### A SIMPLE ENERGY BALANCE MODEL OF ICE SEGREGATION

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### ABSTRACT

Step functions of normal frost front advance in dry and wet soils are reviewed and used to introduce the geometry of the soil frost system. An energy balance equation at the initiation of ice lens growth is developed using simple physical assumptions. This model is then employed to simulate the effects of varied surface temperature depression and water table depth on the timing, position and rate of ice lens growth. Lowering the water table delays the onset of segregation but the initial heave rates are all above half a meter per year in the region examined.

# INTRODUCTION

Step functions in the surface temperature regime have been used to analyze the rate of penetration of a freezing front in soils. These solutions require a knowledge of the thermal conductivities of the unfrozen and frozen soils  $(K_{11}, K_{f})$  the volumetric heat capacity of these materials ( $C_{ii}$ ,  $C_{f}$ ) and the surface temperature  $(T_a)$ . The first step equation to be considered is that of Terzaghi (1952). The material is assumed to be at 0°C at all depths when the surface temperature is lowered to T for all future time. The depth of frost (Z<sub>f</sub>) in this material, which contains no water, and thus is free of latent heat effects, is shown as equation 1 where no subscripts are employed in the thermal property notations as the thermal properties are assumed to be the same in frozen and unfrozen

states. The time length since the initiation of the thermal disturbance is (t) in compatible units.

$$Z_f = \sqrt{12(K/C)t}$$
 Terzaghi Equation (1)

Note that the depth of penetration in the "dry case" is only a function of the thermal diffusivity (K/C) and time being independent of the magnitude of the thermal disturbance.

Terzaghi (1952) modified the equation developed by Stefan (1891) for soils with the addition of the effects of an assumed saturated pore space  $(X_p)$  in combination with latent heat of fusion (L) and the energy withdrawal necessary to undercool the frozen layer. This equation is shown as equation 2.

$$Z_{f} = \sqrt{(2K_{f}|T_{a}|t)/(X_{p}L+C_{f}(|T_{a}|/2))}$$
Stefan Equation
(2)

Note that in this case the depth of penetration is dependent upon both the magnitude and duration of the surface thermal disturbance. In this solution the impact of the heat flow from the unfrozen zone is neglected to achieve a simple algebraic solution. However the Neumann Formulation which requires a numerical solution includes heat flow in the unfrozen zone (Jumikis (1966)). It should be noted that the soil heave is limited to  $\left\lceil \frac{\rho_w}{\rho_1} \right\rceil - \rho_w \right\rceil X_p z_f \text{ in this system where } (\rho_w/\rho_1)$ 

is the ratio of water to ice density.

Lastly, consider segregation ice freezing in which the heave is not theoretically

bounded when it occurs over a stable water

table. The geometry of this system is abstracted in Fig. 1. Arakawa (1966) demonstrated that a segregation efficiency index could be computed based on the heat and water flux rates in the soil. Here we use additional notation  $\mathbf{K}$  and  $\mathbf{H}$  for the hydraulic conductivity of soil water and soil water potential. The notation employed is listed as an appendix.

The formulation of the Arakawa Model is as follows

$$Q_f = K_f \left(\frac{\partial T}{\partial Z}\right)_f$$
 heat flow in frozen zone (3.1 toward surface

$$Q_u = K_u \left(\frac{\partial T}{\partial Z}\right)_u$$
 heat flow in unfrozen (3.2) zone toward freezing plane

$$Q_{W} = LK \left(\frac{\partial H}{\partial Z}\right)_{U}$$
 heat necessary to freeze (3.3) water arriving in the freezing region

$$Q_{\uparrow} = Q_{f} - Q_{u}$$
 total heat loss freezing plane (3.4)

If 
$$Q_{\downarrow} > Q_{\downarrow}$$
 normal freezing (3.5)

If 
$$Q_{\downarrow} \leq Q_{\downarrow}$$
 segregation freezing (3.6)

It is also possible to set up a segregation index (E).

$$E = Q_{W}/Q_{\uparrow} \qquad \text{If } E \ge 1., \text{ ice lensing} \qquad (3.7)$$

In review there are three models for increasingly complex modes of soil frost; these are abstracted in Table 1. The Terzaghi and Stefan solutions are used in civil engineering practice. It will be recalled that  $|T_0|$  •t in the Stefan solution can be expressed as the frost degree-day integral. The Arakawa Model has been used as a conceptual device for setting up computer models of coupled flow with rather awesome complexity (Guymon and Luthin (1974), Harlan (1975)). However numerical models of coupled flow become extremely sensitive to iteration frequency when vapor transport is added to heat and water flux (Outcalt (1979)). Earlier Palmer (1967) produced a simple analytical solution to the ice lensing problem. The

intention of this paper is to produce another extremely simple model of the process using equations 3.5 and 3.6.

### THE STRUCTURE OF THE SIMPLE MODEL

The energy balance geometry sketched in Fig. 1 is somewhat similar to the Stefan Model. However, water flux is added and the equation is set to represent the equilibrium that must exist when ice segregation begins. If equilibrium is a reasonable assumption then the duration of time necessary for the freezing region to deepen to the level at which segregation occurs can be estimated by solving the Stefan equation for time. The heave rate could then be expressed by equation 4.

$$dh/dt = \left(\rho_w/\rho_i\right) + \frac{H_v-H_f}{Z_w-Z_f}$$
 heave rate (4)

The problem is now to specify the state variables (T $_{\rm f}$ , H $_{\rm f}$ ) and the system parameter  ${\rm K}$ (see Appendix). A parsimonious and admittedly expedient approach is to assume that segregation occurs at some slightly subfreezing temperature (-0.1°C). If there is a constant relationship  $(\partial H/\partial T = 12.43 x)$  $10^3$  cm./°C) below 0°C then H(T) and  $\frac{1}{10}$  K(H) can be specified. Then  $\rm H_{\rm f}$  is -1243. cm.  $\rm H_{\rm 2}0$  and the value of  $K(H(-.05^{\circ}C))$  is representative of the value of unsaturated hydraulic conductivity in the region where water is moving toward the segregation region (at -0.1°C) through the freezing region. We used available data for a sandy loam in the calculations ( $K = 1. \times 10^{-7}$  cm. sec. -1). It is assumed further that the water potential gradient is linear between the water table and the "f level." In this simple model, again in the interest of parsimony, no adjustment was made for gravity potential as that involves assumptions concerning the time dependent evolution of the potential gradients which

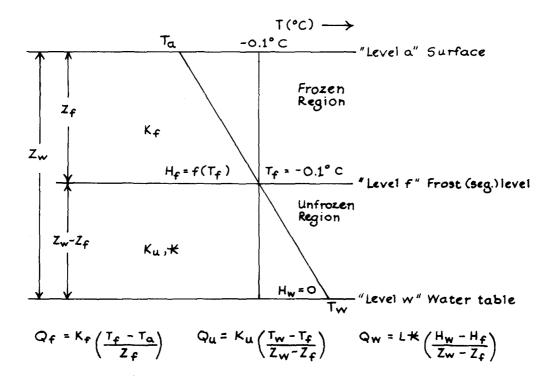


Fig. 1. Geometry of Segregation Ice Formation System at the Onset of Lens Growth  ${\rm Q_f}\text{-}{\rm Q_u}\text{-}{\rm Q_w}$  = 0.0 (see text)

Table 1. Summary of Frost Models

Frost Type	Step Function Model	Soil Heave	Depth of Frost a function of
Dry Frost	Terzaghi	None	K <sub>f</sub> , C <sub>f</sub> , t
Wet Frost	Stefan	Limited to $[(\rho_w/\rho_i)-\rho_w]^{X}_{p}^{Z}_{f}$	$^{\mathrm{K}}_{\mathrm{f}}$ , $^{\mathrm{C}}_{\mathrm{f}}$ , $^{\mathrm{t}}_{\mathrm{a}}$
Segregation Frost	Arakawa	No theoretical limit	Initially as with wet frost then terminated by segregation

may or may not be in hydraulic contact with the water at the -0.1°C isotherm especially during the initial stages of frost advance toward a deep water table. The bulk thermal conductivities of the frozen and unfrozen regions were calculated on the basis of a porosity of 38%, again assuming saturation and total frost or thaw in these zones. These assumptions while being extremely liberal permit the formulation of the equilibrium expression in a nice transcendental form reminiscent of the Neumann solution.

$$\begin{split} & \text{K}_{\mathbf{f}}(\frac{\text{T}_{\mathbf{f}}^{-\text{T}}\text{a}}{\text{Z}_{\mathbf{f}}}) - \text{K}_{\mathbf{u}}(\frac{\text{T}_{\mathbf{w}}^{-\text{T}}\text{f}}{\text{Z}_{\mathbf{w}}^{-\text{Z}}\text{f}}) - \text{LK}(\frac{\text{H}_{\mathbf{w}}^{-\text{H}}\text{f}}{\text{Z}_{\mathbf{w}}^{-\text{Z}}\text{f}}) = \text{f}\left(\text{Z}_{\mathbf{f}}\right) \text{ (5)} \\ & \text{where } \text{f}\left(\text{Z}_{\mathbf{f}}\right) \neq 0.0 \text{ by the Newton method.} \\ & \text{Convergence to within 1 x 10}^{-10} \text{ cal. cm.}^{-2} \\ & \text{sec.}^{-1} \text{ was achieved in under ten iterations} \\ & \text{using initial values for Z}_{\mathbf{f}} \text{ of (0.2 cm.)} \text{ and } \\ & \text{(Z}_{\mathbf{w}} - 0.2 \text{ cm.)}. \end{split}$$

## RESULTS OF SIMULATION TESTS

The behavior of this rather crude analog was analyzed by setting the temperature of the water table to +1°C and varying the water table depth over the range of 10 - 100 cm. in 10 cm. steps and the surface temperature over the range of -2 to -20 °C in 2 °C steps. This process yielded three matrices showing distance from the water table to the lens base distance (cm.), time at which segregation began (days) and heave rate due to segregation (mm./day). The initial heave due to normal in-situ frost is  $\left| \left( \frac{\rho_w}{\rho_{\, {f i}}} \right) \right| - \rho_w \left| X_p Z_{\, {f f}} \right|$ . The matrices produced by these simulations are presented in Table 2. The following generalization can be drawn about model performance from these results.

First the unfrozen thickness at segregation decreases with lowered surface temperature at each water table depth. Therefore normal frost depth decreases with increasing surface temperature. Segregation is thus

delayed by lowered surface temperature until the frozen zone temperature gradient is reduced and the water potential gradient increases as normal frost advances. The length of time until the start of segregation is influenced by both the magnitude of the temperature depression at the surface and the depth of normal freezing (equation 2). These two factors interact in such a way to produce a maximum time to segregation at a temperature depression of approximately -6°C at all water table depths tested. At lowered surface temperatures the magnitude of the temperature depression overrides frost depth effects to reduce the time at which segregation begins. At all water table depths the increased surface temperature depression increases the heave rate. The largest heave rates occur with large temperature depressions. However, it should be noted that over long periods of time even the lowest heave rate (1.67 mm./day) is equivalent to .5 m./yr. Increased water table depths delay the onset of segregation but eventually substantial heave rates are produced.

### CONCLUSIONS

The object of this paper is first to describe the energy budget model for ice lensing in the most simple manner possible. In the process a considerable amount of physical detail was sacrificed on the "altar of parsimony." This lost detail may be significant.

However given laboratory tests it may be possible to calculate values for the "apparent hydraulic conductivity" and utilize the solution methods described here for engineering calculations in much the same manner as the Stefan solution is now employed.

Table 2. <u>Simulation Matrices</u> (see text)

				Suri	face Temp	perature	(°C)						
	-2	-4	-6	-8	-10	-12	-14	-16	-18	-20			
10	7.31	5.69	4.66	3.94	3.41	3.01	2.69	2.42	2.20	2.02			
20	14.61	11.33	9.32	7.88	6.83	6.48	5.37	4.85	4.41	4.04			
30	21.92	17.07	13.97	11.82	10.24	9.01	8.05	7.27	6.61	6.06			
40	29.22	22.76	18.63	15.76	13.65	12.03	10.74	9.69	8.82	8.08			
50	36.53	23.45	23.29	19.70	17.06	15.04	13.43	12.12	11.03	10.10			
60	43.84	34.14	27.95	23.62	20.47	18.04	16.11	14.54	13.23	12.12			
70	51.14	39.83	32.61	27.59	23.89	21.05	18.80	16.96	15.43	14.14			
80	58.45	45.52	37.27	31.53	27.30	24.07	21.48	19.39	17.64	16.15			
90	65.75	51.21	41.93	35.46	30.71	27.06	24.16	21.81	19.84	18.17			
.00	73.06	56.90	46.59	39.40	34.13	30.07	26.35	24.25	22.05	20.20			

THICKNESS OF UNFROZEN ZONE AT SEGREGATION (CM.)

					Sur	face Tem	perature	(°C)						
		-2	-4	-6	-8	-10	-12	-14	-16	-18	-20			
]	10	0.10	0.13	0.14	0.13	0.13	0.12	0.12	0.11	0.11	0.10			
2	20	0.41	0.53	0.55	0.54	0.51	0.46	0.47	0.44	0.42	0.41			
3	30	0.91	1.19	1.23	0.21	1.16	0.11	1.05	1.00	0.95	0.91			
4	40	1.62	2.11	2.19	2.15	2.06	1.96	1.87	1.78	1.69	1.62			
-	50	2.53	3.29	3.42	3.35	3.22	3.07	2.91	2.78	2.65	2.53			
6	50	3.65	4.74	4.93	4.84	4.63	4.41	4.20	4.00	3.81	3.65			
7	70	4.97	6.45	6.71	6.57	6.31	6.01	5.72	5.44	5.19	4.96			
8	30	6.49	8.43	8.77	8.58	8.24	7.84	7.46	7.11	6.78	6.48			
9	90	8.21	10.67	11.09	10.87	10.42	9.93	9.45	8.99	8.58	8.21			
10	00	10.14	13.17	13.69	13.41	12.87	12.26	11.67	11.10	10.59	10.13			

TIME TO START OF LENS GROWTH (DAYS)

Table 2 (continued)

				Suri	face Temp	perature	(°C)			
	-2	-4	-6	-8	-10	-12	-14	-16	-18	-20
10	15.98	20.51	25.06	29.62	34.20	38.83	43.39	48.18	52.95	57.76
20	7.99	10.26	12.53	14.81	17.03	18.00	21.73	24.09	26.48	28.90
30	5.33	6.84	8.35	9.87	11.40	12.96	14.49	16.06	17.65	19.27
30	3.99	5.13	6.27	7.41	8.55	9.71	10.37	12.05	13.24	14.45
50 60	3.20	4.10	5.01	5.92	6.84	7.76	8.69	9.64	10.59	11.56
60	2.66	3.42	4.18	4.94	5.70	6.47	7.25	8.03	8.82	9.63
70	2.23	2.93	3.58	4.23	4.89	5.55	6.21	6.88	7.56	8.26
80	2.00	2.56	3.13	3.70	4.23	4.85	5.43	6.02	6.62	7.23
90	1.73	2.23	2.78	3.29	3.80	4.31	4.83	5.35	5.88	6.42
100	1.60	2.05	2.51	2.96	3.42	3.88	4.35	4.81	5.29	5.73

INITIAL HEAVE RATE AT SEGREGATION (MM./DAY)

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Appendix. Notation

	Notation	Units	Interpretation
Boundary Conditions	T <sub>a</sub>	°C	surface temperature
(const.)	$^{\mathrm{T}}$ w	°C	water temperature at water table
	$^{ m H}_{ m w}$	cm	water potential at water table (0.)
	<sup>T</sup> f	°C	temp. where segregation occurs $(-0.1^{\circ}\text{C})$
	н <sub>f</sub>	cm	water potential at $T_f(-1243.)$
	Z w	cm	distance from surface to water table
Parameters	K <sub>f</sub>	cal cm <sup>-1</sup> sec <sup>-1</sup> °C <sup>-1</sup>	thermal conductivity of frozen region
	K <sub>u</sub>	$\mathrm{cal}~\mathrm{cm}^{-1}\mathrm{sec}^{-1}\mathrm{°C}^{-1}$	thermal conductivity of unfrozen region
	*	cm sec -1	hydraulic conductivity at -0.05°C
Variable	z <sub>f</sub>	cm	distance from surface to "f level"
Const.	L	cal cm <sup>-3</sup>	volumetric heat of fusion (80.)