AN ESR AND X-RAY DIFFRACTION STUDY OF A FIRST-ORDER PHASE TRANSITION IN CdPS3

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The ESR spectra of Mn^{+2} doped single crystals of CdPS $_3$ indicate a first-order phase transition (monoclinic to orthorhombic) at ca. 260 K. This is confirmed directly by an X-ray diffraction examination of the doped and undoped crystals at various temperatures.

Introduction

The transition metal chalcogenophosphates form a series of isomorphous compounds with the general chemical formula MPX $_3$ where M is a row-four transition metal and X is either S or Se. A full crystal structure has been reported only for the room-temperature iron compound of sulfur. Above, the room-temperature lattice parameters and crystal space group have been determined for all compounds in the series. At ambient temperature, the transition metal sulfides are reported to be either monoclinic or trigonal and the selenides, trigonal.

These materials all crystallize in a layered structure with planes of hexa-coordinated metal atoms sandwiched between planes of hexagonally packed chalcogen atoms. The phosphorus atom-pairs occupy sites in the metal atom plane. The metal atom coordination geometry is trigonally-distorted octahedral in both the monoclinic and the trigonal forms. The macroscopic structure is built from repeated layers of the X-M-X sandwich structure held together by weak van der Waals forces between the adjacent chalcogenide planes.

The layered structure of the MPX $_3$ compounds is similar to that of the more extensively studied transition metal dichalcogenides, MX $_2$. The MPX $_3$ structure may be viewed as the MX $_2$ structure with 1/3 of the metal replaced by a P $_2$ atom-pair. The effect of intercalation on the charge transport properties of MX $_2$? Intercalation compounds has been extensively investigated. Relatively little research in this area has been reported for the intercalation compounds of MPX $_3$ materials, to which end the present study is a prelude. Extensive investigations of structural phase transitions in the quasi two-dimensional transition metal dichalcogenides have been reported, while there have been no reports to date of similar transitions in the iso-structural chalcogenophosphates.

It is the purpose of the present communication to report on the observation and characterization of a first-order phase

transition in the compound CdPS3. The phase transition was observed in single crystals using x-ray diffraction techniques and ESR spectroscopy of $\rm Mn^{+2}$ substitutional impurity centers.

Experimental

CdPS $_3$ was synthesized from a powdered mixture containing stoichiometric amounts of CdS, phosphorus (red) and sulfur. Samples for use in the ESR studies were doped with 1 at.% Mn.

the ESR studies were doped with 1 at.% Mn. Large single crystals $(5\times5\times0.1~\text{mm}^3)$ were grown from the vapor phase³ in evacuated Vycor tubes. The tubes were placed in a two-zone vertical tube furnace. The material was sublimed from the hotter to the cooler zone. A temperature gradient of 943 K \rightarrow 873 K produced excellent crystal growth in approximately one week. Elemental analysis of CdPS₃ crystals prepared by this method was in agreement with the empirical formula (weight %Cd=45.89, %P=12.24, %S=41.83).

For ESR investigation, crystals were oriented conoscopically by identifying the optic axes in the ab crystal plane. The b-axis could be identified from the crystal morphology. Single crystals were mounted on a rotating sample stage attached to the microwave cavity of a Brücker model ER200E X-band ESR spectrometer and spectra were recorded as a function of the angle of rotation about the crystallographic a, b and c* (axb) axes of the room-temperature monoclinic structure. The temperature of the crystal could be varied using a Varian Model 4257 variable temperature accessory modified to fit the microwave cavity.

For X-ray diffraction studies, small single crystals were mounted on a Huber model 100 Weissenberg camera fitted with a liquid nitrogen cooling attachment (stability ± 0.2 K). The crystal was oscillated around various crystallographic axes in the ab plane in order to pick up reflections sensitive to the observed structural distortion, in this case the monoclinic angle, β . A small number of such reflections were isolated by appropriate choice of oscillation range and zero-layer Weissenberg screens. In order to record changes occurring at

the phase transition, the X-ray film was translated stepwise (every ~ 10 minutes) as the temperature was varied in 3 K steps. A 1kW sealed-tube source of Cu K radiation was used throughout. In this way, a rather striking record of variations in lattice parameters across the transition is obtained. Further details of the method appear in Ref. 9.

Results

The angular dependence of the ESR spectrum of Mn⁺² centers in CdPS₃ was recorded over the temperature range 140 K to 500 K. Above 300 K, the spectrum was characteristic of a single manganese site (spectrum A, Figure 1); however, as the crystal was cooled below about 220 K, all the lines in the spectrum diminished in intensity and a second set of lines (spectrum B, Figure 1) appeared which gradually became dominant at the lowest temperatures attained. The fractional intensity of the high temperature spectrum is plotted as a function of temperature in Figure 1 where the small arrows indicate the direction of temperature change during measurement. The observed intensity variations are reproducible and exhibit substantial hysteresis.

The marked behavior shown in Figure 1 is also seen, in Figure 2, to be accompanied by strong discontinuities in at least one of the lattice parameters over approximately the same temperature range, confirming that a first-order structural transition is responsible for the observed changes in the ESR spectrum. On cooling from the room-temperature monoclinic phase (see Figure 2(b)) a new set of X-ray reflections appear at T^+ = 248 K. These reflections are indexed on the basis of an orthorhombic cell with the same contents as the original monoclinic cell. Full Weissenberg data of the a*c* plane, taken at T = 215 K, confirms that indeed β changes discontinuously from 107°, its value at room temperature, to precisely 90° at the phase transition. In contrast, only small changes in the other lattice parameters determining interplanar spacings are observed at the phase transition (see Table 1). On heating (Figure 2(a)), the transition is seen to be reversible, occurring at T^{\dagger} = 264 K.

Discussion

The data presented in Figures 1 and 2 suggest the existence of a first-order phase transition centered at about 260 K. Between approximately 220 and 290 K both phases co-exist as is clear from the multiple site nature of the ESR spectrum. The angular dependence of the ESR spectrum of $\rm Mn^{+2}$ centers in CdPS3 obtained both at ambient temperatures and below, could be accurately reproduced by theoretical calculation with the following spin-Hamiltonian, which is suitable for orthorhombic or lower symmetry $\rm ^{10}$

$$H_{s} = \beta_{o} \cdot H \cdot g \cdot S + D[S_{z}^{2} - 1/3S(S+1)] + E(S_{x}^{2} - S_{y}^{2}) + 1/6a[S_{x}^{4} + S_{y}^{4} + S_{z}^{4}]$$
$$- 1/5S(S+1)(3S^{2} + 3S - 1)] + A_{zz}S_{z}I_{z} + A_{xx}S_{x}I_{x} + A_{yy}S_{y}I_{y}.$$

where the first term represents the electronic Zeeman interaction, the second and third terms represent, respectively, the axial and rhombic

Figure 1: Fractional Intensity $(\frac{I_A}{I_A + I_B})$ of the ESR spectrum of Mn(II)-doped CdPS3. Insets show A and B phase spectra and arrows indicate direction of cooling and heating.

parts of the fine structure (zero field splitting), the fourth term represents the cubic part of the crystal field and the last term represents the nuclear Zeeman (hyperfine) interaction S=I=5/2 for Mn(II) (d^5). It was assumed that the electronic Zeeman interaction, the zero-field interaction and the hyperfine interaction all have the same set of principal axes (x, y and z). The y-axis of the spin-Hamiltonian was chosen parallel to the two-fold crystallographic axis (b-axis) of the room temperature monoclinic lattice. The x and z principal axes then lie in the crystallographic ac plane perpendicular to \vec{b} .

The angular variation of the ESR spectra was analyzed by numerical diagonalization of the matrix of the spin-Hamiltonian, Equation 1. The parameters of the spin-Hamiltonian for both the room-temperature and low-temperature phases of CdPS3 are given in Table 1. These were obtained by using an iterative procedure to fit the theoretical spectrum to the experimental data.

The fact that the crystal field parameters (D, E and a) all have non-zero values, and that g shows a small anisotropy for the room temperature phase indicates that Mn $^{+2}$, which substitutionally replaces Cd $^{+2}$, has nearly axial symmetry with a slight rhombic distortion. 11 , 12 This result is consistent with the proposed monoclinic room-temperature structure of CdPS $_3$. The zero

value obtained for the rhombic distortion parameter in the low-temperature phase suggests a rigorously orthorhombic lattice. This is

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Table 1					
Symmetry	Monoclinic (300 K)	Orthorhombic (143 K)			
D E Axx Ayy Azz 9xx 9yy 9zz a	-365±2 -6±1 -69.2±1 -69.4±1 -70.5±1 2.006±0.002 2.007±0.002 2.001±0.002 0.7±0.2	-353±2 0 -69±2 -69±2 -70±2 2.005±0.001 2.005±0.001 2.002±0.001 0.3±0.2			
à + b	6.17Å* 10.67Å	6.18±0.03Å** 10.67±0.05Å			
ċ	6.82Å	6.40±0.02Å			
β	107.1				

Note: The units of D, E, A_{ij} and a are 10^{-4} cm $^{-1}$.

^{*}Ref. 1. **This study.

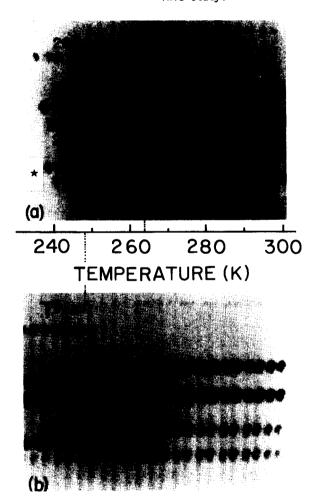


Figure 2: Sequential X-ray oscillation exposures for a single crystal of undoped CdPS3, a) on heating, b) on cooling. * denotes reflections formed by the β -component of the radiation. The indexed spots are from Cu $K\overline{\alpha}$.

confirmed directly by X-ray diffraction in the low-temperature phase. The establishment of a mirror plane parallel to a* at the transition, together with systematic absences of the type $h + k \neq 2n$, identifies Cmm2 and Cmmm as possible low-temperature space groups, the former being distinguished by non-centrosymmetry. At this point, a full structural determination has not been carried out.

The thermal hysteresis of T^{\uparrow} - T^{\downarrow} = 16 K, as shown in Figure 2, is considerably smaller than the value of 40 K given by the ESR results in Figure 1. This difference is due to the sensitivity of the ESR technique to small amounts of the minor coexisting phase which occurs, respectively, at the upper and lower edges of the hysteresis loop in Figure 1. This coexistence is also evident, but less so, in Figures 2(a) and 2(b).

Lattice parameter measurements, listed in Table 1, indicate that the pseudo-hexagonal arrangement of the ab layer is well preserved through the transition (i.e., $b \approx a \sqrt{3}$) and that the structural transition involves principally a sliding motion of these layers, aided by the weak van der Waals bonding. This mechanism is also consistent with the small (2%) decrease in the interlayer spacing at the transition.

There are two unusual features of the phase transition in CdPS3 which should be mentioned. Firstly, the change in monoclinic angle ($\Delta\beta$ = 17.1°) at the transition is surprisingly large; especially since this implies relative displacement of adjacent layers by as much as 2A. Secondly, the structural symmetry appears to increase on cooling, with the loss of the monoclinic-phase unique axis and the establishment of at least one extra mirror plane in the orthorhombic phase. This is the reverse of the familiar cases in which the symmetry is broken by the appearance of an order parameter in the low-temperature phase. A similar transition has been reported 13 in the monoclinic layer compound β -MoTe₂, the low-temperature phase again being orthorhombic. The reason for the instability of these monoclinic compounds is not understood at

this time. There is no evidence for charge density waves or incommensurate behavior in either of these monoclinic structures in the form of electron diffraction satellites.

Conclusion

The ESR spectra of Mn⁺² doped CdPS₃ single crystals and the x-ray diffraction of doped and undoped crystals are consistent with a first-order structural phase transition from the room-temperature monoclinic lattice to a low-temperature orthorhombic lattice occurring at about 260 K. The transition is reversible and hysteretic and involves considerable sliding of adjacent layers across the van der Waals gap.

We have also examined the ESR spectra of the Mg and Zn compounds and have found no evidence of a phase transition in the temperature range 140-500 K. A preliminary analysis of the angular dependence of the ESR spectra of these compounds indicates that the lattices are orthorhombic at room temperature. Previous x-ray diffraction studies of these materials at room temperature have indicated a monoclinic lattice. Thus, it is likely that the Mg and Zn chalcogenophosphates exhibit both monoclinic and orthorhombic modifications.

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