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Final Report

EVALUATION OF WAVE-RESISTANCE COEFFICIENTS
FOR POLYNOMIAL CENTERPLANE SINGULARITY DISTRIBUTIONS

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ABSTRACT

Wave-resistance coefficients for "thin ships" are obtained from the Michell integral by the application of the Birkoff-Kotik transformation. Mathematical expressions derived for the case of polynomial centerplane singularity distributions are shown to be of rather simple form and definable in terms of integral functions possessing nonsingular integrands.

A computer program for the calculation of wave-resistance coefficients has been prepared and computed results for a simple mathematical hull form are given. These are similar to but differ noticeably from results obtained by means of asymptotic expansions. It is expected that the values of the wave-resistance coefficients reported here are closer to being exact because no approximations are made in the derivation of expressions from which these are calculated beyond the approximations made in the linear wave-resistance theory. Furthermore, all functions involved in the computation are well behaved.

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NOMENCLATURE

$f(x,z)$	Singularity distribution function
$g(x,z)$	Equation of the hull surface
$h(h,w)$	Nondimensional singularity distribution function
u,w	Nondimensional coordinates of center plan of hull
x,y,z	Coordinate system fixed in ship
\bar{x},\bar{y}	Function variables
$A_{\alpha\beta}^I, A_{\alpha\beta}^{II}$	Coefficients of hull function polynomial
B	Half-breadth of ship
$C(x,y)$	Michell function
C_w	Wave-resistance coefficient
$[\Delta^{\alpha\beta} C_w]_I,$	Wave-resistance coefficient related to the α - β term of the hull function
$[\Delta^{\alpha\beta} C_w]_{II}$	
D	Depth of singularity distribution
F	Inverse square of Froude No ($= \frac{gL}{V^2}$)
$H(\xi, \zeta)$	Hull function
L	Length of ship
M,N	Maximum power of the variables of the hull function polynomial
$M_{\alpha\beta}^I(\bar{x},\bar{y}),$	Special functions as defined in text
$M_{\alpha\beta}^{II}(\bar{x},\bar{y})$	

NOMENCLATURE (Concluded)

$P_{2n-1}(\bar{x}),$	Havelock P-function
$P_{2n}(\bar{x})$	
$\bar{P}_{2n-1}(\bar{x}, \bar{y}),$	Generalized Havelock P-function
$\bar{P}_{2n}(\bar{x}, \bar{y})$	
R_w	Wave-resistance, lb
V	Speed of ship, fps
α, β	Non-negative integer powers of the hull function polynomial
$(-\alpha)_k$	Factorial function $(= (-\alpha)(-\alpha+1)\dots(-\alpha+k-1))$
$\delta(\bar{x})$	Dirac delta function
ξ, ζ	Variables of Birkhoff-Kotik transformation
$\Psi_{\alpha\beta}(\bar{x}, \bar{y})$	Special function as defined in text.

INTRODUCTION

The analytical representation of the wave-resistance of a ship published by J. H. Michell¹ in 1898 forms in spite of many recent developments the basis for a major portion of the numerical and theoretical studies being conducted today in the field of wave-resistance theory. The work described in this report is in this respect no exception. What makes the "thin ship" linear theory of Michell particularly attractive is the simplicity of the surface on which we consider the singularities representing the ship's hull to be distributed, this surface being in the case of the Michell theory the centerplane of the ship. The authors believe that the "thin ship" theory provides the greatest flexibility for many studies of fundamental character, such as the problem of minimum wave-resistance, and therefore, that developments related to this theory, will be centered interest if they facilitate computational work, increase accuracy and, perhaps more significantly, provide us with additional analytic properties of Michell's integral.

Within limitations of the linear theory the evaluation of Michell's integral may be expected to give pertinent information about the wave-resistance of "thin ships" of general shape. It is a quintuple integral, however, and its numerical evaluation is a formidable task. Furthermore, the integral is improper, requiring by most methods of evaluation the application of numerical integration techniques. This will introduce numerical approximations which may prove to be of rather serious magnitude. Several methods of evaluation of Michell's integral have been proposed in the past. One of the most noteworthy is due to Weinblum² who developed a set of functions suitable for the case of affine hull forms. Because of this geometric feature his method was severely limited in scope. Other researchers made use of asymptotic expressions for the Michell integral. Inui³ in particular has followed such a procedure.

A fresh approach to the problem of formulating the wave-resistance of thin ships was taken by Birkhoff, Korvin-Kroukovsky, and Kotik⁴ in their joint papers presented before the Society of Naval Architects and Marine Engineers in 1954. Based directly on the Michell integral two transformations were proposed, which would lead to a suitable separation of the functional parameters. The first of these transformations revealed a separation of the functional integrand into two functions, one depending only upon hull form parameters and the other upon ship speed and draft. As a consequence it offered the prospect of providing a better mathematical definition of the relationships between hull form and resistance. Furthermore, it appeared that the Birkhoff-Kotik transformation would lead to simplifications in the numerical work. Especially the idea of being able to tabulate wave-resistance coefficients as functions of hull parameters was attractive at a time when the full impact of the high-speed computers had yet to be felt.

The wave-resistance integral resulting from the Birkhoff-Kotik transformation was left by its proponents in a form which required a double numerical integration over a finite rectangular domain. That operation was complicated by the existence of a singularity located on the boundary of this domain. In an attempt to circumvent the numerical difficulties posed by the singularity, Michelsen⁵ assumed the general singularity distribution to be given by a double finite series. Combining this procedure with certain mathematical developments he was able to formally integrate the wave-resistance integral and write the wave-resistance coefficient for each term of the singularity distribution function in terms of an infinite series. These results led to a contract in 1961 with the Bureau of Ships administered by the D.T.M.B. for a research project with the objectives of making a more detailed study of functional properties, and to prepare a computer program needed for the calculation of wave-resistance coefficients which could subsequently be tabulated. In the following are the results of that project.

Because a relatively long time has elapsed following commencement of the study, it may be appropriate to review its history, the reason being that this in itself offers a lesson to be learned. The equations requiring programming were far from simple and did indeed demand the services of an expert programmer. A number of functions had to be calculated from subroutines not available in the computer library. Some of these functions, such as the Confluent Hypergeometric Functions of the second kind, presented very interesting studies in themselves. After about a year of efforts expanded on the writing of suitable subroutines the total program was ready for trial.

Results indicated complete failure. Wave-resistance was for a representative case given as approximately 10^3 times the expected value, and for a different set of conditions the computer dismayingly recorded it as being negative. A careful analysis of each part of the terms entering into the expression for the wave-resistance calculations revealed that the formulation was absolutely impractical for numerical work. No matter what speed was considered one term or another would become very large, and terms of the alternating series would first increase in magnitude making convergence extremely slow. The main problem was simply computer overflow, however. Some terms reached magnitudes of more than 10^{38} .

Several changes took place at this time. Firstly the programmer left the project for the reason of more profitable employment. Secondly, and on the other side of the ledger, a faster and larger computer was being installed. The project was not of very high priority, however, working from the very beginning on a curtailed budget. Application for additional funding was not approved and research results obtained at this point could hardly warrant that any different action be taken. The nature of the project was such that partial success could not be achieved. Either the objectives were fully realized or nothing at all would be gained from our efforts.

Optimism again prevailed when it was found that the mathematical formulation could be modified to avoid some of the difficulties referred to above. Details of these modifications were described by Michelsen⁶ in a paper read at the International Seminar on Theoretical Wave-Resistance in 1963. The second programmer did now leave the project, however, and with all funds expended it became a matter of the authors completing the research in their spare time.

During the ensuing work it became apparent that the formulas given by Michelsen in 1963 were not fully satisfactory. The expression for the wave-resistance coefficients related to the individual terms of the singularity distribution function had an asymptotic series which appeared to lead to some erratic behavior of the total wave-resistance coefficient of a simple singularity distribution.

To the researcher it is always amazing how painful it is to find the most logical path to the solution of his problem. The mathematical derivation of the next section is a case in point. As far as the authors can judge it avoids all the pitfalls of previous work and leads to a simple concise formula for the wave-resistance. Accuracy can easily be checked because the functions involved possess several recurrence relations. More important may be the fact that computer time needed is relatively short so that it should prove useful in the determination of optimum hull forms.

DERIVATION OF EXPRESSIONS FOR THE WAVE-RESISTANCE COEFFICIENTS

The Michell integral for the wave-resistance coefficient for a "thin ship" can, upon the application of the Birkhoff-Kotik transformation,⁶ be written as

$$C_w = \frac{16F^2}{\pi} \left(\frac{D}{L}\right)^2 \int_0^1 d\xi \int_0^2 d\zeta H(\xi, \zeta) C(F\xi, F\frac{D}{L}\zeta) \quad (1)$$

This coefficient has here been defined by

$$C_w = \frac{R_w}{\frac{1}{2} \rho V^2 B^2}$$

The hull function $H(\xi, \zeta)$ is continuous. It is, however, defined by two separate expressions referred to two subregions as follows:

$$\text{Region I: } [H(\xi, \zeta)]_I = \int_{-\frac{1}{2}}^{\frac{1}{2} - \xi} du \int_0^{\zeta} dw h(u, w) h(\xi + u, \zeta - w)$$

$$0 \leq \xi \leq 1 \quad ; \quad 0 \leq \zeta \leq 1 \quad (2a)$$

$$\text{Region II: } [H(\xi, \zeta)]_{II} = \int_{-\frac{1}{2}}^{\frac{1}{2} - \xi} du \int_{\zeta - 1}^1 dw h(u, w) h(\xi + u, \zeta - w)$$

$$0 \leq \xi \leq 1 \quad ; \quad 0 < \zeta \leq 2 \quad (2b)$$

It should be noted that the hull function in this form does not depend upon the uniform depth D of the singularity distribution. For a nondimensional singularity distribution given by

$$h(u, w) = \sum_{m=0}^M \sum_{n=0}^N C_{mn} u^m w^n \quad (3)$$

integrations indicated in Equations (2a, 2b) are readily carried out and the results can be written

$$[H(\xi, \zeta)]_I = \sum_{\alpha=0}^{2M+1} \sum_{\beta=1}^{2N+1} A_{\alpha\beta}^I \xi^\alpha \zeta^\beta \quad (4a)$$

$$[H(\xi, \zeta)]_{II} = \sum_{\alpha=0}^{2M+1} \sum_{\beta=0}^{2N+1} A_{\alpha\beta}^{II} \xi^\alpha \zeta^\beta \quad (4b)$$

Derivation of the coefficients $A_{\alpha\beta}^I$ and $A_{\alpha\beta}^{II}$ are given in Appendix A.

To the approximation of the Michell "thin ship" theory

$$h(u, w) = \frac{L}{B} \frac{d}{dx} g(x, z) = \frac{L}{B} f(x, z)$$

at corresponding points: $u = \frac{x}{L}$; $w = \frac{z}{D}$, where $g(x, z)$ is the equation of the hull surface. The function $f(x, z)$ may be taken to be the singularity distribution function of sources and sinks located on the longitudinal center plane of the hull.

Equation (1) can be evaluated for each term of Equations (4a) and (4b). By designating the wave-resistance coefficient resulting from the terms $[\xi^\alpha \zeta^\beta]_{I, II}$ as $[\Delta^{\alpha\beta} C_w]_{I, II}$, respectively, Equation (1) thus becomes

$$\begin{aligned} C_w &= \sum_{\alpha=0}^{2M+1} \sum_{\beta=1}^{2N+1} A_{\alpha\beta}^I M_{\alpha\beta}^I \left(\frac{D}{L}, F\right) + \sum_{\alpha=0}^{2M+1} \sum_{\beta=0}^{2N+1} A_{\alpha\beta}^{II} M_{\alpha\beta}^{II} \left(\frac{D}{L}, F\right) \\ &= \sum_{\alpha=0}^{2M+1} \sum_{\beta=1}^{2N+1} [\Delta^{\alpha\beta} C_w]_I + \sum_{\alpha=0}^{2M+1} [\Delta^{\alpha\beta} C_w]_{II} \end{aligned} \quad (5)$$

The summation on β in the first term in Equation (5) above starts at $\beta = 1$ because the hull function vanishes linearly as $\zeta \rightarrow 0$. This is property of the Birkhoff-Kotik transformation.

DERIVATION OF EXPRESSIONS FOR $M_{\alpha\beta}^I$ and $M_{\alpha\beta}^{II}$

In order to show the details of Equation (5) consider any one of the terms of the summation. With the Michell function $C(x,y)$ defined by

$$C(x,z) = \frac{1}{2} e^{-y} \int_0^{\infty} e^{-yt} \cos x \sqrt{t+1} \frac{\sqrt{t+1}}{t^{1/2}} dt \quad (6)$$

we obtain from Equation (1)

$$M_{\alpha\beta}^I \left(\frac{D}{L}, F \right) = \frac{16F^2}{\pi} \left(\frac{D}{L} \right)^2 \int_0^1 \int_0^1 \xi^\alpha \zeta^\beta d\xi d\zeta \\ \therefore \frac{1}{2} \int_0^{\infty} e^{-\frac{FD}{L} \zeta(1+t)} \cos (F\xi\sqrt{t+1}) \frac{\sqrt{t+1}}{t^{1/2}} dt \quad (7)$$

Integration by parts with respect to ζ a number of times equal to β Equation (7) becomes:

$$M_{\alpha\beta}^I \left(\frac{D}{L}, F \right) = \frac{16F^2}{\pi} \left(\frac{D}{L} \right)^2 \int_0^1 \xi^\alpha (d\xi) \frac{1}{2} \int_0^{\infty} [\beta! \left(\frac{L}{FD} \right)^{\beta+1} (1+t)^{-(\beta+1)} \\ - \sum_{p=0}^{\beta} (-\beta)_p (-1)^p \left(\frac{L}{FD} \right)^{p+1} (1+t)^{-(p+1)} e^{-\frac{FD}{L} (1+t)}] \\ \therefore \cos (F \sqrt{t+1}) \frac{\sqrt{t+1}}{t^{1/2}} dt \quad (8)$$

The integral

$$I_1 = \int_0^1 \xi^\alpha d\xi \int_0^{\infty} \cos (F\xi\sqrt{t+1}) \frac{dt}{t^{1/2} (1+t)^{\beta+1/2}} \quad (9)$$

can be expressed in terms of Havelock P-functions. This is readily shown by integrating by parts with respect to ξ a number of times equal to α .

Thus

$$I_1 = 2(-1)^\beta F^{-(\alpha+1)} \left[\left\{ \sum_{k=0}^{\alpha} F^{\alpha-k} (-\alpha)_k P_{2\beta+k}(F) \right\} + \alpha! P_{2\beta+\alpha}(0) \delta(1 + (-1)^\alpha) \right] \quad (10)$$

where $\delta(x)$ is the Dirac delta function. The Havelock P-functions are defined as follows

$$P_{2n-1}(x) = \frac{1}{2} (-1)^n \int_0^{\infty} \cos(x\sqrt{t+1}) \frac{dt}{t^{1/2}(1+t)^{n+1/2}} \quad (11)$$

$$P_{2n}(x) = \frac{1}{2} (-1)^n \int_0^{\infty} \sin(x\sqrt{t+1}) \frac{dt}{t^{1/2}(1+t)^{n+1}} \quad (11b)$$

The integral

$$I_2 = \int_0^1 \xi^\alpha d\xi \int_0^{\infty} e^{-\frac{FD}{L}(1+t)} \cos(F\sqrt{t+1}) \frac{dt}{t^{1/2}(1+t)} \quad (12)$$

can similarly be shown to be given by

$$I_2 = 2(-1)^\beta F^{-(\alpha+1)} \left[\left\{ \sum_{k=0}^{\alpha} F^{\alpha-k} (-\alpha)_k \bar{P}_{2\beta+k}\left(F, \frac{FD}{L}\right) \right\} + \alpha! \bar{P}_{2\beta+\alpha}(0) \delta(1+(-1)^\alpha) \right] \quad (13)$$

where the \bar{P} -functions are given by

$$\begin{aligned} \bar{P}_{2n-1}(x,y) &= \frac{1}{2} (-1)^n \int_0^{\infty} e^{-yt} \cos(x\sqrt{t+1}) \frac{dt}{t^{1/2}(1+t)^{n+1/2}} \\ &= (-1)^n \int_0^{\frac{\pi}{2}} e^{-y \tan^2 \theta} \cos^{2n-1} \theta \cos(x \sec \theta) d\theta \end{aligned} \quad (14a)$$

$$\begin{aligned} \bar{P}_{2n}(x,y) &= \frac{1}{2} (-1)^n \int_0^{\infty} e^{-yt} \sin(x\sqrt{t+1}) \frac{dt}{t^{1/2}(t+1)^{n+1}} \\ &= (-1)^n \int_0^{\frac{\pi}{2}} e^{-y \tan^2 \theta} \cos^{2n} \theta \sin(x \sec \theta) d\theta \end{aligned} \quad (14b)$$

If I_2 is written as $\Psi_{\alpha\beta}(F, \frac{FD}{L})$ we note the integral I_1 will be given by $\Psi_{\alpha\beta}(F, 0)$. With this notation the expression for the wave-resistance coefficient becomes very simple in form. Substitution into Equation (8) gives

$$\begin{aligned} M_{\alpha\beta}^I(\frac{D}{L}, F) &= \frac{\delta}{\pi} (\frac{FD}{L})^{1-\beta} \left[\beta! \Psi_{\alpha\beta}(F, 0) \right. \\ &\quad \left. - e^{-\frac{FD}{L}} \sum_{p=0}^{\beta} (-\beta)_p (-1)^p (\frac{FD}{L})^{\beta-p} \Psi_{\alpha p}(F, \frac{FD}{L}) \right] \end{aligned} \quad (15)$$

From these results and Equation (1) it follows that for the region $0 \leq \xi \leq 1$, $1 \leq \zeta \leq 2$ we can immediately write

$$\begin{aligned} M_{\alpha\beta}^{II}(\frac{D}{L}, F) &= \frac{\delta}{\pi} (\frac{FD}{L})^{1-\beta} \left\{ \sum_{p=0}^{\beta} (-\beta)_p (-1)^p (\frac{FD}{L})^{\beta-p} x \right. \\ &\quad \left. \left[e^{-\frac{FD}{L}} \Psi_{\alpha p}(F, \frac{FD}{L}) - (2)^{\beta-p} e^{-\frac{2FD}{L}} \Psi_{\alpha p}(F, \frac{2FD}{L}) \right] \right\} \end{aligned} \quad (16)$$

Equations (15) and (16) are the complete expressions for the wave-resistance coefficients of a "thin ship" based on the Birkhoff-Kotik transformation of the Michell integral. It should be pointed out that all functions involved are well behaved and easy to compute. Details of the computer programs are given in Appendix B.

NUMERICAL EXAMPLE

To check on calculations throughout the development of our mathematical formulations and the associated computer programs it was decided to consider a simple hull form for which the wave-resistance had been previously computed from the "thin ship" theory. An easy choice was the hull form of parabolic W.L. used by Inui and others. Results for this case are readily available in the literature.³

Choosing $\frac{D}{L} = .10$

and $h(u,w) = -8u$ (17)

we obtain from Equations (2a) and (2b)

Region I

$$A_{01}^I = \frac{16}{3} ; A_{11}^{II} = -16 ; A_{31}^I = \frac{32}{3}$$

Region II

$$A_{01}^I = -\frac{16}{3} ; A_{11}^I = 16 ; A_{31}^I = -\frac{32}{3}$$

$$A_{00}^I = \frac{32}{2} ; A_{10}^{II} = -32 ; A_{30}^{II} = \frac{64}{3}$$

With these values for the coefficients of the hull function polynomial as input, the computer provides us with the wave-resistance coefficient. The graph of this coefficient as a function of Froude number is shown in Figure 1. It is noteworthy that results differ significantly from those obtained by means of Inui's asymptotic approximation.

The contributions made to the wave-resistance coefficient by the individual terms of the hull function are given in Figures 2 through 10. A somewhat peculiar behavior was noted in the coefficients for region II in the range of Froude number $\approx .30$ whenever $\beta = 1$. Accuracy of calculations have been verified several times and it has been concluded that the values of the coefficients are indeed as shown in the figures.

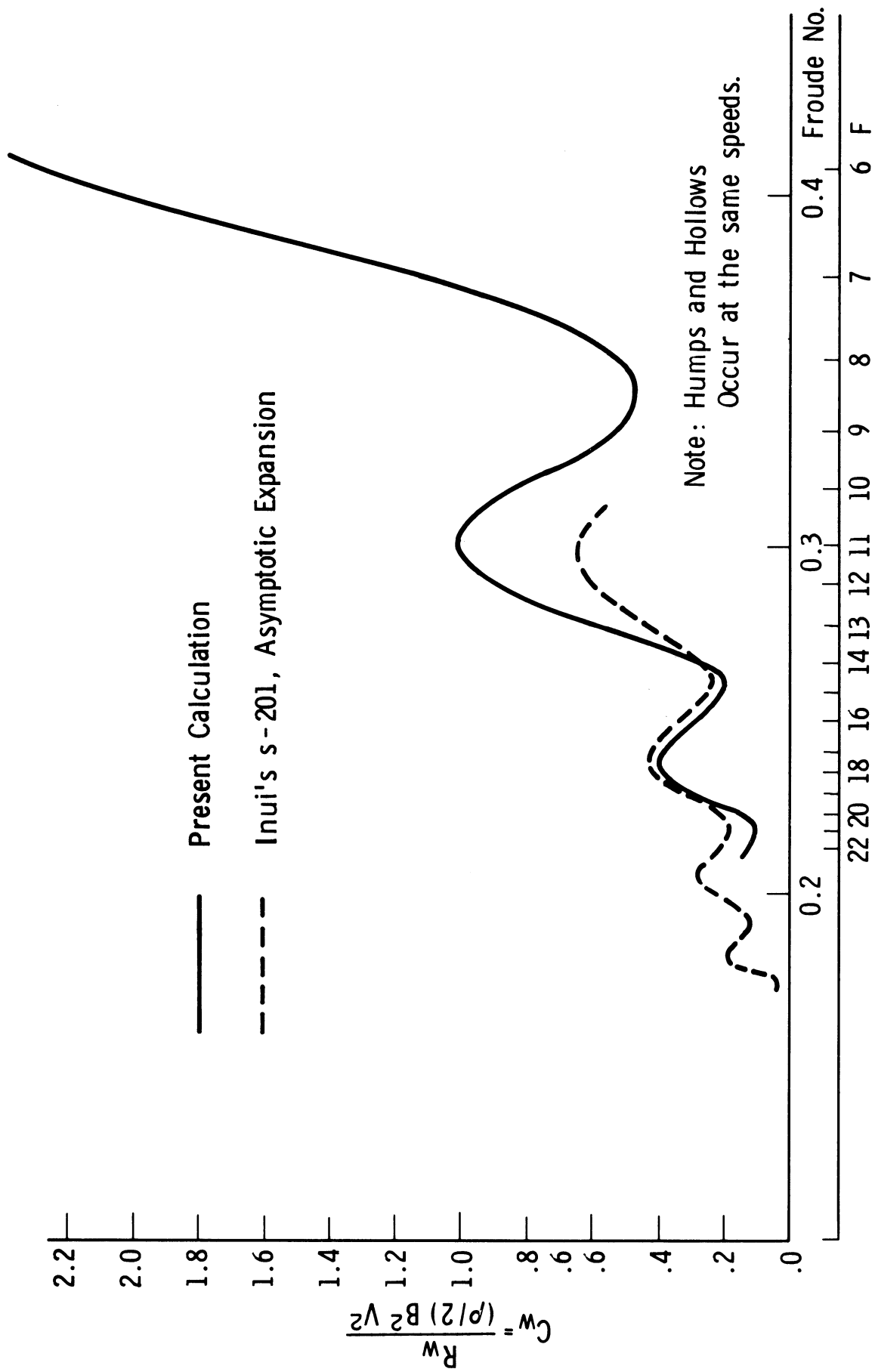


Figure 1. Wave-resistance coefficient for $h(u, w) = -\delta u$, $D/L = .10$

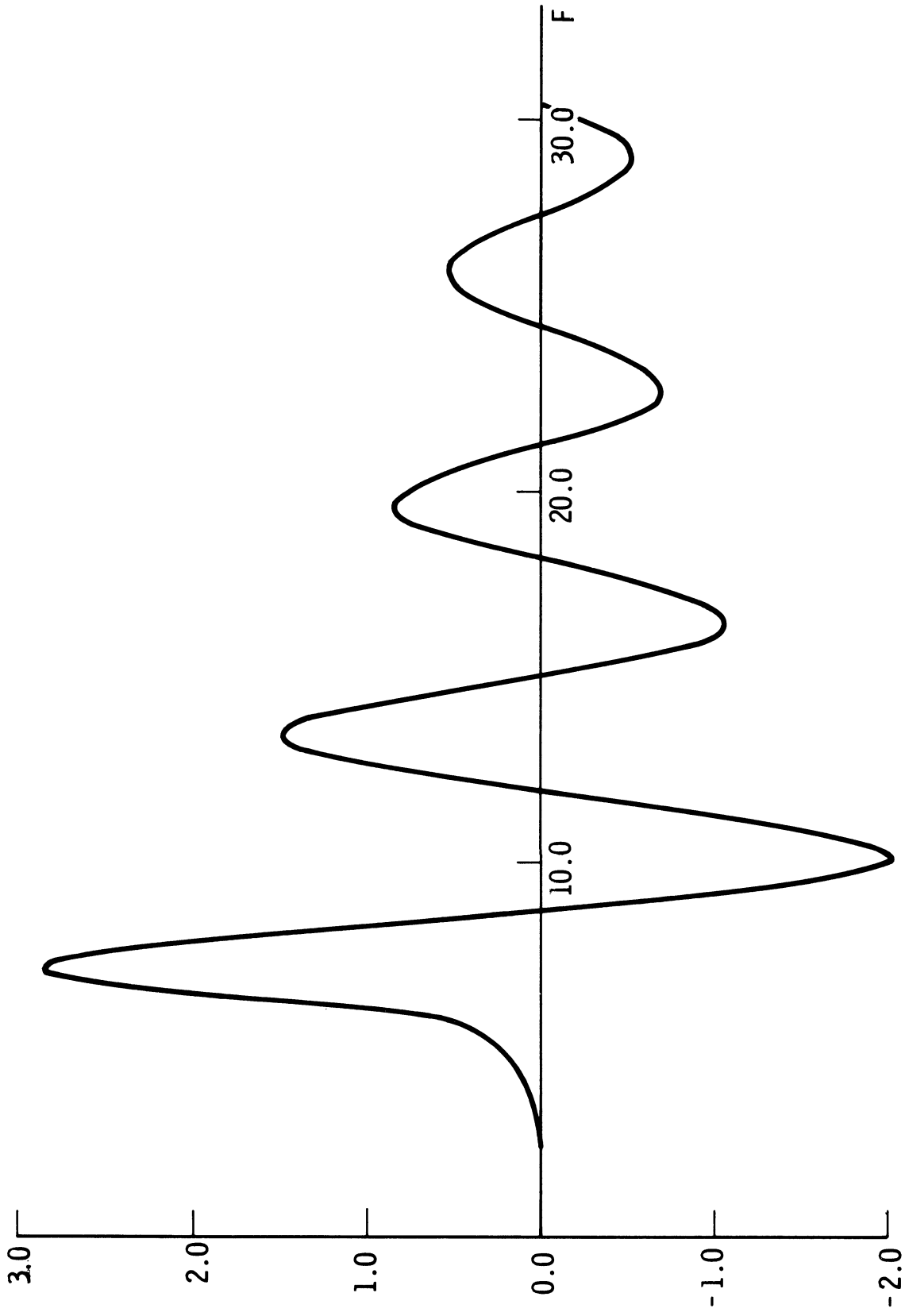


Figure 2. M_{O1}^I x 10 versus F ($\frac{D}{L} = .1$; $h(u,w) = -8u$)

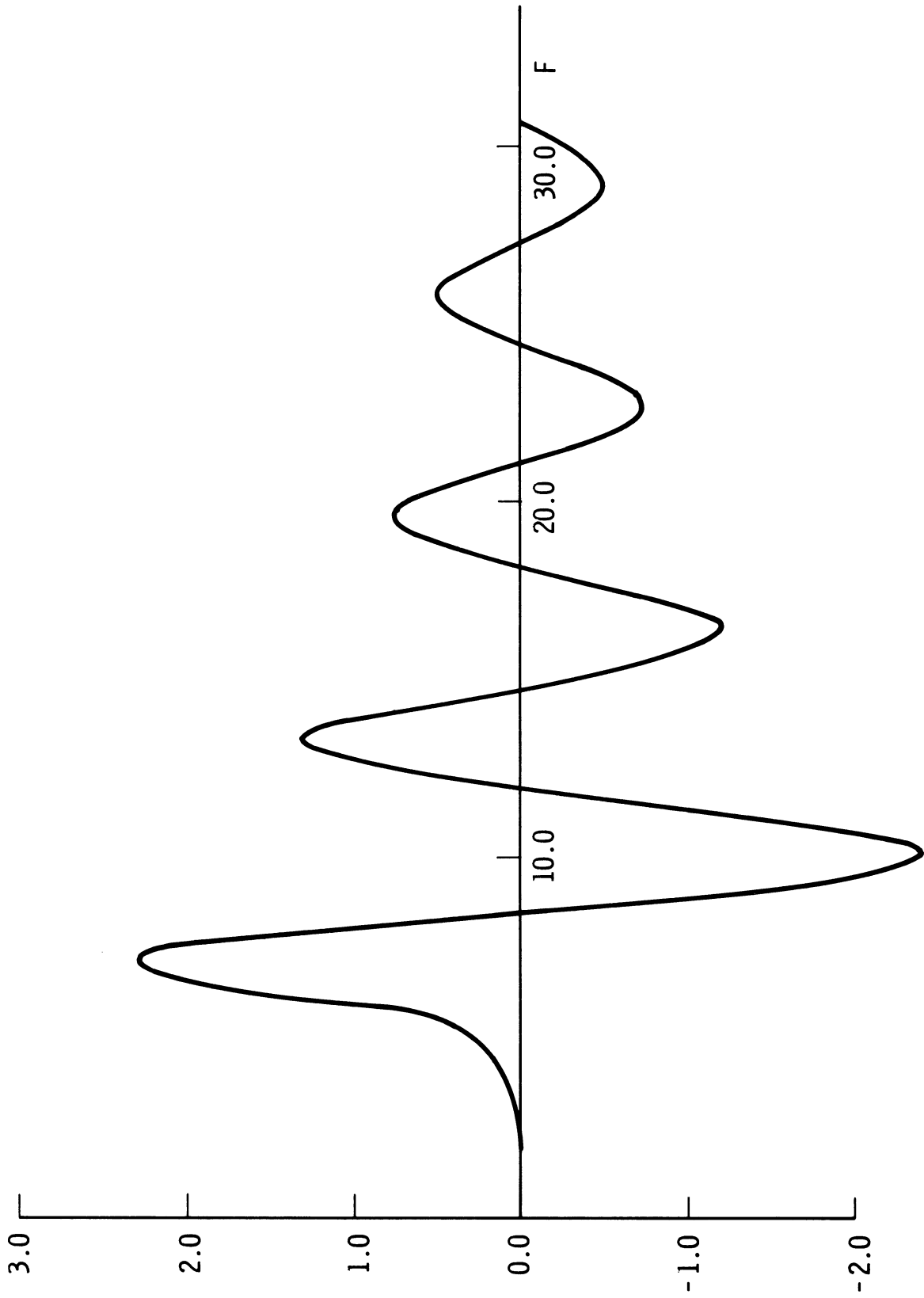


Figure 3. M_{11}^I x 10 versus F ($\frac{D}{I} = .1$; $h(u,w) = -8u$)

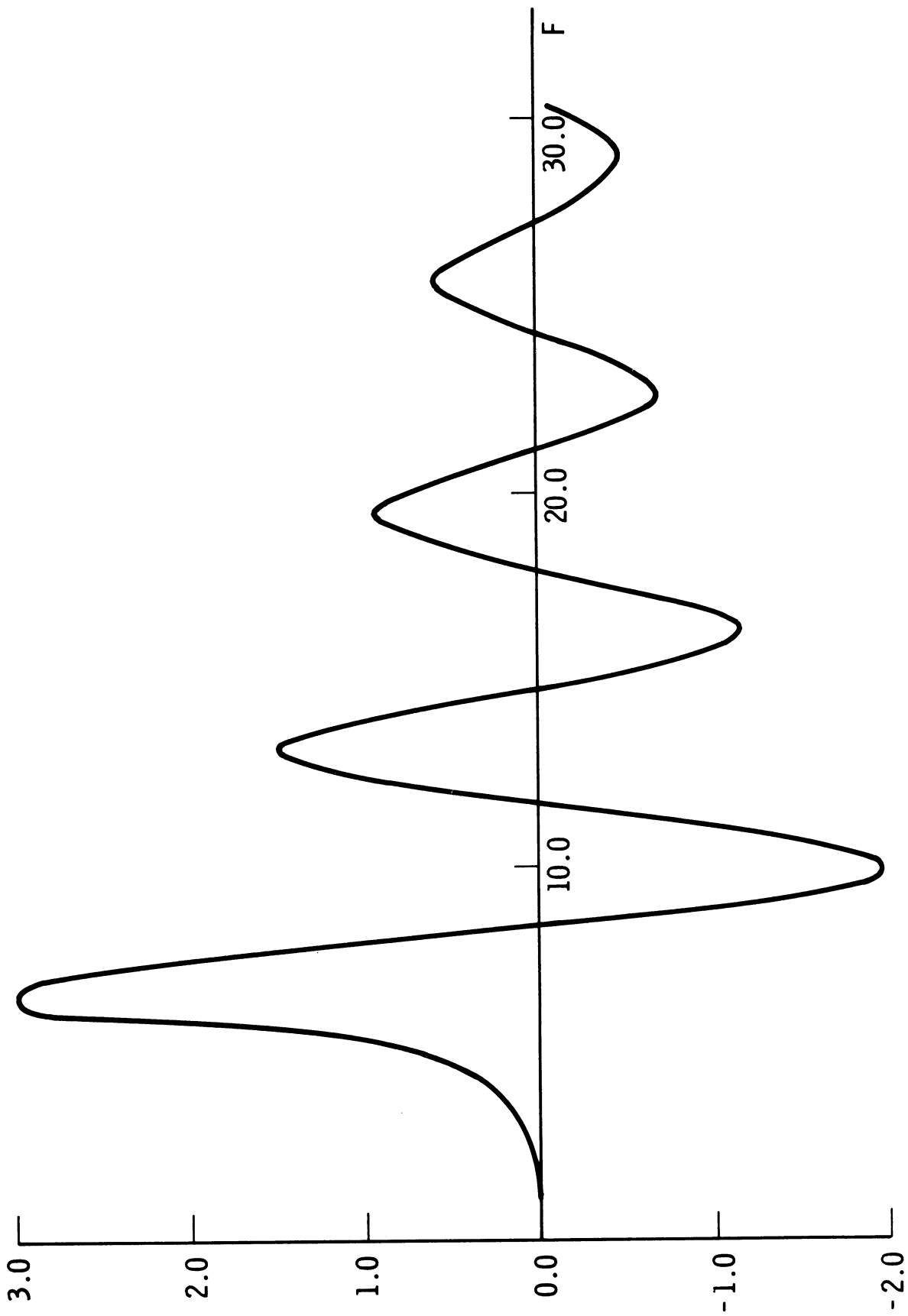


Figure 4. $M_{21}^I \times 10$ versus F ($\frac{D}{L} = .1$; $h(u,w) = -8u$)

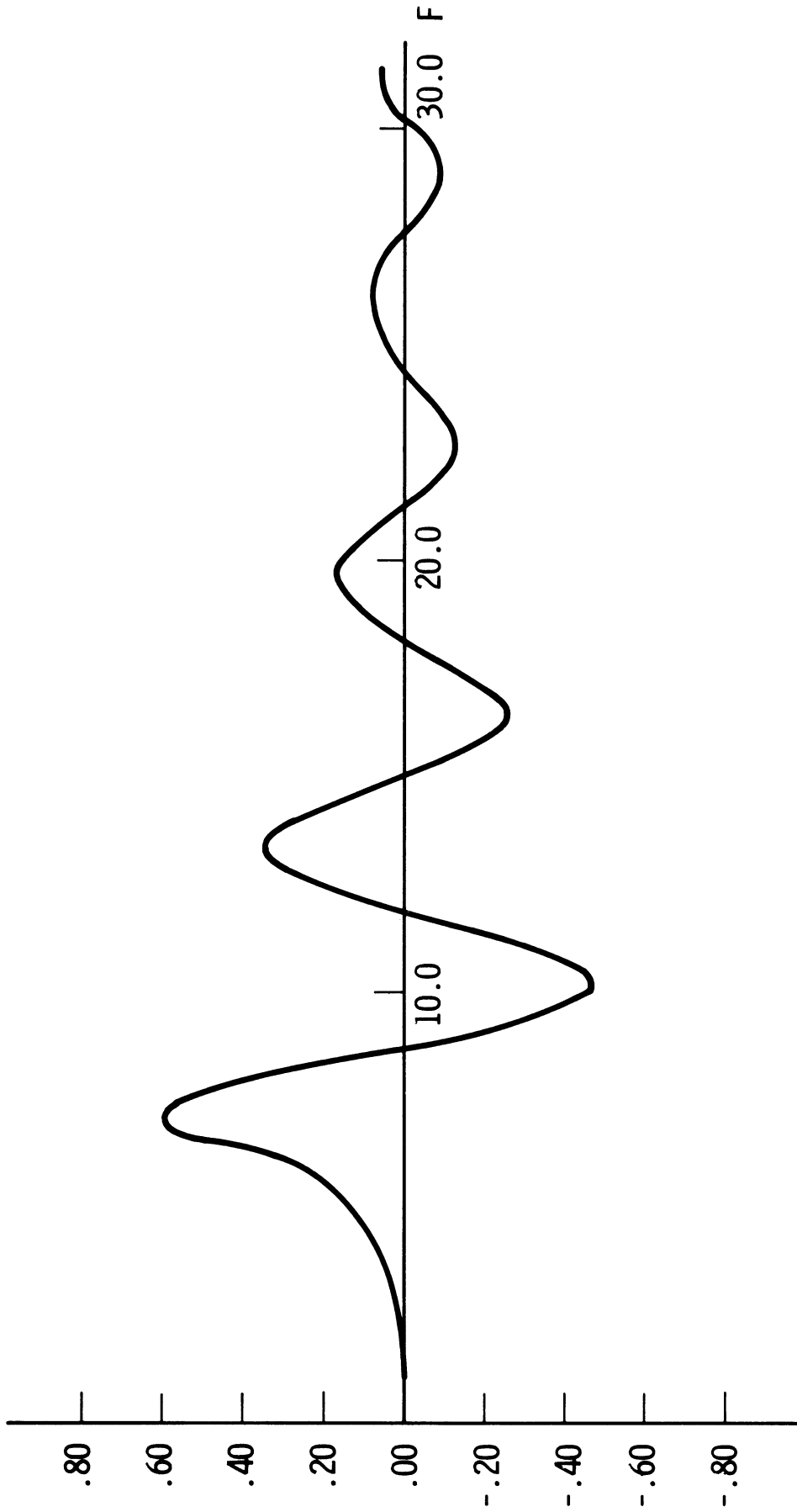


Figure 5. $M_{00}^{II} \times 10$ versus F ($\frac{D}{L} = .1$; $h(u,w) = -\delta u$)

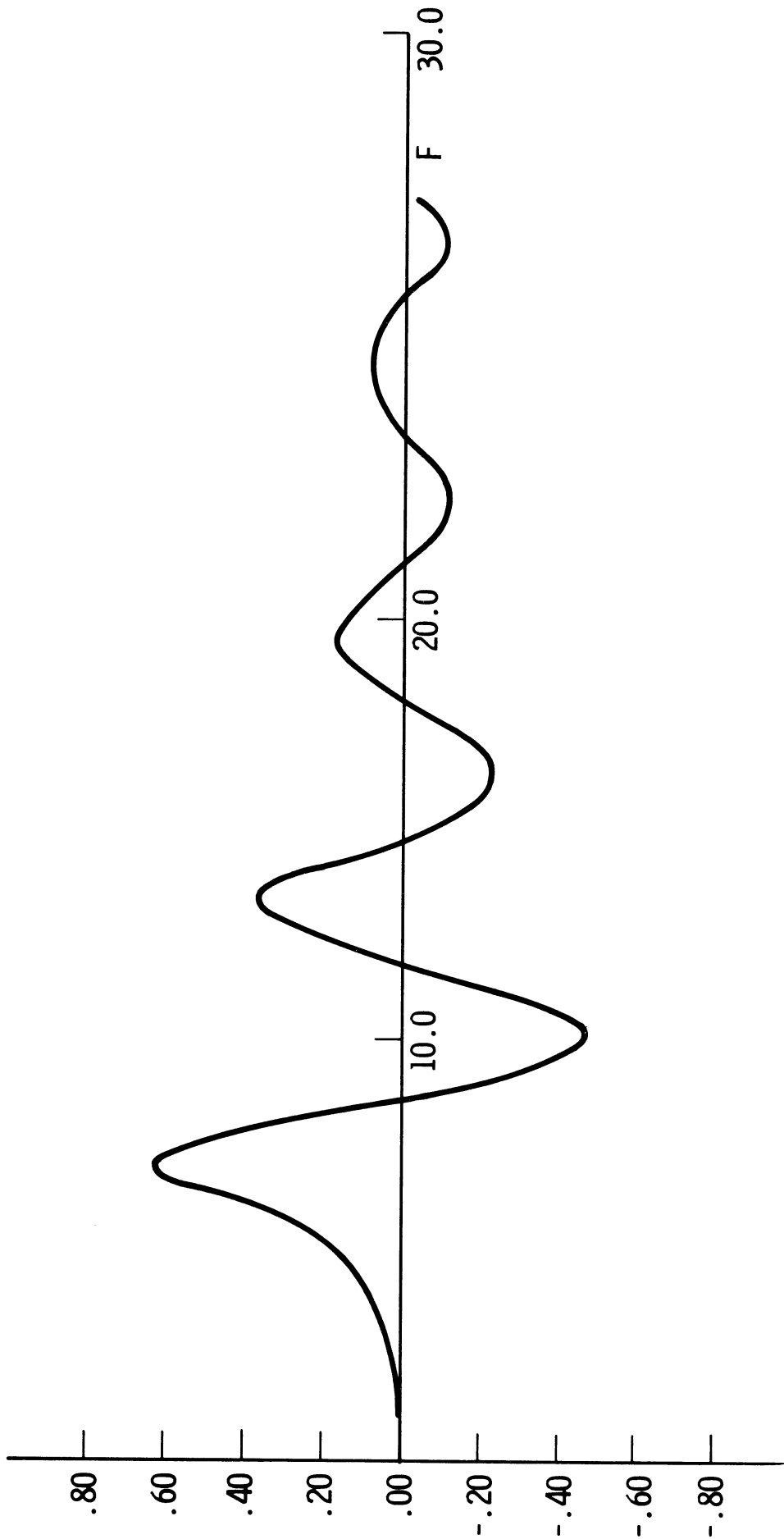


Figure 6. $M_{10}^{II} \times 10$ versus F ($\frac{D}{L} = .1$; $h(u,w) = -\delta u$)

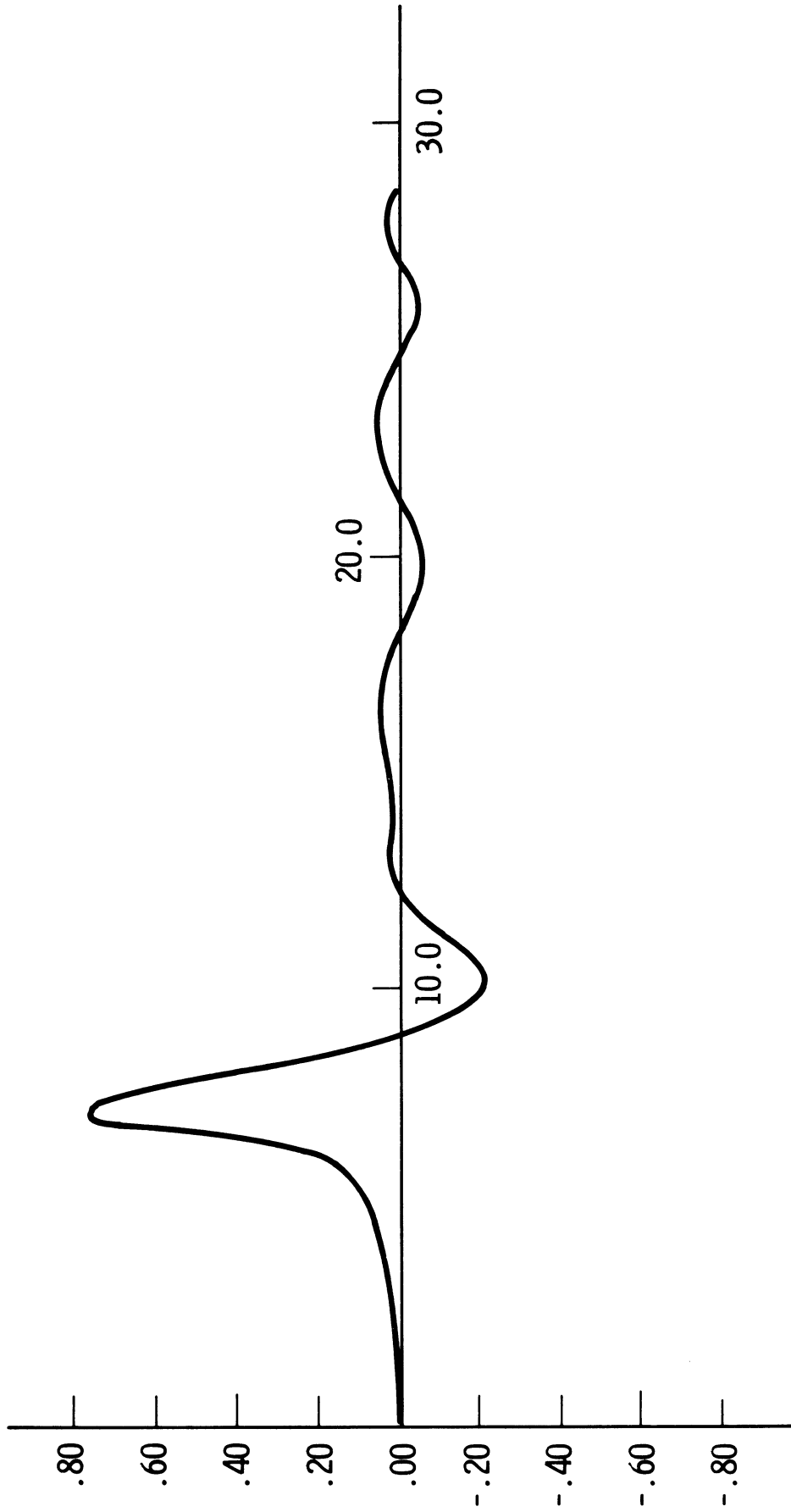


Figure 7. $M_{01}^{II} \times 10$ versus $F \left(\frac{D}{L} = .1 ; h(u, w) = -\delta u \right)$

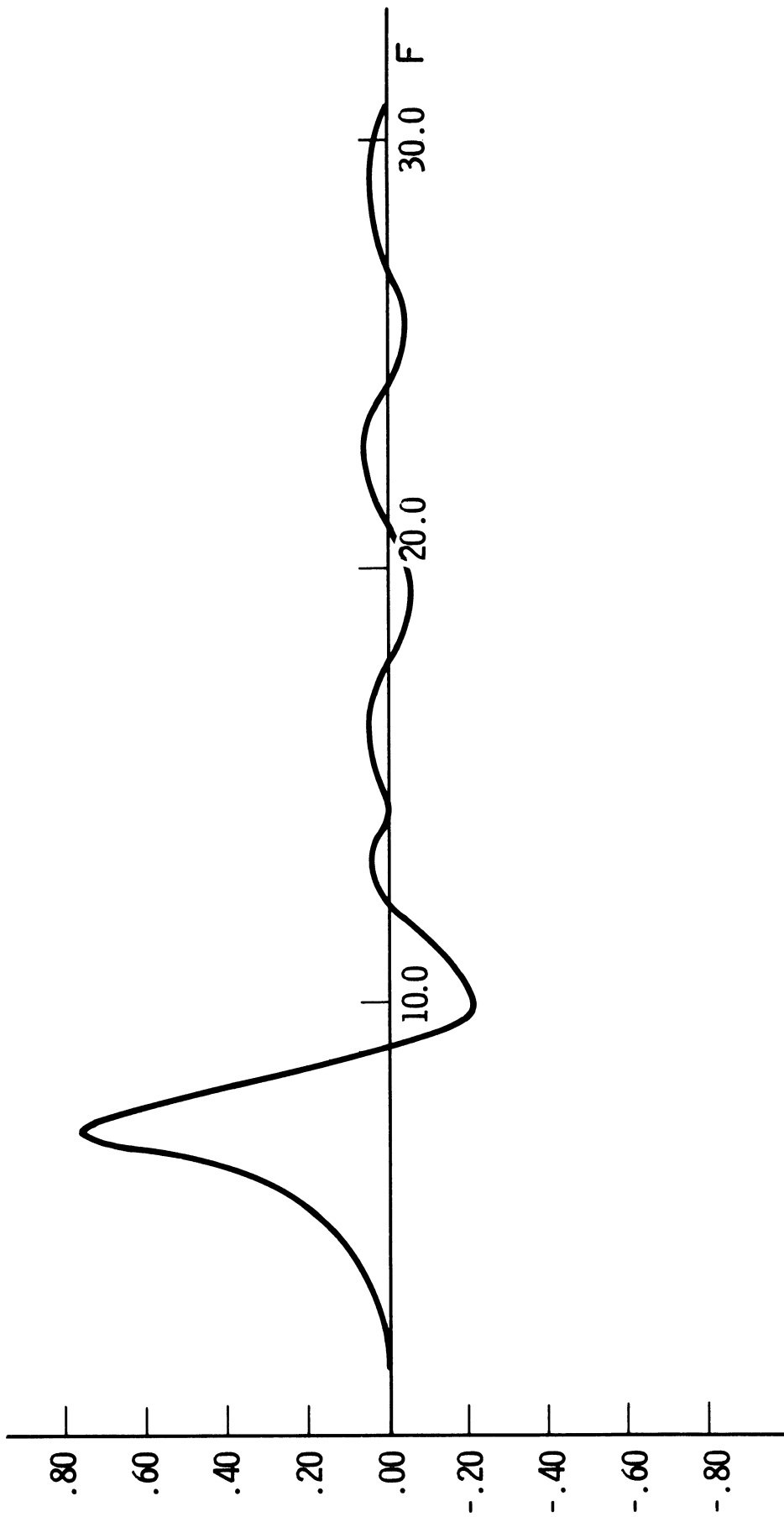


Figure 8. $M_{II}^{II} \times 10$ versus F ($\frac{D}{L} = .1$; $h(u,w) = -8u$)

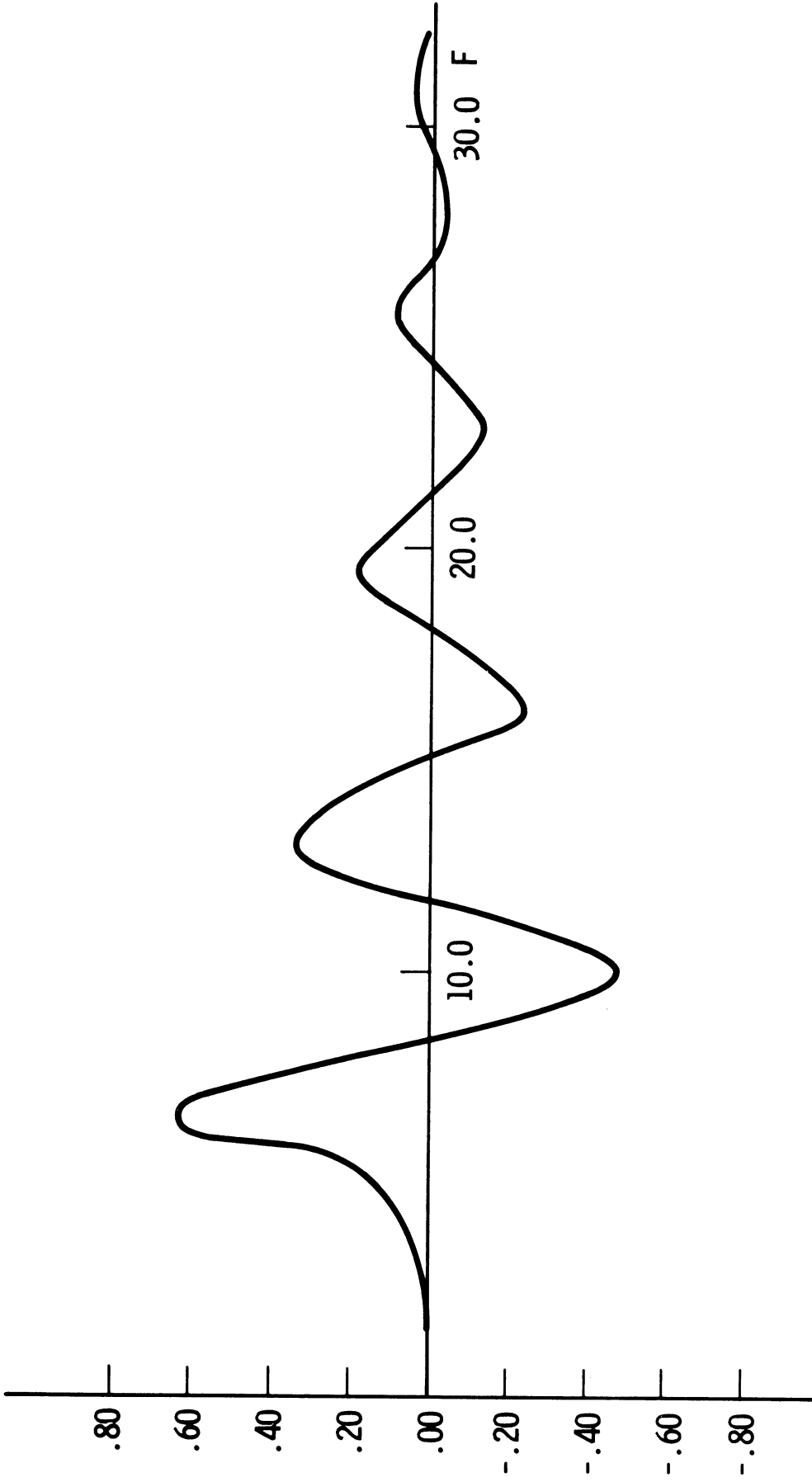


Figure 9. $M_{30}^{II} \times 10$ versus F ($\frac{D}{L} = .1$; $h(u,w) = -8u$)

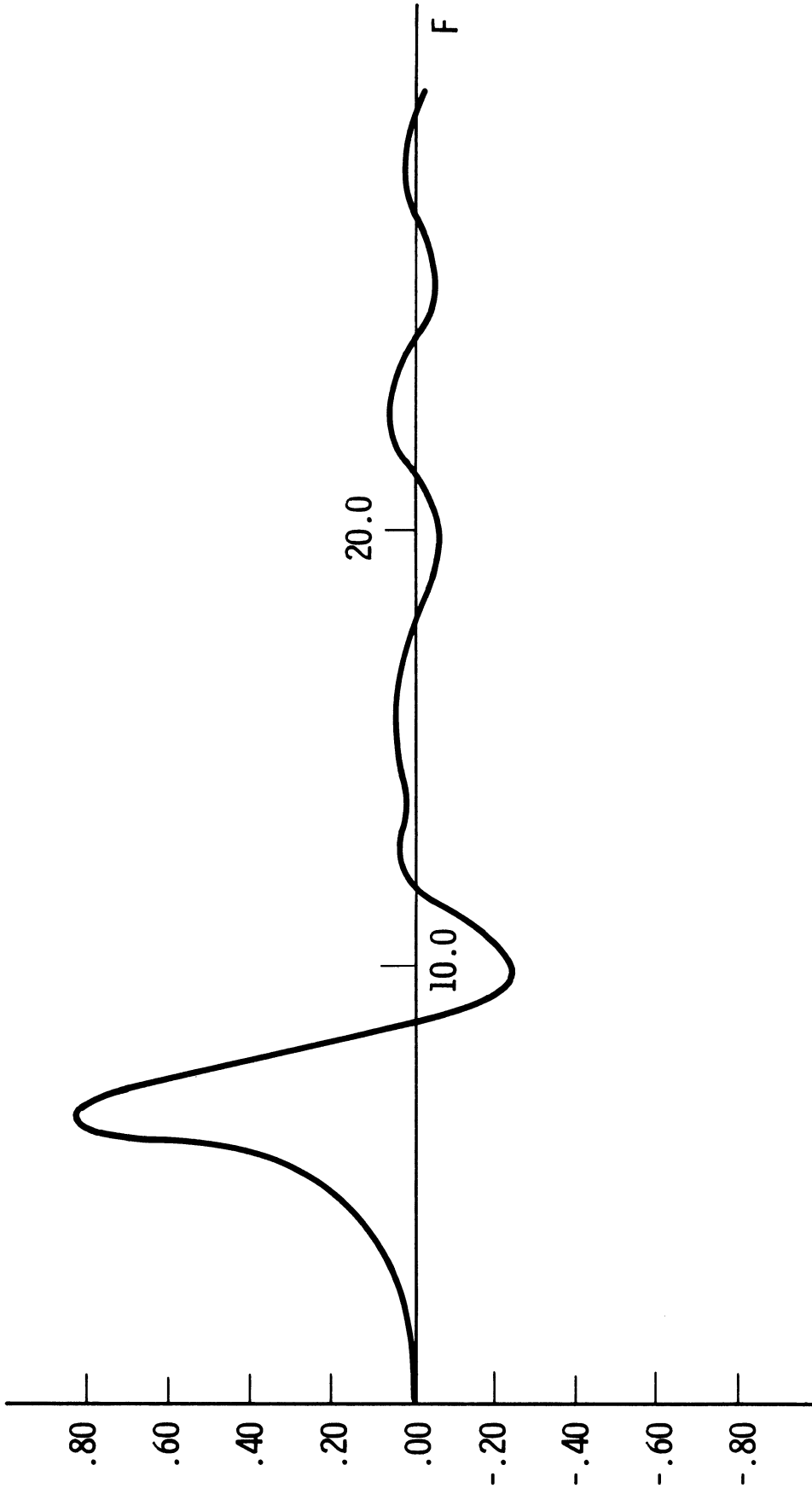


Figure 10. $M_{31}^{II} \times 10$ versus F ($\frac{D}{I} = .1$; $h(u,w) = -8u$)

CONCLUSIONS AND RECOMMENDATIONS

It is believed that the method of evaluation of the Michell integral presented in this report may well prove to be the most general, fastest, and most accurate in existence. Low and high speeds pose no special problems. In fact the middle speed range is where computation time is the longest. Even so it takes only a few seconds to calculate the wave-resistance for the sample hull form considered in the report. At this time we have not taken full advantage of all recurrence relationships of the generalized Havelock \bar{P} -functions and other features which can further reduce computing time. It is therefore recommended that the computer program now available be modified to take advantage of all possible short cuts that can be incorporated.

Of equal importance to the prepared computer programs in the mathematical formulation of the wave-resistance coefficients on which they are based. It is noted that the finite sums of the expression for C_w form a linear system in term of the coefficients $A_{\alpha\beta}^I$ and $A_{\alpha\beta}^{II}$ of the hull function polynomial. These functions are in turn linearly dependent upon the coefficients C_{mn} of the singularity distribution function. The formulation of a minimum wave-resistance problem with suitable constraints is therefore fairly straightforward and should be pursued as soon as possible.

Another extension to the present work worthy of consideration would be the inclusion in the theory of concentrated sources and sinks, as well as singularities of higher orders. Such developments would be necessary to make the formulation of the wave-resistance of the "thin ship" as described in this report as flexible and complete as possible.

ACKNOWLEDGMENTS

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APPENDIX A

COEFFICIENTS OF THE HULL FUNCTION POLYNOMIALS

From Equations (2), (3), and (4) it follows that the hull function can be written

$$\begin{aligned}
 [H(\xi, \zeta)]_I &= \sum_{\alpha=0}^{2M+1} \sum_{\beta=0}^{2N+1} A_{\alpha\beta}^I \xi^\alpha \zeta^\beta \quad (19) \\
 &= \int_{-\frac{1}{2}}^{\frac{1}{2}-\xi} du \int_0^1 dw \left[\sum_{i=0}^M \sum_{j=0}^N C_{ij} u^i w^j \right] \\
 &= \left[\sum_{m=0}^M \sum_{n=0}^N C_{m,n} (u+\xi)^m (\zeta-w)^n \right] \\
 &= \sum_i^M \sum_j^N \sum_m^M \sum_n^N C_{ij} C_{mn} \int_{-\frac{1}{2}}^{\frac{1}{2}-\xi} u^i (u+\xi)^m du \int_0^1 w (\zeta-w)^n dw
 \end{aligned}$$

Considering the integral

$$[I_1]_I = \int_{-\frac{1}{2}}^{\frac{1}{2}-\xi} u^i (u+\xi)^m du \quad (20)$$

By the binomial theorem we have that

$$\begin{aligned}
 [I_1]_I &= \sum_{k=0}^m \binom{m}{k} \xi^{m-k} \int_{-\frac{1}{2}}^{\frac{1}{2}-\xi} u^{i+k} du \\
 &= \sum_{k=0}^m \binom{m}{k} \xi^{m-k} (i+k+1)^{-1} u^{i+k+1} \Big|_{-\frac{1}{2}}^{\frac{1}{2}-\xi}
 \end{aligned}$$

$$\begin{aligned}
&= \sum_{k=0}^m (-1)^{i+k+1} \binom{m}{k} (i+k+1)^{-1} \xi^{m-k} \\
&\left[\sum_{\ell=0}^{i+k+1} \binom{i+k+1}{\ell} (-1)^\ell \left(\frac{1}{2}\right)^\ell \xi^{i+k+1-\ell} - \left(\frac{1}{2}\right)^{i+k+1} \right] \quad (21)
\end{aligned}$$

where $\binom{m}{k} = \frac{m!}{k! (m-k)!}$

Similarly we find that

$$\begin{aligned}
[I_2]_I &= \int_0^1 w^j (\zeta-w)^n dw \\
&= \sum_{p=0}^n \binom{n}{p} \zeta^{n-p} (-1)^p (j+p+1)^{-1} \quad (22)
\end{aligned}$$

Substitution of (21) and (22) into (19) then gives

$$\begin{aligned}
H(\xi, \zeta)]_I &= \sum_i^M \sum_j^N \sum_m^M \sum_n^N \sum_{k=0}^m \sum_{p=0}^n (-1)^{i+k+1+p} \binom{m}{k} \binom{n}{p} \\
&\frac{C_{ij} C_{mn}}{(i+k+1)(j+p+1)} \xi^{m-k} \zeta^{n-p} \left[\sum_{\ell=0}^{i+k+1} \binom{i+k+1}{\ell} (-1)^\ell \left(\frac{1}{2}\right)^\ell \right. \\
&\left. \xi^{i+k+1-\ell} - \left(\frac{1}{2}\right)^{i+k+1} \right] \\
&= \sum_i^M \sum_j^N \sum_m^M \sum_n^N \sum_k^m \sum_{p'}^n \sum_{\ell}^{i+k+1} (-1)^{i+k+p+\ell+1} \frac{\left(\frac{1}{2}\right)^\ell}{(i+k+1)(j+p+1)}
\end{aligned}$$

$$\begin{aligned}
& \binom{m}{k} \binom{n}{p} \binom{i+k+1}{l} \xi^{m+i+1-l} \zeta^{n-p} C_{ij} C_{mn} \\
& + \sum_i \sum_j \sum_m \sum_n \sum_k \sum_p (-1)^{i+k+p} \frac{\left(\frac{1}{2}\right)^{i+k+1}}{(i+k+1)(j+p+1)}
\end{aligned} \tag{23}$$

$$\binom{m}{k} \binom{n}{p} \xi^{m-k} \zeta^{n-p} C_{ij} C_{mn}$$

Following an identical procedure it can be shown that

$$\begin{aligned}
[H(\xi, \zeta)]_{II} &= \sum_{i,j,m,n,k,p,l} C_{ij} C_{mn} (-1)^{i+k+1+p+l} \binom{m}{k} \binom{n}{p} \\
& \binom{i+k+1}{l} \left(\frac{1}{2}\right)^l \frac{1}{(i+k+1)(j+p+1)} \xi^{m+i+1-l} \zeta^{n-p} \\
& + \sum_{i,j,m,n,k,p,l} \sum_{q=0}^{j+p+1} C_{ij} C_{mn} (-1)^{i+k+p+q+l} \binom{m}{k} \binom{n}{p} \binom{i+k+1}{l} \\
& \binom{j+p+1}{q} \left(\frac{1}{2}\right)^l \frac{1}{(i+k+1)(j+p+1)} \xi^{m+i+1-l} \zeta^{n+j+1-q} \\
& + \sum_{i,j,m,n,k,p} C_{ij} C_{mn} (-1)^{i+k+p} \binom{m}{k} \binom{n}{p} \left(\frac{1}{2}\right)^{i+k+1} \frac{1}{(i+k+1)(j+p+1)} \\
& \xi^{m-k} \zeta^{n-p}
\end{aligned}$$

$$\begin{aligned}
& + \sum_{i,j,m,n,k,p,q} C_{ij} C_{mn} (-1)^{i+k+p+q+1} \binom{m}{k} \binom{n}{p} \binom{j+p+1}{q} \\
& \left(\frac{1}{2}\right)^{i+k+1} \frac{1}{(i+k+1)(j+p+1)} \xi^{m-k} \zeta^{n+j+1-q} \quad (24)
\end{aligned}$$

The method used in determining the coefficients $A_{\alpha\beta}^I$ and $A_{\alpha\beta}^{II}$ is to calculate and collect the coefficients of equal powers in ξ and ζ . A computer program written in MAD language that will perform this task is shown in Appendix "B."

APPENDIX B

COMPUTER PROGRAMS

Because programmers at The University of Michigan are generally more familiar with MAD, this language has been used throughout. The programs can readily be translated into other languages, however.

WAVE-RESISTANCE COEFFICIENT

Wave-resistance is calculated from Equations (5), (15), (16) and from the coefficients of Equations (23) and (24). The flow diagram is shown in Figure 11 below followed by the associated program.

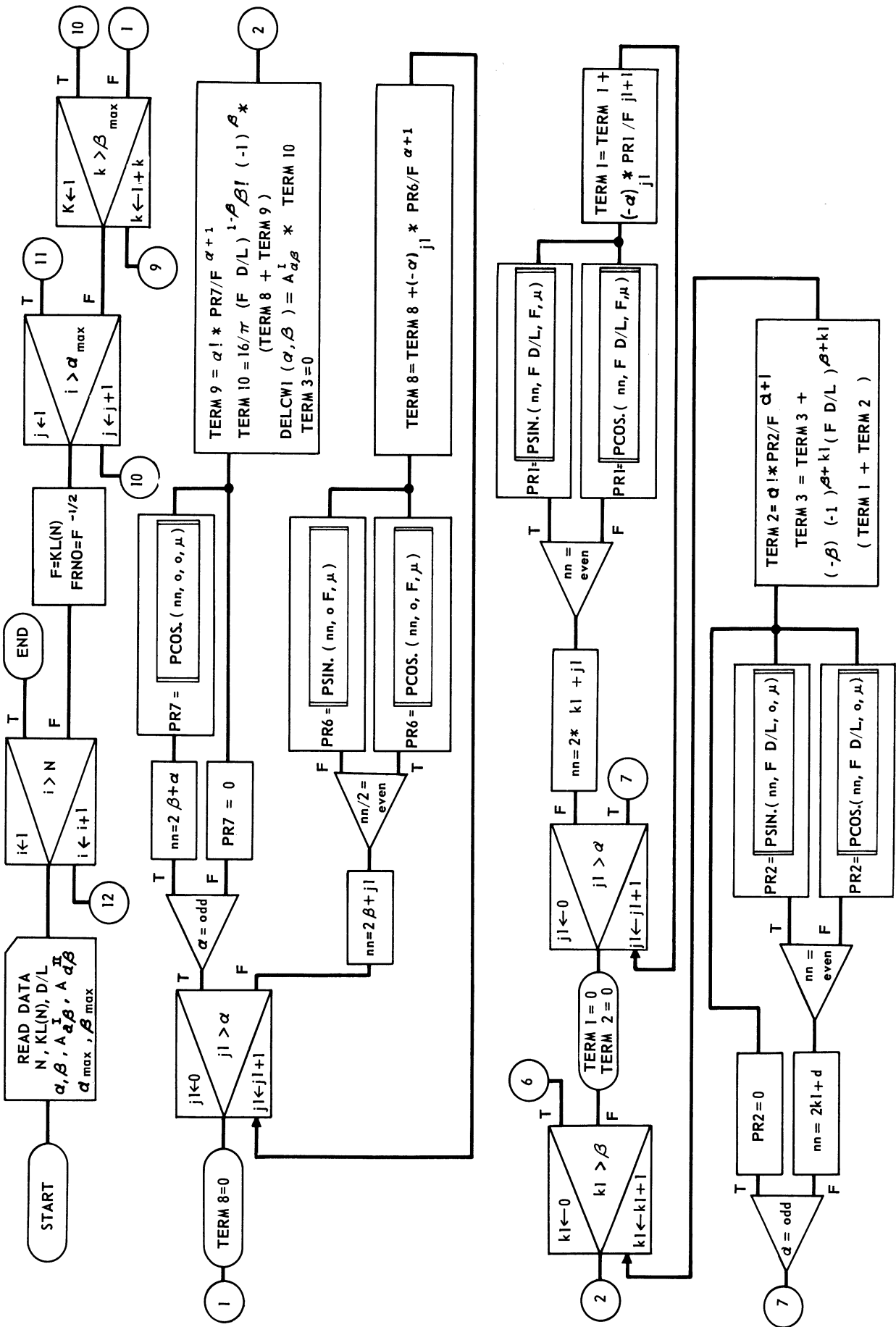


Figure 11. Flow Diagram for Wave-Resistance Calculation

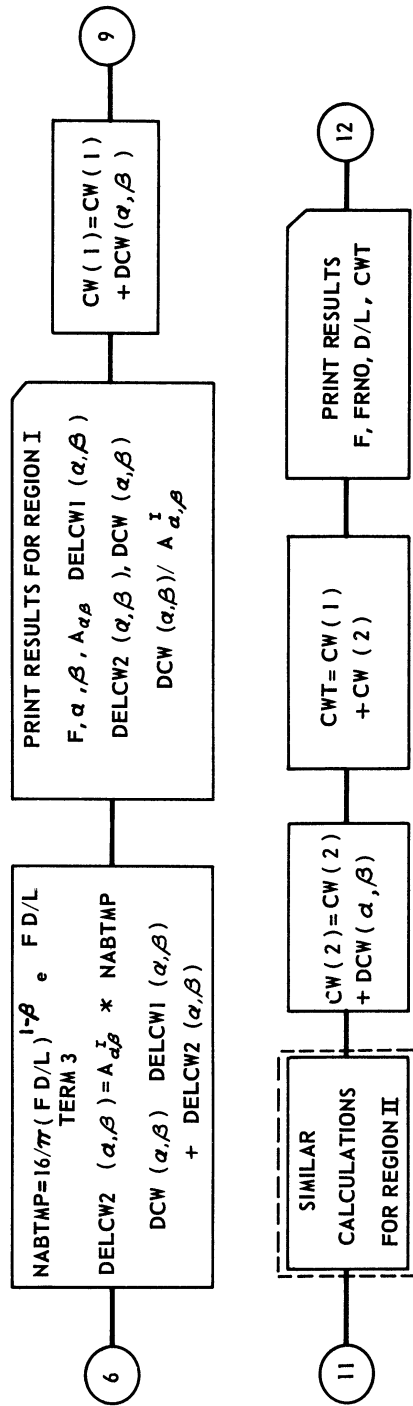


Figure 11 (Concluded)


```

PCNSW6(J)=0B *046 02
PCNSW6(J)=0B *047 02
PCNSW7(J)=0B *048 02
PCNSW7(J)=0B *049 02
PCNSW4(J)=0B *050 02
PCNSW4(J)=0B *051 02
PCNSW5(LJ)=0B *052 02
PCNSW5(J)=0B *053 02
PCNSW5(J)=0B *054 01
PRINT COMMENT$0 RESULTS FOR REGION 1$ *055 01 01
END OF CONDITIONAL *056 01 01
THROUGH L2, FOR J=1,1,J.G, MAX A *057 01 01
ALP=ALPHA(J) *058 01 01
THROUGH L2, FOR K=1,1,K.G, MAX B *059 02 02
BET=BETA(K) *060 02 02
*061 03 03
*062 02 02
*063 03 03
*064 03 03
*065 03 03
*066 04 04
*067 04 04
*068 04 04
WHENEVER NNH=2.E.NN *069 01 01
PR6=PCNS6(NN) *070 02 02
OTHERWISE *071 02 02
PR6=PSIN.(NN,0.,E,MU) *072 02 02
WHENEVER SW(2), PRINT COMMENT$ PSIN$ *073 02 04
WHENEVER SW(2), PRINT RESULTS NN, BET, J1, PR6 *074 02 04
PCNSW6(NN)=1B *075 02 04
PCNS6(NN)=PR6 *076 02 04
END OF CONDITIONAL *077 02 04
OTHERWISE *078 01 04
WHENEVER PCNSW6(NN) *079 01 04
PR6=PCNS6(NN) *080 02 04
OTHERWISE *081 02 04
PR6=PCOS.(NN,0.,E,MU) *082 02 04
WHENEVER SW(2), PRINT COMMENT$ PCOS$ *083 02 04
WHENEVER SW(2), PRINT RESULTS NN, BET, J1, PR6 *084 02 04
PCNSW6(NN)=1B *085 02 04
PCNS6(NN)=PR6 *086 02 04
END OF CONDITIONAL *087 02 04
END OF CONDITIONAL *088 01 04
TERM 8=TERM 8+FACT.(-ALP,J1)*PR6/F.P.(J1+1) *089 04 04
ALP=ALP/2 *090 03 03
WHENEVER ALPH*2.NE.ALPH *091 03 03
NN=2*BET+ALP *092 01 01
NNH=NN/2 *093 01 03
WHENEVER PCNSW7(NN) *094 01 03
PR7=PCNS7(NN) *095 02 03
OTHERWISE *096 02 03
PR7=PCOS.(NN,0.,0.,MU) *097 02 03
WHENEVER SW(2), PRINT COMMENT$ PCOS$ *098 02 03
WHENEVER SW(2), PRINT RESULTS NN, BET,ALP, PR7 *099 02 03
WHENEVER SW(2), PRINT RESULTS PR7, NN *100 02 03
PCNSW7(NN)=1B *101 02 03
PCNS7(NN)=PR7 *102 02 03
END OF CONDITIONAL *103 02 03
OTHERWISE *104 01 03
PR7=0. *105 01 03

```

LOOP

L8

```

END OF CONDITIONAL
TERM 9=Z.(ALP)*PR7/F.P.(ALP+1)
TERM 10=16./13.14159*EDL.P.(BEI-1))*Z.(BEI)*(-1).P.(BEI)*
1 (TERM 8+TERM 9)
DELCL1(ALP,BEI)=AAB(ALP,BEI)*TERM 10
EXECUTE ZERO. (TERM 3)
THROUGH L3, FOR KI=0,1,K1.G. BEI
EXECUTE ZERO. (TERM1,TERM2)
THROUGH L6, FOR JI=0,1,J1.G. ALP
NN=2*K1+J1
WHENEVER (NN/2)*2 .E. NN
WHENEVER PSNSW(NN)
PRI=PSNS1(NN)
OTHERWISE
PRI=PSIN.(NN,FDL,F,MU)
WHENEVER SW(2), PRINT COMMENT$ PSIN$
WHENEVER SW(2), PRINT RESULTS NN,K1,J1,PRI
PSNSW(NN)=1R
PSNS1(NN)=PRI
END OF CONDITIONAL
OTHERWISE
WHENEVER PCNSW(NN)
PRI=PCNS1(NN)
OTHERWISE
PRI=PCOS.(NN,FDL,F,MU)
WHENEVER SW(2), PRINT COMMENT$ PCOS$
WHENEVER SW(2), PRINT RESULTS NN,K1,J1,PRI
PCNSW(NN)=1B
PCNS1(NN)=PRI
END OF CONDITIONAL
END OF CONDITIONAL
TERM1=TERM1+FACT.(-ALP,J1)*PRI/F.P.(J1+1)
ALPH=ALP/2
WHENEVER ALPH#2 .NE. ALP
NN=2*K1+ALP
WHENEVER (NN/2)*2 .E. NN
WHENEVER PSNSW(NN)
PR2=PSNS1(NN)
OTHERWISE
PR2=PSIN.(NN,FDL,F,MU)
WHENEVER SW(2), PRINT COMMENT$ PSIN$
WHENEVER SW(2), PRINT RESULTS NN,K1,ALP,PR2
PSNSW(NN)=1B
PSNS1(NN)=PR2
END OF CONDITIONAL
OTHERWISE
WHENEVER PCNSW(NN)
PR2=PCNS1(NN)
OTHERWISE
PR2=PCOS.(NN,FDL,F,MU)
WHENEVER SW(2), PRINT COMMENT$ PCOS$
WHENEVER SW(2), PRINT RESULTS NN,K1,ALP,PR2
PCNSW(NN)=1B
PCNS1(NN)=PR2
END OF CONDITIONAL
END OF CONDITIONAL
PR2=0.
END OF CONDITIONAL
FORM 2=Z.(ALP)*PR2/F.P.(ALP+1)

```

```

*106 01
*107 03
*108 03
*109 03
*110 03
*111 03
*112 04
*113 04
*114 05
*115 05
*116 05
*117 02
*118 02
*119 02
*120 02
*121 02
*122 02
*123 02
*124 02
*125 01
*126 01
*127 02
*128 02
*129 02
*130 02
*131 02
*132 02
*133 02
*134 02
*135 01
*136 05
*137 04
*138 04
*139 04
*140 04
*141 02
*142 03
*143 03
*144 03
*145 03
*146 03
*147 03
*148 03
*149 03
*150 02
*151 02
*152 03
*153 03
*154 03
*155 03
*156 03
*157 03
*158 03
*159 03
*160 02
*161 01
*162 01
*163 01
*164 04

```

L6


```

PR4=PCNS4(NN)
OTHERWISE
PR4=PCOS.(NN,2.*FDL,F,MU)
WHENEVER SW(2), PRINT RESULTS NN,K1,J1,PR4
PCNS4(NN)=PR4
PCNSW4(NN)=1B
END_OF_CONDITIONAL
END_OF_CONDITIONAL
WHENEVER SW(3), PRINT RESULTS K1,J1,PR1,PR4
TERM1=TERM1+FACT.(-ALP,J1)*PR1/(F.P.J1*NAPE.P.FDL)
TERM 4=TERM 4+FACT.(-ALP,J1)*PR4/(F.P.J1*NAPE.P.(2.*FDL))
ALPH=ALP/2
WHENEVER ALPH*2,NE.,ALP
NN=2*K1+ALP
WHENEVER PCNSH0(NN)
PR2=PCNS0(NN)
OTHERWISE
PR2=PCOS.(NN,FDL,0.,MU)
WHENEVER SW(2), PRINT RESULTS NN,K1,ALP,PR2
PCNSW0(NN)=1B
PCNS0(NN)=PR2
END_OF_CONDITIONAL
WHENEVER PCNSW5(NN)
PR5=PCNS5(NN)
OTHERWISE
PR5=PCOS.(NN,2.*FDL,0.,MU)
WHENEVER SW(2), PRINT RESULTS NN,K1,ALP,PR5
PCNSW5(NN)=1B
PCNS5(NN)=PR5
END_OF_CONDITIONAL
OTHERWISE
PR5=0.
PR2=0.
END_OF_CONDITIONAL
WHENEVER SW(3), PRINT RESULTS K1,J1,PR2,PR5
TERM 2=Z.(ALP)*PR2/(NAPE.P.FDL*F.P.ALPH)
TERM 5=Z.(ALP)*PR5/(NAPE.P.(2.*FDL)*F.P.ALPH)
WHENEVER SW(3), PRINT RESULTS TERM1,TERM2,TERM1+TERM2
WHENEVER SW(3), PRINT RESULTS TERM4,TERM5,TERM4+TERM5
TERM 3=TERM 3+FACT.(-BET,K1)*FDL.P.(BET-K1)*(TERM1+TERM 2)*
1.(-1.)P.(BET+K1)
TERM 6=TERM 6+FACT.(-BET,K1)*(2.*FDL).P.(BET-K1)*(TERM4+
1.TERM 5)*(-1.)P.(BET+K1)
TMP1=16.*TERM3/(3.14159*FDL.P.(BET-1))/F
IMP2=16.*TERM6/(3.14159*FDL.P.(BET-1))/F
DCW1(ALP,BET)=AAB1(ALP,BET)*TMP1
DCW2(ALP,BET)=AAB1(ALP,BET)*IMP2*(-1.)
DCW3(ALP,BET)=DCW1(ALP,BET)+DCW2(ALP,BET)
WHENEVER SW(1)
PRINT COMMENT$ F ALPHA BETA AAB1 THIR
1.D FOURTH SUM COEFF$
PRINT FORMAT OUTPUT,F,ALP,BET,AAB1(ALP,BET),DCW1(ALP,BET),
1 DCW2(ALP,BET),DCW3(ALP,BET), DCW3(ALP,BET)/AAB1(ALP,BET)
END_OF_CONDITIONAL
CW(2) = CW(2) + DCW3(ALP,BET)
CMT=CW(1) + CW(2)
PRINT COMMENT $8$

```

L7

L5

L4

```

L1      PRINT RESULTS F, ERNO, DOVERL,   CWT, CW(1), CW(2)
      PRINT COMMENT$1$
      TRANSFER TO START
*278
*279
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*291
*291
*291
*291
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*297

VECTOR VALUES KLFIRM=$16F5.1*$
VECTOR VALUES MDIM=2,1,10
VECTOR VALUES MDIM1=2,1,10
VECTOR VALUES OUTPUT=$S10,F5.1,2I4,5E15.7*$
VECTOR VALUES SW=2B,0B,0R
VECTOR VALUES NAPE=2.73218281828
INTEGER KK, NNH, KKK
INTEGER N, MAXA, MAXA1, MAXB, MAXR1, I, J, K, J1, K1, ALPHA, ALPHAL,
1  BETA, RETAL, NN, ALP, BET, ALPH
VECTOR VALUES DIM1=2,1,10
VECTOR VALUES DIM2=2,1,10
DIMENSION AAB(1000,MDIM),AAB1(1000,MDIM1), I2AB(1000,DIM1),
1  DELCW1(1000,DIM1), DELCW2(1000,DIM1), DCW(1000,DIM1),
2  DCW1(1000,DIM2), CW(2), ALPHA(50), ALPHAL(50), BETA(50),
3  RETAL(50), KL(100), SW(5)
DIMENSION DCW2(1000,DIM2),DCW3(1000,DIM2)
DIMENSION PSNS1(100),PCNS1(100),PSNSO(100),PCNSO(100),
1  PSNS4(100),PCNS4(100),PSNS5(100),PCNS5(100),
2  PCNSW4(100),PCNSW(100),PSNSW(100),PCNSWJ(100),PSNSW4(100),
3  PCNSW4(100),PCNSW5(100),PCNSW5(100)
DIMENSION PSNSW6(100), PCNSW6(100), PSNSW7(100), PCNSW7(100),
1  PSNS6(100), PCNS6(100), PSNS7(100), PCNS7(100)
BOOLEAN SW, PSNSW, PCNSW, PSNSWJ, PCNSWJ, PSNSW4, PCNSW4, PSNSW5,
1  PCNSW5
POOLEAN PSNSW6, PCNSW6, PSNSW7, PCNSW7
END OF PROGRAM

```

THE FOLLOWING NAMES HAVE OCCURRED ONLY ONCE IN THIS PROGRAM.
 COMPILATION WILL CONTINUE.

START1 *J10

GENERALIZED HAVELock P-FUNCTION

Generalized Havelock P-function in Equations (14a) and (14b) are called PCOS.(KK,FDL,FF, MUMU) and PSIN.(KK,FDL,FF, MUMU) respectively, where

KK: integer order of \bar{P} -function i.e., $2n-1$ or $2n$.

FDL: argument y

FF: argument x

MUMU: a number corresponding to an allowable maximum error in numerical integration.

As θ is increased toward $\frac{\pi}{2}$, the oscillating frequency of the integrand increases but the amplitude decreases making the contribution to the integrand smaller. The method of integration is by Gauss-Legendre 9-point quadrature formulae applied in the intervals between two consecutive zero values of the integrand. Zero values of integrands are found from

$$\frac{\cos}{\sin}(x \sec \theta) = 0 .$$

Whenever the area under a half-cycle is smaller than a specified number (MUMU), the integration is automatically terminated. The error is always smaller than that due to the oscillation.

MAD (24 SEP 1964 VERSION) PROGRAM LISTING

EXTERNAL FUNCTION (KK,FOL,FF,MUMU)
CALCULATION OF GENERALIZED HAVELOCK P FUNCTION BY NUMERICAL

```

INTEGRATION
ENTRY TO PSIN.
NAPE=2.718281828
LOWER=C.
SUM=C.
A=FF/3.1415927
MA=A+C.9999999
THROUGH LOOP A2, FOR M=MA+1,M.G.600
UPPER=ARCCOS.(A/M)
AREA=C.
THROUGH LOOP A1, FOR K=1,1,K.G.9
X=(1.+HG(1.-K))*UPPER+(1.-HG(1.-K))*LOWER)*.5
COSINE=COS.(X)
SINE=SIN.(X)
TANGT=SINE/COSINE
C=FOL*TANGT*TANGT
WHENEVER C.G.30.
AREA=C.
TRANSFER TO E1
END OF CONDITIONAL
EPC=NAPE .P. C
COSPR=COSINE .P. KK
WHENEVER KK.E.0, COSPR=1.0
INTGRD=COSPR/EPC*SIN.(FF/COSINE)
AREA=AREA+IG(K)*INTGRD*(UPPER-LOWER)*C.5
LIMIT=AREA
WHENEVER .ABS. (LIMIT) .LE. MUMU, TRANSFER TO LOOP A3
LOWER=UPPER
SUM=SUM+AREA
SUM=SUM*(1.-1.)P.(KK/2 )
FUNCTION RETURN SUM
ENTRY TO PCOS.
NAPE=2.718281828
LOWER=C.
SUM=C.
A=FF/3.1415927
MA=A+C.4999999
THROUGH LOOP A5, FOR M=MA+1,M.G.600
UPPER=ARCCOS.(A*2./(2.*M+1.))
AREA=C.
THROUGH LOOP A4, FOR K=1,1,K.G.9
X=(1.+HG(1.-K))*UPPER+(1.-HG(1.-K))*LOWER)*.5
COSINE=COS.(X)
SINE=SIN.(X)
TANGT=SINE/COSINE
C=FOL*TANGT*TANGT
WHENEVER C.G.30.
AREA=C.
TRANSFER TO E4
END OF CONDITIONAL
EPC=4APR .P. C

```

```

COSPR=COSINE *P. KK
WHENEVER KK.E.0, COSPR=1.0
INTGRD=COSPR/EPC*COS.(FF/COSINE)
  AREA=AREA+IG(K)*INTGRD*(UPPER-LOWER)*0.5
LIMIT=AREA
WHENEVER ABS.(LIMIT) .LE. MUMU, TRANSFER TO LOOP A6
LOWER=UPPER
SUM=SUM+AREA
SUM=SUM*(-1.) *P.((KK+1)/2 )
FUNCTION RETURN SUM
INTEGER KK,MA,M,K
DIMENSION IG(10), H6(10)
VECTOR VALUES HG(1) = .96816024, .83603111, .61337143, .32425342
  1 .0, -.32425342, -.61337143,
  2 -.83603111, -.96816024
VECTOR VALUES IG(1) = .08127439, .18064816, .26061070,
  1 .31234708, .33023936, .31234708,
  2 .26061070, .18064816, .08127439
END OF FUNCTION

```

```

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*053
*054
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*067

```


HULL FUNCTION

The program below describes the calculation of the coefficients of the polynomials in the variables ξ and ζ given by Equations (23) and (24). The form of these expressions are such that a flow diagram is superfluous.

MAD (09 AUG 1965 VERSION) PROGRAM LISTING

THE HULL FUNCTION POLYNOMIAL
MICHELSEN AND KIM

EXTERNAL FUNCTION FACT.(N,P)
INTERNAL FUNCTION F.(KKK)=FACT.(1,KKK)
INTERNAL FUNCTION FAT.(II,JJ)=F.(II)/(F.(JJ)*F.(II-JJ))

PRINT COMMENT \$1\$ HULL FUNCTIONS\$
PRINT COMMENT \$8
PRINT COMMENT \$8

1 FINN C. MICHELSEN AND HUN CHOL KIMS

PRINT COMMENT \$4\$ GENERAL INPUT\$
PRINT COMMENT \$4\$ READ AND PRINT DATA, MM,NN ,SM,SW(1)
EXECUTE FTRAP.
FV=NN+1

AM=2*MM+1
BM=2*NN+1
AM1=MM
BM1=NN

DIM(1)=NN+2
DIM(2)=NN +1
DIM1(1)=AM+2
DIM1(2)=AM+1

PRINT COMMENT \$1\$ POLYNOMIAL COEFFICIENTS ARE\$
PRINT COMMENT \$8

PRINT RESULTS C(0,0)...C(MM,NN)
THROUGH LOOP 1, FOR A=0,1,A.G.AM
THROUGH LOOP 1, FOR B=0,1, B.G. BM
AAB2(A,B)=0.

LCCP 1
AAB(A,B)=0.
THROUGH LOOP 2, FOR I=0,1, I.G. AM1
THROUGH LOOP 2, FOR J=0,1, J.G. BM1
THROUGH LOOP 2, FOR M=0,1, M.G. AM1
THROUGH LOOP 2, FOR N=0,1, N.G. BM1
WHENEVER C(I,J).E.0..OR.C(M,N).E.0., TRANSFER TCENTR2
THROUGH OOP 2, FOR L=0,1, L.G. (I+1)
THROUGH OOP 2, FOR K=C,1, K.G. M

TEMP=
C(I,J)*F.(I+1)*(-1.)-P.L*
1 .5.P.(I+1-L)/(F.(L)*F.(I+1-L))*C(M,N)*F.(M)*F.(N)/(FACT.(I+1,
2 M+1)*FACT.(J+1,N+1))*FACT.(I+1,K)*.5.P.K/F.(K)
WHENEVER SW(1)
PRINT FORMAT CHK, F.(I+1), (-1.)P.L, .5.P.(I+1-L), F.(L),
1 F.(I+1-L), F.(M), F.(N), FACT.(I+1,M+1), FACT.(J+1, N+1),
2 FACT.(I+1,K), .5.P.K, F.(K)

*C01
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```

VECTOR VALUES CHK=$S2,12E9.3*$
PRINT RESULTS AAB(L-K+M,N+J+1), C(I,J), C(M,N)
END CF CONDITIONAL
LOOP 2
  AAB(L-K+M,N+J+1)=AAB(L-K+M,N+J+1)+TEMP
CONTINUE
CONTINUE
THROUGH LOOP 3, FOR I=C,1,I.G. AMI
THROUGH LOOP 3, FOR J=0,1,J.G. BM1
THROUGH LOOP 3, FOR M=C,1,M.G. AMI
THROUGH LOOP 3, FOR N=0,1,N.G. BM1
  WHENEVER C(I,J).E.O..OR.C(M,N).E.O., TRANSFERTCENTR3
THROUGH OOP 3, FOR K=0,1, K.G. M
THROUGH OOP 3, FOR P=C,1, P.G. K
  TEMP= (-1.)P.I*.5.P.(I+1)*C(I,J)*
1 C(M,N)*F.(M)*F.(N)/(FACT.(I+1,M+1)*FACT.(J+1,N+1))*(-1.)P.P*
2 FACT.(I+1,K)*.5.P./F.(P)*F.(K-P)
OOP 3
ENTR 3
  CONTINUE
  CONTINUE
  CONTINUE
THROUGH LOOP 4, FOR I=0,1,I.G. AMI
THROUGH LOOP 4, FOR J=0,1,J.G. BM1
THROUGH LOOP 4, FOR M=C,1,M.G. AMI
THROUGH LOOP 4, FOR N=0,1,N.G. BM1
  WHENEVER C(I,J).E.O..OR.C(M,N).E.O., TRANSFER ICENTR4
THROUGH OOP 4, FOR K=0,1,K.G. M
THROUGH OOP 4, FOR P=C,1,P.G. N
THROUGH OOP 4, FOR L=C,1,L.G.(I+K+1)
  TEMP= -C(I,J)*C(M,N)*(-1.)P.
1 (I+K+P+L)*FA T.(N,P)*FA T.(M,K)*FA T.(I+K+1,L)*.5.P.L/
2 ((I+K+1)*(J+P+1))
  WHENEVER SW(1)
PRINT RESULTS -C(I,J),C(M,N),(-1.)P.(I+K+P+L),FA T.(N,P),FA
T.(M,K),FA T.(I+K+1,L),.5.P.L/((I+K+1)*(J+P+1))
END CF CONDITIONAL
OOP 4
ENTR 4
  AAB2(M+I-L+1,N-P)=AAB2(M+I-L+1,N-P)+TEMP
CONTINUE
CONTINUE
THROUGH LOOP 5, FOR I=0,1,I.G. AMI
THROUGH LOOP 5, FOR J=0,1,J.G. BM1
THROUGH LOOP 5, FOR M=C,1,M.G. AMI
THROUGH LOOP 5, FOR N=0,1,N.G. BM1
  WHENEVER C(I,J).E.O..OR.C(M,N).E.O., TRANSFERTCENTR5
THROUGH OOP 5, FOR K=0,1,K.G. M
THROUGH OOP 5, FOR P=C,1,P.G. N
THROUGH OOP 5, FOR L=C,1,L.G.(I+K+1)
THROUGH OOP 5, FOR Q=C,1,Q.G.(J+P+1)
  TEMP= C(I,J)*C(M,N)*
1 (-1.)P.(I+K+P+Q+L)*FA T.(N,P)*FA T.(M,K)*FA T.(I+K+1,L)*
2 FA T.(J+P+1,Q)*.5.P.L/((I+K+1)*(J+P+1))
OOP 5
ENTR 5
  AAB2(M+I-L+1,N+J+1-Q)=AAB2(M+I-L+1,N+J+1-Q)+TEMP
CONTINUE
CONTINUE
THROUGH LOOP 6, FOR I=C,1,I.G. AMI
THROUGH LOOP 6, FOR J=C,1,J.G. BM1
THROUGH LOOP 6, FOR M=C,1,M.G. AMI
THROUGH LOOP 6, FOR N=0,1,N.G. BM1
  WHENEVER C(I,J).E.O..OR.C(M,N).E.O., TRANSFERTCENTR6
THROUGH LOOP 6, FOR K=C,1,K.G. M

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*C42 C1
*C43 C1
*C44 01
*C45 C8
*C46 C6
*C47 C6
*C48 C2
*C49 C3
*C50 C4
*C51 C5
*C52 C6
*C53 C6
*C54 07
*C55 C8
*C56 C8
*C57 C6
*C58 C6

*C59 C2
*C60 C3
*C61 C4
*C62 C5
*C63 C6
*C64 C6
*C65 07
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*C67 C5
*C67 C5
*C68 C5
*C69 01
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*C79 C6
*C80 07
*C81 C8
*C82 C5
*C83 10
*C83 C6
*C84 C6
*C85 C6
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      THROUGH OOP 6, FOR P=0,I,P.G.N
      TEMP= C(I,J)*C(M,N)*(-1.)-P.(I+K+P)*
      FAT.(N,P)*FAT.(M,K)*.5.P.(I+K+1)/(I+K+1)*(J+P+1)
      AAB2(M-K,N-P)=AAB2(M-K,N-P)+TEMP
      CONTINUE
      LOOP 6
      LOOP 6
      THROUGH LOOP 7, FOR I=C+1,I.G.AM1
      THROUGH LOOP 7, FOR J=C+1,J.G.BM1
      THROUGH LOOP 7, FOR M=0,1,M.G. AM1
      THROUGH LOOP 7, FOR N=0,1,N.G. BMI
      WHENEVER C(I,J).E.C..OR.C(M,N).E.O., TRANSFERTCENTR7
      THROUGH OOP 7, FOR K=C+1,K.G.M
      THROUGH OOP 7, FOR P=0,1,P.G.N
      THROUGH OOP 7, FOR Q=C+1,Q.G.(J+P+1)
      TEMP= C(I,J)*C(M,N)*(-1.)-P.
      1 (I+K+P+1+Q)*FA T.(N,P)*FA T.(M,K)*FA T.(J+P+1)
      2 (I+K+1)/(I+K+1)*(J+P+1)
      AAB2(M-K,N+J-Q+1)=AAB2(M-K,N+J-Q+1)+TEMP
      CONTINUE
      LOOP 7
      LOOP 7
      PRINT COMMENT $1$
      PRINT COMMENT$0
      RESULT FOR REGION 1$

      WHENEVER SW .AND. AAB(A,B).NE.O.
      PUNCH FORMAT OUT,A,B,AAB(A,B),LL,K
      END OF CONDITIONAL
      PRINT RESULTS LL,K,AAB(A,B)

      PRINT COMMENT$1$
      PRINT COMMENT $0
      RESULTS FOR REGION 2$

      WHENEVER SW .AND. AAB2(A,B).NE.C.
      PUNCH FORMAT OUT1,A,B,AAB2(A,B),LL,K
      END OF CONDITIONAL
      PRINT RESULTS LL,K,AAB2(A,B)

      TRANSFER TO START
      BCLEAN SW
      INTEGER MM,NN,FV,K,I,J,M,N,P,Q,L,A,B,AM,BM,AM1,BM1,KKK
      INTEGER II,JJ,LL,PP,QQ
      FORMAT VARIABLE FV
      VECTOR VALUES OUT=$S10,4HAAB(,I2,1H,,I2,2H)=,E25.8,2I5*$
      VECTOR VALUES OUT1=$S10,5HAAB1(,I2,1H,,I2,2H)=E24.8,2I5*$
      VECTOR VALUES DIM=2,0,10
      VECTOR VALUES DIM1=2,1,10
      VECTOR VALUES SW=0B,0B,0B
      DIMENSION C(100,DIM),AAB(500,DIM1), AAB2(500, DIM1)
      DIMENSION SW(10)
      END OF PROGRAM

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