

# Kinetic-Scale Turbulence in the Venusian Magnetosheath

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This is the author manuscript accepted for publication and has undergone full peer review but has not been through the copyediting, typesetting, pagination and proofreading process, which may lead to differences between this version and the [Version of Record](#). Please cite this article as [doi: 10.1029/2020GL090783](https://doi.org/10.1029/2020GL090783).

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## Abstract

While not specifically designed as a planetary mission, NASA’s Parker Solar Probe (PSP) mission uses a series of Venus gravity assists (VGAs) in order to reduce its perihelion distance. These orbital maneuvers provide the opportunity for direct measurements of the Venus plasma environment at high cadence. We present first observations of kinetic scale turbulence in the Venus magnetosheath from the first two VGAs. In VGA1, PSP observed a quasi-parallel shock,  $\beta \sim 1$  magnetosheath plasma, and a kinetic range scaling of  $k^{-2.9}$ . VGA2 was characterised by a quasi-perpendicular shock with  $\beta \sim 10$ , and a steep  $k^{-3.4}$  spectral scaling. Temperature anisotropy measurements from VGA2 suggest an active mirror mode instability. Significant coherent waves are present in both encounters at sub-ion and electron scales. Using conditioning techniques to exclude these electromagnetic wave events suggests the presence of developed sub-ion kinetic turbulence in both magnetosheath encounters.

## 1 Introduction

Astrophysical environments are often characterized by nonlinear turbulent processes, which transfer energy from large fluid-like scales to kinetic dissipative scales. The relative accessibility of space-plasma environments has driven our understanding of these universal processes (Bruno & Carbone, 2005; Chen, 2016; Verscharen et al., 2019). While properties of large scale magnetohydrodynamic (MHD) turbulence have been studied since the earliest days of space exploration (Coleman, 1968; Matthaeus & Goldstein, 1982), relatively recent advancements in instrumentation have enabled analysis of kinetic scale turbulence (Leamon et al., 1998; Alexandrova et al., 2012; Chen & Boldyrev, 2017).

Evidence for kinetic scale plasma-turbulence largely stems from observations of the terrestrial magnetosphere and solar wind. At ion kinetic scales magnetic spectra steepen, due to some combination of dispersive and dissipative effects, leading to a sub-ion scale energy cascade (Alexandrova et al., 2008; Sahraoui et al., 2009; Alexandrova et al., 2009; Sahraoui et al., 2010; Alexandrova et al., 2012). Kinetic spectra with approximate  $k^{-2.7}$  scaling characterize the solar wind at 1 AU and the inner heliosphere (Sahraoui et al., 2009; Chen et al., 2010; Alexandrova et al., 2012; Sahraoui et al., 2013; Bowen, Mallet, Bale, et al., 2020). The observed steepening is consistent with the dispersion of Alfvénic to kinetic Alfvén wave (KAW) turbulence alongside some intermittency or dissipation (Schekochihin et al., 2009; Boldyrev & Perez, 2012a; Howes et al., 2011; Chen et al., 2013; Franci et al., 2015, 2016). At electron kinetic scales, further spectral steepening is measured (Alexandrova et al., 2009, 2012; Sahraoui et al., 2013; Huang et al., 2014; Chen & Boldyrev, 2017).

Kinetic scale steepening in Earth’s magnetosphere (Dudok de Wit & Krasnoselkikh, 1996; Czaykowska et al., 2001) is likely connected to magnetospheric heating (Sundkvist et al., 2007); however the shape and spectral scaling of magnetospheric turbulence is a topic of significant debate. Commonly observed inertial range turbulence, with approximate Kolmogorov-like  $k^{-5/3}$  scaling, is not universally present in the terrestrial magnetosphere (Czaykowska et al., 2001; Alexandrova et al., 2008); a common interpretation is that shock structure may prevent the formation of fluid scale turbulence in the magnetosheath (Vörös, Zhang, Leubner, et al., 2008; Huang et al., 2017; Chhiber et al., 2018). However, instabilities may serve as a source of turbulent and nonlinear fluctuations, which may vary the inertial range spectrum (Sahraoui et al., 2006). Kinetic range spectra observed in the terrestrial magnetosphere are similar to the 1 au solar wind, consistent with KAW turbulence (Alexandrova et al., 2008; Huang et al., 2014; Chen & Boldyrev, 2017). However, variation in kinetic range scaling of magnetosheath spectra has been reported (Rezeau et al., 1999; Alexandrova et al.,

2008; Huang et al., 2014), possibly attributable to intermittency (Alexandrova, 2008; Boldyrev & Perez, 2012b; Zhao et al., 2016), or dissipation (Howes et al., 2011).

Knowledge of kinetic scale processes of extraterrestrial magnetospheres is limited by the resources required for distant space-missions. Saur (2004) suggest that turbulent dissipation is significant to heating Jupiter’s magnetosphere. Saturn’s magnetosphere has kinetic turbulence with scalings similar to that observed at Earth and inferred turbulent dissipation rates that can account, for magnetospheric heating (von Papen et al., 2014). Hadid et al. (2015) suggest that perpendicular shock geometry may prevent formation of an inertial range at Saturn, though kinetic scales are largely invariant behind both quasi-parallel and quasi-perpendicular shocks. Observations from Jupiter, reveal similar properties such as spectral steepening at kinetic scales, and the lack a  $k^{-5/3}$  inertial range (Tao et al., 2015).

Kinetic-scale turbulence is also observed in the magnetospheres of Mars and Mercury. Uritsky et al. (2011) study kinetic scale turbulence in Mercury’s magnetosphere, observing a fluid-kinetic break, and steep anomalous scaling of inertial range fluctuations, attributed to finite Larmor radius (FLR) effects; the authors highlight potential ion-scale instabilities and the presence of coherent electron scale waves. Huang et al. (2020) suggest that no inertial range forms in Mercury’s magnetosheath, and that heavy exospheric ions contribute to deviation from canonical  $k^{-5/3}$  spectra. Ruhunusiri et al. (2017) demonstrate that spectral energy scaling of turbulence near Mars is well ordered by magnetospheric structure: shallow inertial range spectra are found in the magnetosheath, though kinetic range turbulence seems developed; solar wind-like inertial range and kinetic spectra are observed near the magnetic pileup region, suggesting turbulent processing.

Parker Solar Probe utilizes resonant orbital encounters with Venus to reduce its perihelion altitude (Fox et al., 2016), providing an opportunity for detailed observations of kinetic scale turbulence in the Venusian magnetosphere. At closest approach, PSP will fly within 400 km of the Venusian surface, placing it within Venus’s ionosphere (Zhang et al., 2007; Futaana et al., 2017). Though not designed specifically to study the Venus plasma environment, PSP shares technological heritage with modern magnetospheric missions (McFadden et al., 2008; Wygant et al., 2013; Kletzing et al., 2013). Observations made by PSP during these encounters promise to contribute significantly to understanding the planet’s magnetosphere.

Nonlinear waves and MHD turbulence in Venusian plasma have been studied previously. Vörös, Zhang, Leubner, et al. (2008) demonstrate intermittent turbulence in the Venusian wake and magnetosheath. Based on observations of shallow spectra with Gaussian fluctuations, Vörös, Zhang, Leubner, et al. (2008) suggest that MHD turbulence may not develop uniformly throughout the magnetosphere, in agreement with observations from other planetary environments (Czaykowska et al., 2001; Hadid et al., 2015; Ruhunusiri et al., 2017; Huang et al., 2017; Chhiber et al., 2018). Xiao et al. (2018) show that shock geometry is important in shaping the inertial range, with developed  $k^{-5/3}$  spectra appearing more readily behind quasi-parallel shocks. Xiao et al. (2020) additionally show that day/night asymmetry strongly affects the development of inertial scale turbulence. Many inertial scale nonlinear waves, instabilities, and vortices have been reported near Venus, which are potential drivers of turbulence (Wolff et al., 1980; Amerstorfer et al., 2007; Balikhin et al., 2008; Pope et al., 2009; Walker et al., 2011; Golbraikh et al., 2013; Volwerk et al., 2016; Futaana et al., 2017).

There are relatively few kinetic scale observations of fluctuations at Venus. Dwivedi et al. (2015) suggest that a break exists between MHD and kinetic ranges, and that anomalous inertial range scaling is possibly due to mirror mode structures generated through temperature anisotropy. The authors suggest that kinetic scale fluctuations may be a combination of nonlinearly interacting kinetic turbulence with instability

124 driven modes; however the observations are limited by the 1 Hz magnetometer resolu-  
 125 tion. Kinetic scale wave phenomenon have been studied in detail; with much focus on  
 126 the Venusian ionosphere (Russell et al., 2013). High frequency, electron scale waves,  
 127 likely generated through plasma instabilities, have been well documented in the fore-  
 128 shock, upstream solar wind, and magnetosheath (Strangeway, 2004). Ion scale waves  
 129 have been identified both upstream and downstream the bow shock (Russell et al.,  
 130 2006; Delva et al., 2015).

131 Here, we study signatures of kinetic scale turbulence in the Venusian magne-  
 132 tosheath. We demonstrate differences in spectral energy scalings in the kinetic range,  
 133 likely due to bow-shock geometry, plasma  $\beta$ , and the presence of the mirror insta-  
 134 bility. In addition to kinetic scale turbulence, the sub-ion and electron scales in the  
 135 magnetosheath are characterized by significant wave activity (Page, 2020). The use  
 136 of conditioning (Sorriso-Valvo et al., 1999; Kiyani et al., 2006; Chen et al., 2014) to  
 137 exclude coherent sub-ion scale waves reveals that despite significant differences in spec-  
 138 tral scaling signatures of a developed kinetic cascade are present in both encounters.  
 139 At electron scales the spectrum further steepens, similar to observations from Earth’s  
 140 magnetosphere (Huang et al., 2014; Chen & Boldyrev, 2017).

## 141 2 Data

142 We implement measurements from the electromagnetic FIELDS instrument (Bale  
 143 et al., 2016) as well as the Solar Wind Electron Alpha and Proton (SWEAP) investiga-  
 144 tion (Kasper et al., 2016) during PSP’s first two Venus gravity assists (VGA1 occurring  
 145 Oct 31, 2018 and VGA2 on Dec 26, 2019).

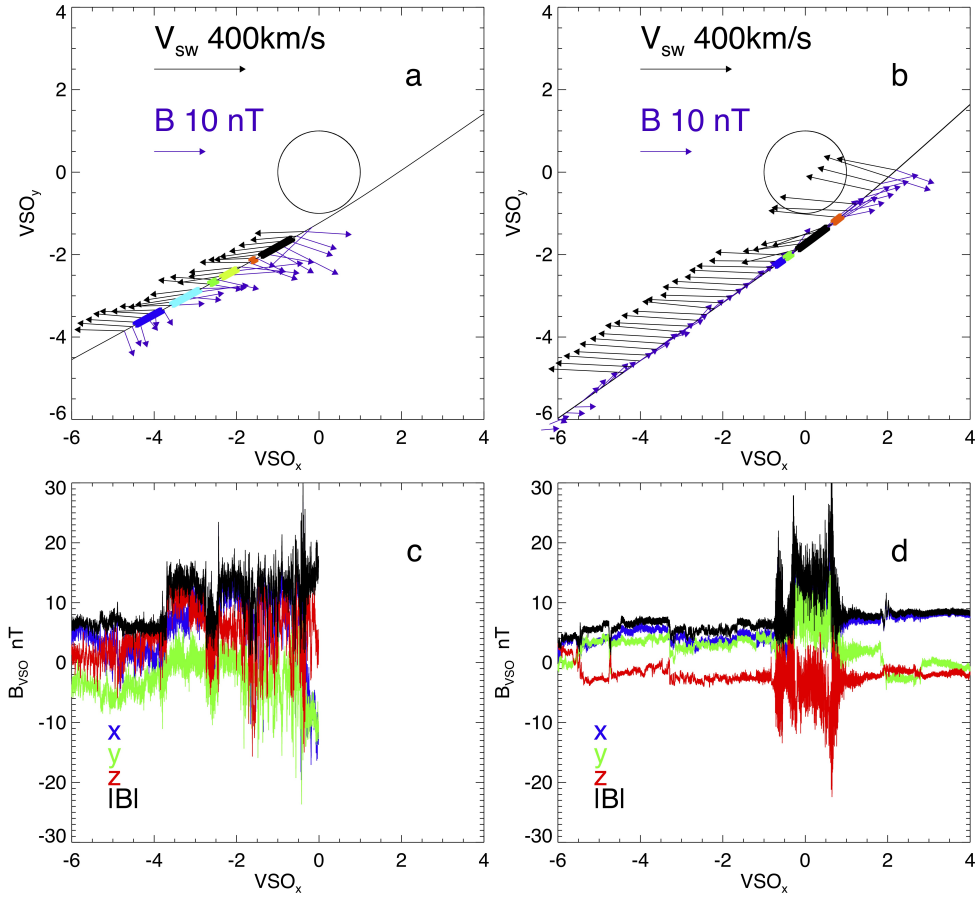
146 FIELDS measures electromagnetic fluctuations, creating a variety of data prod-  
 147 ucts (Bale et al., 2016; Malaspina et al., 2016; Pulupa et al., 2017; Bowen, Bale, et  
 148 al., 2020). The magnetic field is measured by a low frequency fluxgate magnetometer  
 149 (MAG) and an AC coupled search coil magnetometer (SCM). We use merged SCM and  
 150 MAG (SCaM) data, with DC-146 Hz bandwidth (Bowen, Bale, et al., 2020). Following  
 151 the first solar encounter, the SCM sensor  $x$  axis has exhibited significant anomalous  
 152 behavior. Thus, for VGA2 only two component magnetic field measurements (SCM  $y$   
 153 and  $z$ ) are available at kinetic scales.

154 PSP is specifically configured for measuring solar wind plasma in the inner helio-  
 155 sphere (Fox et al., 2016), which can complicate measurements of the Venusian plasma  
 156 environment. During VGA1, the solar limb-sensor (which maintains correct pointing  
 157 during solar encounters) responded to the Venusian albedo, turning off the instruments  
 158 midway magnetospheric transit, Figure 1(a-c). Additionally, SWEAP’s field of view  
 159 (FOV) is designed to measure the solar wind and its aberration in the spacecraft frame  
 160 (Kasper et al., 2016; Case et al., 2020; Whittlesey et al., 2020), leading to issues in  
 161 sampling the planetary plasma.

### 162 2.1 VGA1

163 During VGA1 SWEAP/Solar Probe ANalyzers (SPAN) ion measurements did  
 164 not capture the core proton distribution in its FOV, though electron measurements  
 165 from SPAN were made. During VGA2, PSP was configured with the spacecraft boom  
 166 in sunlight, in order to diagnose temperature dependence of the anomalous SCMx be-  
 167 havior, which unfortunately resulted in noisy SWEAP/Solar Probe Cup (SPC) mea-  
 168 surements. However, SPAN measured distributions of both magnetosheath electrons  
 169 and protons.

170 Figure 1 shows PSP’s trajectory in the VSO  $x - y$  plane during VGA1 (a) and  
 171 VGA2 (b). Magnetic field data are shown in Figure 1(c-d). Five bow-shock cross-



**Figure 1.** (a-b) Trajectory of PSP during VGA1 and VGA2 in VSO  $x$ - $y$  plane. Black arrows show scaled plasma flow; purple arrows show measured magnetic field. (c-d) Vector magnetometer measurements for VGA1 and VGA2 ( $x$ ,  $y$ ,  $z$ /blue, green, red) with the magnitude (black).

ings were recorded during VGA1. Figure 1(a) shows foreshock (FS) regions (blue, green, red), and magnetosheath (MS) regions (teal, yellow, black). Figure 2(a) shows vector magnetic time series for VGA1, with regions demarcated by dashed lines. Upstream quantities are  $B_0=5.9$  nT,  $T_p=5.9$  eV,  $T_e=10.9$  eV,  $n_p=11$  cm<sup>-3</sup>,  $n_e = 31$  cm<sup>-3</sup>,  $V_{sw}=410$  km/s.

We focus on the downstream magnetosheath from 8:34:30-08:38:30, with  $B_0 = 12.7$  nT,  $T_i=11$  eV,  $T_e=14.45$ , eV  $n_p=20$ , cm<sup>-3</sup>,  $n_e=55$  cm<sup>-3</sup>, and  $V_{MS}=380$  km/s. Magnetic coplanarity suggests quasi-parallel shock geometry, with a normal of 175° (Paschmann & Daly, 1998). Significant differences between  $n_e$  and  $n_p$  are observed both upstream and downstream; however the ratio  $n_e/n_i \sim 2.7$  stays constant across the shock. Additionally, a cross shock density ratio, 1.8, is observed for both electrons and protons, suggesting that while error exists in the absolute measurement of density, the relative scaling is physical. Estimates for upstream  $\beta_p$  range between 0.7-2.0; downstream  $\beta_p$  ranges from 0.6-1.5.

Figure 2(b-c) shows trace power-spectra for the FS and MS. Largely non-power-law spectra are observed indicating significant wave activity and instabilities (Burgess et al., 2005). The MS fluctuations show power-law spectra, commonly associated with turbulence. Vertical lines show spacecraft frame frequencies corresponding to  $k\rho_i \sim 1$  and  $k\rho_e \sim 1$ , assuming the Taylor hypothesis  $k = 2\pi f/V_{sw}$ . Magnetosheath kinetic spectra scale as  $k^{-2.9}$ , and no spectral break observed at ion-kinetic scales (e.g.  $k\rho_i \sim 1$ ). The extension of kinetic range spectra into inertial range frequencies has been interpreted as the result of FLR effects (Uritsky et al., 2011); parallel shock dynamics likely affect plasma kinetics in this region, leading to the lack of an observed inertial range (Xiao et al., 2018). Sahraoui et al. (2006) attribute the extension of kinetic range scaling into fluid-scales with the presence of mirror modes. Spectral properties of the three MS regions are similar, though strong electron scale wave activity observed behind the first shock crossing (teal) is seemingly absent from other MS intervals.

## 2.2 VGA2

During VGA2, two (inbound and outbound) shock crossings occurred, Figure 2(d,e) shows separate FS and MS regions. Upstream parameters are  $B_0= 7.8$  nT,  $T_e=16$  eV,  $n_p=21$  cm<sup>-3</sup>,  $n_e= 24$  cm<sup>-3</sup>  $V_{sw}=340$  km/s, due to poor measurements of upstream  $T_i$ , we cannot report upstream  $\beta_i$ .

SPAN resolved the ion distribution in the downstream magnetosheath, characterized by:  $B_0=14$  nT,  $T_p = 92$  eV,  $n_p = 15$  cm<sup>-3</sup>,  $n_e= 57$  cm<sup>-3</sup>. The significant difference between ion and electron densities is likely not physical: absolute ion-density is likely affected by FOV issues. There is decent agreement between  $n_e$  and  $n_p$  from SPC in the upstream solar wind; we set  $n_p = n_e = 57$  cm<sup>-3</sup>. The downstream MS flow is  $V_{MS} = 276$  km/s and  $V_a = B/\sqrt{2\mu_0\rho} = 40$  km/s, such that the Taylor hypothesis is applicable for Alfvén waves. Magnetic coplanarity of the VGA2 bow-shock gives a shock normal of 115 degrees, quasi-perpendicular to the upstream field.

Figure 2(e) shows FS spectra with non-power-law scaling and significant wave activity; the MS spectra, Figure 2f shows power-law scaling. Figure 2(f-g) shows  $k^{-3.4}$  spectrum for scales between  $k\rho_i = 1$  and  $kd_e = 1$ , with further steepening to an approximate  $k^{-6.3}$  spectrum near electron scales. The steepening occurs at a frequency between  $k\rho_i = 1$  and  $kd_e = 1$ , though there are uncertainties in the electron measurements. The observation of a secondary steepening at electron scales is consistent with observations in the terrestrial magnetosphere (Huang et al., 2014; Chen & Boldyrev, 2017).

The spectral index of the MS spectra,  $k^{-3.4}$ , is significantly steeper than in VGA1, or what is typically associated with kinetic Alfvén wave (KAW) turbulence



(Boldyrev & Perez, 2012b; Zhao et al., 2016). Simulations can recover similarly steep spectra, though typically at low  $\beta$  (Franci et al., 2015, 2016). At high  $\beta$ , increased damping may result in enhanced spectral steepening over the kinetic range (Howes et al., 2007, 2011). VGA2 shows an inertial-kinetic scale break around  $k\rho_i = 1$ , which is not evident behind the quasi-parallel shock. The inertial range is possibly less steep than  $k^{-5/3}$ , thought due to the short interval it is difficult to measure with great confidence

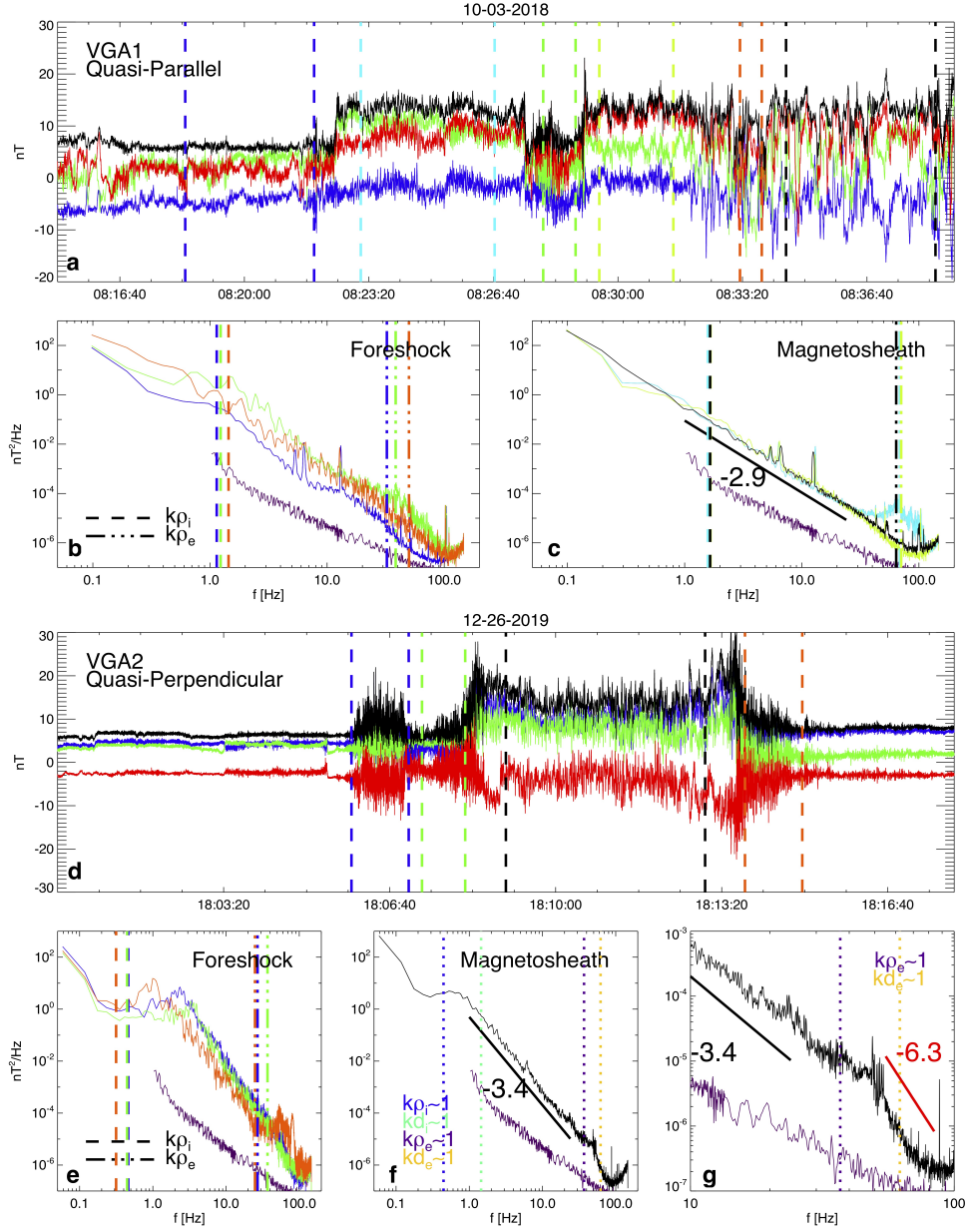
Kinetic Alfvén wave turbulence is commonly associated with a  $k^{-7/3}$  spectrum, with some variation from intermittency or damping (Howes et al., 2007; Boldyrev & Perez, 2012b; Howes et al., 2011). The kinetic spectrum measured with  $d_i < 1/k < d_e$  is significantly steeper than predictions of KAW turbulence (Schekochihin et al., 2009; Howes et al., 2011; Boldyrev & Perez, 2012b; Franci et al., 2015, 2016; Zhao et al., 2016; Grošelj et al., 2018). Notably Rezeau et al. (1999), previously measured  $k^{-3.4}$  scaling behind the terrestrial bow-shock. If the steep  $k^{-3.4}$  spectrum is a signature of significant heating, the measured  $T_i/T_e > 1$  may indicate preferential ion heating through turbulent dissipation via Landau damping, which is observed in simulations at high  $\beta$  (Kawazura et al., 2019).

### 2.3 Temperature Anisotropy

During VGA2, SPAN measured anisotropic temperatures, with  $T_{\perp}/T_{\parallel} \sim 2$ . At high  $\beta$ , significant  $T_{\perp}/T_{\parallel} > 1$  will drive mirror mode or Alfvén ion-cyclotron (AIC) instabilities (Gary, 1992; Hellinger et al., 2006; Bale et al., 2009). Figure 3 shows proton velocity distributions measured by SPAN-Ion during VGA2 in instrument coordinates. Instrumental FOV effects are highlighted by the cutoff in the  $y$  direction. Bi-Maxwellian fits allow for computation of temperature anisotropies  $T_{\perp}/T_{\parallel}$ .

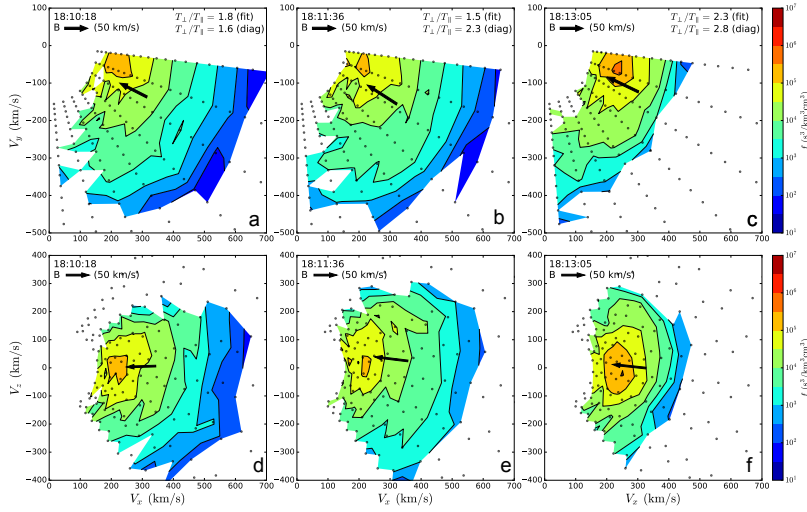
An alternative method of estimating the temperature anisotropy through diagonalizing the measured temperature moment tensor from SPAN verifies this measurement. The temperature tensor is rotated into a frame aligned with the magnetic field; assuming gyrotropy, there are enough degrees of freedom to calculate  $T_{\perp}$  and  $T_{\parallel}$  from well-measured tensor components ( $T_{xx}, T_{zz}, T_{xz}$ ) without using the poorly-measured  $y$ -component. The independent methods of calculating temperature anisotropy provide similar results, and thus confidence in the measurement.

While turbulent heating significantly affects spectral indices, it's likely that the  $T_{\perp}/T_{\parallel} \sim 2$  anisotropy plays a role in the kinetic cascade. Dwivedi et al. (2015) suggest that the kinetic scale spectra at Venus may relate to the growth of these instabilities in the magnetosheath. The growth of the AIC instability is associated with circularly polarized electromagnetic waves at ion scales (Verscharen et al., 2019). Analysis of polarization signatures reveals little significant circular polarization suggesting that a mirror instability may dominate; however, the angle between the mean field and the solar wind flow is  $118^\circ$ , such that quasi-parallel waves may be hard to identify (Bowen, Mallet, Huang, et al., 2020). Volwerk et al. (2008) previously reported mirror modes behind a quasi-perpendicular bow shock at Venus. At  $T_{\perp}/T_{\parallel} \sim 2$  and  $\beta \sim 10$ , growth rates for mirror mode may be large, e. g. as  $0.1 \omega_c$  (Hellinger et al., 2006). For  $f_{ci} \sim 1$  Hz, this corresponds to a growth rate of  $\sim 10$  s. The presence of  $\alpha$  particles and other heavy ions in the magnetosphere can affect instability growth rates (Chen et al., 2016; Verscharen et al., 2019); it has been suggested that heavy ions stabilize the AIC instability (Price et al., 1986). The steep kinetic range spectrum may result from the introduction of KAW with nonlinear interactions with driven non-propagating mirror mode structures. The mirror mode is commonly associated with anti-correlated magnetic and kinetic pressure; however, the SPAN measurement cadence is not sufficient to determine correlations at kinetic scales.



**Figure 2.** (a) Vector magnetic field measurements from VGA1. Color coded lines demarcate three foreshock regions (blue, green, orange) from three magnetosheath regions (teal, yellow, black). (b,c) Color coded power-spectra for intervals shown in (a); dashed/dotted lines correspond to convected ion/electron gyroradius  $k\rho_{i/e}$ . Purple curve shows SCM sensitivity. (d) Vector magnetic field measurements from VGA2. (e) Power spectra from foreshock regions (blue, green, orange) and  $k\rho_{i/e}$ . (f) Magnetosheath spectra with convected ion/electron gyroradius  $k\rho_{i/e}$  and inertial length  $kd_{i/e}$ . (g) Magnetosheath spectra from 10-100 Hz, showing electron scale steepening.





**Figure 3.** (a-c) Proton distributions from VGA2 magnetosheath observed by SPAN at three times in sensor  $x - y$  plane. (d-f) Proton distributions from SPAN for VGA2 magnetosheath in sensor  $x - z$  plane. The magnetic field, in Alfvén units, is shown as a black arrow.

272 Sahraoui et al. (2006) discuss the mirror instability in the terrestrial magne-  
 273 tosheath, demonstrating non-propagating structures characteristic of the mirror mode;  
 274 however they measure an energy spectrum similar to the canonical KAW  $k^{-2.7}$  scaling,  
 275 which extends into scales typically associated with the inertial range. The presence of  
 276 these modes, and other instabilities, likely effects observed signatures of kinetic scale  
 277 turbulence.

### 278 3 Signatures of a Kinetic Cascade

279 Systematically shallow spectra at inertial scales suggest that inertial range magne-  
 280 tosheath turbulence may not always form (Czaykowska et al., 2001; Alexandrova et  
 281 al., 2008). Whether instabilities can drive kinetic scale turbulence in the absence of  
 282 an inertial range cascade is an open question (Hadid et al., 2015). The higher order  
 283 moments of distributions of turbulent fluctuations provide information regarding the  
 284 development and dissipation of turbulence (Matthaeus et al., 2015; Tessein et al., 2013;  
 285 Mallet et al., 2019; Bandyopadhyay et al., 2020).

286 Distributions of turbulent fluctuations are often characterized with statistical mo-  
 287 ments of increments, (Monin & Yaglom, 1971, 1975; Dudok de Wit & Krasnoselkikh,  
 288 1996; Sorriso-Valvo et al., 1999; Hnat et al., 2002; Kiyani et al., 2006). However, incre-  
 289 ments cannot resolve spectral scaling steeper than  $k^{-3}$  (Frisch, 1995; Cho & Lazarian,  
 290 2009). For scalings observed in the Venus magnetosheath, alternative measurements of  
 291 fluctuation amplitudes, such as the continuous wavelet transform (CWT), are required  
 292 to capture higher order properties of kinetic range turbulence (Farge, 1992; Farge &  
 293 Schneider, 2015; Kiyani et al., 2015).

$$\tilde{B}(s, \tau) = \sum_{i=0}^{N-1} \psi\left(\frac{t_i - \tau}{s}\right) B_j(t_i); \quad (1)$$

we use the Morlet wavelet

$$\psi(\xi) = \pi^{-1/4} e^{-i\omega_0\xi} e^{-\frac{\xi^2}{2}},$$

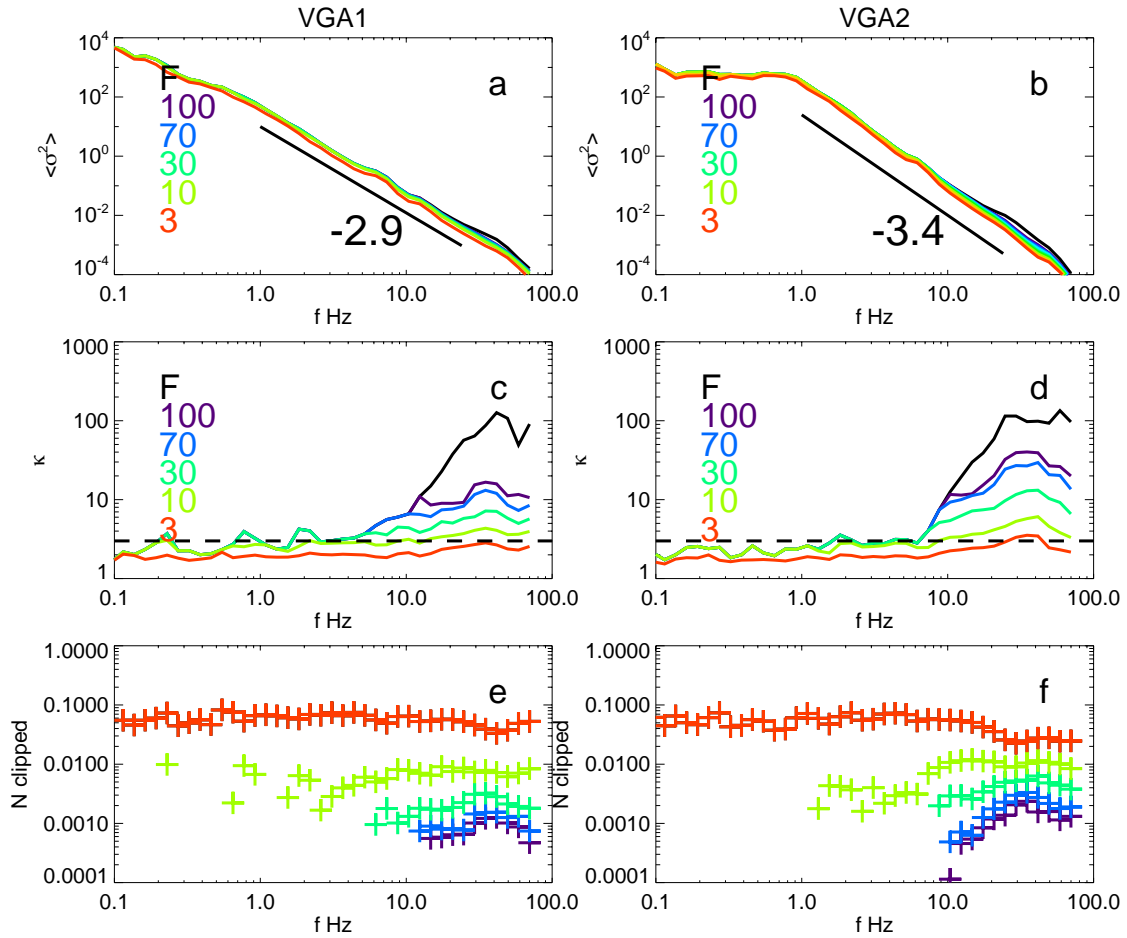
with  $\omega_0 = 6$ .

Figure 4(a-b) shows  $\langle \sigma_s^2 \rangle$  for VGA1 and VGA2. Figure 4(c-d) show the scale dependent kurtosis  $\kappa = \langle |\tilde{B}^4| \rangle / \langle \sigma_s^2 \rangle^2$  = computed for each wavelet scale. Increasing  $\kappa$  is seen in both VGA1 and VGA2 at  $f \gtrsim 10$  Hz.

Excluding outlier fluctuations at a given scale, conditioning, decreases effects of transients, e.g. those observed in VGA1 and VGA2 by Page (2020) and Goodrich (2020), on  $\kappa$  (Kiyani et al., 2006). For each scale, wavelet coefficients with  $\sigma^2 > F \langle \sigma^2 \rangle$  are removed for  $F = 3, 10, 30, 70, 100$ , and  $\langle \sigma^2 \rangle$  and  $\kappa$  are recomputed. Large decreases in  $\kappa$  are observed when removing outliers, while the power is not greatly affected. The conditioning has similar effects for both VGA1 and VGA2, indicating that though the spectral scalings differ, the scaling of kurtosis is similar. In both cases  $F=10$ , removes approximately 1% of fluctuations in sub-ion scales, though the kurtosis remains larger than 3 (expected for Gaussian fluctuations). This indicates the presence of non-Gaussian fluctuations commonly associated with kinetic range turbulence (Kiyani et al., 2009; Hadid et al., 2015; Kiyani et al., 2015). Higher order moments can be difficult to compute accurately for finite sample lengths (Kiyani et al., 2006). Dudok de Wit (2004) suggest requiring explicit convergence of higher order moments, though they derive an approximate required number of samples given by  $\log_{10} N - 1$ . For these 4 minute, ( $N \sim 70000$ ) records,  $\log_{10}(N) - 1 = 3.85$ , suggesting that kurtosis may not be perfectly resolved. While our measurement of kurtosis may lack accuracy, non-Gaussianity of kinetic scale fluctuations is evident in the distributions of wavelet coefficients (not shown). Hadid et al. (2015) show different scaling properties of higher order moments of turbulent amplitudes behind quasi-perpendicular and quasi-parallel shocks at Saturn, implying differences in the kinetic scale intermittency, but do not perform any conditioning.

## 4 Summary

We present measurements of kinetic scale turbulence in the Venusian magnetosphere behind both a quasi-perpendicular and quasi-parallel bow shock. A steep kinetic range spectrum is observed behind the quasi-perpendicular (VGA2) shock with a sub-ion  $k^{-3.4}$  scaling. Observation of significant temperature anisotropy ( $T_{\perp}/T_{\parallel} \sim 2$ ) in  $\beta \sim 10$  plasma suggests that the mirror or Alfvén ion cyclotron instabilities are quite strong; the lack of observed circular polarization suggests a dominant mirror instability (Gary, 1992; Hellinger et al., 2006). The nonlinear generation of mirror modes (Southwood & Kivelson, 1993) may increase nonlinear interaction rates at kinetic scales, steepening the cascade from typically observed  $k^{-8/3}$  spectra (Huang et al., 2014; von Papen et al., 2014; Hadid et al., 2015; Chen & Boldyrev, 2017). The steep spectra may also be associated with preferential ion heating at high  $\beta$  (Kawazura et al., 2019). At  $kd_e = 1$  a secondary kinetic steepening is observed consistent with the observations of the terrestrial magnetosphere (Huang et al., 2014; Chen & Boldyrev, 2017). Behind the quasi-parallel shock a  $k^{-2.9}$  scaling occurs; no measurements of temperature anisotropy were available. Though spectral energy scaling varies between Venus encounters, the kurtosis in either case shows similar signatures of non-Gaussianity, indicating kinetic range developed turbulence. Our results highlight the importance of ion-scale instabilities in shaping kinetic turbulence in planetary environments.



**Figure 4.** (a,b) CWT spectra  $\langle \sigma^2 \rangle$  for VGA1 and VGA2 (black); colors correspond to conditioned spectra. (c,d) Effect of conditioning on wavelet kurtosis for VGA1 and VGA2. (e-f) Percentage of clipped wavelet coefficients at each conditioning level.

## 5 Acknowledgements

Parker Solar Probe FIELDS and SWEAP instrumentation were developed under contract NNN06AA01C. PSP data is publicly available at NASA Space Physics Data Facility (SPDF) <https://cdaweb.gsfc.nasa.gov/>. FIELDS data is also hosted at [sprg.ssl.berkeley.edu/data/psp/data/sci/fields/](http://sprg.ssl.berkeley.edu/data/psp/data/sci/fields/). Discussion of the merged SCM and MAG (SCaM) data may be found in (Bowen, Bale, et al., 2020).

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