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Residual Study: Testing Jupiter Atmosphere Models Against Juno MWR Observations

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Abstract

The Juno spacecraft provides unique close-up views of Jupiter underneath the synchrotron radiation belts while circling Jupiter in its 53-day orbits. The Microwave Radiometer (MWR) on board measures Jupiter thermal radiation at wavelengths between 1.37 cm and 50 cm, penetrating the atmosphere to a pressure of a few hundred bars and greater. The mission provides the first measurements of Jupiter's deep atmosphere, down to ~250 bars in pressure, constraining the vertical distributions of its kinetic temperature and constituents. As a result, vertical structure models of Jupiter's atmosphere may now be tested by comparison with MWR data. Taking into account the MWR beam patterns and observation geometries, we test several published Jupiter atmospheric models against MWR data. Our residual analysis confirms Li et al. 2017's result that ammonia depletion persists down to 50~60 bars where ground-based VLA was not able to observe. We also present an extension of the study that iteratively improves the input model and generates Jupiter brightness temperature maps which best match the MWR data. A feature of Juno's north-to-south scanning approach is that latitudinal structure is more easily obtained than longitudinal, and the creation of optimum two-dimensional maps is addressed in this approach.

This is the author manuscript accepted for publication and has undergone full peer review but habitroduction the copyediting, typesetting, pagination and proofreading process, which may lead to differences between this version and the Version of Record. Please cite this article On daig 127,1208, 120-20 For Space 23 ft started its flybys over Jupiter, allowing the Microwave Radiometer (MWR) instrument [Janssen et al. 2017] to measure the thermal emission of Jupiter's atmosphere from pressure levels of

approximately 0.5 bar to a few hundred bars [Bolton et al., 2017]. A main objective of the MWR is to measure Jupiter's deep water abundance, because it is key to understanding the history of the giant planet's volatile and heavy elements [Helled and Lunine, 2014], and is essential for understanding the meteorology that is observed at the visible cloud level [e.g. Showman, 2007]. Jupiter's brightness temperatures at the MWR wavelengths are highly sensitive to the ammonia distribution in Jupiter's atmosphere [Janssen et al., 2017], the major opacity source in the MWR channels. Prior to the arrival of the Juno spacecraft, the only direct knowledge of Jupiter's deep atmosphere profile (i.e. temperature, water and ammonia abundances) came from the Galileo Probe, which was restricted to in situ measurements during a single plunge into the atmosphere at a longitude of 4.5°W (System III) and a planetocentric latitude of 6.5°N (within a relatively dry and cloudless area at the southern edge of the north equatorial belt, NEB) down to less than 20 bars. Most previous works have focused on similarly clear and dry downwelling regions with depletions of ammonia and water, where the probe entered and where spectroscopic measurements are taken [Bjoraker et al. 1986, Grassi et al. 2017]. Earth-based radio observations [e.g. de Pater et al. 2016, 2019] give global coverage, but require assumptions about limb-darkening because viewing angle is correlated with latitude, and are limited by the foreground synchrotron radiation emitted by high-energy electrons gyrating around Jupiter's intense magnetic field [Burke and Franklin 1955, Santos-Costa et al. 2017]. Juno's orbit takes it beneath the radiation belts, largely alleviating the limitations imposed by synchrotron emission, and allowing finer spatial resolution than most Earth-based radio observations while observing each location from multiple viewing angles.

Li et al. (2017) have retrieved a vertical ammonia distribution from 0.5 bar to 100 bars inverted from Juno MWR observations from Juno's first science pass, perijove (PJ) 1, and Li et al. (2020) used MWR data to derive water and ammonia abundance in a narrow latitude band near the equator. de Pater et al. (2019) also derived a vertical ammonia distribution from ground-based VLA observations at 3-37 GHz, showing compatibility with a subset of the MWR data described by Li et al. (2017). However, interpreting MWR data in this way requires complex data processing and a variety of assumptions about data smoothness, symmetry of the planet, etc. [see, e.g., Oyafuso et al 2020]. Instead, in this paper we apply a model in the forward direction to produce synthetic MWR observations and then compare those predictions with the actual MWR data set. This approach is simpler, requires fewer assumptions, and allows direct comparison of residuals with the known uncertainties of the MWR data set [Janssen et al. 2017].

In Section 2, we begin with a brief description of the MWR observations. In Section 3, we provide details of the modeling approach and present the iterative residual calculation used to generate 2D Jupiter maps with respect to latitude and longitude. In Section 4, we test several Jupiter atmosphere models proposed by previous works against the MWR observations.

2. MWR Observation Description

The Juno/MWR instrument measures the thermal emission of Jupiter's atmosphere at six widely separated wavelengths and multiple emission angles from nadir to greater than 50 degrees. The antenna temperature (T_a) represents the power received by the antenna, which is the convolution of the source brightness temperature (T_b) distribution in the field of view over the broad antenna gain pattern [Janssen et al. 2017]; i.e.,

$$T_a(\theta,\phi) = \int_0^{2\pi} \int_0^{\pi} T_b(\theta',\phi') g(\theta'-\theta,\phi'-\phi) \sin(\theta'-\theta) d\theta' d\phi'$$
(1)

where $(\theta' - \theta, \phi' - \phi)$ is the angular deviation from the beam boresight direction (θ, ϕ) . $T_b(\theta', \phi')$ is the brightness temperature in the direction of the solid angle element $\sin(\theta' - \theta)d\theta'd\phi'$ and $g(\theta' - \theta, \phi' - \phi)$ is the gain of the antenna toward this direction for an antenna pointed in the direction (θ, ϕ) . The gain $g(\theta, \phi)$ [Janssen 1993, Chapter1, Eqn.1.27] is normalized to 1 over the 4π sky as shown in Equation 2 [Janssen et al. 2017], which makes it a factor of 4π smaller than the commonly defined gain in radio astronomy [Rohlfs and Wilson 1999, Chapter 5, Eqn.5.48].

$$\int_0^{2\pi} \int_0^{\pi} g(\theta, \phi) \sin\theta d\theta d\phi = 1$$
⁽²⁾

100 The main beam of the antenna contributes most to the antenna temperature and can be considered as a general 101 estimate of the boresight antenna temperature at (θ, ϕ) ; however, the sidelobes view a brightness distribution that can 102 be considerably different. For example, we need to account for off axis Jovian thermal emission at different 103 emission angles, the empty-sky cosmic microwave background (CMB), galactic emission, and Jovian synchrotron 104 radiation, the latter being especially significant in channels 1 and 2.

The gain pattern for each of the six antennas has been measured over the full sphere pre-launch [Janssen et al., 2017] (see Figure 1 for antenna pattern, relative gain with respect to the peak gain). An arbitrary Jupiter atmosphere model can thus be used to calculate the antenna temperature at all positions on Jupiter from all viewing perspectives. This modeled antenna temperature can then be compared with observed antenna temperatures. The residual is defined as the difference between the two:

$$T_{resid} = T_a^{obs} - T_a^{model} \tag{3}$$

The set of residuals obtained for a given pass is thus a measure of how consistent the model is with the real Jupiter atmosphere. Note also that the radiometric measurements have both systematic calibration uncertainties in addition to thermal measurement noise that must be taken into account, so that interpreting the residuals is not always as simple as minimizing the offset. If the model matched Jupiter exactly, the systematic measurement effect due to absolute calibration error could introduce constant but unknown offsets to all the residuals at each frequency. Brown et al. (2020) estimate this to be as large as 2%. Other calibration errors such as uncertainties in the measured beam patterns are expected to result in residuals of this order or less [Janssen et al, 2017]. Intrinsic receiver noise introduces variations of order 0.1% which leads to scatter in the residuals [Oyafuso et al. 2020]. In practice, an atmospheric model would be considered inconsistent with the MWR measurements if the bulk residuals are larger than 2%. It would also be considered inconsistent if the residuals for any channel varied by more than 0.2%.



Figure 1. Antenna patterns (relative gain in dB with respect to peak gain) for all 6 MWR channels. The main beam half power width is 20° for channels 1 and 2, 12° for channels 3, 4, 5, and 11° for channel 6. $x = \sin(polar \ angle) * \cos(azimuthal \ angle), y = \sin(polar \ angle) * \sin(azimuthal \ angle)$, where x = y = 0 at the boresight.

3. Residual Analysis

134 We developed a residual analysis method to compare any brightness temperature distribution model over the planet 135 to Juno MWR observations or evaluate any atmosphere profile model to determine if it results in the MWR observed 136 brightness temperatures. We describe and test several published Jupiter atmospheric models against MWR data in 137 Section 4. In our analysis, all of the geometries of the spacecraft and the antenna beam coverage are calculated using 138 SPICE kernel information [Acton, 1996]. The synchrotron emission and galactic background are accounted for 139 appropriately as described in Oyafuso et al. 2020 and Adumitroaie at al. 2016. In order to verify any proposed 140 atmosphere model with ammonia, water and temperature profiles, we run the radiative-transfer code JAMRT 141 [Janssen et al. 2013] to simulate the radiative-transfer process and obtain the brightness temperatures at various 142 emission angles and latitudes at all MWR frequencies. For a standard forward model, JAMRT takes in user-143 specified NH₃ and H₂O enrichment, an adiabatic temperature profile is then calculated. In a simple standard model, 144 NH₃ and H₂O abundances are uniform in the deep atmosphere and become saturated above the corresponding cloud 145 level. However, JAMRT also allows user-specified NH₃/H₂O abundances and temperature profiles. With six MWR 146 channels sensitive to different pressure levels down to more than 250 bars, the residual values show directly if the 147 composition distribution produced by the atmospheric model is compatible with the data. The effects of ammonia, 148 temperature and water distributions on brightness temperature are entangled and deciding whether a specific model 149 is the best one may not be possible using MWR data alone. However, our residual analysis process can determine 150 whether the model is possible in the sense of being consistent with the MWR observations. We calculate the residual 151 value of all observations, each corresponding to an observed emission angle, latitude and longitude, which indicates 152 how well the proposed model matches the actual Jupiter observations. At any specific latitude, trends in the residual 153 values with respect to the emission angles reveal how well the limb darkening in the proposed model matches the 154 real case. If at large emission angles, the residual values have a downward trend, a larger limb-darkening than the 155 proposed value is observed by the MWR observations. We will review test cases in Section 4 and explain more 156 about how to interpret the output residual values in terms of Jupiter's atmospheric properties. 157

Based on the residual analysis described above, we also developed an iterative approach that takes any initial brightness temperature model, calculates the residual values, updates the brightness temperature model based on the residual values, and uses the updated model as a new input brightness temperature in the next iteration. The goal of this process is to converge on a best estimate of Jupiter's brightness distribution along with its emission angle dependence. We iterate the process until the residual values converge and approach zero. This iterative process is a good way to investigate small longitudinal and latitudinal structures, and make 2D brightness-temperature maps (with respect to latitude and longitude) of Jupiter at MWR observed pressure levels. At each iteration *i*, we begin with convolving the regularly gridded input model *Model*ⁱ (*emissionangle*, *latitude*, *longitude*) with the MWR beam pattern coverage and generate simulated antenna temperatures $T_{a,model}^{i}(Obs_{emiss}, Obs_{lat}, Obs_{lon})$ for each observation point. The regularly gridded input model *Model*ⁱ (*emissionangle*, *latitude*, *longitude*) is in units of brightness temperature at a series of emission angle grids on each Jupiter's latitude - longitude- grid. We then subtract the modeled antenna temperature from the observed antenna temperature and derive the output residual values $Residual^i(Obs_{emiss}, Obs_{lat}, Obs_{lon}) = T_{a,obs}^{i}(Obs_{emiss}, Obs_{lat}, Obs_{lon}) - T_{a,model}^{i}(Obs_{emiss}, Obs_{lat}, Obs_{lon})$, which serve as a proxy for the difference between the input brightness temperature model and the real Jupiter brightness temperature at the corresponding boresight emission angle, latitude and longitude. We introduce the parameter:

$$T_{best}^{l}(Obs_{emiss}, Obs_{lat}, Obs_{lon}) = Residual^{l}(Obs_{emiss}, Obs_{lat}, Obs_{lon}) + T_{model}^{l}(Obs_{emiss}, Obs_{lat}, Obs_{lon})$$
(4)

where $T_{inotel}^{i}(Obs_{emiss}, Obs_{lat}, Obs_{lon})$ is the brightness temperature at the observation point according to the input model, and $T_{best}^{i}(Obs_{emiss}, Obs_{lat}, Obs_{lon})$ is an updated and better estimation for Jupiter's brightness temperature at this specific observing geometry. Constrained by the finite size of our beam and the fact that we lack perfect knowledge of the brightness temperature distribution in adjacent positions at the beginning of this iteration process, $T_{best}^{i}(Obs_{emiss}, Obs_{lat}, Obs_{lon})$ will tend to approach the real Jupiter brightness temperature value after each iteration. Several iterations are expected to be needed to achieve a value within measurement uncertainty.

A crucial step in this iteration process is to update the Jupiter model at the end of each iteration. For typical Juno MWR orbits, each latitude was observed at multiple emission angles and within a small longitudinal range (see Figure 2, middle panel, as an example of a typical MWR observation track), which is necessary for determining the limb-darkening value and diminishing the longitudinal variation in each latitude bin. This is accomplished with the

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spacecraft spin axis orthogonal to its orbital plane and the MWR antenna beams sweeping north to south (antennas 188 fixed to the spacecraft) as the Juno spacecraft also passes over the planet from north to south. A second and special 189 case is the MWR tilt orbit (Fig. 2 left panel), which is designed to minimize the longitudinal coverage (the 190 spacecraft attitude is slightly tilted to compensate for Jupiter's rotation) and provide data for more accurate analysis of nadir brightness and limb darkening. For typical MWR orbits, we first ignore the longitudinal variation and derive 192 the best-fit brightness model with respect to latitudes and emission angles. After that, with the knowledge of the 193 limb-darkening values, we fix the limb darkening and further derive the best-fit nadir brightness temperature model with respect to latitudes and longitudes in order to reveal small longitudinal structures. The cross-track orbit is yet another special MWR orbit [Janssen et al. 2017] (Fig. 2 right panel), where at each latitude a wide range of longitudes are observed but each longitude was only observed once at a specific emission angle. The spacecraft spin axis is parallel with Jupiter's spin axis and the antenna beams sweep from east to west (from lower to higher west longitudes in System III), across the planet as the spacecraft spins. In these cases, at each latitude the variations due respectively to longitudinal structure and limb darkening are entangled and impossible to distinguish. Therefore, a prior value for the limb darkening at each latitude becomes important. For different data features and also for comparison purposes, we proposed four model-updating methods as described below.



Figure 2. MWR observation tracks for an MWR tilt orbit (left panel, PJ5 as an example), a typical MWR orbit (middle panel, PJ7 as an example) and an MWR Cross-track orbit (right panel, PJ19 as an example).

(1) Method 1: Interpolate and smooth among T_{best}^{i} with respect to emission angles.

For typical MWR orbits and MWR tilt orbits with narrow longitudinal coverage, method 1 is used to derive a brightness temperature model with respect to latitudes and emission angles. The iteration process is not sensitive to the initial model and the residuals converge to near zero usually after 5 iterations. We assume that at iteration *i* the input gridded model is $Model^{i}$ (emissionangle, latitude) covering 90° S to 90° N. In the left panel of Figure 3, blue points show the brightness temperature model $T^{i}_{model}(Obs_{emiss}, Obs_{lat})$ for each observation point within one latitude bin at the beginning of this iteration, which is based on the input gridded model *Model* (*emissionangle, latitude*) (vellow curve). At the end of this iteration, we obtain updated brightness temperatures for each observation point $\tau_{best}^i(obs_{emiss}, obs_{lat}) = Residual^i(obs_{emiss}, obs_{lat}) + \tau_{model}^i(obs_{emiss}, obs_{lat})$ (see red points in left panel of Fig. 3). We introduce the factor $f = \frac{\tau_{best}^i(obs_{emiss}, obs_{lat})}{\tau_{model}^i(obs_{emiss}, obs_{lat})}$ (see blue points in right panel of Fig. 3). Within each latitude bin, we fit a smooth spline curve function func(emissionangle) to factor f with respect to the emission angle (Fig. 3, orange curve in right panel) and update the model with Modelⁱ⁺¹(emissionangle, latitude) = Modelⁱ(emissionangle, latitude) * func(emissionangle). In the left panel of Fig. 3, the green curve shows the updated brightness model, which will be used in the next iteration. Thus, we are able to find an updated Jupiter model that catches the main trend in T_{best}^i variation against emission angle, while minimizing the effect from longitudinal variations.



Figure 3. Example of model updating method 1 at an iteration *i* in one latitude bin of width 0.5° . Left panel: the yellow curve shows the input gridded model at the beginning of this iteration and the green curve is the updated model. Blue points show the brightness-temperature model for each observation and red points are the updated brightness at the end of this iteration. Right panel: the blue points show the value of factor *f* with respect to emission angle, which is then fit with the orange curve, a smooth spline curve function.

The iteration process is not sensitive to the initial guess of the brightness model. In Figure 4, we show the iterative results (residual values with respect to latitudes (left column) and emission angles (right column)) for PJ1 Channel 6 through 5 iterations as an example, starting with a globally uniform 1K brightness temperature, zero limb-darkening model, which has zero Jupiter brightness information. The right bottom panel in Fig. 4 shows the final residual values with respect to emission angle for all latitudes, which reveals how well the limb darkening value matches the real case. After 5 iterations, the tilting trend with respect to emission angle is within error bars determined by the instrument performance model. We used model updating method 1: "Interpolate and smooth among T_{best}^{l} with respect to emission angle" in this test case. Even starting with a model with no real Jupiter brightness temperature information, the residual values approach zero after 5 iterations and obvious structures are minimized in the final residual values with respect to latitude or emission angle. Therefore, we can conclude that the final brightness-temperature model is a good approximation to the real Jupiter brightness temperature according to MWR observations. It is consistent with observations at the 1- σ level, with all latitudes within expected error bars. Figure 5 shows the final best-fit model at four latitudes, where the initial models are 1K at all latitudes and emission angles. In Fig. 4f, we noticed that although the overall offset of the residuals is less than 2% (~ 3K) and the residuals approach the ~0.3K instrument noise at most latitudes after 5 iterations, at certain latitude bins (such as around 5°N to 20°N and around 35°N) the standard deviation is obviously larger than at other latitudes. This can also be seen in Fig. 5 where at 15°N, the brightness temperature vs. emission angle curve has more spread. This can be attributed to small longitudinal or temporal structures, which will be dealt with later using model update method 4 - "Fix limb-darkening values".

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Figure 4. Results from iterative residual analysis on PJ1 channel-6 observations. Panels a-f: residual values with respect to latitude after each iteration, which approach a mean of zero after 5 iterations. Panels g-l: the same residual values with respect to emission angle, which becomes flat after 5 iterations. From the instrument performance model [Janssen et al, 2017], the overall offset of the residuals should be less than 2%, which is ~ 3K. In the final iteration, the residuals approach the ~0.3K instrument noise at most latitudes, with exceptions which we attribute to spatial variations not accounted for in this longitudinally-averaged model.



Figure 5. The best-estimated brightness temperature $T_{best}(obs_{emiss}, obs_{lat})$ for observations in latitude bins (a) 30°S, (b) 0°, (c) 15°N, and (d) 30°N after 5 iterations (blue markers), starting with a uniform initial model (1K at all latitudes and all emission angles).

(2) Other model updating methods for longitudinally-uniform models

For comparison purposes, we introduced two additional model updating methods to derive brightness models with respect to latitudes and emission angles.

(2.1) Method 2: Fit to 3 coefficients

According to theoretical and empirical deduction, the Jupiter limb darkening can be approximated with the 3-coefficient equation [Oyafuso et al. 2020]:

$$T_B(\theta) = [A_0 + A_1(1 - \mu) + A_2(1 - \mu)^2] \cdot f(\theta)$$
(5)

where θ is the emission angle and $\mu = \cos(\theta)$. A_0 is the nadir brightness temperature, and A_1 and A_2 are the limb-darkening coefficients. $f(\theta)$ is an empirical angular profile suggested by Oyafuso et al. (2020) in order to account for the rapid brightness drop-off at larger emission angles [Figure.2 in Oyafuso et al. 2020]. At the end of each iteration and within each latitude bin, this method fits T_{best}^i with respect to μ for all the observations in that latitude bin to the equation above.

(2.2) Method 3: Spline interpolation

With this method, T_{best}^i is fitted with respect to emission angle with a spline function. Though MWR also observes Jupiter at high emission angles, due to the large main-beam coverage as projected onto the planet at high emission angles and significant synchrotron leakage into the beam in Channels 1 and 2, we only use observations with emission angles less than 60° ($\mu > 0.5$) to update the Jupiter model. The spline-interpolation method is able to match the brightness temperature variations at these small to medium emission angles, but fails to imitate the rapid drop off at emission angles close to 90°.

In Figure 6, we show examples for updating a brightness-temperature model at Channel 6 from iteration 1 to 2 (with respect to emission angle) at two latitude bins, 70°S and 0°, respectively, using the three different methods described above. The spline-fit results are not able to match the rapid drop of brightness temperature around 90° emission angle, while the 3-coefficient fit drops fastest there due to the imposed empirical angular profile.



Figure 6. Comparison between three model-updating methods when deriving a brightness-temperature model with respect to latitude and emission angle (Channel 6 data as an example). In each panel, red dots show the T_{model}^1 derived from input regularly gridded model (magenta curve) at the beginning of iteration 1 and blue dots show the T_{best}^1 which is then used to derive the updated model (green curve) for the next iteration.

(3) Method 4: Fix limb darkening.

As described earlier in Section 3, for typical Juno MWR orbits, each latitude was observed at multiple emission angles and within a small longitudinal range. In these cases, we iterate the process with model-update method 1 to achieve a longitudinal average (considering that the longitudinal range is rather narrow) while varying the limb darkening profile at all latitudes for each perijove, and then apply the limb-darkening profile and iterate with model-update method 4 to obtain the longitudinal variation in nadir brightness temperature. We fix the limb-darkening values, and only nadir brightness temperature at various latitudes are updated during these iterations. At the end of each iteration *i*, T_{best}^i is corrected to the nadir direction by:

$$T_{best_nadir}^{i}(0, Obs_{lat}, Obs_{lon}) = T_{best}^{i}(Obs_{emiss}, Obs_{lat}, Obs_{lon}) * \frac{T_{model}^{i}(0, Obs_{lat}, Obs_{lon})}{T_{model}^{i}(Obs_{emiss}, Obs_{lat}, Obs_{lon})}$$
(6)

 $T_{best,nadir}^{i}(0,0bs_{lat},0bs_{lon})$ is then used to update the Jupiter nadir brightness-temperature model on a regular grid. One specific use of this method is to study longitudinal structures and make 2D (latitude, longitude) Jupiter nadir brightness-temperature maps (see Figure 7 as an example for PJ7 Channel 6), which can be compared with visible upper atmosphere features in order to reveal deep structures connected to those features.

As for the special Juno cross-track orbits (such as PJ19), these perijove observations alone are not able to determine longitudinal structures and limb darkening values at the same time. Therefore, we take averaged limb-darkening values from previous perijoves (PJ1 to PJ9) as prior and iterate with model-updating method 4 to make 2D (latitude, longitude) maps [Bolton et al. 2020].



Figure 7. 2D (latitude, longitude) nadir brightness temperature map for PJ7 Channel 6 obtained after the iterative process with model updating method 4, which reveals small longitudinal structures, with a blow up map in latitude range 40° S to 40° N on the right.

4. Test Cases - Test Jupiter atmosphere models against MWR observations

To examine our residual-analysis process and to investigate the science results from it, we test a few previously published ammonia and water concentration profiles against the MWR observations. At microwave frequencies, ammonia, as the main opacity source, affects the brightness temperature most significantly, while on the other hand the water abundance would affect the atmospheric structure through latent heating and molecular weight effects. The behavior of residuals with respect to latitude and emission angle indicates how well a given profile matches the real Jupiter atmosphere and what modification should be made to better match the observations. Note that the variation in the MWR data due to the instrument noise should be quite small, but none of the models shown here attempt to take into account longitudinal or temporal variations. It is therefore not surprising that the residuals in the following test cases show more variation than instrument noise alone.

The bulk oxygen elemental ratio in Jupiter determines the water abundance in its atmosphere assuming a compositionally homogeneous envelope, and is key to discriminating among models for the origin of Jupiter and accretion processes in the planet-forming disk. The water abundance prior to Juno was measured by the Galileo Probe Mass Spectrometer (GPMS) as subsolar down to at least 19 bars [Wong et al.2004]. Wong et al. 2004 reported a mole fraction of H₂O of $(4.1 \pm 1.3) \times 10^{-5}$ at 11.0-11.7 bars and $(4.2 \pm 1.4) \times 10^{-4}$ at 17.6-20.9 bars,

358 which are both averages of data spanning the corresponding pressure range. However, the mixing ratio of H_2O in the 359 deep well-mixed atmosphere was not measured with GMPS and may be higher at deeper levels. Niemann et al. 1998 reported an upper limit to H_2O of 8×10^{-7} at 2.7 bar. More recently, ground-based infrared spectroscopy has constrained water abundance profiles in Great Red Spot (GRS) spectra as well as typical belt and zone spectra [Bjoraker et al. 2015, 2018]. The GRS spectrum requires a water cloud top at $p \ge 5$ bar, inducing a O/H lower limit of 1.1 × solar (corrected to the protosolar O/H ratio of Asplund et al. 2009) [Bjoraker et al. 2018]. They also found that the South Equatorial Belt (SEB) hot spot follows the H₂O profile observed by the Galileo Probe Mass spectrometer (GMPS) and becomes very dry above at P < 4.5 bars. The South Tropical Zone (STZ) has a saturated H₂O profile until reaches its cloud top between 4 and 5 bars [Bjoraker et al. 2018]. On the other hand, the water vapor profile retrieved from the Galileo probe Net-Flux Radiometer (NFR) measurements [Sromovsky et al. 1998] shows significantly sub-saturated water abundance at pressures greater than approximately 1.5 bars. Li et al. 2020 reported on the water abundance in the equatorial region, defined as from 0 to 4° north latitude, to be $2.5^{+2.2}_{-1.5} \times 10^{3}$ ppm, or $2.7^{+2.4}_{-1.6}$ times the protosolar oxygen elemental ratio to H, based on 1.25 to 22 GHz (1.4 to 24 cm) data from Juno MWR probing approximately 0.7 to 30 bars pressure. In Figure 8 we show various profiles of H_2O derived from ground-based, Galileo Probe, and Juno data. Arrows denote the water volume mixing ratio corresponding to O/H enrichments ranging from 0.01 to 10 times solar. Ground-based retrievals are shown for a $5-\mu m$ Hot Spot in the South Equatorial Belt and for the Great Red Spot, which follows a saturated profile above an opaque cloud at 5 bars. Water abundances derived from two investigations on the Galileo Probe are also shown. The continuous profile is from the Net Flux Radiometer; the points with error bars are from the Galileo Probe Mass Spectrometer. Finally, the recent H₂O measurement in the Equatorial Zone by Li et al. (2020) from Juno/MWR is shown in blue along with its error bars in light blue shadowed area. The difference between the derived H₂O abundances from Juno/MWR and Galileo is likely because the 5- μm Hot Spot Galileo probe entered was particularly dry, whereas the MWR measured H₂O in the Equatorial Zone.



Figure 8. Water abundance distribution profiles in and below the cloud condensation region for different assumed bulk abundances [Sromovsky et al. 1998, Bjoraker et al. 2015, Wong et al. 2004]. Blue line with shadowed area shows $2.7^{+2.4}_{-1.6}$ times the protosolar oxygen elemental ratio to H at 0° to 4° North [Li et al. 2020]. The black arrow shows the opaque water cloud at $p \ge 5$ bar at Great Red Spot (GRS), with a saturated H₂O profile above [Bjoraker et al. 2018]. The solar photospheric abundance is according to Asplund et al. 2009, adjusted to protosolar values.

Based on Juno/MWR PJ1 nadir brightness-temperature data, Li et al. (2017) (hereafter, Li17) published an ammonia concentration profile for pressures from 0.5 bar to approximately 200 bars, with a deep ammonia abundance of $\sim 3.5 \times 10^{-4}$. Their profile shows a depletion of ammonia over the specified latitude range down to a hundred bars and a column of high-concentration ammonia gas in the northern equatorial zone (EZ). We apply their ammonia profile in the latitudinal range from 50°S to 50°N in our residual analysis. As their ammonia profiles are presented in latitude bins of about 2°, a linear interpolation is used for intermediate latitudes. The ideal adiabatic temperature

395 profile (see Figure 9) used by Li17 was derived at Equatorial Zone and then applied to all other latitudes [Li et al. 396 2017]. The temperature profile depends on their retrieved deep ammonia and water abundance, with modification in 397 the upper troposphere that are constrained by mid-infrared observations [Fletcher et. al 2009]. The resulting residual 398 plots are shown in Figure 10 for MWR observations from PJ1 to PJ9 and channels 1 through 6. The panels on the 399 left show the residuals for their model in the latitudinal direction. Their model is a local model without large-scale 400 circulation. In channels 4, 5 and 6, although there are still latitudinal structures in the residuals for individual perijoves, when combining all the perijoves together, the mean residual values are near zero (excluding PJ7 402 observations near the Great Red Spot). The variation across perijoves in channels 4, 5, 6 suggest unmodeled 403 temporal or longitudinal structures. In channel 1 through 3 we see deviations around 10°N to 20°N, which shows up in all perijoves, suggesting that a different ammonia volume mixing ratio or temperature profile is needed at 10 bar and deeper at those latitudes.



Figure 9. Temperature profiles used in the test cases. Red: Ideal adiabatic temperature profile assuming 2 solar water abundance in the deep atmosphere. Blue and green: wet and dry adiabatic temperature profile proposed by de Pater et al. 2019. The inserted figure shows a blow-up of these temperature pressure profiles at 1-10 bars in linear scale of pressure.

With the limited knowledge of synchrotron radiation and ammonia/water opacities in the deep atmosphere observed by channel 1, matching the observations becomes especially difficult. On the other hand, a deviation of temperature from ideal adiabat that Li17 assumed is also a possibility. The panels on the right in Fig. 10 show how the residuals distribute with respect to emission angles, which is a very important indication for whether the input atmosphere profile generates the correct limb darkening. In Channels 3 to 6, all the distributions are rather flat and thus we see no systematic problem with the limb darkening. But in Channel 1 and 2, residuals deviate further from 0 at larger emission angles. There is potential confusion between effects from errors in the temperature profile or from greater synchrotron contribution to the antenna temperature at larger emission angles. However, differences in spacecraft pointing from orbit to orbit would lead to systematic, orbit-dependent trends in the high-emission angle residuals if synchrotron emission were the dominant source of error. Instead, the residuals point to errors at high emission angle that are consistent from orbit to orbit. Due to the high residuals at large emission angles, for channel 1 and 2, Li17 focused on data with emission angles smaller than 40°. A more detailed reexamination of the assumptions they made in their model would be helpful in better understanding the MWR data.

Soon after Li17, de Pater et al. 2019 analyzed VLA observations (hereafter, dePater19) between 3 to 37 GHz. Their longitude-smeared spectra show high NH₃ abundance of vmr 4.1×10^{-4} in the deep atmosphere (P > 8 - 10bars), decreasing at higher altitudes (see black line in Figure 13 for their profile at 6.5°N). As shown in Fig. 13, the deep ammonia abundance derived by dePater19 and Li17 are rather close, however, dePater19 shows ammonia 432 depletion only to 8-10 bars (at most 20 bars at NEB), while Li17 showed such depletion much deeper, persisting 433 down to 50-60 bars. The VLA observations were not able to see the atmosphere deeper than 10 - 20 bars because of strong synchrotron radiation in the foreground at longer wavelength. Both works show a high ammonia concentration in the northern Equatorial Zone (~ 0° - 5° N). Another big difference is that Li17's results show a unique ammonia concentration with slight increase with altitudes starting from about 7 bars to 2 bars and a relatively high abundance just below the NH₃ cloud layer, while the dePater19's ammonia concentration profile monotonically decreases with altitude, from well-mixed in the deep atmosphere to the ammonia cloud bottom. This pressure range is most sensitive in Juno/MWR Channels 3 through 6. To calculate the residuals of the dePater19 profile in the latitude range from 75° S to 75° N, we employed their dry and wet adiabatic temperature profiles (Fig. 9) and 4 times solar water abundance in the deep atmosphere as they proposed [de Pater et al. 2019]. As we compared the residual results from the dePater19 model (Figure 11) with the residual results from Li17 (Fig. 10), in channel 5 and 6, which is sensitive to P < 4 bars and most sensitive around 1 -2 bars for channel 5, 0.5 - 0.6 bars for channel 6, both dePater19 and Li17 fit the MWR observations from PJ1 to PJ9 comparatively well at low latitudes. However, in channel 5 dePater19 starts to deviate from MWR data from mid- to high- latitudes. In Figure 12, we show the mean and standard deviation of residuals in 2° latitude bins (left column) and 2° emission angle bins (right column) from both Li17 (orange curve) and dePater19 (blue curve). The values are derived from residuals in Fig.10 for Li17 and in Fig.11 for dePater19 combining all PJ1 to PJ9 results. In Channel 6, at low- to mid- latitudes, both Li17 and dePater19 have mean residuals close to zero while both miss some small latitudinal structures. In Channel 5, both models match the observations at low latitudes, while dePater19 misses more small latitudinal structures. At mid- to high-latitudes, Li17 matches the observations better than dePater19 in channel 5. The mean residuals with respect to emission angles are flatter for the Li17 model, with an offset less than 1K for channel 6 and 5 within the 2% absolute calibration uncertainty. However, there is a downward tilt at large emission angle for dePater19 in Channel 6, suggesting a larger limb darkening is required by the observations. In Channel 4, which is sensitive to P < 10 bars and especially around 3 - 5 bars, both dePater19 and Li17 models match the observations from 20S to 10N. However, outside this latitude, residuals from dePater19 start to deviate from zero. dePater19 residuals are significantly off-zero in Channel 3, which is sensitive to P < 40 bars and most sensitive around 8 - 10 bars. For channel 1 to 3, the residual vs emission angle panel for dePater19 is tilted downwards at larger emission angles, suggesting the limb darkening in dePater19 is rather small compared to the real Jupiter atmosphere. At the same time, Li17 residuals in channel 3 still remain around 0 and the limb darkening also matches the data well. As we go deeper into the atmosphere with Channels 1 and 2, Li17 residuals are significantly closer to zero as compared to the dePater19 residuals. In Channels 1, 2 and 3, dePater19 residuals are far above 0, which indicates that the MWR observed brightness temperature is much higher than predicted by their model and could be a sign of too much ammonia concentration in their model at P > 10 bars, where VLA (dePater19) were not able to observe.

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Residual (K)



Figure 10. Results for residual analysis on ammonia distribution from Li et al. 2017 using PJ1 to PJ9 observations from channel 1 to channel 6. Each perijove is plotted with different colors; blue: PJ1; orange: PJ3; green: PJ4; red: PJ5; purple: PJ6; brown: PJ7; pink: PJ8; grey: PJ9. The panels on left show the residuals of their model in the latitudinal direction; the panels on the right show how the residuals distribute with respect to emission angle.

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Figure 11. Results for residual analysis on ammonia distribution from de Pater et al. 2019 using PJ1 to PJ9 observations from channel 1 to channel 6. The same color code is used as given in the caption for Fig. 10. The panels on left show the residuals their model in the latitudinal direction; the panels on the right show how the residuals distribute with respect to emission angle.

Residual (K)



Figure 12. Comparison of mean and standard deviation values of residuals in 2° latitude bins (left column) and 2° emission angle bins (right column) between ammonia distributions from Li. et al. 2017 (orange) and de Pater et al. 2019 (blue). The mean latitudes and emission angles in each latitude bin have been shifted 0.5° between the two models in the plot for clearer view.

Before the arrival of the Juno spacecraft at Jupiter, Galileo Probe, which entered at a latitude of 6.5° N (the southern edge of NEB), provided the only knowledge of Jupiter's deep atmosphere profile (i.e. temperature, water and ammonia abundances below the water cloud base). The Galileo Probe Mass Spectrometer (GPMS) estimated an NH₃ VMR of $5.7 \pm 2.2 \times 10^{-4}$ in the 8.9–11.7 bar region [Wong et al. 2004], and Galileo probe-to-orbiter signal attenuation [Folkner et al. 1998] interpreted with updated ammonia opacities [Hanley et al. 2009] found a maximum

NH₃ VMR of 8.4 \pm 0.6 \times 10⁻⁴ at 9.7 bar (with lower NH₃ VMR at both higher and lower pressures along the probe entry path). The net flux radiometer (NFR) inferred a NH₃ VMR of 2.5×10^{-4} at P > 5 bars, dropping to about 1.5×10^{-4} at 2.5 bars and then decreasing more rapidly with altitudes [Sromovsky et al., 1998]. de Pater et al. 2001 presented disk-averaged brightness temperature from Very Large Array (VLA) observations showing ammonia decreasing at pressure $P \leq 4$ bar and a global depletion of ammonia in the region around 2 bars, reaching subsolar (\lesssim 0.5). All derived ammonia VMR profiles from all of these works are summarized in Fig. 13, together with the ammonia concentration profiles of Li17 and dePater19 at 6.5°N for comparison. Except for Li17 that provides ammonia distribution down to thousands of bars, all other ammonia profiles are considered reaching uniform deep ammonia abundance below ~20 bars. There are a few cases where two lines have the same color and style, indicating estimates of the upper- and lower- limits: the two blue dashed lines indicate the analysis of microwave data by de Pater and Massie (1985), which show NH₄SH and NH₃ cloud formation at about 1.4 and 0.5 bars with two UV photolysis effect limits in the stratosphere; the yellow dot-dashed lines represent the NH₃ profile inside (lower mixing ratio) and outside (higher mixing ratio) a hot spot from Fouchet et al. 2000 ISO-SWS observations; the cyan dotted lines show the low- and high- limits from the Galileo probe radio due to the uncertainty in their observations [Folkner et al., 1998]. In these test cases, we treat as if these ammonia abundances are globally representative, which is only true in the disk-averaged VLA results [de Pater and Massie, 1985].



Figure 13. Ammonia abundance distribution profiles from all previous observations around 6.5° N. Except for Li17, all other ammonia profiles are considered reaching uniform deep ammonia abundance below ~20 bars.

We found that the deep ammonia abundance in Li17 is very close to the ISO/SWS value outside the hot spot and the disk-averaged VLA results from de Pater and Massie (1985). When calculating the residuals for these profiles, we used the ideal adiabatic temperature profile derived by Li17 at Equatorial Zone. For each ammonia profile, we tested a range of possible water-concentration profiles as shown in Fig. 8. We show here only the results for a deep water abundance of 2 times solar value. Results for the other water profiles look similar. All of the residuals are plotted in Figure 14 (a) to (f) using MWR PJ1 data. Since only the region around 6.5°N is valid for most of these profiles, we have limited the range of latitudes from 0°N to 17°N. For comparison, we also plot the residuals for Li17 and dePater19. As we focus at 6.5°N, in the deep atmosphere (Channels 1 and 2), the de Pater and Massie 1985, ISO/SWS and Li17 models fit the MWR antenna temperature best, due to their fairly close deep ammonia

abundances. The spread of residuals at one latitude for certain models indicates a limb-darkening mismatch. In Channels 3, 4, and 5, the disk-averaged VLA results (de Pater and Massie, 1985), Li17 and dePater19 models all result in the smallest residuals at 6.5°N. ISO/SWS residuals start to deviate far from zero at channel 4 due to their ammonia condensation at relatively high pressures in their abundance profile. In channel 6, at the troposphere, all models coincide with the MWR observation fairly well except the Galileo Probe NFR result and the disk-averaged VLA [de Pater and Massie, 1985] result with the upper-limit estimate of the effect of UV photolysis on the profile in the stratosphere. Despite the difference in the ammonia profile between disk-averaged VLA (de Pater and Massie, 1985) result with the lower-limit estimate and Li17, they are both consistent with the MWR observations at 6.5°N for channels 1 - 6. However, due to the finite size of the MWR main beam, a good-fit at one single latitude doesn't guarantee that the model is close to the real Jupiter atmosphere. Even when the MWR beam is centered at 6.5° N, some of the received emission originates at other latitudes.



Figure 14. Results for residual analysis on previous published ammonia distributions as listed in Fig.13 using PJ1 535 observations around 6.5°N from channel 6 (a) to channel 1 (f). For each channel, residuals from all 11 models are 536 plotted in two panels for easier viewing. 537

8 10 atitude (degree)

15 10

-10 -1-20 20 10 -10 -20 30 20 10 -50 60 40 20

539 5. Conclusions

We provide a direct method to compare any Jupiter atmosphere model against the MWR observations through residual analysis. By investigating the residual values at various latitudes, emission angles and channels (sensitive to conditions at pressures between ~ 1 and 100 bars), we are able to tell which part of the Jupiter model deviates from the observations and often what kind of modification is required. This technique may not provide a unique solution, but is useful to identify flaws in existing models. For the two latitudinally resolved models we investigated (Li17 and dePater19), they are mostly consistent with MWR observations at low- to mid- latitudes from Channel 4 to 6, with overall mean residual less or close to the absolute calibration error ~2%. However, our residual analysis from Channel 1 to 3, where ground-based VLA was not able to observe, favors Li17's ammonia profile with ammonia depletion persisting deeper down to 50-60 bars and with a slightly smaller deep ammonia abundance, while dePater19 overall mean residuals start to be over 4% and reach as high as 12% in Channel 1. Additionally, the iterative process is able to generate Jupiter brightness temperatures with respect to emission angles, latitudes and longitudes that best match the MWR observed antenna temperatures. The 2D Jupiter maps with respect to latitude and longitude at all MWR observed pressure levels are of particular interest, as they will be very useful when comparing with upper-atmosphere observations at UV, near-IR and mid-IR in order to reveal the deep structures of upper atmosphere features [Fletcher et al. 2020].

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https://pds-atmospheres.nmsu.edu/data_and_services/atmospheres_data/JUNO/microwave.html

Data files can be found at:

https://pds-atmospheres.nmsu.edu/PDS/data/jnomwr_1100/data_calibrated/.

All datasets for this research are available in the citation reference: Zhang, Zhimeng, Adumitroaie, Virgil, Allison, Michael, Arballo, John, Atreya, Sushil, Bjoraker, Gordon, ... Wong, Michael. (2020). Dataset for Residual Study: Testing Jupiter Atmosphere Models Against Juno MWR Observations [Data set]. Zenodo. http://doi.org/10.5281/zenodo.3936065

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Figure1.



Figure2.



Figure3.



Figure4.





Figure5.



Figure6.



Figure7.





Figure9.



Figure10.



Figure11.



Figure12.



Figure13.



Figure14.









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PJ1 Channel 6



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