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DETERMINATION OF DROP TRAJECTORIES BY MEANS
OF AN EXTENSION OF STOKES' LAW

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SUMMARY

Trajectory data were determined for drops in air flowing over a cylinder, a sphere, a ribbon, and several airfoils, by reduction of the number of parameters previously used. One trajectory data curve for each body was thus obtained where an entire family of curves was previously necessary.

This reduction of the number of parameters was suggested by Langmuir¹.

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The equations for two-dimensional motion of spherical drops in flowing air (considered incompressible) may be written in non-dimensional form as follows:

$$Kv_x \frac{dv_x}{dx} = \frac{C_{DR}}{24} (u_x - v_x) \quad (1)$$

$$Kv_y \frac{dv_y}{dy} = \frac{C_{DR}}{24} (u_y - v_y) \quad (2)$$

$$v_x = \frac{dx}{dt}, \quad v_y = \frac{dy}{dt}, \quad (3)$$

where:

$$K = \frac{2a^2 U \rho_d}{9L \mu} = \frac{\lambda_s}{L}. \quad (4)$$

λ is the distance the drop will travel if projected into still air with velocity U (in the absence of gravity) if the drag force obeys Stokes' Law.

Langmuir and Blodgett¹ found that results obtained by the differential analyzer solution of (1) and (2) were in good agreement with values resulting from calculations based on K' , an adjusted value of K , and Stokes' Law, where (for flow around a cylinder):

$$K'_o - 1/8 = (K - 1/8) \frac{\lambda}{\lambda_s}. \quad (5)$$

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λ is the distance a droplet will travel if projected into still air with an initial velocity of U when the drag coefficient follows values given in reference 1. λ/λ_s is obtained from:

$$\frac{\lambda}{\lambda_s} = \frac{1}{R_U} \int_0^{R_U} \frac{24}{C_D R} dR . \quad (6)$$

λ_s/λ is therefore an average value of $C_D R/24$ for a drop projected into still air with an initial velocity U and final velocity zero. If, for the case of still air, Eqs. (1) and (2) are solved by using an average value of $C_D R/24K$, there is a K_o for which we can apply Stokes' Law and obtain the same trajectory. K_o is defined by:

$$\frac{1}{K} \left(\frac{C_D R}{24} \right)_{\text{average}} = \frac{1}{K_o} , \quad (7)$$

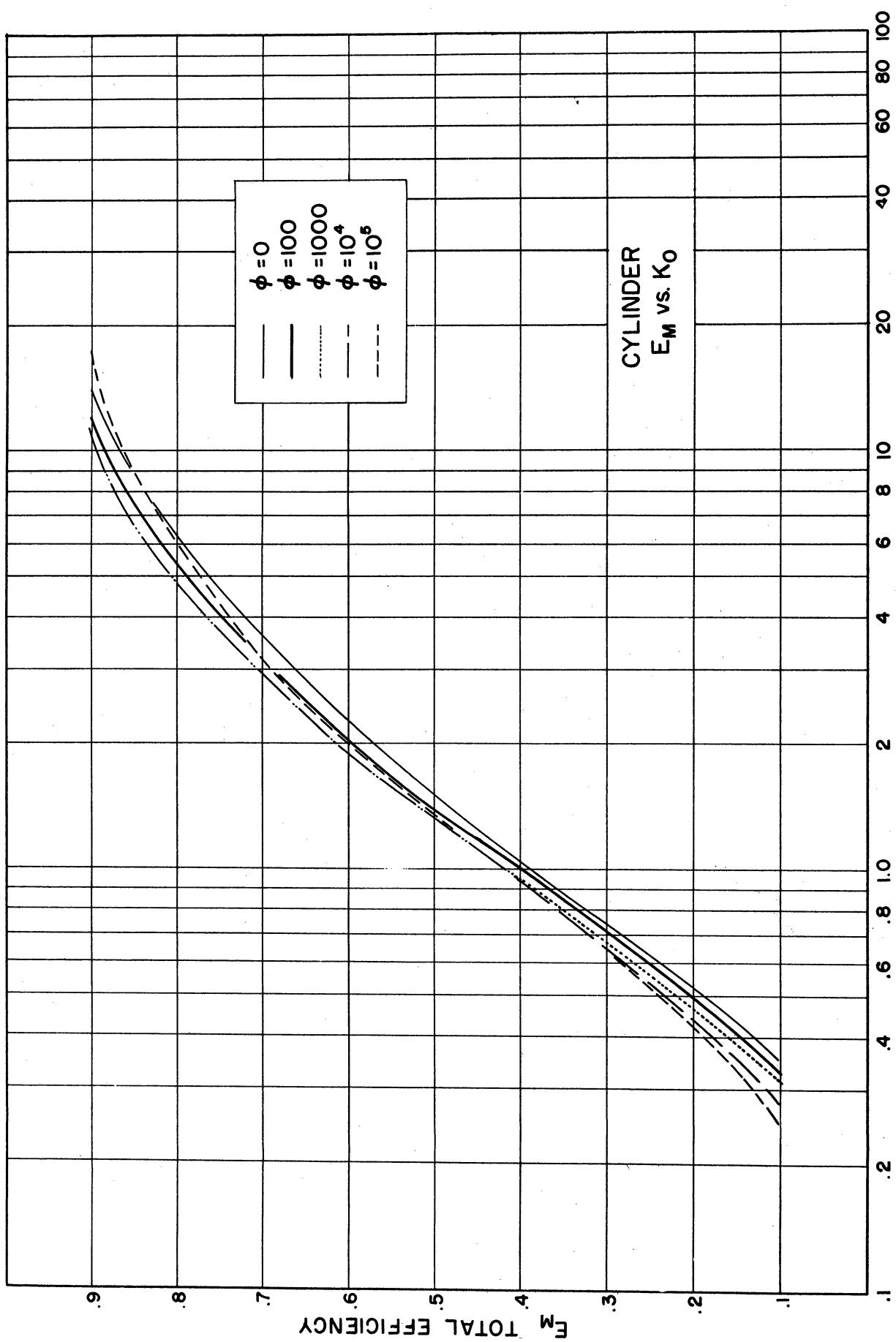
or:

$$K_o = \frac{\lambda}{\lambda_s} K . \quad (8)$$

If this adjustment of K is extended to the case of a varying air-velocity field, the curves of Figures 1-20 result.

An examination of the curves suggests that for each body, instead of a wide spread of curves as shown in the Guibert² and Langmuir and Blodgett reports, only one curve is then necessary.

FIGURE I



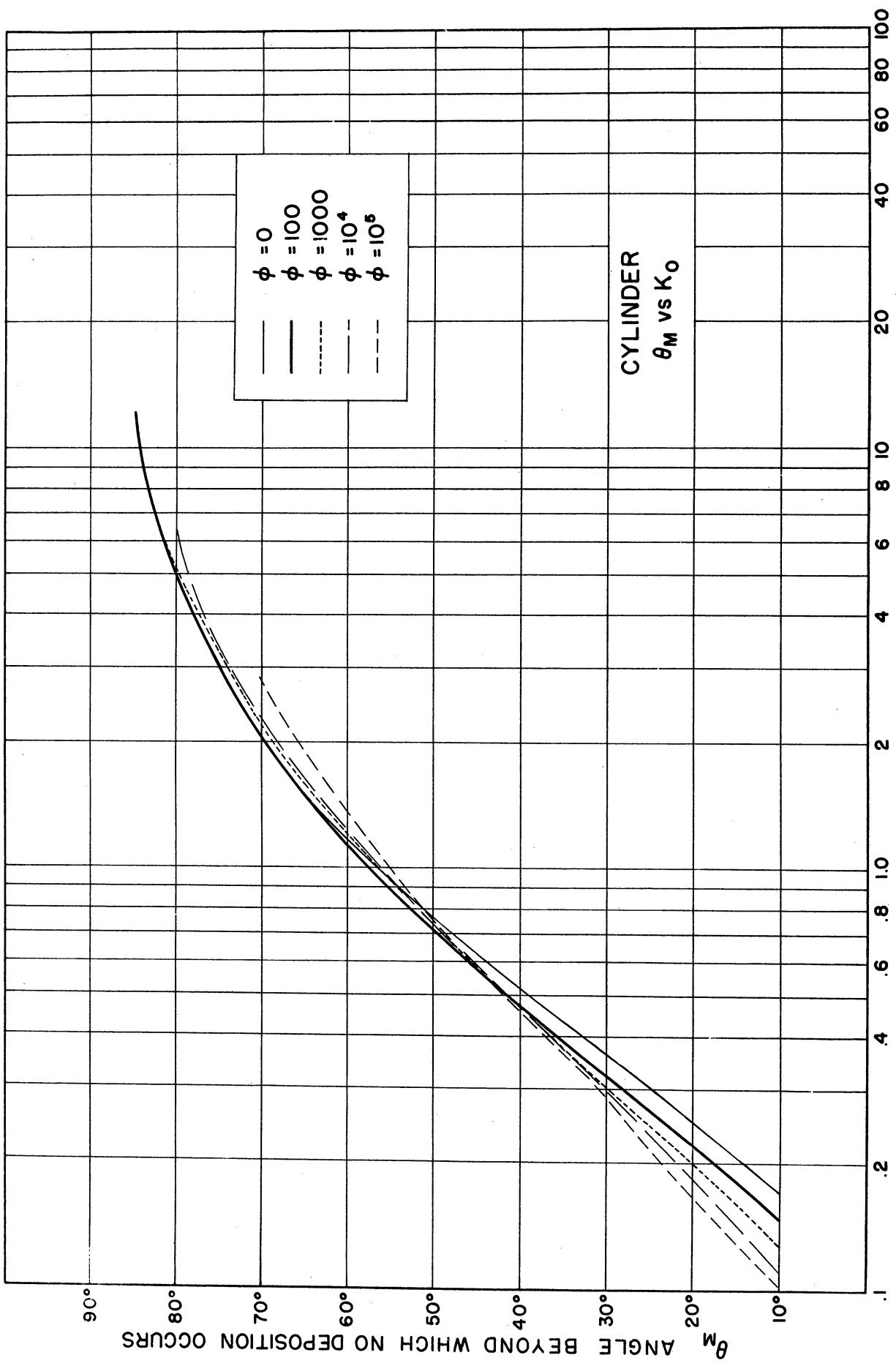
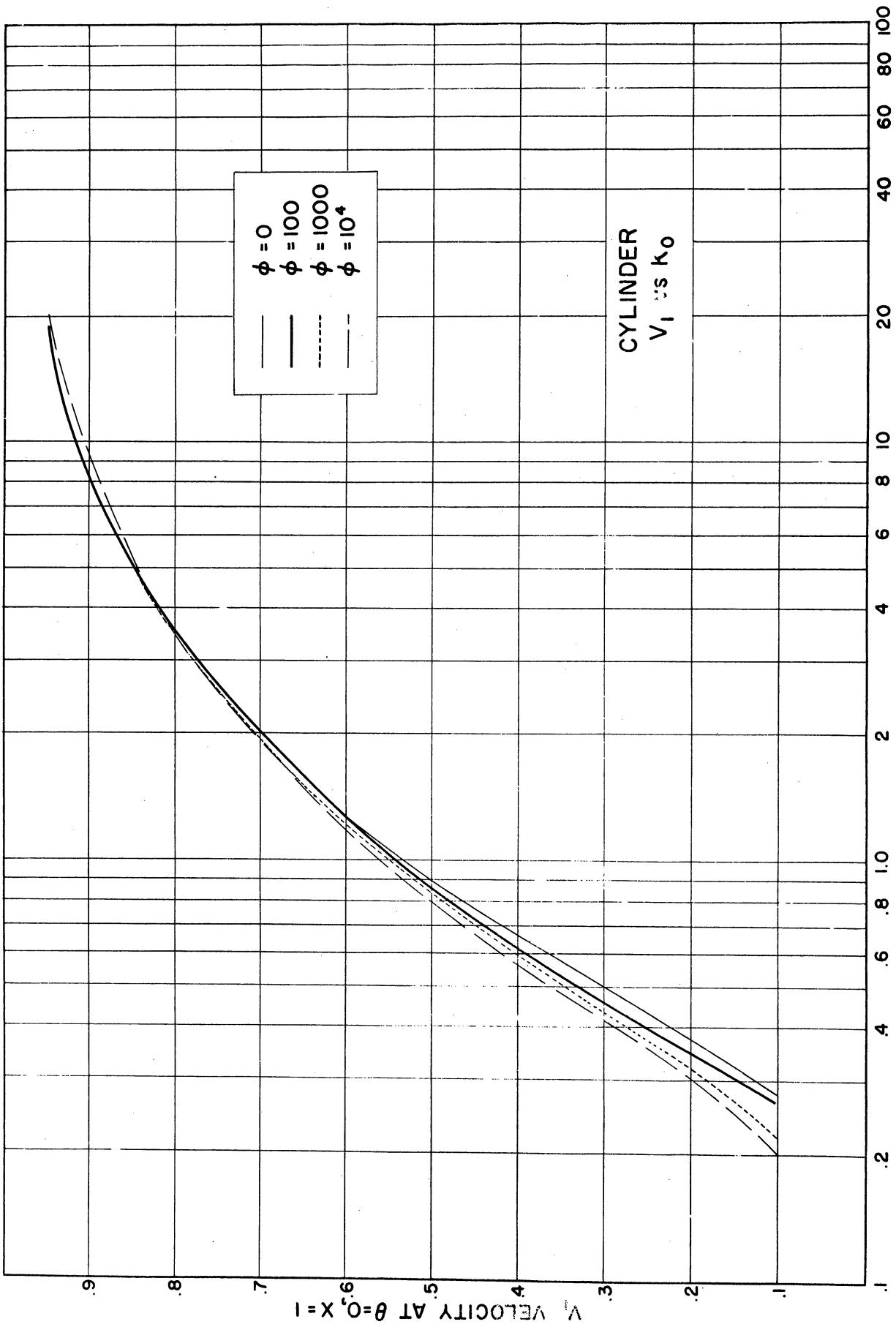


FIGURE 2

FIGURE 3



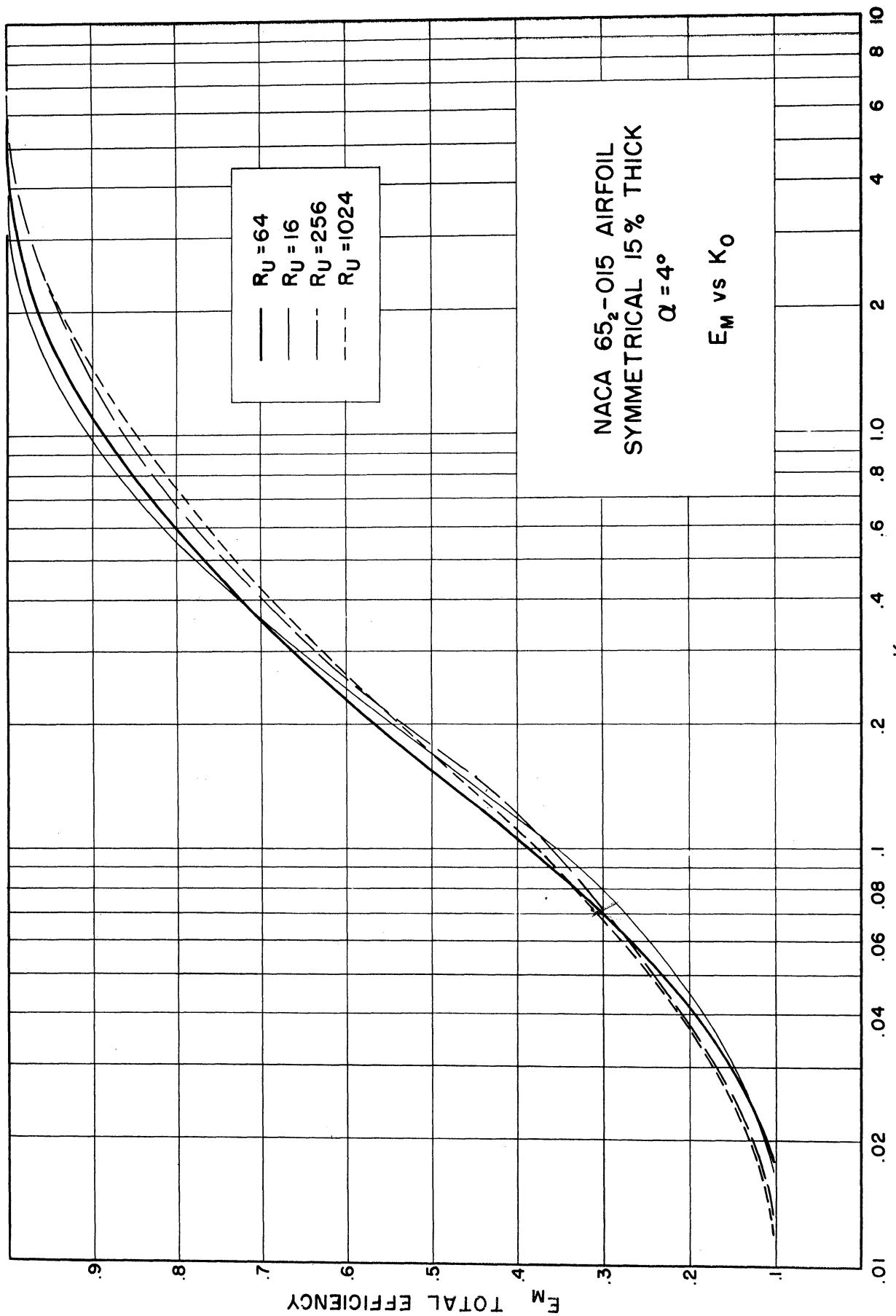


FIGURE 4

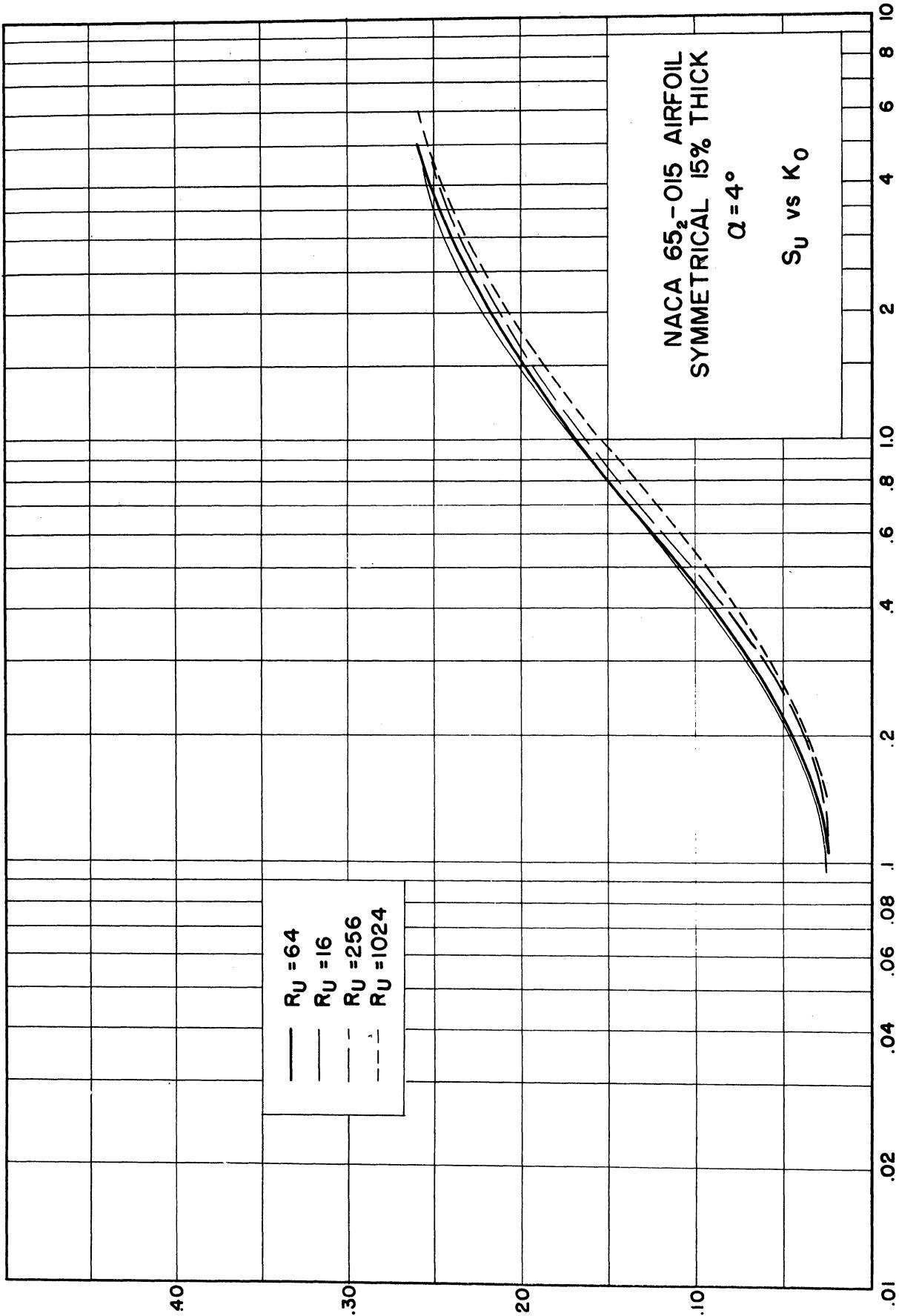


FIGURE 5

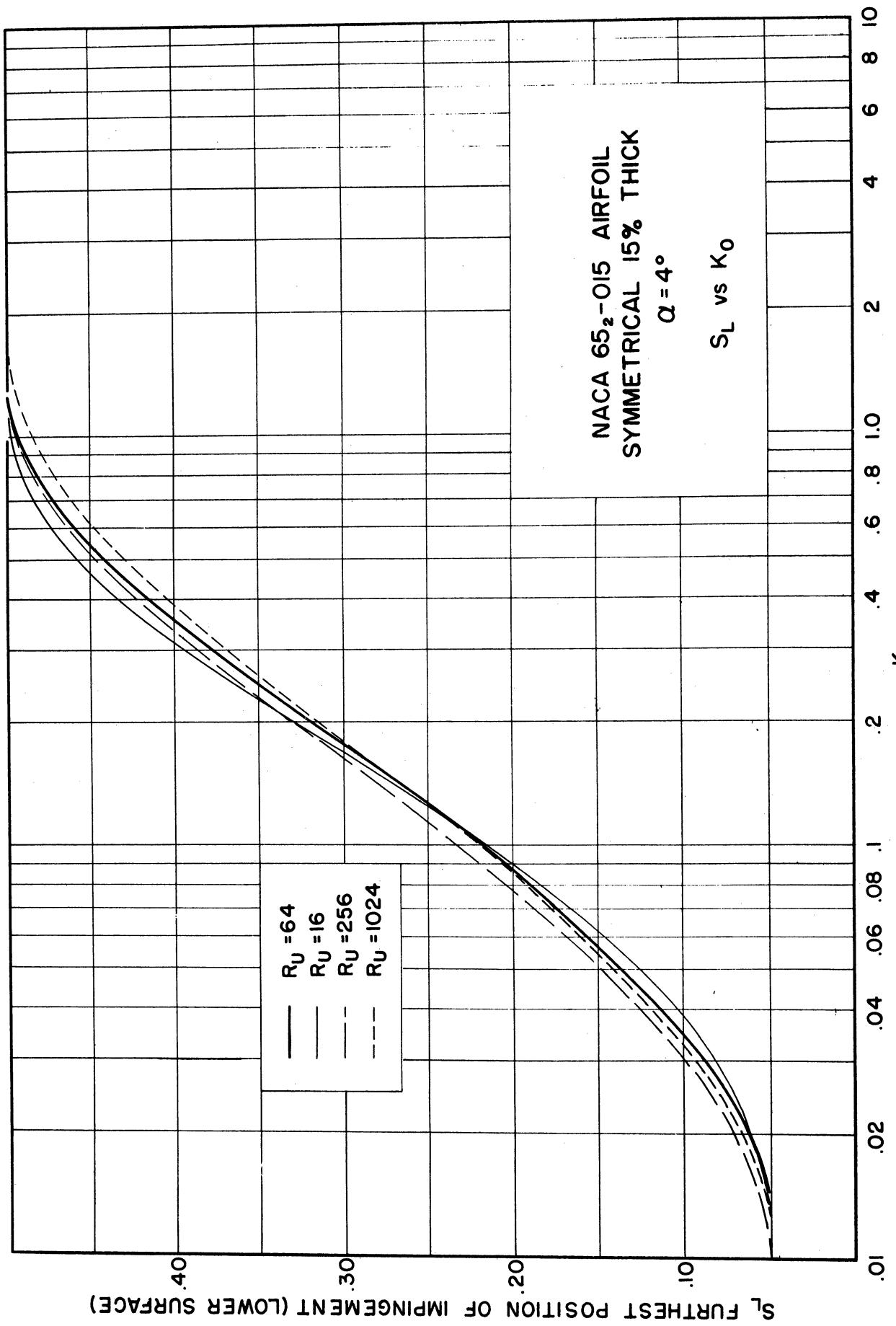


FIGURE 6

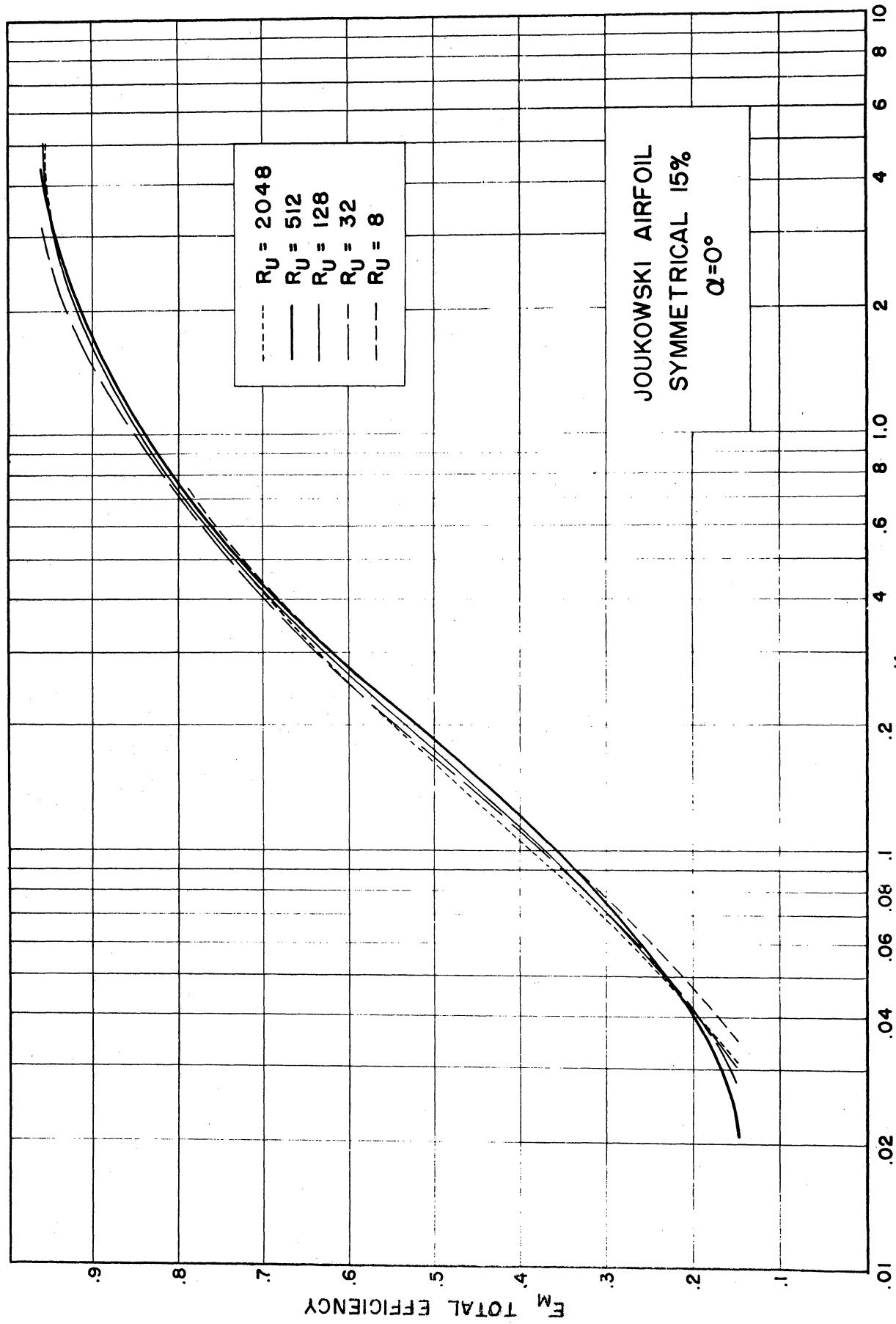


FIGURE 7

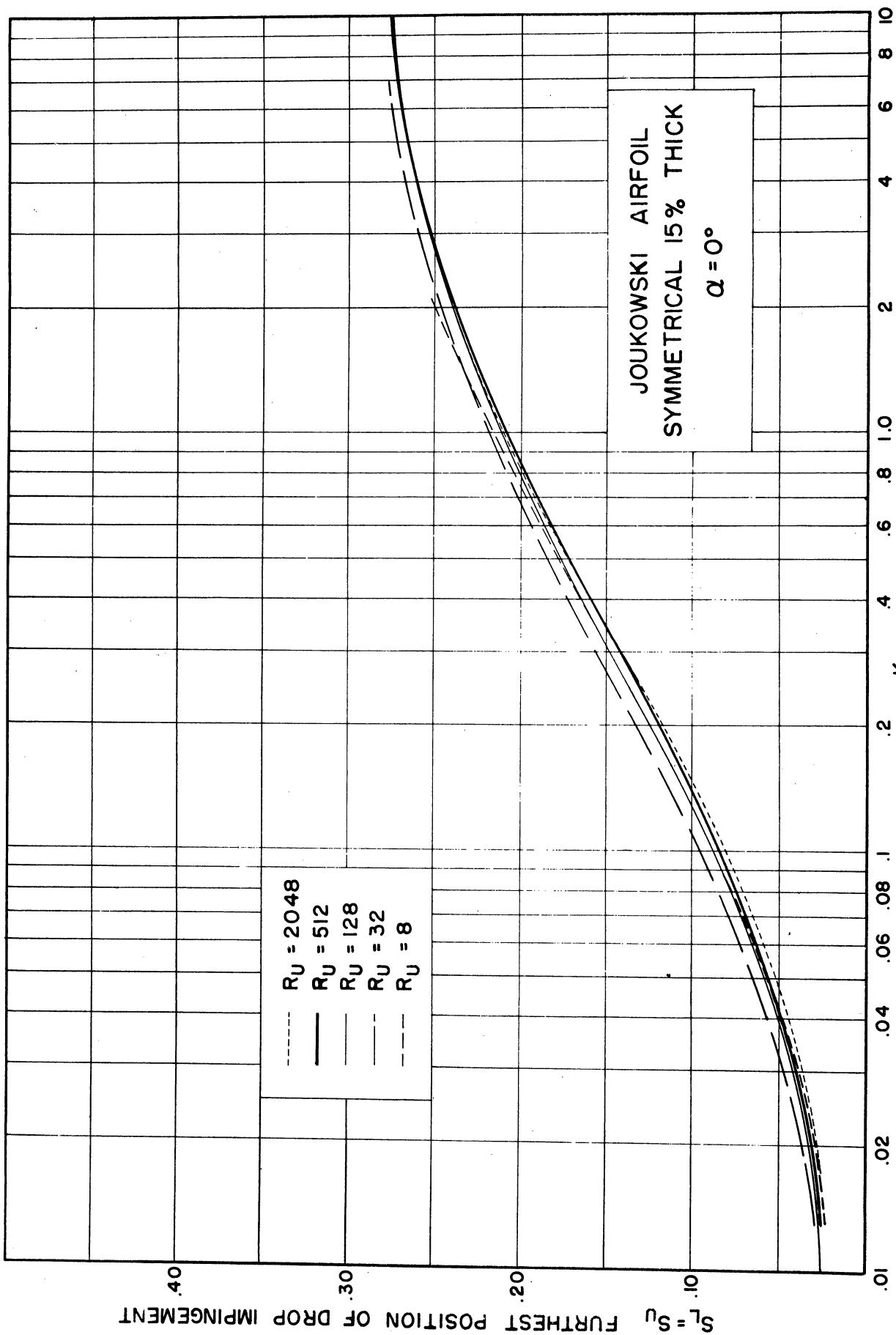


FIGURE 8

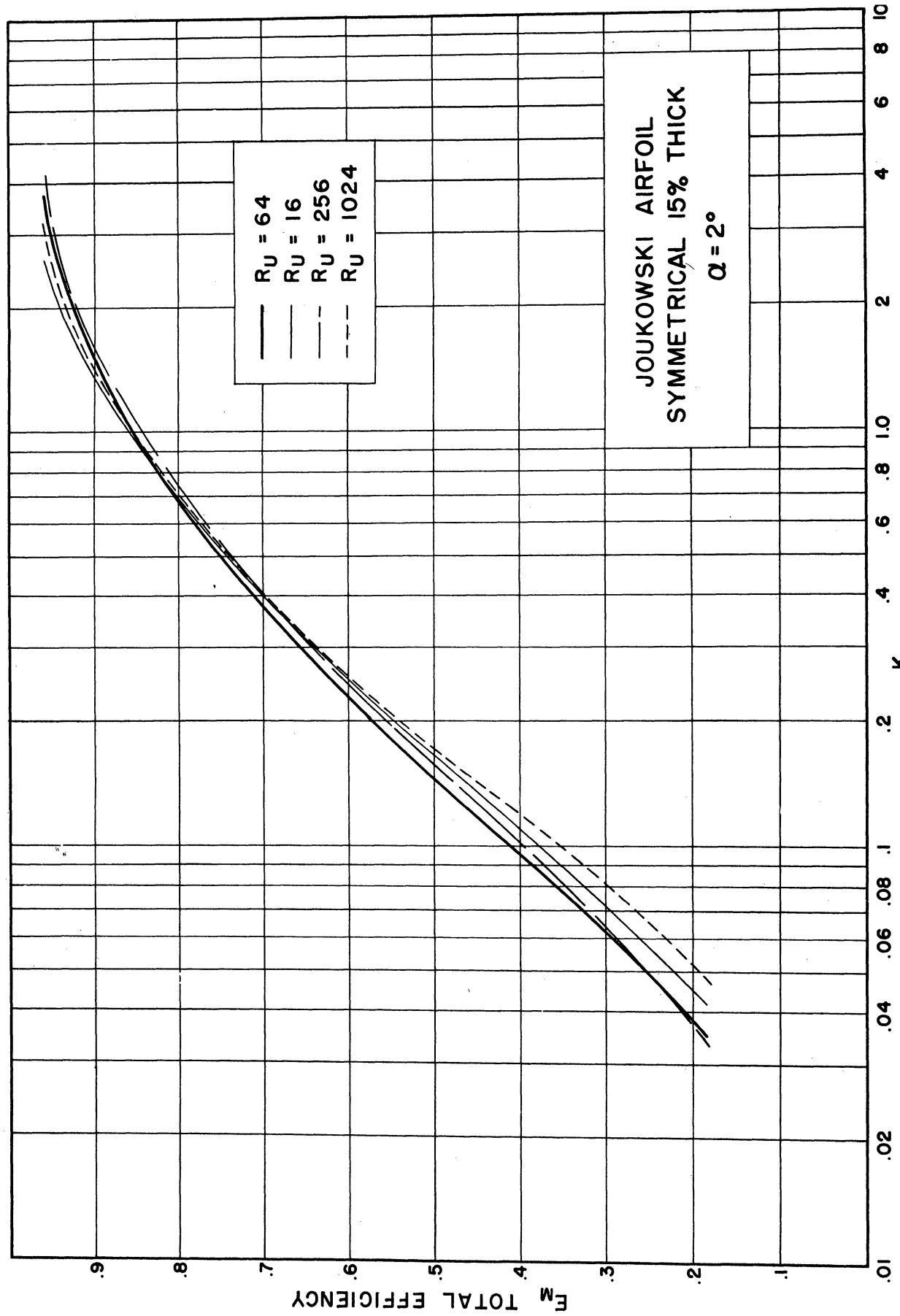


FIGURE 9

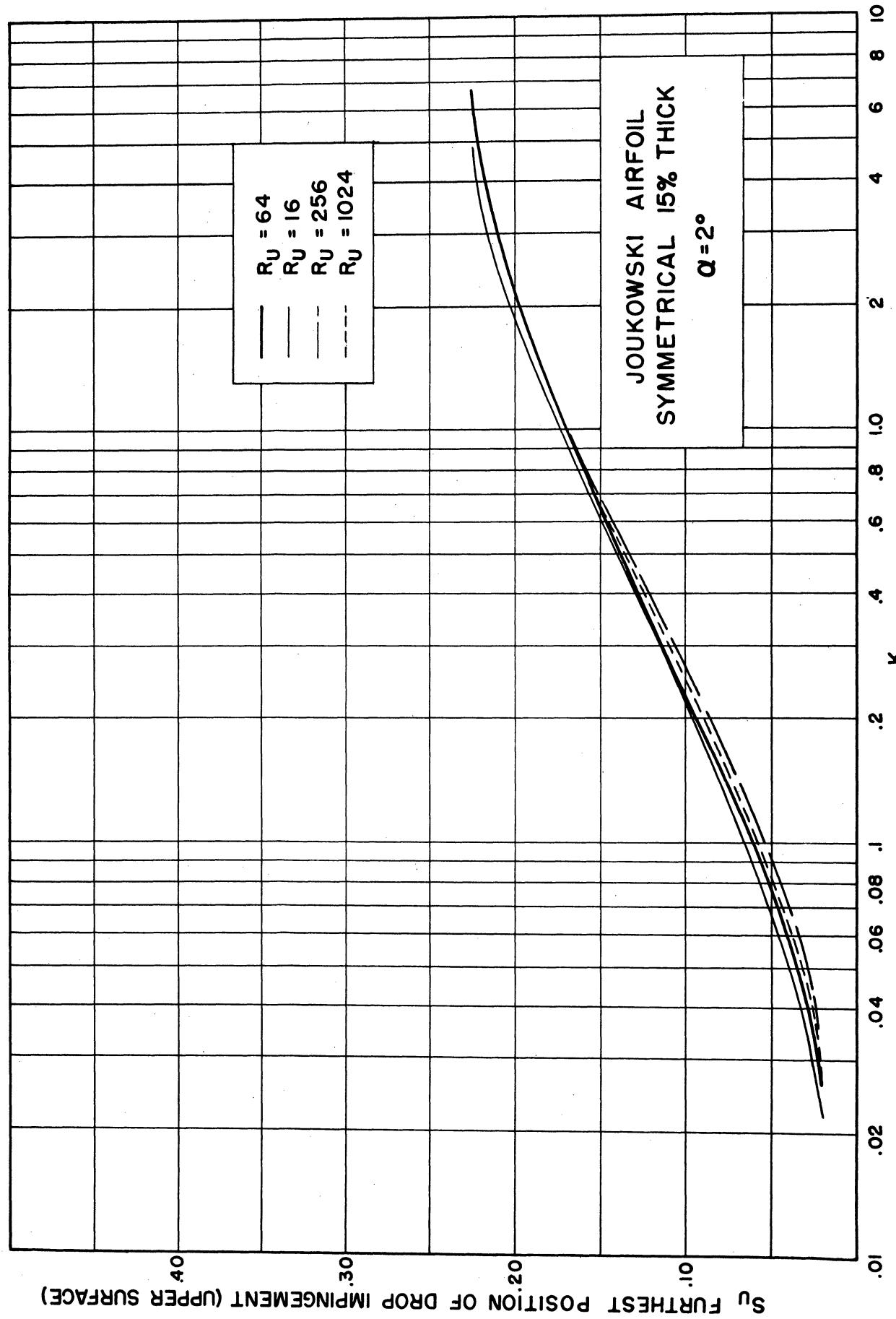


FIGURE 10

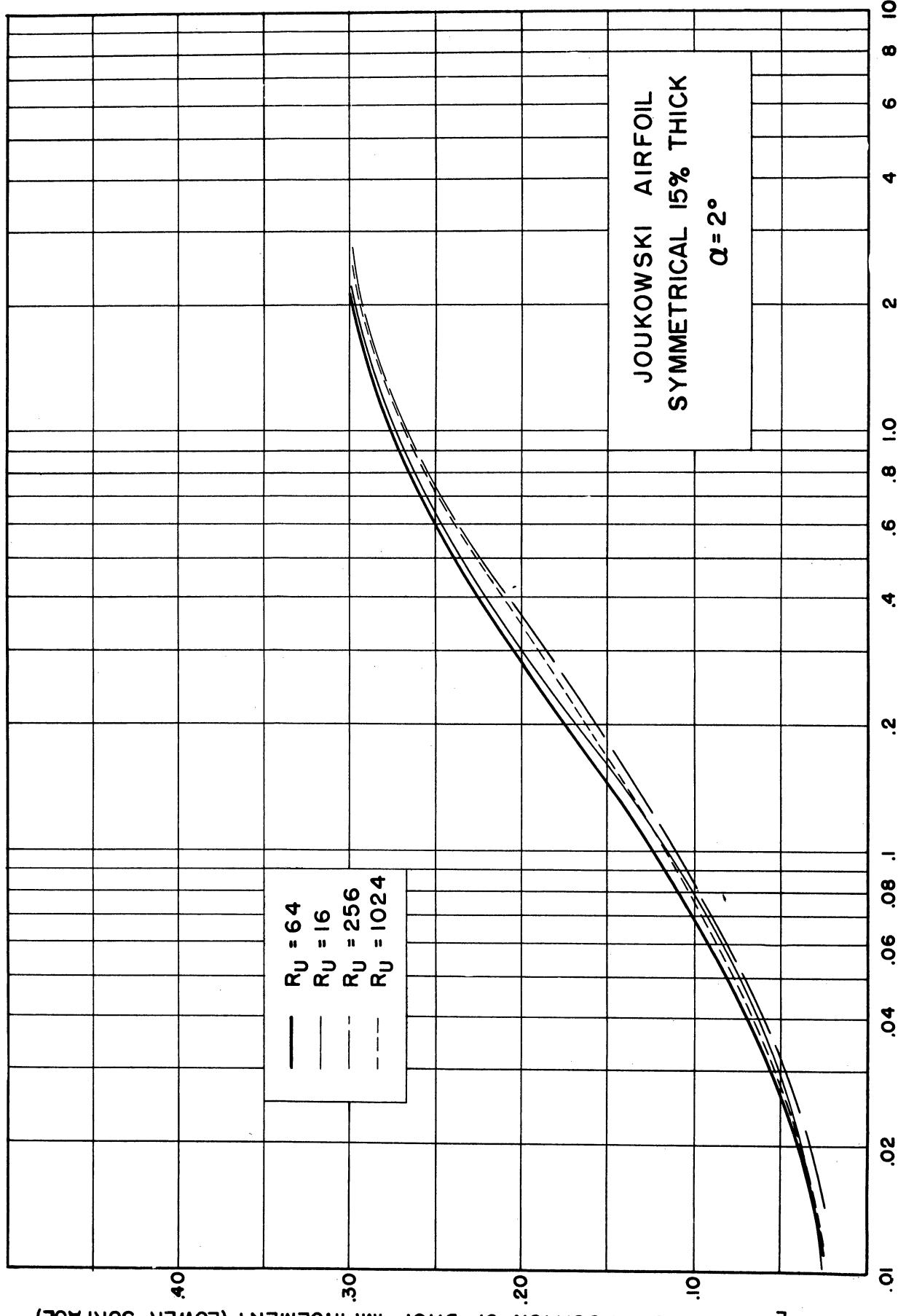


FIGURE 11

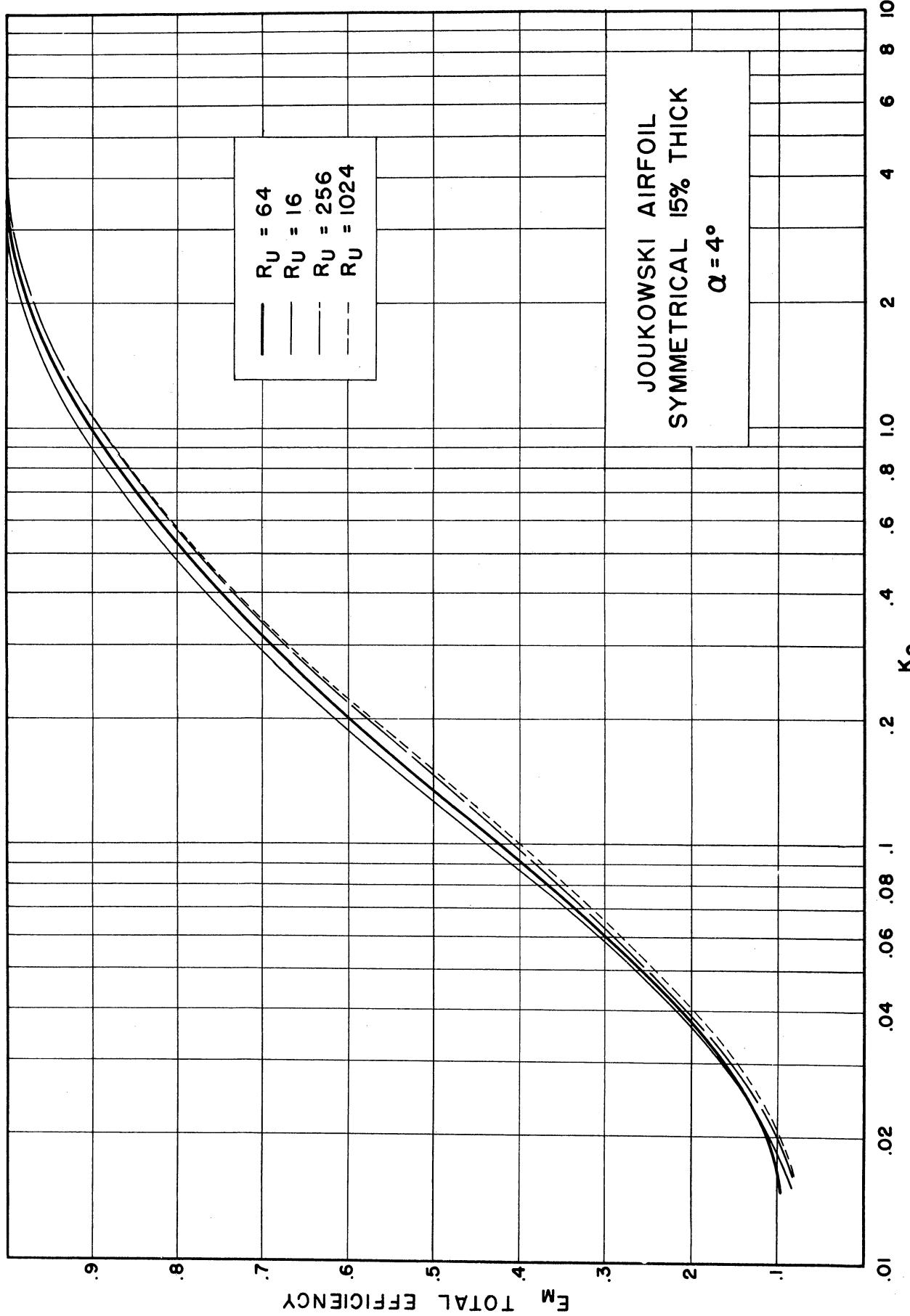


FIGURE 12

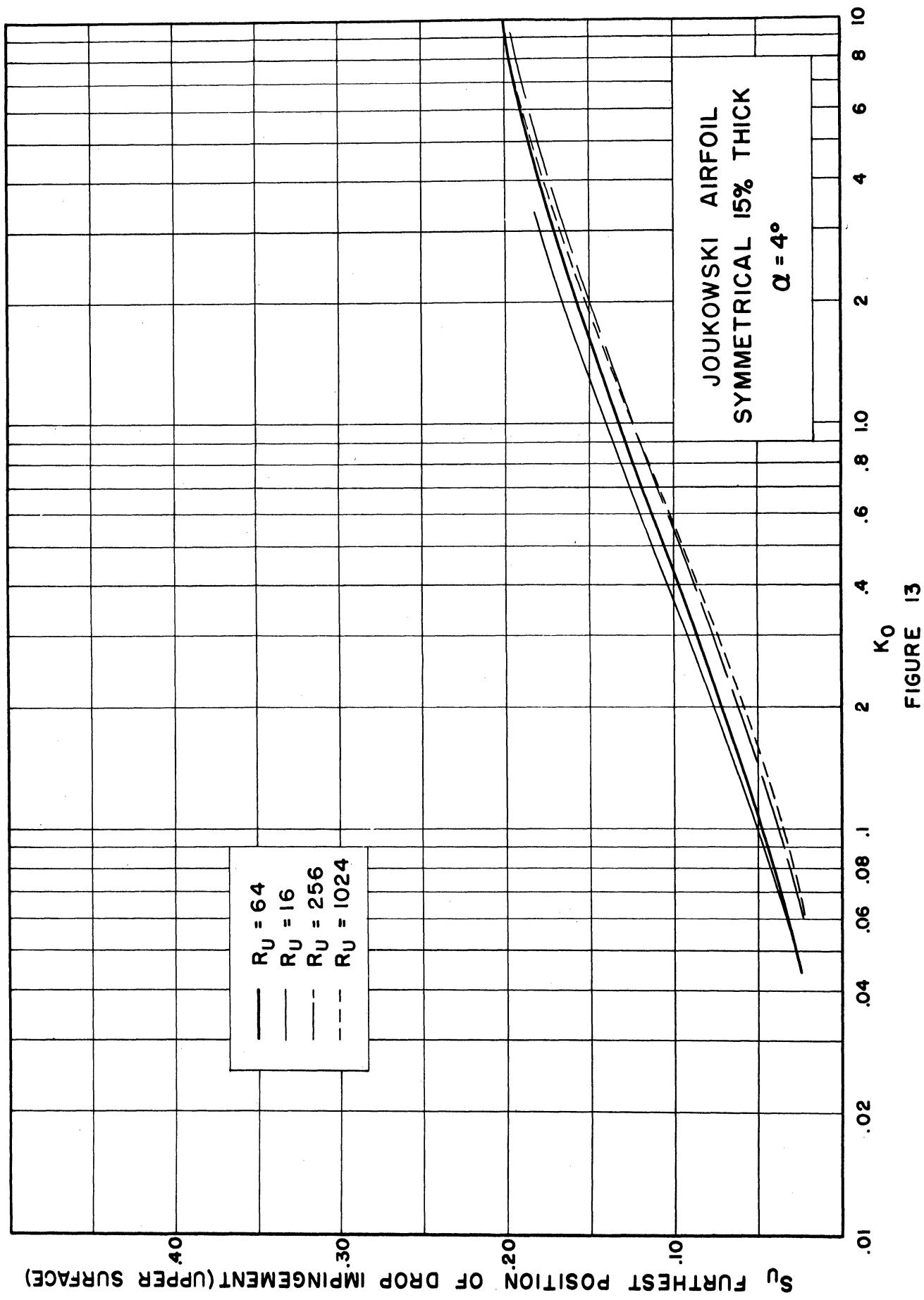


FIGURE 13

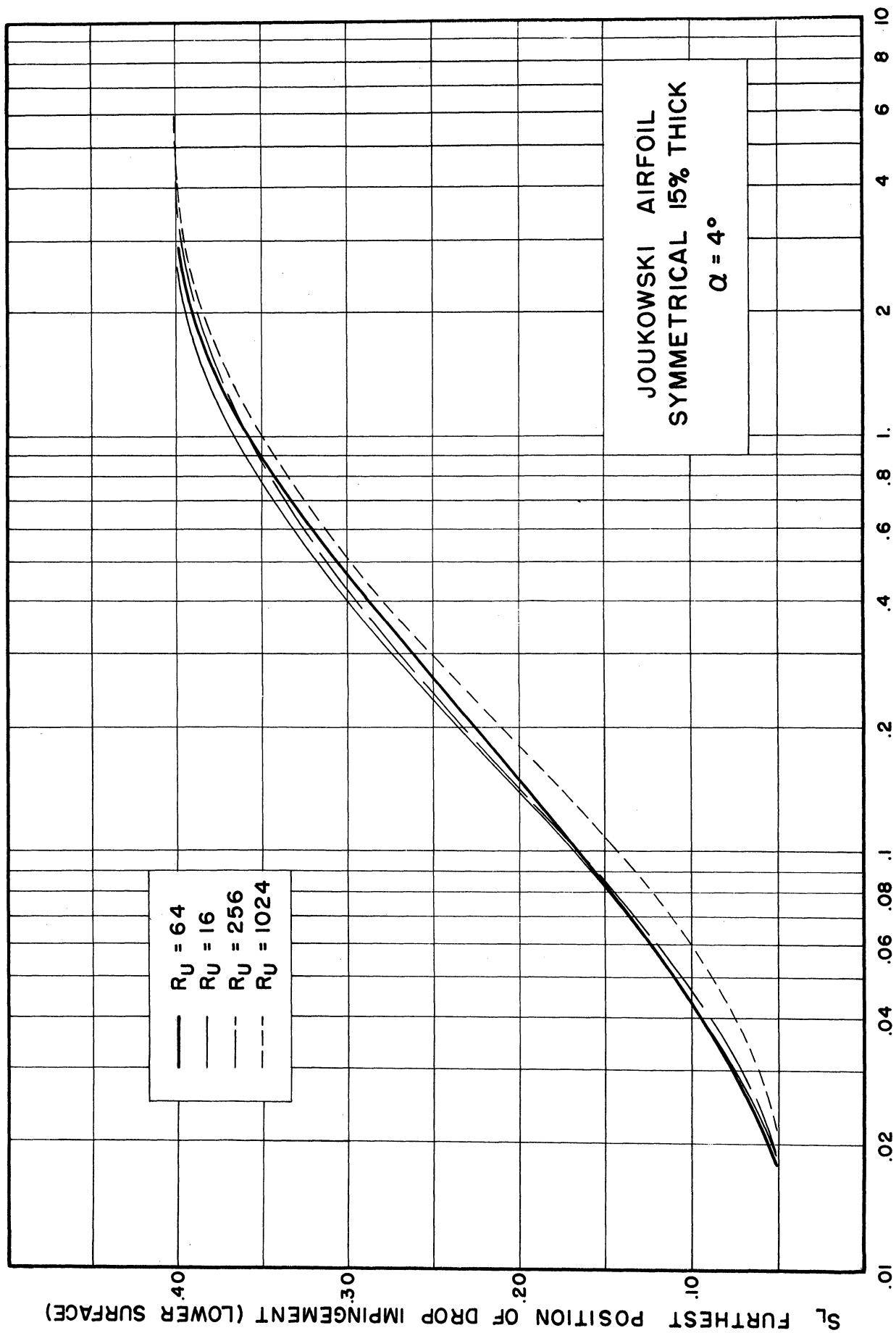


FIGURE 14

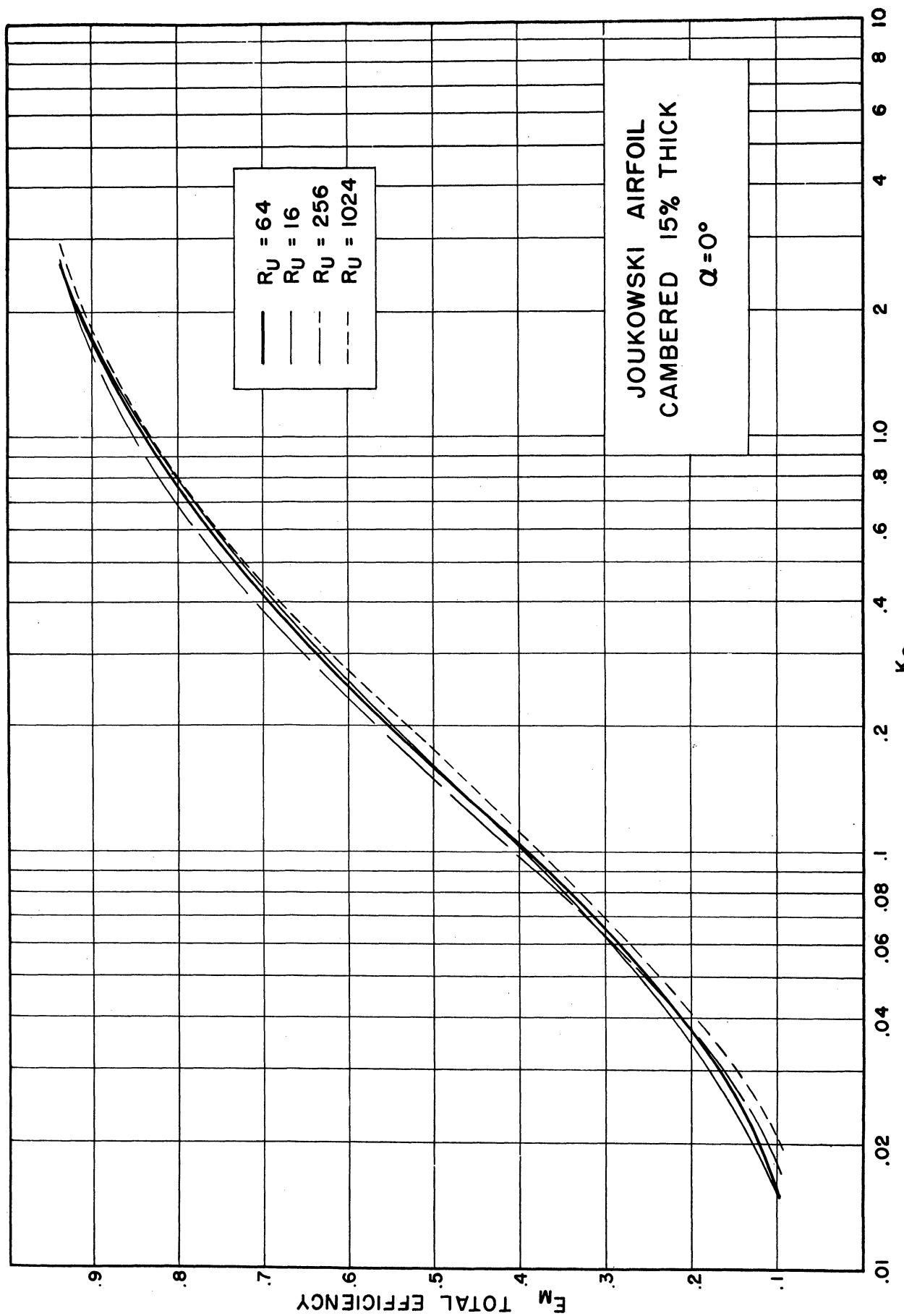


FIGURE 15

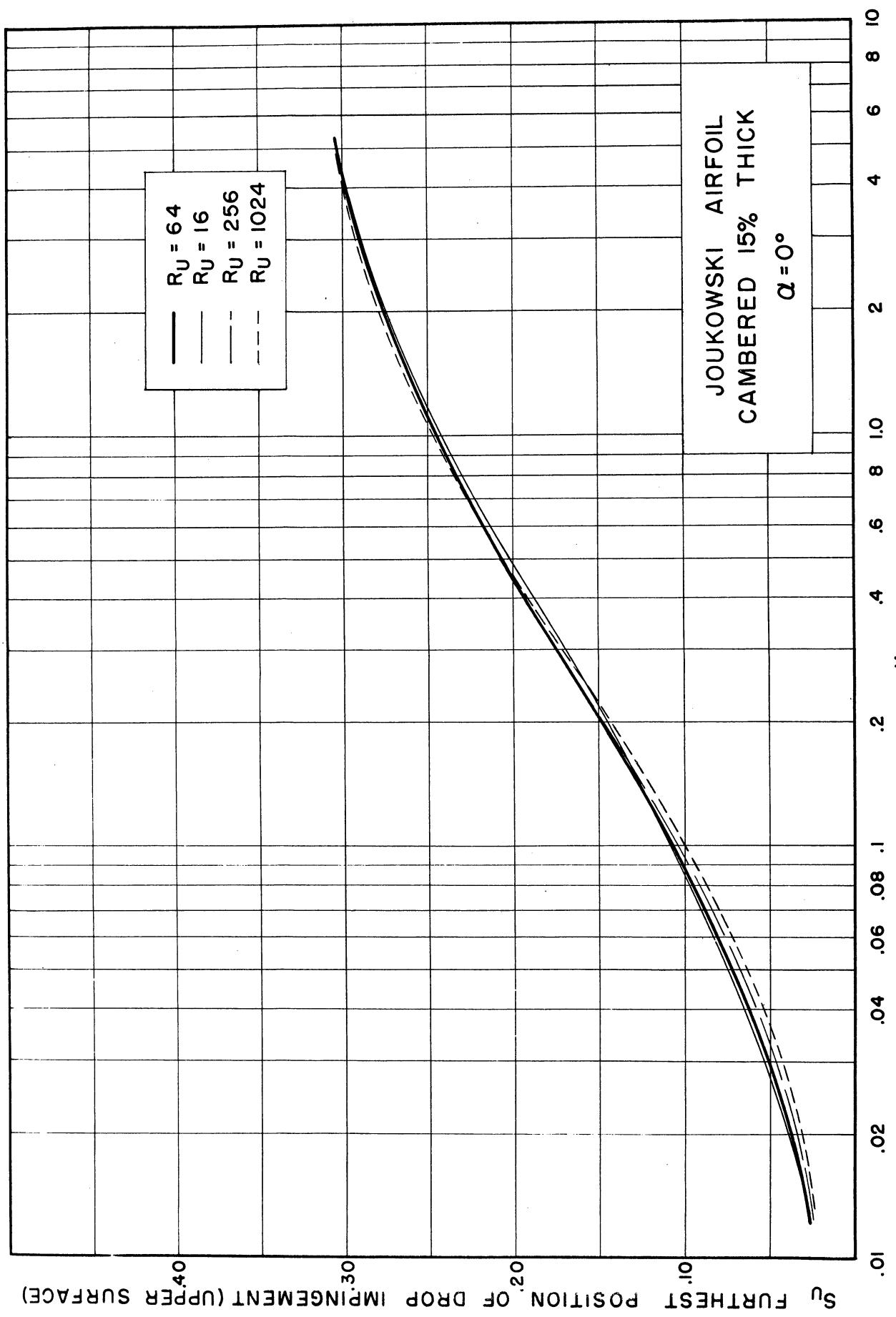


FIGURE 16

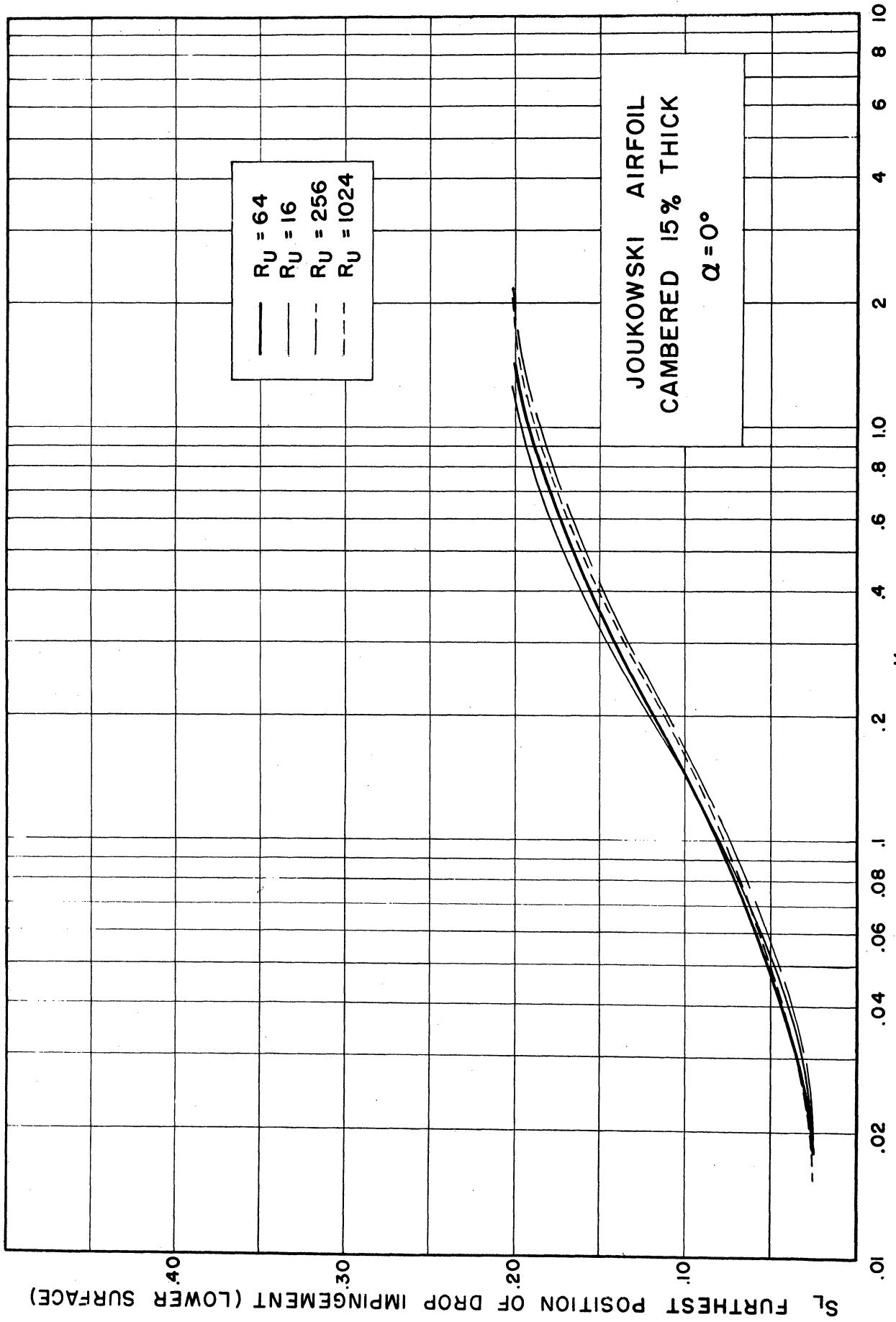


FIGURE I7

E_M TOTAL EFFICIENCY K_O

FIGURE 18

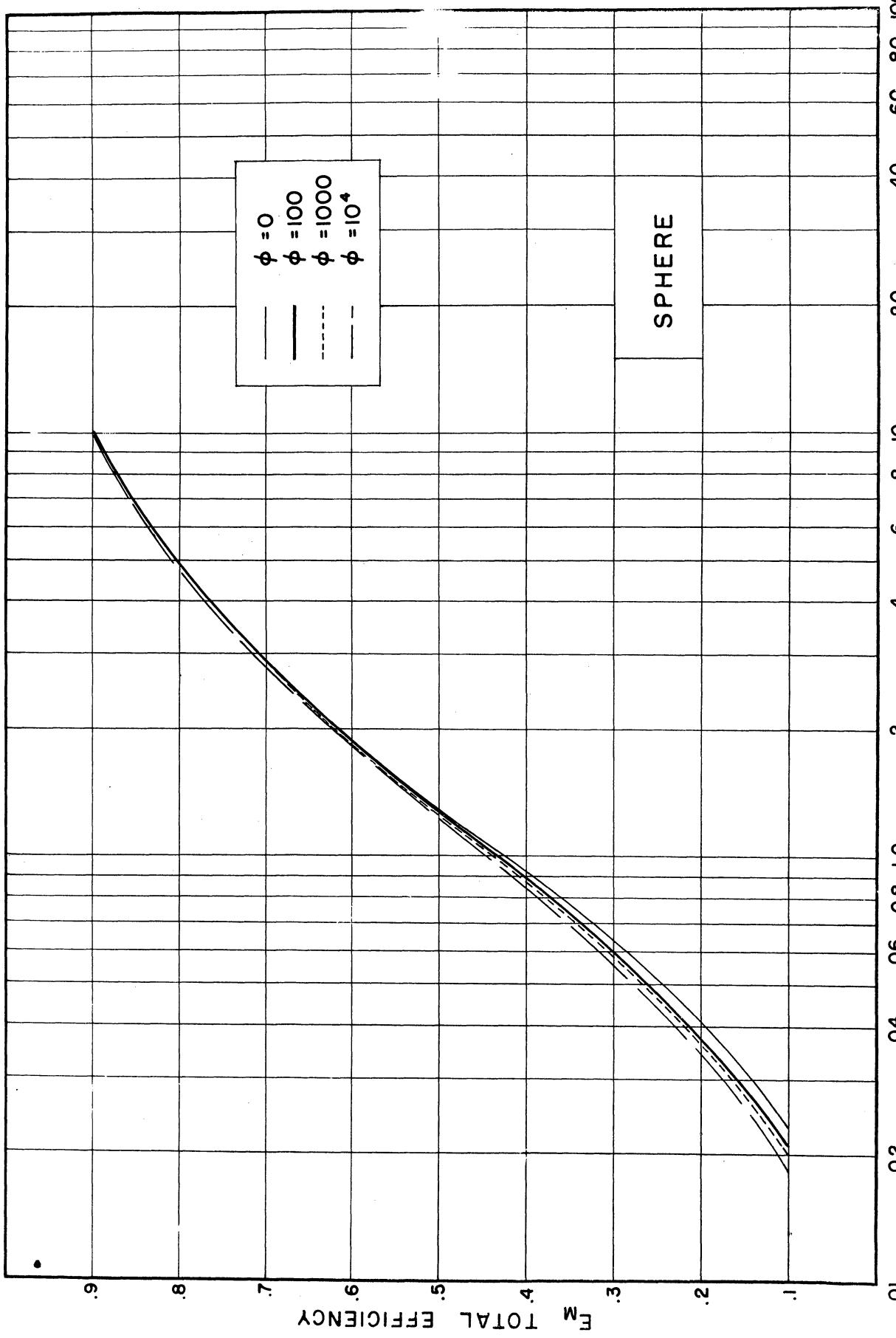
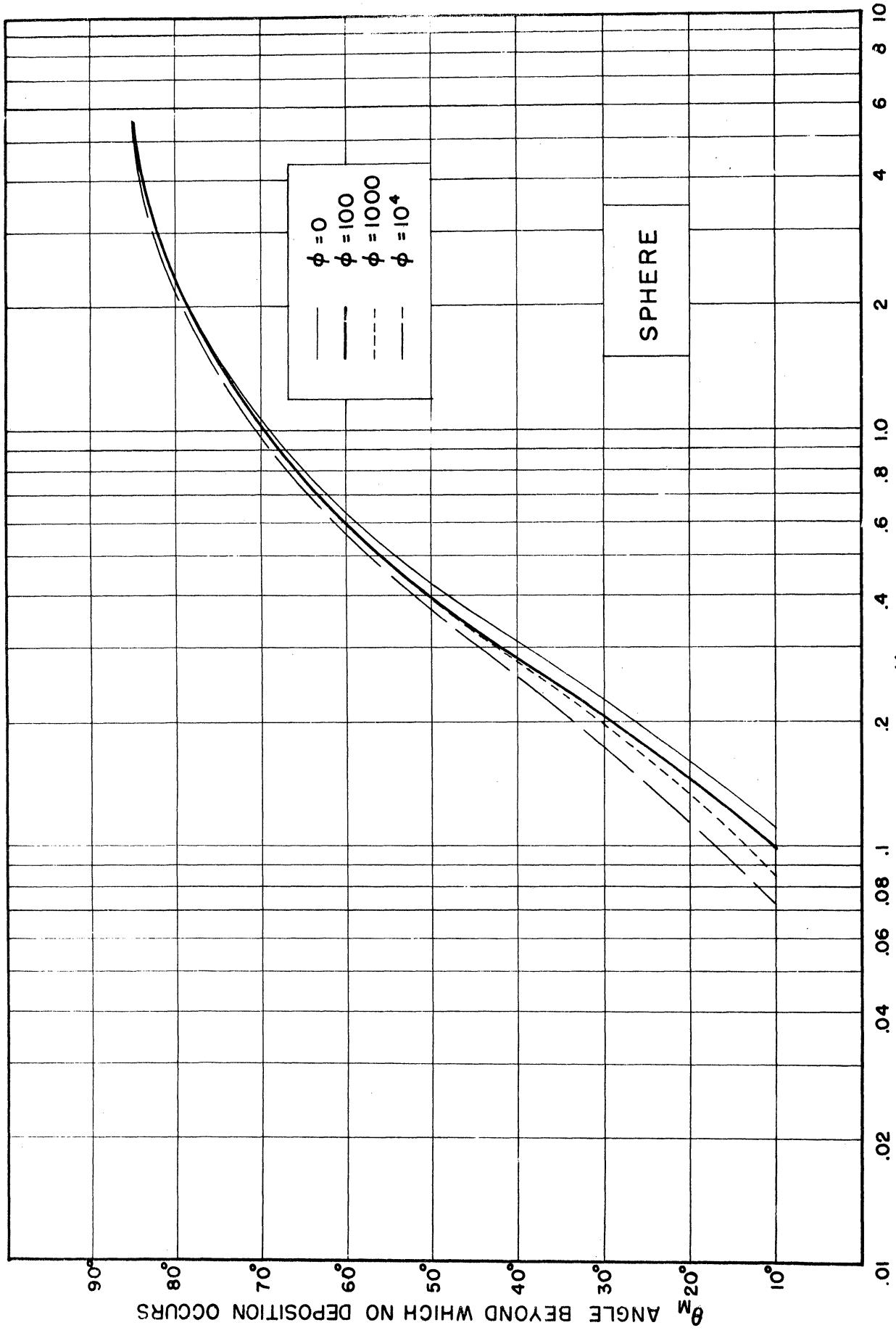


FIGURE 19



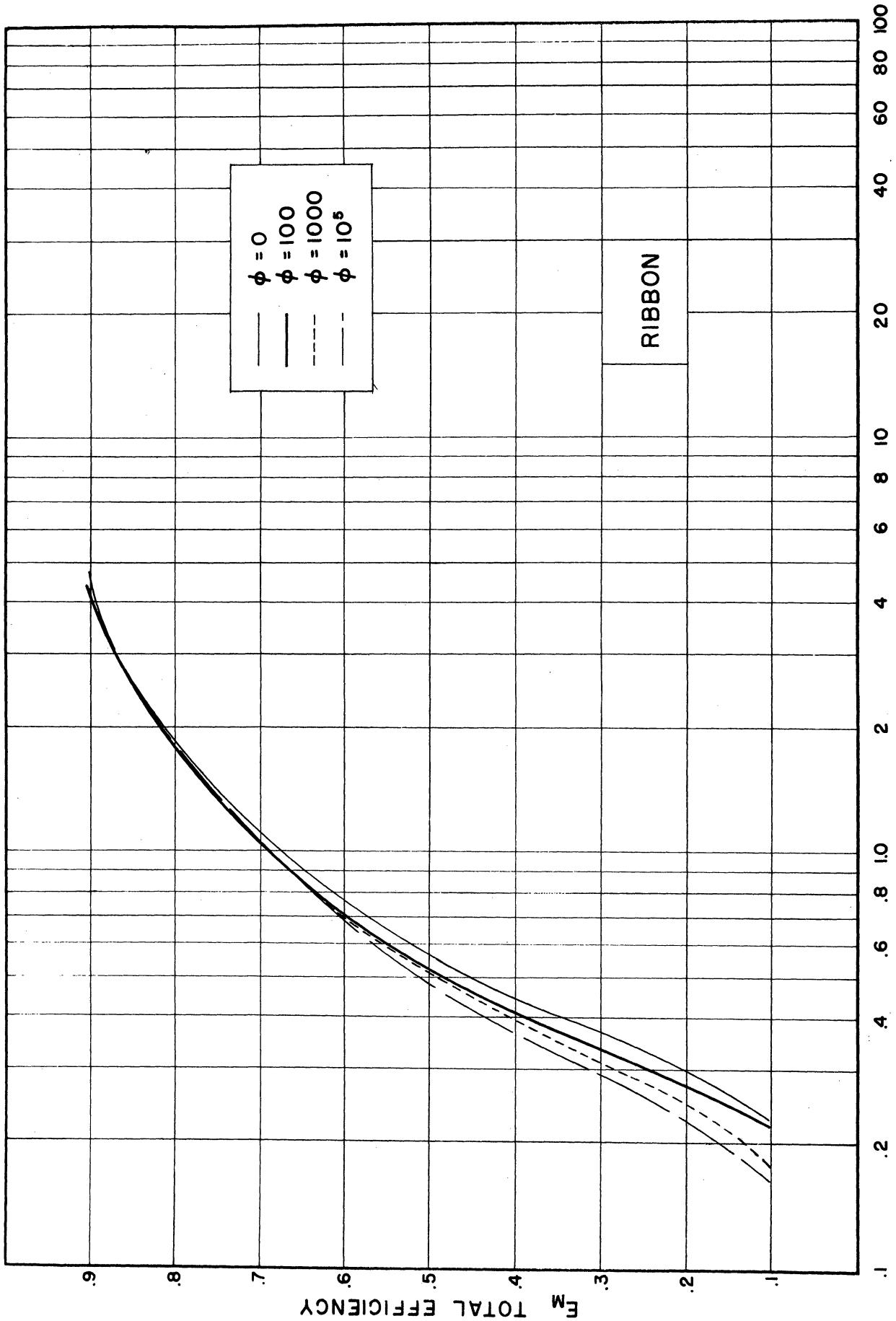


FIGURE 20

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The curves of Figures 1-20 were computed in the following manner: For a given airfoil a fixed value of R_U was chosen corresponding to the differential analyzer curves in the two reports by Guibert. Then for selected values of E_M (or S_U or S_L) the values of ψ were obtained from these curves and K computed from the formula $K = R_U/\psi$. The ratio λ/λ_s was taken from Table I of the Langmuir and Blodgett report, in each case corresponding to the value of R_U under consideration. Finally K_o was obtained from the formula $K_o = K \lambda/\lambda_s$. This results in a table of values giving E_M as a function of K_o . Several values of R_U were chosen for each body so that a family of curves of E_M versus K_o with R_U as parameter results.

The curves pertaining to the sphere, cylinder, and ribbon were obtained in the following manner. A value for ϕ was chosen for a particular body and several values of E_M (or θ_M or V_1 , as the case may be) and corresponding values of K were taken from the Langmuir and Blodgett curves. From $R_U = (\phi K)^{1/2}$ an R_U was calculated for the chosen value of ϕ and each value of K . The ratio λ/λ_s was then obtained from Table I of the Langmuir and Blodgett report. K_o was then calculated from $K_o = K \lambda/\lambda_s$ and E_M (or θ_M or V_1) as a function of K_o was plotted. This procedure was repeated for several values of ϕ for each body.

The following curves are examples of those used in the above calculations.

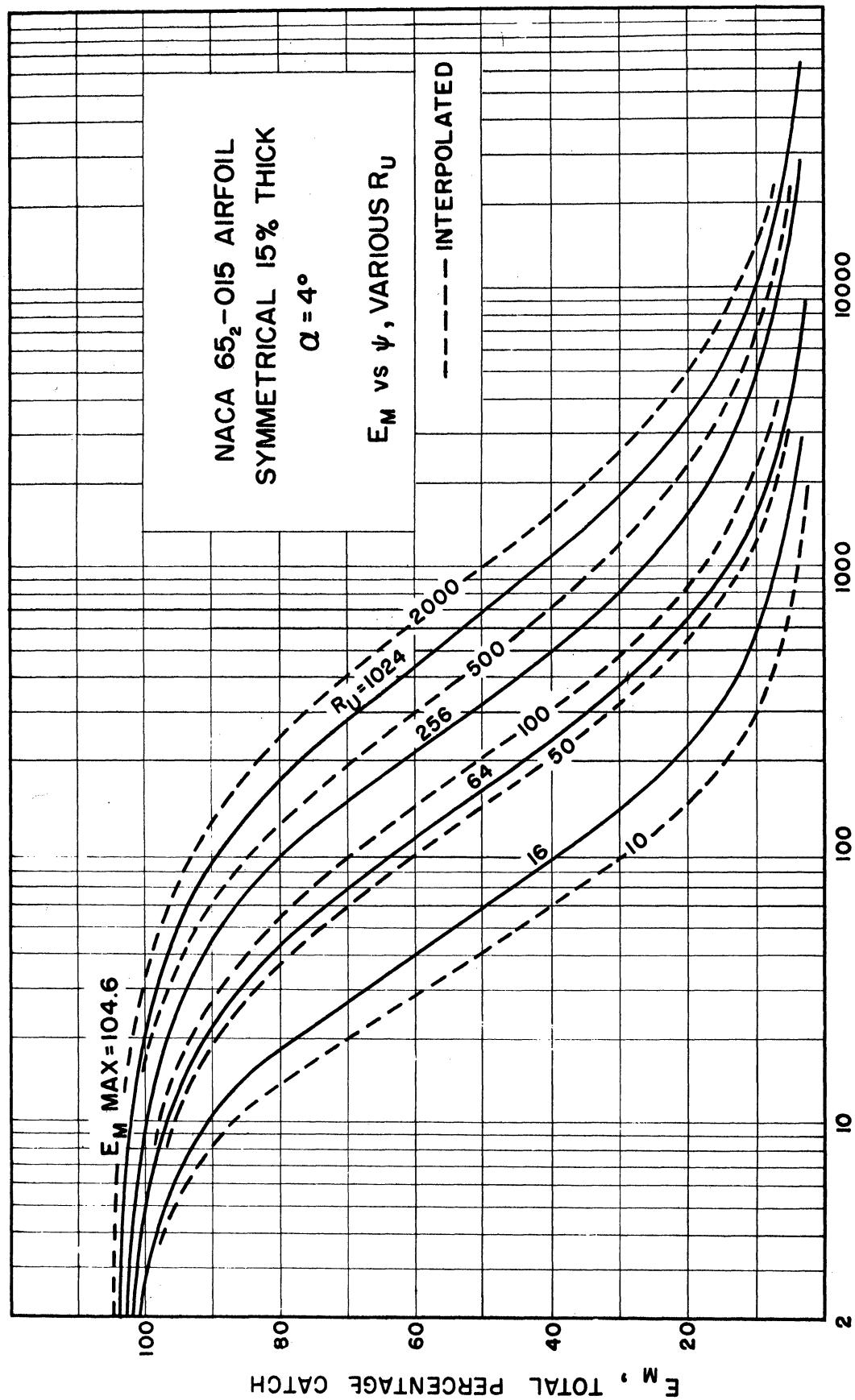
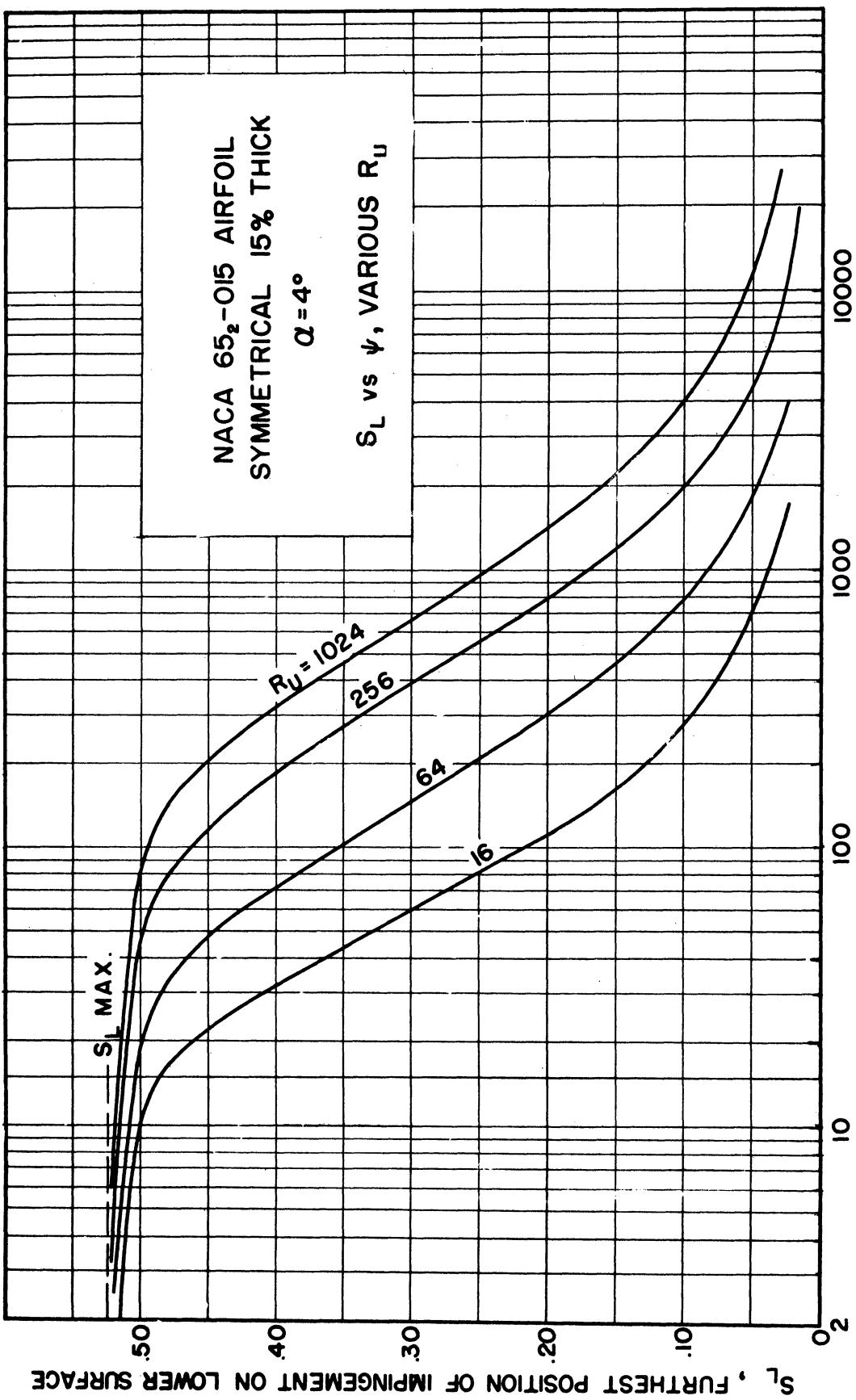


FIGURE 21 REPRODUCED FROM REFERENCE 2



ψ , SCALE MODULUS
 FIGURE 22 REPRODUCED FROM REFERENCE 2.

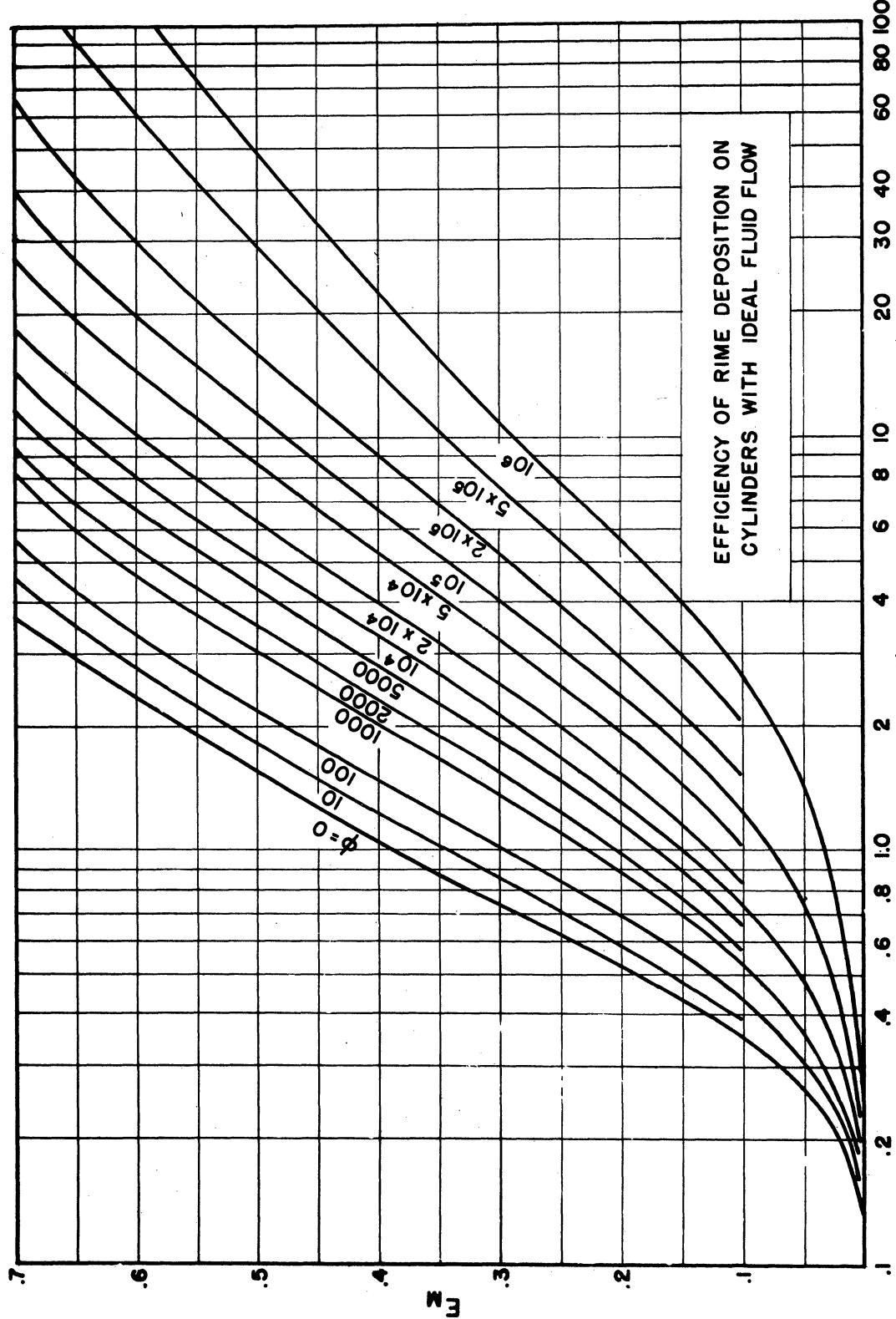


FIGURE 23 REPRODUCED FROM REFERENCE 1

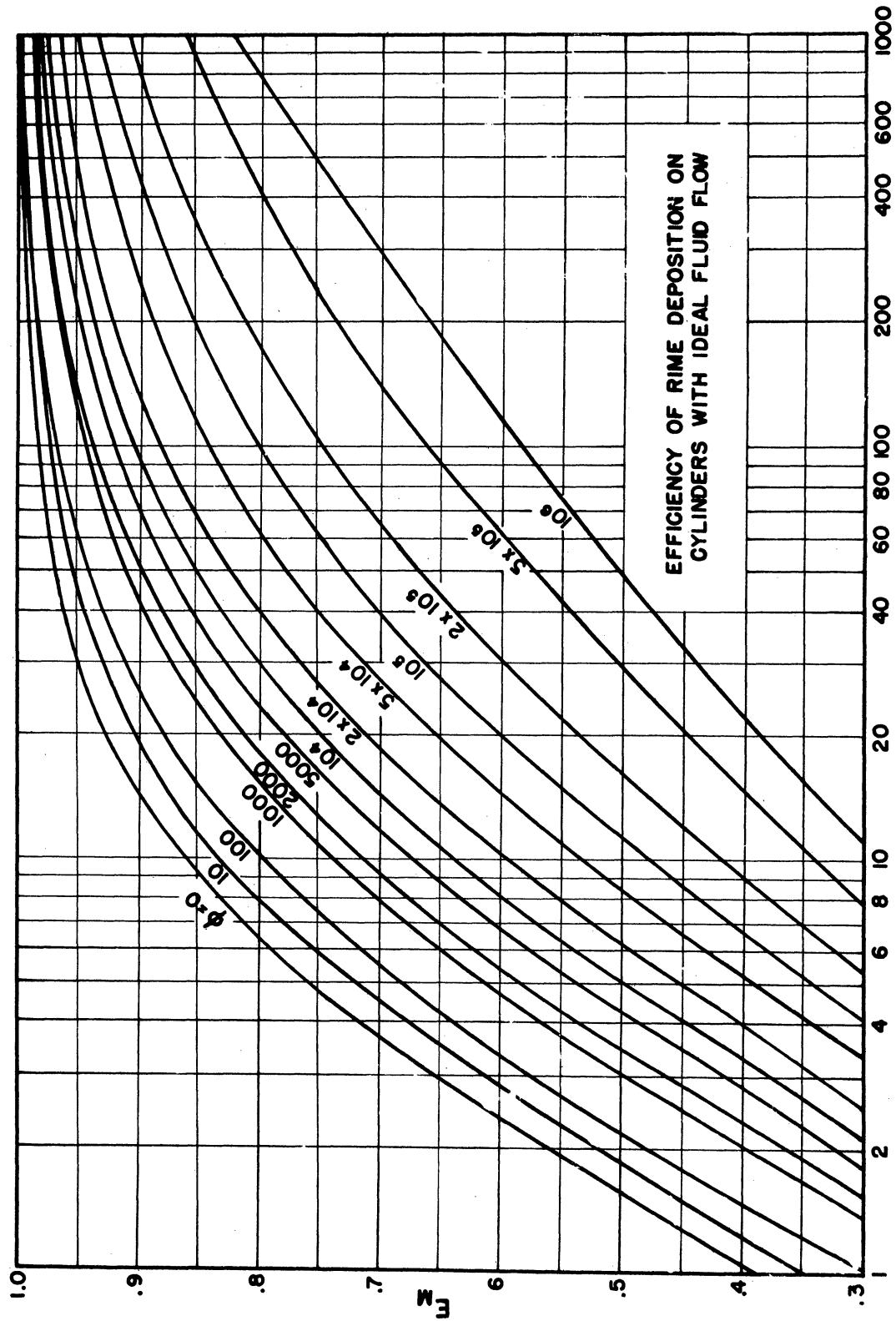


FIGURE 24 REPRODUCED FROM REFERENCE 1.

REFERENCES

1. I. Langmuir and K. B. Blodgett, "A Mathematical Investigation of Water Droplet Trajectories", General Electric Company Report, 1945 (also A.A.F. Tech. Report 5418).
2. A. G. Guibert, E. Jansen, and W. M. Robbins, "Determination of the Rate, the Area, and the Distribution of Impingement of Water Drops on Various Airfoils from Trajectories obtained on the Differential Analyzer", University of California, Department of Engineering, Sept. 1948.
A. G. Guibert Addendum I to above, April, 1949.

SYMBOLS

| | |
|-------------|----------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------|
| a | radius of drop (ft.) |
| c_D | drag coefficient* |
| E_M | total efficiency or total percentage catch; ratio of y at $x = \infty$ to radius of cylinder or sphere; ratio of distance between the initial positions of the upper and lower tangent trajectories to maximum thickness of airfoil* |
| K | defined by (4)* |
| K'_o | defined by (5)* |
| K_o | defined by (8)* |
| L | Characteristic length (ft.); radius of cylinder or sphere; chord length of airfoil |
| R_U | Reynolds Number based on velocity at ∞ , $(2aU\rho_a/\mu)^*$ |
| R | Reynolds Number based on relative velocity of drop in air $(R_U u-v /U)^*$ |
| S_u, S_L | furthest position of drop impingement on surface of airfoil, measured from chord line in units of chord length, on upper and lower surface respectively* |
| t | time in units of U/L^* |
| U | velocity of air at infinity (ft./sec.) |
| u_x, u_y | air velocity components in x and y directions respectively, in units of U^* |
| u | air velocity $(u_x^2 + u_y^2)^{1/2}*$ |
| v_x, v_y | drop velocity components in x and y directions respectively, in units of U^* |
| v | drop velocity $(v_x^2 + v_y^2)^{1/2}*$ |
| V_1 | drop impact velocity at stagnation point of cylinder in units of U^* |
| x, y | coordinates in units of L^* |
| α | angle of attack of airfoil (degrees) |
| λ_s | Stokes' Law range of drop in still air (ft.) |

* dimensionless

| | |
|------------|----------------------------------------------------------------------------------------------|
| λ | range of drop in still air (ft.) |
| μ | absolute viscosity of air (lbs/ft-sec) |
| ϕ | R_U^2/K^* |
| ρ_a | density of air (lbs/ft ³) |
| ρ_d | density of drop (lbs/ft ³) |
| θ | angle measured from x axis of sphere or cylinder (degrees) |
| θ_M | angle measured from x axis of cylinder or sphere beyond which no deposition occurs (degrees) |
| ψ | R_U/K^* |

* dimensionless

