

LABORATORY OBSERVATIONS OF ELECTRON TEMPERATURE IN THE WAKE OF A SPHERE IN A STREAMING PLASMA

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Abstract—A parametric study was performed of electron temperature variation in the wake of a conducting sphere in a streaming plasma. The flow conditions were varied as follows: the ambient electron temperatures in the range 850–2450 K; the ambient electron densities in the range 5×10^4 – $7 \times 10^6/\text{cm}^3$; and body potentials relative to plasma potential in the range of +1.7 to –2.8 V for an ion beam energy of ~ 4 eV. Electron temperature enhancements were observed which ranged up to 200 per cent above ambient in the nearest proximity of the body surface. The magnitude of the enhancement depends upon the ambient density, temperature and body potential.

INTRODUCTION

The interaction between a spacecraft and its environmental space-plasma, in particular the charged particle distribution around ionospheric satellites, has been the topic of numerous theoretical investigations in the past decade (for a recent review, see Gurevich *et al.*, 1969). The available *in situ* results, on the other hand, contain a fragmentary and limited amount of information on the electron distributions at specific locations in the near wake zone. Even less information is available on the ion distributions (for a recent review, see Samir, 1973). Angular distributions of electron temperature in the closest vicinity of some spacecraft have also been presented by Berthelie and Sturges (1967), Medved (1969) and Samir and Wrenn (1972).

Recently, variations of electron temperature at the surface of a body-probe were studied by Illiano and Storey (1974) in the laboratory. Preliminary transverse profiles of electron temperature in the near wake region behind a sphere in laboratory streaming plasmas were reported by Samir *et al.* (1974). Here we present a parametric study of the dependence of the electron temperature in the wake upon ambient electron temperature, density, and body potential.

DESCRIPTION OF THE EXPERIMENT

Transverse and axial profiles of electron temperature, T_e , behind a stainless steel spherical target

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body of radius $R_0 = 4$ cm were measured in an Ar^+ plasma beam at a background pressure of 4×10^{-6} Torr. The plasma is produced with a modified Kauffman ion thruster whose small orifice effectively yields a “cold” ion beam. (The transverse ion temperature is estimated at ~ 300 K.) Additional details of the experimental set up can be found in Oran *et al.* (1974).

The measurements were made with a stainless steel, cylindrical Langmuir probe, 1.5 cm long and 0.1 cm dia., oriented perpendicular to the flow direction. The probe was bombarded with a high energy ion beam before each run to remove surface contamination. In addition, prior to and at the conclusion of each run, the i - v characteristics of the probe were checked to insure that there was no “hysteresis” present.

It should be noted that the term “electron temperature” used herein refers to the numerical values obtained in the conventional manner by a straightline fit to the current-voltage characteristics. The decrease of I_e in the wake region was compensated to some degree (in recording the i - v curves) by increasing an amplifier gain. Hence, a straightline-fit was obtained over typically 1.5 decades both for the ambient and the wake measurements.

The authors are aware of the possibility that the electron distribution in the wake may not be truly represented by a Maxwellian. Hence, the use of the term “temperature” in the conventional manner may not be justified. However, no such deviations were detected in the slopes of the current-voltage characteristics used in the present study which may indicate no significant departures from a Maxwellian. A more detailed discussion on the above is given in Samir and Wrenn (1972) and Illiano and Storey (1974).

RESULTS AND DISCUSSION

We present profiles of the ratio $R(T_e) = T_{ew}/T_{eo}$, where T_{ew} = electron temperature in the wake and T_{eo} = ambient electron temperature.

Ambient conditions were determined by mapping the flow field with the target body removed from the flow. All distances are normalized to the sphere radius, and the values of the flow conditions are shown in the figures. Here N_o is the ambient plasma density, E_i the ion beam energy, ϕ_s the body potential relative to plasma potential, and $\lambda_D = 6.9(T_{eo}/N_o)^{1/2}$ cm is the Debye length. The experimental error for each value of $R(T_e)$ is $\lesssim 20$ per cent.

Figure 1 presents examples of axial profiles of $R(T_e)$ at different locations downstream from the centre of the sphere (r/R_o) for body potentials ϕ_s in the range $+0.8$ – -1.2 V. It can be seen that an electron temperature enhancement significantly above the experimental error exists in the very near wake.

Examining the general influence of ϕ_s on $R(T_e)$, we see that at a distance of $1.25 R_o$ from the centre of the sphere (on the wake axis) the enhancements can be as great as 3 times T_{eo} for the more negative body potentials (see Fig. 1). For these potentials, however, the $R(T_e)$ enhancement seems to be confined to distances (r/R_o) $\lesssim 2$. For more positive values of ϕ_s , the enhancement appears to extend further downstream, although the maximum value of the enhancement is much smaller.

Figure 2 presents examples of transverse profiles

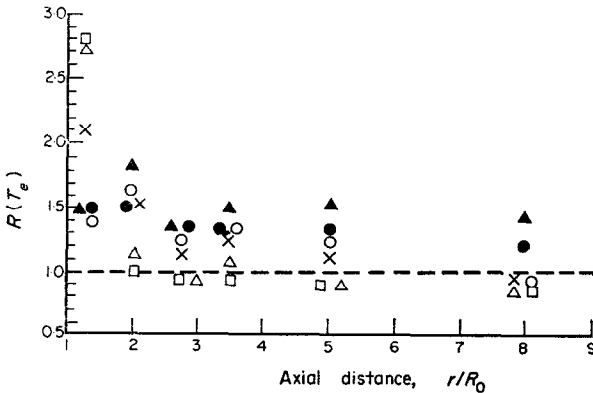


FIG. 1. AXIAL PROFILES OF $R(T_e) = T_{ew}/T_{eo}$ AT DIFFERENT DOWNSTREAM DISTANCES (r/R_o) BEHIND THE SPHERE AND FOR DIFFERENT BODY POTENTIALS, ϕ_s : $\blacktriangle \sim +0.8$ V, $\bullet \sim +0.3$ V, $\circ \sim 0$ V, $\times \sim -0.3$ V, $\triangle \sim -0.7$ V, AND $\square \sim -1.2$ V.

The ambient plasma conditions are: electron temperature (T_{eo}) ~ 1050 K, density (N_o) $\sim 6 \times 10^4/\text{cm}^3$, ion beam energy (E_i) ~ 3.2 eV and Debye length (λ_D) ~ 0.9 cm.

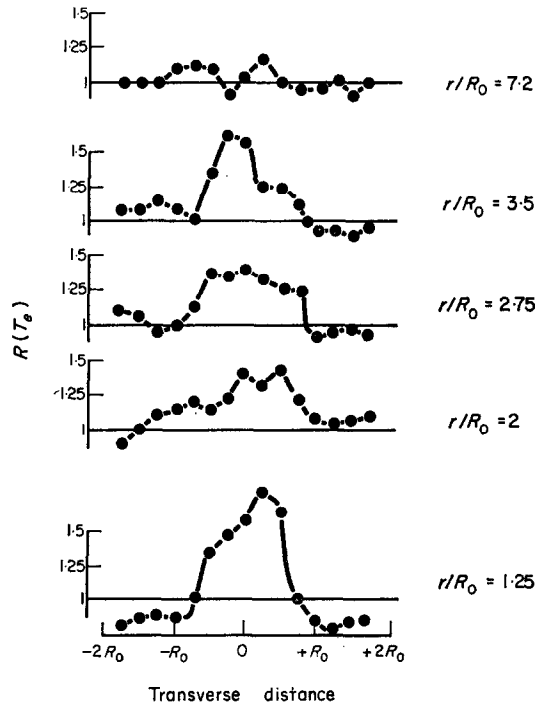


FIG. 2. TRANSVERSE PROFILES OF $R(T_e) = T_{ew}/T_{eo}$ AT DIFFERENT DOWNSTREAM DISTANCES (r/R_o) BEHIND THE SPHERE WHICH WAS AT FLOATING POTENTIAL ($\phi_s \sim -0.3$ V).

The ambient plasma conditions are: electron temperature (T_{eo}) ~ 1200 K, density (N_o) $\sim 7.5 \times 10^4/\text{cm}^3$, ion beam energy (E_i) ~ 5.3 eV and Debye length (λ_D) ~ 0.9 cm. After Samir *et al.* (1974).

of $R(T_e)$ at different downstream locations. The plasma flow conditions are similar to those in Fig. 1, although in this instance the body was at floating potential (~ -0.3 V). Significant electron temperature enhancements are also observed in the near wake. The transverse extent of the enhancement appears to be on the order of the sphere diameter.

Figure 3 presents axial profiles of $R(T_e)$. In this case, the ambient density N_o is larger by about an order of magnitude compared with the density given in Fig. 1. The rest of the flow conditions are similar. Here no enhancement in $R(T_e)$ can be observed to within the experimental error. Hence the dependence of T_{ew} on N_o is clearly demonstrated. Figure 3 also indicates that at the ambient density $7 \times 10^5/\text{cm}^3$, the T_e enhancement is relatively insensitive to the variations of ϕ_s over the range of -1.2 – $+0.8$ V. This is in contrast to the strong ϕ_s dependence of $R(T_e)$ for $N_o = 6 \times 10^4/\text{cm}^3$ (see Fig. 1), in the near wake.

Figure 4 presents axial profiles of $R(T_e)$ where

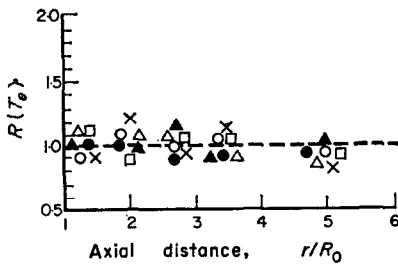


FIG. 3. AXIAL PROFILES OF $R(T_e) = T_{ew}/T_{e0}$ AT DIFFERENT DOWNSTREAM DISTANCES (r/R_0) BEHIND THE SPHERE FOR DIFFERENT BODY POTENTIALS, ϕ_s : $\blacktriangle \sim +0.8$ V, $\bullet \sim +0.3$ V, $\circ \sim 0$ V, $\times \sim -0.3$ V, $\triangle \sim -0.7$ V AND $\square \sim -1.2$ V.

The ambient plasma conditions are: electron temperature (T_{e0}) ~ 850 K, density (N_0) $\sim 7 \times 10^9/\text{cm}^3$, ion beam energy (E_i) ~ 3.6 eV and Debye length (λ_D) ~ 0.25 cm.

T_{e0} has been increased by more than a factor of 2 compared with that of Fig. 1, while the ambient density is basically the same. It is clearly seen that the effect of increasing the ambient temperature by a factor of ~ 2 is equivalent to increasing the ambient density by a factor of ~ 10 .

In addition, it is interesting to note that in the case of Fig. 3, the value of λ_D is decreased compared to that of Fig. 1, while the case of Fig. 4 represents

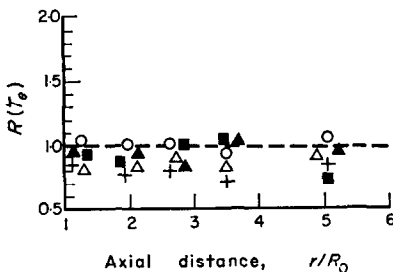


FIG. 4. AXIAL PROFILES OF $R(T_e) = T_{ew}/T_{e0}$ AT DIFFERENT DOWNSTREAM DISTANCES (r/R_0) BEHIND THE SPHERE AND FOR DIFFERENT BODY POTENTIALS, ϕ_s : $\blacksquare \sim +1.7$ V, $\blacktriangle \sim +0.8$ V, $\circ \sim 0$ V, $\triangle \sim -0.7$ V AND $\times \sim -2.8$ V.

The ambient plasma conditions are: electron temperature (T_{e0}) ~ 2450 K, density (N_0) $\sim 5 \times 10^8/\text{cm}^3$, ion beam energy (E_i) ~ 4.7 eV and Debye length (λ_D) ~ 1.55 cm.

an increase of λ_D . This seems to indicate that the variation of the temperature enhancement can better be described in terms of its dependence on the status of the plasma (e.g. T_{e0} , N_0) rather than its dependence on the parameter λ_D .

Our main finding, i.e. $R(T_e) > 1$ in the very near wake at negative body potentials, is in agreement with *in situ* results. Moreover, the Ariel I and Explorer 31 results (Henderson and Samir, 1967; Samir and Wrenn, 1972) pointed towards the possibility that the enhancements may be confined to the very near wake. This is also seen in our results (Fig. 1 and 2).

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