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ON THE CONSTRUCTION OF THE LATTICE OF SP PARTITIONS

by

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In this communication I present a method for constructing the lattice of SP partitions for any given finite-state machine via a state-pair graph constructed directly from the machine's transition function and a homomorphism which reduces a set of subgraphs to the full SP partition lattice for the machine. The method is a graphical extension of the two-state generator procedure of Hartmanis and Stearns.

The initial steps in developing this approach were made at a time when I was trying to innovate alternative methods for programming the SP partition lattice calculations on a computer. The elegant graphical structure of some of the examples I tried immediately attracted my attention and moved me to develop the theory to the point presented in this paper; I include one such example in this report.

As it turned out I elected to use another method in the actual computer program—a method I felt would be more thrifty regarding memory requirements; details on the program and its method are to appear in another report—Piatkowski, Thomas F. , Computer Programs Dealing with Finite-State Machines: Part II—soon to be published by the Department of Electrical Engineering, Systems Engineering Laboratory, The University of Michigan.

In the presentation that follows I will assume that the reader is familiar with the theory and notation of Hartmanis and Stearns.

1) In the manner of Hartmanis and Stearns let $M = \langle S, I, \delta \rangle$ be a finite-state machine.

2) Construct directed graph $G = \langle V, E \rangle$ from M such that

a) V , the set of vertices = $\{\{s_i, s_j\} \mid s_i, s_j \in S, s_i \neq s_j\}$

i. e. V is the set of unordered pairs of distinct states in S ;

and

b) E , the set of edges = $\left\{ (u_k, u_\ell) \mid \begin{array}{l} u_k, u_\ell \in V, \\ u_k \neq u_\ell \\ \exists x \in I, \text{ s.t. } u_\ell = \delta(u_k, x) \end{array} \right\}$

where $u_k = \{s_i, s_j\} \Rightarrow \delta(u_k, x) = \{\delta(s_i, x), \delta(s_j, x)\}$

i. e. E is the set of state pair transitions induced by M on V

with no distinction made by input and with state-pair-merging

transitions and self-loops ignored.

3) Define \mathcal{H} as the following set of subsets of V :

$$\mathcal{H} = \{H \mid H \subseteq V, R_G(H) = H\}$$

where $R_G(H) = \{\delta(v, x) \mid v \in H, x \in I^*, \delta(v, x) \in V\}$,

i. e. $R_G(H)$ is the set of all vertices in G reachable from H .

Note in particular that $\phi, V \in \mathcal{H}$.

4) A partial ordering is induced on \mathcal{H} by set inclusion; i. e.

$H_m < H_n$ iff $H_m \subset H_n$. Furthermore, two binary operations $+$ and \cdot can be defined on \mathcal{H} corresponding to set union and intersection; i. e.

$$H_m + H_n = H_m \cup H_n$$

and

$$H_m \cdot H_n = H_m \cap H_n.$$

These operations are closed in \mathcal{H} ; this can be shown as follows:

First we show that $+$ is closed in \mathcal{H} . For all $G = \langle V, E \rangle$ and any $H_m, H_n \subseteq V$ it is always true that

$$R_G(H_m \cup H_n) = R_G(H_m) \cup R_G(H_n).$$

In addition, $H_k \in \mathcal{H} \implies R_G(H_k) = H_k$

and thus

$$\forall H_m, H_n \in \mathcal{H}, R_G(H_m \cup H_n) = R_G(H_m) \cup R_G(H_n) = H_m \cup H_n.$$

In other words, $\forall H_m, H_n \in \mathcal{H}, H_m \cup H_n \in \mathcal{H}$.

To show that \cdot is closed in \mathcal{H} we note that for all $G = \langle V, E \rangle$ and any $H_m, H_n \subseteq V$ it is always true that $R_G(H_m \cap H_n) \supseteq H_m \cap H_n$ via the definition of R_G . In addition

$$\forall H_m, H_n \in \mathcal{H}, R_G(H_m \cap H_n) \subseteq R_G(H_m), R_G(H_n) = H_m, H_n$$

which implies that

$$\forall H_m, H_n \in \mathcal{H}, R_G(H_m \cap H_n) \subseteq H_m \cap H_n.$$

Thus $\forall H_m, H_n \in \mathcal{H}, R_G(H_m \cap H_n) = H_m \cap H_n$

or in other words, $H_m \cap H_n \in \mathcal{H}$.

5) Claim: $\langle \mathcal{H}, +, \cdot \rangle$ is a lattice.

Proof: $\langle \mathcal{H}, +, \cdot \rangle$ satisfies one of the several equivalent definitions for a lattice, namely \mathcal{H} is non-empty, and $+$ and \cdot are binary operations satisfying the following postulates for any $H_m, H_n, H_p \in \mathcal{H}$:

$$(i) \quad H_m \cdot H_m = H_m \quad ; \quad H_m + H_m = H_m$$

$$(ii) \quad H_m \cdot H_n = H_n \cdot H_m \quad ; \quad H_m + H_n = H_n + H_m$$

$$(iii) \quad H_m \cdot (H_n \cdot H_p) = (H_m \cdot H_n) \cdot H_p ; \quad H_m + (H_n + H_p) = (H_m + H_n) + H_p$$

$$(iv) \quad H_m \cdot (H_m + H_n) = H_m \quad ; \quad H_m + (H_m \cdot H_n) = H_m$$

Q. E. D.

6) Let $\langle P, +, \cdot \rangle$ be the lattice of partitions on S with substitution property relative to M .

7) Claim: $\langle P, + \rangle$ is a homomorphic image of $\langle \mathcal{H}, + \rangle$.

[Note: It is not the claim that $\langle P, +, \cdot \rangle$ is a homomorphic image of $\langle \mathcal{H}, +, \cdot \rangle$].

Proof: For each $H \in \mathcal{H}$, define

$$\pi(H) = \sum_{\{s_i, s_j\} \in H} \tau_{ij}$$

where \sum denotes the usual partition summation and where τ_{ij} denotes the

(n-1)-block partition on S in which each block is a singleton except for one doubleton block that identifies s_i and s_j . We must show

(a) that $\pi: \mathcal{H} \xrightarrow{\text{onto}} \mathcal{P}$ and (b) that π preserves the structure of the + operator.

a) $\forall H \in \mathcal{H}$, $\pi(H)$ has SP; this can be shown as follows:

$$s_i \equiv s_j(\pi(H)) \iff$$

\exists a string of elements in H which can be arranged in the following pattern

$$\{s_i, s_{a_1}\} \{s_{a_1}, s_{a_2}\} \{s_{a_2}, s_{a_3}\} \dots \{s_{a_k}, s_j\}$$

but

$$\forall x \in I^*, \{\delta(s_i, x), \delta(s_{a_1}, x)\} \{\delta(s_{a_1}, x), \delta(s_{a_2}, x)\} \dots \{\delta(s_{a_k}, x), \delta(s_j, x)\}$$

(with any identical pairs removed) is also a string of elements in H (since $R_G(H) = H$) with the above pattern; thus

$\delta(s_i, x) \equiv \delta(s_j, x)(\pi(H))$ for $\forall x \in I^*$. In other words, $\pi(H)$ has SP.

Thus $\forall H \in \mathcal{H}$, $\pi(H) \in \mathcal{P}$ which means that π at least maps \mathcal{H} into \mathcal{P} .

Furthermore for $\forall p \in \mathcal{P} \exists H \in \mathcal{H} . \ni . \pi(H) = p$; namely,

$$H = \{\{s_i, s_j\} \mid s_i \neq s_j, s_i \equiv s_j(p)\}.$$

That $H \in \mathcal{H}$ follows from the fact that $H \subseteq V$ and

$$R_G(H) = \left\{ \left. \begin{array}{l} \{\delta(s_i, x), \delta(s_j, x)\} \\ \{s_i, s_j\} \in H; \\ x \in I^*; \\ \{\delta(s_i, x), \delta(s_j, x)\} \in V \end{array} \right\} = H$$

since $\{s_i, s_j\} \in H \implies s_i \equiv s_j(p)$

$\implies \forall x \in I^*, \delta(s_i, x) \equiv \delta(s_j, x)(p)$

$\implies \{\delta(s_i, x), \delta(s_j, x)\} \in H$ if $\delta(s_i, x) \neq \delta(s_j, x)$.

Thus every $p \in P$ has a pre-image under π in \mathcal{H} ; i. e.

$\pi: \mathcal{H} \xrightarrow{\text{onto}} P$.

b) $\forall H_m, H_n \in \mathcal{H},$

$$\begin{aligned} \pi(H_m + H_n) &= \sum_{\{s_i, s_j\} \in H_m \cup H_n} \tau_{ij} = \sum_{\{s_i, s_j\} \in H_m} \tau_{ij} + \sum_{\{s_i, s_j\} \in H_n} \tau_{ij} \\ &= \pi(H_m) + \pi(H_n); \end{aligned}$$

thus π preserves the structure of the $+$ operator.

Q. E. D.

8) Observation: The π mapping will not, in general, yield a homomorphism from $\langle \mathcal{H}, +, \cdot \rangle$ to $\langle P, +, \cdot \rangle$ since the structure of the \cdot operator is not preserved. (See section 12 for such an example.)

9) Claim: $\forall H_m, H_n \in \mathcal{H}, H_m > H_n \Rightarrow \pi(H_m) \geq \pi(H_n).$

Proof: $H_m > H_n \Rightarrow H_m \supset H_n \Rightarrow H_m = H_n \cup (H_m - H_n)$

$$\begin{aligned} \Rightarrow \pi(H_m) &= \sum_{\{s_i, s_j\} \in H_n \cup (H_m - H_n)} \tau_{ij} \\ &= \sum_{\{s_i, s_j\} \in H_n} \tau_{ij} + \sum_{\{s_i, s_j\} \in (H_m - H_n)} \tau_{ij} \\ &= \pi(H_n) + \sum_{\{s_i, s_j\} \in (H_m - H_n)} \tau_{ij} \\ \Rightarrow \pi(H_m) &\geq \pi(H_n) \end{aligned}$$

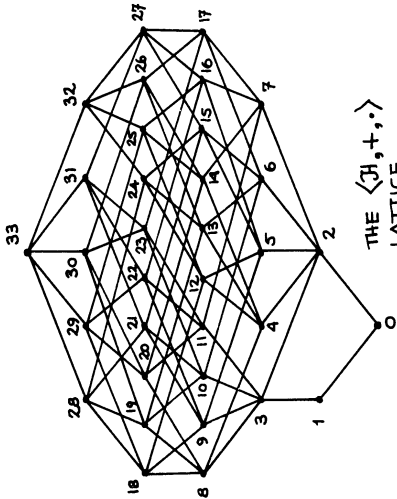
Q. E. D.

- 10) Observation: $\langle P, +, \cdot \rangle$ can be constructed directly from $\langle \mathcal{A}, +, \cdot \rangle$ using the π mapping. First of all $\pi(\mathcal{A}) = P$ and the structure of $\langle P, +, \cdot \rangle$ can certainly be deduced from P itself; however the fact that π also preserves the structure of the $+$ operator and some aspects of the $>$ relation can be used to good effect in determining the structure of $\langle P, +, \cdot \rangle$. For example: $\pi(\phi) = 0$, $\pi(V) = I$, and every lattice atom in P must have a pre-image under π in the lattice atoms of \mathcal{A} , etc. . The last mentioned observation follows from the fact that for every atom $p \in P$ $\exists H \in \mathcal{A} . \exists . \pi(H) = p$; every atom $H' \leq H$ must be mapped by π into p or some lesser point; but zero is the only point less than p and $\pi(H') \neq \text{zero}$ via the definition of π ; therefore $\pi(H') = p$; i. e. atom p has a pre-image under π which is a lattice atom in \mathcal{A} .
- 11) Observation: the strong components (the maximal strongly connected subgraphs) of G can be collapsed to single nodes with multiple associated state pairs and all of the results given in this paper will still obtain.

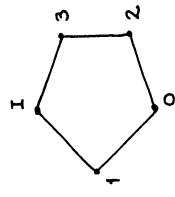
12) Example: Machine C from Hartmanis and Stearns, p. 107

$H \in \mathcal{H}$	NODES OF ASSOCIATED NODES OF G'	$\Pi(H)$	$p \in P$
0	\emptyset	ZERO	0
1	$\{v_1\}$	$0\bar{1}, 1\bar{5}, 2\bar{6}, 3\bar{4}$	1
2	$\{v_2\}$	$0\bar{1}, 2\bar{3}, 4\bar{5}, 6\bar{7}$	2
3	$\{v_1, v_2\}$	I	I
4	$\{v_2, v_3\}$	I	I
5	$\{v_2, v_4\}$	$0\bar{1}, 2\bar{3}, 4\bar{5}, 6\bar{7}$	3
6	$\{v_2, v_3\}$	I	I
7	$\{v_2, v_6\}$	$0\bar{1}, 2\bar{3}, 4\bar{5}, 6\bar{7}$	3
8	$\{v_1, v_2, v_3\}$	I	I
9	$\{v_1, v_2, v_4\}$	I	I
10	$\{v_1, v_2, v_5\}$	I	I
11	$\{v_1, v_2, v_6\}$	I	I
12	$\{v_2, v_3, v_4\}$	I	I
13	$\{v_2, v_3, v_5\}$	I	I
14	$\{v_2, v_3, v_6\}$	$0\bar{1}, 2\bar{3}, 4\bar{5}, 6\bar{7}$	3
15	$\{v_2, v_4, v_5\}$	I	I
16	$\{v_2, v_4, v_6\}$	I	I
17	$\{v_2, v_5, v_6\}$	I	I
18	$\{v_1, v_2, v_3, v_4\}$	I	I
19	$\{v_1, v_2, v_3, v_5\}$	I	I
20	$\{v_1, v_2, v_3, v_6\}$	I	I
21	$\{v_1, v_2, v_4, v_5\}$	I	I
22	$\{v_1, v_2, v_4, v_6\}$	I	I
23	$\{v_1, v_2, v_5, v_6\}$	I	I
24	$\{v_2, v_3, v_4, v_5\}$	I	I
25	$\{v_2, v_3, v_4, v_6\}$	I	I
26	$\{v_2, v_3, v_5, v_6\}$	I	I
27	$\{v_1, v_2, v_3, v_4, v_5\}$	I	I
28	$\{v_1, v_2, v_3, v_4, v_6\}$	I	I
29	$\{v_1, v_2, v_3, v_5, v_6\}$	I	I
30	$\{v_1, v_2, v_3, v_4, v_5, v_6\}$	I	I
31	$\{v_2, v_3, v_4, v_5, v_6\}$	I	I
32	$\{v_1, v_2, v_3, v_4, v_5, v_6\}$	I	I
33	$\{v_1, v_2, v_3, v_4, v_5, v_6\}$	I	I

NOTE: $H_{24} \cdot H_{26} = H_{13}$
 BUT
 $\Pi(H_{24}) \cdot \Pi(H_{26}) = I \cdot I = I$
 $\neq \Pi(H_{13})$



THE $\langle H, +, \cdot \rangle$ LATTICE

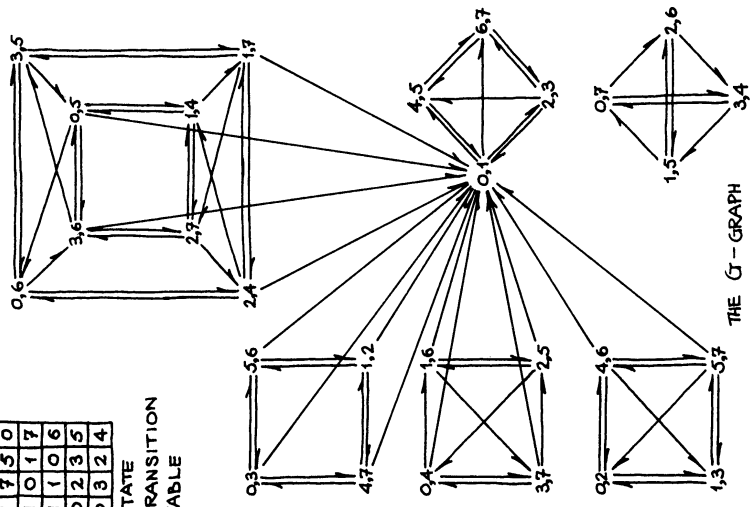


THE $\langle P, +, \cdot \rangle$ LATTICE

NODES OF G'	ASSOCIATED STATE PAIRS
v_1	$\{0,7\}, \{1,5\}, \{2,6\}, \{3,4\}$
v_2	$\{0,1\}, \{2,3\}, \{4,5\}, \{6,7\}$
v_3	$\{0,2\}, \{1,3\}, \{4,6\}, \{5,7\}$
v_4	$\{0,4\}, \{1,6\}, \{2,5\}, \{3,7\}$
v_5	$\{0,5\}, \{1,2\}, \{4,7\}, \{5,6\}$
v_6	$\{0,5\}, \{0,6\}, \{1,4\}, \{1,7\}, \{2,4\}, \{2,7\}, \{5,3\}, \{3,6\}$

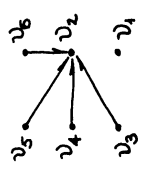
a	b	c	d
0	0	4	6
1	1	5	7
2	0	6	4
3	1	7	5
4	1	0	1
5	1	1	0
6	0	2	3
7	0	3	2

STATE TRANSITION TABLE



THE G' -GRAPH

G' , THE GRAPH OBTAINED FROM G BY COLLAPSING ALL STRONG COMPONENTS



- 13) Bibliography: Hartmanis, J. and Stearns, R. E. , Algebraic Structure Theory of Sequential Machines (Prentice-Hall, Englewood Cliffs, New Jersey, 1966).

