

# ADVANCED ELECTRONIC MATERIALS

## Supporting Information

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Controlling Defect Formation of Nanoscale AlN: Toward  
Efficient Current Conduction of Ultrawide-Bandgap  
Semiconductors

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**Section 1. Theoretical computational details**

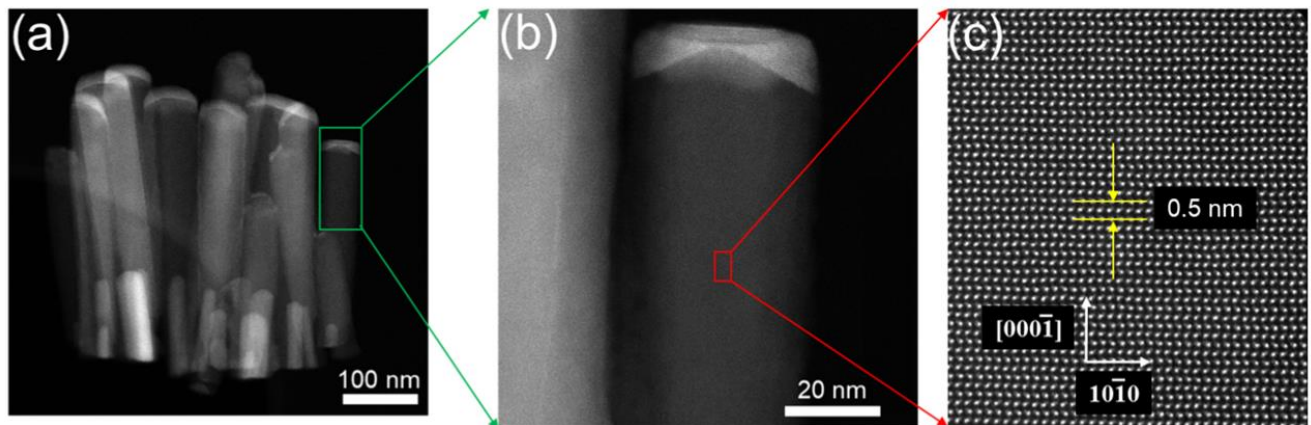
To calculate the defect formation energy in AlN, first-principles density functional theory calculations were performed with the projector-augmented wave method [1, 2] as implemented in the Vienna ab initio Simulation Package (VASP). [3] PBE pseudopotentials (PAW\_PBE Al\_GW and N\_GW\_new) [4] were used for Al and N, with a plane-wave cutoff energy of 500 eV to converge the total energy to 1 meV/atom. The HSE06 hybrid density functional [5] with a mixing parameter of 0.33 was employed to obtain a band gap of 6.2 eV for AlN, in close agreement with the experiment. [6] Defect calculations were performed for AlN using 96-atom orthorhombic supercells [7] with a  $2 \times 2 \times 2$   $\Gamma$ -centered Brillouin-zone-sampling grid. All structures were relaxed by allowing atoms to move until the force on ions is less than 0.02 eV/Å, with spin-polarization included for unpaired electrons. The defect formation energies were calculated by the following equation:

$$E^f(X^q) = E_t(X^q) - E_t(\text{AlN}) + \sum_i n_i \mu_i + q(E_F + E_{VBM}) + E_{corr}.$$

$E_t(X^q)$  denotes the total energy of the supercell with a defect X in charge state  $q$  and  $E_t(\text{AlN})$  is the total energy of the perfect supercell.  $n_i$  is the number of atoms that the supercell exchanges with the chemical reservoir and  $\mu_i$  is the chemical potentials.  $E_F$  is the Fermi energy referenced

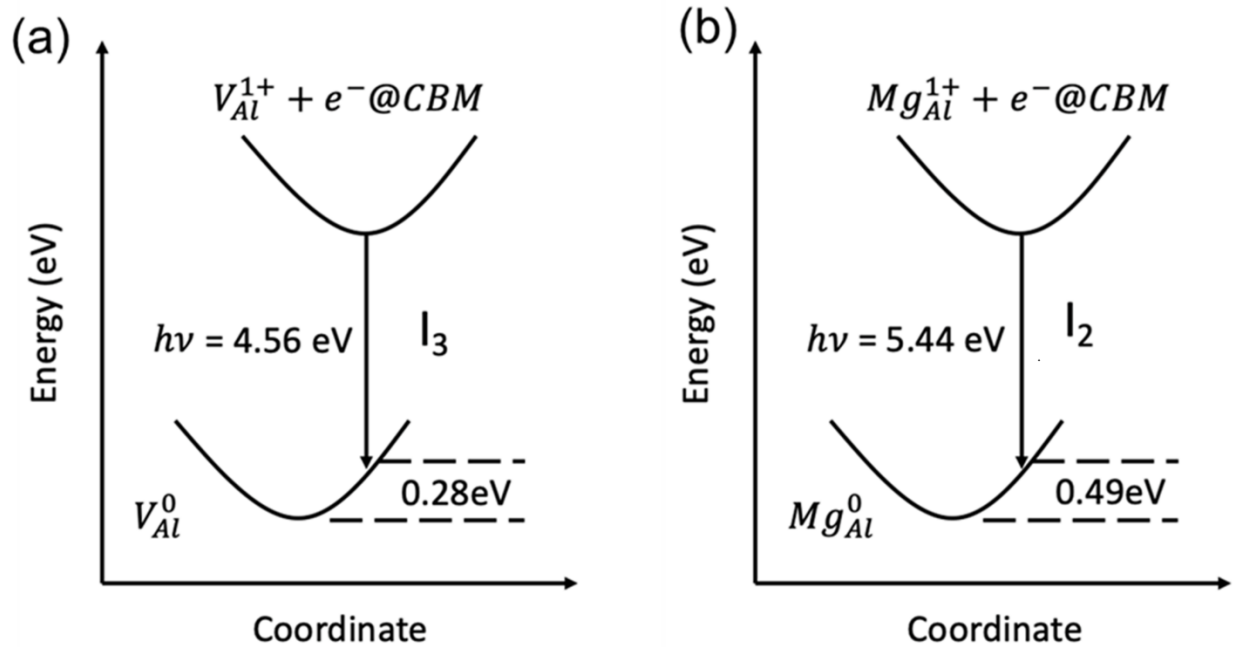
to the valence band maximum of the AlN.  $E_{corr}$  accounts for the correction to the artificial interaction between charged defect and its image charge. Here, we adopt the approach of Freysoldt *et al.* to evaluate the correction. [8] We assume Mg-rich conditions and  $\mu_{Mg}$  is bounded by the formation of secondary  $Mg_3N_2$  compound. Our results on defect formation energies and charge-transition levels in AlN are in good agreement with previous theoretical calculations. [9, 10]

## Section 2. STEM images of as-grown AlN nanowire *p-i-n* structure



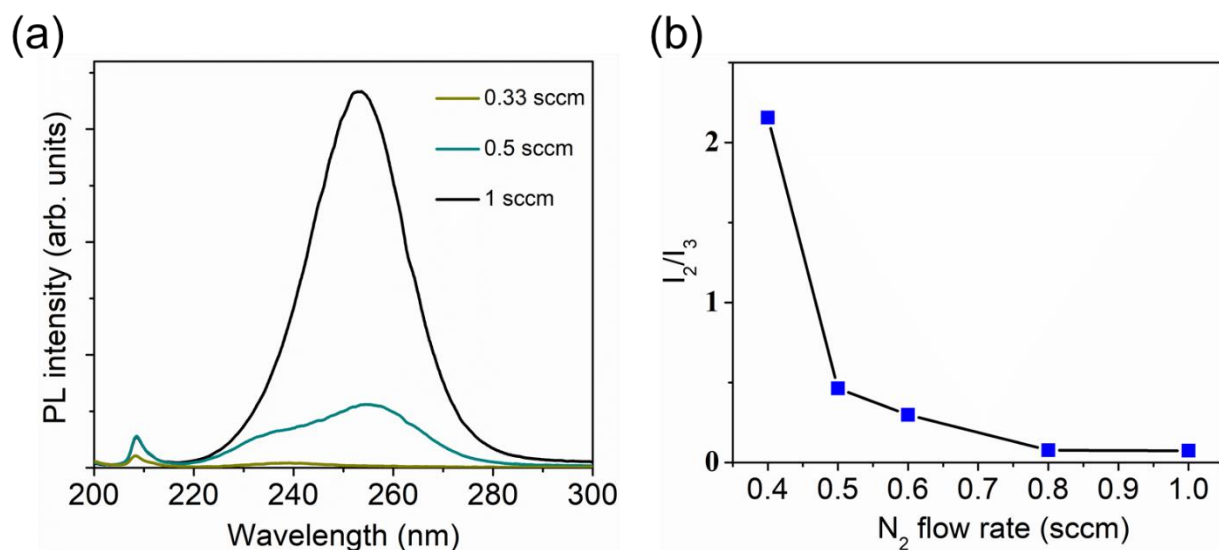
**Figure S1.** Microstructures of AlN nanowires: (a) Low-magnification STEM-HAADF image showing several AlN nanowires; (b) Magnified STEM-HAADF image shows that the top end of the nanowire outlined in (a) has a thin AlGaN as contact layer; (c) High-resolution STEM-HAADF image taken from the outlined i-AlN active region in (b) showing crystal structure free from dislocations and stacking faults.

**Section 3. Configuration-coordinate diagram illustrating optical processes related to Mg dopant and Al vacancy in AlN**



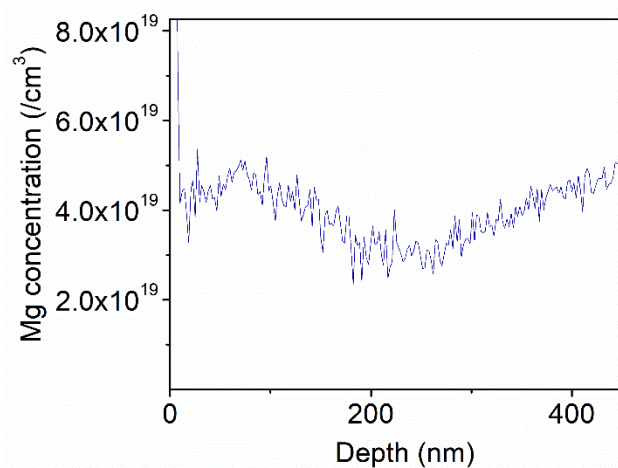
**Figure S2.** Configuration-coordinate diagram illustrating optical processes related to Mg dopant and Al vacancy in AlN. (a) Recombination of an electron at the CBM with  $V_{Al}^{1+}$  to form  $V_{Al}^0$  leads to an emission peak at 4.56 eV, explaining origin of  $I_3$ . (b) Recombination of an electron at the CBM with  $Mg_{Al}^{1+}$  to form  $Mg_{Al}^0$  leads to an emission peak at 5.44 eV, explaining origin of  $I_2$ .

**Section 4. Non-normalized photoluminescence spectra and intensity ratio of  $I_2$  to  $I_3$  versus the  $N_2$  flow rate**



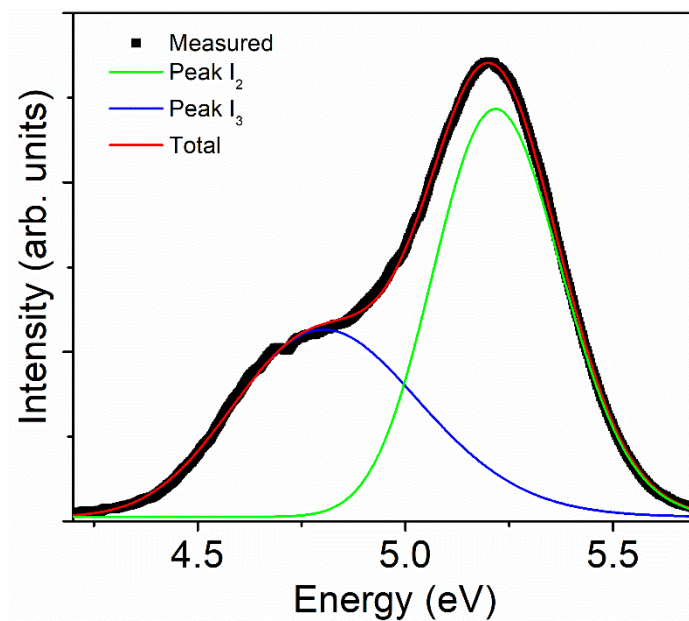
**Figure S3.** (a) Non-normalized photoluminescence spectra of samples from Group B. All samples were grown at a substrate temperature of 810 °C and Mg BEP of  $3 \times 10^{-9}$  Torr while the nitrogen flow rate was varied among different samples. (b) Intensity ratio of  $I_2$  to  $I_3$  versus the  $N_2$  flow rate derived from photoluminescence spectra of samples from Group B.

**Section 5. Eurofins EAG Materials Science secondary ion mass spectrometry (SIMS) depth profile of Mg-doped AlN nanowires grown on Si substrate**



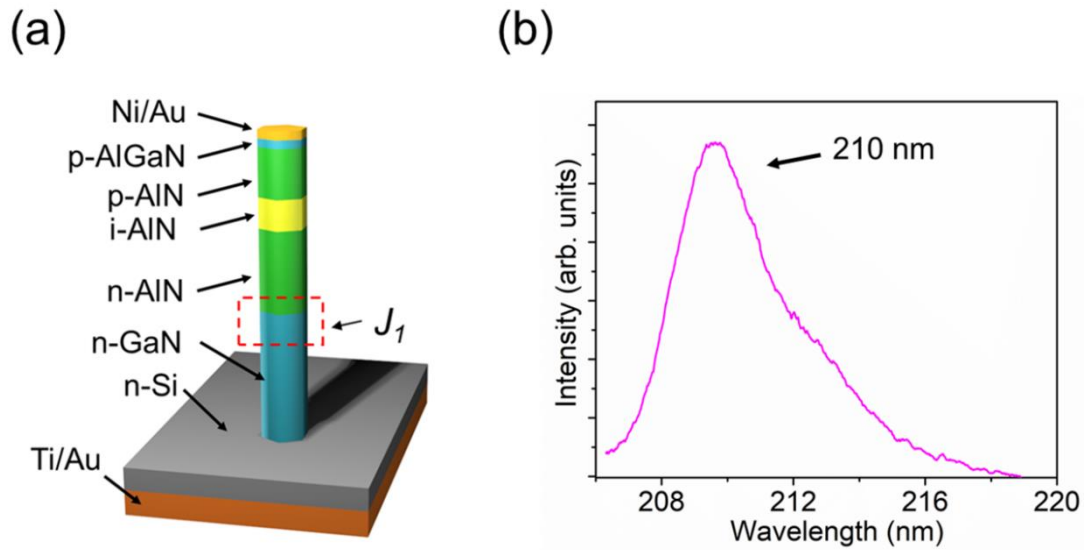
**Figure S4.** SIMS depth profile of Mg-doped AlN nanowires grown on Si substrate at 865 °C substrate temperature, 1 sccm N<sub>2</sub> flow rate and 7×10<sup>-9</sup> Torr Mg BEP. A fill factor of 30% for the nanowires was used to derive the average Mg concentration.

**Section 6. Photoluminescence spectrum and Gaussian profile analysis of Mg-doped AlN nanowires**

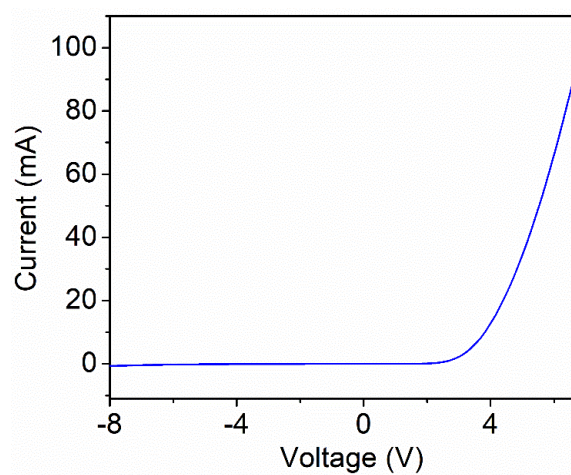


**Figure S5.** Photoluminescence spectrum and Gaussian profile analysis of Mg-doped AlN nanowires measured under  $0.78 \text{ W/cm}^2$  excitation power. The sample was grown under a substrate temperature of  $865 \text{ }^\circ\text{C}$ ,  $\text{N}_2$  flow rate of  $1 \text{ sccm}$  and Mg BEP of  $7 \times 10^{-9} \text{ Torr}$ .

## Section 7. Schematic and electroluminescence spectrum of AlN nanowire-based LED



**Figure S6.** (a) Schematic of AlN nanowire based LED on Si(111) substrate.  $J_1$  points to the heterojunction between  $n$ -GaN and  $n$ -AlN. (b) Room-temperature electroluminescence spectrum measured under 30 mA current injection from as-fabricated AlN LEDs.

Section 8. Current-voltage characteristics of  $n$ -GaN/ $n$ -AlN unipolar heterojunction



**Figure S7.** *I-V* characteristics of *n*-type doped unipolar GaN/AlN heterojunction. 80 nm Ti/20 nm Au was used for metal contact with the nanowires.

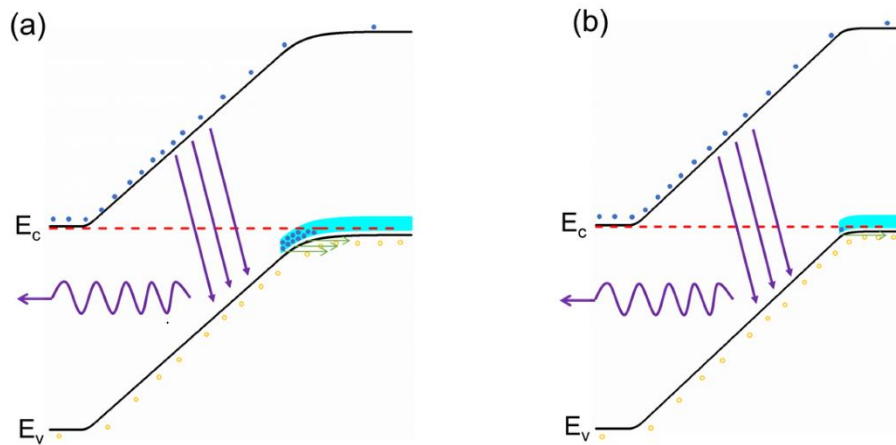
**Table II:** Minimum ideality factor and corresponding voltages of AlN nanowire LEDs with different Mg concentrations in the *p*-AlN layer.

	Sample I	Sample II	Sample III
Mg concentration [cm <sup>-3</sup> ]	1×10 <sup>19</sup>	2×10 <sup>19</sup>	6×10 <sup>19</sup>
Minimum ideality factor	14.9	4.6	3.58
Corresponding voltage [V]	3.6	2.4	2.4

### Section 9: Analysis of the charge carrier (hole) tunneling process

The hole tunneling process was further explained by calculating the energy band diagram of the AlN *p-i-n* homojunction using two different Mg concentrations as shown in Figure S8. For Mg concentration of 1×10<sup>19</sup>/cm<sup>3</sup>, the calculated depletion width into the *p*-AlN region is ~27 nm. Under equilibrium, despite the Fermi-level is flat throughout the *p-i-n* structure, the built-in potential in the *p*-AlN region causes the band structure as well as the impurity band to bend downwards, wherein the deep level energy states below the Fermi-level will be filled by electrons. Under forward bias, these captured electrons can tunnel from one state to adjacent empty states in the valence band, contributing to the tunneling component in Eq. 1 and increasing the ideality factor. When  $N_A$  increases to 6×10<sup>19</sup>/cm<sup>3</sup>, hole carrier concentration in the valence band increases due to higher Mg concentration and broadened Mg impurity band, leading to a significantly reduced depletion width into *p*-AlN of ~8 nm and a suppressed band bending of *p*-AlN as shown in Figure S8b. As the number of available electrons inside the impurity band reduces, the electron tunneling process is significantly suppressed. It's worth

noting that the deep-level assisted tunneling process will also be suppressed at higher measurement temperature as shown in Figure 4d, wherein hole concentration increases due to a larger portion of ionized Mg-impurities.



**Figure S8.** Schematic illustration of the energy band diagram of AlN nanowire *p-i-n* junction with Mg concentration of (a)  $1 \times 10^{19}/\text{cm}^3$  and (b)  $6 \times 10^{19}/\text{cm}^3$ . Mg impurity band bending and deep level assisted electron tunneling process are also illustrated.

## Supplementary Reference

- [1] P. E. Blöchl, *Phys. Rev. B* **50**, 17953 (1994)
- [2] G. Kresse, D. Joubert, *Phys. Rev. B* **59**, 1758 (1999)
- [3] G. Kresse, J. Furthmüller, *Phys. Rev. B* **54**, 11169 (1996)
- [4] J. P. Perdew, K. Burke, M. Ernzerhof, *Phys. Rev. Lett.* **77**, 3865 (1996)
- [5] J. Heyd, G. E. Scuseria, M. Ernzerhof, *J. Chem. Phys.* **118**, 8207 (2003)
- [6] H. Yamashita, K. Fukui, S. Misawa, S. Yoshida, *J. Appl. Phys.* **50**, 896 (1979)
- [7] C. G. Van de Walle, S. Limpijumnong, J. Neugebauer, *Phys. Rev. B* **63**, 245205 (2001)
- [8] C. Freysoldt, J. Neugebauer, C. G. Van de Walle, *Phys. Status Solidi B* **248**, 5, 1067–1076 (2011)
- [9] J. L. Lyons, A. Janotti, C. G. Van de Walle, *Phys. Rev. Lett.* **108**, 156403 (2012)
- [10] Q. Yan, A. Janotti, M. Scheffler, C. G. Van de Walle, *Appl. Phys. Lett.* **105**, 111104 (2014)