

Supporting Information

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Optimizing the Proton Conductivity with the Isokinetic  
Temperature in Perovskite-Type Proton Conductors  
According to Meyer–Neldel Rule

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## Proton conductivity

Typical impedance spectra for proton conducting perovskites show two semicircles: at high-frequency for grain interior (referred as bulk), and at low-frequency for from grain boundaries.

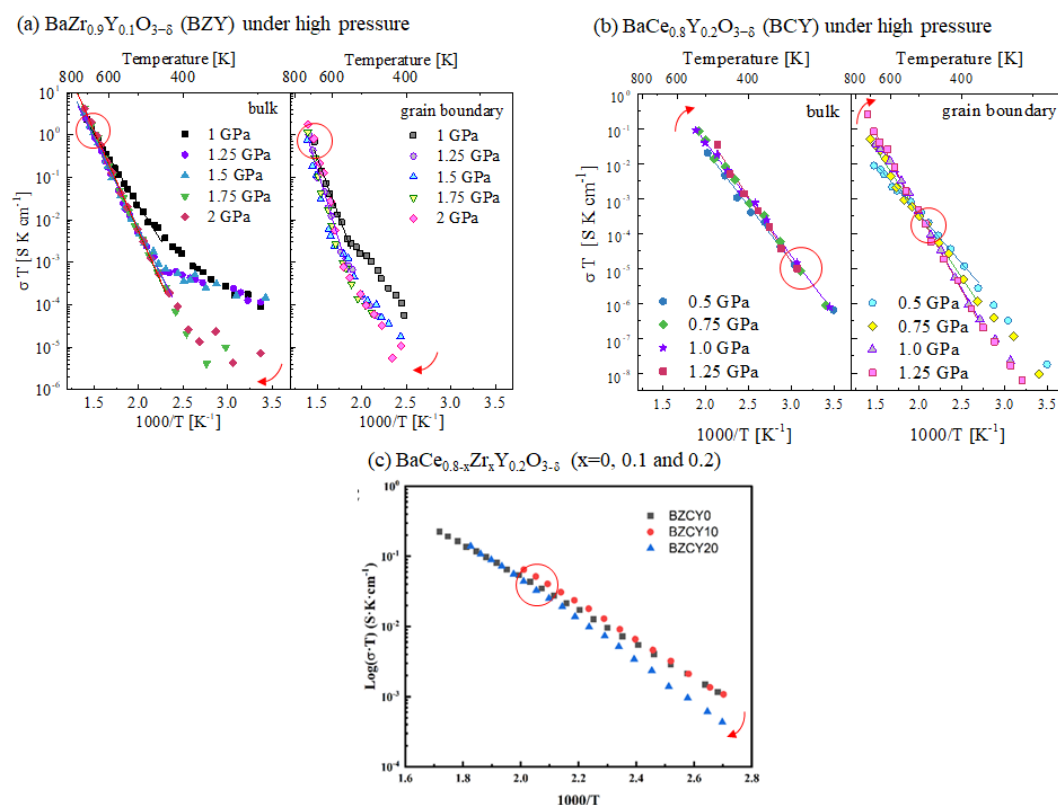


Figure S1. The proton conductivity of BZY <sup>[1, 2]</sup> a), BCY <sup>[3]</sup> b) under high pressure, and BaZr<sub>0.8-x</sub>Ce<sub>x</sub>Y<sub>0.2</sub>O<sub>3- $\delta$</sub>  (x=0, 0.1 and 0.2, denoted as BZCY0 to 20) with varying Ce content <sup>[4]</sup> c). Solid lines show the range that the activation energies and prefactors were extracted. Red circles indicate the point that the proton conductivities intersect.

## Raman peak analysis

### 1. Baseline subtraction

An iterative polynomial algorithm was used to fit the baselines of each Raman spectra in this work and the baseline was subtracted, thereof. A fifth-order of polynomial was used in the fitting procedures.<sup>[5]</sup> Spectral region from  $100\text{ cm}^{-1}$  to  $1000\text{ cm}^{-1}$  of the spectra was processed.

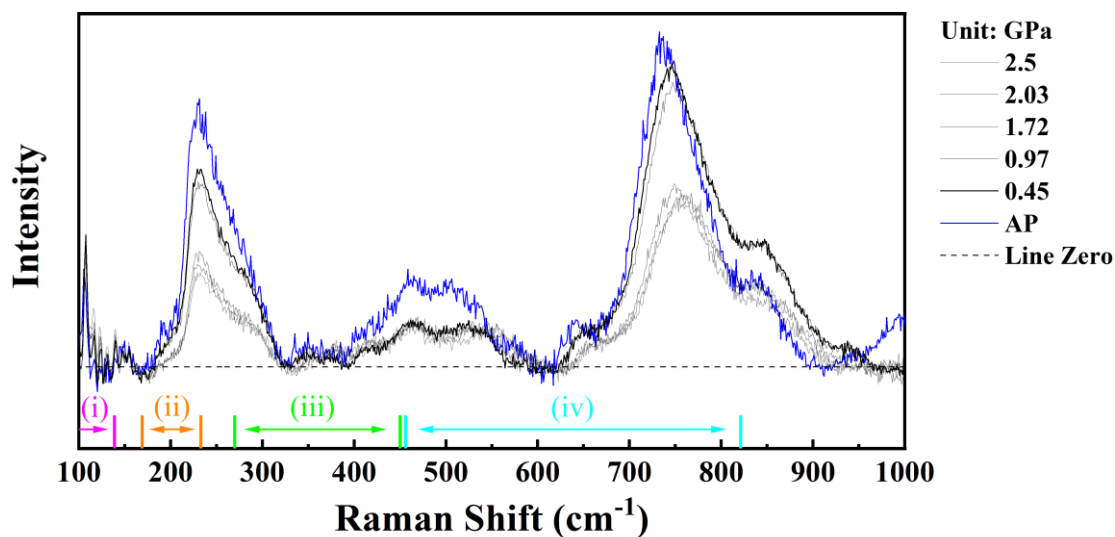


Figure S2. Pressure-dependent Raman spectra (after baseline subtraction) for BZY. (i) Ba-ZrO<sub>6</sub> stretching, (ii) ZrO<sub>6</sub> torsion, (iii) O-Zr-O bending, (iv) Zr-O stretching.

## 2. Multi-peak fitting

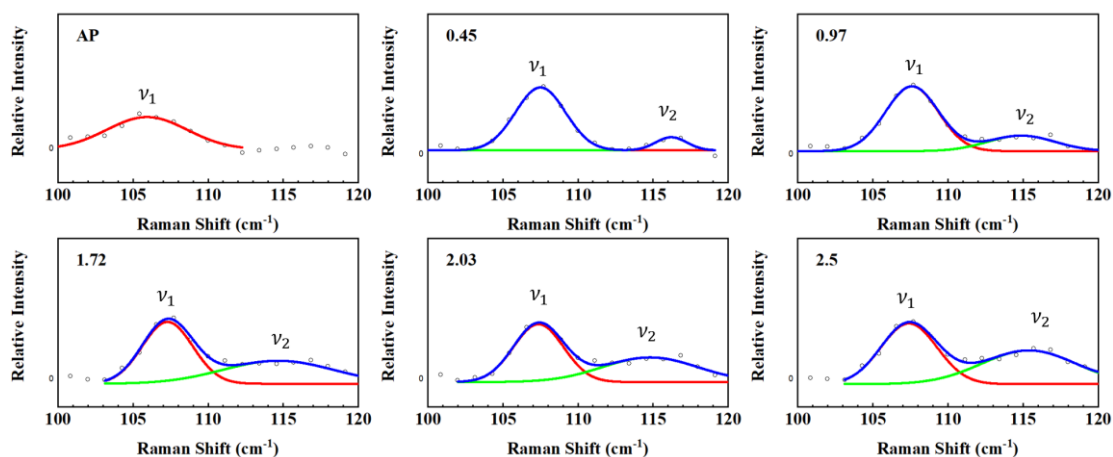


Figure S3. Raman peak fitting diagram from ambient pressure (AP) to 2.5 GPa for the Ba-ZrO<sub>6</sub> stretching motion of BZY from 100 to 150 cm<sup>-1</sup>. Black open circles: experimental data; blue line: cumulative fit for the Raman peak; red line:  $\nu_1$ ; green line:  $\nu_2$ . Two Gaussian functions were used in the peak fitting procedure.

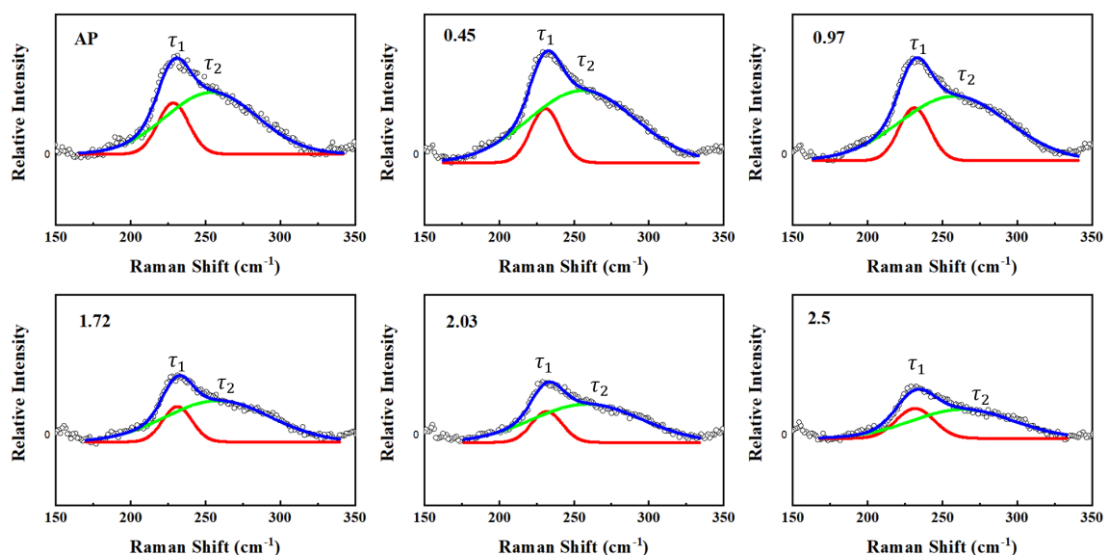


Figure S4. Raman peak fitting diagram from AP to 2.5 GPa for the ZrO<sub>6</sub> torsional motion of BZY from 200 to 350 cm<sup>-1</sup>. Black open circles: experimental data; blue line: cumulative fit; red line:  $\tau_1$ ; green line:  $\tau_2$ . Two Gaussian functions were used in the peak fitting procedure.

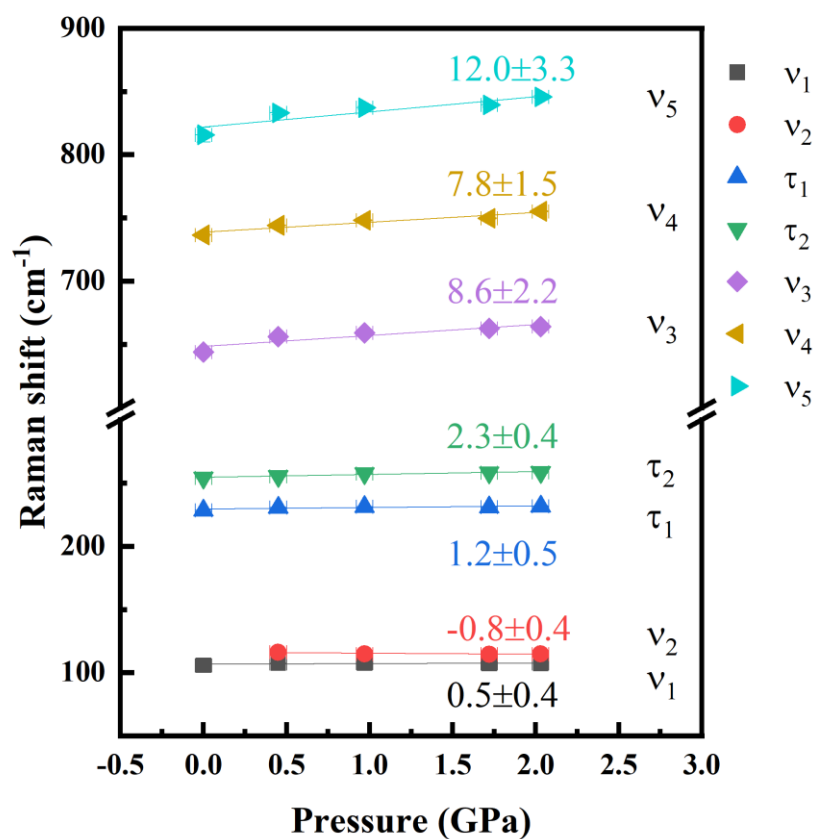


Figure S5. Raman shift of the vibrational modes in BZY as a function of pressure. Error bars are within the size of symbols.

**Evolution of lattice parameters under pressure**

According to the XRD structural refinement in Figure S6a, we obtained the pressure-dependent lattice volume and the Zr-O bond length, as shown in Figure S6b. The Birch–Murnaghan equation of state was used to fit the pressure-volume (P-V) diagram, with the first-pressure derivative of the bulk modulus  $K'_0$  taken as 4<sup>[6]</sup>:

$$P = \frac{3K_0}{2} \left[ \left( \frac{V}{V_0} \right)^{-\frac{7}{3}} - \left( \frac{V}{V_0} \right)^{-\frac{5}{3}} \right] \quad (\text{S1.})$$

where  $K_0$  is the isothermal bulk modulus and,  $V_0$  is the unit cell volume at ambient pressure.

From the fit to the P–V curve, we obtain the bulk modulus  $K_0 = 151.9 \pm 24$  GPa, and  $V_0 =$

$74.6 \pm 0.12$  Å<sup>3</sup>. The relationship between Zr-O bond length and pressure was then determined.

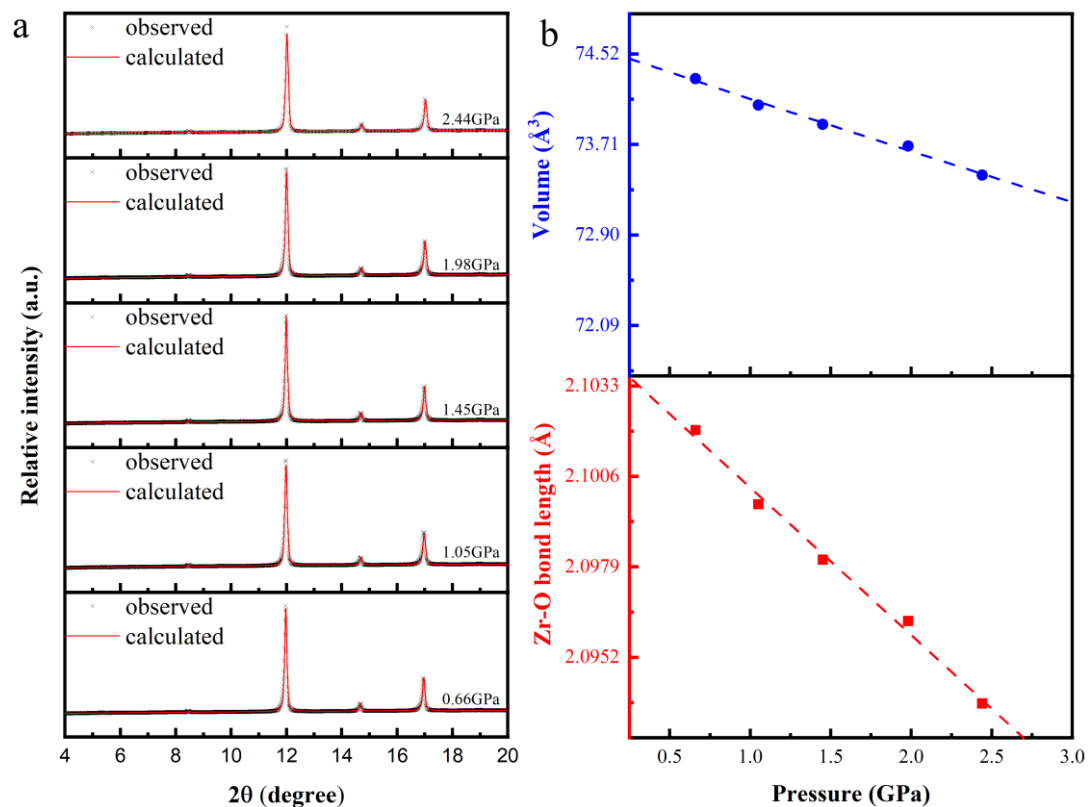


Figure S6. Pressure-dependent X-ray diffractogram for BZY: a) Structural refinement of XRD pattern from 0.66 to 2.44 GPa (with  $R_f$ -factor  $< 4\%$ , Bragg  $R$ -factor  $< 5\%$ ); b) Upper: pressure–volume diagram of BZY, Lower: pressure-dependent Zr-O bond length. Dashed lines represent the fit according to the Birch–Murnaghan equation of state with  $K'_0 = 4$ .



### Contribution of specific Raman modes

According to Fig. 3,  $E_a$  can be considered as a function of all the vibration modes, i.e.

$E_a(\nu_1, \nu_2, \tau_1, \tau_2, \nu_3, \nu_4, \nu_5)$ . Take partial differential of  $E_a$  over all the frequencies,

$$dE_a = \left(\frac{\partial E_a}{\partial \nu_1}\right) d\nu_1 + \left(\frac{\partial E_a}{\partial \nu_2}\right) d\nu_2 + \left(\frac{\partial E_a}{\partial \tau_1}\right) d\tau_1 + \left(\frac{\partial E_a}{\partial \tau_2}\right) d\tau_2 + \left(\frac{\partial E_a}{\partial \nu_3}\right) d\nu_3 + \left(\frac{\partial E_a}{\partial \nu_4}\right) d\nu_4 + \left(\frac{\partial E_a}{\partial \nu_5}\right) d\nu_5$$

If we consider all the vibration modes were equally weighted, that is, all coefficients were approximately equal, e.g.

$$\left(\frac{\partial E_a}{\partial \nu_1}\right) \approx \left(\frac{\partial E_a}{\partial \nu_2}\right) \approx \left(\frac{\partial E_a}{\partial \tau_1}\right) \approx \left(\frac{\partial E_a}{\partial \tau_2}\right) \approx \left(\frac{\partial E_a}{\partial \nu_3}\right) \approx \left(\frac{\partial E_a}{\partial \nu_4}\right) \approx \left(\frac{\partial E_a}{\partial \nu_5}\right)$$

From Fig. 3, we can see that the changes in  $\nu_3, \nu_4, \nu_5$  ( $d\nu_3, d\nu_4$  and  $d\nu_5$ ) are much larger than the changes in the other vibration modes, suggesting that these modes contribute the most to changes in  $E_a$ .

### Raman frequency for $\text{BaZr}_{0.8-x}\text{Ce}_x\text{Y}_{0.2}\text{O}_{3-\delta}$ ( $x=0, 0.1$ and $0.2$ )

$\text{BaZr}_{0.8}\text{Y}_{0.2}\text{O}_{3-\delta}$ ,  $\text{BaZr}_{0.7}\text{Ce}_{0.1}\text{Y}_{0.2}\text{O}_{3-\delta}$  and  $\text{BaZr}_{0.6}\text{Ce}_{0.1}\text{Y}_{0.2}\text{O}_{3-\delta}$  were prepared using the same method as  $\text{BaZr}_{0.9}\text{Y}_{0.1}\text{O}_{3-\delta}$  described in this work. Raman Scattering spectra were collected with a Senterra R200-L Dispersive Confocal Raman Spectrometer using 10mW 532nm laser, a 50 $\times$  magnification objective, and with the accumulation of 1s.

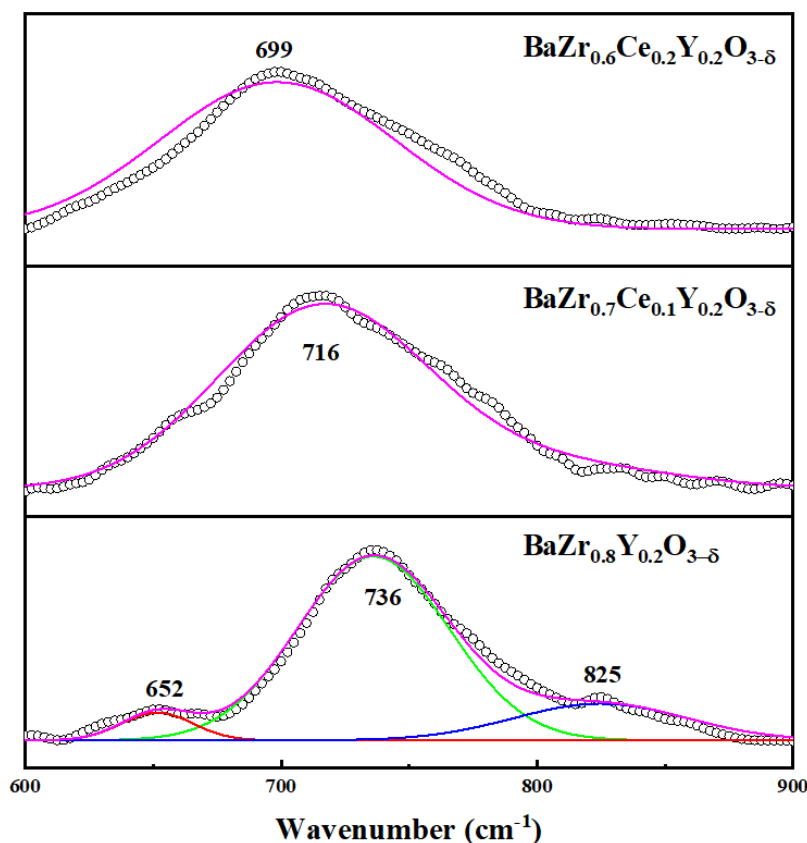


Figure S7. Raman peak fitting diagram from 600 to 900  $\text{cm}^{-1}$  for  $\text{BaZr}_{0.8}\text{Y}_{0.2}\text{O}_{3-\delta}$ ,  $\text{BaZr}_{0.7}\text{Ce}_{0.1}\text{Y}_{0.2}\text{O}_{3-\delta}$  and  $\text{BaZr}_{0.6}\text{Ce}_{0.1}\text{Y}_{0.2}\text{O}_{3-\delta}$ . Black open circles: experimental data; purple line: cumulative fit; red line:  $\nu_3$ ; green line:  $\nu_4$ ; blue line:  $\nu_5$ . Three Gaussian function were used in the peak fitting procedure for  $\text{BaZr}_{0.8}\text{Y}_{0.2}\text{O}_{3-\delta}$ ; for  $\text{BaZr}_{0.7}\text{Ce}_{0.1}\text{Y}_{0.2}\text{O}_{3-\delta}$  and  $\text{BaZr}_{0.6}\text{Ce}_{0.1}\text{Y}_{0.2}\text{O}_{3-\delta}$ , the three bands merge into one, therefore one Gaussian function was used in fitting.

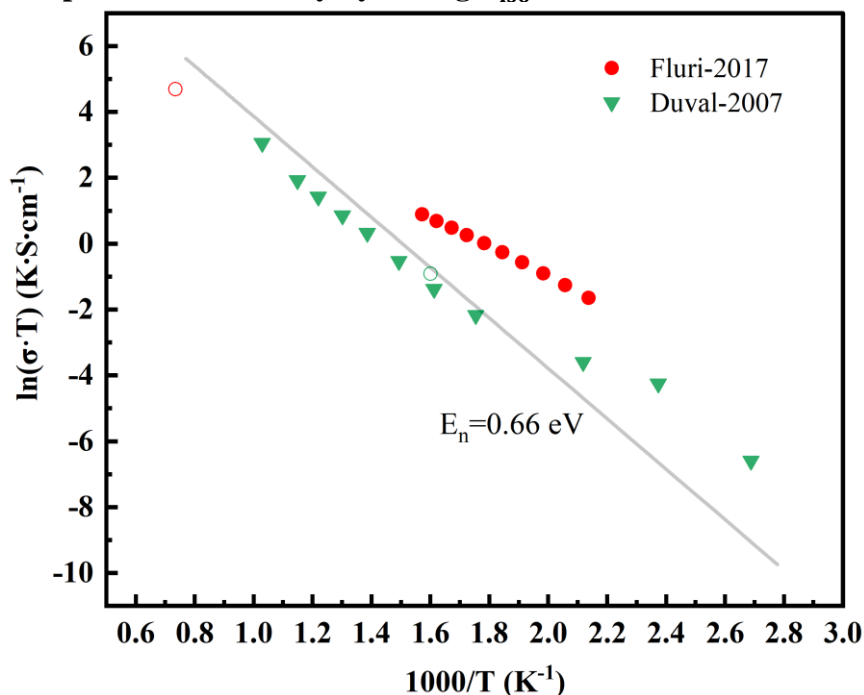
Enhancement of proton conductivity by tuning  $T_{\text{iso}}$ 

Figure S8. The bulk proton conductivity of  $\text{BaZr}_{0.9}\text{Y}_{0.1}\text{O}_{3-\delta}$  (BZY10) reported in ref. [7] for sintered pellet and in ref. [8] with epitaxial strain, respectively. The green open circle shows the isokinetic-temperature of BZY10 reported in this work. The red open circle is the isokinetic temperature for BZY10 epitaxial film. The solid line indicates the critical energy determined by Equation (5).

## References

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