# URANIUM CHALCOGENIDES-III

## HEAT CAPACITIES AND THERMODYNAMIC PROPERTIES OF US<sub>1.0</sub> AND USe, FROM 5 TO 350°K\*

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Abstract – Heat capacities have been measured and derived thermodynamic properties calculated for the structurally related US<sub>1.9</sub> and USe<sub>2</sub> phases from 5 to 350°K. Values of the heat capacity  $(C_p)$ , entropy  $(S^{\circ} - S_0^{\circ})$  and Gibbs energy function  $[-(G^{\circ} - H_0^{\circ})/T]$  at 298·15°K in cal mole<sup>-1</sup> °K<sup>-1</sup> are 17·68, 25·91 and 13·82 for US<sub>1.9</sub> and 18·92, 31·98 and 17·86 for USe<sub>2</sub>, respectively. Schottky-type anomalies are present in US<sub>1.9</sub>; a  $\lambda$ -type thermal anomaly with  $\Delta St = 0.19$  cal mole<sup>-1</sup> °K<sup>-1</sup> occurs at 13·1 °K in USe<sub>2</sub>.

THE PRESENCE of paramagnetism in the uranium chalcogenides makes a study of their thermal properties seem of special interest for exploration of the mechanisms by which spin order is approached. For  $US_{1.9}$  and  $USe_2$ , both with tetragonal structures, magnetic susceptibility measurements [1-3] indicate the presence of 2 unpaired 5*f* electrons with comparatively weak atomic interactions between neighbors since the Curie-Weiss law is fulfilled at higher temperatures with a Weiss constant of about  $-50^{\circ}$ K. A structural study [4] shows that some of the uranium atoms in the tetragonal  $US_{2-x}$  structure are distributed at random; this might cause further anomalies in the behavior of the heat capacity or in the zero point entropy.

### **EXPERIMENTAL**

The preparation of the US<sub>1.9</sub> and USe<sub>2</sub> samples was carried out by reacting elements of previously described [1] purity in evacuated and sealed quartz tubes in the right proportions. In the case of the sulfide, the reaction was performed by heating the mixture to 970°C over a period of 3 weeks and keeping it at this temperature for 2 days. After cooling, the sample was then crushed and homogenized at 800°C for 1 week. After the heat capacity determinations were completed, chemical analyses were made on the US<sub>1.9</sub> sample by K. Jensen of the Argonne National Laboratory through the cooperation

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- 1. F. Grønvold, H. Haraldsen, T. Thurmann-Moe and T. Tufte, J. inorg. nucl. Chem. 30, 2117 (1968).
- 2. W. Trzebiatowski and W. Suski, Bull. Acad. pol. Sci. Sér. Sci. chim. 9, 277 (1961).
- 3. V. I. Chechernikov, A. V. Pechennikov, E. I. Yarembash, L. Martynova and V. K. Slavyanskikh, Soviet Phys. JETP 26, 328 (1968).
- 4. R. C. L. Mooney Slater, Z. Kristallogr. 120, 278 (1964).

of H. E. Flotow. The results indicated 79.78  $\pm$ 0.10 per cent by weight of uranium, 20.15  $\pm$ 0.07 per cent by weight of sulfur by a distillation process and 20.42  $\pm$ 0.09 per cent by weight by gravimetry. (Theoretical: U = 79.62, S = 20.38). The reported results have been corrected for high recovery on standards. Hydrogen, oxygen, and nitrogen were  $1.3 \pm 0.4$  ppm, 470  $\pm$  30 ppm, and 25  $\pm$ 6 ppm respectively according to L. F. Krout of the same laboratory. Spectrochemical analysis by J. Lech indicated in (ppm by weight, factor of two precision): Al ~ 10, As < 20, Cr ~ 4, Fe ~ 20, K < 100, Mg ~ 5, Mo < 20, P < 100, Si ~ 60, Ti < 200 and Zn < 50. Spark spectral analyses from the Chemical Engineering Department of ANL indicated in (ppm atomic): Cr ~ 24, Fe ~ 220, K ~ 21, Mn ~ 20, P ~ 67 and Sc ~ 20. Mass spectrographic results showed Al ~ 13, Mg ~ 33, Si ~ 1500 in the same units.

The selenide was prepared in the same way, but the temperature of homogenization was 840°C. Analyses similarly performed at the Argonne National Laboratory on the diselenide sample indicated

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Uranium: 59.92 \pm 0.10 per cent by weight, 3 determinations (theoretical U = 60.12)
Selenium: 39.79 \pm 0.08 per cent by weight, 6 determinations (theoretical Se = 39.88)
Hydrogen: 5.6 \pm 0.3 ppm
Oxygen: 780 \pm 20 ppm
Nitrogen: 12 \pm 1 ppm.
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By spectrochemical analysis (in ppm by weight, factor of 2 precision): Al ~ 10, As < 20, B ~ 2, Cr ~ 4, Cu ~ 10, Fe ~ 30, K ~ 100, Mg ~ 30, Mo < 20, P < 100, Si ~ 40, Ti < 200 and Zn < 50. By spark spectrometry (in ppm atomic): Cr = 12, Fe = 220, Mn = 20, P = 67, Sc = 7. By mass spectrometry (in pom atomic): Al ~ 20, Mg ~ 25, Na ~ 67, Si ~ 1100.

Measurements were made in the Mark III vacuum cryostat[5] provided with an electronic adiabatic-shield control system consisting of three separate channels of recording circuitry with proportional, rate, and reset action. These kept the temperature differences smaller than a millidegree and thereby reduced heat exchange with the calorimeter to a magnitude negligible in comparison to other sources of error. The gold-plated copper calorimeters used (laboratory designation W-29 for US<sub>1.9</sub> and W-28 for USe<sub>2</sub>) had capacities of 42 and 92 cm<sup>3</sup>. The heat capacities of the empty calorimeters were determined separately and small corrections were applied for differences in the amounts of helium gas, indium-tin solder, and Apiezon-T grease for the loaded and empty calorimeters. The heat capacity of the US<sub>1.9</sub> sample (146·551 g), represented about 92 per cent of the total below 10°K and decreased to 66 per cent above 90°K. For the USe<sub>2</sub> sample (196·420 g), the corresponding values are 92 per cent and 60 per cent, respectively. All measurements of mass, temperature, resistance, voltage, and time were based upon calibrations or standardizations made by the U.S. National Bureau of Standards.

#### **RESULTS AND DISCUSSION**

Experimental heat capacity values for the  $US_{1.9}$  and  $USe_2$  samples are presented in Table 1 in chronological order for the mean temperatures of the determinations after correction for curvature of the heat capacity. The data are given in terms of the defined thermochemical calorie, equal to  $4 \cdot 1840$  J, and an ice point of 273 · 15°K. They are considered to have probable errors decreasing from about 3 per cent at 5°K to 0.5 per cent at 10°K and to less than 0.1 per cent above 20°K. For US<sub>1.9</sub>, the shape of the heat capacity curve (see the comparison with US<sub>2</sub>[6] in Fig. 1) is nearly normal, while for USe<sub>2</sub> a  $\lambda$ -shaped anomaly is present, having its maximum at 13.1 °K (see Fig. 2).

The smoothed heat capacities and the thermodynamic functions derived from them by means of a digital computer using a previously described program[7]

<sup>5.</sup> E. F. Westrum, Jr., J. chem. Educ. 39, 443 (1962).

<sup>6.</sup> E. F. Westrum, Jr. and F. Grønvold, J. inorg. nucl. Chem. 30, 2127 (1968).

<sup>7.</sup> B. H. Justice, Doctoral Dissertation, University of Michigan (1961); U.S.A.E.C. Rep. TID-12722 (1961).

Т	C <sub>p</sub>	Т	$C_p$	Т	$C_p$	
Substoi	chiometric i	uranium disi	ulfide [US <sub>1</sub> .	<sub>8</sub> , 1 mole =	299·00g]	
Seri	ies I	276-54	17.42	21·93	2·1 <b>99</b>	
		285.59	17.53	24.25	2.629	
196-93	15.89	294.70	17.63	26.98	3.032	
209.68	16-21	304.22	17.74	30.10	3.383	
		314.12	17.84	32.80	3.676	
Seri	es H	324-06	17.92	36.58	4.093	
Series II		332.32	18.01	40.65	4.554	
138.71	13.605	338.91	18.06	40.41	4.519	
147.00	14.040	346.14	18-11	44.56	4.997	
155-02	14.414	01011	10 11	48.95	5.509	
153.54	14-348	Series III		\$3.60	5.077	
161.82	14.706	Serie	Series III		5.572	
170.56	15.051	6.05	0.114	64.31	7.700	
170.16	15.247	6.40	0.106	70.66	7.090	
107.50	15.60	0.45	0.141	70.00	0 700	
107.30	15.02	7.30	0.220	04.04	8.720	
190.00	12.00	8.34	0.239	84·24	9.403	
204.00	10.08	9.51	0.327	91.32	10-146	
213.37	10.31	10.62	0.451	98.46	10.698	
222-13	16.48	11-92	0.569	104.43	11-229	
231.14	16.68	13-34	0.737	111.64	11.829	
240.20	16.87	14.77	0-942	118.87	12.355	
<b>249∙0</b> 7	17.02	16.30	1.172	126-23	12.846	
<b>258</b> ·23	17.17	18-04	1.463	133-92	13-331	
267-42	17.31	19-92	1.808	141-85	13.771	
	Uranium di	iselenide [U	Se <sub>2</sub> , 1 mole	= 395·99g]		
Seri	ies I	Serie	Series III		Series VI	
71.10	11.478	$\Delta H$ R	lun A	6.14	0.021	
75·90	12.041	11.33	1·028	6.94	0.111	
82·30	12.746			7 <b>·96</b>	0.302	
89.97	13.476	Serie	s IV	9.21	0.539	
<b>98</b> ·11	14.069			10-05	0.736	
106.98	14.647	5.17	0.032	<b>10·7</b> 1	0.864	
116-02	15-17	5.43	0.030	11.27	1· <b>0</b> 16	
125-20	15.64	5.92	0.026	11.77	1.138	
134.57	1 <b>6·06</b>	6.73	0.087	12.21	1.293	
143-68	16· <b>4</b> 0	7.43	0.152	12.60	1.446	
1 <b>52·70</b>	1 <b>6</b> ·71	8.12	0.307	1 <b>2·96</b>	1.619	
161.75	16-96	9.20	0.529	13·29	1.602	
1 <b>70</b> ·75	17·19	10.52	0.833	13.65	1.354	
17 <b>9·73</b>	17.38	11· <b>89</b>	1.190	14.05	1.283	
188-86	17.56	13.44	1.455	14.55	1.330	
198-22	17· <b>75</b>	15-27	1.436	15.12	1.409	
207·59	17· <b>93</b>			15-74	1.504	
216-95	18.08	Serie	es V	16- <b>48</b>	1.624	
226-32	18-22			17-24	1.753	
235· <b>60</b>	18·34	5·50	0.026	17· <b>95</b>	1.875	
2 <b>44·90</b>	18-43	5-97	0.021			
254·19	1 <b>8</b> -53	6.70	0.099	Serie	s VII	
263-53	18-61	7·50	0.150			

Table 1. Heat capacities of US<sub>1.9</sub> and USe<sub>2</sub>\*

T	<i>C</i> <sub>p</sub>	Т	<i>C</i> <sub>p</sub>	T	$C_p$	
272.90	18.71	8.22	0.349	5.55	0.046	
282·23	18.78	<b>9</b> ·27	0·551	$\Delta H R$	un <i>B</i>	
<b>291</b> ·58	18-87	$\Delta Ht$	Run B	19·29	2.118	
300-86	18·99	16·30	16.30 1.595		Series VIII	
310-11	19.01	18-31	1-944			
319-33	19.08	20.34	2.312	12-49	1.432	
328.53	19.14	22.71	2.761	12.63	1.445	
337.59	19.17	25.40	3.283	12.75	1.481	
346.00	19-21	28.62	3.925	12.87	1.541	
		31·87	4.589	1 <b>2·99</b>	1.664	
Serie	es II	34.30	5.096	13.10 1.749		
		36.42	5.537	13.22	1.648	
9.28	0.494	39.22	6.107	13-35	1.550	
9.65	0.620	42.92	6.849	13-51	1.454	
10.40	0.809	47.11	7.668	13-68	1.362	
11-46	1.061	51.58	8-495	13.87	1.326	
12.60	1.455	56-48	9.345	14·06	1.301	
14.16	1.367	61·90	10.233			
15-93	1.533	67.83	11.084			
17.57	1-810	74.34	1 <b>1-871</b>			

Table 1 (Contd.)

\*Units: cal, mole, °K.



Fig. 1. Comparison of the heat capacities of  $US_2[6]$  and  $US_{1:9}$ .



Fig. 2. The heat capacity of USe<sub>2</sub> showing the lambda-type anomaly.

are given in Table 2. The thermodynamic functions are considered to have a precision corresponding to a probable error of less than 0.1 per cent above 100°K.

Since uncertainty exists with regard to the degree of order of spins at 5°K, the entropies obtained by extrapolation in the range  $0-5^{\circ}$ K might not be correct even for use in chemical thermodynamic calculations in which nuclear spin and isotopic mixing contributions cancel. Consideration should also be given to the fact that two (1.9) of the ten uranium atoms in the tetragonal unit cell are reported to be in a four-fold position[4]. If this is also the case at low temperatures, US<sub>1.9</sub> has zero-point entropy of approximately

$$\Delta S = (1/5)R \ln 2 = 0.28 \text{ cal mole}^{-1} \circ \mathrm{K}^{-1}.$$

For these reasons, the entropy values are described as  $(S^{\circ} - S_{0}^{\circ})$  and the Gibbs energy function values must be used with caution.

Approximately one-fifth of the uranium atoms in US<sub>1.9</sub> are surrounded by eight sulfur atoms at the corners of a right antiprism, and the remainder (fourfifths) are surrounded by six sulfur atoms in a less regular arrangement [4]. The molal heat capacity of US<sub>1.9</sub> is somewhat higher than that of US<sub>2</sub> in the range  $5-30^{\circ}$ K, even though there are more atoms and a Schottky transition with a maximum at 25°C in the latter compound [6]. To account for the excess heat capacity in US<sub>1.9</sub>, it appears that one-fifth of the uranium atoms have a singlet ground state and a doublet excited state corresponding to  $\Delta E/k = 56^{\circ}$ K; the other four-fifths have a double ground state and an excited singlet state separated by  $\Delta E/k = 56^{\circ}$ K.

T	$C_p$	S°-S°	H°-H° -	-(G°-H°)/T
S	ubstoichior	netric uran	ium disulfide	2
	[US <sub>1.9</sub> ,	1  mole = 2	99·00g]	
5	0.071	0.023	0.085	0.006
10	0.366	0.144	1.039	0.040
15	0.968	0.396	4.258	0.113
20	1.832	0.788	11.17	0.230
25	2.740	1.297	22.66	0.391
30	3.389	1.859	38.11	0-588
35	3.915	2.420	56.34	0.810
40	4.481	2.981	7 <b>7</b> ·40	1.046
45	5.052	3.542	101.2	1.292
50	5.631	4.104	127.9	1.545
60	6·790	5.233	1 <b>90·0</b>	2.066
70	7.921	6.365	263-6	2.599
80	8.994	7·494	348.3	3.141
90	9.989	8.611	443.2	3.687
100	1 <b>0·90</b>	9.712	547.7	4.234
110	11.71	10·79	660.9	4.781
120	12.44	11.84	781.7	5.326
130	13·09	12.86	909-4	5.867
140	13.67	13.85	1043	6.402
150	14-18	14.81	1183	6-931
160	14.63	15-74	1327	7-453
170	15.03	16.64	1475	7·967
180	15-38	17.51	1627	8·473
190	15.70	18.35	1782	8·971
200	15-98	19-16	1941	9·461
210	16.23	19.95	2102	9.942
220	16.45	20.71	2265	10.41
230	16.66	21.45	2431	1 <b>0·88</b>
240	16.85	22.16	2598	11-33
250	17.03	22.85	2768	11.78
260	17.19	23.52	2939	12.22
270	17.34	24.17	3112	12.65
280	17.47	24.81	3286	13.07
290	17.59	25.42	3461	13.49
300	17.70	26.02	3637	13.90
310	17.80	26.60	3815	14-30
320	17.89	27.17	3993	14.69
330	17.98	27.72	4173	15·08
340	18.07	28-26	4353	15.46
350	18.13	28.78	4534	15-83
273-15	17-38	24.37	3166	12-78
<b>298</b> ·15	17:68	25-91	3605	13-82

Table 2. Thermodynamic properties of  $US_{1.9}$  and  $USe_2^*$ 

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Table 2 (Contd.)

T	Cp	S°-S°	$H^{\circ}-H^{\circ}_{0}-(0)$	$G^{\circ}_{,-}-H^{\circ}_{,0})/T$		
Urani	Uranium diselenide [USe2, 1 mole = 395-99g]					
5	0.029	0.010	0.035	0.003		
10	0.676	0.158	1.289	0.029		
15	1.392	0.641	7.570	0.137		
20	2.250	1.155	16-60	0-325		
25	3.204	1·758	30-21	0.550		
30	4.205	2.430	48-71	0-806		
35	5.242	3·1 <b>56</b>	72.32	1.090		
40	6.267	3.923	101-10	1 396		
45	7.260	4·719	1 <b>34-94</b>	1.720		
50	8.206	5.533	173-6	2.061		
60	9.914	7.185	264-5	2.777		
70	11.343	8.824	371-0	3.524		
80	12.504	10.417	<b>490</b> ·4	4.287		
90	13.439	11· <b>946</b>	620·3	5.054		
100	14.202	13-402	7 <b>58</b> ∙6	5-816		
110	1 <b>4·840</b>	14.787	903·9	6-569		
120	15-38	16-102	1055·1	7· <b>309</b>		
130	15.86	17-352	1211-4	8-034		
140	16.26	18.543	1372-0	<b>8·74</b> 3		
150	16.62	19-677	1536-5	9.434		
160	1 <b>6</b> ·91	20.759	17 <b>04</b> ·2	10·1 <b>0</b> 8		
170	17.17	21.793	1874-6	10·7 <b>66</b>		
180	17.40	22.781	2047-5	11 <b>·406</b>		
190	17 <b>·60</b>	23.727	2222.5	12.030		
200	17.78	24.634	2399-4	12-637		
210	17 <b>·95</b>	25:506	2578-1	13-229		
220	18.12	26.345	2758·4 <sup>,</sup>	13- <b>80</b> 7		
230	18.26	27.153	2940·3	1 <b>4·369</b>		
240	18-39	27 <del>:93</del> 3	3123-6	14·918		
250	18-50	28.69	3308-0	15-45		
260	1 <b>8</b> ·59	<b>29</b> ·41	3493.5	15-98		
270	18.68	30-12	3679.8	16·49		
280	18.76	30-80	3867-0	16·99		
290	18.85	31.46	4055·1	17-47		
300	18-94	32.10	4244-0	17 <b>·95</b>		
310	1 <b>9·02</b>	32.72	4433-8	18-42		
320	19-09	33-33	4624-4	18-87		
330	19·14	33-91	<b>48</b> 15·5	19·32		
340	19.18	34.49	<b>5007</b> ·1	19·76		
350	1 <b>9·24</b>	35.04	5199·2	20.19		
273-15	1 <b>8·70</b>	30-33	3739	16-65		
<b>298</b> -15	18.92	31-98	4209	17 <b>·86</b>		

\*Units: cal, mole, °K.

On subtracting the heat capacity contributions of these two Schottky-type excitations from the observed heat capacities, points are obtained which compare reasonably well (see Fig. 3) with the non-transitional (or "lattice") heat capacity of  $US_{1.9}$ . The latter is estimated on the basis of data for US and  $US_3$  (see Ref. [6]). A more accurate evaluation of the lattice heat capacity is obviously needed. It might well show that the estimated lattice contribution is too high and the Schottky transitions near 25°K in both  $US_{1.9}$  and  $US_2$  have singlet ground states.



Fig. 3. Resolution of the excess heat capacity of US<sub>1.9</sub>. The contribution from the fifth of the atoms with  $\Delta E/k \triangleq 68^{\circ}$ K is shown by - -, that of the other four-fifths by  $-\cdot - \cdot -$  and the sum by -----. The circles and the solid line represent the excess heat capacity obtained by subtracting a lattice contribution estimated by fitting the observed heat capacity at higher temperatures to a suitable Debye function (after trial removal of the Schottky tail), from the observed heat capacity.

In USe<sub>2</sub>, which has the same type of structure as US<sub>1.9</sub>, only a cooperative transition is evident, probably because of our present inability to resolve the lattice heat capacity accurately. The entropy and enthalpy of the  $\lambda$ -type transition were estimated after fitting a lattice contribution through this region.

So far no magnetic or crystallographic work has been carried out below the transition, but the negative Weiss constant indicates that the compound probably becomes anti-ferromagnetic at low temperatures. The observed entropy increment, 0.19 cal mole<sup>-1</sup> °K<sup>-1</sup>, is particularly small and indicates that only a fraction of the uranium atoms are involved in the transition. The magnetic measurements indicate that at least the majority of the uranium atoms are in a triplet state. Assuming that this obtains also for the one-fifth of the uranium atoms symmetrically surrounded by selenium atoms, these would be expected to contribute an entropy increment to the cooperative transition of

$$(1/5) R \ln 3 = 0.41 \text{ cal mole}^{-1} \,^{\circ}\mathrm{K}^{-1}$$

or if the ground state is split into an energetic singlet and a low-lying doublet,

$$(1/5) R \ln 2 = 0.28 \text{ cal mole}^{-1} \, {}^{\circ}\mathrm{K}^{-1}.$$

The latter alternative apparently agrees best with the observations.

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