

T H E U N I V E R S I T Y O F M I C H I G A N

COLLEGE OF ENGINEERING

Department of Electrical Engineering

Space Physics Research Laboratory

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MEASUREMENTS OF ATMOSPHERIC PRESSURE, TEMPERATURE, DENSITY,  
AND COMPOSITION AT VERY HIGH ALTITUDES

Prepared for the project by

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## 1.0 INTRODUCTION

This is the sixth report in a series which outlines a research effort whose object is the determination of the ambient pressure, temperature, density and composition of the earth's atmosphere at altitudes where the mean free path of the neutral particles is appreciably greater than the dimensions of the measuring object.

As noted in previous reports, the effort is devoted to the following tasks:

- (a) a theoretical study of the general measurement problem, and several associated problems;
- (b) development of suitable sensors;
- (c) development of associated instrumentation to permit fruitful employment of the sensors; and
- (d) development of an ultra-high-vacuum system capable of achieving pressures as low as the state of the art permits, with the final objective of sensor calibration and testing.

The following sections describe the work done in these areas since the last report.

## 2.0 THEORETICAL STUDY

In previous reports it has been mentioned that an investigation directed towards the determination of the optimum orifice size with reference to the main spherical body and its motion is being carried out. The result of this work is included in this report as an appendix.

Using Eq. (7) from the appendix, values of the orifice area have been computed for a response time of  $10^{-3}$  sec and a chamber volume of  $10 \text{ cm}^3$  at various altitudes. They are given in Table I. These values indicate that for this experiment an orifice area of  $0.5 \text{ cm}^2$  will be a good compromise.

TABLE I  
ORIFICE AREA FOR A RESPONSE TIME OF  $10^{-3}$  SEC,  
AND CHAMBER VOLUME OF  $10 \text{ CM}^3$  AS A FUNCTION OF ALTITUDE

Height (km)	Temperature (K°)	Orifice Area ( $\text{cm}^2$ )
100	199	1.031
125	571	0.609
150	1031	0.453
175	1359	0.395
200	1404	0.388
225	1414	0.387
250	1415	0.387

### 3.0 SENSOR DEVELOPMENT

The "prototype" omegatron (Serial No. 8) which was outlined in the last report has been built and is presently undergoing extensive tests. It is a metallic tube consisting of two main parts, as shown in Fig. 3.1: the envelope, which has the polepieces brazed into position, and the base plate, which supports the tube elements mounted on four sapphire rods. The "plates" are made from nonmagnetic stainless-steel mesh to permit free particle movement. The electron emission is produced by a short .003-in. tungsten filament supported by two molybdenum rods. These rods are glass-coated and are attached to a platinum sheet which is, in turn, spot-welded to the stainless-steel mesh. The anode which collects the electrons is mounted in a similar fashion on the opposite side of the chamber. Electric connections to the tube elements are made by molybdenum rods pushed through a 1/2-in. Teflon gasket supported by the bottom plate. A Teflon O-ring is used to vacuum-seal the base plate to the envelope.

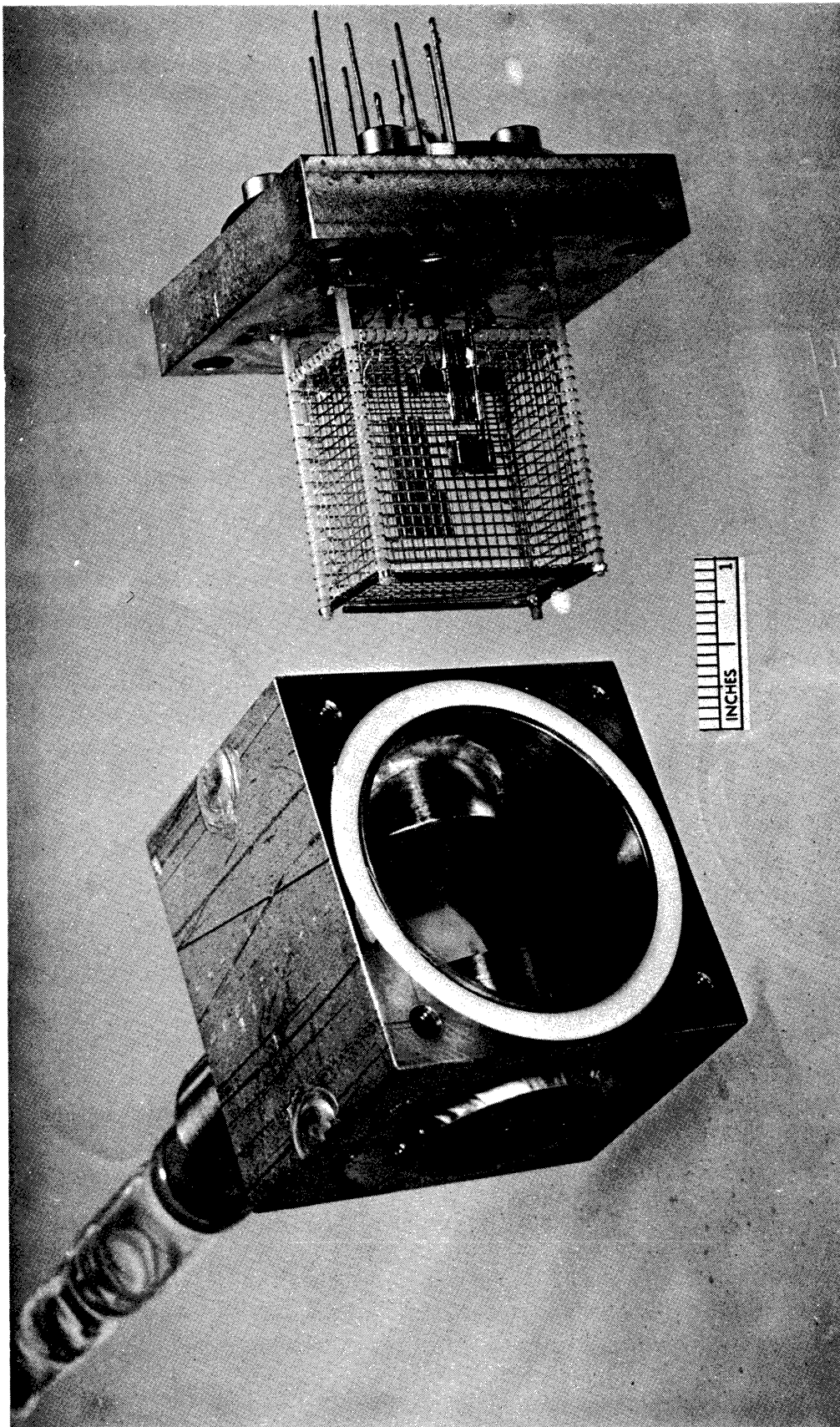


Fig. 3.1. Omegatron, Serial No. 8.

## 4.1 SPHERE

The Vector TRPT 250 all-transistor 1/3 Watt telemetry transmitter has been chosen as the most suitable for this experiment, taking into consideration such factors as size, availability, and power requirements.

As mentioned in the last report, the electrometer system necessary to convert the output current of the omegatron and the ionization gauge to suitable voltage levels for telemetry has been developed and built for another experiment and will be flight-tested in the near future.

The filaments of both the omegatron and the ionization gauge must maintain a constant emission over a wide range of ambient conditions. To achieve this, a filament regulator is necessary. For testing purposes in the laboratory, a series type regulator is satisfactory; such a device has been designed and built, the circuit diagram of which is shown in Fig. 4.1. While the series regulator is a good enough laboratory apparatus, it is not practical for flight use as it is very inefficient. A large fraction of the power from the battery is lost in the series transistor, thus increasing the ampere-hour capacity requirement of the battery and also presenting a secondary problem in heat dissipation.

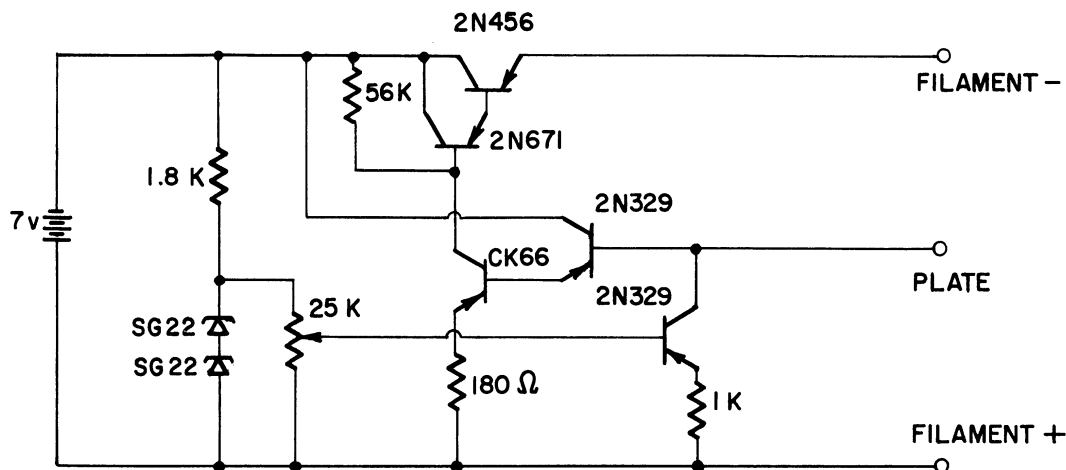


Fig. 4.1. Circuit diagram of the series type filament regulator.

It has been decided that a switching type regulator circuit using transistors would be the most suitable for our purposes. Such a regulator was designed, built, and tested. Figures 4.2 and 4.3 show the completed printed circuit board and the circuit diagram, respectively. To provide the proper heat sink for the power transistor, it is mounted separately from the board. The regulator is basically an asymmetric free-running multivibrator with controllable on-off periods. The operation of the system can be described as follows. Transistors  $T_4$ ,  $T_5$ , and  $T_6$  make up the emitter coupled multivibrator. The filament of the device to be regulated is used as the common emitter resistor. Transistors  $T_5$  and  $T_6$  are coupled together so as to be able to handle the power requirements of the filament. To obtain maximum battery economy, the power supply is made up of two batteries connected in

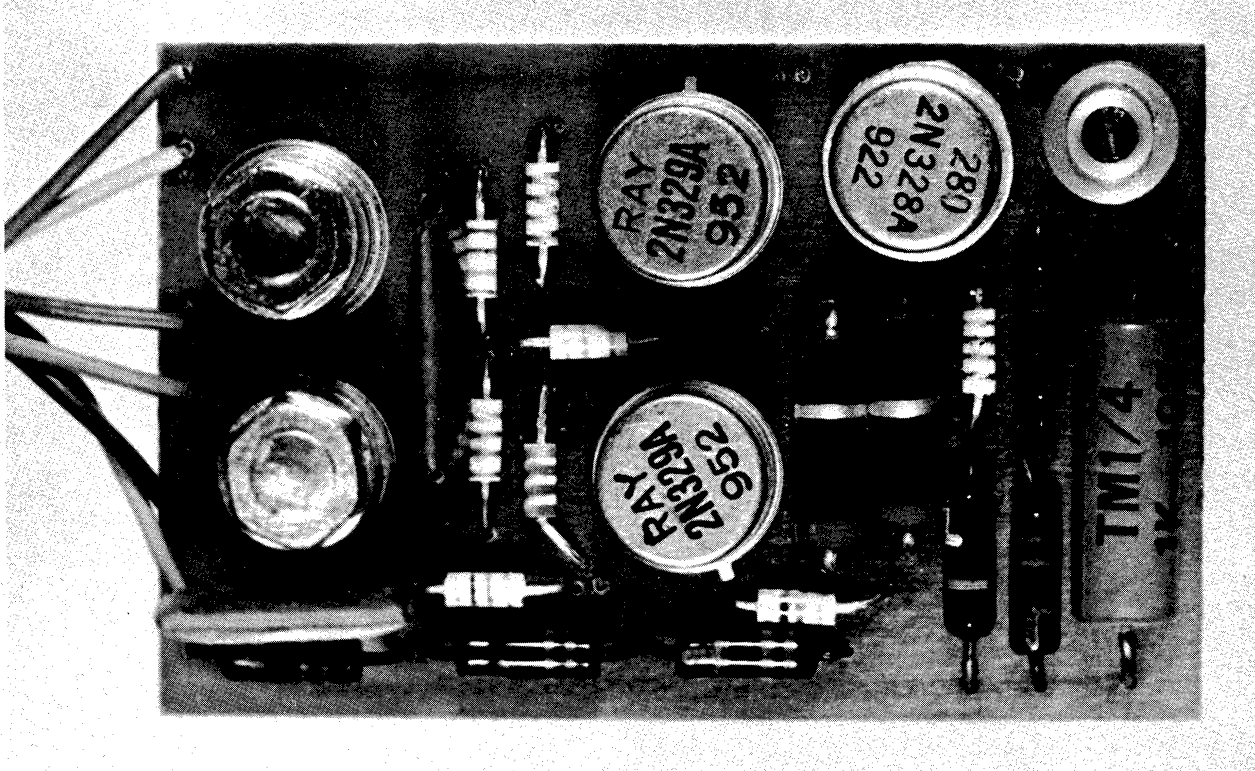
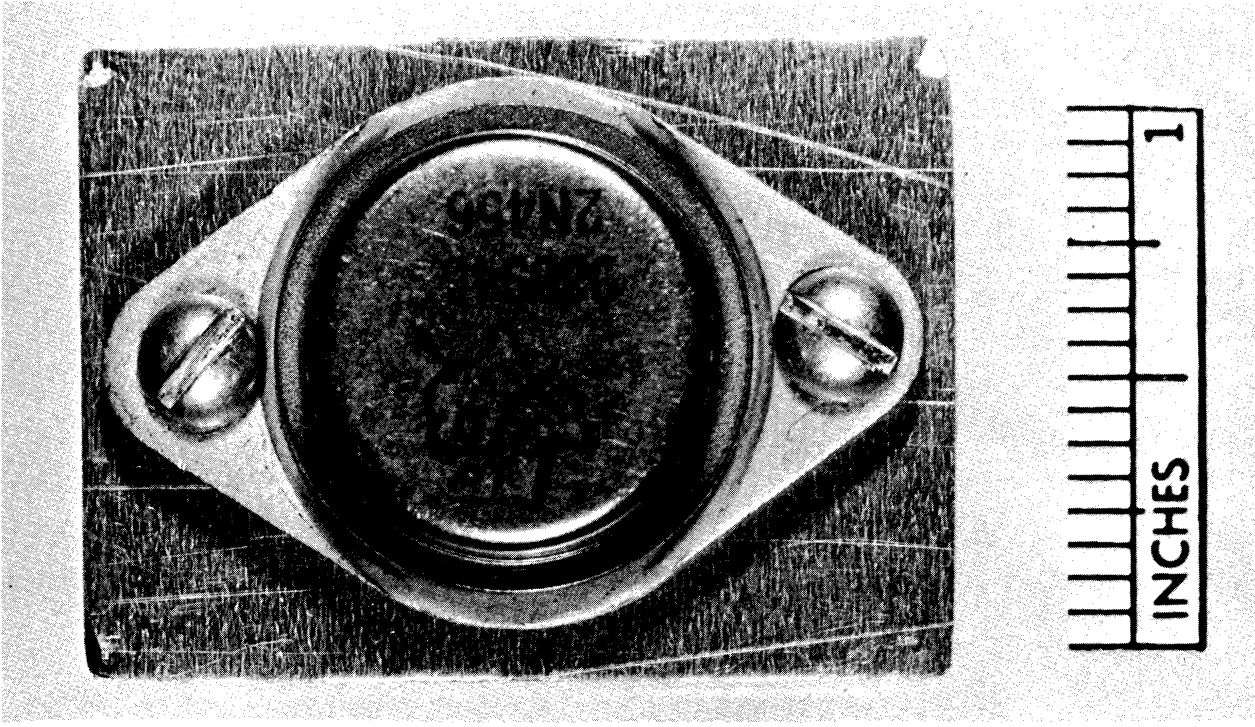


Fig. 4.2. The switching type filament regulator.



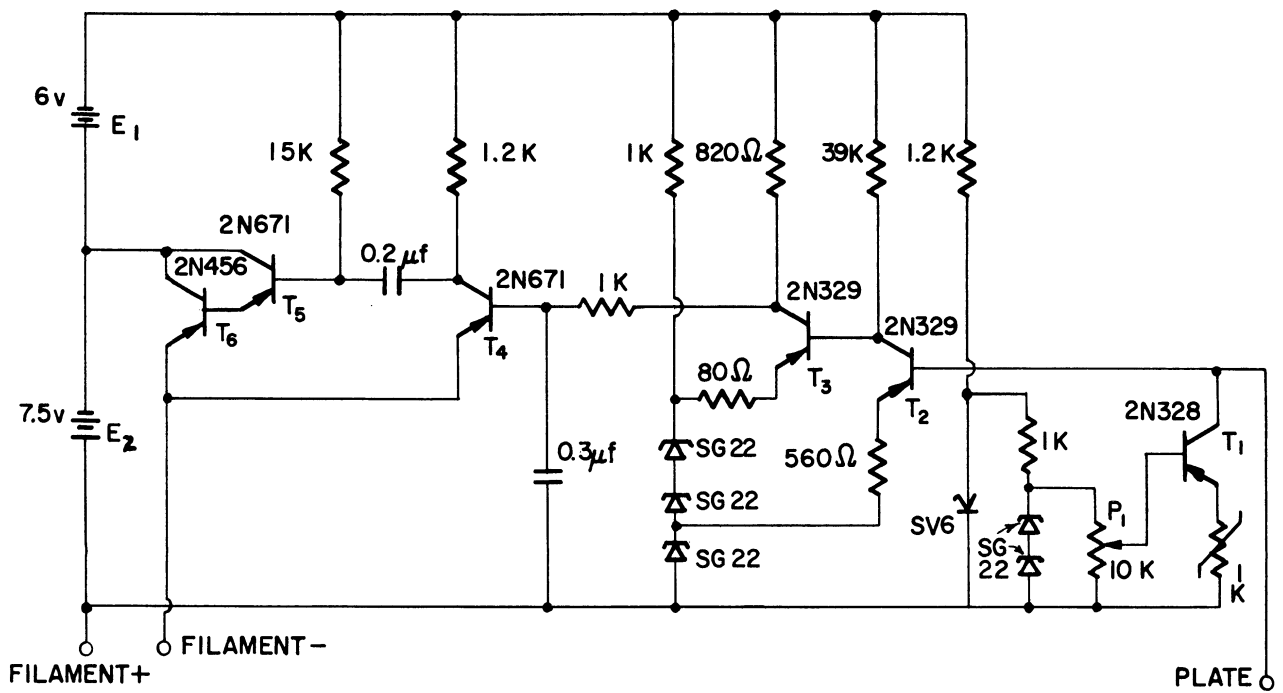


Fig. 4.3. Switching type filament regulator circuit.

series. There is a 7.5-volt high-current-capacity SilverCell providing the power for the filament, and a 6-volt mercury battery connected in series with it, providing a total of 13.5 volts to supply the rest of the circuit. Variation in the on-off rate of the  $T_5$  and  $T_6$  half of the multivibrator is accomplished by varying the base potential of  $T_4$ . This potential is set by  $T_3$  and  $T_2$  which are connected as conventional common emitter amplifiers. To get maximum sensitivity, the base current of  $T_1$  is the algebraic difference between the current from the constant current supply  $T_7$  and the emission current of the device to be regulated. Potentiometer  $P_1$  permits adjustment in the emission level of the filament. The emission current stays constant, to within 1% when  $E_1$  changes from 2.5 to 7.5 volts, and within 4% when filament supply  $E_2$  changes from 6.5 to 9 volts.

Extensive temperature tests have been carried out and it was established that the only component requiring temperature compensation is the transistor  $T_1$  of the reference supply. It has been pointed out<sup>25</sup> that temperature drifts in silicon amplifiers, properly biased and operating at temperatures below  $80^\circ\text{C}$ , are due primarily to changes in the base to emitter voltage  $V_{be}$ . These changes are stable with time, and are a linear function of temperature, so they can be compensated for by positive temperature coefficient resistors. This was done in our circuit by placing a Texas Instrument 1 K $\Omega$  sensor in series with the emitter of transistor  $T_1$ . With this compensation in the circuit the variation in operating level was less than 1% up to temperatures of  $60^\circ\text{C}$ .

There is available from the electrometer system a 400 c/s square wave with an amplitude of 70 volts peak to peak, which is used after rectification to provide the high voltages necessary to operate the omegatrons and the thermionic gauges. Figures 4.4 and 4.5 show the printed circuit board and the circuit diagram of the high-voltage supply, respectively. It consists of two voltage doublers in series with a floating center point. Six zener diodes are connected in series to provide a regulated output voltage of approximately 110 volts.

The omegatron also needs a stable rf oscillator. Such a transistor oscillator is being developed and is ready for testing. It will be described in the next report.

## 4.2 ROCKET NOSE CONE

It was noted in the last report that for proper functioning of the experiment a stable platform is required for launching the sphere. It was also pointed out that, if proper precautions are taken, the rocket can be expected to remain essentially upright as it passes the 100- to 125-km level. At these heights the rocket would normally be spinning at about 1-3 rps, but a simple system employing the pressurizing helium has now been developed by Aerojet which despins the rocket at the end of the powered flight. This is done by diverting the helium to a number of small jets situated in the front end of the rocket which are controlled by a roll gyro. As this system is readily available, it will not be necessary to rotate either the sphere or the nose cone as was originally anticipated. The nose cone carrying the sphere and some associated equipment will be vacuum-sealed at a pressure of approximately  $10^{-5}$  mm Hg to provide vacuum cleanliness and to avoid contamination.

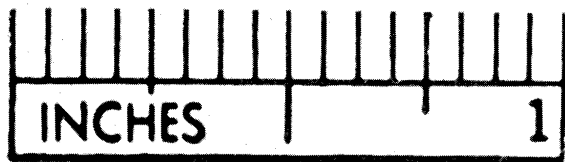
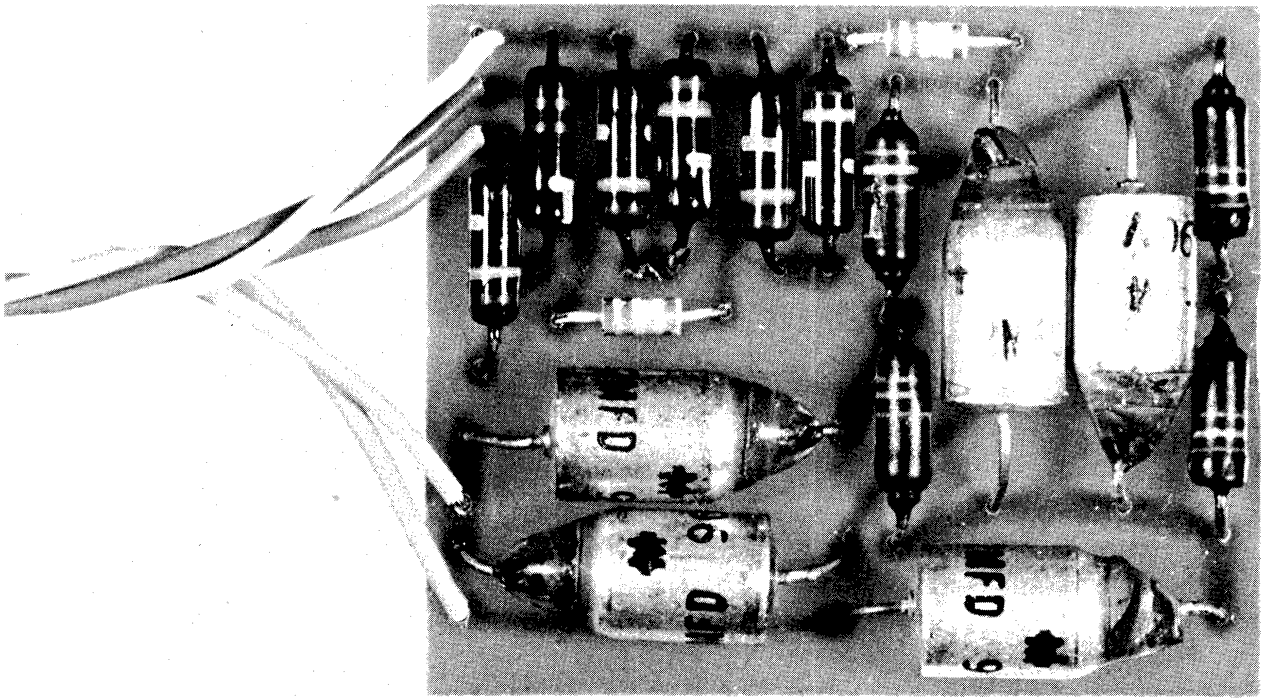


Fig. 4.4. The high-voltage supply.

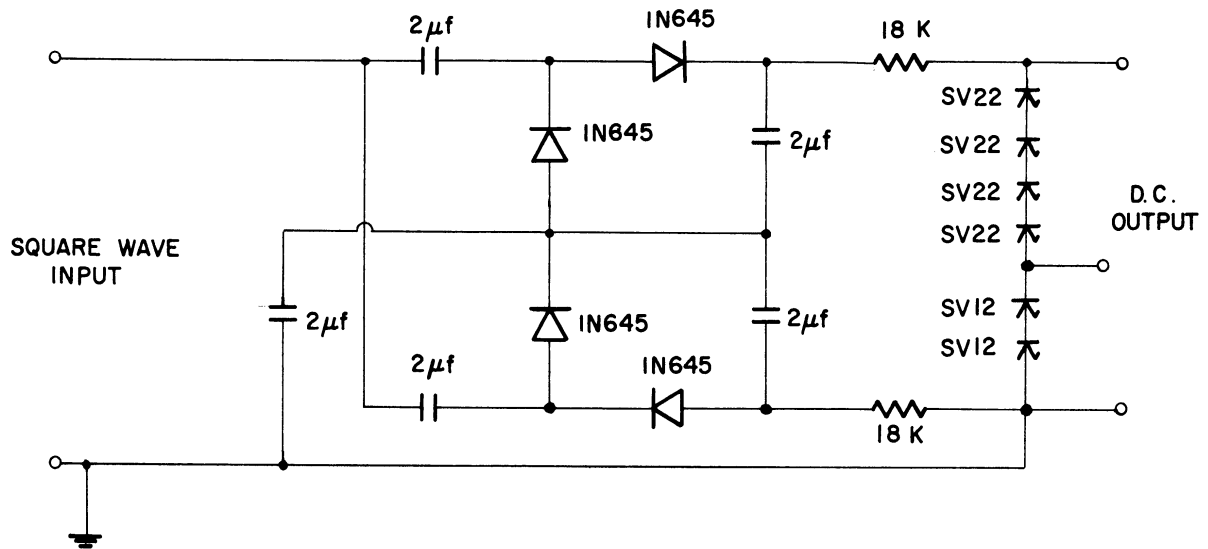


Fig. 4.5. High-voltage supply circuit.

## 5.0 ULTRA-HIGH-VACUUM SYSTEM

No further effort has been devoted to the UHV system. It has been reported earlier that the system is capable of attaining a pressure of  $\sim 5 \times 10^{-11}$  mm Hg and is available for use.

## 6.0 FUTURE WORK

The extended contract is scheduled to terminate on July 31, but as the task has not yet been completed, a further no-cost extension was requested. The funds remaining under this project are expected to cover the launching and associated expenses. However, for the completion of the development and construction of the nose-cone system, the sensors and associated instrumentation, some of which has been developed for separately financed work, support from other sources is being used. It is now planned that the launching will take place late in the fall of this year.

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## APPENDIX

Prepared by Mr. Madhoo Kanal for  
inclusion in this report.



## NOMENCLATURE

- $\Sigma$  = Area of the chamber hole.
- $\Delta t_0$  = Response time of the detector.
- $\Delta t$  = Time interval
- $N_i$  = Initial number density of the gas in the chamber.
- $V$  = Volume of the chamber.
- $C_{m_i}$  = Most probable velocity of the gas particles.
- $N_t$  = Number density of the gas in the chamber at any instant.
- $N_{i0}$  = Number density of the gas in the chamber at  $t = 0$ .

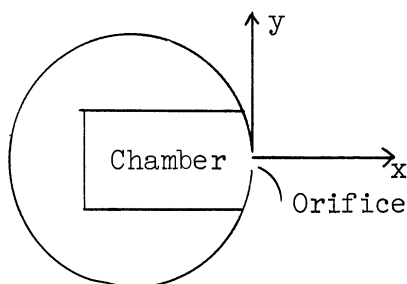
## INTRODUCTION

This section will deal with the study of the relation between the orifice size and the response time of the detector. In the experiment a sphere of radius  $R$  is ejected from the rocket in the upper atmosphere. The sphere moves through this atmosphere and collects certain samples of the gas from it. The gas is collected in a chamber, which is inside the sphere. The chamber has an orifice through which the flow of the gas occurs. Initially the chamber contains a gas at pressure  $P_i$  and temperature  $T_i$  in equilibrium with the outside medium. The equilibrium implies that the number of particles flowing into the chamber is equal to the number of particles flowing out into the surrounding medium in a certain time interval  $\Delta t$ . It is assumed that the gas, the walls of the chamber, and the detector are at the same temperature. A sudden change of density in the outside medium will unbalance the equilibrium between the number of particles flowing in and out of the chamber. If such sudden changes occur in short intervals of time, the orifice should be of that size which will allow the detector to sense these changes in the time interval in which they occur.

## ORIFICE SIZE AND THE RESPONSE TIME OF THE DETECTOR

We begin by deriving the simple relation between the orifice size  $\Sigma$  and the time interval during which any sudden unbalance of equilibrium is brought to equilibrium again. Assume the gas inside and outside of the chamber to have the same law of distribution of velocities at every point of space.

Take the center of the orifice as the origin and draw from the point a system of lines to represent in magnitude and direction the velocities of the different molecules of the gas. Referred to the orthogonal axes, the coordinates of the extremity of any line will be  $u_x, u_y, u_z$ , the components of velocity of the corresponding molecules. Velocity of the sphere will add no components to the components of the particle velocity because of the way the coordinates are chosen.



Since only the particles with the x-component of velocity have the probability of getting out of the chamber, the only direction taken into consideration will be the x-direction. Let  $N_i$  be the number density of the gas inside the chamber and  $C_{m_i}$ , their most probable velocity. Then the fraction of the number of particles per unit volume with the velocity component between  $u_x$  and  $u_x + du_x$  is given by the Maxwellian distribution law,

$$\frac{dN_i}{N_i} = \frac{1}{\sqrt{\pi}} \exp\left(-\frac{U_x^2}{C_{m_i}^2}\right) d\left(\frac{U_x}{C_{m_i}}\right). \quad (1)$$

The number of particles,  $G_i$ , leaving the chamber in time  $\Delta t$  is then given by

$$G_i = \frac{\Sigma \Delta t N_i C_{m_i}}{\sqrt{\pi}} \int_0^{\infty} \frac{U_x}{C_{m_i}} \exp\left(-\frac{U_x^2}{C_{m_i}^2}\right) d\left(\frac{U_x}{C_{m_i}}\right),$$

from which we get

$$G_i = \frac{\Sigma \Delta t N_i C_{m_i}}{2\sqrt{\pi}}. \quad (2)$$

Let the number density of the particles in the chamber at any instant, in the state of unbalanced equilibrium, be  $N_t$ , and let  $N_i$  be the density when the equilibrium is established. Then the net flow of particles at the orifice in infinitesimal time  $\delta t$  is given by

$$\delta NV = - \frac{\sum C_{m_i} (N_t - N_i)}{2\sqrt{\pi}} \delta t , \quad (3)$$

where  $\delta N$  is the infinitesimal change in density of the gas in the chamber during time  $\delta t$ . The minus sign in Eq. (3) signifies that, if  $N_t > N_i$ , there is a net flow of particles from the chamber to the outside medium and hence a decrease in density. From Eq. (3) we get

$$\lim_{\delta t \rightarrow 0} \frac{\delta N}{\delta t} = \frac{dN}{dt} = - \frac{\sum C_{m_i} (N_t - N_i)}{2\sqrt{\pi} V} . \quad (4)$$

The amount of time required for equilibrium to establish is determined by solving the simple differential Eq. (4), with the condition that at  $t = 0$ ,  $N_t = N_{i0}$ , where  $N_{i0}$  is the initial density, i.e., before the equilibrium was unbalanced. Thus from Eq. (4) we get

$$\int \frac{dN}{(N_t - N_i)} = - \frac{\sum C_{m_i}}{2\sqrt{\pi} V} \int dt$$

or

$$\ln (N_t - N_i) = - \frac{\sum C_{m_i}}{2\sqrt{\pi} V} t + C , \quad (5)$$

Where  $C$  is the constant of integration.

Now

$$\text{at } t = 0 \quad N_t = N_{i0}$$

$$\therefore C = \ln(N_{i0} - N_i) .$$

Substituting the value of  $C$  in Eq. (5), we get

$$\ln \left[ \frac{(N_t - N_i)}{(N_{i0} - N_i)} \right] = - \frac{\sum C_{m_i} t}{2\sqrt{\pi} V}$$

$$\therefore N_t = N_i + (N_{i0} - N_i) \exp \left( - \frac{\sum C_{m_i} t}{2\sqrt{\pi} V} \right) \quad (6)$$

From Eq. (6) let us define the time constant  $\Delta t_0$ :

$$\Delta t_0 = \frac{2\sqrt{\pi}V}{\sum C_{m_i}} \quad (7)$$

Substituting  $\Delta t_0$  for  $\frac{2\sqrt{\pi}V}{\sum C_{m_i}}$  in Eq. (6), we get

$$N_t = N_i + (N_{i0} - N_i) \exp\left(-\frac{t}{\Delta t_0}\right) . \quad (8)$$

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