Polarized Raman Spectra of Oriented Films of α-Helical Poly(L-alanine) and its N-Deuterated **Analogue**

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Polarized Raman spectra were obtained on well oriented films of α-helical poly(L-alanine) (α-PLA) and its Ndeuterated analogue in order to obtain definitive assignments of the E2-species modes. Such assignments were aided by an analysis of factors that can lead to 'cross-talk' between A, E_1 and E_2 spectra, viz. reflection of the incident beam and non-perfect lateral and axial chain orientation in the sample. The convincing assignment of at least eleven significant E₂ modes in α-PLA has led to the refinement of a new empirical force field and a more detailed normal mode analysis. © 1998 John Wiley & Sons, Ltd.

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INTRODUCTION

Poly(L-alanine) (PLA) is the simplest L-type polypeptide, with a methyl group as its side chain. It is well known that this molecule can take on a well defined right-handed α -helix or an antiparallel-chain pleated β sheet, depending on the sample preparation method. Much attention has been paid to this molecule as a first step in understanding the structures and functions of polypeptides and proteins. In this effort there have been extensive studies of its infrared 1-8 and Raman 3,9-15 spectra and several normal mode calculations on α-PLA have also been done.4,16-19

In general, the vibrational normal modes of an infinite regular helix are characterized by A, E_1 and E_2 symmetry species, 20,21 where A and E_1 species are infrared active and A, E_1 and E_2 species are Raman active. For a complete vibrational analysis of α-helical PLA, it is essential to separate all three symmetry species experimentally. Separation of A and E_1 species is usually possible through polarized infrared radiation with an oriented sample, and extensive work on α-PLA was done over 40 years ago by Elliott. Since the E_2 -species modes are optically active only in Raman spectra, the separation of the E_2 - from the A- and E_1 -species modes is indispensable for a complete vibrational analysis. About 30 years ago, Fanconi et al.9 obtained the explicit forms of Raman tensors corresponding to the optically active Raman A, E_1 and E_2 modes, and partially succeeded in separating each symmetry species by using polarized laser Raman scattering. However, owing to difficulties in preparing an oriented sample, there has

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been no further advance in polarized Raman studies on α-PLA.

We have been able to obtain good polarized Raman spectra of oriented films of α -PLA and its N-deuterated analogue. Polarized infrared and far-infrared spectra of such (essentially uniaxial) oriented films have also been measured,22 and we have been able to separate convincingly vibrational modes belonging to each of the three symmetry species. Based on these new experimental results, a new empirical force field has been refined and a rigorous normal mode analysis has been done for α-PLA.22

We report here the results of the polarized Raman studies on oriented α-PLA and its N-deuterated analogue (α -PLA-ND). The detailed vibrational analysis will be presented elsewhere.²²

SELECTION RULES FOR A REGULAR HELICAL POLYPEPTIDE

The Raman activities of a molecule can be determined from non-zero projections of the polarizability tensor, α , on each symmetry species, viz. the Raman tensors: if the projection to a particular symmetry species has any non-zero tensor element then a vibrational transition belonging to that symmetry species is Raman active in its fundamental.23 For an infinite regular helical polymer there are in general three Raman tensors:9

$$[\alpha]_{A} = \frac{1}{2} \begin{pmatrix} \alpha_{11} + \alpha_{22} & \alpha_{12} - \alpha_{21} & 0\\ \alpha_{21} - \alpha_{12} & \alpha_{11} + \alpha_{22} & 0\\ 0 & 0 & 2\alpha_{33} \end{pmatrix}$$
(1)
$$[\alpha]_{E_{1}} = \frac{1}{2} \begin{pmatrix} 0 & 0 & \alpha_{13} - i\alpha_{23}\\ 0 & 0 & \alpha_{23} + i\alpha_{13}\\ \alpha_{31} - i\alpha_{32} & \alpha_{32} + i\alpha_{31} & 0 \end{pmatrix}$$
(2)

$$\alpha]_{E_1} = \frac{1}{2} \begin{pmatrix} 0 & 0 & \alpha_{13} - i\alpha_{23} \\ 0 & 0 & \alpha_{23} + i\alpha_{13} \\ \alpha_{31} - i\alpha_{32} & \alpha_{32} + i\alpha_{31} & 0 \end{pmatrix} (2)$$

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$$[\alpha]_{E_2} = \frac{1}{4} \begin{pmatrix} \alpha_{11} - \alpha_{22} - i(\alpha_{12} + \alpha_{21}) \\ \alpha_{12} + \alpha_{21} + i(\alpha_{11} - \alpha_{22}) \\ 0 \end{pmatrix}$$

$$\alpha_{12} + \alpha_{21} + i(\alpha_{11} - \alpha_{22}) \quad 0$$

$$\alpha_{22} - \alpha_{11} + i(\alpha_{12} + \alpha_{21}) \quad 0$$

$$0 \quad 0 \quad 0$$
(3)

where the helix axis is taken as the z-axis corresponding to the Cartesian coordinate index 3 and vibrational modes belonging to the E_1 and E_2 symmetry species are doubly degenerate with $[\alpha]_{E_1^*} = ([\alpha]_{E_1})^*$ and $[\alpha]_{E_2^*} = ([\alpha]_{E_2})^*$, respectively. For a helical structure with an exact threefold rotational symmetry, as in an ideal 3₁₀helix, the E_1 and E_2 modes are degenerate in such a way that $[\alpha]_{E_2} = [\alpha]_{E_1*}$ and they are identical with each other.

In using Eqns (1)–(3) for an oriented helix, we can plausibly assume that the polarizability tensor α is symmetric and also satisfies

$$\alpha_{11} = \alpha_{22} \neq \alpha_{33} \tag{4}$$

since there is no preference between the x-axis and the y-axis. Under these assumptions, the above Raman tensors become

$$[\alpha]_{A} = \begin{pmatrix} \alpha_{11} & 0 & 0 \\ 0 & \alpha_{11} & 0 \\ 0 & 0 & \alpha_{33} \end{pmatrix}$$

$$[\alpha]_{E_{1}} = \frac{1}{2} \begin{pmatrix} 0 & 0 & \alpha_{13} - i\alpha_{23} \\ 0 & 0 & \alpha_{23} + i\alpha_{13} \\ \alpha_{31} - i\alpha_{32} & \alpha_{32} + i\alpha_{31} & 0 \end{pmatrix}$$

$$[\alpha]_{E_{2}} = \frac{1}{2} \begin{pmatrix} -i\alpha_{12} & \alpha_{12} & 0 \\ \alpha_{12} & +i\alpha_{12} & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$(5)$$

$$[\alpha]_{E_2} = \frac{1}{2} \begin{pmatrix} -i\alpha_{12} & \alpha_{12} & 0\\ \alpha_{12} & +i\alpha_{12} & 0\\ 0 & 0 & 0 \end{pmatrix}$$
 (7)

These tensors can be used to separate all the optically active Raman modes of the A, E_1 and E_2 symmetry species for an oriented helical sample with perfect axial and lateral alignment. If we assume the helix axis (sample orientation direction) as the z-axis and the direction in which the scattered light is observed as the

Table 1. Experimental set-ups related to Fig. 1 and corresponding selection of symmetry species from polarized Raman spectra of oriented helical fibers

Experime			
(a)	(b)	Symmetry species	
$Z(xx)Y-\alpha_{11}^{b}$	$X(zz)Y-\alpha_{33}$	A	
$Z(xz)Y-\alpha_{31}$	$X(zx)Y-\alpha_{13}^a$	E ₁	
$Z(yz)Y-\alpha_{32}$	$X(yz)Y-\alpha_{32}$	E ₁	
$Z(yx)Y-\alpha_{12}^{b}$	$X(yx)Y-\alpha_{12}^{b}$	E_2	

a Just for simplicity, this experimental arrangement was used in our E_1 spectral measurements. For the selection of α_{32} component [X(yz)Y arrangement] a half-wave plate is required so as to change the polarization direction of incident light.

^b In uniaxially oriented helical fibers, in which a random orientation exists about the helix axis, the spectrum from the Z(xx)Y arrangement is essentially identical with those from $X(y_x)Y$ or $Z(y_x)Y$ arrangements (see text).

y-axis (Fig. 1), the experimental separation of the A, E_1 and E_2 species can be achieved as given in Table 1. Each experiment is designated in a particular way, e.g. X(yx)Y; X designates the direction in which the incident laser beam propagates and Y denotes the direction in which the scattered light is observed; y denotes the polarization direction of the incident laser beam and x denotes that of the scattered light. Since the induced dipole moment, $\vec{\mu}^{ind}$, due to the incident electric field, $\vec{\mathbf{E}}^{\text{inc.}}$, is given by

$$\vec{\mu}_j^{\text{ind.}} = \sum_{k=1}^3 \alpha_{jk} \vec{E}_k^{\text{inc.}} + \cdots$$
 (8)

X(yx)Y is related to the α_{12} component and selects only the E_2 symmetry species through Eqn (7) by using the experimental set-up (b) (Fig. 1). Similarly, X(zz)Y and X(zx)Y select only A and E_1 symmetry species, respectively. In the experimental set-up (a), it is impossible to place the oriented sample parallel to the incident laser beam direction, and usually the sample has to be tilted slightly, resulting in mixing between the spectra of $Z(xx)Y-\alpha_{11}$ and $Z(yx)Y-\alpha_{12}$. Therefore, the experimental set-up (b) was used in our spectral measurements, and in the foregoing discussion our A, E_1 and E_2

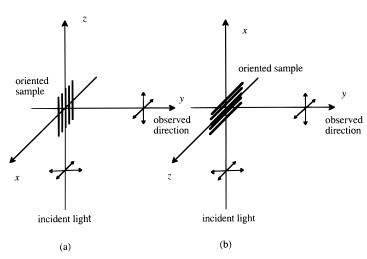


Figure 1. Schematic experimental set-up for polarized Raman spectra. Incident laser beam (a) parallel and (b) perpendicular to the oriented fiber axes.

spectra represent X(zz)Y, X(zx)Y and X(yx)Y arrangements, respectively. Of course, if the helix axis orientation is not perfect, some 'cross-talk' between these spectra will occur (see below).

EXPERIMENTAL

The PLA samples used in this study were purchased from Sigma Chemical [molecular weight based on viscosity determination ≈ 23000 , degree of polymerization (DP) \approx 324)]. When the initially powdered sample is dissolved in dichloroacetic acid (DCA), it is well known² that the resulting PLA is in the right-handed α-form. The concentrated solutions of PLA in DCA were injected into diethyl ether. The resulting precipitated cluster of swollen samples was kneaded on a glass plate by using a stainless-steel micro-spatula until it became fairly sticky. Oriented α-PLA films were made by stroking this sticky sample with a smooth-edged fine metal piece. The kneading process was crucial to achieving a well oriented film by stroking. For the Raman spectra, oriented samples were made on a glass slide and sample films were separated later from the glass. The residual solvent (DCA) was removed from the film by extraction with diethyl ether or pure water followed by drying in a vacuum oven at 40 °C. During the extraction, the sample film usually separated by itself from the glass slide. For Raman spectra the sample was mounted again on a glass slide for easy sample holding.

Polarized infrared spectra²² showed a very high degree of orientation in such a film: for example, in the best samples the 908(A), $892(E_1)$ cm⁻¹ pair¹⁹ were clearly separated, with barely detectable shoulders seen in the unexpected polarization, and similarly for the 1518(A), $1548(E_1)$ cm⁻¹ pair.¹⁹ The degree of orientation was less perfect for the somewhat thicker Raman samples but, as the spectra show, it was sufficiently high to permit a separation of the symmetry species.

The N-deuterated analogue of an oriented α -PLA was made by using O-deuterated DCA. In order to increase the degree of N-deuteration, oriented films made from O-deuterated DCA were soaked in D₂O for about 1 day at 65 °C. The O-deuterated DCA was distilled from the mixed solution of D₂O (95%) and undeuterated DCA (5%). Distillation was repeated with the predominantly O-deuterated DCA a second time, and more than 95% of the DCA was O-deuterated.

Raman spectra were recorded using a Spex 1403 double monochromator with excitation by 514.5 nm radiation from an argon ion laser and 2 cm⁻¹ resolution. Oriented film samples of reasonable thickness were mounted on a small glass slide (thickness ca. 1 mm). Data were obtained with an entrance slit size of 280 µm, an accumulation time of 1 s in each increment of 1 cm⁻¹ and 60 scans at a laser power of 0.16 W for undeuterated samples (50 scans at a laser power of 0.15 W for the N-deuterated sample). Since the PLA sample prepared from DCA can easily be burnt, it is important to defocus the incident laser beam at the sample to an adequate extent. In order to reduce the background fluorescence, the sample had been exposed to a laser beam of 0.07 W for about 7-15 h before the actual data were recorded.

RESULTS AND DISCUSSION

The observed polarized Raman spectra of α -PLA are shown in Fig. 2 and those of α -PLA-ND in Fig. 3. The band wavenumbers, plus clear spectral species assignments, are listed in Table 2. Since the assignment of A and E_1 modes is unambiguous from the polarized infrared spectra, 22 we focus this discussion primarily on the identification of the E_2 modes.

In this connection, a few points need to be considered carefully to interpret the observed spectra correctly. First, a detailed look into the experimental arrangement X(yx)Y (Fig. 4), which was used to record our E_2 spectra and was believed to select out only the α_{12} component, shows that this set-up can indeed additionally select out the α_{11} component corresponding to the X(xx)Y arrangement, which can give rise to A-species bands through Eqn (5) due to the reflected light from the glass slide or from the back sample surface contacting the glass slide. The reflected light acts like another incident beam having some component polarized in the x-direction. The resulting appearance of Abands from the α_{11} component, in addition to the dominant E_2 bands from α_{12} , cannot be completely eliminated even though the glass slide holding the oriented sample is removed. As a result of this effect, a vibrational mode that can give an intense A-species Raman

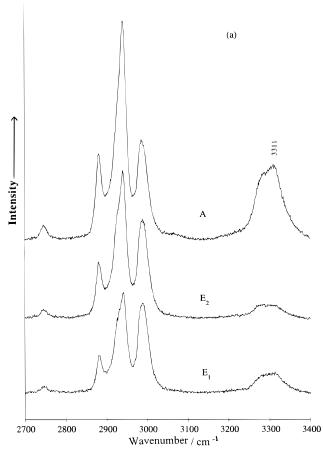


Figure 2. Polarized Raman spectra of oriented α -PLA-NH: A = X(zz)Y, $E_1 = X(zx)Y$ and $E_2 = X(yx)Y$ spectra. Intensity scale is arbitrary. (a) 2700–3400 cm⁻¹ region; (b) 800–1800 cm⁻¹ region; (c) 100–800 cm⁻¹ region; (d) 30–100 cm⁻¹ region.

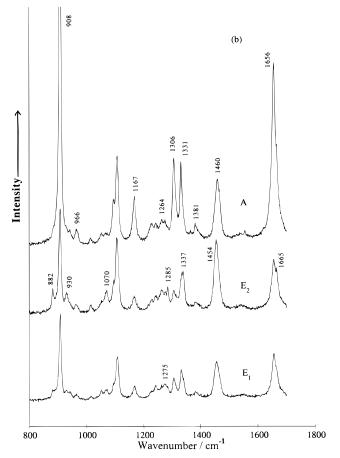


Figure 2. Continued

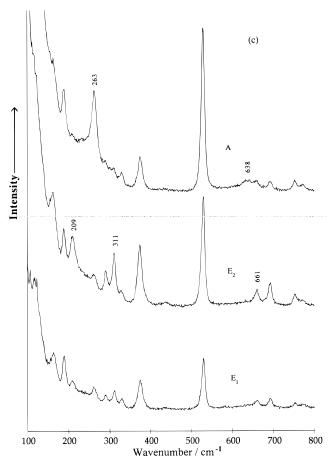


Figure 2. Continued

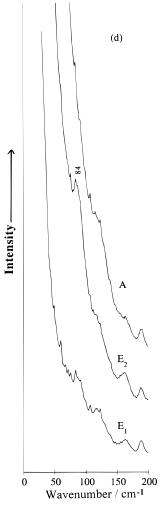


Figure 2. Continued

band from α_{11} can also appear with significant intensity in the E_2 spectra.

A more important problem arises from the fact that all helices do not have the same orientation about the helix axis, i.e. the sample orientation method only gives uniaxial orientation. This means that observationally α_{11} will be equivalent to α_{12} . In fact, when we measured a $Z(xx)Y-\alpha_{11}$ (A) spectrum [using experimental set-up (a) and fixing the sample orientation axis 5-10° from the incident beam], we found it to be the same as the $X(yx)Y-\alpha_{12}$ (E_2) spectrum. Therefore, in contrast to what was previously thought,⁹ in the case of a uniaxially oriented sample, i.e. with random orientation about the helix axis, we cannot independently separate modes due to α_{11} (A-species) from those due to α_{12} (E₂species). However, for a uniaxially oriented sample those bands in the E_2 spectrum that have significantly different wavenumbers from corresponding bands in the $A(X(zz)Y-\alpha_{33})$ and E_1 spectra can be confidently assigned to the E_2 species, while some E_2 bands of the same wavenumber as those in the A spectrum might be due to α_{11} .

Finally, we must consider the effect due to the nonperfect axial orientation in the sample. In order to give a reasonable Raman intensity the sample has to be fairly thick, and it is more difficult to obtain a high

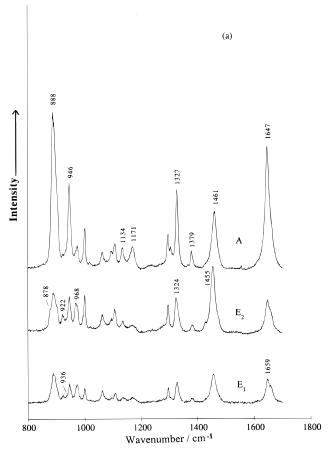


Figure 3. Polarized Raman spectra of oriented α -PLA-ND: A = X(zz)Y, $E_1 = X(zx)Y$ and $E_2 = X(yx)Y$ spectra. Intensity scale is arbitrary. (a) 800–1800 cm⁻¹ region; (b) 100–800 cm⁻¹ region; (c) 30–100 cm⁻¹ region.

degree of axial orientation in such a thick sample. Nevertheless, we still think that the majority of the PLA helix axes are parallel to the stroking direction. Such a small non-perfect orientation in the sample [where Eqn (4) no longer holds] leads to slight additional E_1 (from α_{13} or α_{31}) spectral components in the A spectra, A (from α_{33} or α_{11}) and E_2 (from α_{11} ; see Eqn (3)) spectral components in the E_1 spectra and E_1 (from α_{32}) spectral components in the E_2 spectra. Although the degree of orientation in the sample for Raman scattering is less than that in the sample for infrared, our polarized spectra overall show that the samples prepared by the stroking method are oriented well enough to separate some key bands belonging to each of the three symmetry species.

In the amide A region we observe two NH stretch (s) bands centered around 3311ms(A) and 3288m(A) cm⁻¹, respectively. This doublet nature, which is more prominent in the infrared spectra, turned out to be dependent on the amount of residual solvent DCA in the oriented PLA sample. Detailed analysis of this region will be presented elsewhere, ²⁴ but it shows that the medium strong band around 3311(A) cm⁻¹ arises from the α -helical structure while the band around 3288m(A) cm⁻¹ is due to an additional structure in the PLA sample. Since the CH s modes (3000–2800 cm⁻¹) can be involved in extensive Fermi resonances with CH bend

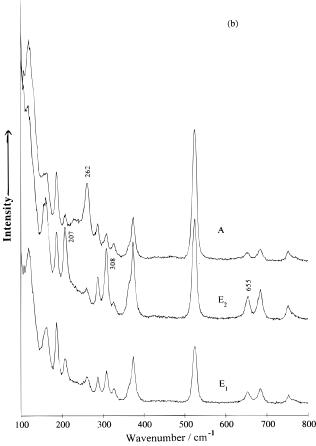


Figure 3. Continued

modes,²⁵ they will not be discussed further here.

In addition to the typical amide I band at 1656vs(A) cm⁻¹ [on N-deuteration at 1647vs(A) cm⁻¹], we observed a clear E_2 spectral band at 1665mw cm⁻¹ (no significant corresponding band could be seen in the polarized infrared spectrum). In view of the expectation that in this spectrum there will also be some A mode contribution from the α_{11} component, the evident presence of the 1665 cm⁻¹ band must be considered compelling, and we newly assign it to the E_2 amide I mode. Its wavenumber is significantly higher than the value of 1642 cm⁻¹ previously predicted from a perturbation transition dipole coupling analysis, ¹⁹ but is consistent with a treatment incorporating direct interaction force constants. ²²

The strong 1454 cm⁻¹ band in the E_2 spectrum, with a definite shoulder at 1460 cm⁻¹ corresponding to the A-species mode, indicates that it should be assigned to an E_2 -species mode. Similarly, the 1337m cm⁻¹ band in the E_2 spectrum is distinctly separable from the 1331s(A) cm⁻¹ band, which is clearly an A-species mode from the infrared spectrum and whose contribution to the E_1 and E_2 spectra we have anticipated. We therefore assign the 1337 cm⁻¹ band to the E_2 species.

The early assignment of the amide III modes by Chen and Lord¹⁴ is well confirmed by the observed bands at 1264mw(A) and 1275mw(E_1) cm⁻¹. The clear presence of an E_2 -spectral band at 1285mw cm⁻¹ fits such an assignment and is consistent with earlier normal mode

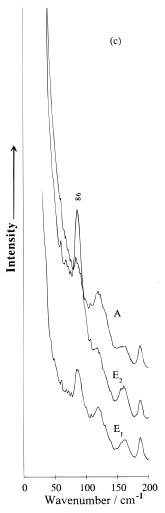


Figure 3. Continued

calculations of its expected position.¹⁹ The corresponding amide III' modes are observed at 946s(A), $973m(E_1)$ and $968m(E_2)$ cm⁻¹ in the spectrum of α -PLA-ND. The lower wavenumber of the E_2 mode compared with the

 E_1 mode in this case is evident from the E_2 spectrum and is well reproduced by the newly refined force field.²²

The bands near 1106s cm⁻¹ in the A, E_1 and E_2 spectra may well be present in all three species, and normal mode calculations indeed show this to be the case.²² The E_2 spectrum, however, unambiguously shows that bands at 1070mw, 930mw and 882mw cm⁻¹ should be assigned exclusively to this species. The last two seem to have counterparts at 922w and 878sh cm⁻¹, respectively, in the E_2 spectrum of α -PLA-ND.

A few bands appear to present a problem in assignment to a species, in that they seem to be too strong in the E_2 spectrum in view of the fact that the normal mode calculations, 22 and in some cases the polarized infrared spectra, place them in the A species. Thus, the 752mw(A) cm⁻¹ band, while it could also be expected to show up from the experimental reflected beam, seems to be unusually strong in the E_2 spectrum. Of course, the significant presence in the E_2 spectrum of strong A-species bands at 1655 and 908 cm⁻¹, where neither is expected from the calculation²² is consistent with this observation and may just indicate that these modes are associated with large α_{11} components. The 693 cm $^{-1}$ band is clearly an A-species mode from the infrared spectrum^{19,22} and its appearance in the E_2 spectrum may be due to a similar effect. Note that all these bands have nearly the same wavenumbers in A and E_2 spectra. The 661mw cm $^{-1}$ band clearly belongs to the E_2 species, as does its counterpart at 655m cm⁻¹ in α-PLA-ND. This new information, together with the shift of the 693 cm⁻¹ band to 685 cm⁻¹ in α -PLA-ND, justifies reassigning these bands to skeletal modes²² rather than the previous assignment of the ca. 660 cm⁻¹ band to an amide V mode. 19 The 638mw(A) cm⁻¹ band, which disappears on N-deuteration and was not previously noted, ¹⁹ will also be seen²² to be assignable to the A-species amide V mode.

The 530 and 374 cm⁻¹ bands could be present in all species. The 311ms cm⁻¹ band clearly belongs to the E_2 species, as do those at 209m and 84m cm⁻¹. The 290

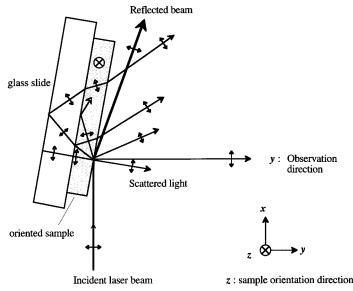


Figure 4. Schematic projection along the sample orientation direction of an experimental set-up for the polarized Raman E_2 spectra. The incident laser beam is polarized in the y-axis.

Table 2. Observed wavenumbers $(\tilde{\nu})$ of polarized Raman bands of α -poly(L-alanine)

	α-PLA-NHª			$\alpha\text{-PLA-ND}^a$	
A	E_1	E_2	Α	E ₁	E_2
\tilde{v}/cm^{-1}	$\tilde{v}/cm-1$	\tilde{v}/cm^{-1}	\tilde{v}/cm^{-1}	\tilde{v}/cm^{-1}	\tilde{v}/cm^{-1}
<i>3311</i> ms					
3288m					
2989s	2989s	2988s			
	2970sh				
2941 vs	2941s	2941s			
2932sh	2928sh	2931sh			
2925sh	2926sh	2924sh			
2882m	2882m	2881m			
2747mw	2747w	2745w			
1665sh	1663sh	<i>1665</i> mw	1662sh	<i>1659</i> ms	1656sh
<i>1656</i> vs	1656ms	1655ms	1647vs	1647ms	1649m
1555vw			1556vw		
1465sh			1472sh	1471sh	1471sh
<i>1460</i> ms	1456ms	1460sh	<i>1461</i> ms	1457s	
	1454sh	<i>1454</i> s	1453sh		<i>1455</i> s
				1444sh	
			1429sh		1429w ^b
<i>1381</i> mw	1381w	1382w	<i>1379</i> mw	1381w	1382mw
	1338sh	<i>1337</i> m			
	1334sh	1335sh			
<i>1331</i> s	1331m		<i>1327</i> s	1327ms	1327sh
				1325sh	<i>1324</i> ms
	1313sh	1313sh			
<i>1306</i> s	1306m	1306mw	1305w°	1305vw°	
			1296m	1296m	1296m
	1284sh	<i>1285</i> mw			
1274mw	<i>1275</i> mw	1276w			
<i>1264</i> mw	1262vw	1262mw			
1242w	1242w	1243w	1239vw		
1229w	1229vw	1229w	1229vw		
<i>1167</i> ms	1168mw	1166mw	<i>1171</i> m	1169w	1170w
				1142sh	
			<i>1134</i> m	1134w	1135w
1106s	1108ms	1106s	1107m°	1107m°	1107mw°
1094w	1094sh	1095sh	1095mw	1094vw	1094w
1070vw	1071w	<i>1070</i> mw			
1051w	1051w	1051sh	1062mw	1062mw	1062mw
1015w	1015vw	1014w	1017vw		1021vw
			1001ms	999mw	1000ms
			975mw	973m	973sh
			970sh	970sh	<i>968</i> m
<i>966</i> mw	964w	962w	<i>946</i> s	946m	946ms
942vw	942vw	222	000.1	<i>936</i> sh	000
927sh	928vw	<i>930</i> mw	922sh	202.15	<i>922</i> w
<i>908</i> vvs	908s	908s	900sh°	900sh°	903sh°
	892sh	222	222	000	000
	882w	<i>882</i> mw	<i>888</i> vvs	888s	889ms
770	774	770	700-1	878sh	<i>878</i> sh
773w	771vw	770w	763sh	764w	763w
752mw	753w	753mw ^b 693mw ^b	751w	753mw	753mw ^b 685m ^b
693mw	693w		685w	685mw 654w	
659w	661w	<i>661</i> mw	653w	034W	<i>655</i> m
<i>638</i> mw	F20	E20-	E24	E2E	E25
530vs	530ms	530s	524vs	525vs	525vs
375m	374m	374ms	374m	374ms	374ms
220	220	220	363sh	365sh	365sh
329mw	329w	328vw	326w	326w	326w
309w 290vw	311mw 289w	<i>311</i> m 290mw ^b	309w 287w	308mw 287mw	<i>308</i> ms 289 m ^b
	289W 260mw				
263ms		261w	<i>262</i> ms	260mw	259mw
208vw 189m	208w 189m	<i>209</i> m 188 m	208w 187m	207mw 187ms	<i>207</i> m 187m
163w	163m	162m	162w	162m	161m
139sh	103111	102111	139sh	102111	101111
119sh	119mw	119sh	1398n 120m	119m	118w
113811					
	86mw	<i>84</i> m	85vw	86ms 60sh	<i>86</i> ms
	60sh	60sh		60sh	60sh

^a s = Strong; m = medium; w = weak; v = very; and sh = shoulder. The intensity comparison between bands is recommended only within one symmetry species. The wavenumber values of the clearly separable (only from Raman spectra) fundamentals are in italics.

Raman spectra) fundamentals are in italics.

^b The intensity of this band seems to arise mostly from the α_{11} component rather than from α_{12} due to the random orientation about the helix axis (see text). After trials, ²² this band was better assigned to the A species rather than to the E_2 species, and in some cases the observed infrared A-species band clearly supported this assignment.

^e Band due to non-perfect N-deuteration.

cm⁻¹ band is another that is clearly an A-species band from the infrared spectrum,^{19,22} so its relative strength in the E_2 spectrum may also be due to a large α_{11} component.

CONCLUSIONS

A complete vibrational analysis of α -PLA, including a convincing normal mode analysis, requires experimental assignments of the observed infrared and Raman bands, i.e. the determination of the symmetry species to which each observed mode belongs. For the general regular helix, the infrared- and Raman-active A and E_1 species are readily assigned from polarized infrared spectra of oriented samples. The E_2 -species modes, however, are optically active only in the Raman spectra. In this study we have improved on previous work and present a more definitive experimental assignment of most of the E_2 -species Raman bands.

The keys to achieving this assignment were twofold: obtaining more highly oriented samples and analysing

the sources of 'cross-talk' between the spectra. In the first case, we believe we have obtained better sample orientation by refining the techniques for stroking out cast films. The resulting polarized laser Raman spectra give evidence of good chain orientation. In the second case, we have carefully examined factors affecting selection rules, such as reflection of the incident beam and non-perfect lateral and axial chain orientation in the sample, that lead to bands of one species appearing in spectra of another. This has permitted us to make secure assignments despite some 'cross-talk' between the spectra.

The new E_2 -species assignments, plus polarized infrared and far-infrared spectra, ²² have permitted the refinement of a new empirical force field for α -PLA and a more complete normal mode analysis of this structure. ²²

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