

## PRODUCTION OF $\gamma$ , $\Lambda^0$ , $K_S^0$ AND $\bar{\Lambda}^0$ IN pp COLLISIONS AT 102 GeV/c<sup>\*</sup>

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We have measured cross sections for  $\gamma$ ,  $K_S^0$ ,  $\Lambda$  and  $\bar{\Lambda}$  production at 102 GeV/c and find:  $\sigma(\gamma) = 170 \pm 16$  mb.,  $\sigma(K_S^0) = 4.6 \pm 0.5$  mb.,  $\sigma(\Lambda) = 3.2 \pm 0.4$  mb., and  $\sigma(\bar{\Lambda}) = 0.23 \pm 0.10$  mb. Both  $\langle n_{\pi^0} \rangle$  and  $\langle n_{K_S^0} \rangle$  appear to rise linearly with  $n_-$  while the ratio  $\langle n_{K_S^0} \rangle / \langle n_{\pi^0} \rangle$  is approximately independent of  $n_-$ . The integrated invariant cross section as a function of  $x$  as well as  $d\sigma/dy$  and  $d\sigma/dp_T^2$  are presented and compared with other data.

Using a 30 000 picture exposure of the 30-inch liquid hydrogen bubble chamber at the National Accelerator Laboratory to 102 GeV/c protons we have measured the inclusive production of  $\gamma$ ,  $K_S^0$ ,  $\Lambda$  and  $\bar{\Lambda}$ . In order to find all events with an associated  $V^0$  ( $K_S^0$ ,  $\Lambda$  or  $\bar{\Lambda}$ ) or  $\gamma$ , two independent scans of the film were made and all conflicts between the two scans were resolved. Within a restricted fiducial volume a total of 505  $V^0/\gamma$ 's were found to be associated with beam track interactions and 488 of these events were successfully measured<sup>†1</sup>. These events were geometrically reconstructed and kinematically fitted using the TVGP-SQUAW program. A requirement that the mass of the  $e^+e^-$  pair be less than 20 MeV/c<sup>2</sup> was used to select  $\gamma$  candidates. The  $K^0/\Lambda$  ( $\bar{\Lambda}$ ) ambiguities were resolved through ionization information when possible or through a selection on the decay angle of the  $\pi^-$  with respect to the line of flight of the  $K_S^0$  in the  $K_S^0$

rest frame<sup>†2</sup>. In addition all neutral particles were restricted to be in the backward hemisphere in the pp c.m. system. After all acceptance criteria were imposed there remained 124  $\gamma$ 's, 105  $K_S^0$ 's, 76  $\Lambda$ 's, and 6  $\bar{\Lambda}$ 's with average weights (inverse detection efficiencies) of 76.6, 2.39, 2.70 and 3.4 respectively<sup>†3</sup>.

In table 1 we list the inclusive cross sections and the average number of particles observed per inelastic pp interaction as a function of charged multiplicity for  $\pi^0$ ,  $K_S^0$  and  $\Lambda$  production. We have assumed that all  $\gamma$ 's come from  $\pi^0$  decay and that  $\sigma(\pi^0) = \frac{1}{2}\sigma(\gamma)$ . These total inclusive cross sections are in general agreement with the trends reported in other high energy pp experiments [1], and lend support to the observation that the  $\Lambda$  production cross section changes very slowly between 69 and 303 GeV/c.

<sup>†2</sup> All events ambiguous between  $K_S^0$  and  $\Lambda$  ( $\bar{\Lambda}$ ) interpretations were taken as  $\Lambda$  ( $\bar{\Lambda}$ ) events if the cosine of the angle between the  $\pi^-$  and the direction of the  $K_S^0$ , measured in the  $K_S^0$  rest frame, was in the interval  $-0.94 \leq \cos \theta \leq -0.86$  ( $0.88 \leq \cos \theta \leq 0.92$ ). This selection introduces essentially no bias into the experimental spectra.

<sup>†3</sup> These weights do not include the additional factor of 2 required to correct for events produced in the forward hemisphere in the pp c.m. but they do include  $V^0$  neutral branching values.

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<sup>†1</sup> All cross sections have been corrected for unmeasurable events.

Table 1  
Cross sections for  $pp \rightarrow \text{neutral} + n \text{ charged} + \text{anything}$

$n$ charged	$\sigma(\pi^0)$ (mb.)	$\langle n_{\pi^0} \rangle$	$\sigma(K_S^0)$ (mb.)	$\langle n_{K_S^0} \rangle$	$\sigma(\Lambda)$ (mb.)	$\langle n_\Lambda \rangle$
2	$7.0 \pm 2.6$	$1.5 \pm 0.6$	$0.36 \pm 0.13$	$0.07 \pm 0.03$	$0.33 \pm 0.13$	$0.07 \pm 0.03$
4	$14.7 \pm 3.8$	$1.9 \pm 0.5$	$0.76 \pm 0.18$	$0.10 \pm 0.02$	$1.15 \pm 0.22$	$0.15 \pm 0.03$
6	$28.0 \pm 5.3$	$3.7 \pm 0.7$	$1.26 \pm 0.23$	$0.17 \pm 0.03$	$0.71 \pm 0.17$	$0.09 \pm 0.03$
8	$14.1 \pm 3.8$	$2.4 \pm 0.6$	$1.02 \pm 0.21$	$0.17 \pm 0.03$	$0.57 \pm 0.17$	$0.09 \pm 0.03$
10	$11.1 \pm 3.3$	$2.9 \pm 0.9$	$0.68 \pm 0.18$	$0.18 \pm 0.04$	$0.41 \pm 0.14$	$0.11 \pm 0.04$
12	$6.2 \pm 2.5$	$3.7 \pm 1.5$	$0.36 \pm 0.13$	$0.21 \pm 0.07$	$0.05 \pm 0.05$	$0.03 \pm 0.03$
14	$1.4 \pm 1.2$	$2.1 \pm 1.8$	$0.04 \pm 0.04$	$0.06 \pm 0.06$	—	—
16	$1.1 \pm 1.1$	$5.0 \pm 5.0$	$0.05 \pm 0.05$	$0.24 \pm 0.24$	—	—
18	$1.5 \pm 1.5$	$27.3 \pm 27.3$	$0.04 \pm 0.04$	$0.77 \pm 0.77$	—	—
Total	$85.0 \pm 8.1$	$2.62 \pm 0.25$	$4.58 \pm 0.46$	$0.141 \pm 0.014$	$3.22 \pm 0.37$	$0.099 \pm 0.012$

In fig. 1 we plot the average number of neutrals observed per inelastic pp interaction as a function of charged topology. The ratio  $\sigma(K_S^0)/\sigma(\pi^0) \sim 0.05$  is approximately independent of the associated charged multiplicity. The approximate linear rise of  $\langle n_{\pi^0} \rangle$  is observed in all experiments at or above 69 GeV/c, in contrast to lower energy pp data [2] where  $\langle n_{\pi^0} \rangle$  is approximately constant as a function of  $n_c$ . The dashed

curve in fig. 1b is given by  $\langle n_{\pi^0} \rangle = n_c$ , a form to which high energy data has been compared. A better parameterization of the data at 102 GeV/c is

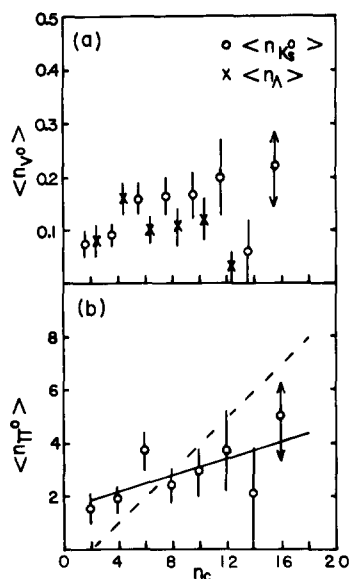


Fig. 1. (a) Average number of  $K_S^0$  (circles) and  $\Lambda$  (crosses) per inelastic pp interaction and (b) average number of  $\pi^0$  per inelastic pp interaction as a function of charged multiplicity. The curves are described in the text.

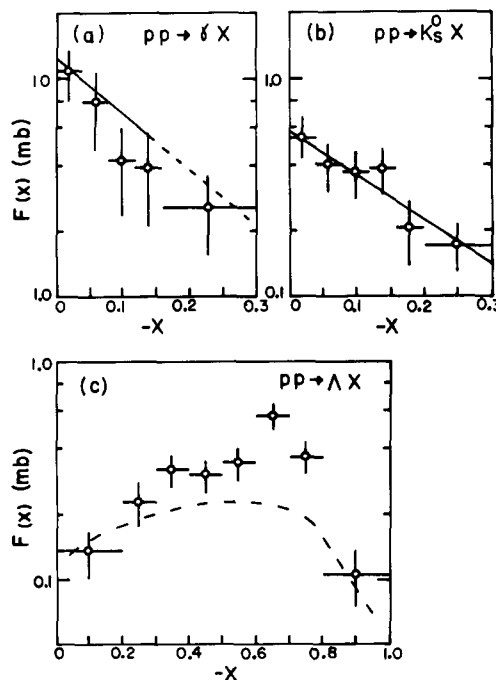


Fig. 2. Invariant cross sections  $F(x) = (2/\pi\sqrt{s}) \int E \times (d^2\sigma/dx dp_T^2) dp_T^2$  as a function of  $x$  for (a)  $pp \rightarrow \gamma X$ , (b)  $pp \rightarrow K_S^0 X$ , and (c)  $pp \rightarrow \Lambda X$ .  $E$ ,  $p_L$ , and  $p_T$  are the energy, the longitudinal momentum, and transverse momentum of the particle in the pp center of mass system and  $x = 2p_L/\sqrt{s}$ . The curves are described in the text.

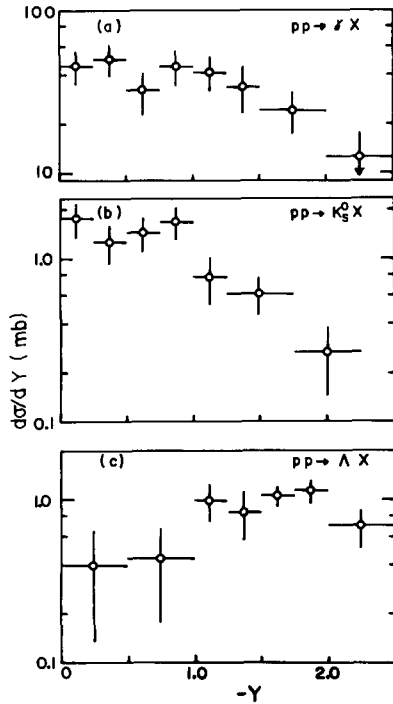


Fig. 3. The cross section  $d\sigma/dy$  versus  $y = (1/2) \ln [(E + p_L)/(E - p_L)]$  for (a)  $pp \rightarrow \gamma X$ , (b)  $pp \rightarrow K_S^0 X$ , and (c)  $pp \rightarrow \Lambda X$ .

$\langle n_{\pi^0} \rangle = (1.8 \pm 0.5) + (0.31 \pm 0.17) n_-$  (solid curve). The rise of  $\langle n_{\pi^0} \rangle$  with  $n_-$  is not in agreement with the predictions of multiperipheral models in which single pions are independently emitted [3]. The total  $\pi^0$  production cross section at 102 GeV/c,  $\sigma(\pi^0) = 85 \pm 8$  mb., is comparable to the total  $\pi^-$  production cross section,  $\sigma(\pi^-) = 66.9 \pm 1.3$  mb.

In fig. 2 we plot the invariant cross section integrated over  $p_T^2$ .

$$F(x) = \frac{2}{\pi\sqrt{s}} \int E \frac{d^2\sigma}{dx dp_T^2} dp_T^2$$

for  $\gamma$ ,  $K_S^0$  and  $\Lambda$  production. The curve in fig. 2a is an integral over  $p_T^2$  of an interpolation formula suggested by Neuhofer et al. [4] as a possible parametrization of  $\gamma$  production data at equivalent lab momenta of 500 GeV/c, 1100 GeV/c, and 1500 GeV/c. The small systematic difference observed in the applicable range of the formula (solid line) may indicate that the invariant cross section for  $\gamma$  production does not scale in this x region.

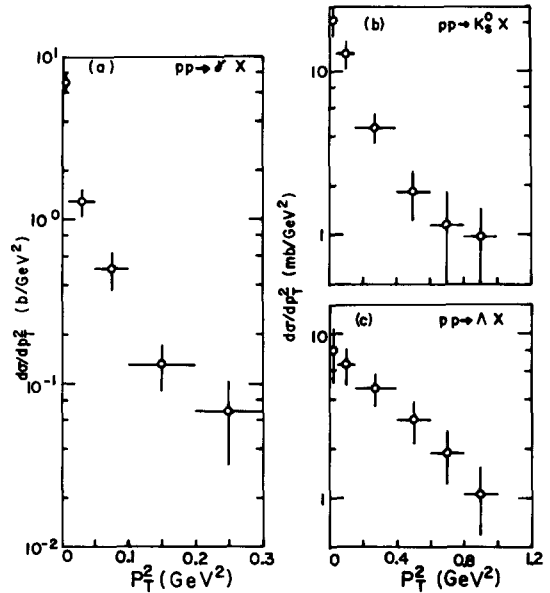


Fig. 4. The cross section  $d\sigma/dp_T^2$  versus  $p_T^2$  for (a)  $pp \rightarrow \gamma X$ , (b)  $pp \rightarrow K_S^0 X$ , and (c)  $pp \rightarrow \Lambda X$ . The  $p_T^2$  distributions have been normalized to account for events from both the forward and the backward hemispheres in the pp c.m. system.

The invariant cross section for  $K_S^0$  production, displayed in fig. 2b, shows an exponential fall off, typical of meson distributions, with a slope of  $4.7 \pm 2.0$  (solid curve). This slope is compatible with that observed at 205 GeV/c and 303 GeV/c. The data on  $F(x)$  for  $\Lambda$  production is similar in all experiments above 69 GeV/c, however, when compared to the 24 GeV/c data of Muck et al. [5] (dashed curve) a rise is seen in the proton fragmentation region ( $-x \gtrsim 0.6$ ).

In fig. 3 we plot  $d\sigma/dy$  as a function of  $y$  (c.m. rapidity) for the above three reactions. Both  $\gamma$  and  $K_S^0$  production are characterized by a plateau whose half

Table 2  
Transverse momenta for particles produced in pp collisions at 102 GeV/c.\*

Particle	$\langle p_T \rangle$ (GeV/c)	$\langle p_T^2 \rangle$ (GeV/c) <sup>2</sup>
$\gamma$	$0.175 \pm 0.020$	$0.050 \pm 0.009$
$K_S^0$	$0.424 \pm 0.043$	$0.246 \pm 0.038$
$\Lambda$	$0.541 \pm 0.060$	$0.364 \pm 0.052$
$\pi^-$	$0.339 \pm 0.010$	$0.171 \pm 0.010$

\* Data are given for  $p_T < 1.5$  GeV/c.

width is approximately one unit in  $y$ . Distributions in transverse momentum are shown in fig. 4 where we plot  $d\sigma/dp_T^2$  as a function of  $p_T^2$ . A typical rapid fall off is observed for all particle production with the steepness being a function of the mass of the produced particle.

In table 2 we summarize the parameters of the  $p_T$  spectra for  $\gamma$ ,  $K_s^0$ ,  $\Lambda$  and  $\pi^-$  production at 102 GeV/c [6].

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