

Infrared Spectra of High Polymers. II. Polyethylene*

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The infrared spectrum of polyethylene has been obtained between about 3000 cm^{-1} and 70 cm^{-1} , polarization measurements on oriented specimens having been obtained to about 350 cm^{-1} . Assignments of the fundamentals are made with the help of a group theory analysis. The assignment of the controversial CH_2 wagging mode is discussed in detail and especially in terms of new evidence from the spectra of *n*-paraffins, both as single crystals and as polycrystalline aggregates. It is shown that this mode is to be assigned to a weak band at 1369 cm^{-1} . A satisfactory determination of the bands which arise from CH_2 groups is also made possible by a study of the paraffin spectra. The splitting of bands in the spectrum is conclusively shown to arise from interactions between molecules in the crystalline phase. The nature of this interaction is discussed in terms of recent theories.

INTRODUCTION

THE preceding paper in this series¹ dealt with the general theory which can be applied to the interpretation of the infrared spectra of high polymers. This paper considers in detail the interpretation of the infrared spectrum of polyethylene. It is natural to begin with this polymer since it is structurally the simplest and has a relatively uncomplicated spectrum. In spite of these helpful features, however, only a partial interpretation has been achieved in the past, primarily as a result of lack of pertinent experimental data on *n*-paraffins and of a rigorous theoretical discussion. We have obtained new experimental data on polyethylene which include (a) an extension of the spectrum from 700 cm^{-1} to 70 cm^{-1} , (b) a study of the spectra of crystalline *n*-paraffins and the changes in spectrum resulting from a change in state, and (c) an examination of the spectrum of a fully deuterated *n*-paraffin. With the help of these added data, and the utilization of group theory methods in the analysis, we have been able to achieve a satisfactory interpretation of the spectrum of polyethylene.

EXPERIMENTAL

The spectroscopic techniques used in obtaining the spectra have been described already.¹ The samples of polyethylene used in this investigation were unoriented sheets obtained from E. I. du Pont de Nemours and Company. Orientation was produced by cold-drawing as fully as possible. It is known that this treatment produces uniaxial orientation of the crystallites, the crystal *c*-axis, or chain axis, becoming parallel to the stretching direction.^{2,3} In addition, a partial orientation of the chains in the amorphous regions occurs, such that they too are parallel to the stretching direction.⁴ The samples of pure *n*-paraffins were made available

to us by Carbide and Carbon Chemicals Company ($\text{C}_{36}\text{H}_{74}$ and $\text{C}_{64}\text{H}_{130}$) and the Shell Development Company ($\text{C}_{27}\text{H}_{58}$, $\text{C}_{28}\text{H}_{58}$, $\text{C}_{29}\text{H}_{60}$, $\text{C}_{30}\text{H}_{62}$). The solid state spectra of these hydrocarbons were obtained by melting a suitable thickness onto a rocksalt plate. A heated microcell was used to examine the spectra in the melted state. Some of these paraffins were studied as single crystals, which were grown by evaporation of a solution of the paraffin in a suitable solvent. The orientation and perfection of the crystals were determined from x-ray diffraction photographs. The fully deuterated *n*-paraffin was purchased from Tracerlab. It does not consist of a single species but closely approximates $\text{C}_{100}\text{D}_{202}$ in composition. Its spectrum was obtained as a thin film melted onto a rocksalt plate.

THE INFRARED SPECTRUM OF POLYETHYLENE

The infrared spectrum of oriented polyethylene between 700 and 100 cm^{-1} has been reported on.⁵ We have recently extended the spectrum to about 70 cm^{-1} , and such a spectrum is shown in Fig. 1. Various specimen thicknesses (indicated on the diagram) had to be used in order to obtain all of the bands under optimum conditions. The dichroism was measured only at frequencies above about 350 cm^{-1} . The polarization is indicated on the spectrum, the solid curve representing radiation polarized with the electric vector perpendicular to the stretching direction and the dashed curve radiation polarized parallel to the stretching direction. The broad band at about 550 cm^{-1} was found to be unpolarized. Below about 400 cm^{-1} interference fringes appeared in the spectrum (corresponding to a refractive index of about 1.46) and somewhat obscured the exact transmission curve. The best value, obtained from samples in which the interference was reduced, is shown by the dotted curve below 400 cm^{-1} . The frequencies of the observed band maxima, their approximate relative intensities, and their dichroism with respect to the direction of stretch are given in Table I. Also shown there are the assignments, which will be discussed later.

⁵ S. Krimm and G. B. B. M. Sutherland, Symposium on Molecular Structure and Spectroscopy, Columbus, Ohio, June, 1952.

* Presented in part at the American Physical Society Meeting, Ann Arbor, March, 1954.

¹ Liang, Krimm, and Sutherland, *J. Chem. Phys.* **25**, 543 (1956).

² C. W. Bunn, *Trans. Faraday Soc.* **35**, 482 (1939).

³ C. W. Bunn and T. C. Alcock, *Trans. Faraday Soc.* **41**, 317 (1945).

⁴ S. Krimm, *J. Chem. Phys.* **22**, 567 (1954).

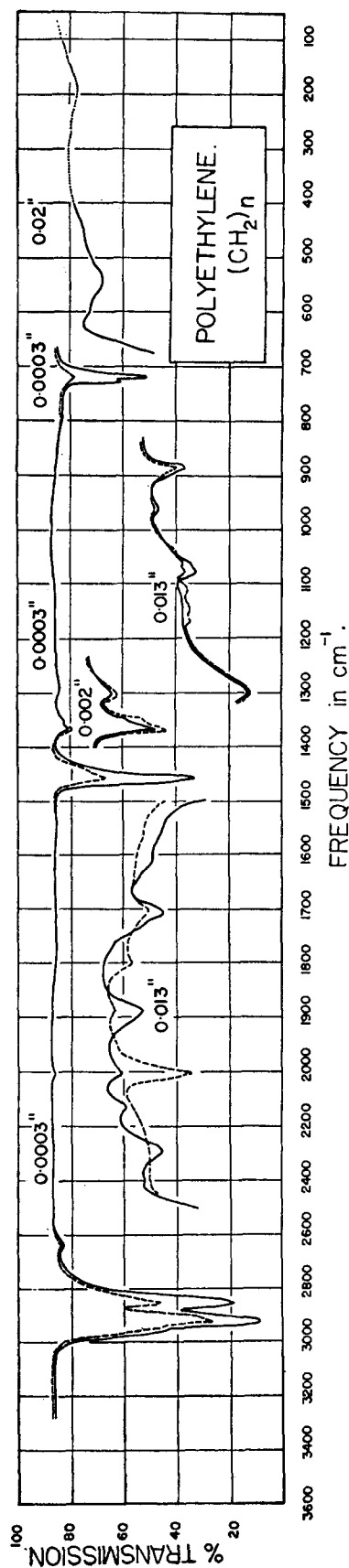


FIG. 1. Infrared spectrum of oriented polyethylene. —: radiation with electric vector polarized perpendicular to stretching direction. - - -: radiation with electric vector polarized parallel to stretching direction.

The observed bands in the prism region ($3200\text{--}700\text{ cm}^{-1}$) are in satisfactory agreement with the results obtained by previous workers⁶⁻⁸ the difference in band maxima never being greater than about 2 cm^{-1} . The spectrum has also been examined by means of a grating spectrometer in the 2900 cm^{-1} and 1400 cm^{-1} regions.⁹ The following differences have been observed. The grating spectra show a weak shoulder at 2875 cm^{-1} which is not resolved in the prism spectra. They also indicate that the band which we have observed at 1460 cm^{-1} is composite, with a strong component at 1463 cm^{-1} , a weaker one at 1473 cm^{-1} , and a weak shoulder at 1456 cm^{-1} . Our value is apparently the mean of these components, which are not resolved in the prism spectra. The three maxima observed by us at 1375 cm^{-1} , 1369 cm^{-1} , and 1353 cm^{-1} are reported at 1378 cm^{-1} , 1370 cm^{-1} , and 1353 cm^{-1} , respectively. As a result of our grating spectrometer observations in the frequency region below 700 cm^{-1} , we have found⁵ two new bands at 600 cm^{-1} (with a weak shoulder at 543 cm^{-1}) and at about 200 cm^{-1} . The 600 cm^{-1} band has also been observed since^{9,10} in prism spectra. Some of the other weaker bands are somewhat variable in their appearance in spectra obtained by different observers. This is due to the fact that they arise from impurities and thus vary from specimen to specimen, as for example the bands at 1710 cm^{-1} (associated with carbonyl groups¹¹) and in the 900 cm^{-1} to 1000 cm^{-1} region (arising from unsaturated carbon-carbon bonds¹²). The spectrum in Fig. 1 may be considered as representative of that of polyethylene.

INTERPRETATION OF THE SPECTRUM

A. Predicted Spectrum of Isolated Polyethylene Chain

Before considering the interpretation of the spectrum of polyethylene, it is necessary to know the structure of the molecule. Although polyethylene may be considered to a first approximation as an infinite chain of CH_2 groups, the chains are found in essentially two different configurations. In the crystalline regions, whose structure is known from x-ray diffraction studies,² the chain is found in the planar zig-zag configuration. In the amorphous regions, which can constitute up to half of the specimen,¹³ the chain configuration is an essentially random one restricted only by the conservation of bond angles and distances.¹⁴ The random con-

⁶ J. J. Fox and A. E. Martin, Proc. Roy. Soc. (London) **A175**, 208 (1940).

⁷ H. W. Thompson and P. Torkington, Proc. Roy. Soc. (London) **A184**, 3 (1945).

⁸ Elliott, Ambrose, and Temple, J. Chem. Phys. **16**, 877 (1948).

⁹ Rugg, Smith, and Wartman, J. Polymer Sci. **11**, 1 (1953).

¹⁰ E. Borello and C. Mussa, J. Polymer Sci. **13**, 402 (1954).

¹¹ Rugg, Smith, and Bacon, J. Polymer Sci. **13**, 535 (1954).

¹² Cross, Richards, and Willis, Discussions Faraday Soc. **9**, 235 (1950).

¹³ S. Krimm and A. V. Tobolsky, J. Polymer Sci. **7**, 57 (1951).

¹⁴ A. Charlesby, J. Polymer Sci. **10**, 201 (1953).

figuration in the amorphous regions becomes a partially oriented one when the sample is stretched.⁴ It is thus very important to keep in mind that the usual spectrum of polyethylene represents a superposition of the spectra of two different types of chain configurations. The spectra of these two components are not identical, and failure to recognize this has led to erroneous assignments.

The spectrum of the polymer in the amorphous regions is not subject to any selection rules,¹ whereas the high symmetry of the chain configuration in the crystalline region does give rise to selection rules in its spectrum. We will now consider the spectrum to be expected from the crystalline polyethylene. In particular, although the complete analysis (to be discussed later) requires consideration of the two molecules in the unit cell, the essential features of the spectrum can be derived from a single isolated molecule, and we discuss this problem first. As pointed out previously,¹ the transition to the full unit cell can generally only have the effect of splitting the frequencies of the isolated molecule, usually by a small amount, and, depending on the symmetry in the space lattice, possibly changing some selection rules.

The repeating unit in a single infinitely long polyethylene chain consists of two CH₂ groups. This unit is shown in Fig. 2 between wavy vertical lines. The symmetry elements which reproduce the repeating unit, a translation along the chain being considered as identity, are also shown in the figure. The screw axis is denoted by C_2^s and the glide plane by σ_g . This set of symmetry elements defines a factor group, in terms of which the spectrum can be analyzed.¹⁵ This mathematical analysis has the physical implication that only those fundamental modes can be active in infrared or Raman spectra in which corresponding atoms in all unit cells move in phase. The factor group for a single polyethylene chain is found^{15,16} to be isomorphous with the point group D_{2h} , and the character table for this point group is given in Table II. Also shown in Table II, (and derived from the factor group analysis) are the number of normal modes and the selection rules in the infrared and Raman spectra under each symmetry species. The number of normal modes, n_i , under any given symmetry species are also shown as they are distributed among pure translations, T , translatory type vibrations, T' , rotatory type vibrations, R' , (which includes a pure rotation R_z) and internal vibrations, n_i' . The table shows that of the $3 \times 6 - 4 = 14$ normal modes of the repeating unit of polyethylene, 5 will give rise to infrared active fundamental frequencies, 8 to Raman active fundamentals, and one will be inactive in both infrared and Raman spectra. The frequencies which are active in infrared are inactive in Raman spectra, and vice versa; that is, the mutual exclusion rule holds.

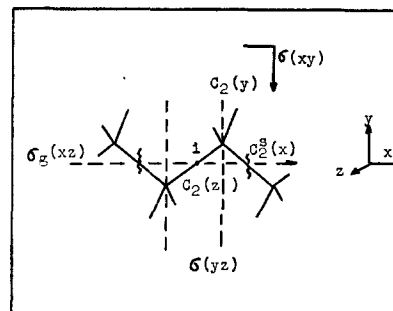


Fig. 2. Portion of polyethylene chain with associated symmetry elements; C_2 —diad, C_2^s —screw diad, σ —mirror plane, σ_g —glide plane, i —center of inversion. The small V's denote CH₂ groups, the C being at the apex and the H's at the extremities of the V.

Furthermore, the polarization of a given infrared fundamental can be determined from the translation belonging to that symmetry species. Thus, modes under species B_{1u} and B_{2u} are expected to show perpendicular polarization with respect to the chain axis, and modes under B_{3u} , parallel polarization. The factor group analysis is therefore immediately useful in predicting the number and type of normal modes which can be infrared active.

We can carry this information one step further and deduce the approximate physical nature of the normal vibrations from the symmetry properties of the various

TABLE I. Infrared spectrum and assignments for polyethylene.

Frequency, cm ⁻¹	Intensity	Polarization	Assignment ^a
2959	w (sh)	...	$\nu_a(\text{CH}_3)$
2925	vs	σ	$\nu_a(\text{CH}_2)^c$
2874 ^b	vw	...	$\nu_s(\text{CH}_3)$
2853	s	σ	$\nu_s(\text{CH}_2)^d$
2640	w	σ	
2295	vw	σ	
2130	vvw	σ	
2010	vw	π	1303+721=2024
1890	vw	σ	
1805	vvw	π	
1710	vw	σ	
1470 ^b	s	σ	} $\delta(\text{CH}_2)$
1460	s	σ	
1456 ^b	vw	...	$\delta_a(\text{CH}_3)$
1375	m	$\pi(?)$	$\delta_s(\text{CH}_3)$
1369	w	π	$\gamma_w(\text{CH}_2)$
1353	w	π	$\gamma_w(\text{CH}_2)$ amorphous regions
1303	w	π	$\gamma_w(\text{CH}_2)$ amorphous regions
1170	vvw	π	
1150	vvw	π	
1110	vvw	π	
1080	vw	σ	$\nu(\text{CC})$ amorphous regions
1065	vw(sh)	σ	$\nu(\text{CC})$
965	vvw	...	
888	vw	σ	$\gamma_r(\text{CH}_3)$
731	m	σ	} $\gamma_r(\text{CH}_2)$
721	s	σ	
600	vw	...	
543	w	...	
200	vw	...	

^a ν = stretching, δ = bending, γ_w = wagging, γ_r = rocking.

^b Observed in grating spectra of polyethylene.⁹

^c Split in n -paraffin spectra, with components at 2924 cm⁻¹ and 2899 cm⁻¹.

^d Split in n -paraffin spectra, with components at 2857 cm⁻¹ and 2850 cm⁻¹.

¹⁵ S. Bhagavantam and T. Venkatarayudu, *The Theory of Groups and Its Application to Physical Problems* (Andhra University, Waltair, India, 1951).

¹⁶ M. C. Tobin, *J. Chem. Phys.* **23**, 891 (1955).

TABLE II. Character table, number of normal modes, and selection rules for a single isolated polyethylene chain.

D_{2h}	E	$C_2(z)$	$C_2(y)$	$C_2^*(x)$	i	$\sigma(xy)$	$\sigma(xz)$	$\sigma(yz)$	n_i	T	T'	R'	n_i'	IR	R
A_g	1	1	1	1	1	1	1	1	3	0	1	0	2	f ^a	p
B_{1g}	1	1	-1	-1	1	1	-1	-1	2	0	1	1	0	f	d
B_{2g}	1	-1	1	-1	1	-1	1	-1	1	0	0	1	0	f	d
B_{3g}	1	-1	-1	1	1	-1	-1	1	3	0	R_x	1	1	f	d
A_u	1	1	1	1	-1	-1	-1	-1	1	0	0	1	0	f	f
B_{1u}	1	1	-1	-1	-1	-1	1	1	3	T_y	0	1	1	a	f
B_{2u}	1	-1	1	-1	-1	1	-1	1	3	T_z	0	0	2	a	f
B_{3u}	1	-1	-1	1	-1	1	1	-1	2	T_x	0	1	0	a	f

^a f: forbidden, a: active, p: polarized, d: depolarized.

species. This implies knowledge of the normal vibrations of a single CH_2 group, which we possess from the results of detailed studies on small molecules. The normal modes of a polyethylene chain, obtained in this way, are shown in Fig. 3. It will, of course, be realized that the relative amplitudes of motion of the various atoms are not determinable from symmetry considerations alone, but depend upon the particular masses and on the intramolecular force field. The normal modes of the component CH_2 groups are, however, well known, as is also the relative phase of the motion in neighboring groups,^{1,17} viz., 0 or π . The modes can therefore be very closely approximated by the above procedure. The derivation of the normal modes of the chain from those of the component CH_2 groups leads to a natural classification of the physical nature of the chain modes in terms of the corresponding group modes. This classification is shown in Table III, and refers to the modes pictured in Fig. 3. The expected polarization of the infrared active modes is indicated by σ , for perpendicular, and π , for parallel. The four zero frequencies, corresponding to the three translations and the rotation about the chain axis, are indicated respectively by T_x , T_y , T_z , and R_x in Fig. 3 and Table III. It will be noted that, of the 5 infrared active fundamentals, 4 are expected to have perpendicular polarization, $\nu_s(\text{CH}_2)$ (B_{2u}), $\nu_a(\text{CH}_2)$ (B_{1u}), $\delta(\text{CH}_2)$ (B_{2u}), and $\gamma_r(\text{CH}_2)$ (B_{1u}) and one is expected to have parallel polarization, $\gamma_w(\text{CH}_2)$ (B_{3u}). The skeletal modes A_g and B_{1g} are inactive in the infrared, as is the twisting frequency, $\gamma_t(\text{CH}_2)$.

We have been concerned so far with the fundamental frequencies of an infinitely long chain of CH_2 groups. In reality, the polyethylene chain is not infinitely long, and there is also evidence for some branching. Both of these effects account for the presence of CH_3 groups, which can and do contribute appreciably to the spectrum.⁶ It will therefore be necessary to know which bands in the spectrum arise from CH_3 groups. Since the CH_3 groups are essentially randomly located in the specimen, it is most probable that there are no sharp selection rules for their fundamental frequencies. We would therefore expect to find in the spectrum fre-

quencies assignable to the stretching, bending, wagging, and rocking motions of the hydrogen atoms in the CH_3 group. These will be considered in greater detail in a succeeding section.

B. Assignments of Active CH_2 Fundamentals of Polyethylene Chain

We now consider the correlation of the bands in the observed spectrum with those predicted from the theory. As mentioned previously, we will at present neglect the effects on the spectrum of neighboring chains in the crystal, since these effects are expected to be of secondary importance with respect to the assignment of fundamental frequencies. The results of considering these interactions will be discussed later. In what follows we shall deal only with the infrared active fundamentals.

The assignments for most of the bands in the spectrum have already been satisfactorily made. Thus, a detailed investigation of the $\nu_a(\text{CH}_2)$ and $\nu_s(\text{CH}_2)$ internal modes of the CH_2 group in hydrocarbons and polyethylene⁶ has established the frequencies of these modes at 2926 cm^{-1} and 2853 cm^{-1} , respectively. Similarly, it is well established^{18,19} that the $\delta(\text{CH}_2)$ vibration gives rise to absorption at about 1460 cm^{-1} . It may therefore be assumed that the strong band at about 1460 cm^{-1} in polyethylene is due mainly to this bending mode of the CH_2 group. The rocking mode of the CH_2 group, $\gamma_r(\text{CH}_2)$, has been shown^{20,21} to give rise to the remaining intense band in the spectrum near 725 cm^{-1} . (The doubling of this band in the spectrum will be considered later.) It will be noted that these frequencies have been predicted to be perpendicularly polarized and are observed to be so. (See Fig. 1.)

The assignment of the remaining active fundamental, the CH_2 wagging mode, $\gamma_w(\text{CH}_2)$, has not been as unequivocal as that of the four fundamentals just considered. Before discussing this assignment in detail, it may be well to review the reasons for the controversy

¹⁸ G. Herzberg, *Infrared and Raman Spectra* (D. Van Nostrand Company, Inc., New York, 1945), p. 195.

¹⁹ L. Kellner, Proc. Phys. Soc. (London) **A64**, 521 (1951).

²⁰ N. Sheppard and G. B. B. M. Sutherland, Nature **159**, 739 (1947).

²¹ G. B. B. M. Sutherland and A. Vallance Jones, Nature **160**, 567 (1947).

¹⁷ Liang, Sutherland, and Krimm, J. Chem. Phys. **22**, 1468 (1954).

over it. As can be seen from our analysis, the $\gamma_w(\text{CH}_2)$ mode is expected to show parallel polarization in the spectrum of an oriented sample. Two prominent parallel bands were found⁸ in the spectrum, at about 1375 cm^{-1} and 1303 cm^{-1} . On the basis of a calculation of the expected normal frequencies of polyethylene, the more intense 1375 cm^{-1} band had been assigned^{22,23} to the $\gamma_w(\text{CH}_2)$ mode. This assignment has been criticized,²⁴ however, on the grounds that there is a well-established correlation⁷ between the 1375 cm^{-1} band and the CH_3 group content of the specimen. These authors²⁴ have accordingly preferred assigning the 1375 cm^{-1} band to a symmetrical CH_3 bending mode, $\delta_s(\text{CH}_3)$ and therefore assigning the 1303 cm^{-1} band to the $\gamma_w(\text{CH}_2)$ mode. Although other authors²⁵⁻²⁷ have favored the former assignment, in spite of the CH_3 group correlation, there has been presented no conclusive evidence for either choice. We will now discuss in detail an unambiguous assignment of the $\gamma_w(\text{CH}_2)$ fundamental.

To begin with, it has been noted²⁸ that the assignment of the 1303 cm^{-1} band to the $\gamma_w(\text{CH}_2)$ fundamental is unsatisfactory because this band is present only when the molecule is in a disordered phase. This is illustrated

TABLE III. Classification of normal modes of polyethylene chain.

Type of motion	Group mode	Phase between groups	Chain mode
Internal hydrogen vibrations of CH_2 group	CH_2 symmetric stretching, $\nu_s(\text{CH}_2)$	0	A_g
		π	$B_{2u}(\sigma)$
	CH_2 asymmetric stretching, $\nu_a(\text{CH}_2)$	0	B_{3g}
		π	$B_{1u}(\sigma)$
	CH_2 bending, $\delta(\text{CH}_2)$	0	A_g
External hydrogen vibrations of CH_2 group	CH_2 wagging, $\gamma_w(\text{CH}_2)$	0	$B_{2u}(\sigma)$
		π	$B_{3u}(\pi)$
	CH_2 rocking, $\gamma_r(\text{CH}_2)$	0	B_{3g}
		π	$B_{1u}(\sigma)$
	CH_2 twisting, $\gamma_t(\text{CH}_2)$	0	A_u
		π	B_{2g}
Skeletal vibrations (CH_2 group moving as a unit)	x direction	0	$B_{3u}(T_x)$
		π	B_{1g}
	y direction	0	A_g
		π	$B_{2u}(T_y)$
	z direction	0	$B_{3g}(R_x)$
		π	$B_{1u}(T_z)$

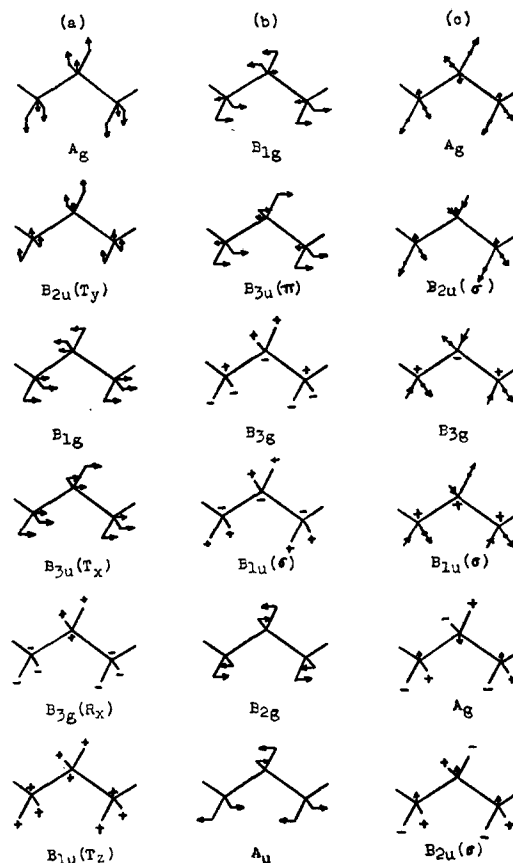


FIG. 3. Normal modes of a single polyethylene chain. Column (a): skeletal vibrations, (b): external hydrogen vibrations of CH_2 group, (c): internal hydrogen vibrations of CH_2 group.

in Fig. 4, which shows the spectrum of the n -paraffin $\text{C}_{64}\text{H}_{130}$, (a) in the solid state (as a polycrystalline aggregate deposited from the melt), (b) at its melting point, and (c) about 15° above the melting point. It is seen that the 1303 cm^{-1} band is not present in the spectrum of the crystalline material, (a), which it should be if it arises from the $\gamma_w(\text{CH}_2)$ mode of a planar zig-zag chain. (It will be shown later that this mode is also active when the full crystalline arrangement is considered instead of just a single chain.) Rather, the 1303 cm^{-1} band appears only when the sample is melted, thus indicating that it is not correlated with the planar zig-zag form of the chain but with the disordered configuration. Furthermore, the assignment of the 1303 cm^{-1} band to $\gamma_w(\text{CH}_2)$ is not in agreement with the results of our examination of the spectra of a large variety of polymers which contain CH_2 but no CH_3 groups. We always find a band, which can be attributed to the $\gamma_w(\text{CH}_2)$ mode as a result of its parallel polarization in stretched specimens, in the region of $1350\text{--}1370\text{ cm}^{-1}$, never near 1300 cm^{-1} . It therefore seems fairly conclusive that the 1303 cm^{-1} band cannot be assigned to the $\gamma_w(\text{CH}_2)$ mode of a planar zig-zag chain.

Although, in the discussion of this assignment various authors^{7,8,22-27} have referred to the band at 1375 cm^{-1} ,

²² T. Simanouti, J. Chem. Phys. **17**, 734 (1949).

²³ T. Simanouti and S. Mizushima, J. Chem. Phys. **17**, 1102 (1949).

²⁴ Brown, Sheppard, and Simpson, Discussions Faraday Soc. **9**, 261 (1950).

²⁵ G. M. Barrow, J. Chem. Phys. **19**, 345 (1951).

²⁶ W. B. Person and G. C. Pimentel, J. Am. Chem. Soc. **75**, 532 (1953).

²⁷ H. Primas and H. H. Gunthard, Helv. Chim. Acta **36**, 1791 (1953).

²⁸ S. Krimm, Phys. Rev. **94**, 1426 (1954).

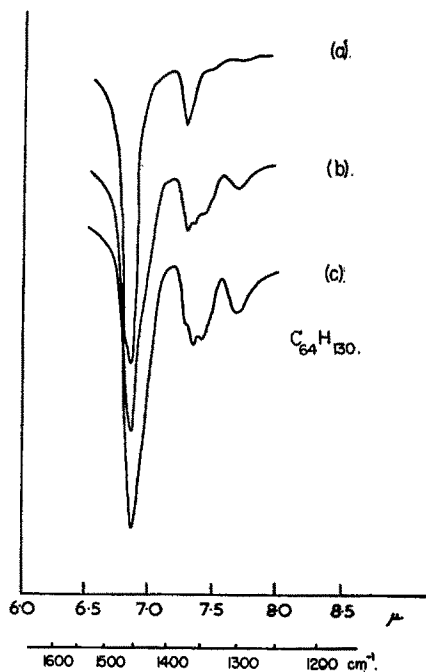


FIG. 4. Spectra of $n\text{-C}_{64}\text{H}_{130}$: (a) as a solid at room temperature, (b) as a liquid at the melting point, (c) about 15° above the melting point.

it has become evident¹² that this band is in reality a composite one. This has also been demonstrated especially clearly in the grating spectra of polyethylene,⁹ which show a structure similar to that in Fig. 4b. Peaks are found at 1375, 1369, and 1353 cm^{-1} . It has also been noted¹² that only the intensity of the 1375 cm^{-1} component (when measured in the liquid state) can be correlated with the CH_3 group content, and it was therefore suggested that the two components at 1369 and 1353 cm^{-1} are associated with CH_2 groups. Recent results²⁹ on terminally deuterated n -octane confirm that the 1369 cm^{-1} band is associated with a CH_2 group vibration and the results of our studies on single crystals of n -paraffins conclusively show that the $\gamma_w(\text{CH}_2)$ mode of a planar zig-zag chain is to be assigned to this band.

In Fig. 5 we show the spectra of the crystalline n -paraffins $\text{C}_{29}\text{H}_{60}$ and $\text{C}_{36}\text{H}_{74}$. By means of x-ray diffraction photographs, it was ascertained that the single crystals of $\text{C}_{29}\text{H}_{60}$ were in the usual orthorhombic form,³⁰ while those of $\text{C}_{36}\text{H}_{74}$ were in the monoclinic form.³¹ The spectra under (a) were obtained from a solidified melt, and therefore represent a random aggregate of small crystals. The spectra shown under (b) were obtained from single crystals of the respective paraffins. (These curves are separated for purposes of clarity.) For the $\text{C}_{29}\text{H}_{60}$, the solid and broken curves in Fig. 5(b) represent spectra with the electric vector

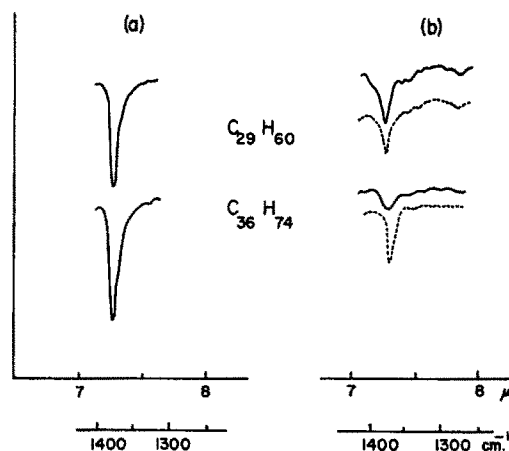


FIG. 5. Spectra of $n\text{-C}_{29}\text{H}_{60}$ and $n\text{-C}_{36}\text{H}_{74}$. (a) Sample solidified from the melt. (b) Single crystals: — radiation with electric vector polarized along a -axis, - - - radiation with electric vector polarized along b -axis.

polarized along the a and b crystal axes, respectively. In this case, therefore, the radiation was directed along the c -axis of the single crystal, and the electric vector was always perpendicular to the chain axis. In the case of the $\text{C}_{36}\text{H}_{74}$, although the complete structure is not yet known, it appears from preliminary results³² that the chain axis and the beam direction would make an angle of about 30° with each other in our crystal, and therefore the electric vector could have a component along the chain axis direction. The solid and broken curves in Fig. 5b represent in this case spectra with the electric vector polarized, respectively, along the crystal axis corresponding to the a and the b axes of the orthorhombic crystal.

The spectra in Fig. 5 show the following features. In the (a) spectra, corresponding to a random array of crystals, there is in both cases a strong peak at 1375 cm^{-1} with a weak shoulder near 1369 cm^{-1} . No band is found near 1353 cm^{-1} or 1303 cm^{-1} . The absence of the 1303 cm^{-1} band in crystalline paraffins, and therefore the rejection of its assignment to $\gamma_w(\text{CH}_2)$, has been noted earlier. In the (b) spectra, we find for $\text{C}_{29}\text{H}_{60}$ only a single unpolarized band at 1375 cm^{-1} , whereas for $\text{C}_{36}\text{H}_{74}$ we find two bands, one at 1375 cm^{-1} and the other at 1369 cm^{-1} , the latter being more intense than the former. These results indicate that the 1369 cm^{-1} band is unquestionably the only band in this region which is polarized along the chain axis, and therefore it is the only possible one which can be assigned to $\gamma_w(\text{CH}_2)$. The fact that this band is single, in accordance (as will be seen later) with the predicted spectrum based upon the full unit cell treatment, further substantiates this assignment.

The 1375 cm^{-1} band is undoubtedly the $\delta_s(\text{CH}_3)$ frequency. We can also readily see that, as a result of the chain inclination, the 1369 cm^{-1} band should be

²⁹ G. C. Pimentel and W. A. Klemperer, *J. Chem. Phys.* **23**, 376 (1955).

³⁰ A. Müller, *Proc. Roy Soc. (London)* **A120**, 437 (1928).

³¹ A. E. Smith, *J. Chem. Phys.* **21**, 2229 (1953).

³² A. E. Smith (private communication).

relatively stronger in the spectrum of the $C_{36}H_{74}$ single crystal. The clarification of the apparent polarization of these two bands in the $C_{36}H_{74}$ single crystal spectrum must await a complete structure determination of the monoclinic form of *n*-paraffins.

These results, as well as those on melted paraffins shown in Fig. 4, indicate that the presence in the polyethylene spectrum of bands at 1353 cm^{-1} and 1303 cm^{-1} must be attributed to the fact that the material is partially amorphous. In fact, it seems quite probable that at least the 1353 cm^{-1} band is to be assigned to a CH_2 wagging mode on chains in the amorphous regions. The shift in frequency from 1369 cm^{-1} is not unreasonable since the crystalline and amorphous chains represent different rotational isomers, which it is known^{33,34} can have different frequencies. The presence of the 1369 cm^{-1} band in the spectrum of the liquid is not too surprising, since undoubtedly there are segments of chains which are essentially in a planar zig-zag configuration. This is particularly true near the melting point.³⁵ The assignment of the 1353 cm^{-1} band to a CH_2 wagging mode is substantiated by the observed parallel polarization of this band in stretched polyethylene, in which it is known⁴ that the amorphous regions are partially oriented with the chain axis parallel to the stretching direction. The fact that the 1303 cm^{-1} band is also found to have parallel polarization indicates that it too may be assignable to a CH_2 wagging mode in the disordered phase. Another possibility is, however, the assignment to a CH_2 twisting mode on an amorphous chain. We will discuss this later. It may be noted that the grating spectra⁹ show almost no polarization for the 1375 cm^{-1} peak. The parallel polarization which has been attributed to this band⁸ undoubtedly arises from the overlapping parallel band at 1369 cm^{-1} . It is therefore unnecessary to postulate⁸ any special orientation of the CH_3 groups in the crystalline regions of polyethylene.

There are some puzzling features about the intensities of the bands we have been discussing upon which it is worth remarking. In the first place, it is not immediately obvious why the intensity of the CH_2 wagging mode should be so much weaker than that of the other CH_2 group frequencies. This is apparently related to the magnitude of the dipole moment change during the vibration, since the weak intensity persists in the liquid state. It presumably was the impression that this mode should not be a weak one which led to the erroneous assignment of the 1375 cm^{-1} band to this vibration. We have furthermore observed that the 1375 cm^{-1} band, assigned to $\delta_s(CH_3)$, is itself of variable intensity in the *n*-paraffins. An examination of the spectra of the *n*-paraffins $C_{27}H_{56}$, $C_{28}H_{58}$, $C_{29}H_{60}$, and $C_{30}H_{62}$ as poly-

crystalline solid films deposited from the melt reveals the following facts. In the $C_{27}H_{56}$ and $C_{29}H_{60}$ spectra there is a band at 1375 cm^{-1} with a shoulder on the low frequency side at 1369 cm^{-1} , as indicated in Fig. 4. The ratio of the optical density at the peak of the 1460 cm^{-1} $\delta(CH_2)$ band to that at the peak of the 1375 cm^{-1} band is in both cases about 3.3:1. In the $C_{28}H_{58}$ and $C_{30}H_{62}$ spectra, on the other hand, there is only a single symmetrical peak at 1369 cm^{-1} with no component at 1375 cm^{-1} , and the ratio of the 1460 to 1369 peak optical densities in this case is about 9.5:1. In other words, it would appear from these results that the $\delta_s(CH_3)$ vibration is considerably weaker in crystalline paraffins with an even number of carbon atoms in the chain than in those with an odd number. This has also been observed⁹ in the spectrum of $C_{18}H_{38}$. (This conclusion is not altered by making allowance in the prism spectra for the small $\delta_a(CH_3)$ contribution to the 1460 cm^{-1} band.) It would seem natural to try to account for these observations in terms of the differing interactions between adjacent CH_3 groups on neighboring chains in the crystal. This is not unreasonable, since the relative orientation of two such methyl groups is different for chains with an even number of carbon atoms than for chains with an odd number of carbon atoms. This explanation, however, cannot be the complete one, since in the case of the even paraffins $C_{36}H_{74}$ and $C_{64}H_{130}$ we also find the 1375 cm^{-1} band (see Figs. 4 and 5). The corresponding optical density ratio in both cases is about 9.5:1. A similar observation has been noted⁹ for $C_{48}H_{98}$. These results indicate that other factors, perhaps the crystal form, are also important. They point up the fact that our understanding of the factors determining the intensities of bands in the crystalline state is very incomplete. Indeed, until we understand more fully the factors affecting the intensity of the $\delta_s(CH_3)$ vibration, it may be desirable to exercise caution in interpreting the results of CH_3 group determinations in polyethylene obtained by means of the intensity of this band.^{12,36}

Finally, we will consider the assignments of the CH_3 group frequencies, which make a significant contribution to the polyethylene spectrum. In doing so we will assume that the CH_3 groups are randomly distributed in the sample. We may note that, contrary to this view, it has been suggested,³⁷ on the basis of the decrease in intensity of the 1375 cm^{-1} band on melting, that the CH_3 groups are in the crystalline regions of the polymer. This conclusion, however, may not be completely unambiguous when we consider that we yet do not understand the factors which determine the intensity of this fundamental. We have already noted that even the intensity variation of this band in crystals is not entirely simple. It is, therefore, justifiable to consider that we are dealing with independent

³³ J. K. Brown and N. Sheppard, *Discussions Faraday Soc.* **9**, 144 (1950).

³⁴ J. K. Brown and N. Sheppard, *Trans. Faraday Soc.* **50**, 535 (1954).

³⁵ B. E. Warren, *Phys. Rev.* **44**, 969 (1933).

³⁶ W. M. D. Bryant and R. C. Voter, *J. Am. Chem. Soc.* **75**, 6113 (1953).

³⁷ R. B. Richards, *J. Appl. Chem. (London)* **1**, 370 (1951).

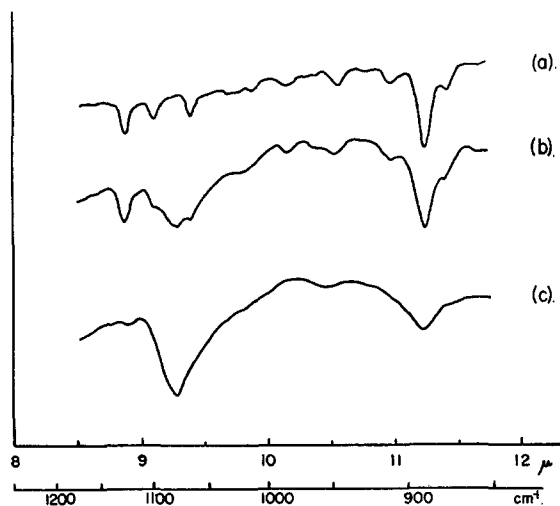


FIG. 6. Spectra of n - $C_{36}H_{74}$: (a) as polycrystalline solid at room temperature; (b) as a liquid at the melting point; (c) about 15° above the melting point.

CH_3 groups, such as would be present in the amorphous phase. We thus expect that to each CH_3 group can be associated 7 fundamentals: the symmetrical, $\nu_s(CH_3)$, and asymmetrical, $\nu_a(CH_3)$, internal hydrogen stretching modes; the symmetrical $\delta_s(CH_3)$, and asymmetrical, $\delta_a(CH_3)$, internal hydrogen bending modes; and the external rocking, $\gamma_r(CH_3)$, wagging $\gamma_w(CH_3)$, and twisting, $\gamma_t(CH_3)$, modes.

The $\nu_s(CH_3)$ and $\nu_a(CH_3)$ frequencies have been established by Fox and Martin⁶ at 2872 and 2962 cm^{-1} respectively. In our polyethylene spectrum we only find a shoulder at 2959 cm^{-1} , which is presumably the $\nu_a(CH_3)$ mode. The $\nu_s(CH_3)$ mode is less intense, and we do not detect it in our prism spectra. Both bands have been found, however, in grating spectra,⁹ where they occur at 2874 cm^{-1} and 2959 cm^{-1} respectively. The two internal bending modes are also well established, the $\delta_s(CH_3)$ mode as we have seen, at 1375 cm^{-1} and the $\delta_a(CH_3)$ mode^{18,19} at about 1460 cm^{-1} .

Of the external deformation frequencies of the CH_3 group, we expect that the twisting mode will be very

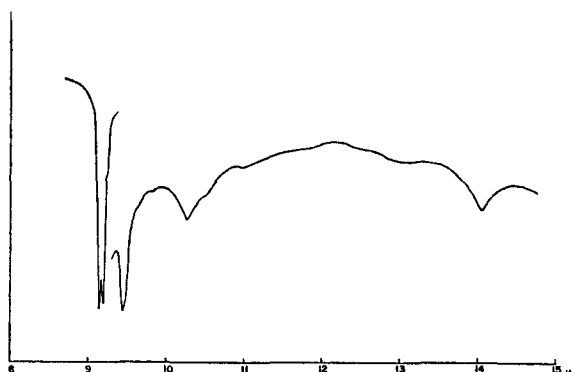


FIG. 7. Spectrum of deuteroparaffin (approximate composition $C_{100}D_{202}$).

weak, if not inactive. We have assumed that it is not detectable in the spectrum. We will consider now the assignment of the rocking and wagging modes. In terms of the plane of the carbon chain adjoining a CH_3 group, we term the rocking mode as the deformation taking place perpendicular to this plane and the wagging mode as the one which occurs in the plane. We will consider first the location of these two frequencies from an examination of the spectra of the n -paraffins. These frequencies are expected to occur in the range between 800 cm^{-1} to 1200 cm^{-1} , which is shown in Fig. 6 for the n -paraffin $C_{36}H_{74}$. Curve (a) is the spectrum of the polycrystalline solid at room temperature, curve (b) is of the material as a liquid at its melting point, and curve (c) is of the material at about 15° above its melting point. The prominent band at 890 cm^{-1} , found in polyethylene at 888 cm^{-1} , can immediately be assigned to a CH_3 frequency for the following reasons: (1) In our spectrum of the deuteroparaffin $C_{100}D_{202}$, shown in Fig. 7, this band has shifted to 711 cm^{-1} , just about the amount expected for a CH_3 frequency; (2) A correlation is found¹² between the intensity of this band and the CH_3 group content of the hydrocarbon; (3) The band is relatively unchanged, compared to neighboring bands, on passing from the solid to the liquid state, as can be seen from Fig. 6. The significance of the latter point is that skeletal frequencies of the planar zig-zag carbon chain, which are also expected to occur in this region, are most likely to be affected by a change of state, since the skeletal configuration is thereby altered. This has also been observed in the Raman spectra of n -paraffins.³⁸ We therefore assign as CH_3 frequencies those bands which are least affected by a change in state. (It might be noted that although the skeletal frequencies are predicted to be inactive for an infinite chain, they can show activity in finite chains.¹⁹)

As to the other CH_3 frequency, the above argument about band constancy with change in state, and knowledge³⁹ of the expected region for such modes, leads us to assign the 1128 cm^{-1} band in $C_{36}H_{74}$ to a CH_3 frequency. In all of the spectra of n -paraffins which we have examined, a band is found within 2 cm^{-1} of this value. There also appears to be some support for this assignment from studies on terminally deuterated n -octane.²⁹ Thus we assign the CH_3 modes to bands at about 890 cm^{-1} and 1128 cm^{-1} . It has been claimed¹⁰ that the absorption band in polyethylene at 543 cm^{-1} is correlated with the CH_3 content of the sample. We have been unable to detect this band in suitably thick samples of the n -paraffin $C_{36}H_{74}$, and therefore think that it is unlikely that it represents a CH_3 frequency. Furthermore, it does not seem to be in an appropriate region, even for a torsional mode.⁴⁰

The assignment of the two CH_3 frequencies at 890

³⁸ S. Mizushima and T. Simanouti, *J. Am. Chem. Soc.* **71**, 1320 (1949).

³⁹ R. S. Rasmussen, *J. Chem. Phys.* **16**, 712 (1948).

⁴⁰ K. S. Pitzer, *J. Chem. Phys.* **12**, 310 (1944).

TABLE IV. Character table, number of normal modes, and selection rules for crystalline polyethylene.

D_{2h}	E	$C_2^*(a)$	$C_2^*(b)$	$C_2^*(c)$	i	$\sigma_a(bc)$	$\sigma_b(ac)$	$\sigma(ab)$	n_i	T	T'	R'	n_i'	IR	R
A_g	1	1	1	1	1	1	1	1	6	0	2	1	3	f ^a	p
B_{1g}	1	1	-1	-1	1	1	-1	-1	3	0	1	2	0	f	d
B_{2g}	1	-1	1	-1	1	-1	1	-1	3	0	1	2	0	f	d
B_{3g}	1	-1	-1	1	1	-1	-1	1	6	0	2	1	3	f	d
A_u	1	1	1	1	-1	-1	-1	-1	3	0	1	2	0	f	f
B_{1u}	1	1	-1	-1	-1	-1	1	1	6	T_a	1	1	3	a	f
B_{2u}	1	-1	1	-1	-1	1	-1	1	6	T_b	1	1	3	a	f
B_{3u}	1	-1	-1	1	-1	1	1	-1	3	T_c	0	2	0	a	f

^a See reference a in Table II.

and 1128 cm^{-1} to the rocking and wagging motions has now to be considered. If the disposition of CH_3 groups in oriented polyethylene were known, we could use the observed dichroism of these bands to assist in making the assignment. Unfortunately, we cannot assume this knowledge. Similarly, complete knowledge of the structure of the monoclinic form of paraffins would be helpful, since we observe that in a single crystal of orthorhombic $\text{C}_{29}\text{H}_{60}$ the 890 cm^{-1} band appears with about the same relative intensity as in the polycrystalline solid, whereas in a single crystal of monoclinic $\text{C}_{36}\text{H}_{74}$ (which is viewed at an angle to the chain axis) this band is considerably weaker. It seems likely, nevertheless, just from consideration of the restoring forces involved in the two motions, that we should assign the 890 cm^{-1} band to the out-of-plane mode, $\gamma_r(\text{CH}_3)$, and the 1128 cm^{-1} band to the in-plane mode, $\gamma_w(\text{CH}_3)$. This would be in agreement with conclusions reached from a study of short chain n -paraffins.⁴¹

D. Predicted Spectrum of Crystalline Polyethylene

We have already noted that although the main features of the spectrum of polyethylene can be predicted from consideration of a single chain, a complete interpretation must take account of the actual crystal structure of the material. We will consider now this more complete analysis. A comparison of the two treatments should serve as a guide to what may be expected in the spectra of other polymers.

The procedure for the group theory analysis is similar to that discussed for the case of the single chain.^{2,15,16} A cross section of the unit cell of crystalline polyethylene,² perpendicular to the c -axis (chain axis) and bisecting a C—C bond, is shown in Fig. 8. The CH_2 groups lie in planes parallel to the ab plane, and in Fig. 8 they are indicated as being above (solid line) and below (dashed line) the plane of the section. The structure belongs to the space group D_{2h}^{16} , which has a factor group isomorphous with the point group D_{2h} . The character table, number of normal modes, and selection rules are given in Table IV, the symbols having the same meaning as those in Table II. As in the case

of the single chain, approximate normal modes can be derived from the symmetry species. These are shown in Figs. 9, 10, and 11, and they are classified in Table V. The motions indicated by arrows take place parallel to the ab plane, those indicated by plus and minus signs take place perpendicular to this plane. As can be seen from Table IV, B_{1u} and B_{2u} modes will be polarized perpendicular to the chain axis while B_{3u} modes will be polarized parallel to this axis.

The factor group analysis predicts 12 infrared active fundamentals in the spectrum of crystalline polyethylene. In comparing these with the fundamentals predicted for a single chain, the following differences are to be noted: (1) All of the single chain modes which exhibited polarization perpendicular to the chain axis are split into two components in the crystal. Of these, as will be readily seen from Figs. 10 and 11, the B_{1u} modes are polarized along the a -axis of the crystal and the B_{2u} modes are polarized along the b -axis. The extent of the splitting, and therefore the possibility of observing the two components in the spectrum, cannot of course be determined from symmetry considerations alone. It depends upon the nature and magnitude of the forces which are involved. (2) The parallel polarized $\gamma_w(\text{CH}_2)$ mode is not split in the infrared spectrum of the crystal. (3) A CH_2 twisting mode, inactive in a single chain, is predicted to be active in the crystal.

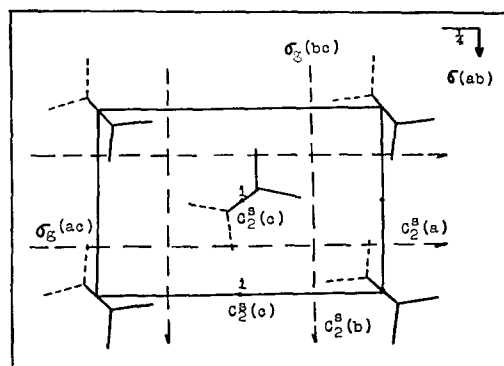


FIG. 8. Cross section of unit cell of polyethylene perpendicular to c -axis and bisecting C—C bond. C_2^s —screw diad, σ —mirror plane, σ_g —glide plane (σ_g is a diagonal glide plane, with translation $(b+c)/2$), i —center of inversion. CH_2 groups are indicated as above (solid V) and below (dashed V) the plane.

⁴¹ Brown, Sheppard, and Simpson, Trans. Roy. Soc. (London) A247, 35 (1954).

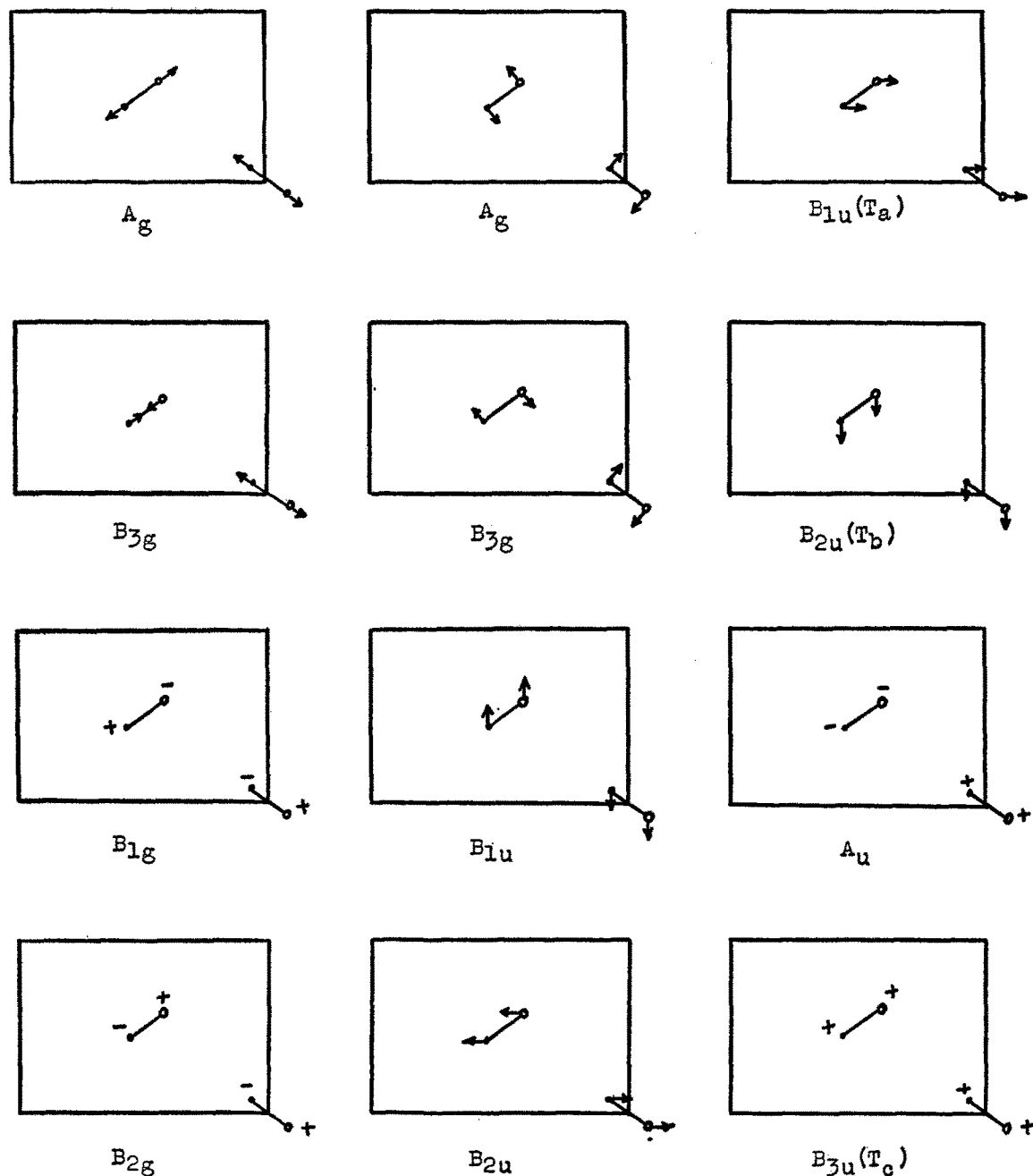


FIG. 9. Normal modes of polyethylene unit cell: translations (T) and translatory-type vibrations (T').

(4) Two lattice modes, B_{1u} and B_{2u} , are predicted to be active in the crystal. Since these modes are expected to be at a very low frequency, they have probably not been observed in the range covered by the present work.

Thus, as a result of the transition from the single chain to the crystal, we find a splitting of some of the fundamental frequencies and activity predicted for some previously inactive frequencies. The interactions which lead to a splitting of frequencies in the crystal will not drastically change the positions of the funda-

mental frequencies.⁴² It is for this reason that we can carry through a satisfactory assignment based on the analysis of a single chain. The activity in the crystal spectrum of modes which are inactive in a single chain is a result of the change in the symmetry of the potential field of the molecule. Although such activity can result, it will rarely, if ever, happen that a fundamental which is predicted to be active in the single chain

⁴² D. F. Hornig, Discussions Faraday Soc. 9, 115 (1950).

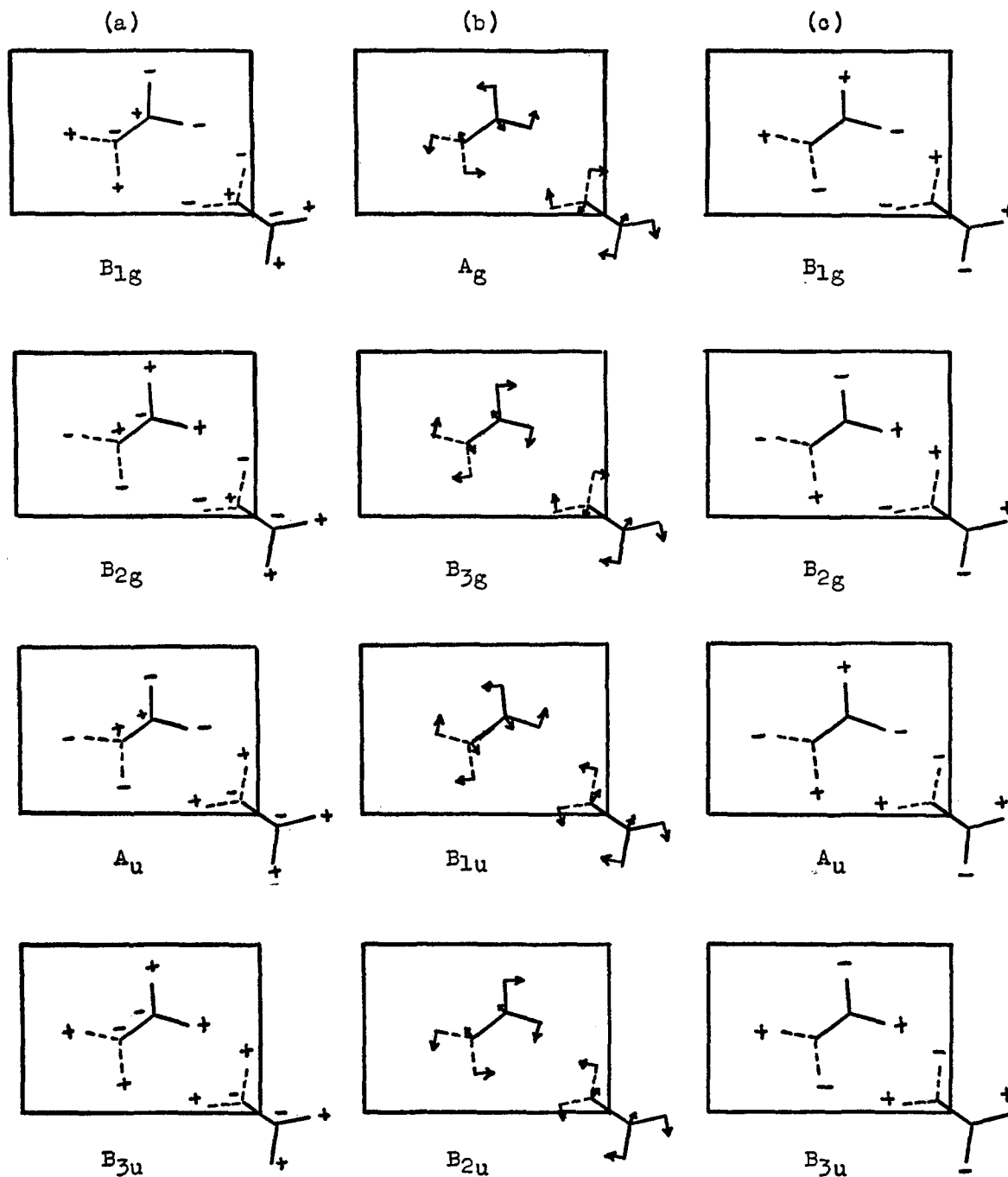


FIG. 10. Normal modes of polyethylene unit cell: rotatory-type vibrations (R').

becomes inactive in the crystal.¹ This further justifies the use of the single chain analysis.

E. Assignments of Frequencies in Crystalline Polyethylene

We wish to consider now those features of the spectrum of polyethylene which cannot be accounted for on the basis of a single chain analysis. Most prominent among these is the evident doubling in the band at

about 725 cm^{-1} , the $\gamma_r(\text{CH}_2)$ mode, and the similar effect obtained at 1460 cm^{-1} in the $\delta(\text{CH}_2)$ mode.^{43,44} We will show that these effects are associated with the crystalline regions of polyethylene, and follow from the results of the factor group analysis.

It has been suggested^{43,44} that the splitting of the

⁴³ R. S. Stein and G. B. B. M. Sutherland, *J. Chem. Phys.* **21**, 370 (1953).

⁴⁴ R. S. Stein and G. B. B. M. Sutherland, *J. Chem. Phys.* **22**, 1993 (1954).

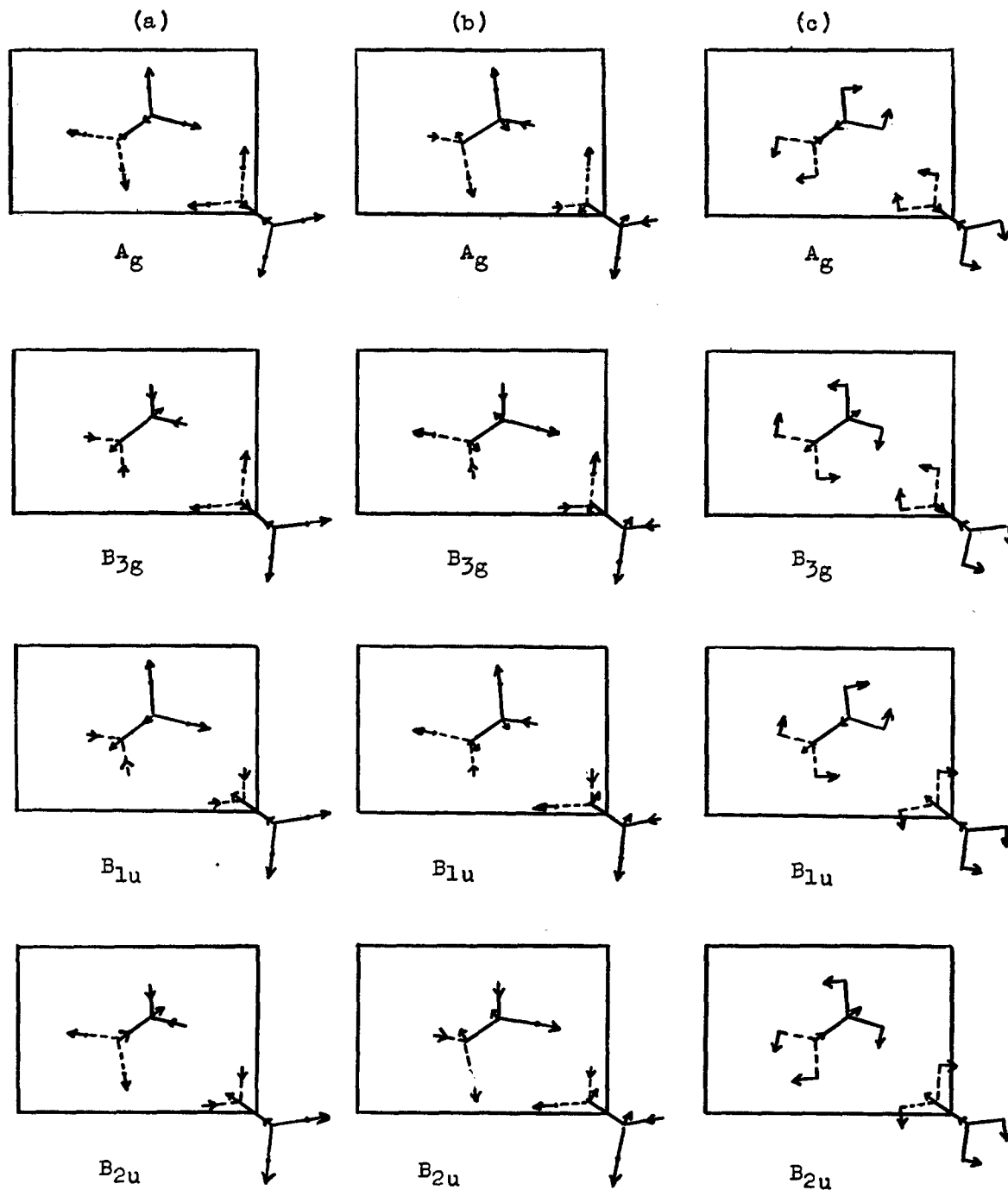


FIG. 11. Normal modes of polyethylene unit cell: internal vibrations (ν_i').

$\gamma_r(\text{CH}_2)$ mode into two components at 721 cm^{-1} and 731 cm^{-1} arises from interactions between adjacent molecules in the crystalline phase which results in both an in-phase and an out-of-phase rocking mode. This is also predicted from the group theory analysis, which shows that the B_{1u} mode will be polarized along the a -axis of the crystal and the B_{2u} mode along the b -axis. An obvious test of this hypothesis would be to examine the polarization of these components in a single crystal,

with radiation directed along the c -axis. This has been done,⁴ and we indeed find such specific polarization for the components, *viz.*, the 721 cm^{-1} band polarized along the b -axis and the 731 cm^{-1} band polarized along the a -axis. The crystal which was examined,⁴ *viz.*, $\text{C}_{36}\text{H}_{74}$, was in the monoclinic form. We have also examined a single crystal of orthorhombic $\text{C}_{29}\text{H}_{60}$ and find exactly the same results.

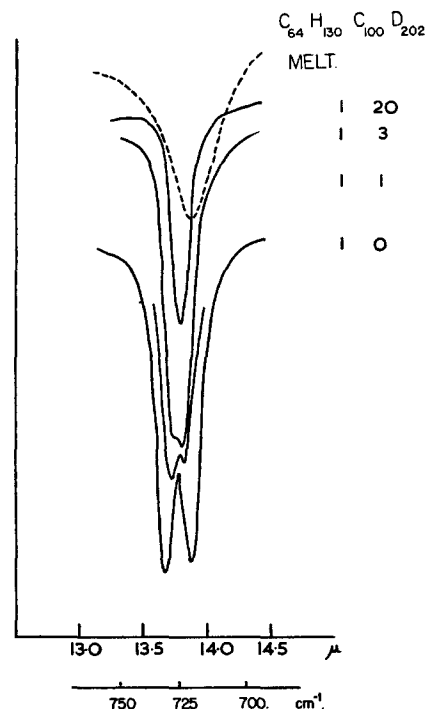
Further evidence that this splitting arises from an

TABLE V. Classification of normal modes of crystalline polyethylene.

Type of motion	Species	Illustrated in Figure	Column
Translations	B_{1u}, B_{2u}, B_{3u}	9	c
Lattice modes			
(a) Translatory	B_{1u}, B_{2u}	9	b
	A_u	9	c
(b) Rotatory	A_g, B_{3g}	9	b
Skeletal vibrations	$A_g, B_{3g}, B_{1g}, B_{2g}$	9	a
External CH_2 modes			
(a) CH_2 wagging	$B_{1g}, B_{2g}, A_u, B_{3u}$	10	a
(b) CH_2 rocking	$A_g, B_{3g}, B_{1u}, B_{2u}$	10	b
(c) CH_2 twisting	$B_{1g}, B_{2g}, A_u, B_{2u}$	10	c
Internal CH_2 modes			
(a) CH_2 sym. stretching	$A_g, B_{3g}, B_{1u}, B_{2u}$	11	a
(b) CH_2 asym. stretching	$A_g, B_{3g}, B_{1u}, B_{2u}$	11	b
(c) CH_2 bending	$A_g, B_{3g}, B_{1u}, B_{2u}$	11	c

interaction between chains in the crystal is provided by a study of solid solutions of a paraffin in a deuteroparaffin.²⁸ In Fig. 12 are shown the spectra in the 725 cm^{-1} region of a series of solid solutions of $\text{C}_{64}\text{H}_{130}$ in a deuteroparaffin whose formula is approximately $\text{C}_{100}\text{D}_{202}$. The broken curve represents a melt of pure $\text{C}_{64}\text{H}_{120}$. A number of observations can be made. (1) In a dilute enough solution, such as 1:20 by weight, the $\gamma_r(\text{CH}_2)$ mode is a single band. This would be in support of the hypothesis that the splitting results from an interaction between molecules, since this interaction would diminish when neighboring chains in the unit cell are no longer equivalent. We are of course assuming (which does not seem unreasonable) that the solid solution has the same crystal structure as the normal paraffin. (2) The frequency of this single band in the solid solution is 725 cm^{-1} , which is different from the value of 721 cm^{-1} for this frequency in the liquid state. This probably represents the difference between the $\gamma_r(\text{CH}_2)$ frequency in the planar zig-zag chain and in the rotational isomers of it which are found in the liquid. This is analogous to our assignment of the 1353 cm^{-1} band to a $\gamma_w(\text{CH}_2)$ mode on a chain in the disordered phase. (3) It appears that the magnitude of the splitting is a function of the concentration. While in the pure paraffin the components are at 721 and 731 cm^{-1} , in the 1:1 solution they are at 722 and 728 cm^{-1} and in the 1:3 solution they occur at 724 and 727 cm^{-1} . This is difficult to understand on the basis of an interaction only between adjacent molecules,^{43,45} and would seem to imply a more extended interaction between molecules in the crystal.

On the basis of a van der Waals interaction between the hydrogen atoms on adjacent molecules, an attempt has been made⁴⁵ to account for the magnitude of the observed splitting at 725 cm^{-1} and the polarization of the components. There are two substantial discrepancies between the predictions of this theory and observation. (1) The theory predicts that the two components of the doublet should both be at a higher frequency


 FIG. 12. Spectra of solid solutions of $\text{C}_{64}\text{H}_{130}$ in $\text{C}_{100}\text{D}_{202}$. Numbers give weight ratio.

than the mode in an isolated molecule. As we have seen from the study of the solid solutions this does not seem to be the case. The single frequency of the unperturbed molecule is between the values of the doublet in the perturbed case. There is a possibility that the value in the unperturbed molecule is nevertheless different from that in an isolated molecule, as a result of the effect of the crystalline field.⁴² Usually however, the frequency in the solid is lower than in the isolated molecule. The discrepancy requires clarification. (2) A more serious difficulty is in the predicted polarization of the two components. The theory predicts that for the $721\text{--}731 \text{ cm}^{-1}$ doublet the high frequency component should be polarized along the a -axis (as is found), while for the $\delta(\text{CH}_2)$ mode doublet, which we find at 1461 and 1470 cm^{-1} , the high frequency component should be polarized along the b -axis. This is not the case for monoclinic $\text{C}_{36}\text{H}_{74}$,⁴ as can be seen from Fig. 13. We have recently obtained the spectrum of orthorhombic $\text{C}_{29}\text{H}_{60}$ and find the same situation here, i.e., both at 725 cm^{-1} and 1460 cm^{-1} the high frequency component is polarized along the a -axis of the crystal. It is difficult to reconcile this with the proposed theory. In fact, the progressive splitting of the 725 cm^{-1} band in solid solutions of a deuteroparaffin makes it difficult to believe that the interaction could be the simple one proposed between adjacent molecules.

We have also observed that in a single crystal of $\text{C}_{36}\text{H}_{74}$ the CH_2 stretching frequencies are split, in accordance with the group theory analysis. For the $\nu_a(\text{CH}_2)$ mode, the high frequency component at 2924

⁴⁵ R. S. Stein, J. Chem. Phys. 23, 734 (1955).

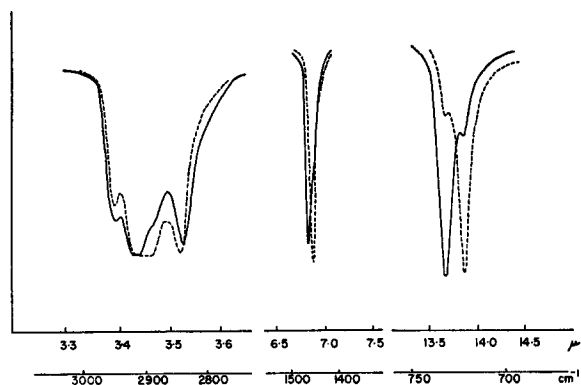


FIG. 13. Spectrum of single crystal of monoclinic n - $C_{36}H_{74}$. — radiation with electric vector polarized along a -axis, - - - radiation with electric vector polarized along b -axis.

cm^{-1} is polarized along the a -axis, the low frequency component at 2899 cm^{-1} being polarized along the b -axis. For the $\nu_s(CH_2)$ mode, the high frequency component at 2857 cm^{-1} is polarized along the b -axis, the low frequency component at 2850 cm^{-1} being polarized along the a -axis. This reversal of polarization of the high frequency component for $\nu_s(CH_2)$ is difficult to account for on the simple interaction theory.⁴⁵

We may consider now the $\gamma_t(CH_2)$ mode, which is predicted to be active in crystalline polyethylene and to exhibit parallel polarization. Studies on n -paraffins⁴¹ have suggested that this frequency should occur at about 1300 cm^{-1} . It should be infrared active only for the even-numbered carbon chains. In our examination of the n -paraffins $C_{26}H_{58}$, $C_{30}H_{62}$, $C_{36}H_{74}$, and $C_{64}H_{130}$ we find a very weak band at about 1295 cm^{-1} which is common to all of these spectra. It is possible that this is the $\gamma_t(CH_2)$ mode of the crystalline phase. The relative intensity of this band diminishes with increasing chain length, and we would therefore expect that it is probably not observable in the spectrum of polyethylene. The only band found in this region in polyethylene is that at 1303 cm^{-1} . As we have already seen, this band arises from the amorphous component and therefore cannot be assigned to $\gamma_t(CH_2)$ in the crystal. The question arises, however, as to whether this band might be a $\gamma_t(CH_2)$ mode of an amorphous chain. The evidence does not strongly favor this possibility. In the deuteroparaffin spectrum, a band of similar shape and intensity, which also increases in intensity upon melting, is found at about 975 cm^{-1} . The ratio of frequencies, 1.34, does not seem appropriate to a pure twisting mode. In addition, the observed parallel polarization of the 1303 cm^{-1} band in oriented samples is not readily interpretable in terms of a $\gamma_t(CH_2)$ mode. If the CH_2 groups in the amorphous regions are considered as independent units of symmetry C_{2v} , then a twisting mode is forbidden in the infrared. If we take a CH_2-CH_2 unit, of symmetry C_{2h} , as the independent group, then a $\gamma_t(CH_2)$ mode is permitted but its polarization is predicted to be perpendicular to the carbon chain axis, as in the case of the crystalline n -paraffins.⁴¹ It is therefore

probable that if this frequency occurs in a partially oriented amorphous chain it would show perpendicular polarization. These two facts lead us to assign the 1303 cm^{-1} band most probably to a $\gamma_w(CH_2)$ mode of an amorphous chain.

The previous discussion of the spectrum of crystalline polyethylene has been based on the assumption that the chain is infinitely long. This is, of course, not actually the case, and as a result it is expected¹⁹ that the skeletal vibrations, which are inactive for an infinite chain, can now become active. We wish to consider briefly the possible occurrence of such bands in the spectra of polyethylene and of paraffins. A calculation¹ of the values of the two skeletal frequencies A_g and B_{1g} of a single chain (see Fig. 3(a)) indicates that they should be found at about 1070 cm^{-1} and 1137 cm^{-1} , respectively. An examination of the spectra of the various n -paraffins studied in this work shows a fairly uniform grouping of bands at about 1065 cm^{-1} and 1100 cm^{-1} . (No such bands were found in $C_{64}H_{130}$, which probably indicates that this molecule is behaving very much like an infinite chain.) An example of these bands is shown in Fig. 6a, where they occur at 1068 cm^{-1} and 1100 cm^{-1} . It seems probable that they are to be associated with the above-mentioned skeletal modes. This is substantiated by their behavior on melting the sample. As Fig. 6b and c show, they disappear and are replaced by a single stronger band at 1078 cm^{-1} . The latter probably represents the mean value of the skeletal stretching frequency in the amorphous chains. We would therefore assign the weak band at 1080 cm^{-1} in polyethylene to a skeletal frequency in the amorphous phase.

As for the remaining weak bands in the spectrum of polyethylene, little can be said beyond the statement that they probably are combinations and overtones. In one case this seems to be fairly well substantiated, *viz.*, for the band at 2010 cm^{-1} in polyethylene. This band is found at 2015 cm^{-1} in polycrystalline $C_{36}H_{74}$, and is very weak. Its relative intensity increases, however, when the sample is melted, much as does the 1303 cm^{-1} band. In fact the band at 2010 cm^{-1} could be a combination involving the 1303 and 721 cm^{-1} bands. The other cases are not so clear-cut, and may very well arise from combinations between active and inactive modes. The very weak band at 965 cm^{-1} in polyethylene is also found in the spectra of the n -paraffins, but we have been unable to assign it.

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