface drifters. On the other hand,

120 in contrast to those datasets, moorings offer the ability to both look beneath the surface
121 and to examine a broad range of frequencies, encompassing a broad range of dynamical
122 regimes.

123 In this study, we compare temperature variance and KE in three global simulations
124 of MITgcm and two global simulations of HYCOM to moored historical observations. By
125 examining different frequency bands, we test the ability of these simulations to model mo-
126 tions ranging from mesoscale eddies to the internal gravity wave continuum. Our study
127 is quasi-global, in that about 2,000 records, distributed over the major ocean basins, are
128 used. Our study builds upon our model-data comparison of low-frequency eddy available
129 potential energy, in *Luecke et al.* [2017], to examine temperature variance over a wider
130 range of frequencies. Both the HYCOM and MITgcm simulations are performed at dif-
131 ferent resolutions, thus allowing us to study the impact of resolution on model-data com-
132 parisons. Additionally, in order to obtain denser spatial coverage of motions in the low-
133 frequency band dominated by mesoscale eddies, we compare surface ocean geostrophic
134 eddy kinetic energy from both MITgcm and HYCOM to values computed from the Archiving,
135 Validation, and Interpretation of Satellite Oceanographic data (AVISO) gridded satel-
136 lite altimeter product [*Ducet et al.*, 2000]. Our study is unique in being a quasi-global
137 comparison of models and moored in-situ data, across five simulations of varying res-
138 olutions performed with two different models, of both temperature variance and kinetic
139 energy (KE), across a range of frequencies running from the mesoscale eddy band to the
140 internal gravity wave continuum.

141 **2 Models, Observations, and Methods**

142 We use temperature variance and KE spectra from HYCOM simulations of two res-
143 olutions ($1/12.5^\circ$ and $1/25^\circ$), and MITgcm simulations of three resolutions ($1/12^\circ$, $1/24^\circ$,
144 and $1/48^\circ$). More general information about HYCOM can be found in *Chassignet et al.*
145 [2009], and more general information about MITgcm can be found in *Marshall et al.* [1997].
146 Temperature variance and KE spectra computed from the models are compared to spec-
147 tra calculated from a database of moored historical observations (MHO) obtained from
148 the Global Multi-Archive Current Meter Database [GMACMD; *Wright et al.* [2014]]. For
149 our comparison of surface ocean geostrophic eddy kinetic energy in HYCOM, MITgcm,
150 and the AVISO satellite altimeter product, we use the AVISO and MITgcm results from
151 Figure (1) of *Qiu et al.* [2018]. Qiu et al. used the global SSH anomaly product created

152 by Ssalto/Duacs and distributed by AVISO with support from CNES (<http://marine.copernicus.eu/>). The AVISO product we use here is reported on a 0.25 degree horizontal grid with a 7-day temporal increment, and is generated via the spatial-temporal interpolation of altimetric data [Ducet *et al.*, 2000]. The surface ocean geostrophic eddy KE field in HYCOM was computed by us following the same methodology used in Qiu *et al.* [2018] for the MITgcm results, as confirmed in a 2018 personal communication with Bo Qiu. The HYCOM and MITgcm results are both computed by differentiating the SSH field computed from spatially and temporally averaging hourly model output, from the high resolution model grid onto the AVISO grid and temporal increment. This is done so that model results are qualitatively consistent with the spatial-temporal smoothing inherent in the creation of the AVISO product.

163 2.1 Moored Historical Observations

164 Time-series of temperature and velocity from moored historical observations (MHO) are obtained from the GMACMD and can be found at <http://stockage.univ-brest.fr/~scott/GMACMD/gmacmd.html>. These records span from 1974 to 2008, and are not contemporaneous with the model runs. Only locations with seafloor depths greater than 168 500 meters are included. We select MHO time-series that are longer than 90 days, and exclude records containing gaps in the time-series. The remaining temperature and velocity records are then visually inspected and quality controlled for instrument errors and other problems such as severe discretization.

172 Our selection criteria yield a total of 1,711 instrumental records of temperature and 2,102 instrumental records of velocity, some of which are co-located, distributed around the global ocean. The geographical locations are shown in the top panel of Figure 1. The spatial coverage of both temperature and velocity is sparse and uneven, with some basins such as the North Atlantic and North Pacific having a higher density of observations than others. The bottom panel of Figure 1 shows the distribution of instruments by depth for velocity (right) and temperature (left) records. There is a clear preference for measurements in the upper 1000m.

2.2 HYCOM

Two simulations of HYCOM are used in this comparison; HYCOM12, with a 1/12.5 degree horizontal resolution, and HYCOM25, with a 1/25 degree horizontal resolution. The spatial scales of the model grids are approximately 8km and 4km at the equator respectively. The HYCOM output is saved hourly, and HYCOM12 output spans October 1, 2011 through April 1, 2012, while HYCOM25 output spans January 2014 through December 2014. Both simulations contain 41 hybrid layers in the vertical direction. Atmospheric fields, such as pressure, buoyancy, and wind forcing, used in both HYCOM simulations, are taken from the U.S. Navy Global Environmental Model, NAVGEM [Hogan *et al.*, 2014]. HYCOM12 is forced by NAVGEM on three hour intervals, while HYCOM25 is forced hourly. The HYCOM simulations are forced using a 0.5° application grid interpolated from the NAVGEM primary 37 km grid.

The HYCOM simulations are additionally forced by the astronomical tidal potential of the two largest diurnal constituents (K_1 and O_1) along with the three largest semidiurnal constituents (M_2 , S_2 , and N_2) [Cartwright, 1999]. These five tidal constituents account for 97 percent of the global variance in the 10 largest tidal constituents in the Global Ocean Tide Model [GOT99.2; Ray, 1999]. The self-attraction and loading (SAL) [Hendershott, 1972] term is taken from the altimeter-constrained TPXO model of Egbert *et al.* [1994] and Egbert and Erofeeva [2002]. An Augmented State Ensemble Kalman Filter is implemented in both simulations to reduce the global RMS error of M_2 barotropic tidal elevations with respect to TPXO, in waters deeper than 1 km, to approximately 2.6 cm [Ngodock *et al.*, 2016].

A Smagorinsky scheme is employed for horizontal viscosity and a Laplacian scheme is utilized for horizontal diffusivity, while a KPP scheme [Large *et al.*, 1994] is used for both vertical diffusivity and viscosity. The parameterized topographic wave drag field taken from Jayne and St. Laurent [2001] is tuned to minimize barotropic tidal errors with respect to TPXO [Egbert *et al.*, 1994; Egbert and Erofeeva, 2002]. The wave drag strength employed in HYCOM25 was half the strength of the drag used in HYCOM12. A description of the wave drag tuning can be found in Buijsman *et al.* [2015], and more information on the impact of the wave drag on barotropic and baroclinic tides can be found in Ansong *et al.* [2015] and Buijsman *et al.* [2016].

2.3 MITgcm

Three global ocean simulations of MITgcm [Marshall *et al.*, 1997], at horizontal resolutions of 1/12, 1/24, and 1/48 degree, are used in this comparison and are referred to as MITgcm12, MITgcm24, and MITgcm48 respectively. All three MITgcm simulations have 90 z-levels in the vertical, with level thicknesses varying from 1 m at the surface to 480 m near the bottom at the maximum model depth of 7 km. Bathymetry is taken from the Smith and Sandwell [1997] Version 14.1 and IBCAO Version 2.23 Jakobsson *et al.* [2008]. MITgcm is forced at the surface with atmospheric fields from the 0.14° European Center for Medium-Range Weather Forecasts (ECMWF) atmospheric operational model analysis at six-hour intervals. Starting in 2011, atmospheric forcing is converted to surface fluxes using the bulk formulae of Large and Yeager [2004]. In regions of ice coverage, ocean surface fluxes are computed using the sea ice model of Losch *et al.* [2010].

The MITgcm runs employ the full luni-solar tidal potential of Weis *et al.* [2008], which is applied to MITgcm as additional atmospheric pressure forcing [Ponte *et al.*, 2015]. No parameterized topographic wave drag was applied to the MITgcm simulations.

A Leith scheme is used for horizontal diffusivity and a KPP scheme is used for vertical diffusivity. The MITgcm12 simulation is initialized on 1 January 2010 from a data-constrained 1/6° simulation provided by the Estimating the Circulation and Climate of the Ocean, Phase II project (ECCO2) [Menemenlis *et al.*, 2008] and integrated for 1 year without tides with ERA-Interim [Dee *et al.*, 2011] surface boundary conditions. Tidal forcing and atmospheric boundary conditions from the 0.14° ECMWF analysis are applied starting on 1 January 2011. The MITgcm24 simulation is initialized from MITgcm12 fields on 17 January 2011. The MITgcm48 simulation is initialized from MITgcm24 fields on 10 September 2011.

We use ~ 7 months of hourly model output from the three MITgcm simulations for the following periods: 1 January 2012 through 20 July 2012 for MITgcm12, 2 October 2012 through 22 April 2013 for MITgcm24, and 28 January 2012 through 22 August 2012 for MITgcm48. The MITgcm records are only 7 months in length due to the fact that, at the time we extracted the model output, it had only been run for this long (after a suitable spinup period).

2.4 Methods

In order to compute frequency spectra, a linear trend and mean are removed from each time series of temperature $T(t)$, zonal velocity $U(t)$ and meridional velocity $V(t)$, where t denotes time. Each detrended time series is multiplied by a Tukey window having a taper to constant ratio of 0.2. Approximately 10-15 percent of the total variance is lost as a result of the application of the Tukey window.

The frequency spectra are computed from each time series for each MHO instrument, and from the output of the corresponding nearest horizontal neighbor model grid points interpolated to the depth of the MHO instrument. A discrete Fourier Transform is defined as:

$$\widehat{\text{Field}}(\omega) = \sum_{t=0}^{N-1} \text{Field}(t)e^{-i\omega t}, \quad (1)$$

where ‘Field’ denotes either temperature or velocity, ω denotes frequency, and N denotes the total number of points in the time series. In the case of temperature, variance is calculated as

$$\text{Temperature variance} = \frac{2\delta t}{N} \int_{\omega_{min}}^{\omega_{max}} |\widehat{T(\omega)}|^2 d\omega \quad (2)$$

and for kinetic energy, half the total velocity variance,

$$\text{KE} = \frac{\delta t}{N} \int_{\omega_{min}}^{\omega_{max}} |(\widehat{U(\omega)})^2 + (\widehat{V(\omega)})^2| d\omega, \quad (3)$$

where δt is the temporal sampling interval, and ω_{min} and ω_{max} represent the lower and upper bounds of the frequency band of interest. We integrate over five bands, supertidal (12-2.06 cpd), semidiurnal (2.05-1.86 cpd), diurnal (1.05-0.87 cpd), subtidal (0.7-0.1 cpd) and mesoscale (0.09-0.01 cpd). In the case of MHO records shorter than 100 days in length, the mesoscale band is integrated to the lowest possible frequency. Additionally, for KE only, we integrate over the near-inertial band defined as $0.9 - 1.1f$, where f is the local Coriolis frequency. This represents a ‘double counting’ of KE in regions where f and the diurnal tidal frequencies overlap.

We employ several statistical metrics to quantify differences between temperature variance and kinetic energy in the models and observations. A linear regression in the form $\log_{10}(\text{model}) = A * \log_{10}(\text{data}) + b$, for both temperature variance and kinetic energy

269 is calculated. A ratio γ between the mean of the model variances and the mean of the
 270 observational variances is defined as:

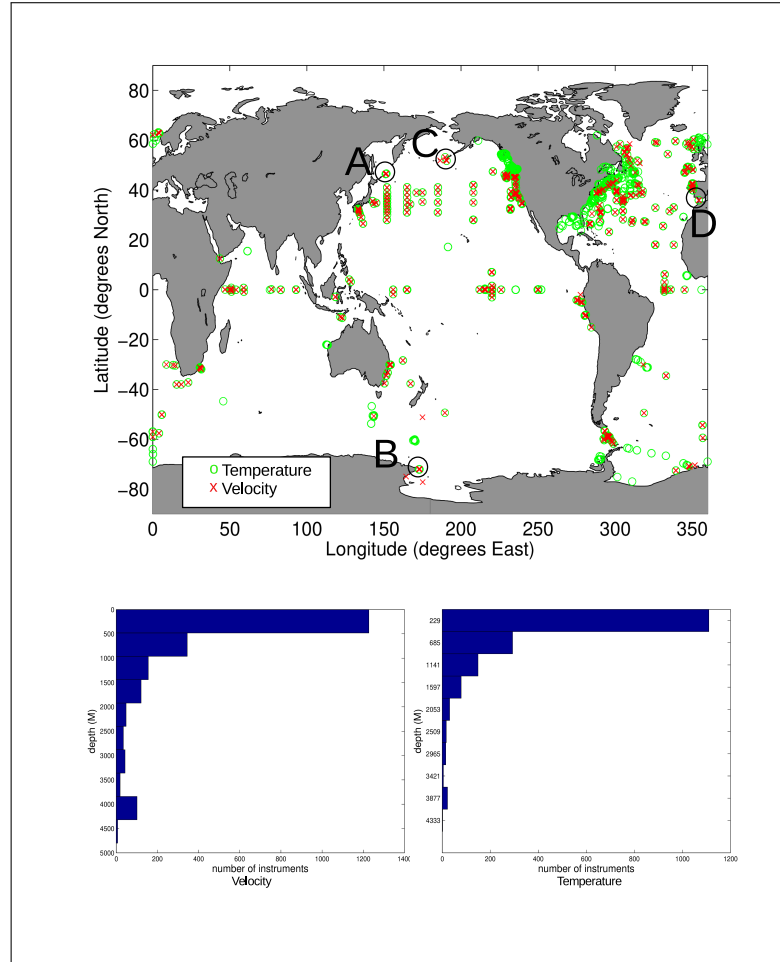
$$271 \quad \gamma = \frac{\sum_{i=1}^n \text{Variance}_{\text{model}}}{\sum_{i=1}^n \text{Variance}_{\text{observed}}}, \quad (4)$$

272 where i is an instrument index and n is the total number of instruments used in the calcu-
 273 lation. In cases where a horizontal location contains MHO instruments at multiple depths,
 274 the instruments at each depth are counted as separate instruments. Finally, a Spearman
 275 correlation coefficient r_s [Kokoska and Zwillinger, 2000] is calculated between model and
 276 data variance across the MHO locations. The Spearman correlation has an advantage of
 277 being less impacted by individual strong outliers in the data. Additionally, because Spear-
 278 man is a rank correlation, it provides an understanding of how the models perform from
 279 low to high energetics, even when the relationship between model bias and variance is not
 280 linear. The ideal values expressing a perfect comparison are equal to one for all of the
 281 metrics A , γ , and r_s , and zero for b .

282 In order to calculate our AVISO-derived eddy kinetic energy (EKE) values from
 283 model output, we average the modeled SSH data into a weekly time series on a 0.25 de-
 284 gree longitude Mercator grid and compute the EKE from the geostrophic velocity field as
 285 in Qiu *et al.* [2018] where:

$$286 \quad u = -\frac{g}{f} \frac{\partial \eta}{\partial y}, \quad \text{and} \quad v = \frac{g}{f} \frac{\partial \eta}{\partial x}. \quad (5)$$

287 Here g is the local acceleration due to gravity, f is the Coriolus frequency, and x and y
 288 are the zonal and meridional spatial coordinates. Equation (5) gives the geostrophic veloc-
 289 ities for u and v respectively, and $\text{EKE} = \frac{1}{2}(u^2 + v^2)$.

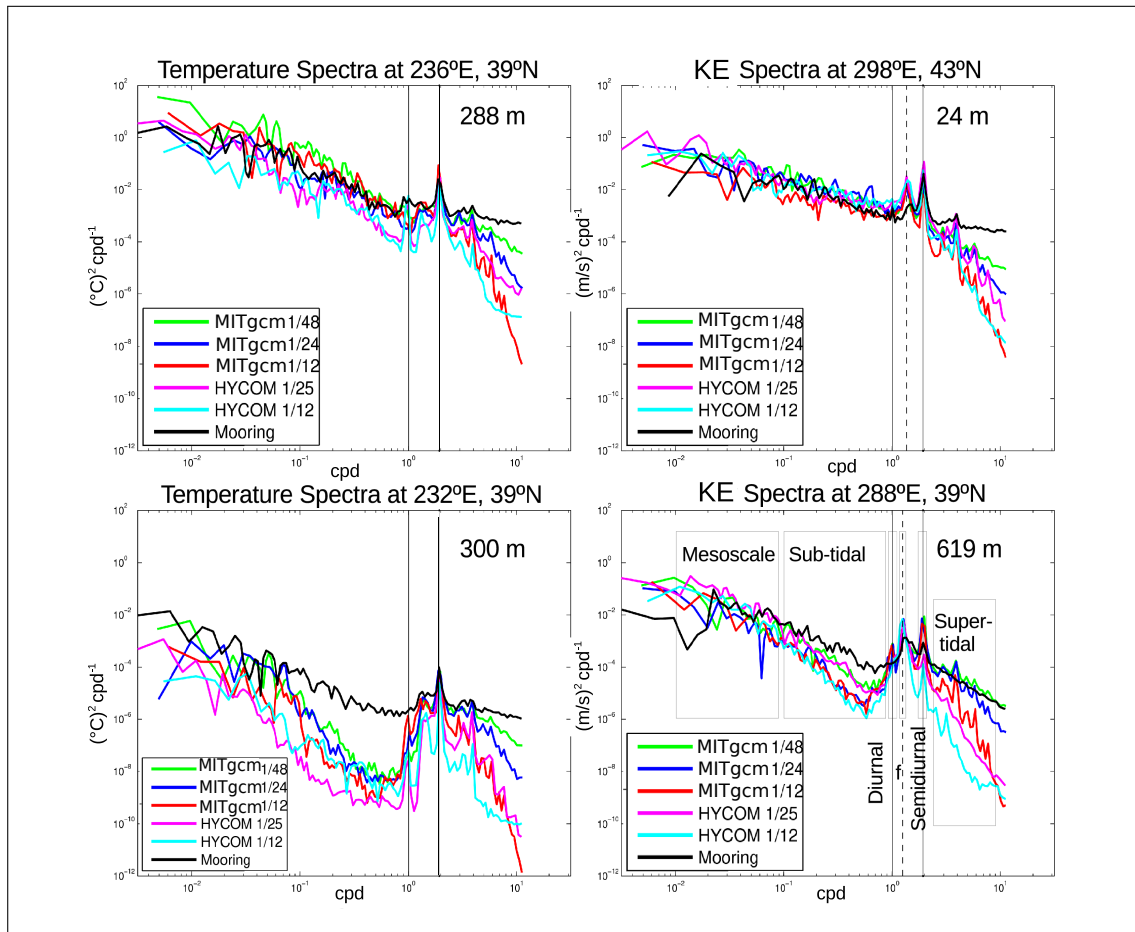


290 **Figure 1.** Geographical locations and depths of MHO records used in the model-MHO comparison. Tem-
 291 perature observations are shown with green circles, while velocity observations are denoted by red x's. The
 292 bottom panels show the depth distribution for velocity (left) and temperature (right) instrument record loca-
 293 tions. The labels “A”, “B”, “C”, and “D” will be referred to later in the text.

294 3 Model-MHO comparisons

295 In this section, we present global model-MHO comparisons of temperature vari-
 296 ance and KE. Results for temperature variance and KE have been grouped into the six
 297 frequency bands defined above: supertidal, semi-diurnal, diurnal, near-inertial, subtidal,
 298 and mesoscale, where again the near-inertial band is defined only for KE, not for tempera-
 299 ture variance. Before we summarize and discuss the model vs. MHO comparisons in each
 300 frequency band, we present sample frequency spectra of the temperature variance and KE,
 301 for four individual instrument locations shown in Figure 2. The left two panels in Fig-
 302 ure 2 depict typical temperature variance spectra, while the right two panels show typical

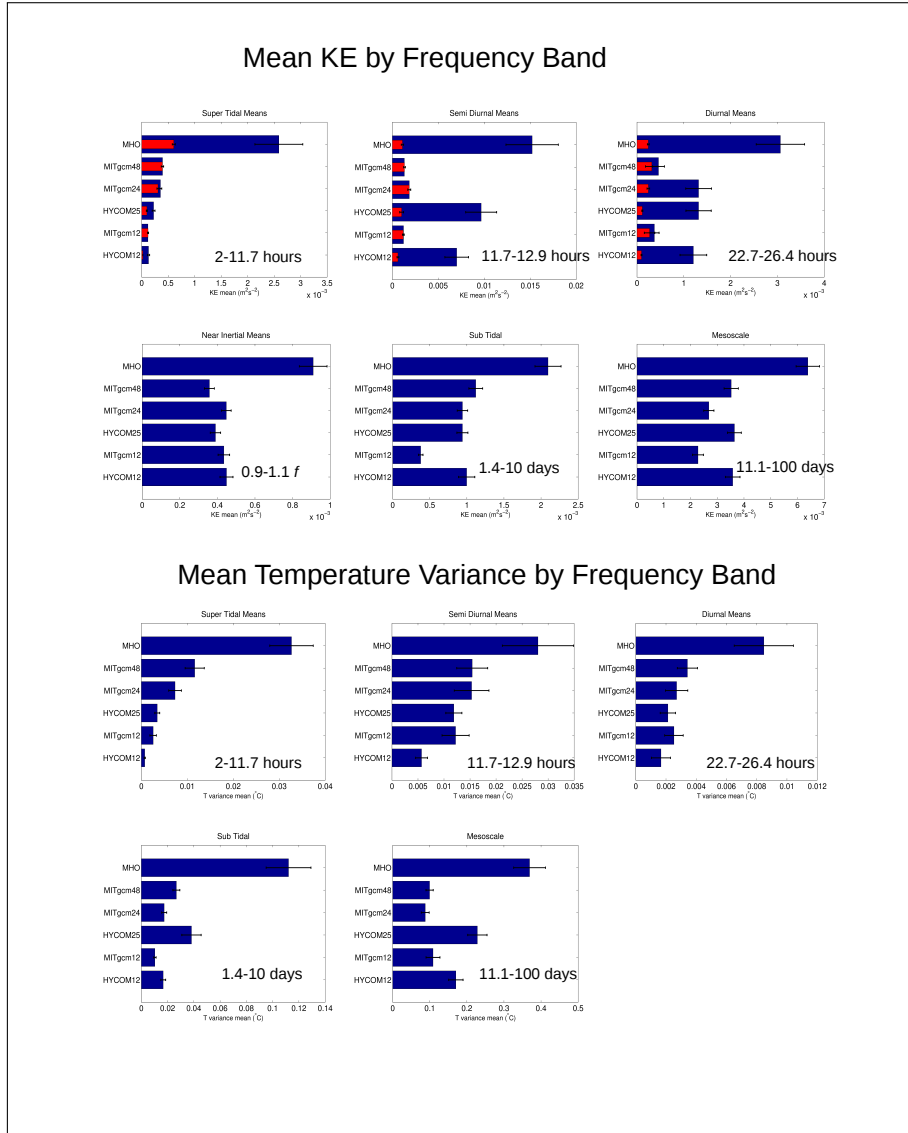
303 KE spectra. At supertidal frequencies, the highest resolution model (MITgcm48) clearly
 304 has more variance, and matches the observations more closely than the lower resolution
 305 models. The lowest resolution simulations (MITgcm12 and HYCOM12) have the least
 306 variance at supertidal frequencies. Also of note is a model deficiency in both KE (bottom
 307 right panel) and temperature variance (bottom left panel) that occurs in between 0.1 and 1
 308 cpd. The lack of model energy in this band for some instruments has been noted in other
 309 model-data comparisons, for instance in *Savage et al.* [2017b], and prompted our division
 310 of lower-frequency motions into subtidal and mesoscale.



311 **Figure 2.** Sample spectra of temperature variance (in the Eastern Pacific) and KE (Northeast Atlantic) for
 312 all 5 simulations. Instrument locations and depths are given in the subplots. The solid vertical lines show the
 313 diurnal (left) and semi-diurnal (right) tidal frequencies, and the dashed vertical lines show the local Coriolis
 314 frequency. The approximate bounds of each frequency band discussed in this paper are delineated by boxes in
 315 the lower right subplot.

3.1 Summary of Mean Temperature Variance and KE in all Frequency Bands

Here we summarize the model-MHO comparisons, using the statistical metrics described in Section 2.4, computed over all instruments, across all five simulations and all six frequency bands. Figure 3 shows the spatial mean KE (top) and temperature variance (bottom) within each frequency band, for all five simulations and for the MHO, computed over all instruments at the MHO locations and depths. As seen in the differences between the red and blue bars in the top panel, the spatial means are sensitive to inclusion of a small number of continental margin/marginal sea locations, denoted by “A”, “B”, “C”, and “D” in figure 1. We will return to this point later in the paper. In some frequency bands, within each model, increased resolution improves the comparison between models and data. However, this does not hold universally across all simulations and frequency bands.

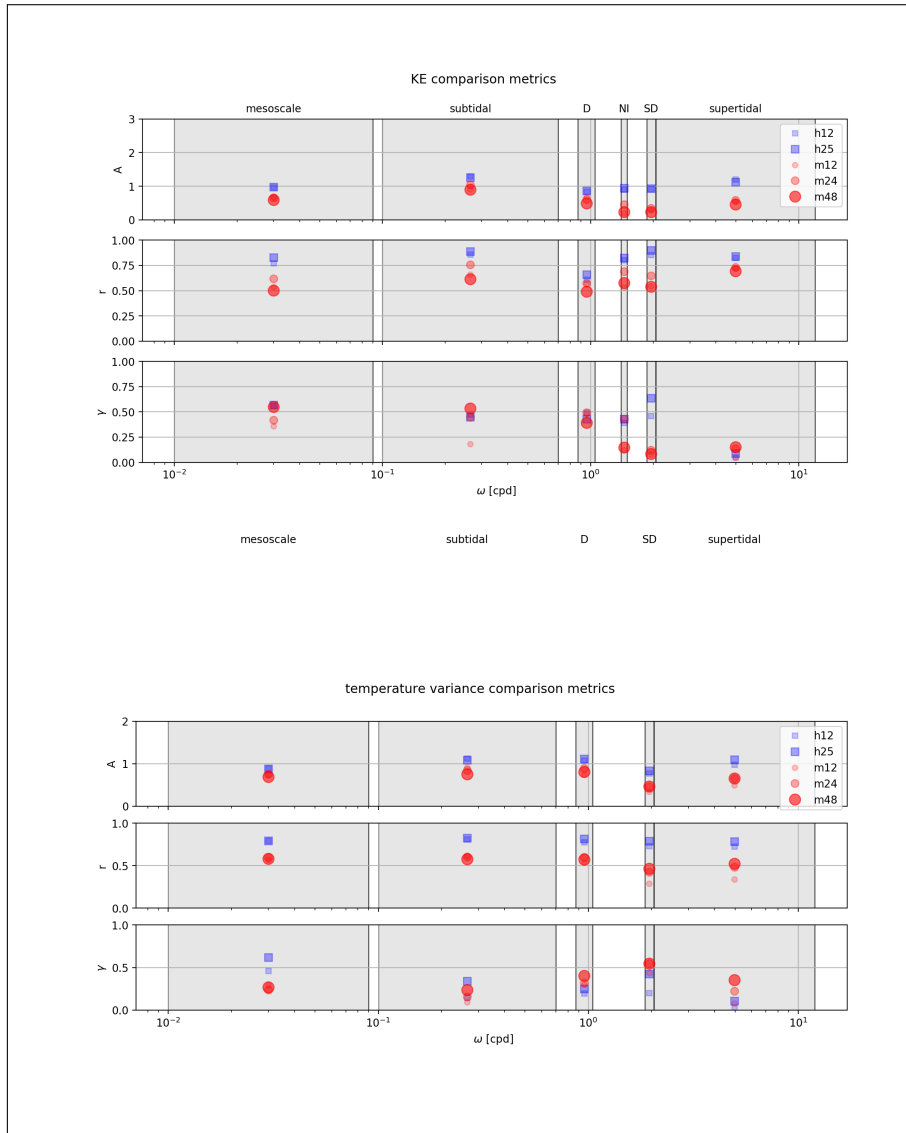


327 **Figure 3.** (top) Spatial mean KE ($\text{m}^2 \text{s}^{-2}$), computed over the MHO instrument locations and depths,
 328 across the frequency bands defined in the text. (Bottom) Spatial mean temperature variance ($^\circ\text{C}^2$) separated
 329 by frequency band. The red bars denote averages for which results at a small number of continental margin/
 330 margin/marginal sea locations (denoted by “A”, “B”, “C”, and “D” in figure 1) have been removed. Note that the
 331 near-inertial band is only computed and displayed for KE.

332 Model improvement with increased resolution is particularly noticeable in the case
 333 of both supertidal KE and temperature variance, where increased model resolution leads
 334 to more total variance, and closer comparisons with MHO results for both MITgcm and
 335 HYCOM. Increased resolution also appears to improve temperature variance in both MIT-
 336 gcm and HYCOM within both the semidiurnal and diurnal bands, but not the KE within

337 these bands. Within the near-inertial band, changes in resolution appear to make little dif-
338 ference, possibly because of the relatively large horizontal length scales associated with
339 low-mode near-inertial waves [Simmons and Alford, 2012]. Within the sub-tidal frequency
340 band, both MITgcm KE and HYCOM temperature variance show closer agreement in
341 model-MHO comparisons as model resolution is increased. Within the mesoscale band,
342 both models improve with resolution with respect to KE, but only HYCOM improves with
343 respect to temperature.

344 Figure 4 shows the statistical comparison metrics for all model runs across all fre-
345 quency bands and across all instruments, without any removal of anomalous results at
346 continental margin/marginal sea locations. Across all frequency bands, for both KE and
347 temperature variance, HYCOM has a universally higher spatial correlation r_s with the
348 MHO data than does MITgcm. The linear regressions tend to be closer to their ideal val-
349 ues within the HYCOM simulations than in MITgcm, likely as a result of the smaller
350 amount of scatter in HYCOM. As is seen in Figure 3, increases in model resolution tend
351 to increase γ for both temperature variance and kinetic energy, with the tidal and super-
352 tidal temperature variance showing the largest changes with model resolution, and with the
353 near-inertial band KE being a notable exception. The dramatic increase in both KE and
354 temperature variance within the supertidal frequency band suggests that the increase in
355 resolution in both HYCOM and MITgcm is resulting in a large increase in energy within
356 the IGW continuum. Within the semi-diurnal and diurnal frequency bands, it appears as
357 if differences between HYCOM and MITgcm, such as the presence or absence of param-
358 eterized topographic wave drag, make a larger difference than resolution with respect to
359 model-data agreement. As will be described shortly, due to the wide range of KE and
360 temperature variance in the MHO instruments, the overall means and statistics are sensi-
361 tive to extreme values. More nuanced details concerning the model-MHO comparisons in
362 individual frequency bands will be discussed below.



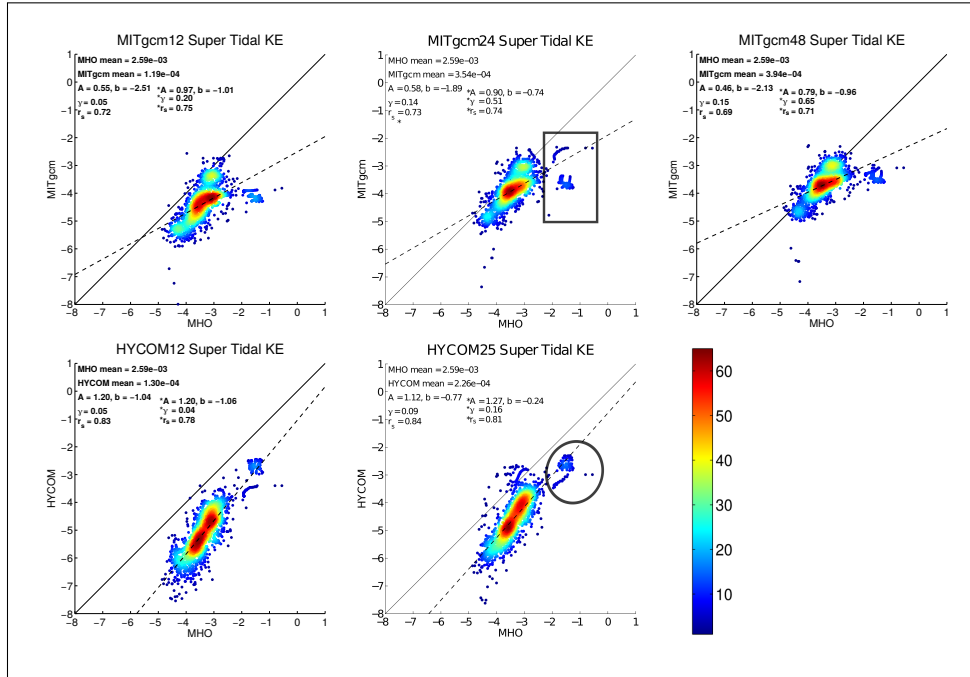
363 **Figure 4.** (top) Statistical comparison metrics, computed over all instrument locations and depths, for KE
 364 in all five simulations, and separated by the frequency bands defined in the text. (bottom) Same but for tem-
 365 perature variance, with the near-inertial band excluded. Note that for this Figure only, we use “D” and “SD”
 366 to denote diurnal and semidiurnal bands respectively, “h12” and “h25” to denote HYCOM12 and HYCOM25,
 367 respectively, and “m12”, “m24”, and “m48” to denote MITgcm12, MITgcm24, and MITgcm48, respectively.
 368 Note that in reality, the near-inertial frequency varies with latitude, and is plotted here at a fixed frequency for
 369 visualization purposes.

3.2 Supertidal Frequency Band

Figure 5 shows scatterplots of band-integrated MITgcm and HYCOM supertidal KE values against MHO values. Comparisons are shown for varying model resolutions at the same MHO locations and depths given in Figure 1. As noted previously, for the supertidal band especially, increased model resolution yields increased KE in both MITgcm and HYCOM. From visual inspection, HYCOM12 (bottom left) is biased lowest compared to MHO, while MITgcm48 (upper right), although still somewhat lower than observations, lies closest to the one-to-one line. The comparison metrics, which are given on each individual subplot of Figure 5, confirm this, as HYCOM12 has both the lowest ratio of the means, γ , while MITgcm48 has the highest values. The closer model-MHO comparison for MITgcm48 relative to HYCOM25 in the supertidal band is likely due in part to the higher horizontal and vertical resolutions in MITgcm48, but may also partly result from the overly energetic internal tides in the MITgcm simulations, which will be discussed in section 3.3. The correlation coefficients are slightly higher for HYCOM than for MITgcm, implying that the supertidal KE has a geographical distribution that is slightly more accurate in HYCOM than in MITgcm, even if the HYCOM values are generally too low.

There is a distinct group of points representing energies around ($10^{-1}\text{m}^2\text{s}^{-2}$) in the MHO data. The anomalous nature of this group of points is clearly visible in all five subplots of Figure 5. The group has been enclosed with an ellipse in the HYCOM25 panel of Figure 5, and enclosed with a box in the MITgcm24 panel. The geographical locations of these anomalous points correspond to the locations in circles labeled “A”, “B”, “C” and “D” in Figure 1, all lying in continental margin/marginal sea areas. At these locations, MITgcm is systematically under-energetic with respect to the MHO results, while HYCOM lies closer to the observations. The statistical metrics computed upon removal of the “A”, “B”, “C” and “D” values are also given and are marked with an asterisk in the subplots of Figure 5 (note that these locations were also removed in the calculation of the results marked in red in the upper panel of Figure 3). The metric values change quantitatively, sometimes significantly, when the anomalies are removed; the spatial correlation r_3 becomes slightly lower for HYCOM and slightly higher for MITgcm, with MITgcm correlation values still being lower than the HYCOM values. The linear regression values fall closer to the one-to-one line, and the ratio of means γ become larger for all of the MITgcm simulations shown in Figure 5. While the linear regressions in MITgcm are more heavily dominated by scatter than in HYCOM, it is clear for both models that increased

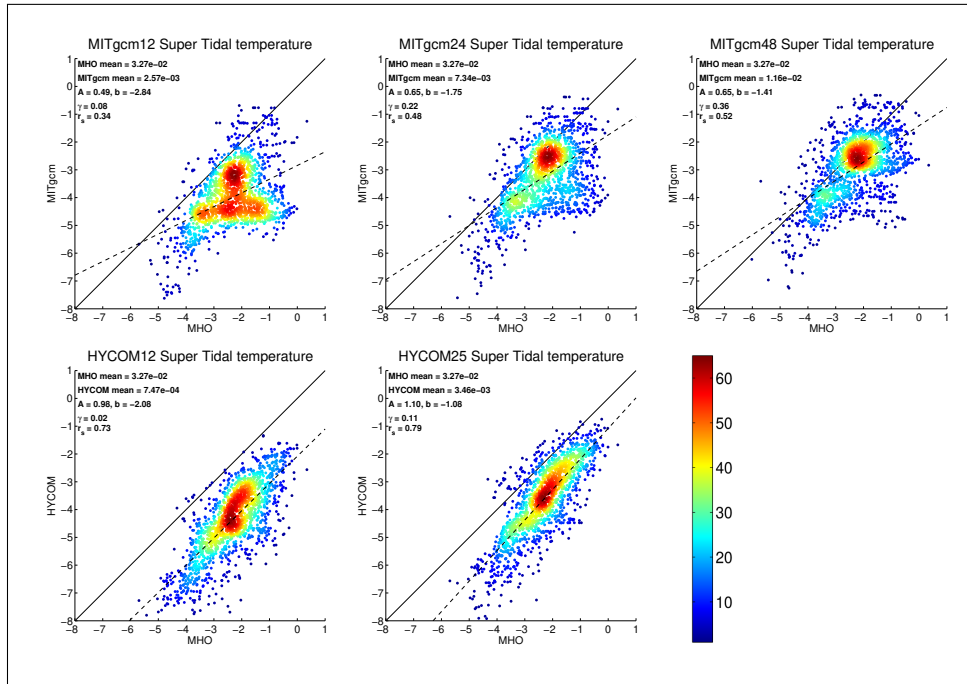
403 resolution results in closer agreement with the data, especially after the anomalies in MIT-
 404 gcm are accounted for. The general trends in the metrics noted previously, before removal
 405 of the anomalous values— higher γ values in MITgcm, higher r_s values in HYCOM—still
 406 hold when the anomalies are removed.



407 **Figure 5.** A point-to-point comparison of supertidal kinetic energy [$\log_{10}(\text{m}^2 \text{s}^{-2})$] between model and
 408 MHO. The top panels show MITgcm12, MITgcm24, and MITgcm48, while the bottom panels show HY-
 409 COM12 and HYCOM25. Population density is given by color, with the most tightly grouped data shown
 410 in red, and the sparsest data in blue. The one-to-one line is shown in solid black, and the linear regression
 411 is shown as a dashed line. Spatial means of MHO and model values, and statistical metrics A , b , γ , and r_s
 412 between each model and the MHO data, are printed on the upper left of each subplot. The second, right-most,
 413 statistical values marked with an asterisk, are computed after the anomalous high-velocity values, enclosed by
 414 a box in the MITgcm24 subplot and an ellipse in the HYCOM25 subplot, have been removed.

415 As with KE, supertidal temperature variance shows improvement with model resolution
 416 in both HYCOM and MITgcm (Figure 6). As with KE, HYCOM12 (bottom left)
 417 is biased lowest compared to MHO, while MITgcm48 (upper right), although still some-
 418 what lower than observations, lies closest to the one-to-one line. Again, the comparison
 419 metrics confirm this, as HYCOM12 has both the lowest ratio of the means, γ , while MIT-
 420 gcm48 possesses the highest. In both MITgcm24 and MITgcm48, the linear regressions

421 are again more dominated by the relatively large scatter in supertidal temperature variance,
 422 however both models improve with an increase in resolution. The correlation coefficients
 423 r_s for supertidal temperature variance are higher in HYCOM than in MITgcm. There are
 424 no visually obvious anomalies in the supertidal temperature variance plots as there were
 425 on the KE plots.



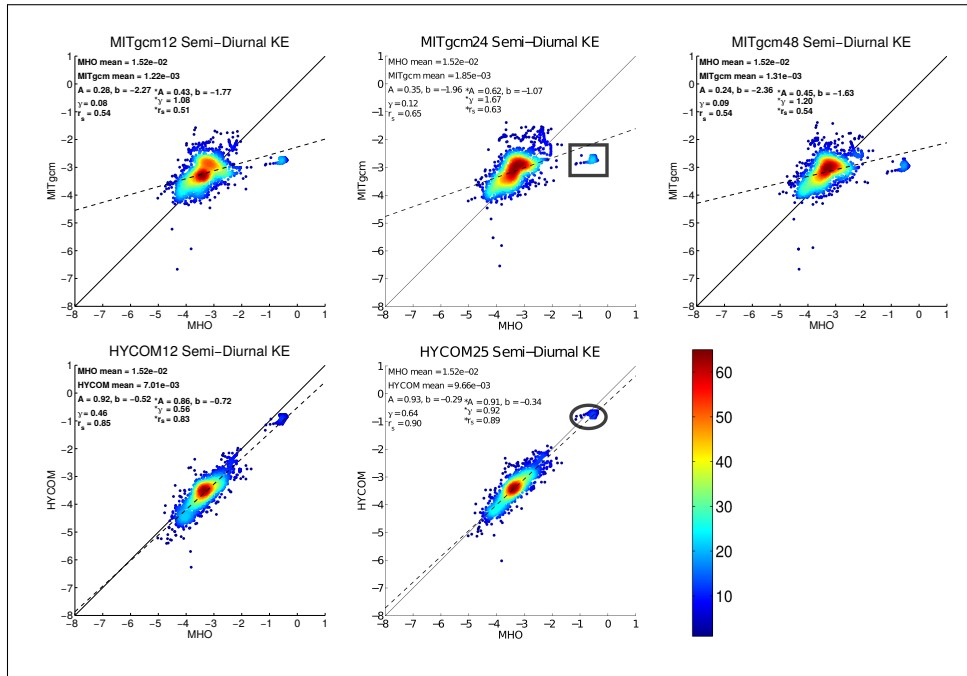
426 **Figure 6.** A point-to-point comparison of supertidal temperature variance [$\log_{10}(^{\circ}C^2)$] between model
 427 and MHO. The top panels show MITgcm12, MITgcm24, and MITgcm48, while the bottom panels show HY-
 428 COM12 and HYCOM25. Population density is given by color, with the most tightly grouped data shown in
 429 red, and the sparsest data in blue. The one-to-one line is shown in solid black. Spatial means of MHO and
 430 model values, and statistical metrics A , b , γ , and r_s between each model and the MHO data, are printed on the
 431 upper left of each subplot.

432 3.3 Semi-Diurnal Band

433 Scatterplots of HYCOM and MITgcm semi-diurnal KE values against MHO values
 434 are shown in Figure 7. From visual inspection, all three MITgcm runs appear to be biased
 435 high compared to the MHO values, except for at a small number of high-velocity locations
 436 in the MHO, where, as in the supertidal KE comparison, MITgcm is biased much too low.

437 As a result, γ values computed including these high-velocity values are lower for MITgcm
438 than for HYCOM, and the linear regressions are similarly skewed.

439 The high-velocity locations, representing energies around ($10^{-1}\text{m}^2\text{s}^{-2}$) in the MHO
440 data, have been highlighted in the HYCOM25 and MITgcm24 panels of Figure 7 as was
441 done in Figure 5. The high-velocity points are clearly visible in the model-data compar-
442 isons across all resolutions. These instrument locations correspond to the circle labeled
443 “D” in Figure 1. At this location (in the Strait of Gibraltar), MITgcm is under-energetic
444 for all three resolutions, while HYCOM lies much closer to observations. When these in-
445 struments are removed, the resulting γ values, given in Figure 7, reveal an overall over-
446 energetic bias in the great bulk of the MITgcm locations. In contrast, both with and with-
447 out the inclusion of the high-velocity locations, the HYCOM γ values are close to, but
448 less than, one. The linear regressions in HYCOM are much closer to one-to-one than in
449 MITgcm, even with the omission of the points discussed above. It is likely that the to-
450 pographic wave drag employed in HYCOM plays an important role in damping HYCOM
451 tidal KE to realistic levels. This would be consistent with discussions in *Ansong et al.*
452 [2015], who showed that an extra damping, such as parameterized topographic wave drag,
453 is needed to make modeled internal tide SSH signatures agree with altimeter observations
454 (see *Buijsman et al.* [2016] for related discussions on the impact of wave drag on tidal en-
455 ergetics). MITgcm does not employ a wave drag, which likely explains its overly large KE
456 outside of special regimes such as the Strait of Gibraltar. More analyses of the MITgcm
457 results are currently underway in order to further test this hypothesis of damping sensitiv-
458 ity. The high sensitivity of the comparison metrics to outliers, particularly in the case of
459 the linear regression coefficient and γ , illustrates some of the difficulties inherent in bulk
460 model-data comparisons such as this. Finally, we note that, both with and without the in-
461 clusion of the high-velocity values, the HYCOM spatial correlation values r_s are always
462 higher than the MITgcm r_s values.



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Figure 7. As in Figure 5 but for semi-diurnal KE.

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Figure 8 shows scatterplots of HYCOM and MITgcm semi-diurnal temperature variance values against MHO values. Visually, there is more improvement between MITgcm12 and MITgcm24, with the values becoming more tightly clustered around the one-to-one line, than there is between MITgcm24 and MITgcm48, as is evident in both the linear regression and correlation coefficient. Both HYCOM12 and HYCOM25 display lower variance levels than their MITgcm counterparts, although HYCOM25 is nearly centered on the one-to-one line. The ratios of the means, γ , are closer to one for MITgcm than for HYCOM. As in all previous comparisons, the correlation coefficients r_s are higher in HYCOM than in MITgcm, reflecting the tighter correlations seen in the scatterplots.

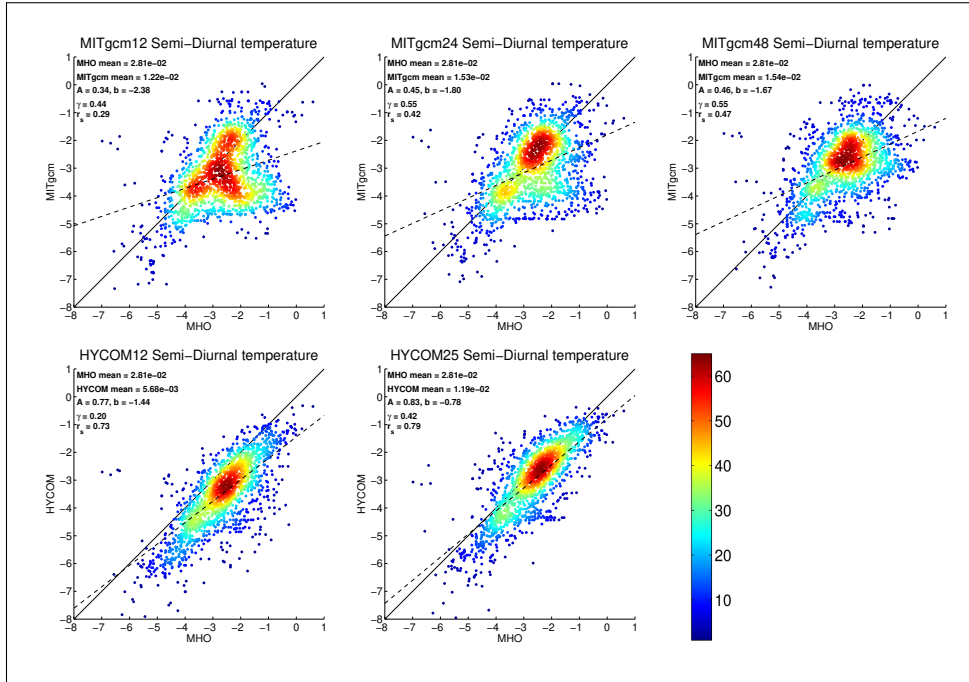


Figure 8. As in Figure 6 but for semi-diurnal temperature variance.

3.4 Diurnal Band

Figure 9 shows scatterplots of HYCOM and MITgcm diurnal KE values against MHO values. Interestingly, it visually appears as if MITgcm24 is better performing than both MITgcm12, and MITgcm48, and this is confirmed by the statistical metrics, A , b , γ , and r_s . Because extreme values tend to dominate the statistics, as discussed below, this result should be treated with some amount of skepticism. The majority of locations visually display a similar amount of scatter in MITgcm and HYCOM here compared with earlier plots, although outliers for high-velocity MHO values are again more apparent in MITgcm. The HYCOM results suggest a slight weak bias for the great bulk of values. This is likely due to over-damping of diurnal tides due to the fact that wave drag in HYCOM is optimally tuned for semidiurnal rather than diurnal tides, as discussed in *Timko et al.* [2013]; see also discussions in *Skiba et al.* [2013], a study that focused solely on diurnal tides and that employed a wave drag optimally tuned for diurnal tides. The optimal wave drag for diurnal tides is found to be weaker than the optimal drag for semidiurnal tides.

Again, there are several distinct groups of locations representing energies beyond ($10^{-2} \text{m}^2 \text{s}^{-2}$) in the MHO data that are anomalously under-energetic in MITgcm. The in-

490 instrument locations associated with these anomalous values correspond to the circles la-
 491 beled “A”, “B”, “C”, “D”, as well as instruments located in regions near the coastal North
 492 West Atlantic, West Pacific, near Taiwan, East Pacific off the Californian coast, the coastal
 493 Antarctic, the Ross sea, and off the eastern extent of the Weddell sea. These locations are
 494 contained in the large rectangle in the central top panel of Figure 9, and corresponding lo-
 495 cations have been covered with a circle and an ellipse in the HYCOM25 subplot of Figure
 496 9. At these locations, MITgcm is systematically under-energetic, while HYCOM appears
 497 to do better in its representation of observations, many of which are located in marginal
 498 seas or continental margins. Additionally, an anomalous group of points representing en-
 499 ergies beyond ($10^{-2} \text{m}^2 \text{s}^{-2}$) can be seen in the MITgcm24 output (and analogous clusters
 500 in both MITgcm12 and MITgcm48). These locations are marked with a smaller rectan-
 501 gle in the central top panel of Figure 9. When the instruments enclosed in the rectangles
 502 shown in Figure 9 (central top panel) are removed, the statistical metrics change slightly
 503 for HYCOM12 and HYCOM25 and MITgcm24, while the changes are more pronounced
 504 for MITgcm12, and MITgcm48. This illustrates that even within a single model, locations
 505 of outliers can vary between simulations having different resolutions. Once again, with
 506 and without the anomalous high-velocity points included, the spatial correlation coeffi-
 507 cients r_s are higher for HYCOM than for MITgcm.

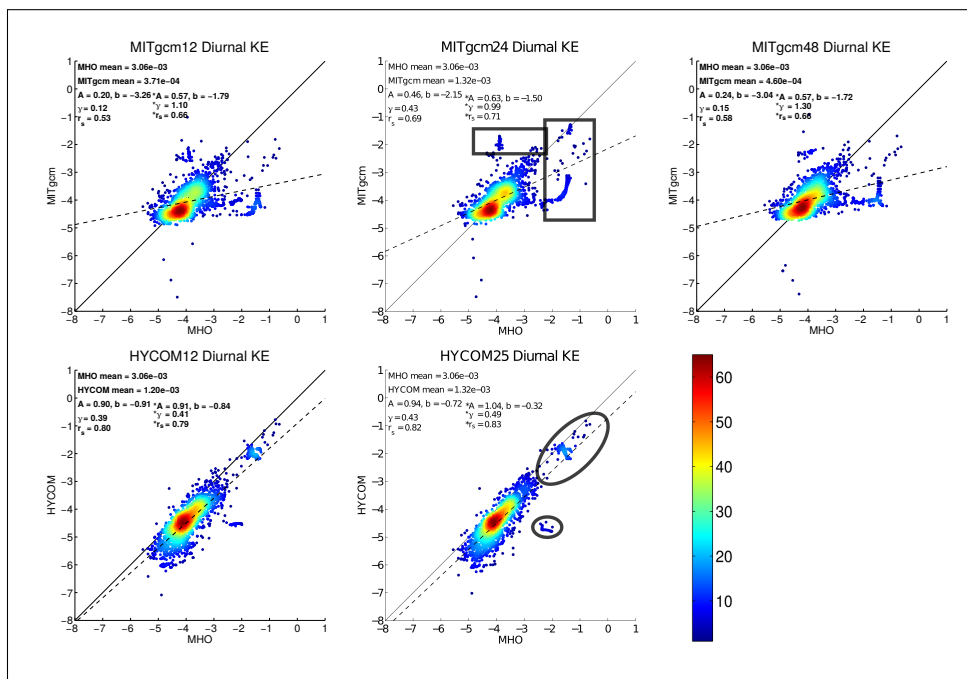
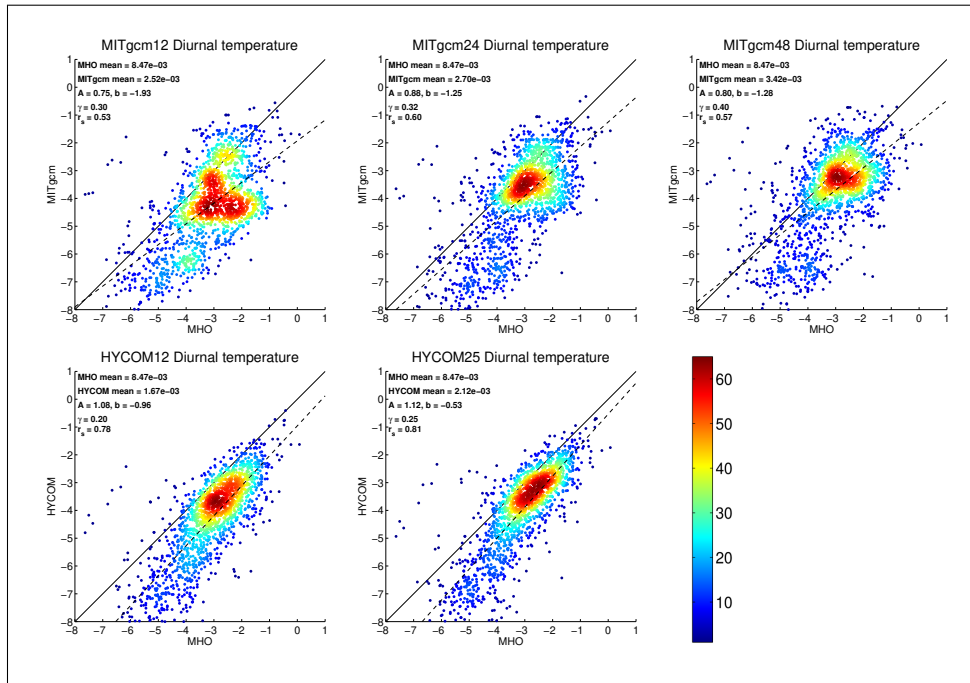


Figure 9. As in Figure 5 but for diurnal KE.

509 Figure 10 shows scatterplots of HYCOM and MITgcm diurnal temperature variance
 510 values against MHO values. Both HYCOM12 and HYCOM25 are weaker on average than
 511 MITgcm, as is seen in lower values of γ , and as discussed in the previous section, we sug-
 512 gest that this may be due to over-damping of diurnal tides by the wave drag employed in
 513 HYCOM. Interestingly, while HYCOM and MITgcm as a whole have opposite linear re-
 514 gression trends, both models have more or less equal performance, with each improving
 515 slightly with increased resolution. Once again, HYCOM has higher correlation values r_s
 516 than MITgcm.

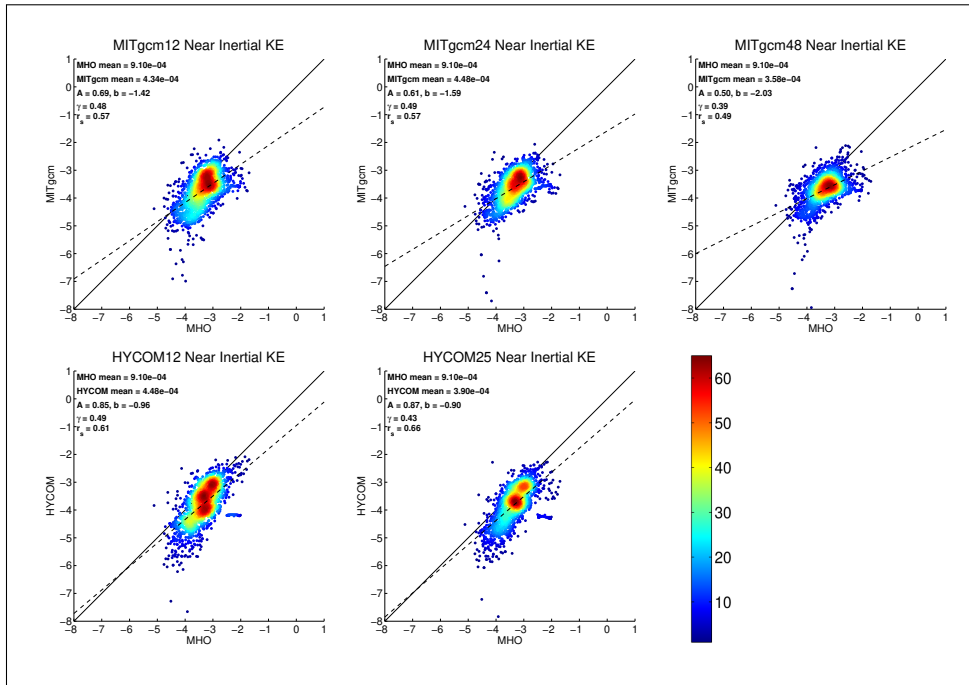


517 **Figure 10.** As in Figure 6 but for diurnal temperature variance.

518 3.5 Near-Inertial Band

519 As near-inertial motions do not, to first order, do not have associated vertical dis-
 520 placements, we do not compare temperature variances in the near-inertial band. Scatter-
 521 plots of HYCOM and MITgcm near-inertial KE values against MHO values are displayed
 522 in Figure 11. Some scatter between models and observations is expected, as the mooring
 523 data spans many different years, and the atmospheric forcing such as storms that generate
 524 near-inertial motions will inherently vary from year to year. The performance of MITgcm
 525 and HYCOM is fairly similar in this band as seen in all four metrics (A , b , γ , and r_s), and

526 is also not greatly sensitive to horizontal resolution, consistent with the fairly large hori-
 527 zontal scales of near-inertial waves (e.g., *Simmons and Alford* [2012]). As in all bands, the
 528 HYCOM spatial correlation values r_s are higher than in MITgcm. However, in this band
 529 the HYCOM r_s values are only slightly higher, consistent with the visual appearance of
 530 Figure 11.

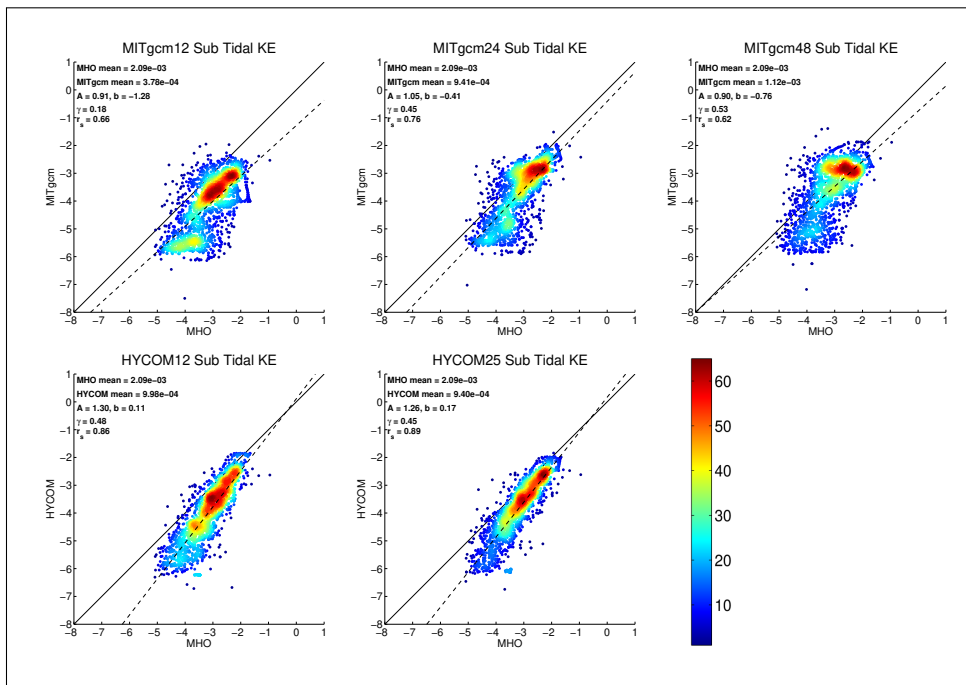


531 **Figure 11.** As in Figure 5 but for near-inertial KE.

532 3.6 Subtidal Frequency Band

533 The subtidal band covers frequencies from 0.7 cpd, a little less than once per day, to
 534 one cycle every 10 days, and is the frequency band most likely to contain sub-mesoscale
 535 eddy motions, which have shorter timescales and smaller spatial extent than mesoscale
 536 motions [*Su et al.*, 2018]. One might expect an increase in model resolution to enhance
 537 the dynamics within these time scales, as higher spatial resolutions precipitate frontal
 538 instabilities and other submesoscale motions [*Capet et al.*, 2008]. Figure 12 shows scat-
 539 terplots of HYCOM and MITgcm subtidal kinetic energy values against MHO values.
 540 From visual inspection, HYCOM12 and HYCOM25 behave similarly, with both runs being
 541 about a factor of 2 lower than the MHO mean ($\gamma = 0.48, 0.45$ for HYCOM12 and HY-
 542 COM25 respectively), and the linear regressions suggest that this bias is lowest at lower

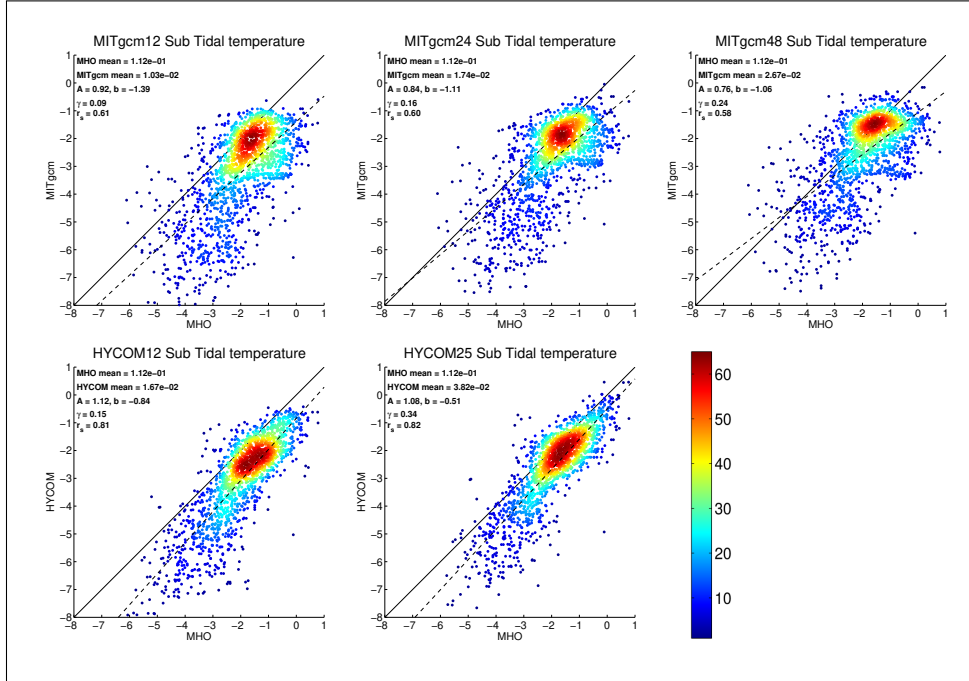
543 energies. MITgcm however shows a steady increase in variance as resolution is increased
 544 ($\gamma = 0.18, 0.45, 0.53$ for MITgcm12, MITgcm24, and MITgcm48 respectively), and the
 545 linear regressions have slopes slightly closer to one. Although MITgcm48 has the high-
 546 est variance, it is only slightly larger than the MITgcm24 and HYCOM variances. Thus
 547 we cannot conclude with any confidence that better resolution of the submesoscale is re-
 548 sponsible for the changes seen with higher resolution in Figure 12. As in other frequency
 549 bands, the HYCOM scatterplots have a tighter visual appearance, and accordingly, some-
 550 what higher r_s values.



551 **Figure 12.** As in Figure 5 but for subtidal KE.

552 The subtidal temperature variance scatterplots are displayed in Figure 13. Visu-
 553 ally, HYCOM12 and HYCOM25 are similar, with both runs having lower means than the
 554 MHO. There is marked improvement in some statistics with $\gamma = 0.15, 0.34$ for HYCOM12
 555 and HYCOM25 respectively. Similarly, MITgcm also increases temperature variance as
 556 resolution increases ($\gamma = 0.09, 0.16, 0.24$ for MITgcm12, MITgcm24, and MITgcm48 re-
 557 spectively). The linear regressions for subtidal temperature variance are very similar to
 558 KE within the same frequency band. However, as the submesoscale is just beginning to
 559 become evident at the ~ 2 km grid spacing in MITgcm48 [Capet *et al.*, 2008], it is again

560 difficult to be confident that resolution of the submesoscales is responsible for the changes
 561 with model horizontal resolution seen in Figure 13. Once again the r_S values for subtidal
 562 temperature variance are higher for HYCOM than for MITgcm.



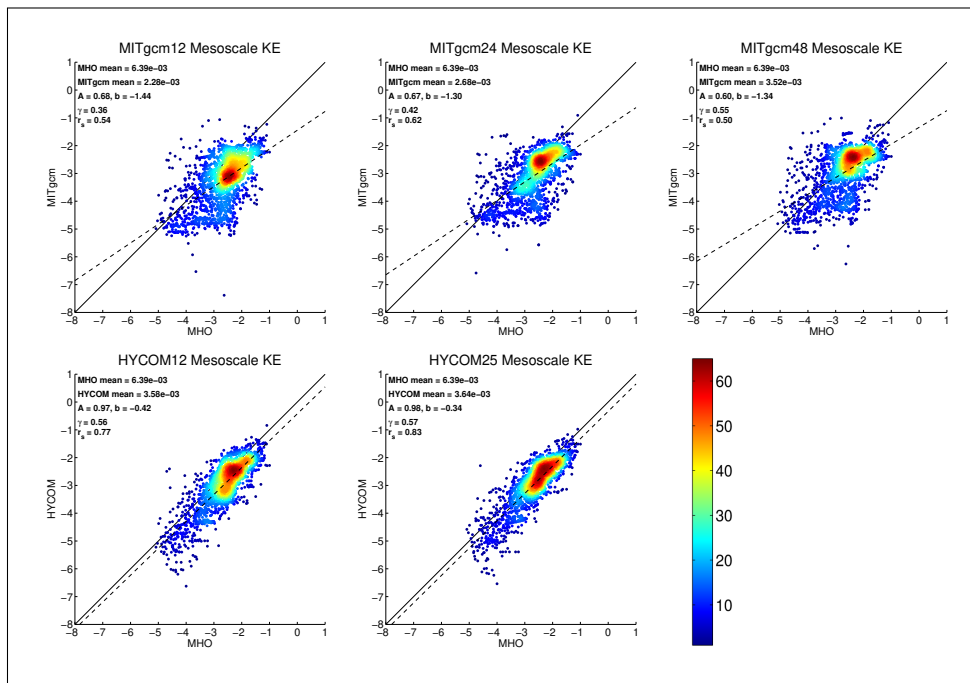
563 **Figure 13.** As in Figure 6 but for subtidal temperature variance.

564 3.7 Mesoscale Frequency Band

565 The mesoscale is the final, and lowest frequency, band that we examine in the model
 566 vs. MHO comparisons. Figure 14 shows scatterplots of HYCOM and MITgcm mesoscale
 567 kinetic energy values against MHO values. From visual inspection, MITgcm12 is biased
 568 lowest compared to MHO, while HYCOM25 is the closest to the one-to one line as con-
 569 firmed by the linear regression, although still somewhat lower than observations. Within
 570 both models, higher resolution runs generally contain more KE variance, ($\gamma = 0.36, 0.42, 0.55$
 571 for MITgcm12, MITgcm24, and MITgcm48 respectively, and $\gamma = 0.56, 0.57$ for HY-
 572 COM12 and HYCOM25 respectively). HYCOM is more correlated than MITgcm, with
 573 $r_s = 0.77, 0.83$ for HYCOM12 and HYCOM25 versus $r_s = 0.54, 0.62, 0.50$ for MITgcm12,
 574 MITgcm24, and MITgcm48.

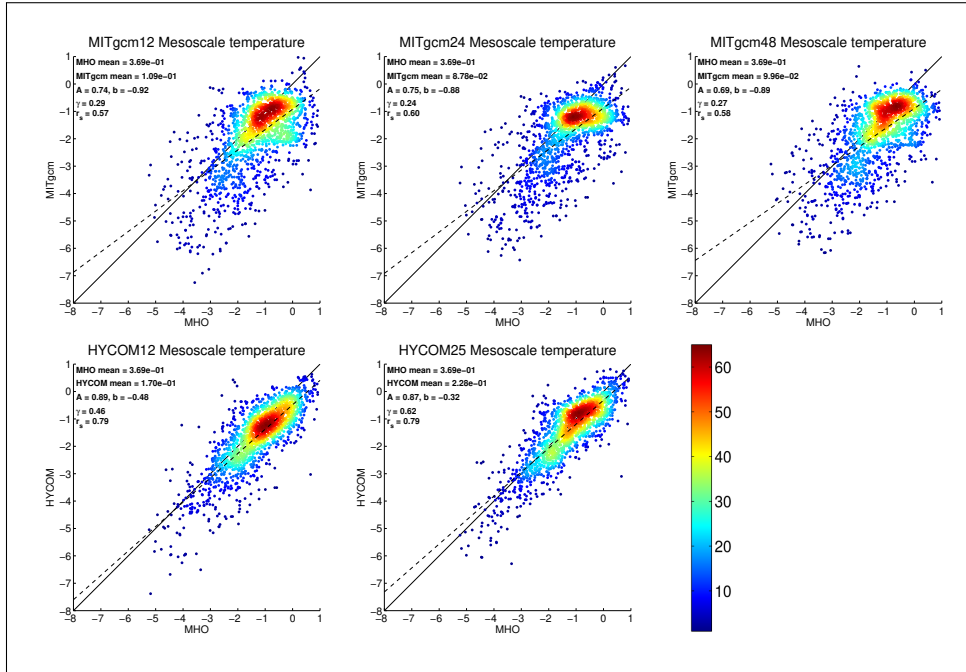
575 The mesoscale temperature variance scatterplots are given in Figure 15. As in all
 576 other model-MHO comparisons presented here, the HYCOM scatterplots have a visually

577 tighter appearance, and accordingly are associated with larger values of the correlation
 578 coefficient r_s . The ratios of the means, γ , are also closer to 1 in HYCOM. The γ values
 579 move closer to one when resolution is increased in HYCOM, but less so for the MITgcm
 580 simulations. As with mesoscale KE, the HYCOM mesoscale temperature variances are
 581 larger than the MITgcm variances, and are closer to observations, as measured by the γ
 582 values. The HYCOM γ , and r_s values computed from the HYCOM runs used here are
 583 similar to those computed in the analysis of low-frequency eddy available potential energy
 584 in the older HYCOM runs used in *Luecke et al.* [2017].



585

Figure 14. As in Figure 5 but for mesoscale KE.



586

Figure 15. As in Figure 6 but for mesoscale temperature variance.

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3.8 Summary of Model-MHO Comparisons

588

Our model-MHO comparisons can be summarized with a few key points:

589

- Increasing the horizontal resolution of the ocean models makes a larger difference in some frequency bands than it does in others, with the supertidal IGW band seeing the largest effects.
- MITgcm48 has the most realistic supertidal IGW variance levels, likely in part due to its higher vertical and horizontal resolution and perhaps due in part to an overly energetic internal tide band.
- On the other hand, in some continental margin/marginal sea regions, MITgcm has too little KE in the high-frequency bands (diurnal, semidiurnal, and supertidal), whereas HYCOM lies closer to observations in such areas.
- HYCOM has higher levels of mesoscale KE and temperature variance than MITgcm.
- Of all the frequency bands examined, the near-inertial band is most similar across the five simulations examined here; the lack of sensitivity to model resolution may

601

602 reflect the relatively large horizontal scales of low-mode near-inertial motions.
603 [*Simmons and Alford, 2012*].

- 604 • HYCOM displays higher spatial correlations r_s than MITgcm across all frequency
605 bands, with the near-inertial band showing the smallest differences between the
606 simulations.

607 The comparison of the MITgcm and HYCOM simulations with the MHO hints at the im-
608 portance of wave drag, in that the internal tide kinetic energy in the MITgcm simulations
609 sits higher than that of the moorings, while HYCOM is closer. This tendency for MIT-
610 gcm to have overly energetic internal tides is seen more clearly in comparisons with ob-
611 servations that have a denser spatial sampling (hence better spatial statistics) than our
612 current meter archive—for instance, in the comparison with the global drifter dataset [*Yu*
613 *et al.*, 2019]; in-preparation work led by Jonathan Brasch shows that the HYCOM inter-
614 nal tide kinetic energies lie closer to the drifters, primarily because of the inclusion of
615 wave drag. An in-preparation paper by Joseph Ansong likewise shows that the internal
616 tides in several models that are run without wave drag are overly energetic when compared
617 to along-track altimetry; a conclusion that is consistent with *Ansong et al.* [2015], who
618 considered HYCOM simulations with and without wave drag. The strength of the model-
619 MHO comparisons lies in the ability to compare different frequency bands, which separate
620 different classes of oceanic motions. The weakness of the model-MHO comparisons lies
621 in the relatively sparse spatial distribution of the $\sim 2,000$ MHO instruments, which im-
622 plies that statistical metrics are sensitive to a small number of extreme values, often tak-
623 ing place in continental margin/marginal sea locations. In the next section we compare
624 surface ocean geostrophic KE fields in the models to AVISO, which provides a spatially
625 dense global product. However, this product focuses on only one class of motions (low-
626 frequency mesoscale flows).

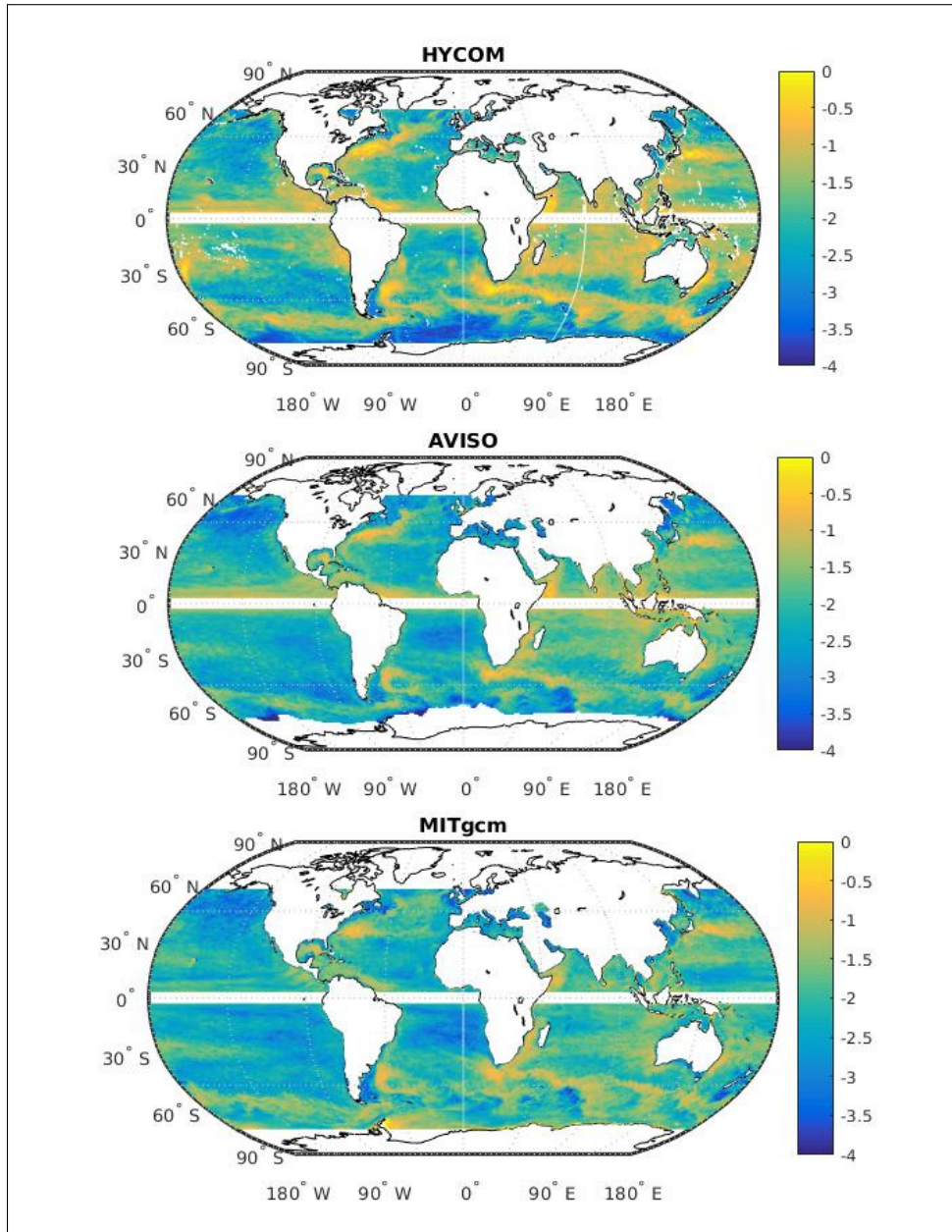
627 **4 Global Model vs. AVISO Comparisons of Surface Ocean Geostrophic Eddy Ki-** 628 **netic Energy**

629 Prompted by differences between MITgcm and HYCOM in our MHO compari-
630 son, we present a global comparison of surface ocean geostrophic eddy KE between HY-
631 COM, the AVISO satellite altimeter product, and MITgcm, shown in Figure 16. The spa-
632 tial structure in both MITgcm48 (bottom panel) and HYCOM25 (top panel) matches the
633 structure seen in AVISO (center) reasonably well. However, there are several regions of

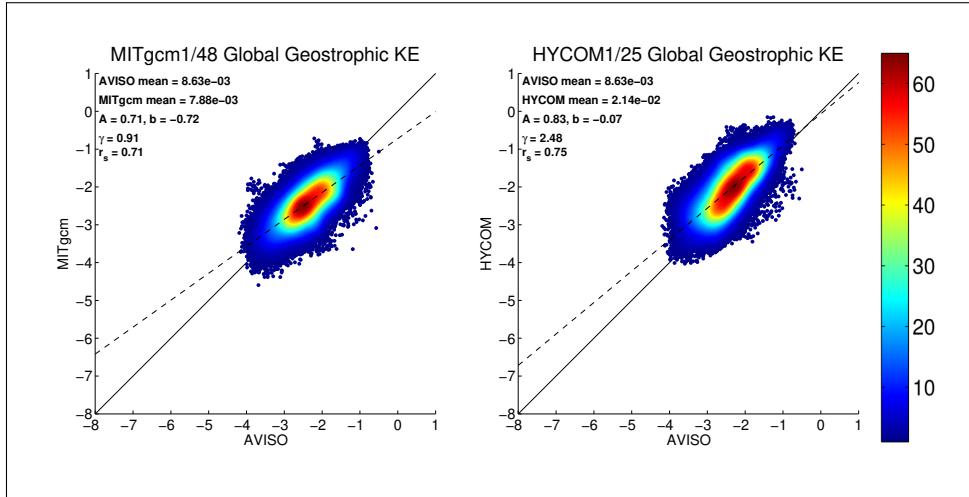
634 poor model-AVISO comparison. For instance, HYCOM25 appears to have more energy
635 than AVISO in the South Atlantic, Southern Indian Ocean, South Eastern Pacific, many
636 tropical regions, and the Antarctic Circumpolar Current, while MITgcm48 appears closer
637 to AVISO in most regions. Conversely HYCOM25 represents the Gulf Stream more accu-
638 rately than MITgcm48.

639 As was done in our comparisons at MHO locations, we can also examine the point-
640 to-point comparisons between the models and AVISO. Figure 17 shows the point-to-point
641 surface ocean geostrophic eddy KE between MITgcm48 and AVISO (left), and between
642 HYCOM25 and AVISO (right), taken from the data shown in Figure 16. Because of con-
643 cerns about the effect of sea ice on the comparisons, points poleward of 55 degrees have
644 been masked out of the comparison. In addition we eliminate all AVISO for which one
645 or more points in the local $1/4^\circ$ by $1/4^\circ$ neighborhood over which we average the model
646 values, are land points. The elimination of points near land or sea ice results in a more re-
647 stricted set of points than was used in *Qiu et al.* [2018]. Globally, the linear regression A
648 is higher in HYCOM25 than in MITgcm48, indicating a HYCOM bias toward higher ener-
649 gies. Examination of the correlation r_s reveals that while both models correlate fairly well,
650 HYCOM25 has a slightly larger r_s value than MITgcm, consistent with the results of our
651 MHO analysis.

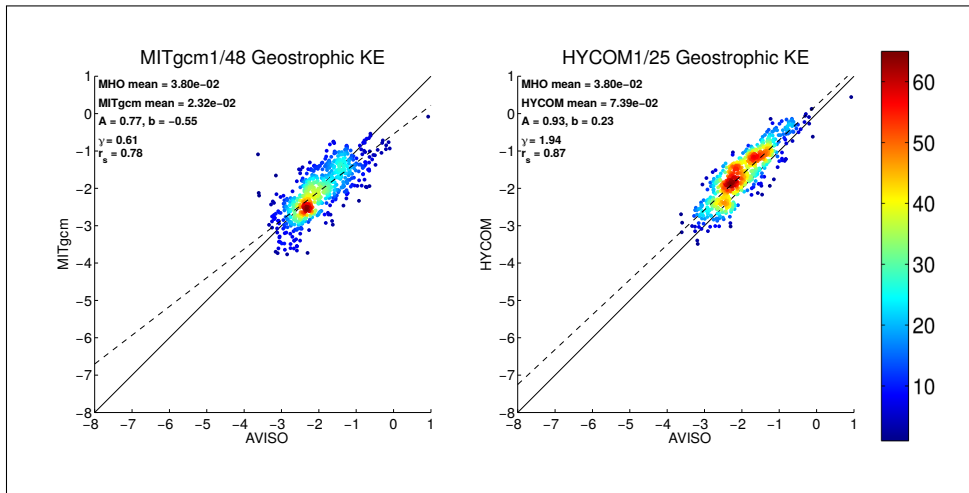
652 Finally, as a test of the sampling bias in the MHO comparisons, we show, in Fig-
653 ure 18, scatterplots of the surface ocean geostrophic eddy KE between MITgcm48 and
654 AVISO (left), and between HYCOM25 and AVISO (right) at the MHO locations used in
655 this work. (For this plot, in order to duplicate the MHO analysis as closely as possible, at
656 MHO locations with instruments at multiple depths, the model-AVISO pairs are included
657 multiple times, once for each MHO instrument). The linear regression A is higher in HY-
658 COM25 than in MITgcm48, consistent with our other results indicating higher energies
659 in the HYCOM surface ocean geostrophic eddy KE field. As measured by γ values, at
660 MHO locations MITgcm25 slightly under-predicts AVISO eddy KE, while HYCOM25
661 over-predicts the eddy KE by a factor of about two. Consistent with previous results of the
662 MHO comparison, the correlation r_s in HYCOM25 is higher than MITgcm48. The results
663 of Figure 18 indicate once again that the HYCOM simulations achieve a higher spatial
664 correlation with respect to observational products than the MITgcm simulations.



665 **Figure 16.** Global surface ocean geostrophic eddy kinetic energy ($\log_{10}(\text{m}^2\text{s}^{-2})$) in HYCOM25 (top),
666 AVISO (center), and MITgcm48 (bottom). Maps are given on a 0.25 degree grid.



667 **Figure 17.** As in Figure 5 but for surface ocean geostrophic eddy KE between MITgcm48 and AVISO
 668 (left), and between MITgcm48 and HYCOM25 and AVISO (right), over the global values taken from the maps of Figure 16.
 669 Locations poleward of 55 degrees and near-land locations (see text for details), have been omitted.



670 **Figure 18.** As in Figure 17 but for surface ocean geostrophic eddy KE between MITgcm48 and AVISO
 671 (left), and between MITgcm48 and HYCOM25 and AVISO (right), at the locations of the MHO intruments.

672 5 Summary and discussion

673 We have assessed the frequency content of temperature variance and kinetic energy
 674 in multiple resolutions of two global ocean general circulation models with embedded
 675 tides (MITgcm 1/12, 1/24 and 1/48 degree, and HYCOM 1/12 and 1/25 degree) using a
 676 database of moored historical observations (MHO). We also compared geostrophic sur-

677 face ocean eddy kinetic energy in the models and in the AVISO satellite altimeter product.
678 The AVISO comparisons can only be done for one class of motions, mesoscale currents
679 and eddies, but offer much denser spatial coverage over the globe. The MHO comparisons
680 allow a separation into different frequency bands, corresponding to different classes of
681 oceanic motions, but for some bands the statistical metrics are highly sensitive to behav-
682 iors at a small number of locations, often in continental shelves and marginal seas. The
683 impact of horizontal model resolution on KE and temperature variance differs between
684 frequency bands. For instance, within the near-inertial band, model performance was not
685 dramatically improved by an increase in resolution. However in the supertidal, subtidal,
686 and mesoscale bands (and to a lesser extent semi-diurnal and diurnal bands), models more
687 closely match the observations as resolution is increased. Particularly within the supertidal
688 band, increasing the horizontal model resolution plays an important role in transferring
689 energy into the IGW continuum.

690 Both HYCOM and MITgcm simulations show good and poor agreement with obser-
691 vations in some respects. For instance the HYCOM simulations have higher spatial cor-
692 relations with the MHO observations than the MITgcm simulations do, at all resolutions
693 and frequency bands. The HYCOM simulations also show a higher correlation with ob-
694 servations in the AVISO surface ocean geostrophic eddy KE comparison. Both the MHO
695 and AVISO comparisons display a tendency for the MITgcm simulations to perform poorly
696 in some continental margin/marginal sea locations, where the high-frequency motions in
697 the MITgcm simulations are often under-energetic relative to observations. The HYCOM
698 simulations do not display this tendency for under-energetic high-frequency motions in
699 near-land locations. The MITgcm simulations, particularly the highest resolution MITgcm
700 simulations, have more realistic energy levels in the IGW continuum. Interestingly, how-
701 ever, the HYCOM simulations have a more energetic mesoscale band than the MITgcm
702 simulations.

703 These important caveats aside, in some respects the highest resolution runs of both
704 MITgcm and HYCOM perform well compared to both observational datasets. These new
705 models capture some of the complexities of oceanic variability from mesoscale eddies,
706 which have timescales of 100 days or so, down to the supertidal continuum, which has
707 timescales of a few to several hours. As high-resolution GCMs that partially resolve an
708 internal gravity wave continuum become increasingly used both to study dynamics, and

709 as a tool for informing observational oceanography, model-data validation across a wide
710 range of frequencies becomes increasingly important.

711 We end by clarifying the obvious, that this paper has reported on particular simu-
712 lations of HYCOM and MITgcm. Our conclusions above apply to these simulations only,
713 and are not meant to apply generally to these two models. It is likely that many of the
714 statements above would change if new simulations of either model with improved param-
715 eterizations, numerics, resolution, etc. were to be developed. This being said, we are con-
716 fident that many of the findings regarding resolutions and parameterizations are broadly
717 applicable. We speculate that the higher spatial correlations of the HYCOM simulations
718 with respect to observations may be due, at least in part, to the long history of tuning HY-
719 COM simulations to be accurate enough for Navy operational purposes. Given time and
720 resources, it is likely that the MITgcm simulations could be brought to a similar state.
721 Similarly, if HYCOM simulations were to be run at higher resolution, they would likely
722 develop a supertidal IGW continuum spectrum with energy levels closer to observations,
723 as in the higher resolution MITgcm simulations reported on here. Introduction of a topo-
724 graphic wave drag into the MITgcm simulations could lead to a better agreement of the
725 modeled internal tide with observations, as has been shown using HYCOM [Ansong *et al.*,
726 2015]. Alternatively, more advanced parameterizations of internal wave energy loss might
727 be developed, that allow us to move away from using topographic wave drag. Develop-
728 ing better mixing parameterizations is the subject of much present and ongoing research,
729 e.g. MacKinnon *et al.* [2017]. The relatively poor performance of the particular MITgcm
730 simulations used here in continental margin/marginal sea regions is interesting. We do not
731 have an explanation for it, and we again caution against drawing general conclusions about
732 MITgcm performance in such regions from the particular solutions examined here.

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750 NRL-JA-7320-2019-4447 has been approved for public release. Distribution is unlimited.
751 Hourly output from the 1/12, 1/24, and 1/48-degree MITgcm simulations is available on
752 the ECCO data portal: <https://data.nas.nasa.gov/ecco/>. A subset of MHO time
753 series that contains the observations used in this paper, along with MITgcm and HYCOM
754 profiles at MHO locations is available on the University of Michigan Library archival
755 system: <https://doi.org/10.7302/dbfp-s644>. One year of sea surface height, and
756 surface velocity fields, from a 1/25th degree HYCOM simulation, can also be accessed
757 through the OSiRIS system at the University of Michigan. Co-authors Arbic and Shriver
758 can be contacted for details on how to access the OSiRIS system.

759 References

- 760 Ansong, J. K., B. K. Arbic, M. C. Buijsman, J. G. Richman, J. F. Shriver, and A. J. Wall-
761 craft (2015), Indirect evidence for substantial damping of low-mode internal tides in
762 the open ocean, *Journal of Geophysical Research: Oceans*, 120(9), 6057–6071, doi:
763 10.1002/2015JC010998.
- 764 Ansong, J. K., B. K. Arbic, M. H. Alford, M. C. Buijsman, J. F. Shriver, Z. Zhao, J. G.
765 Richman, H. L. Simmons, P. G. Timko, A. J. Wallcraft, and L. Zamudio (2017), Semid-
766 iurnal internal tide energy fluxes and their variability in a global ocean model and
767 moored observations, *Journal of Geophysical Research: Oceans*, 122(3), 1882–1900,
768 doi:10.1002/2016JC012184.
- 769 Arbic, B. K., A. J. Wallcraft, and E. J. Metzger (2010), Concurrent simulation of the ed-
770 dying general circulation and tides in a global ocean model, *Ocean Modelling*, 32, 175–
771 187, doi:doi:10.1016/j.ocemod.2010.01.007.

- 772 Arbic, B. K., J. G. Richman, J. F. Shriver, P. G. Timko, E. J. Metzger, and A. J. Wallcraft
773 (2012), Global modeling of internal tides within an eddying ocean general circulation
774 model, *Oceanography*, 25, 20–29, doi:<https://doi.org/10.5670/oceanog.2012.38>.
- 775 Arbic, B. K., M. H. Alford, J. K. Ansong, M. C. Buijsman, R. B. Ciotti, J. T. Farrar,
776 R. W. Hallberg, C. E. Henze, C. N. Hill, C. A. Luecke, D. Menemenlis, E. J. Metzger,
777 M. Müller, A. D. Nelson, B. C. Nelson, H. E. Ngodock, R. M. Ponte, J. G. Richman,
778 A. C. Savage, R. B. Scott, J. F. Shriver, H. L. Simmons, I. Souopgui, P. G. Timko, A. J.
779 Wallcraft, L. Zamudio, and Z. Zhao (2018), *New Frontiers in Operational Oceanog-*
780 *raphy*, chap. 13: A primer on global internal tide and internal gravity wave contin-
781 uum modeling in HYCOM and MITgcm, pp. 307–392, GODAE OceanView, doi:
782 doi:10.17125/gov2018.ch13.
- 783 Buijsman, M., B. Arbic, J. Green, R. Helber, J. Richman, J. Shriver, P. Timko, and
784 A. Wallcraft (2015), Optimizing internal wave drag in a forward barotropic model with
785 semidiurnal tides, *Ocean Modelling*, 85, 42 – 55, doi:[https://doi.org/10.1016/j.ocemod.](https://doi.org/10.1016/j.ocemod.2014.11.003)
786 2014.11.003.
- 787 Buijsman, M. C., J. K. Ansong, B. K. Arbic, J. G. Richman, J. F. Shriver, P. G. Timko,
788 A. J. Wallcraft, C. B. Whalen, and Z. Zhao (2016), Impact of parameterized internal
789 wave drag on the semidiurnal energy balance in a global ocean circulation model, *Jour-*
790 *nal of Physical Oceanography*, 46(5), 1399–1419, doi:10.1175/JPO-D-15-0074.1.
- 791 Capet, X., J. C. McWilliams, M. J. Molemaker, and A. F. Shchepetkin (2008), Mesoscale
792 to submesoscale transition in the california current system. part iii: Energy balance and
793 flux, *Journal of Physical Oceanography*, 38(10), 2256–2269, doi:10.1175/2008JPO3810.
794 1.
- 795 Cartwright, D. E. (1999), *Tides : a scientific history*, 210 pp., Cambridge University Press
796 Cambridge ; New York.
- 797 Chassignet, E. P., and X. Xu (2017), Impact of horizontal resolution (1/12 to 1/50) on gulf
798 stream separation, penetration, and variability, *Journal of Physical Oceanography*, 47(8),
799 1999–2021, doi:10.1175/JPO-D-17-0031.1.
- 800 Chassignet, E. P., H. E. Hurlburt, E. J. Metzger, O. M. Smedstad, J. A. Cummings, G. R.
801 Halliwell, R. Bleck, R. Baraille, A. J. Wallcraft, C. Lozano, H. L. Tolman, A. Sriniv-
802 asan, S. Hankin, P. Cornillon, R. Weisberg, A. Barth, R. He, F. Werner, and J. Wilkin
803 (2009), US GODAE: Global ocean prediction with the HYbrid Coordinate Ocean
804 Model (HYCOM), *Oceanography*, 22(2), 64–75, doi:[dx.doi.org/10.5670/oceanog.2009](https://doi.org/10.5670/oceanog.2009).

805 39.

- 806 Dee, D. P., S. M. Uppala, A. J. Simmons, P. Berrisford, P. Poli, S. Kobayashi, U. Andrae,
807 M. A. Balmaseda, G. Balsamo, P. Bauer, P. Bechtold, A. C. M. Beljaars, L. van de
808 Berg, J. Bidlot, N. Bormann, C. Delsol, R. Dragani, M. Fuentes, A. J. Geer, L. Haim-
809 berger, S. B. Healy, H. Hersbach, E. V. Hólm, L. Isaksen, P. Kállberg, M. Köh-
810 ler, M. Matricardi, A. P. McNally, B. M. Monge-Sanz, J.-J. Morcrette, B.-K. Park,
811 C. Peubey, P. de Rosnay, C. Tavolato, J.-N. Thépaut, and F. Vitart (2011), The ERA-
812 interim reanalysis: configuration and performance of the data assimilation system, *Quar-*
813 *terly Journal of the Royal Meteorological Society*, 137(656), 553–597, doi:10.1002/qj.
814 828.
- 815 Doherty, K. W., D. E. Frye, S. P. Liberatore, and J. M. Toole (1999), A moored profiling
816 instrument, *Journal of Atmospheric and Oceanic Technology*, 16, 1816–1829.
- 817 Ducet, N. P., Y. L. Traon, and G. Reverdin (2000), global high-resolution mapping of
818 ocean circulation from topex/poseidon and ers-1 and -2, *Journal of Geophysical Re-*
819 *search: Oceans*, 105(C8), 19,477–19,498, doi:10.1029/2000JC900063.
- 820 Egbert, G. D., and S. Y. Erofeeva (2002), Efficient inverse modeling of barotropic ocean
821 tides, *Journal of Atmospheric and Oceanic Technology*, 19(2), 183–204, doi:10.1175/
822 1520-0426(2002)019<0183:EIMOBO>2.0.CO;2.
- 823 Egbert, G. D., A. F. Bennett, and M. G. G. Foreman (1994), TOPEX/POSEIDON tides
824 estimated using a global inverse model, *Journal of Geophysical Research: Oceans*,
825 99(C12), 24,821–24,852, doi:10.1029/94JC01894.
- 826 Garrett, C., and E. Kunze (2007), Internal tide generation in the deep ocean, *Annual Re-*
827 *view of Fluid Mechanics*, 39(1), 57–87, doi:10.1146/annurev.fluid.39.050905.110227.
- 828 Hecht, W. M., and H. Hasumi (2008), *Ocean Modeling in an Eddying Regime*, vol. Geo-
829 physical Monograph; 177, American Geophysical Union, 2000 Florida Avenue N. W.,
830 Washington, D.C.
- 831 Hendershott, M. C. (1972), The effects of solid earth deformation on global ocean tides,
832 *Geophysical Journal of the Royal Astronomical Society*, 29(4), 389–402, doi:10.1111/j.
833 1365-246X.1972.tb06167.x.
- 834 Hogan, T. F., M. Liu, J. A. Ridout, M. S. Peng, T. R. Whitcomb, B. C. Ruston, C. A.
835 Reynolds, S. D. Eckermann, J. R. Moskaitis, N. L. Baker, P. McCormack, J. L. C.
836 Viner, J. G. McLay, M. K. Flatau, L. Xu, C. C., and S. W. Chang (2014), The navy
837 global environmental model, *Oceanography*, 27(3), 116–125.

- 838 Jakobsson, M., R. Macnab, L. Mayer, R. Anderson, M. Edwards, J. Hatzky, H. W.
839 Schenke, and P. Johnson (2008), An improved bathymetric portrayal of the arctic ocean:
840 Implications for ocean modeling and geological, geophysical and oceanographic analy-
841 ses, *Geophysical Research Letters*, 35(7), doi:10.1029/2008GL033520, 107602.
- 842 Jayne, S. R., and L. C. St. Laurent (2001), Parameterizing tidal dissipation over rough to-
843 pography, *Geophysical Research Letters*, 28(5), 811–814, doi:10.1029/2000GL012044.
- 844 Kokoska, S., and D. Zwillinger (2000), *Standard probability and statistics tables and for-*
845 *mulae.*, Boca Raton: Chapman & Hall-CRC.
- 846 Kunze, E. (2017), The internal-wave-driven meridional overturning circulation, *Journal of*
847 *Physical Oceanography*, 47(11), 2673–2689, doi:10.1175/JPO-D-16-0142.1.
- 848 Large, W., and S. Yeager (2004), *Diurnal to Decadal Global Forcing for Ocean and Sea-ice*
849 *Models: The Data Sets and Flux Climatologies*, NCAR technical notes, National Center
850 for Atmospheric Research.
- 851 Large, W. G., J. C. McWilliams, and S. C. Doney (1994), Oceanic vertical mixing: A re-
852 view and a model with a nonlocal boundary layer parameterization, *Reviews of Geo-*
853 *physics*, 32(4), 363–403, doi:10.1029/94RG01872.
- 854 Losch, M., D. Menemenlis, J.-M. Campin, P. Heimbach, and C. Hill (2010), On the for-
855 mulation of sea-ice models. Part 1: Effects of different solver implementations and
856 parameterizations, *Ocean Modelling*, 33(1), 129 – 144, doi:https://doi.org/10.1016/j.
857 ocemod.2009.12.008.
- 858 Luecke, C. A., B. K. Arbic, S. L. Bassette, J. G. Richman, J. F. Shriver, M. H. Alford,
859 O. M. Smedstad, P. G. Timko, D. S. Trossman, and A. J. Wallcraft (2017), The global
860 mesoscale eddy available potential energy field in models and observations, *Journal of*
861 *Geophysical Research: Oceans*, 122(11), 9126–9143, doi:10.1002/2017JC013136.
- 862 MacKinnon, J. A., Z. Zhao, C. B. Whalen, A. F. Waterhouse, D. S. Trossman, O. M.
863 Sun, L. C. St. Laurent, H. L. Simmons, K. Polzin, R. Pinkel, A. Pickering, N. J.
864 Norton, J. D. Nash, R. Musgrave, L. M. Merchant, A. V. Melet, B. Mater, S. Legg,
865 W. G. Large, E. Kunze, J. M. Klymak, M. Jochum, S. R. Jayne, R. W. Hallberg, S. M.
866 Griffies, S. Diggs, G. Danabasoglu, E. P. Chassignet, M. C. Buijsman, F. O. Bryan,
867 B. P. Briegleb, A. Barna, B. K. Arbic, J. K. Ansong, and M. H. Alford (2017), Climate
868 process team on internal wave-driven ocean mixing, *Bulletin of the American Meteorolo-*
869 *gical Society*, 98(11), 2429–2454, doi:10.1175/BAMS-D-16-0030.1.

- 870 Maltrud, M. E., and J. L. McClean (2005), An eddy resolving global 1/10 ocean simula-
871 tion, *ocean modelling*, 8(1-2), 31–54, doi:<https://doi.org/10.1016/j.ocemod.2003.12.001>.
- 872 Marshall, J., A. Adcroft, C. Hill, L. Perelman, and C. Heisey (1997), A finite-volume, in-
873 compressible navier stokes model for studies of the ocean on parallel computers, *Jour-
874 nal of Geophysical Research: Oceans*, 102(C3), 5753–5766, doi:10.1029/96JC02775.
- 875 Melet, A., S. Legg, and R. Hallberg (2016), Climatic impacts of parameterized lo-
876 cal and remote tidal mixing, *Journal of Climate*, 29(10), 3473–3500, doi:10.1175/
877 JCLI-D-15-0153.1.
- 878 Menemenlis, D., J.-M. Campin, P. Heimbach, C. Hill, T. Lee, A. Nguyen, M. Schodlok,
879 and H. Zhang (2008), Ecco2: High resolution global ocean and sea ice data synthesis,
880 *Mercator Ocean Quarterly Newsletter*, 31.
- 881 Müller, M., B. K. Arbic, J. G. Richman, J. F. Shriver, E. L. Kunze, R. B. Scott, A. J.
882 Wallcraft, and L. Zamudio (2015), Toward an internal gravity wave spectrum in
883 global ocean models, *Geophysical Research Letters*, 42(9), 3474–3481, doi:10.1002/
884 2015GL063365, 2015GL063365.
- 885 Munk, W., and C. Wunsch (1998), Abyssal recipes ii: Energetics of tidal and wind mix-
886 ing, *Deep-Sea Research I*, 45.
- 887 Ngodock, H. E., I. Souopgui, A. J. Wallcraft, J. G. Richman, J. F. Shriver, and B. K.
888 Arbic (2016), On improving the accuracy of the m_2 barotropic tides embedded in a
889 high-resolution global ocean circulation model, *Ocean Modelling*, 97, 16 – 26, doi:
890 <https://doi.org/10.1016/j.ocemod.2015.10.011>.
- 891 Penduff, T., B. Barnier, J.-M. Molines, and G. Madec (2006), On the use of current meter
892 data to assess the realism of ocean model simulations, *Ocean Modelling*, 11, 399–416.
- 893 Ponte, R. M., A. H. Chaudhuri, and S. V. Vinogradov (2015), Long-period tides in an at-
894 mospherically driven, stratified ocean, *Journal of Physical Oceanography*, 45(7), 1917–
895 1928, doi:10.1175/JPO-D-15-0006.1.
- 896 Qiu, B., S. Chen, P. Klein, J. Wang, H. Torres, L.-L. Fu, and D. Menemenlis (2018), Sea-
897 sonality in transition scale from balanced to unbalanced motions in the world ocean,
898 *Journal of Physical Oceanography*, 48(3), 591–605, doi:10.1175/JPO-D-17-0169.1.
- 899 Ray, R. D. (1999), *A Global Ocean Tide model from TOPEX/POSEIDON altimetry:
900 GOT99.2*, 1 v. pp., National Aeronautics and Space Administration, Goddard Space
901 Flight Center ; National Technical Information Service, distributor Greenbelt, M,
902 Springfield, Va.

- 903 Rocha, C. B., T. K. Chereskin, S. T. Gille, and D. Menemenlis (2016a), Mesoscale to sub-
904 mesoscale wavenumber spectra in Drake Passage, *Journal of Physical Oceanography*,
905 *46*(2), 601–620, doi:10.1175/JPO-D-15-0087.1.
- 906 Rocha, C. B., S. T. Gille, T. K. Chereskin, and D. Menemenlis (2016b), Seasonality of
907 submesoscale dynamics in the kuroshio extension, *Geophysical Research Letters*, *43*(21),
908 11,304–11,311, doi:10.1002/2016GL071349.
- 909 Savage, A. C., B. K. Arbic, J. G. Richman, J. F. Shriver, M. H. Alford, M. C. Buijs-
910 man, J. Thomas Farrar, H. Sharma, G. Voet, A. J. Wallcraft, and L. Zamudio (2017a),
911 Frequency content of sea surface height variability from internal gravity waves to
912 mesoscale eddies, *Journal of Geophysical Research: Oceans*, *122*(3), 2519–2538, doi:
913 10.1002/2016JC012331.
- 914 Savage, A. C., B. K. Arbic, M. H. Alford, J. K. Ansong, J. T. Farrar, D. Menemenlis,
915 A. K. O'Rourke, J. G. Richman, J. F. Shriver, G. Voet, A. J. Wallcraft, and L. Za-
916 mudio (2017b), Spectral decomposition of internal gravity wave sea surface height in
917 global models, *Journal of Geophysical Research: Oceans*, *122*(10), 7803–7821, doi:
918 10.1002/2017JC013009.
- 919 Scott, R. B., B. K. Arbic, E. P. Chassignet, A. C. Coward, M. Maltrud, W. J. Merryfield,
920 A. Srinivasan, and A. Varghese (2010), Total kinetic energy in four global eddying
921 ocean circulation models and over 5000 current meter records, *Ocean Modelling*, *32*,
922 157–169.
- 923 Shriver, J. F., B. K. Arbic, J. G. Richman, R. D. Ray, E. J. Metzger, A. J. Wallcraft,
924 and P. G. Timko (2012), An evaluation of the barotropic and internal tides in a high-
925 resolution global ocean circulation model, *J. Geophys. Res.*, *117*(C10024), doi:10.1029/
926 2012JC008170.
- 927 Silverthorne, K. E., and J. M. Toole (2009), Seasonal kinetic energy variability of near-
928 inertial motions, *Journal of Physical Oceanography*, *39*(4), 1035–1049, doi:10.1175/
929 2008JPO3920.1.
- 930 Simmons, H. L., and M. H. Alford (2012), Simulating the long-range swell of internal
931 waves generated by ocean storms, *Oceanography*, *25*.
- 932 Skiba, A. W., L. Zeng, B. K. Arbic, M. Müller, and W. J. Godwin (2013), On the reso-
933 nance and shelf/open-ocean coupling of the global diurnal tides, *Journal of Physical*
934 *Oceanography*, *43*(7), 1301–1324, doi:10.1175/JPO-D-12-054.1.

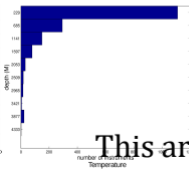
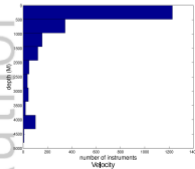
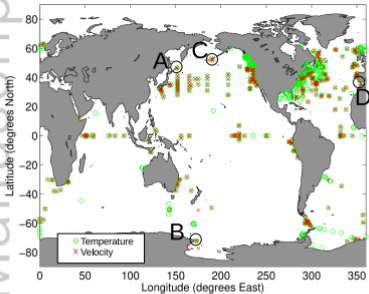
- 935 Smith, W. H. F., and D. T. Sandwell (1997), Global sea floor topography from satellite al-
936 timetry and ship depth soundings, *Science*, 277(5334), 1956–1962, doi:10.1126/science.
937 277.5334.1956.
- 938 Stammer, D., R. D. Ray, O. B. Andersen, B. K. Arbic, W. Bosch, L. Carrère, Y. Cheng,
939 D. S. Chinn, B. D. Dushaw, G. D. Egbert, S. Y. Erofeeva, H. S. Fok, J. A. M. Green,
940 S. Griffiths, M. A. King, V. Lapin, F. G. Lemoine, S. B. Luthcke, F. Lyard, J. Morison,
941 M. Müller, L. Padman, J. G. Richman, J. F. Shriver, C. K. Shum, E. Taguchi, and Y. Yi
942 (2014), Accuracy assessment of global barotropic ocean tide models, *Reviews of Geo-*
943 *physics*, 52(3), 243–282, doi:10.1002/2014RG000450.
- 944 Su, Z., J. Wang, P. Klein, A. F. Thompson, and D. Menemenlis (2018), Ocean subme-
945 soscales as a key component of the global heat budget, *Nature Communications*, 9(1),
946 775, doi:10.1038/s41467-018-02983-w.
- 947 Thoppil, P. G., J. G. Richman, and P. J. Hogan (2011), Energetics of a global ocean cir-
948 culation model compared to observations, *Geophysical Research Letters*, 38(15), doi:
949 10.1029/2011GL048347.
- 950 Timko, P. G., B. K. Arbic, J. G. Richman, R. B. Scott, E. J. Metzger, and A. J. Wallcraft
951 (2012), Skill tests of three-dimensional tidal currents in a global ocean model: A look
952 at the North Atlantic, *J. Geophys. Res: Oceans*, 117(C8), doi:10.1029/2011JC007617.
- 953 Timko, P. G., B. K. Arbic, J. G. Richman, R. B. Scott, E. J. Metzger, and A. J. Wallcraft
954 (2013), Skill testing a three-dimensional global tide model to historical current meter
955 records, *J. Geophys. Res: Oceans*, 118(12), 6914–6933, doi:10.1002/2013JC009071.
- 956 Torres, H. S., P. Klein, D. Menemenlis, B. Qiu, Z. Su, J. Wang, S. Chen, and L.-L. Fu
957 (2018), Partitioning ocean motions into balanced motions and internal gravity waves:
958 A modeling study in anticipation of future space missions, *Journal of Geophysical Re-*
959 *search: Oceans*, 123(11), 8084–8105, doi:10.1029/2018JC014438.
- 960 Wang, J., L.-L. Fu, B. Qiu, D. Menemenlis, J. T. Farrar, Y. Chao, A. F. Thompson, and
961 M. M. Flexas (2018), An observing system simulation experiment for the calibration
962 and validation of the surface water ocean topography sea surface height measurement
963 using in situ platforms, *Journal of Atmospheric and Oceanic Technology*, 35(2), 281–
964 297, doi:10.1175/JTECH-D-17-0076.1.
- 965 Wang, J., L.-L. Fu, H. S. Torres, S. Chen, B. Qiu, and D. Menemenlis (2019), On the
966 spatial scales to be resolved by the surface water and ocean topography ka-band radar
967 interferometer, *Journal of Atmospheric and Oceanic Technology*, 36(1), 87–99, doi:

968 10.1175/JTECH-D-18-0119.1.

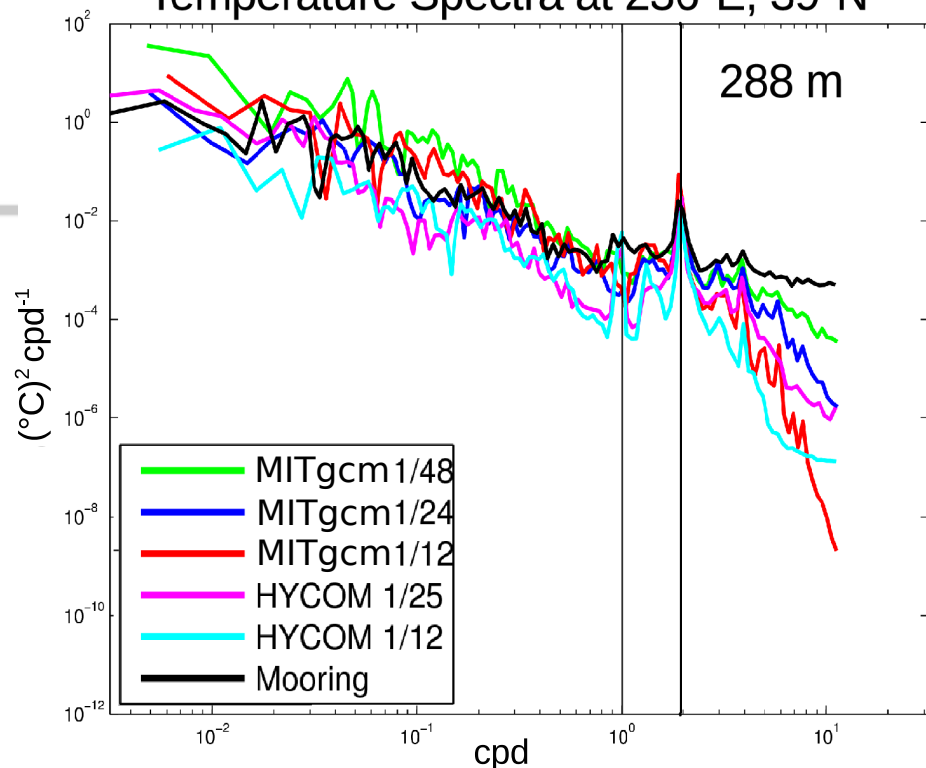
969 Weis, P., M. Thomas, and J. Sündermann (2008), Broad frequency tidal dynamics simu-
970 lated by a high-resolution global ocean tide model forced by ephemerides, *Journal of*
971 *Geophysical Research: Oceans*, 113(C10), doi:10.1029/2007JC004556.

972 Wright, C. J., R. B. Scott, P. Ailliot, and D. Furnival (2014), Lee wave generation rates
973 in the deep ocean, *Geophysical Research Letters*, 41(7), 2434–2440, doi:10.1002/
974 2013GL059087.

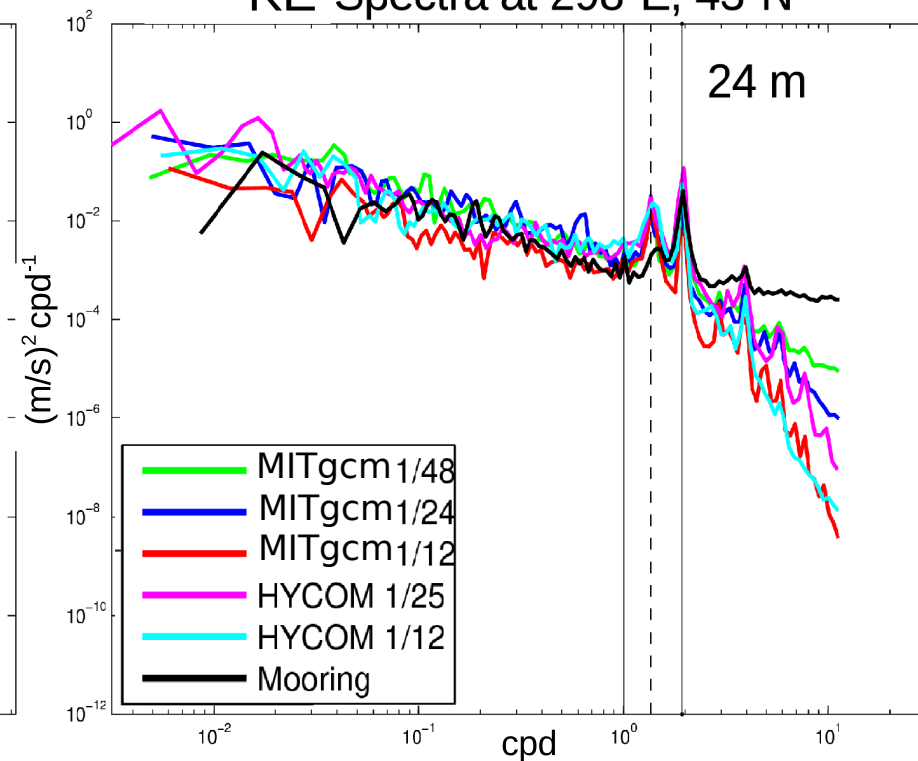
975 Yu, X., A. L. Ponte, S. Elipot, D. Menemenlis, E. D. Zaron, and R. Abernathey (2019),
976 Surface kinetic energy distributions in the global oceans from a high-resolution numeri-
977 cal model and surface drifter observations, *Geophysical Research Letters*, 46(16), 9757–
978 9766, doi:10.1029/2019GL083074.



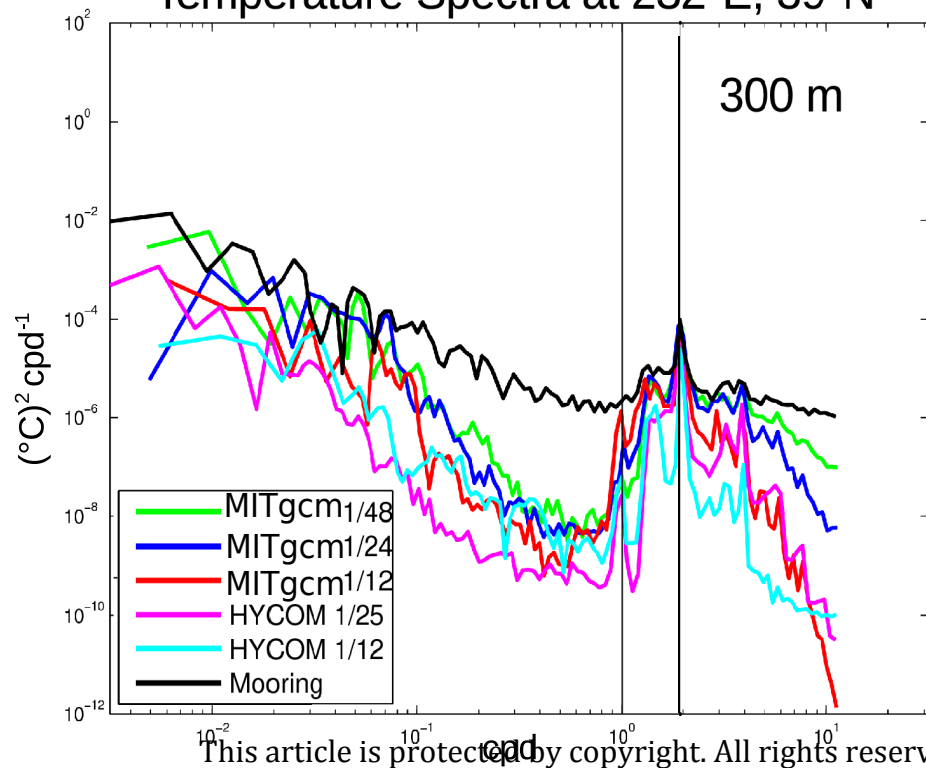
Temperature Spectra at 236°E, 39°N



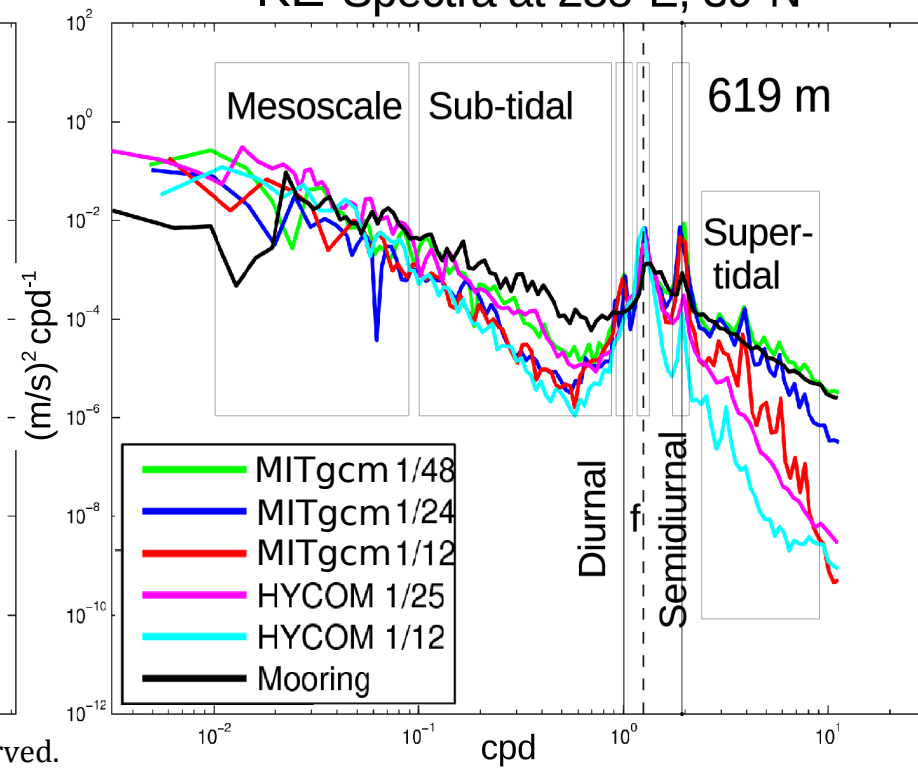
KE Spectra at 298°E, 43°N



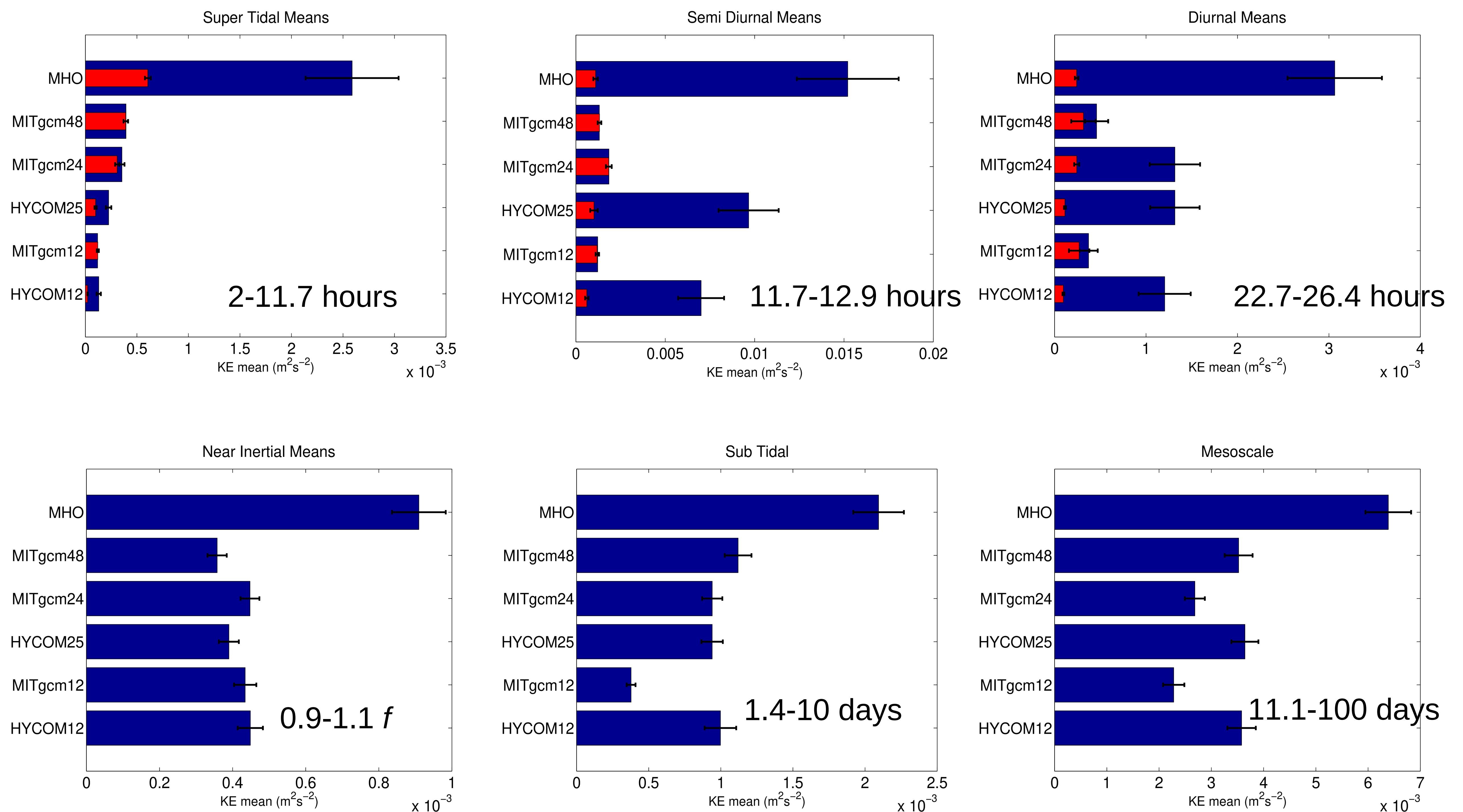
Temperature Spectra at 232°E, 39°N



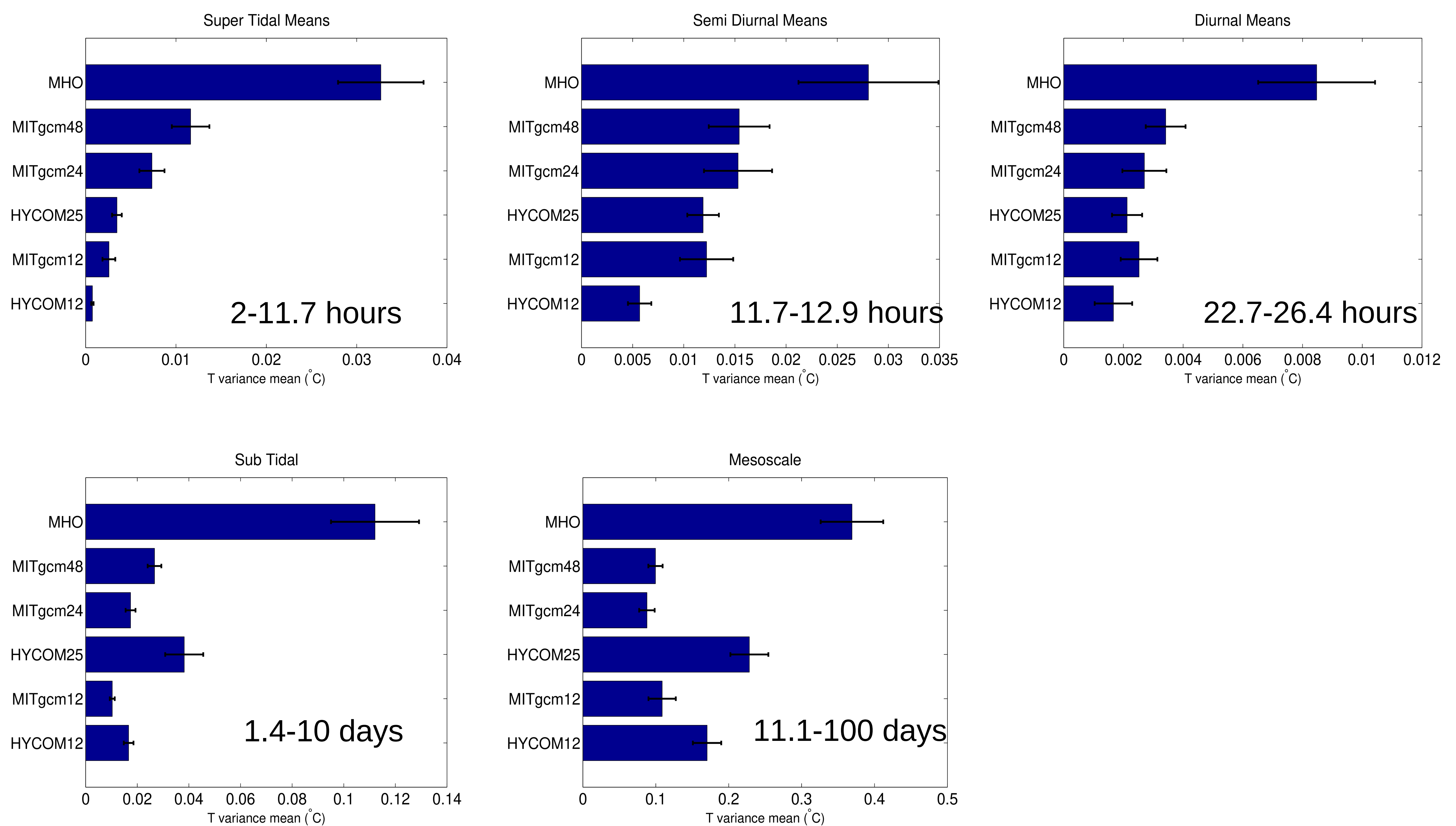
KE Spectra at 288°E, 39°N



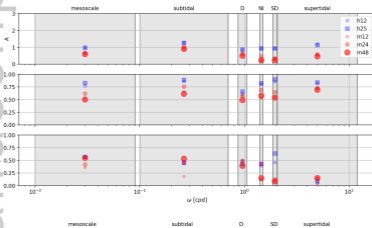
Mean KE by Frequency Band



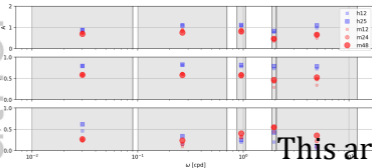
Mean Temperature Variance by Frequency Band



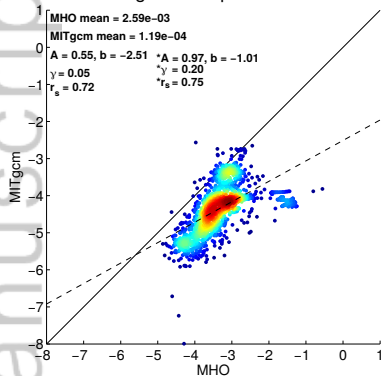
KE comparison metrics



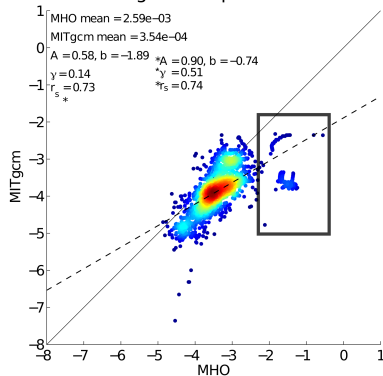
temperature variance comparison metrics



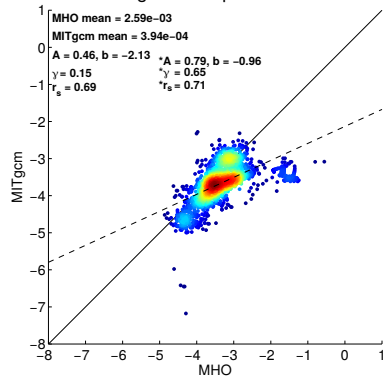
MITgcm12 Super Tidal KE



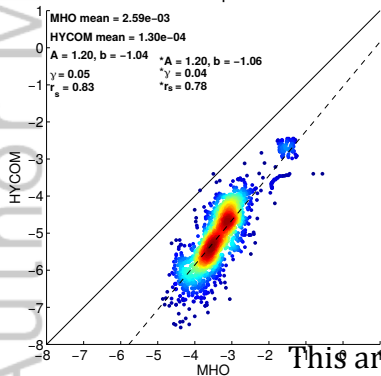
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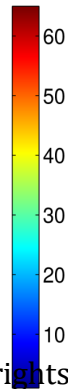
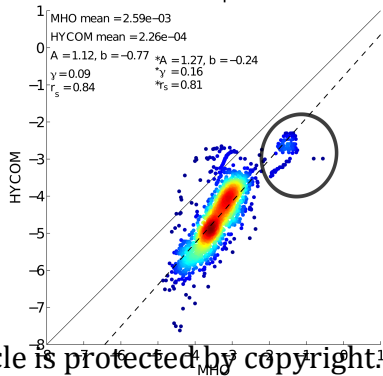
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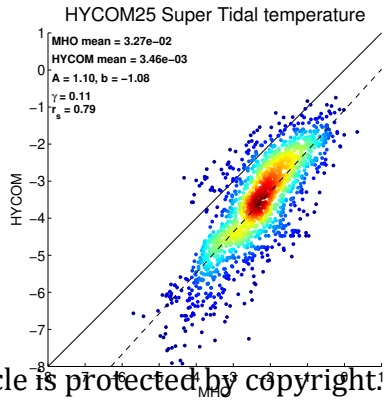
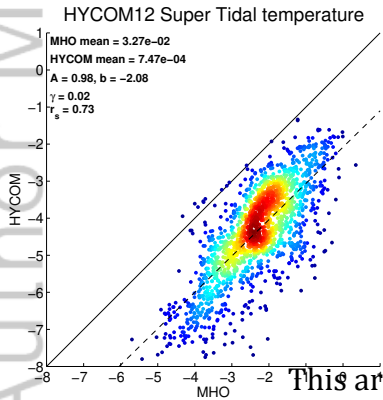
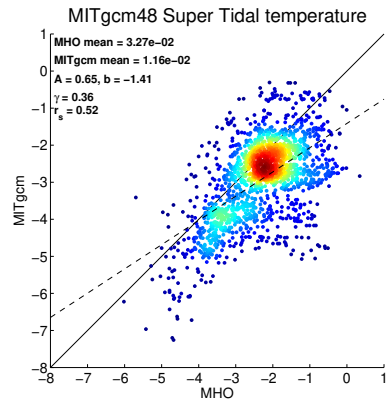
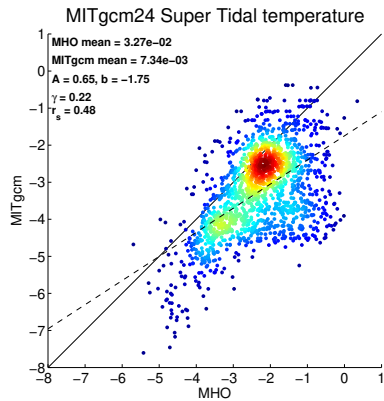
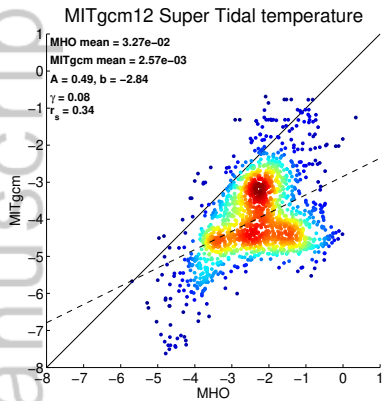


HYCOM12 Super Tidal KE

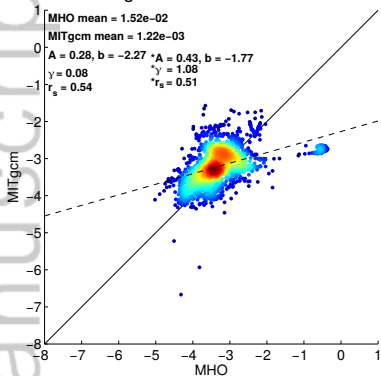


HYCOM25 Super Tidal KE

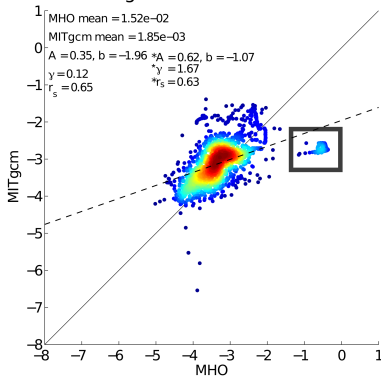




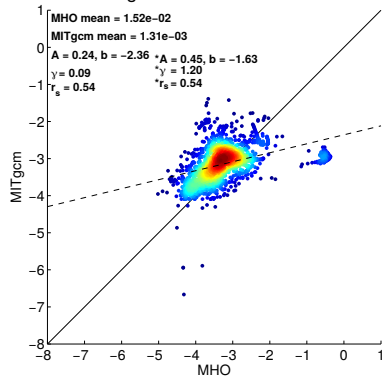
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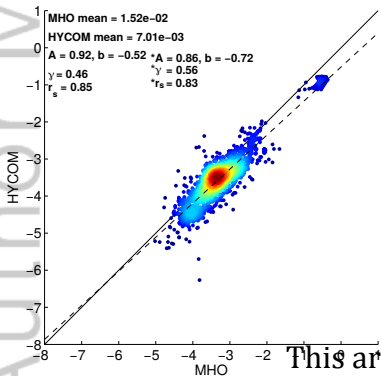
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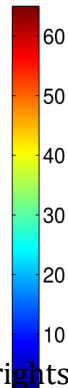
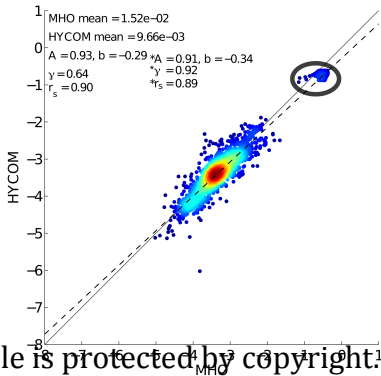
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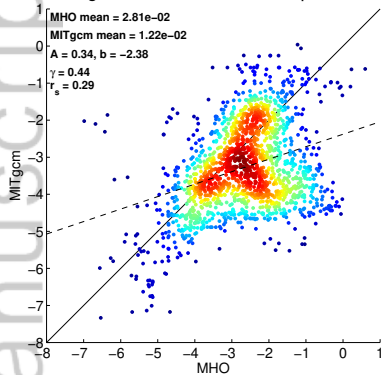
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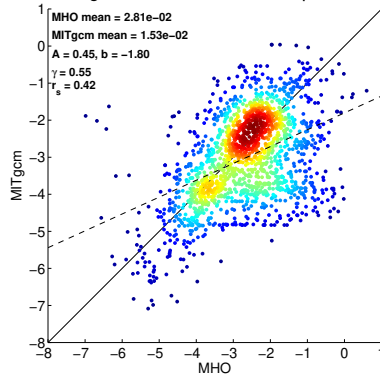
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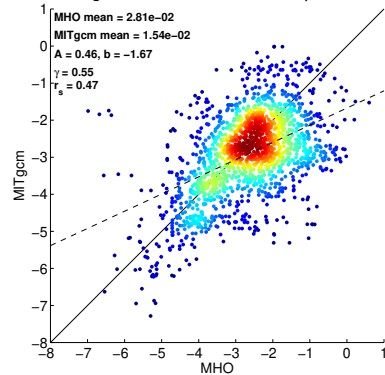
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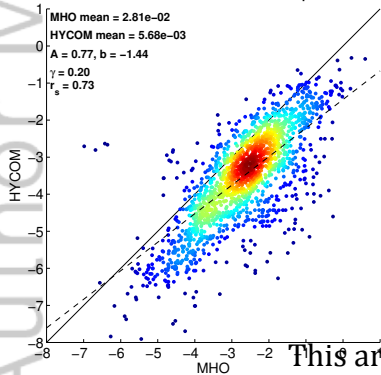
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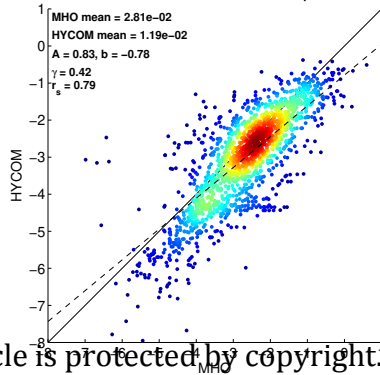
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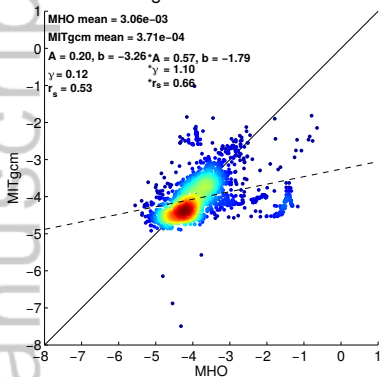
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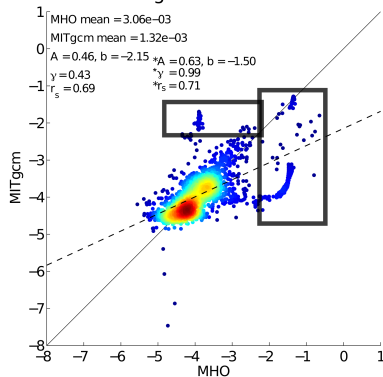
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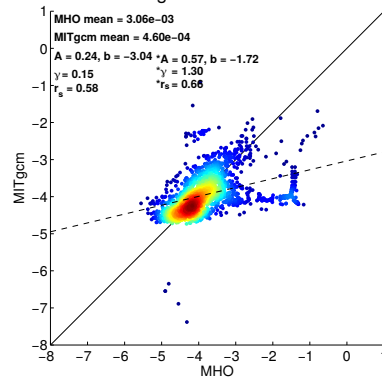
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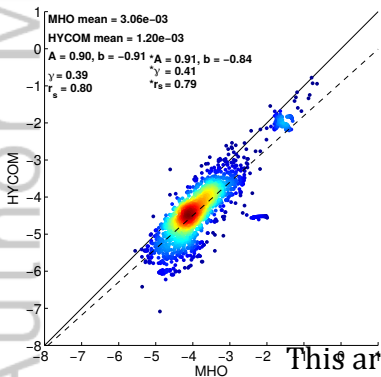
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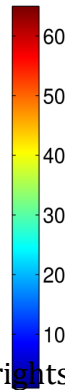
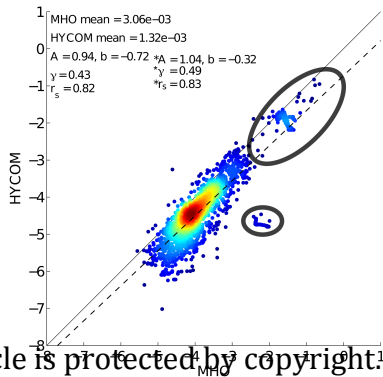
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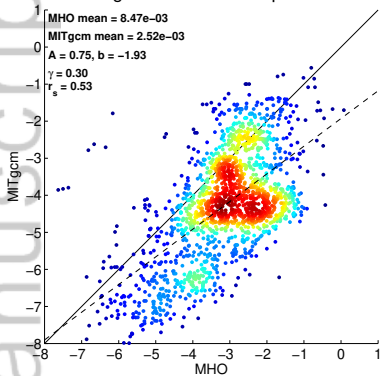
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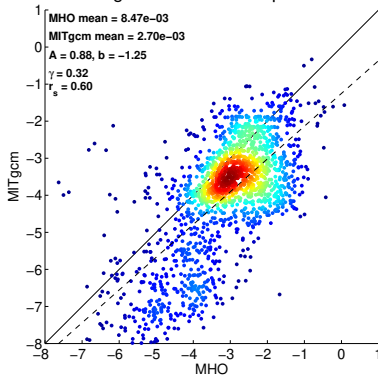
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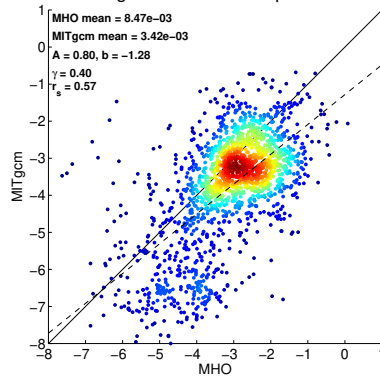
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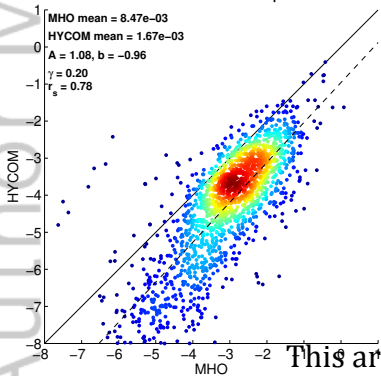
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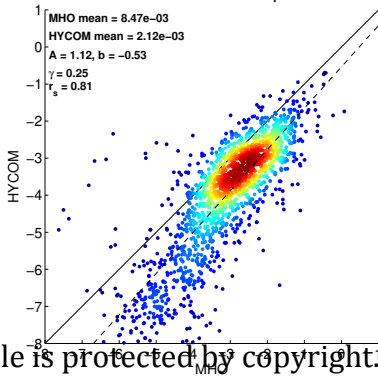
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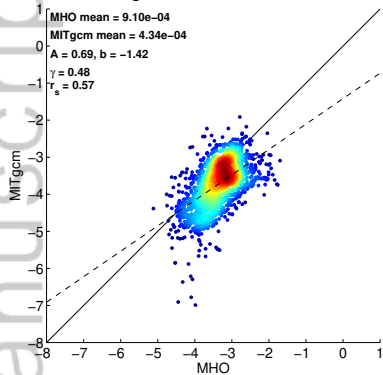
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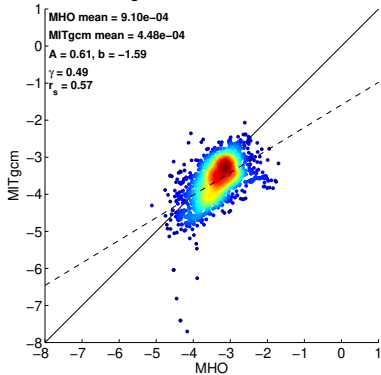
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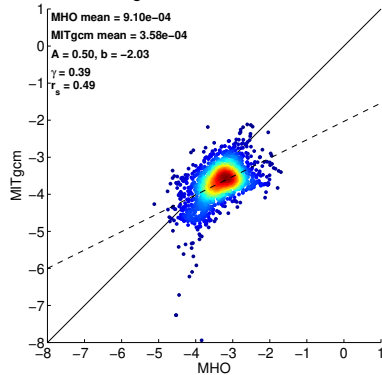
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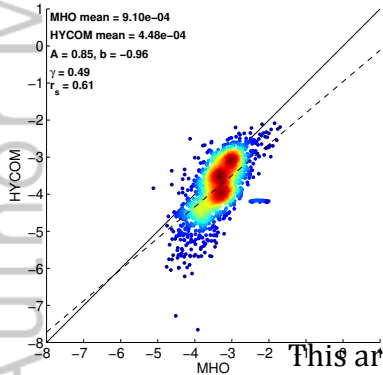
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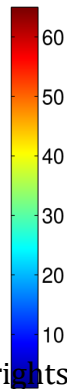
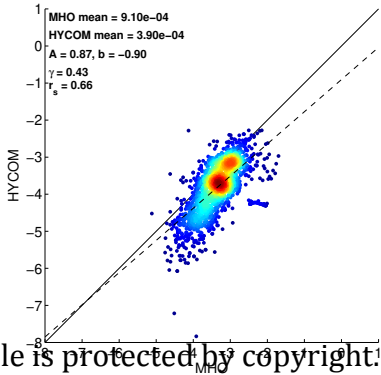
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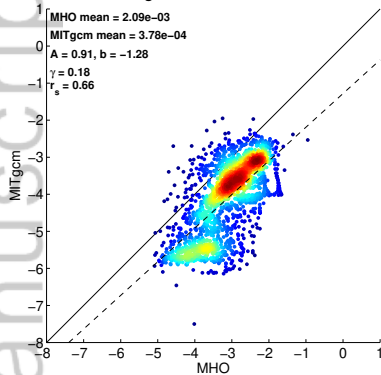
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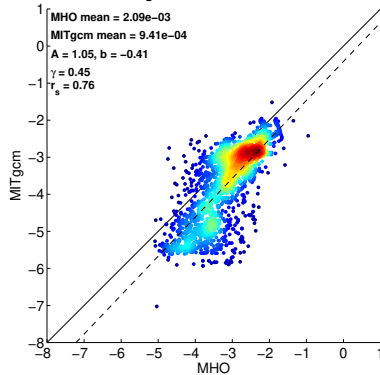
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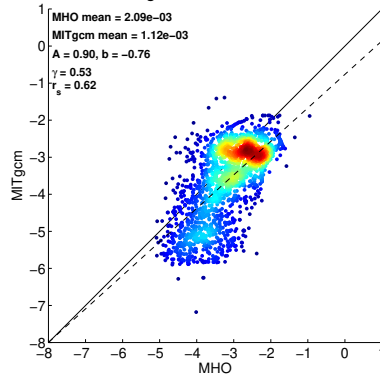
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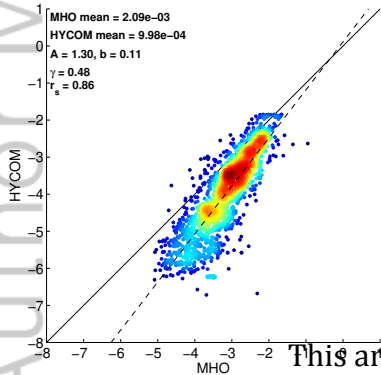
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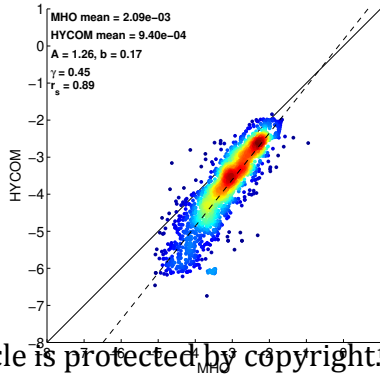
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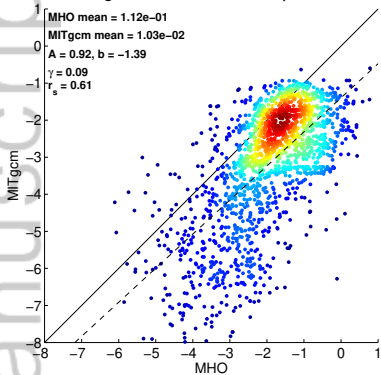
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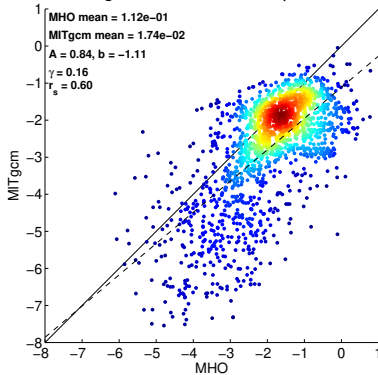
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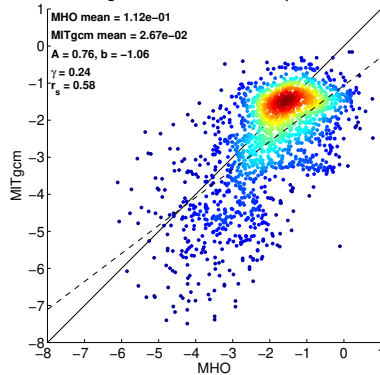
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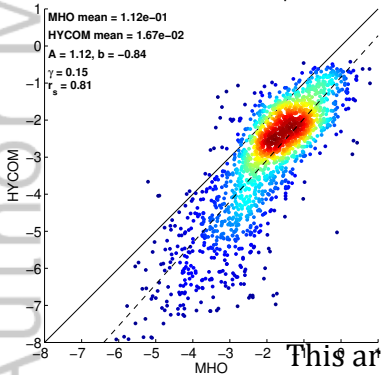
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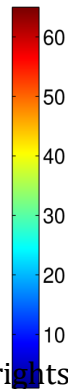
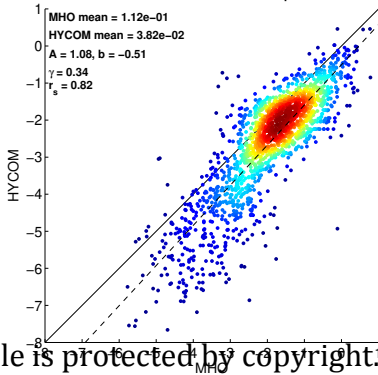
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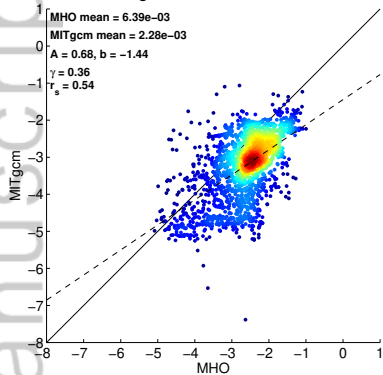
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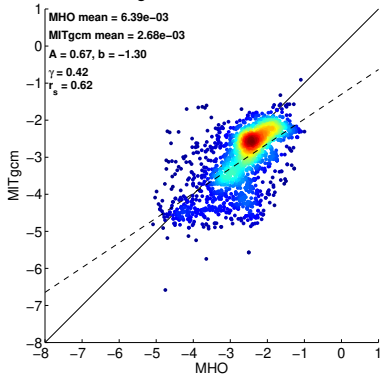
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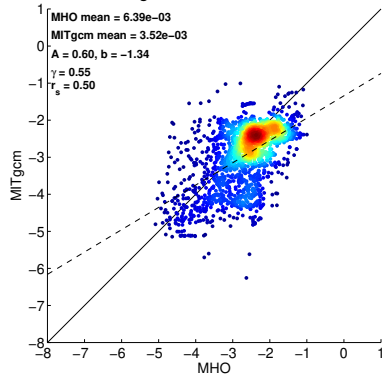
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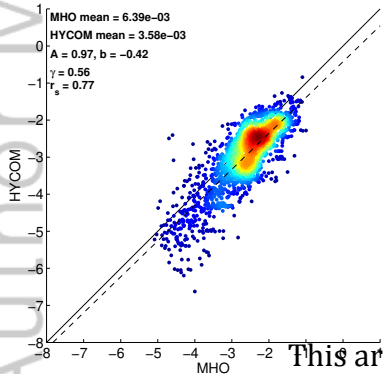
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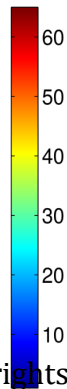
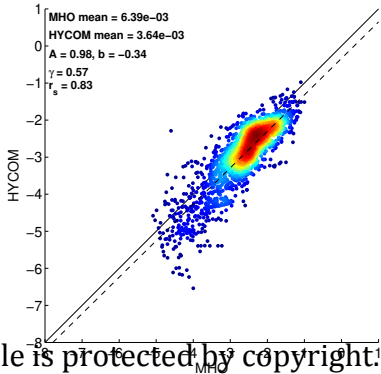
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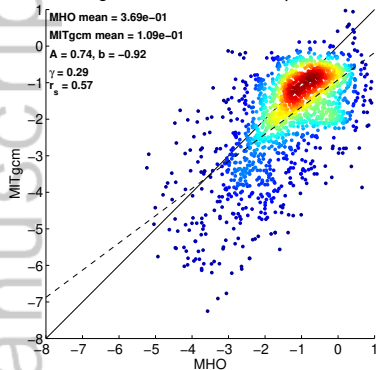
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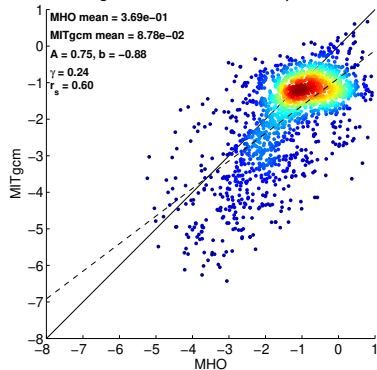
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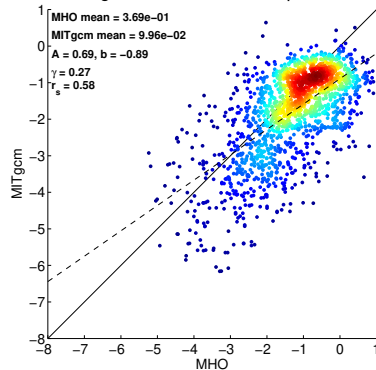
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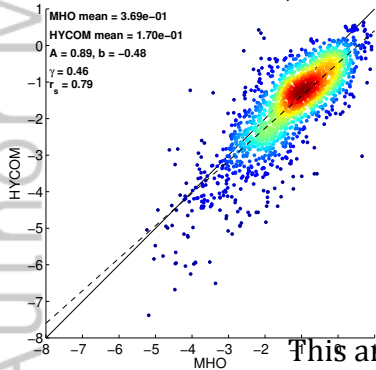
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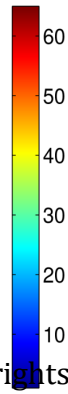
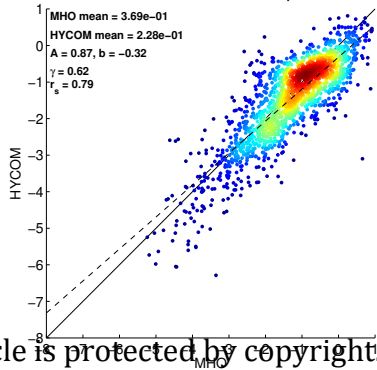
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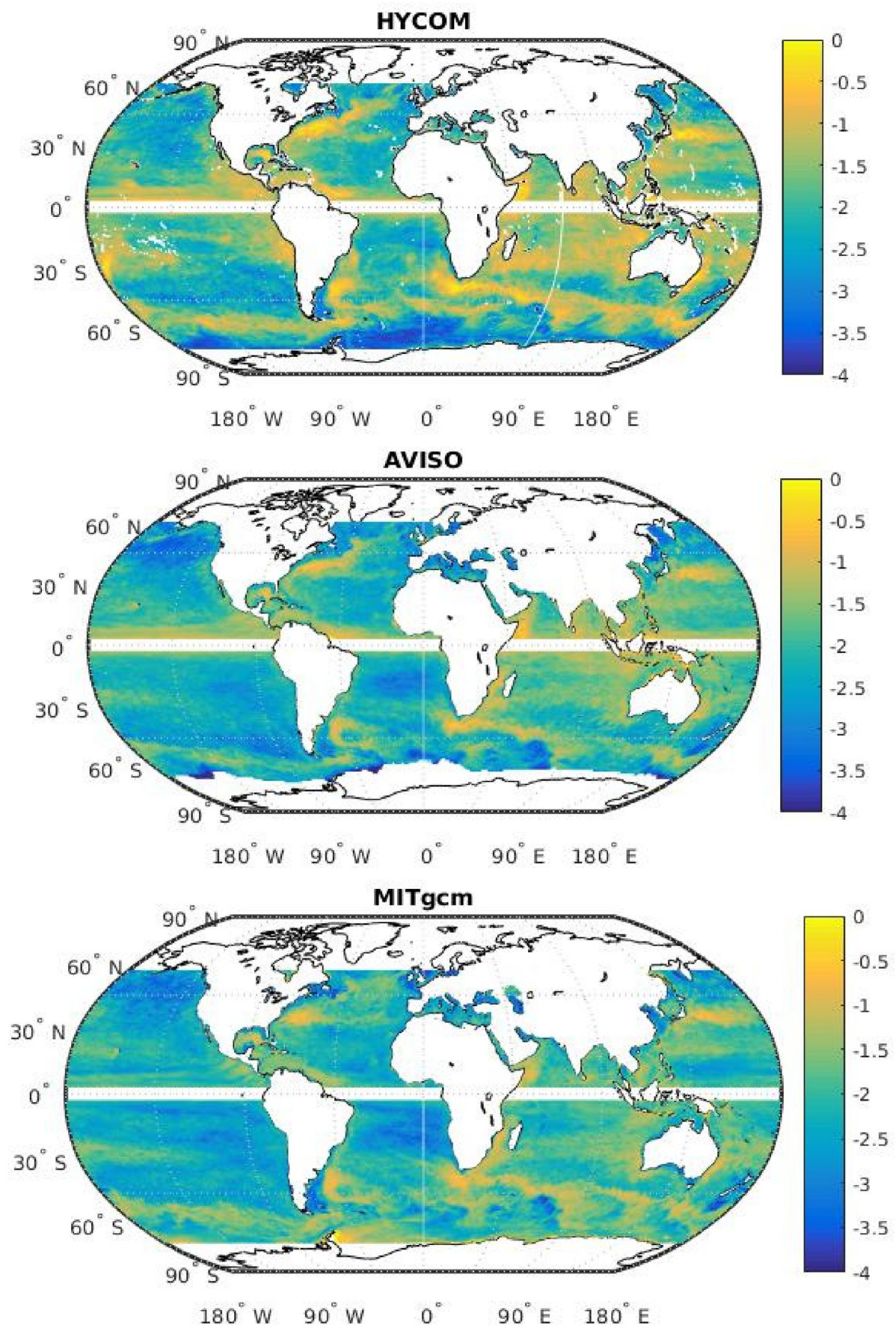


HYCOM12 Mesoscale temperature



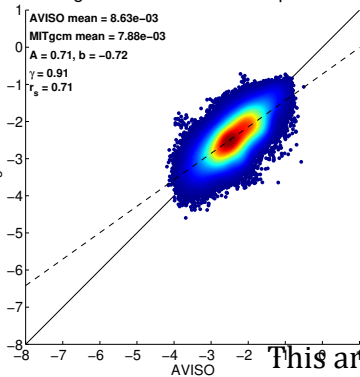
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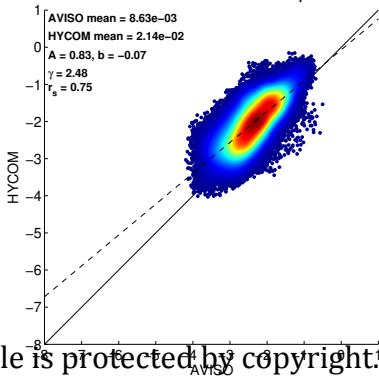


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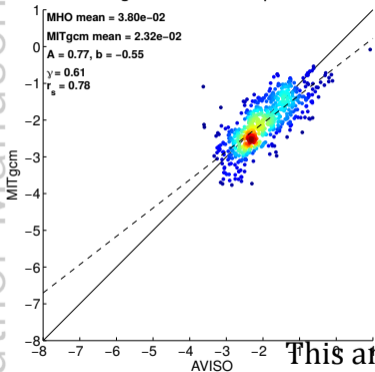
MITgcm1/48 Global Geostrophic KE



HYCOM1/25 Global Geostrophic KE



MITgcm1/48 Geostrophic KE



HYCOM1/25 Geostrophic KE

