CRITICAL EXCITON ANNIHILATION: DIFFUSION, PERCOLATION OR ANDERSON TRANSITION?*

P. W. Klymko and R. Kopelman^{*} Department of Chemistry, University of Michigan, Ann Arbor, Mi 48109, U.S.A.

Guest-Guest Triplet-Triplet annihilation in highly purified mixed crystals of naphthalene (guest) naphthalene-d₈(host) was studied at 1.8 K via spectrally and time-resolved phosphorescence and delayedfluorescence. We find significant <u>deviations</u> from current diffusion theories and homogeneous kinetics but consistency with local heterogeneities (clusterization) and dynamic exciton percolation, as well as with long-range annihilation (via superexchange). There is negative evidence for an Anderson transition.

The transport of excitations in disordered media has several interesting aspects: 1) Potential observation of an Anderson-Mott transition [1]. 2) Tests of current theories on diffusion in disordered media [2-4]. 3) Search for the effects of local heterogeneities ("clusters") and their relation with critical phenomena [5]. 4) Simple, reproducible, systems for studying the principles of heterogeneous kinetics and interface reactions. 5) The possibility of obtaining biomimetic systems. 6) The study of energy upconversion via long-range annihilation (exciton fusion).

The steady-state excitation (Xe lamp) of highly purified isotopic mixed naphthalene crystals ($C_{10}H_{\rm B}$ in $C_{10}D_{\rm B}$ at 1.8 K) is shuttered-off and spectrally resolved phosphorescence (0-0) and delayed fluorescence (0"-512")decays are monitored as a function of time, light intensity guest ($C_{10}H_{\rm B}$) concentration etc. Crystals with 0.1 % mole guest show no delayed fluorescence (time resolution ≈ 1 ms) and just normal triplet decay curves ($\tau = 2.7$ s). All host homofusion, heterofusion, prompt fluorescence, radiative trapping, intersystem crossing etc. are over within 1 ms and no guest homofusion takes place. However, guest-guest homofusion (triplet-triplet annihilation) is the phenomenon studied at higher guest concentrations. The negligible importance of any third channel of decay (i.e. supertrapping by impurities or X-traps) was demonstrated both spectroscopically and kinetically (Fig. 1). The absolute and relative delayed fluorescence (DF) and phosphorescence (P) rates were studied as a function of light (3 orders of magnitude) and time. An explicit example of a DF decay curve

- * Supported by NIH Grant 2 ROI NS 08116-10A1
- + Visiting Professor, Physik, Inst. 3, Univ. of Stuttgart, D - 7000 Stuttgart-80, W.-Germany

457

0 022-2313/81/0000-0000/\$02.75 © North-Holland



 $\label{eq:fig.1: Sample Delayed Fluorescence and "differential" Phosphorescence under identical conditions (A = measured phosph. decay rate <u>minus</u> "natural" one).$



Similar curves have been obtained for different light intensities and guest concentrations (4-20 %), with a constant (\pm 10 %) ratio of 2/(DF). The "differential" phosphorescence decay rates 2 (Fig. 1) would be the ones observed in the limit of zero "natural" decay rate rather than the real one (0.37 s⁻¹) obtained from the low quest concentration samples (0.1 %). Fig. 2 shows the relative DF and P rates at various points in time for given samples. Upon extrapolation of DF (\triangle) and P to t = 0 one obtains in Fig. 3 the effective steady-state annihilation rate.



Fig. 3:

Guest Concentration Dependence of Annihilation, at constant excitation per guest site (probability from phosphorescence and "differential" phosphorescence decays extrapolated to t = 0). The quasicritical behavior of Fig. 3 is interesting but similar to previous energy transport work measured via supertrapping [5]. However, the most striking result is found in Fig. 2 for the samples at and below the "critical" guest concentration (C₂). While for the 20 % guest sample (Fig. 2) one gets (DF) \propto P², as has been standard in previous triplet-annihilation work [5,6], for the 4 % sample the result in the high exciton density range is clearly (DF) \propto P² where X > 2. We note that P is linear with the <u>exciton</u> density (in the guest) due to the monomolecular nature of the "natural decay", and, similarly, DF has usually been expected to be quadratic with <u>exciton</u> density due to the bimolecular nature of the annihilation process [5,6].

In the limit of long times $(t \rightarrow \infty)$, or for steady-state conditions, one gets from current theories of diffusion in disordered media essentially the same answer as from diffusion theories in homogeneous lattices, namely X = 2. On the other hand, the clusterization ("percolation") approach allows for both X = 2 and X > 2, depending on specific conditions, as has recently been demonstrated with the aid of Monte-Carlo calculations [7].

The dependence of C (the "critical guest concentration", where (DF) \cong P) on excitation intensity (Fig. 3) and on time cannot easily be reconciled with an Anderson transition [1]. This is especially true in view of the further shifts in C observed [5] in supertrapping experiments on the same systems ($C_{10}H_8/^{\rm C}C_{10}D_8$). The experiments are also consistent with long-range annihilation (via superexchange [8,9]).

In conclusion, the annihilation experiments are not well described by current theories of diffusion in heterogeneous media. Our system can be considered as a model for heterogeneous kinetics. As both excitation transfer and annihilation processes are most probably limited to the <u>ab</u> crystal plane [5], this is actually a model system for interfacial reactions in synthetic and biological systems. We further believe that cluster-type exciton annihilation may be important for the study of photosynthetic systems.

- [1] Klafter, J. and Jortner, J., Chem. Phys. Lett. 49 (1977) 410.
- [2] Blumen, A. and Silbey, R., J.Chem.Phys. 70 (1979) 3707.
- [3] Blumen, A., Klafter, J. and Silbey, R., J.Chem.Phys. 72 (1980) 5320.
- [4] Godzik, K. and Jortner, J., J.Chem.Phys. 72 (1980) 4471.
- [5] Francis, A.H. and Kopelman, R., Excitation Dynamics in Molecular Solids, in Laser Spectroscopy of Solids, ed. Yen W.M. and Selzer P.M. (Springer, Berlin, 1981) P. 241-302.
- [6] Merrifield, R.E. Pure Appl.Chem. 27 (1971) 481; Suna, A., Phys.Rev. B <u>1</u> (1970) 1716.
- [7] Newhouse, J., Hoshen, J. and Kopelman, R. (unpublished).
- [8] Sternlicht, H., Nieman, G.C. and Robinson, G.W., J.Chem.Phys. 38 (1963), 1326; 39 (1963), 1610.
- [9] Jortner, J. and Kopelman R. (unpublished).