# ENHANCEMENT OF BOILING SURFACES USING NANOFLUID PARTICLE DEPOSITION

by

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# **DEDICATION**

Dedicated to the ones I love.

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#### **ABSTRACT**

To meet increasing power demands across several industries, advanced thermal management systems based on boiling heat transfer have been proposed. Furthermore, nanofluids, a relatively new class of coolants created by suspending 1-100 nm sized particles in a base fluid, have been shown to improve a fluid's thermal properties. This research focuses on two methods using nanofluids to deposit nanoparticles for the creation of enhanced surfaces for boiling heat transfer. Since many of these thermal management systems require a non-conductive fluid, the electrical conductivity of nanofluids is also studied.

Pool boiling studies of nanofluids have demonstrated either enhanced or diminished boiling heat transfer, yet have been unable to distinguish the contributions of increased surface roughness and suppression of bubble transport by suspended particles. This uncertainty is resolved by studying the boiling performance of a surface exposed to a series of boiling tests that alternate between water and a water-based nanofluid. The boiling performance of the coated surfaces increases significantly with each cycle. The measured surface roughness of the intervening nanoparticle layers is used with a model to explain the measured increase in performance. The results demonstrate that the effect of increased surface roughness due to nanoparticle layering can enhance boiling for the base fluid.

A novel method to create enhanced boiling surfaces is electrophoretic deposition of nanoparticles from a nanofluid. A surface was coated using electrophoretic deposition

from a ZnO-propylene glycol based nanofluid. With adequate coating time, such a surface modification method can increase the boiling heat transfer coefficient by about 200%, which was correlated to an increase in the nucleation site density.

In addition, on chip cooling techniques require low conductivity coolants. However, the electrical conductivity of nanofluids has not been widely studied. The particle size and concentration effects on nanofluid electrical conductivity were experimentally investigated and compared to a model based on colloidal suspensions in a salt-free medium. The results showed the electrical conductivity increased with increasing volume fraction and decreasing particle size. At higher volume fractions, the increase of electrical conductivity begins to level off, which is attributed to ion condensation effects in the high surface charge regime.

#### **CHAPTER 1**

#### INTRODUCTION

#### 1.1 MOTIVATION

The need for advanced thermal management systems based on boiling heat transfer with next-generation coolants has been driven by the increase in chip power density and an increase in the amount of electronics onboard a vehicle [1]. Gordon Moore first predicted in 1965 that the number of transistors per chip would double every year [2]. This trend held and now has been observed in other aspects of the semiconductor industry. With this increase in power and decrease in size of transistors, the temperature limit of the current Si chip material is being approached. Currently the power density of chips in high end applications exceeds 1000 kW/m² but must maintain a junction temperature below 105°C [3]. Insulated Gate Bipolar Transistor (IGBT) motor controllers can dissipate up to 400 kW/m² while being held below 125°C in a much harsher automotive environment. Projections have power densities increasing towards 2500 kW/m² over the next few years [4].

There are a number of air cooling methods for power electronics, with the most basic being the natural convection of the chip, which can be further aided by adding a heat sink. In other systems liquid can be flowed through heat sinks to increase the performance. Pin-finned heat exchangers can be easily cast, however, these heat sinks tend to be bulky and significantly increase the weight of the system. To overcome these limitations thermosyphons have been explored as an advanced thermal management

system. A thermosyphon removes heat by conducting it away from the source to boil a fluid in the evaporator section, from which vapor rises to the condensing section. The vapor condenses and the condensate returns to the evaporator. This allows heat from a concentrated source, such as an electronic chip, to be spread out to a larger area using the latent heat of vaporization of the fluid. Thermosyphons rely on natural convection and do not require external power, making them inexpensive and quiet.

Advanced cooling techniques are based on boiling heat transfer since it is an efficient method of removing large heat fluxes through the phase change of the liquid. Pool boiling is boiling of a heated surface submerged in a pool of liquid. Fig. 1.1(a) shows a characteristic curve of boiling heat transfer over four distinct regimes, and Fig. 1.1(b) illustrates the bubble formation in these regimes [5]. The heat flux at the surface (vertical axis) and the superheat (horizontal axis) are measured during boiling. First, heat is removed from the surface through natural convection and no bubbles are formed. Next in the nucleate boiling regime, the liquid at the surface begins to evaporate and bubbles form at nucleation sites. This regime is characterized by a very high heat transfer rate at low surface temperatures. Once the critical heat flux (CHF) has been reached, the population of bubbles obstructs the path of incoming liquid and forms an insulating vapor layer. This is referred to as the transition boiling regime. At this point there is a rapid decrease in the heat that can be removed and a rapid increase of the surface temperature. Finally in the film boiling regime, a stable vapor film is formed and there is no fluidsurface contact. When boiling heat transfer (BHT) is enhanced, the curve of Fig. 1.1(a) shifts up or to the left. This means more heat can be removed from the surface at the same temperature, or at the same heat flux, the surface temperature is lower.

Enhancement of pool boiling heat transfer can accomplished through (a) modifying the surface, (b) sub-cooling the returning liquid, and/or (c) increasing the operating pressure [6]. Enhancement of boiling performance can refer to increases in boiling heat transfer coefficient (HTC), increases in CHF, reduction in the boiling incipience temperature, or decreases in boiling hysteresis. This research focuses on surface modifications to improve the boiling HTC compared to a plain surface.

#### 1.2 BACKGROUND

The thermal properties of coolants are inherently poor and the idea of adding solids to fluids to enhance these properties was first described by Maxwell in 1904 [1]. The problem with using solid particles in fluids is stability; suspensions of µm and larger sized particles are likely to settle. However, it is thought that nanofluids, suspensions of nanometer sized particles dispersed in a base fluid specifically to enhance the thermal properties of the fluid [8], are more stable since the smaller particle sizes allow for a decrease in the sedimentation rate and charges developed on the particle surface provide electrostatic repulsion. Nanofluids have shown promise as an advanced coolant for a variety of applications. One application of this research is the improvement of boiling heat transfer for thermal management of power electronics that control the electric motor and the fuel cell stack in advanced hybrid vehicles.

Nanofluids can be created with particle materials, such as metals, oxides, and nitrides, in a range of base fluids, including water, propylene glycol, oils, and monomers. Nanofluids are often opaque and take on the color of the particle; for example, a white Al<sub>2</sub>O<sub>3</sub>/water-based nanofluid is shown in Fig. 1.2. Two common practices, the one-step and two-step methods, are used to produce nanofluids. The one-step method produces

metallic nanofluids by directly condensing an evaporated metal vapor into a base fluid to generate dispersed nanoparticles [9,10]. The two-step method first produces nanoparticles, typically by inert gas condensation, as a nanopowder, which is then dispersed into a base fluid [11]. Since oxide nanopowders are commercially available, the two-step method is simpler and less expensive for producing nanofluids. These methods allow for nanoparticles of a wide range of shapes, sizes, and types to be dispersed in aqueous and non-aqueous fluids. The amount of nanoparticles in a nanofluid is typically reported as the volume fraction of nanoparticles in the suspension, calculated from the density and mass of nanopowder dispersed. The size of the nanoparticle is usually reported as the size of the primary particle dispersed and can be measured using light scattering techniques.

Nanofluids were first described in 1995 [8] and since then their thermal properties have been widely studied. Most researchers have focused on measuring and modeling the thermal conductivity enhancements. The change in thermal conductivity of nanofluids compared to the base fluid has been widely studied and reviewed by Özering et al. [12]. There is a wide range in the measured increase in the thermal conductivity of nanofluids compared to the base fluid. For example, Lee et al. [13] reported a 10% increase of 38.4 nm Al<sub>2</sub>O<sub>3</sub> particles at 4.4% volume fraction in water, while Das et al. [14] measured a 24% increase for the same particles at 4.0% volume fraction. CuO nanoparticles have had the largest measured increase in thermal conductivity reported for metal oxide based nanofluids [15]. Despite the discrepancies, overall there is a substantial increase in thermal conductivity compared to the base fluid. There have been a variety of models proposed to explain the increase in thermal conductivity. With a

large collaboration over 30 different sites, Buongiorno et al. [16] determined that the increase in thermal conductivity could be modeled with the effective medium theory. Recently, some of the focus has shifted towards nanofluids in other applications, such as forced convection and boiling. In forced convection applications, as review by Godson et al. [17], there have been measured increases in heat transfer coefficient for nanofluids.

Since nanofluids have shown improved thermal properties, they have been studied for boiling applications. However, nanofluid boiling is not well understood. There has been a consensus that nanofluids improve CHF compared to water [18], but there are conflicting reports on nanofluids' influence on the boiling HTC. Taylor and Phelan [19] provide a comprehensive review of nanofluids in pool boiling. A summary of the nanofluid pool boiling literature is presented in Table 1.1 [19-40]. There have been several researchers [20-27] which have measured a decrease in HTC, several [19, 28-36] that have observed an increase and others [37-40] who did not see any changes. Shi et al. [31] measured a nearly 60% increase in HTC for Al<sub>2</sub>O<sub>3</sub>/water-based nanofluids on a copper block, while Bang and Chang [20] measured a 20% decrease in HTC for similar Al<sub>2</sub>O<sub>3</sub>/water-based nanofluids on a plain surface. It was noted in nearly every study that there was some kind of deposition of nanoparticles on the surface. Taylor and Phelan [19], Das et al. [21], and Narayan [30] all concluded that nanoparticles coating the surface was an important factor, with Das et al. [21] and Narayan et al. [30] defining a surface parameter as the ratio of particle size to roughness to explain the change in HTC. Chopkar et al. [28] boiled ZrO<sub>2</sub>/water nanofluids multiple times on the same surface and saw an increase in performance for the first two runs and a decrease after the third time.

Overall, the coating formed during nanofluid boiling and its impact on boiling HTC is not well understood.

For automotive power electronics and fuel cell thermal management applications, the coolant must also be electrically non-conductive. Strong dielectric coolants are typically in these applications. There has not been substantial research on the electrical conductivity of nanofluids. Several researchers have measured increases in [41-44], but a model to predict the increase and particle size effects are not well understood.

#### 1.3 RESEARCH OBJECTIVES

The focus of this research is on two novel methods to create deposits of nanoparticles from nanofluids on surfaces undergoing nucleate pool boiling and to study the electrical conductivity properties of nanofluids. There is lack of understanding on how the layer formed during nanofluid boiling impacts boiling performance. Alternative methods to coat boiling surfaces with nanofluids have not been studied. Therefore, electrophoretic deposition of nanoparticles is investigated. Furthermore, there is a need to explore how the electrical conductivity changes with particle size and concentration and how to model these effects.

The specific tasks in this research include:

(1) To quantify the impact of nanoparticles that attach to a surface during boiling.

The review of nanofluid boiling literature reveals two competing phenomena: surface coating of nanoparticles and bubble suppression due to the particles. To highlight the effects of the surface coating, a novel process is applied to create multiple coatings on a surface by alternating boiling studies between water and

- nanofluids. The surface is investigated to determine BHT enhancement compared to an uncoated surface.
- (2) To use electrophoretic deposition to coat surfaces and investigate their impact on boiling performance. Since BHT can be enhanced by coating surfaces with particles from boiling nanofluids and the nanoparticles in suspension have a surface charge, a novel method is used to enhance boiling surfaces with electrophoretic deposition of particles.
- (3) To measure and model the effect of particle size and concentration on the electrical conductivity. The electrical properties of nanofluids are not well understood, yet it is an important property of coolants in advanced thermal management systems. Three different particle sized nanofluids prepared with similar methods are experimentally investigated over a range of volume fractions. The electrical conductivity is modeled as a colloidal suspension a salt-free medium.

Completion of these objectives will provide a better understanding of how nanofluids can be used to improve the performance of advanced thermal management systems.

#### 1.4 ORGANIZATION OF THE DISSERTATION

This dissertation is presented in a multiple manuscript format. Chapters 2, 3, and 4 are written as individual research papers including the abstract, main body, and references.

Chapter 2 studies the impact of nanoparticle layers formed on a surface while boiling nanofluids. Chapter 3 studies the impact of nanoparticle layers created from

electrophoretic deposition of particles from nanofluids. Chapter 4 investigates the electrical conductivity of non-aqueous nanofluids. Finally, Chapter 5 provides the conclusions of this research and summarizes original contributions of the dissertation. Further, several topics are proposed for future research.

Table 1.1. Review of nanofluid boiling literature

Author(s)	Heating Surface	Nanofluid	Change in HTC	Particle Deposition
Bang and Chang [20]	4 mm x 100 mm	Al <sub>2</sub> O <sub>3</sub> /water	Decrease, up to 20%	Yes
Das et al. [21]	Different roughness cartridge heaters	Al <sub>2</sub> O <sub>3</sub> /water	Decrease, 10-40%	Yes
Jackson et al. [22]	Cu block	Au/water	Decrease, up to 25%	Yes
Kathiravan et al. [23]	Stainless steel (30 x 30 mm)	Cu/water	Decrease, up to 75%	Yes
Kim et al. [24]	stainless steel wire ( $\emptyset$ 0.38 mm)	Al <sub>2</sub> O <sub>3</sub> , ZrO2, SiO <sub>2</sub> /water	Decrease	Yes
Liu et al. [25]	Cu bar (Ø20 mm)	SiO <sub>2</sub> , CuO/water	Decrease	Yes
Milanova and Kumar [26]	NiCr wire	SiO <sub>2</sub> /water	Decrease	Yes/no, depends on particle charge
Sajith et al. [27]	NiCr wire (Ø0.19 mm)	Al <sub>2</sub> O <sub>3</sub> , Cu /water	Decrease	Yes
Chopkar et al. [28]	Cu plate (Ø50.8 mm)	ZrO <sub>2</sub> /water	Increase	Yes
Coursey and Kim [29]	Cu block	$Al_2O_3$	Increase, up to 50%	Yes
Narayan et al. [30]	Vertical stainless steel tube	Al <sub>2</sub> O <sub>3</sub> /water	Increase, up to 70% Decrease, up to 20%	Yes
Shi et al. [31]	Cu bar (Ø60 mm)	Al <sub>2</sub> O <sub>3</sub> /water	Increase, up to 60%	Yes
Taylor and Phelan [19]	Wire (Ø0.26 mm)	$Al_2O_3$	Increase, up to 40%	Yes
Truong [32]	Stainless steel wire	Al <sub>2</sub> O <sub>3</sub> /water	Increase, up to 68%	Yes
Tu et al. [33]	Ti heater (26 x 40 mm)	Al <sub>2</sub> O <sub>3</sub> /water	Increase, up to 64%	Yes
Wen and Ding [34]	Stainless steel disc (Ø150 mm)	Al <sub>2</sub> O <sub>3</sub> /water	Increase, 40%	Yes
Wen et al. [35]	Stainless Steel disc (Ø150 mm)	TiO <sub>2</sub> /water	Increase, up to 50%	No
Witharana [36]	Cu bar (Ø100 mm)	Au, SiO <sub>2</sub> /water	Increase	Not studied
Kim et al. [37]	Cu surface	Al <sub>2</sub> O <sub>3</sub> /water	Negligible	Yes
Kwark et al. [38]	Cu block (1 x 1cm)	Al <sub>2</sub> O <sub>3</sub> /water	Negligible/Slight Degradation	Yes
Vasallo et al. [39]	NiCr wire (Ø0.4	SiO <sub>2</sub> /water	Negligible	Yes
You et al. [40]	mm) Cu surface	Al <sub>2</sub> O <sub>3</sub> /water	Negligible	Not studied

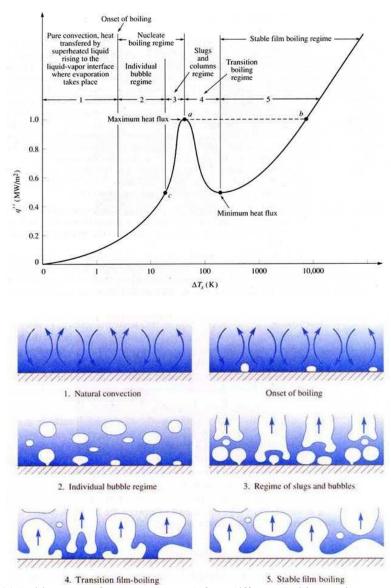


Figure 1.1. (a) Boiling curve for water through four different boiling regimes as a function of superheat and (b) visualization of the different boiling regimes. [5]



Figure 1.2. Al<sub>2</sub>O<sub>3</sub>/water based nanofluid.

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#### **CHAPTER 2**

## EFFECTS OF NANOPARTICLE LAYERING ON NANOFLUID AND BASE FLUID POOL BOILING HEAT TRANSFER FROM A HORIZONTAL SURFACE UNDER ATMOSPHERIC PRESSURE

#### **ABSTRACT**

Previous heat transfer studies of nanofluids have shown that suspended nanoparticles can affect thermal properties within a fluid and furthermore can affect surface roughness by depositing on a heater surface. Pool boiling studies of nanofluids have demonstrated either enhanced or diminished heat transfer, yet have been unable to distinguish the contributions of increased surface roughness and suppression of bubble transport by suspended particles, because they have used base fluids on a clean boiling surface as a comparison. We resolve this uncertainty by studying the boiling performance of a surface exposed to a series of boiling tests that alternate between water and a waterbased nanofluid with suspended 40 nm ZnO nanoparticles. We find that the performance for the water tests increases significantly, showing a 62% enhancement after four cycles. This increase correlates well with a surface roughness model for boiling that uses AFMmeasured surface data to quantify the layering of nanoparticles in intervening nanofluid boiling tests. We find that the performance of the ZnO nanofluid initially shows a 24% enhancement versus water on a clean (unroughened) surface, but then steadily declines in later tests as nanoparticle layering occurs, showing a measured trend that is opposite that of water. We ascribe this decrease to the suppression of bubble formation and motion by the suspended particles. The results demonstrate that the effect of increased surface

roughness due to nanoparticle layering can be twofold, greatly enhancing boiling for the base fluid and slightly decreasing performance for the nanofluid.

#### 2.1 INTRODUCTION

Boiling heat transfer is important to applications such as power electronics cooling, enabling high performance thermal management systems such as thermosyphons, heat pipes, and spray cooling. The boiling performance of a system is affected by the properties of the fluid, such as the thermal conductivity, surface tension, and density; the properties of the boiling surface, such as the material, geometry, orientation, wettability, and surface roughness; the properties of the system, such as system pressure; the fluid-surface interaction; and the mechanisms of bubble growth, departure, and movement [1]. Previous research has aimed to enhance boiling by manipulating the fluid and the surface.

One method that has been proposed to increase fluid heat transfer performance is using nanoparticle suspensions known as nanofluids, which have demonstrated improved thermal conductivity [2-8] and forced convection heat transfer coefficient [9,10] compared to the base fluid. However, in boiling studies, prior results have been inconclusive regarding whether measured enhancements in boiling coefficient are due to surface roughening caused by nanoparticle layering or the effect of nanoparticle motion on bubble nucleation and transport.

It is well known that surface microstructure can heavily influence boiling properties [1]. Microporous coatings [11,12], modulated surfaces [13], and fabricated microstructures [14] have all been used to enhance boiling. It has also been shown that suspended particles may obstruct bubble movement and liquid stirring, leading to a

reduction in boiling performance [15]. These two processes in nanofluid boiling may therefore have opposite effects on the boiling coefficient, as supported by previous studies in which boiling performance initially increases and then decreases as the volume fraction of suspended particles increases [15-17].

Several studies have investigated the boiling of nanofluids on a heated wire and attributed the measured increase in boiling performance to the increase in surface roughness caused by the attachment of the nanoparticles to the wire [18,19]. Kim et al. [20] ascribed measured increases in the critical heat flux to an improved surface wettability caused by the layering of nanoparticles, but also measured a decrease in nucleate boiling coefficient that they attributed to surface effects. For nanofluid boiling on cylindrical and flat surfaces, results have been conflicting. Several researchers have measured an increase in boiling heat transfer using nanofluids [21-23], while others have seen a deterioration of performance [24,25]. Narayan et al. [26] measured a decrease in performance for smooth surfaces ( $R_a$ =48 nm) and an increase in performance for rough surfaces ( $R_a$ =524 nm). In an attempt to reconcile these conflicting measurements, they introduced a surface interaction parameter defined as the ratio of the surface roughness to the particle size, and correlated this parameter with measured boiling performance. Das et al. [26,27] also applied this surface interaction parameter to their data as well as other data reported in the literature. They theorized that the number of nucleation sites is reduced if the roughness of the surface is close to the particle diameter, and increased if the size ratio is larger.

Previous studies, such as those outlined above, have used base fluids on a clean surface as a comparison when investigating nanofluid boiling. Because of this, they have

been unable to distinguish the potential contributions of nanoparticle layering on the boiling surface from the potential suppression of bubble nucleation and motion by suspended particles.

In this study, we use a series of boiling tests that alternate between water and a water-based nanofluid to resolve existing uncertainties regarding competing physical mechanisms in nanofluid boiling. By comparing the performance of water and nanofluid boiling on the same roughened surface, we are able to separate the effect of suspended particles from the effect of surface roughness. We use surface roughness data measured by atomic force microscopy (AFM) in conjunction with a surface roughness correlation model to explain measured improvements in boiling coefficient of water for successive layering of nanoparticles on the boiling surface.

#### 2.2 EXPERIMENTAL SETUP

This section provides details regarding the nanofluids used in the experiments, the design of the test apparatus, and the AFM imaging used to analyze the boiling surface.

## 2.2.1 Nanofluid Preparation

We used deionized (DI) water-based nanofluids with 40nm diameter ZnO nanoparticles. The volume fraction of the ZnO nanoparticles was 2.3%, and the nanofluids contained no dispersant; they were prepared by Nanophase Technologies (Romeoville, IL). The particles have a specific surface area of 33 m<sup>2</sup>/g and a crystal phase of zincite (hexagonal) with an elongated morphology.

### 2.2.2 Boiling Apparatus

The system, as shown in Figure 2.1, consists of a heated surface where boiling occurs, a housing to contain the fluid, and a condensing section. The heat source is a half-sheath, chrome-steel disc heater with a diameter of 82.6 mm and power of 600 W (~120 kW/m<sup>2</sup>). The heater is controlled by a Staco Energy Products Co. (Dayton, OH) 0-120 V, 10 A, variable AC power supply. The heater is mated to a Grade 316L stainless steel disc, 82.6 mm in diameter and 12.7 mm thick, using DOW Corning® 340 silicone heat sink compound to improve heat transfer across the interface. This assembly is embedded in a thermal insulator housing composed of polytetrafluoroethylene (PTFE) 114.3 mm in diameter and 114.3 mm thick. This PTFE housing is packed with glass wool insulation. The working fluid is contained in a polycarbonate housing that has a chamber 88.9 mm in diameter and 114.3 mm tall. During the experiments the vapor condensed onto the surface of an aluminum chamber at the top. The temperature of the boiling surface was extrapolated from the measurements of 3 type-T thermocouples embedded 6.35 mm below the surface and the heat flux into the system. The saturation temperature was measured in the fluid using a type-T thermocouple suspended from the condenser. This configuration is similar to that used by Wen and Ding [23]. The instrumentation was interfaced with a National Instruments CompactDAQ data acquisition system with LabVIEW 8.20 software.

#### 2.3 EXPERIMENTAL PROCEDURE

#### 2.3.1 ZnO layering study

For the boiling tests, heat fluxes of 24, 28, and 32 kW/m<sup>2</sup> were used by setting the input voltage to 60, 65, and 70 V. The system was allowed to reach steady state, defined

as when the surface temperature did not change by more than 0.1°C for 20 minutes. First water was boiled on a clean surface until it reached steady state (approximately 45 minutes). The water was then removed and the ZnO nanofluid was boiled until it also reached steady state (approximately 45 minutes). The ZnO nanofluid was drained, and then without cleaning the surface of deposited nanoparticles, fresh DI water was used as the boiling fluid. This process was repeated to create four layers of ZnO on the boiling surface, with measurement of the boiling coefficient taken during each boiling event. The cycling study is described in Table 2.1.

# 2.3.2 Scanning electron microscopy (SEM) and atomic force microscopy (AFM) of nanoparticle layers

To investigate the heated surface after each boiling cycle,  $12.7 \times 25.4 \times 3.2$  mm stainless steel tabs were fitted to the stainless steel boiling surface using 2.8 mm bolts. After each cycle, one tab was removed and replaced by a new clean tab. The tabs were removed in a sequential order, numbered accordingly, and then imaged using a Hitachi S-4800 Field Emission SEM at 3,000x and 100,000x magnification. To measure the surface roughness of each tab, a Molecular Imaging PicoLE Scanning Probe Microscope with an AFM head and Molecular Imaging PicoScan software was used. The tapping mode was used to image the surface topography over a 5  $\mu$ m x 5  $\mu$ m area. The average surface roughness,  $R_a$  was calculated using Gwyddion 2.16 data visualization and analysis software. AFM malfunction prevented the measurement of surface roughness for the sample from Test 5.

#### 2.4 BOILING PERFORMANCE CHARACTERIZATION

This section provides details on the analysis and uncertainty used to calibrate the system

and compare different tests.

## 2.4.1 Analysis

The heat flux at the heated surface, q, was calculated from the measured total heat into the system, Q, and the heated area, A. The power input,  $Q=V^2/R$ , was calculated from the applied voltage, V, and resistance, R, of the heater. The heat flux is given by:

$$q = \eta \frac{V^2}{RA} \tag{2.1}$$

where  $\eta$  is the fraction of energy that reaches the heated surface after accounting for losses to the environment. The superheat,  $\Delta T_s$ , is the difference between the measured saturation temperature,  $T_{sat}$ , and the boiling surface temperature,  $T_w$ , which was extrapolated from the heat flux, q, and the temperature,  $T_m$ , measured by the thermocouple a distance of l=6.35 mm below the boiling surface:

$$\Delta T_s = T_w - T_{sat} \tag{2.2}$$

$$T_w = T_m - ql/k \tag{2.3}$$

where k = 17 W/mK is the thermal conductivity of stainless steel. The boiling coefficient, h, was used to characterize the boiling:

$$h = \frac{\Delta T_s}{q} \tag{2.4}$$

For pure water undergoing nucleate pool boiling, Rohsenow [28] developed the relationship:

$$\frac{c_p(T_w - T_{sat})}{h_{fg} \operatorname{Pr}^s} = C_{sf} \left[ \frac{q}{\mu h_{fg}} \sqrt{\frac{\sigma}{g(\rho - \rho_v)}} \right]^{1/3}$$
(2.5)

where  $c_p$  is the fluid's specific heat,  $h_{fg}$  is the latent heat of vaporization,  $\mu$  is the fluid viscosity, Pr is the Prandtl Number defined as  $\Pr = \mu c_p/k$ , k is the fluid thermal conductivity,  $\sigma$  is the fluid surface tension,  $\rho$  is the fluid density,  $\rho_v$  is the vapor density, and g is the acceleration due to gravity. For a stainless steel surface and water [1], the fluid surface interaction constants are  $C_{sf} = 0.013$  and s = 1. This relationship was used in a calibration measurement to determine the value of  $\eta$  for the test apparatus. Figure 2.2 shows the Rohsenow correlation and experimental data for DI water for the test apparatus. Least-squares analysis was used find the value of  $\eta$  that minimized the error between the experimental results and the Rohsenow correlation, yielding  $\eta = 0.83$ .

## 2.4.2 Uncertainty

To quantify the uncertainty associated with the measured heat flux, superheat, and boiling coefficient, the sum of squares method was used. The uncertainty  $u_q$  of the heat flux q as described in Eq. (2.1) can be written as:

$$u_{q} = \sqrt{\left(\frac{\delta q}{\delta V}u_{V}\right)^{2} + \left(\frac{\delta q}{\delta R}u_{R}\right)^{2} + \left(\frac{\delta q}{\delta A}u_{A}\right)^{2}}$$
(2.6)

where  $u_V$  is the uncertainty of V,  $u_R$  is the uncertainty of R, and  $u_A$  is the uncertainty of A. The error associated with the voltage measurement was  $\pm 0.2$  V. The resistance of the heater was measured to be 24  $\Omega$   $\pm$  2% over the temperature range of this study. The diameter of the stainless steel disk was machined to 82.6  $\pm 0.03$  mm; therefore the area of

the disk is 6207  $\pm 4$  mm<sup>2</sup>. Using Eq. (2.6), the maximum heat flux uncertainty  $u_q/q$  was calculated to be less than 3%.

The uncertainty of the superheat,  $\Delta T_s$ , as defined in Eq. (2.2) is  $u_s$ :

$$u_{s} = \sqrt{\left(\frac{\delta \Delta T_{s}}{\delta T_{m}} u_{T_{w}}\right)^{2} + \left(\frac{\delta \Delta T_{s}}{\delta k} u_{k}\right)^{2} + \left(\frac{\delta \Delta T_{s}}{\delta L_{i}} u_{L_{i}}\right)^{2} + \left(\frac{\delta \Delta T_{s}}{\delta q} u_{q}\right)^{2} + \left(\frac{\delta \Delta T_{s}}{\delta T_{m}} u_{T_{sat}}\right)^{2}}$$
(2.7)

where  $u_{Tw}$  is the uncertainty of  $T_{w}$ ,  $u_{Tsat}$  is the uncertainty of  $T_{sat}$ ,  $u_{k}$  is the uncertainty of k, and  $u_{Li}$  is the uncertainty of  $L_{i}$ . The error due to the thermocouple measurements was calculated by using the sum of squares method with the design stage uncertainty and the measurement uncertainty. The design stage uncertainty was 0.1 K and the measurement uncertainty was calculated as twice the standard deviation of the temperature measurements (typically  $\leq 0.1$  K). The thermal conductivity of stainless steel was assumed constant over the temperature range of the disk. The distance between the thermocouple and the heater surface had an uncertainty of  $\pm 0.03$  mm. Using Eq. (2.7), the maximum superheat uncertainty  $u_{s}/\Delta T_{s}$  was calculated to be less than 3%.

The uncertainty of the boiling coefficient,  $u_h$ , can be written as:

$$u_{h} = \sqrt{\left(\frac{\partial h}{\partial \Delta T_{s}} u_{s}\right)^{2} + \left(\frac{\partial h}{\partial q} u_{q}\right)^{2}}$$
(2.8)

The maximum boiling coefficient uncertainty  $u_h/h$  was calculated to be less than 10%.

#### 2.5 BOILING PERFORMANCE RESULTS

First we compare the boiling performance of the ZnO nanofluid (Test 1) and water on a clean surface (Test 0). Figure 2.3 shows the demonstrated enhancements of the nanofluid, for which the boiling coefficient is increased over that of water by 24.8%, 16.1%, and 17.7% at input heat fluxes of 24, 28, and 32 kW/m², respectively. Based on the analysis below, we ascribe these increases to the enhanced thermal properties of the nanofluid, such as the thermal conductivity, rather than the surface roughness produced by nanoparticle layering.

Next we examine the boiling performance of water on a surface that has been successively modified by intervening cycles of nanofluid boiling. Figure 2.4 shows the measured increases in boiling coefficient (h) for each input heat flux as ZnO nanoparticles are successively coated over a number of trials. After the first coating, there is little increase of h from the baseline (Test 0), but a trend of increasing h becomes apparent in later cycles as more layers of ZnO particles build up on the surface. A 62% increase in the boiling coefficient for water from the first to the last cycle is measured. In the next section we analyze the origins of this increase by examining the effect of successive ZnO nanoparticle layers on surface roughness.

The boiling performance of ZnO nanofluids is also compared for successive tests. Figure 2.5 shows an initial measured increase in h for all heat fluxes when compared to the baseline (Test 0) of water on a clean surface, as mentioned above. The performance then slightly decreases and remains nearly constant for each additional cycle of nanofluid boiling. The origin of this behavior is discussed below in the context of the coating of the surface.

## 2.6 SEM AND AFM INVESTIGATION OF BOILING SURFACE AND DISCUSSION

Scanning electron microscopy (SEM) images of the sample surface after each of the four nanofluid boiling tests are shown in Figure 2.6. The first image, taken after Test 1, shows a layer of ZnO nanoparticles that does not completely cover the surface, with islands of nanoparticle coverage visible in the 3,000x image (circled in red). We ascribe the relatively small change in boiling coefficient relative to the baseline measured for Test 2 (see Figure 2.4) to this incomplete coverage. Later nanofluid boiling tests yield a completely coated surface, with the 100,000x SEM images showing that the nanoparticles have a similar size distribution on the surface. For these surfaces, we see significantly increased boiling performance.

To explain this measured increase in water boiling performance, we use AFM to quantify the surface topography. Figure 2.7 shows AFM images of the sample surface after several nanofluid boiling tests, from which the surface roughness parameter,  $R_a$ , was determined. The arithmetic average roughness grew after each successive nanofluid boiling test, increasing from an initial value of 0.06  $\mu$ m for the baseline (Test 0) to 0.44  $\mu$ m after the final test.

Gorenflo developed a correlation [29] of boiling coefficient based on the reduced pressure of the fluid and roughness of the boiling surface:

$$h_b = h_0 F_{PF} (q / q_0)^n (R_a / R_{a0})^{0.133}$$
(2.9)

where  $h_0$  is the fluid specific heat transfer coefficient,  $q_0$  is the fixed reference conditions for the heat flux, and  $R_{a0}$  is the reference value for the surface roughness. The pressure

correction factor,  $F_{PF}$ , and the heat flux exponent, n, are functions of the reduced pressure,  $p_r$ . This equation is valid for  $0.005 \le p_r \le 0.95$ .

A normalized boiling coefficient,  $\bar{h}_b$ , can be defined for water on each nanoparticle-roughened surface at each heat flux by dividing by the boiling coefficient of water on the initial clean (Test 0, i.e. unroughened) surface at the same heat flux. Here we note an important advantage of our measurement: because we use water in each case,  $h_0$  is expected to remain constant and thus cancel out in this normalization. Previous studies, which compare nanofluid boiling on roughened surfaces to base fluid boiling on clean surfaces, are unable to distinguish the effects of the suspended nanoparticles on surface roughness ( $R_a$ ) and fluid specific heat transfer coefficient ( $h_0$ ).

We therefore adapt the Gorenflo correlation to yield:

$$\overline{h}_{b} = \frac{h_{b}}{h_{b-clean}} = \left(R_{a} / R_{a-clean}\right)^{0.133}$$
 (2.10)

where  $h_{b\text{-}clean}$  is the boiling coefficient of water on the clean surface and  $R_{a\text{-}clean}$  is the surface roughness of the initial clean (Test 0) surface. Figure 2.8 compares the measured data for the normalized boiling coefficient of water (averaged for the three different heat fluxes) to the prediction of Eq. (2.10).

The model accounting for surface roughness somewhat under-predicts the measured enhancements in boiling coefficient, but still falls within the experimental error. We note that the model does not account for effects such as nanoparticle layer porosity, which is expected to contribute to boiling coefficient enhancement. The model also does not account for the finite thermal conductance of the ZnO nanoparticle layer, which may yield an effective surface temperature that is slightly lower than that extrapolated using

Eq. (2.3). Finally, the model does not account for the fluidization of surface-layered nanoparticles, which may occur during water boiling.

For comparison purposes, we also calculate the averaged normalized boiling coefficient of the ZnO nanofluid. As for the water, we divide the measured boiling coefficient at each heat flux by the boiling coefficient of water on the initial clean (Test 0) surface at the same heat flux. As shown in Figure 2.8, we measure an increase in boiling coefficient of 24% for the incompletely covered surface (Test 1), and then a slight decrease for later tests with a completely coated surface. This suggests that the initial increase is independent of surface roughness and due to thermal properties of the nanofluid, such as increased thermal conductivity, which has been measured for ZnO nanofluids [6]. A portion of the increase could also be due to the disruption of the thermal boundary layer by suspended nanoparticles [15].

For later nanofluid boiling tests, the interaction of the nanoparticles with the coating counteracts not only the enhanced thermal properties of the nanofluid but also the mechanism of enhanced boiling coefficient seen for the water in Tests 2, 4, 6, and 8 (see Figure 2.4). Noting the relatively high volume fraction of the nanofluids and the horizontal flat boiling surface geometry, we attribute this reduced performance for the fully coated surface to the blocking of bubble transport and liquid stirring by the particles in suspension, as described by Iida *et al.* [15] for fluidized beds.

Note that an alternative formulation of the normalized boiling coefficient for the nanofluid could divide measured data by the Test 1 value (rather than Test 0) in order to isolate the effect of the surface/nanofluid interaction for different surface roughnesses. This alternative formulation would simply move the curve down in Figure 2.8 and not

change the trend of reduced performance for increasing surface roughness. We choose the Test 0 normalization in order to present water and nanofluid data on a common basis.

#### 2.7 CONCLUSIONS

The boiling performance of a heated surface exposed to a series of boiling tests that alternate between water and a water-based nanofluid with suspended ZnO nanoparticles has been studied. The creation of a nanoparticle layer on the heated surface is shown to dramatically increase the boiling performance of water and correlate well with a model that uses AFM-measured surface roughness data. Nanofluid boiling performance is shown to increase initially (compared to water) before the deposited nanoparticle layer is completely formed; we attribute this increase to enhanced thermal properties of the nanofluid. After the nanoparticle layer fully covers the heater surface, the nanofluid boiling coefficient declines, which we ascribe to a suppression of bubble nucleation and transport by the suspended nanoparticles.

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**Table 2.1. Cyclic Boiling Experiment Order** 

Test No.	Fluid Boiled
0	DI water (baseline)
1	ZnO nanofluid
2	DI water
3	ZnO nanofluid
4	DI water
5	ZnO nanofluid
6	DI water
7	ZnO nanofluid
8	DI water

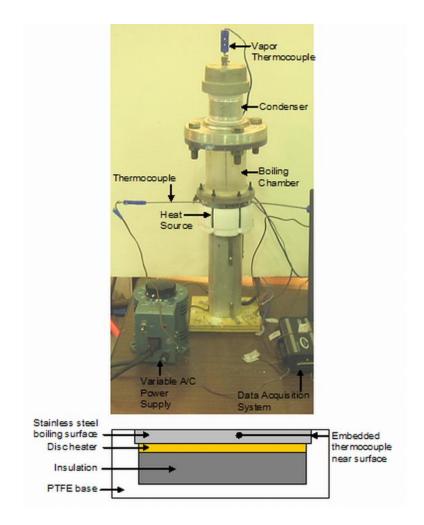


Figure 2.1. Picture of boiling system and schematic of heat source.

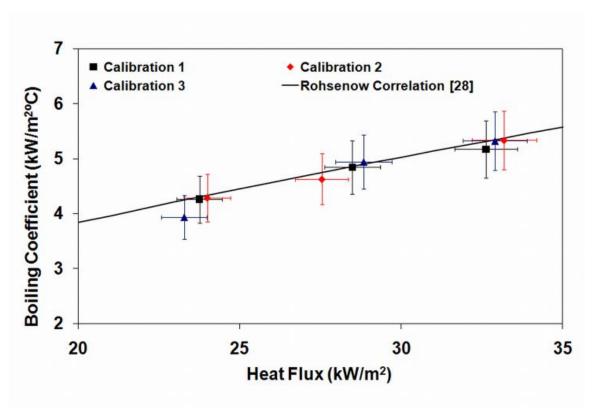


Figure 2.2. Calibration data for water compared to the Rohsenow [28] correlation showing good agreement with a heat transfer efficiency  $\eta$ =0.83.

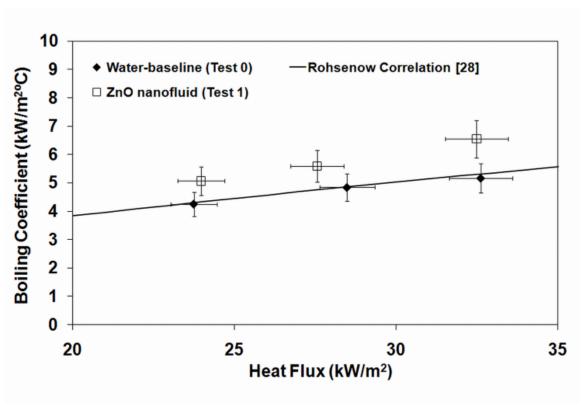


Figure 2.3. Experimental data showing that ZnO nanofluids have a larger boiling coefficient at the same heat flux when compared to the boiling of water on the same clean (unroughened) surface.

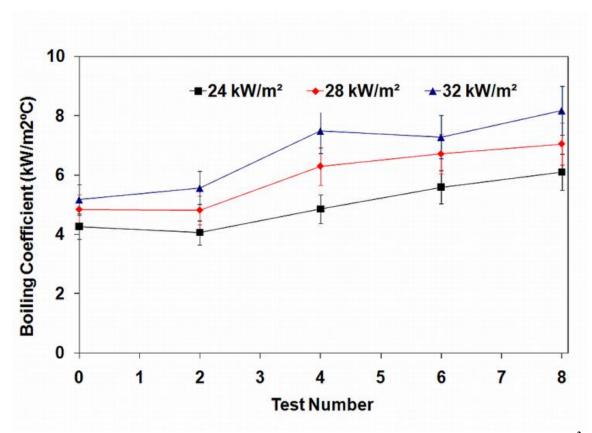


Figure 2.4. Experimental data showing that for each heat flux into the system  $(24, 28, \text{ and } 32 \text{ kW/m}^2)$  the boiling coefficient of <u>water</u> increases with successive coatings of ZnO nanoparticles.

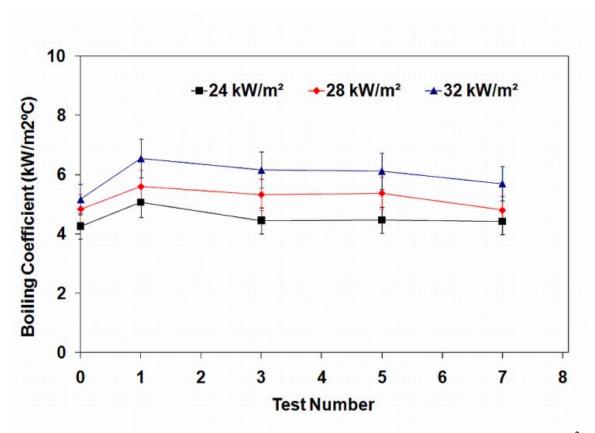


Figure 2.5. Experimental data showing that at each heat flux into the system (24, 28, and 32 kW/m²) the boiling coefficient of the ZnO nanofluid initially increases from the value for water on a clean surface (Test 0) then remains nearly constant after successive coatings of ZnO nanoparticles.

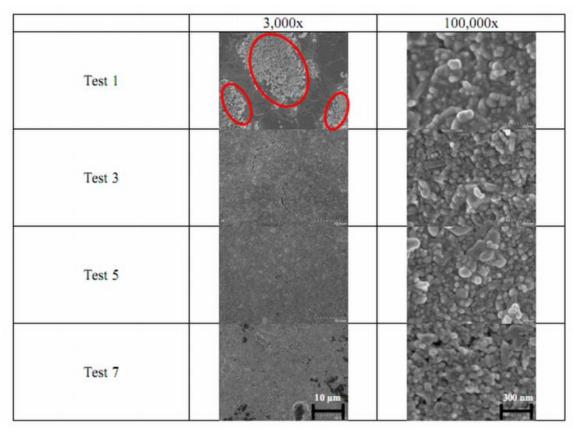


Figure 2.6. SEM images of deposited ZnO nanoparticle layer on stainless steel heater surface after the referenced boiling test. Test 1 yields incomplete surface coverage, while subsequent nanofluid boiling tests yield full coverage.

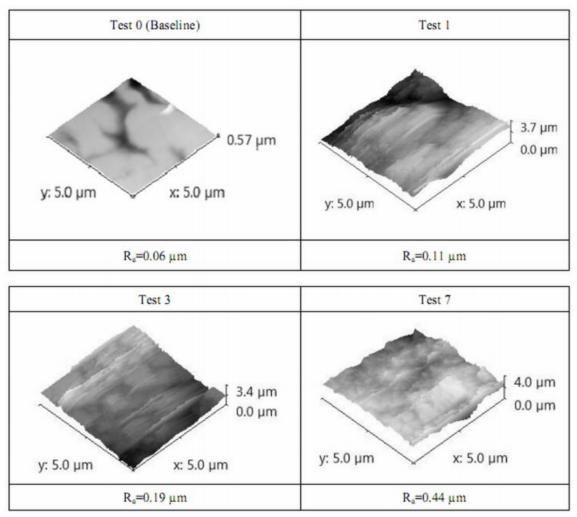


Figure 2.7. AFM images of heater surface samples after the referenced boiling test, showing increasing surface roughness.

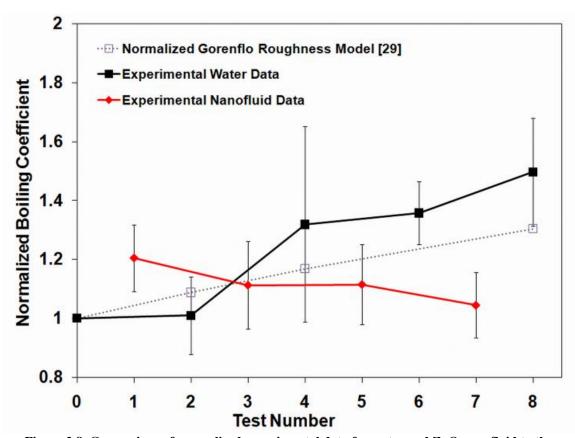


Figure 2.8. Comparison of normalized experimental data for water and ZnO nanofluid to the normalized Gorenflo [29] model which accounts for surface roughness. Water boiling tests show significantly increasing performance due to increasing surface roughness, while nanofluids show moderately decreasing performance that we ascribe to suppression of bubble transport by suspended nanoparticles.

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#### **CHAPTER 3**

# BOILING SURFACE ENHANCEMENT BY ELECTROPHORETIC DEPOSITION OF PARTICLES FROM A NANOFLUID

#### **ABSTRACT**

Boiling heat transfer can be enhanced with a variety of surface treatments. One method that has not been widely studied is electrophoretic deposition of nanoparticles from a nanofluid. A stainless steel surface was coated using electrophoretic deposition of 40 nm ZnO particles from a ZnO-propylene glycol based nanofluid. With adequate coating time, such a surface modification method can increase the boiling heat transfer coefficient by about 200%. A surface with improved performance had an increase of active nucleation site density by 2.9 times as determined by the Mikic-Rohsenow Correlation. This was supported using an analysis of the cavities present in a SEM image of the surfaces which showed a 2.4 times increase in cavities.

#### 3.1 INTRODUCTION

Boiling heat transfer (BHT) is the basis for thermal management in several industries including refrigeration, power systems and electronics. Among a number of strategies to enhance boiling heat transfer, surface modification has proven to be an attractive means to improve the boiling heat transfer coefficient (HTC). Increases of 100% or more in HTC have been demonstrated by simply roughening the surface [1] and through the use of microscale pin fins [2-4], microporous coatings [3], microchannels [5], and micro-drilling of nucleation cavities [6]. Nanoscale surface modification through the

growth of carbon nanotubes [7] and nanorods [8] have demonstrated HTC improvements of 450% and 3000% respectively, while combinations of micromachining and carbon nanotube growth have demonstrated improvements of nearly 900% [9].

Although nanostructuring or microstructuring of a boiling surface has been shown to lead to large improvements in performance, in practice the methods used to fabricate such surface structures (e.g. chemical vapor deposition or micromachining) are often inconvenient or expensive. Surface treatments utilizing nanofluids offer a potentially convenient means for surface modification; a number of studies have investigated the boiling of nanofluids as a means to coat a surface with nanoparticles for enhanced HTC or critical heat flux [10, 11]. However, results of this method have been mixed [12] presumably due to the fact that the deposition process during boiling is not well understood and difficult to control.

Electrophoretic deposition (EPD), on the other hand, is used in a variety of industrial applications, including automotive and appliance coatings, and can be implemented with a high level of automation [13]. In this process, charged particles are attracted to a conducting surface due to an applied electric field. EPD coatings adhere better at high densities than dipped or sprayed coatings, and EPD process parameters such as voltage and duration are relatively simple to control.

Recognizing that nanoparticles in suspension (nanofluids) have a surface charge, several previous studies have examined the EPD coating of surfaces with Al<sub>2</sub>O<sub>3</sub> [11], TiO<sub>2</sub> [14], and ZnO [15] nanoparticles. However, heat transfer measurements have shown only marginal improvements due to poor coverage or low adherence of the coating.

Here we explore the use of high particle volume fraction nanofluids in EPD coating as a means to overcome the coverage and adherence difficulties associated with prior nanofluid EPD studies and demonstrate a significant improvement in boiling performance. An electric field is applied to coat stainless steel (SS) surfaces with ZnO nanoparticles using propylene glycol (PG) based nanofluids to create enhanced boiling surfaces. The surface roughness and contact angle of the surfaces are compared and the surfaces investigated using SEM. The improved performance of the EPD-deposited surface is shown to be related to the increase in the nucleation site density, which is estimated using both the Mikic-Rohsenow correlation and SEM image analysis of the surface.

#### 3.2 EXPERIMENTAL METHOD

## 3.2.1 Nanofluid preparation and characterization

PG-based nanofluids were prepared by Nanophase Technologies (Romeoville, IL) with 40 nm diameter ZnO nanoparticles at 7% volume fraction with no dispersants. The particles had a specific surface area of 33 m $^2$ /g and a crystal phase of zincite (hexagonal) with an elongated morphology. The zeta potential was measured to be 48.6  $\pm$  0.1 mV using a Brookhaven Instruments (Holtsville, NY) ZetaPlus Zeta Potential Analyzer.

## 3.2.2 Coating procedure

1 inch diameter disks made from 20 gauge 304 SS were used as boiling surfaces. To create a consistent surface finish after the machining process, the disk surfaces were roughened with 400 ANSI grit sandpaper. Two SS disks acting as electrodes were separated by a 6.4 mm Teflon spacer and connected through wires to a 0-30V DC power

supply (B&K Precision, Yorba Linda, CA, Model 1660). Since a positive zeta potential was measured for the particles, the negative voltage terminal was attached to the surface to be coated. The electrode assembly was submerged in the nanofluid for a given amount of time while voltage was applied (see Fig. 3.1). Four samples (denoted as Surfaces #1 to #4) were coated by applying a voltage of 20 V for a range of times, as summarized in Table 3.1. The applied voltages and times were similar to those given in Santillan et al. [14] and Miao et al. [15] in which coatings were created, but the boiling performance was not studied. Coated samples were allowed to dry for a 24 hour period before boiling tests.

## 3.2.3 Boiling characterization

Experiments were conducted in an atmospheric pressure boiling test apparatus (Fig. 3.2) consisting of heating, evaporating, and condensing sections. One-dimensional heat flow was provided to the boiling surface through a 1 inch diameter, 10 inch long copper rod with an embedded cartridge heater (Fast Heat Inc., Elmhurst, IL) that was controlled by a variable AC power supply (Staco Energy Products, Inc., Dayton, OH). To limit radial heat losses, the copper rod was encased by a 2.5 inch diameter, 10 inch long polytetrafluoroethylene (PTFE) cylinder. The PTFE insulator had a 4 inch diameter, 1 inch thick shoulder with a groove for an O-ring, which was used to create a seal between the heating and evaporating sections. The coated samples were mounted to the end of the copper rod using a fastener.

The evaporating section consisted of a 200 ml transparent polycarbonate housing that was bolted to the PTFE heating section and sealed with an O-ring. The condensing section had a water-cooled aluminum finned condenser bolted and sealed at the top of the

fluid housing. The condenser had an opening to allow for a thermocouple and to vent the system to atmospheric pressure.

The temperature of the heat source was calculated using values measured by three type-E thermocouples (Omega Engineering, Inc., Stanford, CT, Model SA1-E-SRTC) that were spaced 0.75 inches apart axially, with the top thermocouple located 1.75 inches below the top (sample) surface. Temperatures were recorded (using an Omega OMB-DAQ-56 data acquisition system) until the system reached steady state, which was defined as a change of less than 0.2 K in a 20 minute period. This was typically achieved in approximately 45 minutes. Once at steady state, temperatures were recorded for an additional 10 minutes. A type-T thermocouple (Omega Model TQSS-18G-6) was suspended from the condensing section to monitor the vapor temperature. The heat flux and surface temperature were calculated from the temperature measurements.

Deionized (DI) water was used as the working fluid. The boiling performances of Surfaces #1 - #3 were measured at a heat flux of 181 kW/m<sup>2</sup>. The boiling performance of the baseline surface and Surface #4 were measured over a range of 160-240 kW/m<sup>2</sup>. The experiment was performed twice for each surface: once after the surface was coated ("1<sup>st</sup> run") and afterwards with fresh DI water ("2<sup>nd</sup> run"). Surfaces were not cleaned between these tests.

## 3.2.4 Calculation of heat flux and measurement uncertainty

The heat flux (q) of the system is defined as:

$$q = -k\frac{dT}{dx} = k\frac{T_3 - T_1}{L} \tag{3.1}$$

where  $T_3$  is the temperature measured at location 3 (furthest from the surface, as shown in Fig. 3.2),  $T_1$  is the temperature at location 1 (closest to the surface), k is the thermal conductivity of copper, and L is the distance between the first and third thermocouples. In this study, k = 396 W/m·K and L = 1.5 inches.

Using the root sum of squares method [16], the uncertainty of  $q(u_q)$  is defined as:

$$u_q = \sqrt{\left(\frac{\delta q}{\delta T_1} u_{T_1}\right)^2 + \left(\frac{\delta q}{\delta T_3} u_{T_3}\right)^2 + \left(\frac{\delta q}{\delta k} u_k\right)^2 + \left(\frac{\delta q}{\delta L} u_L\right)^2}$$
(3.2)

where  $u_{TI}$  is the uncertainty of  $T_1$ ,  $u_{T3}$  is the uncertainty of  $T_3$ ,  $u_k$  is the uncertainty of k due to temperature dependence, and  $u_L = 0.005$  inches is the uncertainty of L associated with manufacturing tolerances.

The superheat ( $\Delta T_s$ ) is the difference between the measured saturation temperature ( $T_s$ ) and the boiling surface temperature ( $T_w$ ) extrapolated from the derived heat flux:

$$\Delta T_s = T_w - T_s \tag{3.3}$$

$$T_{w} = T_{1} - qL_{1}/k \tag{3.4}$$

The uncertainty of the super heat,  $u_{Ts}$ , is:

$$u_{T_s} = \sqrt{\left(\frac{\delta \Delta T_s}{\delta T_1} u_{T_1}\right)^2 + \left(\frac{\delta \Delta T_s}{\delta k} u_k\right)^2 + \left(\frac{\delta \Delta T_s}{\delta L_1} u_{L_1}\right)^2 + \left(\frac{\delta \Delta T_s}{\delta q} u_q\right)^2}$$
(3.5)

where  $u_{Li}$  is the uncertainty of  $L_i$ .

The error for each thermocouple measurement was calculated as twice the standard deviation of the measurement. Using Eq. (3.2), the maximum heat flux

uncertainty proportion  $u_q/q$  was calculated to be less than 3%, and using Eq. (3.5), the maximum superheat uncertainty proportion  $u_s/\Delta T_s$  was calculated to be less than 3%.

The boiling heat transfer coefficient (HTC) was used to characterize the boiling performance:

$$HTC = {}^{q}/_{\Delta T_{c}} \tag{3.6}$$

The uncertainty of HTC  $(u_h)$  can be written as:

$$u_h = \sqrt{\left(\frac{\delta \text{HTC}}{\delta \Delta T_s} u_{T_s}\right)^2 + \left(\frac{\delta \text{HTC}}{\delta q} u_q\right)^2} \tag{3.7}$$

The maximum HTC uncertainty proportion  $u_h$ /HTC was calculated to be less than 10%.

#### 3.2.5 Surface characterization

Boiling surfaces were characterized by fluid contact angle measurements, surface roughness measurements, scanning electron microscopy (SEM), and elemental analysis. The static contact angle between the liquid and solid phases was measured using the sessile droplet method with deionized water and a Ramé-Hart goniometer with Drop Image Standard software (Netcong, NJ). The contact angle was measured at five different randomly selected locations on the surface. The surface roughness of the boiling surface was measured using a Form TalySurf 50 Profilometer 112/2564-1862 (Taylor Hobson, Leicester, UK). The arithmetic average of absolute roughness values,  $R_a$ , was measured three times at three randomly selected locations for each surface. SEM images and energy dispersive x-ray spectroscopy (EDX) elemental analysis of the surfaces were captured using a Philips XL30 SEM (FEI Company, Hillsboro, OR).

#### 3.3 BOILING MODELS

Some of the most widely used boiling models are the Rohsenow [18] and Mikic-Rohsenow [19] correlations, which depend strongly on fluid properties and the fluid-surface interaction.

#### 3.3.1 Rohsenow correlation

For pure water undergoing nucleate pool boiling, Rohsenow [18] developed the relationship:

$$\frac{c_{p(T_W-T_S)}}{h_{fg}\operatorname{Pr}^s} = C_{sf} \left[ \frac{q}{\mu h_{fg}} \sqrt{\frac{\sigma}{g(\rho-\rho_v)}} \right]^{1/3} \tag{3.8}$$

where  $c_p$  is the fluid specific heat,  $h_{fg}$  is the fluid latent heat of vaporization,  $\mu$  is the fluid viscosity,  $\sigma$  is the fluid surface tension,  $\rho$  is the fluid density,  $\rho_v$  is the vapor density, g is the acceleration due to gravity, Pr is the Prandtl number of the fluid (defined as Pr =  $\mu c_p/k_l$ ), the exponent s is 1 for water [18], and  $k_l$  is the fluid thermal conductivity. The coefficient  $C_{sf}$  takes into account the fluid-surface interaction. For example, it has determined empirically to fall in the range of 0.0065 to 0.018 for nanofluids [12].

#### 3.3.2 Nucleation site density

The Mikic-Rohsenow model [19] assumes that the main mechanism of heat transfer is transient heat conduction, such that:

$$q = \frac{1}{2} (\pi k_l \rho c_p)^{1/2} f^{1/2} D_d^2 n_a (T_w - T_s)$$
(3.9)

where  $n_a$  is the active nucleation site density, and the bubble departure diameter  $(D_d)$  and frequency (f) are defined by:

$$D_d = A_1 \left[ \frac{\sigma}{(\rho - \rho_v)} \right]^{1/2} \left( \frac{\rho C_p T_s}{\rho_v h_{fg}} \right)^{5/4} \tag{3.10}$$

$$fD_d = A_2 [\sigma(\rho - \rho_v)/\rho_v^2]^{1/4}$$
(3.11)

where A1=1.5 x 10-4 s/m $^{1/2}$  and A2=0.6 m $^{1/4}$ s $^{3/2}$  for water [20].

Since the bubble departure diameter and frequency are modeled as functions fluid properties and there are no models to incoprate the surface effects [21], we assume that they are constant for the different surfaces and ascribe any changes in boiling performance to a change in active nucleation site density. To estimate the active nucleation site density for the surfaces in this study, we rearrange the Mikic-Rohsenow correlation:

$$n_a = \frac{q}{\frac{1}{2} (\pi k_l \rho c_p)^{1/2} f^{1/2} D_d^2 (T_W - T_s)}$$
(3.12)

In this study, the following properties of water at 373.15 K were used:  $k_l = 0.67$  W/m·K,  $h_{fg} = 2257$  kJ/kg,  $c_p = 4219$  J/kg·K,  $\rho = 958$  kg/m³,  $\rho_v = 0.597$  kg/m³,  $\mu = 0.281$  cP, and  $\sigma = 58.9$  mN/m.

## 3.4 RESULTS AND DISCUSSION

### 3.4.1 Impact of coating time on boiling performance

Fig. 3.3 shows the boiling performance of the baseline surface and four coated surfaces. Surface #1 (which had a coating time of 1 minute) and Surface #2 (which had a

coating time of 5 minutes) exhibited decreases in HTC of 26% and 24%, respectively, when compared to the baseline surface. However, Surfaces #3 and #4 (which each had a coating time of 10 minutes) showed an average increase in HTC of approximately 200% compared to the baseline surface. Such an improvement in boiling HTC is large, and demonstrates that adequate coating time is crucial for optimized surface modification. To compare the EPD coated surface to results from nanofluid boiling which has particle deposition,

### 3.4.2 Residual ZnO layer

Fig. 3.4 shows Surface #4 in three stages: before coating, after coating, and after boiling. Before coating (Fig. 3.4(a)), the surface is clean, with sanding marks visible on the surface. After coating (Fig. 3.4(b)), the surface is white due to complete coverage by a thick layer of ZnO nanoparticles. After boiling (Fig. 3.4(c)), the bulk nanoparticle layer has flaked off and left behind residual surface coating that is evident from the white tint. We observed that for each surface, the coating flaked off. An EDX analysis of the surfaces performed after the bulk layer flaked off indicated the presence of both Zn and O as well as the elements of SS (Fe and Cr) as seen in Fig. 3.4(d). These observations of the EPD process for ZnO nanoparticle coating are similar to those of prior nanofluid boiling studies [22] in which a sorption layer that formed during boiling could be cleaned off with a water jet, leaving behind a thin layer of nanoparticles trapped in the surface cavities.

The HTC of the  $1^{st}$  run (after coating) and  $2^{nd}$  run (after  $1^{st}$  run, using fresh DI water) at around  $181 \text{ kW/m}^2$  for the four coated surfaces are reported in Table 1. The change in HTC from the  $1^{st}$  to the  $2^{nd}$  run is small for all four coated surfaces and

approximately within the error of the measurement (10%). The flaking off of the bulk layer occurred early on in the boiling experiment, before the system reached steady state. We therefore conclude that boiling performance is primarily governed by the residual ZnO layer and independent of the bulk layer, which flakes off.

## 3.4.3 Surface roughness

The average and standard deviation of the surface roughness ( $R_a$ ) for the baseline surface and four ZnO coated surfaces after boiling are shown in Table 1. For all coating cases, the residual nanoparticle coating slightly lowered the surface roughness after boiling, compared to the baseline surface. Because the HTC is expected to be proportional to  $R_a^{1.33}$  [23], we conclude that changes in roughness should lead to an HTC decrease of less than 7% and thus do not significantly contribute to the measured change in HTC.

## 3.4.4 Scanning electron microscopy

Surfaces were investigated using SEM at 50X, 200X, and 800X magnifications, as shown in Fig. 3.5. For the baseline surface, the lines created by sanding are visible. Surfaces #1 and #2, which did not exhibit improved boiling performance, have a similar quantity of attached ZnO particles which appear as white specks at all magnifications. Surface #4 has a coating before boiling that appears smooth with some cracking. However, the coating is observed to flake off during the boiling process. After the bulk layer flakes off, Surface #4, which had a large increase in boiling performance, still shows the presence of ZnO clusters (circled in red) are spread across the surface.

# 3.4.5 Estimation of active nucleation site density based on the Mikic-Rohsenow correlation

Figure 3.6 shows the estimated nucleation site density based on experimental data using Eq. (3.12). For the baseline surface, the nucleation site density was found to be between  $1.26 \times 10^4$  and  $1.57 \times 10^4$  sites/m<sup>2</sup> for heat fluxes between 160 and 240 kW/m<sup>2</sup>. For this same range of heat fluxes, Eq. (3.12) predicts between  $3.82 \times 10^4$  and  $4.18 \times 10^4$  sites/m<sup>2</sup> for Surface #4, an increase of 2.9 times.

### 3.4.6 Contact angle

Changes in the contact angle between a fluid and boiling surface can impact the number of active nucleation sites. The static contact angles of the different surfaces measured after boiling are shown in Table 1. Surfaces #1 and #2, which did not have an increase in boiling performance, had a small change in contact angle compared to the baseline surface. Surfaces #3 and #4, which had an increase in boiling performance, had increases of 24.4° and 11.5°, respectively. An increase in contact angle has been shown to increase boiling heat transfer [24]. When a surface has an increase in contact angle (decreased wettability), smaller cavities can become active nucleation sites, such that:

$$n_a \propto (1 - \cos(\theta)) \tag{15}$$

Therefore, the changes in active nucleation site density compared to the baseline surface predicted for Surfaces #3 and #4 are +71% and +33% respectively, and for Surfaces #1 and #2 the changes are -7% and +6%, respectively. Part of the increase in active nucleation site density for Surfaces #3 and #4 can therefore be attributed to the increase in contact angle for the coated surfaces.

## 3.4.7 Cavity analysis from SEM image processing

Analysis of dark spots in SEM images of surfaces has been shown to provide a convenient means to estimate the number of cavities available for nucleation [25]. To corroborate the nucleation site calculations made above, we processed SEM images of several surfaces in our study, utilizing an algorithm that reassigned pixels above or below a threshold grayscale intensity to have a white or black value respectively. The algorithm then outlined black spots (cavities) to yield images such as shown in Fig. 3.7 and counted the number of cavities with size greater than  $30 \ \mu m^2$ .

This method was applied to three different areas of the baseline surface, Surface #1, and Surface #4, yielding an average number of cavities of 56, 79, and 137, respectively. Surface #4 had a 140% increase in the number of cavities compared to the baseline surface, while Surface #1 had only a slight increase in the number of cavities. These percent increases did not vary significantly when choosing different grayscale threshold intensities in the cavity counting algorithm. The increase for Surface #4 is similar to the 190% increase in active nucleation site density calculated above using the Mikic-Rohsenow correlation, further suggesting that the increase in HTC due to the EPD coating is due to an increase in nucleation site density.

#### 3.5 CONCLUSIONS

This study demonstrated that the boiling performance of a surface can be significantly enhanced by electrophoretically depositing a layer of nanoparticles from a nanofluid. When coating times were sufficient, a nearly 200% increase in HTC was measured. The coating slightly lowered the measured surface roughness, but it was not substantial enough to impact the HTC. The coating increased the contact angle and

cavities available for nucleation, resulting in an increase in boiling HTC. The change in active nucleation site density calculated from the experimental data (2.9 X) matched well with the image analysis to determine the change in the number of cavities (2.4 X). This demonstrates that the increase in nucleation site density is likely to be the cause of the improved boiling performance. For future work, an investigation of the bubble departure diameter and frequency would help to broaden the understanding of boiling on an EPD surface, since it was assumed that different surfaces did not impact these values.

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Table 3.1. Coating time, HTC, surface roughness and contact angle results

coating time, I	Coating	HTC (kW/m <sup>2</sup> K) *		$R_a (\mu m)$		Contact Angle		
Surface	Time (min)	1 <sup>st</sup> Run	2 <sup>nd</sup> Run	change from 1 <sup>st</sup> to 2 <sup>nd</sup> Run	Average	Standard deviation	Average	Standard deviation
Baseline	N/A	18.2	-	-	0.19	0.01	65.3°	2.2°
#1	1	12.8	13.5	+5%	0.15	0.01	62.7°	4.7°
#2	5	13.9	16.0	+15%	0.12	0.01	67.4°	1.5°
#3	10	53.2	50.4	-5%	0.18	0.05	89.7°	6.4°
#4	10	53.3	52.3	-2%	0.16	0.01	76.8°	2.7°

<sup>\*</sup> for  $q = 181 \text{ kW/m}^2$ 

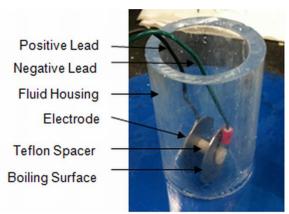


Figure 3.1. Electrophoretic configuration used to coat ZnO nanoparticles on a SS surface.

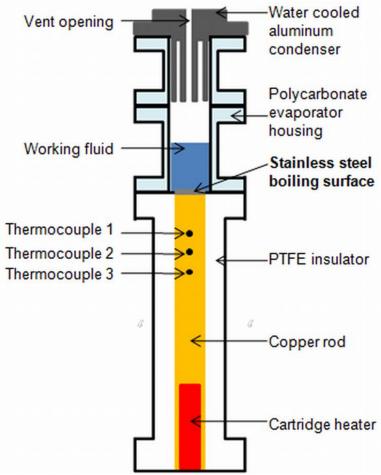


Figure 3.2. Schematic of boiling test system.

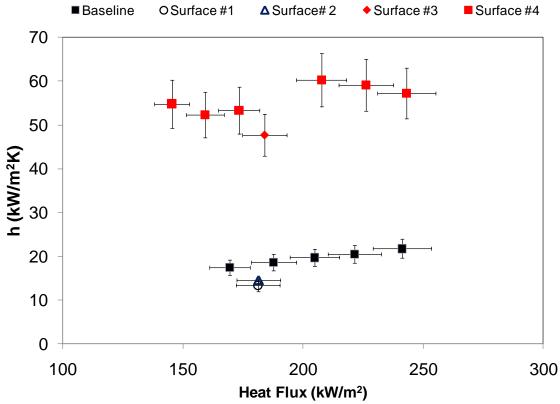


Figure 3.3. Experimental boiling results comparing the baseline surface to the electrophoretically coated Surfaces #1-#4.

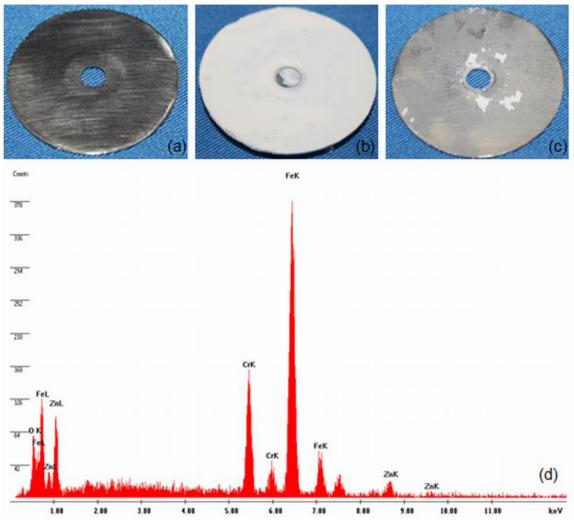


Figure 3.4. Surface #4 shown (a) before coating, (b) after coating/before boiling, and (c) after boiling (when the bulk layer has flaked off). (d) EDX analysis of the surface after boiling, showing that ZnO is still present.

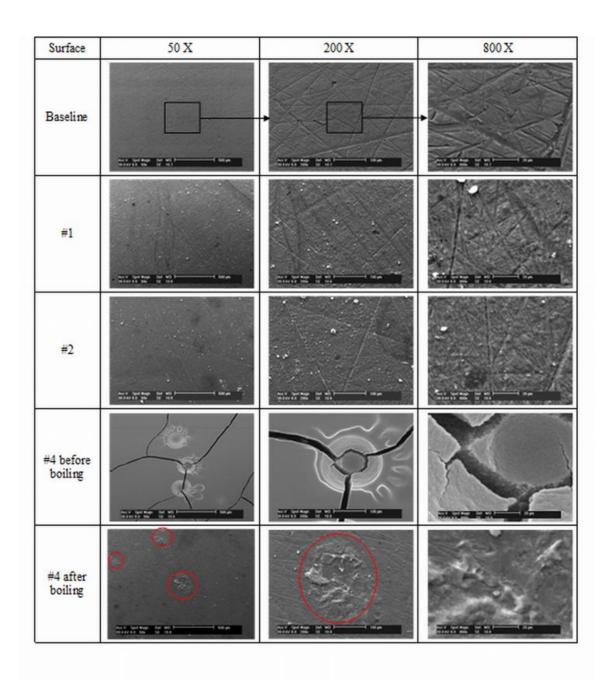


Figure 3.5. SEM images of different surfaces investigated at three magnifications.

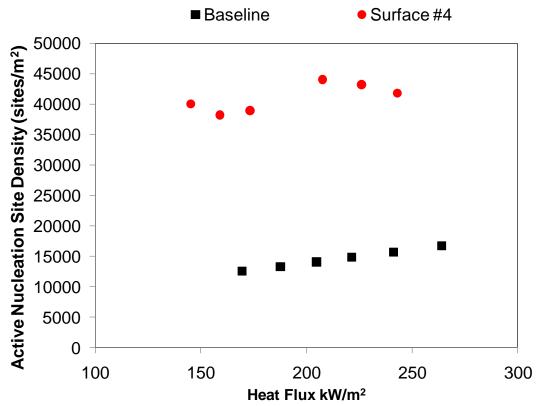


Figure 3.6. Comparison of active nucleation site density calculated from the Mikic Rohsenow correlation for baseline (uncoated) surface and (coated) Surface #4.



Figure 3.7. Representative images produced by thresholding algorithm that are used to analyze the number of cavities available for nucleation in (a) baseline, (b) Surface #1, and (c) Surface #4.

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## **CHAPTER 4**

# INVESTIGATION OF THE ELECTRICAL CONDUCTIVITY OF PROPYLENE GLYCOL BASED ZNO NANOFLUIDS

## **ABSTRACT**

The electrical conductivity is an important property of nanofluids that has not been widely studied. To study the effects of particle size and concentration, the electrical conductivity of propylene glycol based ZnO nanofluids with 20, 40, and 60 nm diameter particles is measured. Since conventional models, such as the Maxwell Model, do not account for the charge formed on the particle's surface, the data was compared to a model based on the electrical conductivity of a colloidal suspension in a salt-free medium with reasonable agreement. Both the model and experimental results showed the ZnO nanofluids increased the electrical conductivity with increasing volume fraction and decreasing particle size compared to that of the base fluid. Furthermore, at higher volume fractions, the increase of electrical conductivity begins to level off, which is attributed to ion condensation effects in the high surface charge regime.

# **4.1 INTRODUCTION**

Nanofluids are created by suspending nanometer size particles into a base fluid [1] and allow for the engineering of the fluid properties by changing the type, size, and amount of particle. Nanofluids have been proposed for advanced heat transfer applications, such as fuel cell thermal management and power electronics cooling. For

these cooling applications, a low electrical conductivity fluid is required. While the fluid thermal properties have been thoroughly investigated [2], the interrelated electrical properties have not.

Nanofluids typically have metal oxide nanoparticles, such as ZnO, TiO<sub>2</sub>, and Al<sub>2</sub>O<sub>3</sub>. When dispersed into a fluid, a surface charge is formed due to the protonation and deprotonation of a surface group such as a hydroxyl ligand (-OH) [3]. This causes a layer of ions called the electric double layer (EDL) to form. This layer may cause polar liquids to form columns along the surface of the particle, and more site densities and larger surface charges may lead to stronger structures [3]. The surface charge can be adjusted by altering the pH of the suspension [4,5], or chemical treating the surface of the particles [5]. The charge on the surface allows for electrophoresis of the fluids. Chang et al. [6] applied an electric field to the SiC nanofluids and measured increases in the hydrodynamic radius of the particles, aggregation, and precipitation of the aggregates. They also noted that the nanoparticles coated the electrode oppositely charged of the nanoparticles.

The electrical conductivity of a suspension can either increase or decrease depending on the background electrolyte, the particle size, and the charge of the particle [7]. Fang and Zhang [8] studied the concentration effects (<0.2%) on electrical conductivity of gold nanoparticles in chloroform and toluene suspensions. Cruz et al. [9] studied high volume fraction (1-35%), 520 nm Al<sub>2</sub>O<sub>3</sub> suspensions in different water-salt solutions. For pure water the conductivity steadily increases with the addition of nanoparticles. Wong et al. [10] also measured a linear increase in electrical conductivity with volume fraction (up to 8.5%) for water-based 36 nm Al<sub>2</sub>O<sub>3</sub> nanofluids. Ganguly et

al. [11] reported increases in electrical conductivity of 13 nm Al<sub>2</sub>O<sub>3</sub> nanofluids with both increasing temperature and volume fraction (up to 3%). Lisunova et al. [12] measured electrical conductivity increases for carbon nanotube (CNT) based nanofluids and measured large increases in electrical conductivity with increasing volume fraction (up to 2%). The previous nanofluid electrical conductivity studies showed increases of up to 150X [11], while conventional models, such as the Maxwell Model [13], only predict an increase of less than 1.5X. The review shows that there is a lack of understanding of the particle size effect and mathematical modeling on nanofluid electrical conductivity.

Several other models have been proposed for the electrical conductivity colloidal suspensions. Carrique et al. [14] numerically studied the effects of overlapping EDLs on the electrical conductivity of concentrated spherical colloidal suspensions. Carrique et al. [15] investigated the role of boundary conditions on the electrical conductivity of salt free suspensions using numerical methods. Carrique and Ruiz-Renna [16] explored the effects of water dissociation and CO<sub>2</sub> contamination on the electrical conductivity of salt-free concentrated suspensions. While these models are valid for a variety of cases, Ohshima [17] developed a simplified analytical model of the electrical conductivity of spherical colloidal particles in a salt-free medium. However, these models have not been used to describe the electrical conductivity metal oxide nanofluids. The analytical model proposed by Ohsima [17] is used in this study to understand the electrical conductivity of non-aqueous nanofluids.

The goal of this research is to experimentally measure and analytically model the particle size and concentration effects on nanofluid electrical conductivity. The electrical conductivity of ZnO nanofluids are investigated for 20, 40, and 60 nm diameter

nanoparticles dispersed in propylene glycol, which allowed for up to a 7% volume fraction without the addition of dispersants. An analytic model for colloidal particles in a salt-free medium is used with the properties of the base fluid and nanoparticle to explain the changes in electrical conductivity.

## 4.2 NANOFLUID PREPARATION AND CHARACTERIZATION

# 4.2.1 Nanofluid Preparation

Propylene glycol (PG) based nanofluids with 20, 40, and 60 nm diameter ZnO nanoparticles were used. The fluids were prepared by Nanophase Technologies (Romeoville, IL) at 7% volume fraction without any dispersants. The PG-based suspensions were diluted to 1, 3, and 5% for all particle sizes. The particles have a crystal phase of zincite (hexagonal) with an elongated morphology. Using PG as a base fluid allowed for higher volume fractions of nanoparticles to be used without the use of dispersants compared to water-based nanofluids.

## 4.2.2 Zeta Potential Measurement

The zeta potential of the nanofluids was measured with a Brookhaven Instruments Zeta Plus (Holtsville, NY) which uses electrophoretic light scattering (ELS). An electric field is applied to the fluid that moves the charged particles towards the positive or negative pole. The direction and speed of the particles are measured to determine the zeta potential by a laser beam that is scattered proportionally to the velocity of the particles. The samples were temperature controlled to 25°C and the zeta potential was measured in 5 separate runs.

# **4.2.3** Electrical Conductivity Measurement

The electrical conductivity of the nanofluids was measured using a Model 72 hand held conductivity meter from Engineered Systems and Designs, Inc (Newark, DE). The meter operated using two electrodes insulated from each other. With a specific size and spacing between them, the meter can measure the conductance across a known volume. The meter has a range of 0.2 to 20,000  $\mu$ S/cm, a resolution of 0.1%, and an accuracy of  $\pm$  2.5%. The meter was calibrated using salt solutions with known electrical conductivities of 10, 74, 714, 2,000, 6,668, and 58,640  $\mu$ S/cm. For the measurements, approximately 200 ml of the fluid was placed in a beaker which was then submerged in a temperature bath of 25 °C. The probe was rinsed in tap water and then dipped into three different beakers of distilled water to prevent any contamination between samples. The probe was then dipped in the sample beaker and stirred until there was a stable reading. The rinsing process was repeated before moving to the next sample. The tests were repeated three times for each sample.

## 4.3 NANOFLUID ELECTRICAL CONDUCTIVITY MODELING

In a salt-free medium, ions are added from the charging of the particle suspended in the base fluid, such as a non-aqueous medium [18]. These ions are referred to as counter-ions. PG, or 1, 2 propanediol, is a polar medium with the formula HO-CH<sub>2</sub>-CHOH-CH<sub>3</sub>, thus allowing the available hydroxyl group to form on the surface when the ZnO is dispersed in the fluid.

Kuwabara's cell model [19], as shown in Fig. 4.1, represents the interaction of the particle in the medium and is used to analyze the electrokinetic properties of colloidal

suspensions. Each particle of radius a is surrounded by a virtual shell of radius b such that the volume fraction,  $\phi = (a/b)^3$ . The EDL thickness is represented by  $\kappa^{-1}$ .

Ohshima [17] derived an analytical expression for the electrical conductivity in a salt-free medium for two cases, one with low surface charge and the other with high surface charge. The critical value of the surface charge of the nanoparticle, Q, was related to the volume fraction by:

$$\frac{Q}{4\pi\varepsilon_{\nu}\varepsilon_{0}a}\frac{ze}{kT} \le \ln(1/\phi) \tag{4.1}$$

where  $\varepsilon_r$  is the relative permittivity of the medium,  $\varepsilon_0$  is the permittivity of a vacuum, z is the valence of the counter-ion, e is the elementary electric charge, k is the Boltzmann constant, and T is the temperature in K. The universal constants used were:  $\varepsilon_0 = 8.85 \text{ x}$   $10^{-12} \text{ F/m}$ , e = 1.60 x  $10^{-19} \text{ C}$ , and k = 1.38 x  $10^{-23} \text{ J/K}$ . In this study, the relative permittivity of PG was 32 and the temperature was held at 298.15 K. The valence of the counter-ion was 2 [20]. The zeta potential,  $\zeta$ , is used to determine Q [17]:

$$Q = 4\pi\varepsilon_r\varepsilon_0\alpha\zeta\tag{4.2}$$

From this Ohshima [17] approximated the electrical conductivity, K, for the two cases as follows. For Case 1, Eq. (4.1) is true, the particle has low surface charge and

$$K = \frac{3zeQ}{4\pi b^3 \lambda} = \frac{3ze\,\varepsilon_r\,\varepsilon_0\phi\zeta}{4\pi\,a^2\lambda} \tag{4.3}$$

For Case 2, Eq. (4.1) is false, the particle has high surface charge and

$$K = \frac{3\varepsilon_r \varepsilon_0 akT}{b^3 \lambda} \ln(1/\phi) = \frac{3\varepsilon_r \varepsilon_0 \phi kT}{a^2 \lambda} \ln(1/\phi)$$
(4.4)

where  $\lambda$  is the drag coefficient of the counter-ion is determined from [21]:

$$\lambda = \frac{N_A e^2 |z|}{\Lambda_c^0} \tag{4.5}$$

where  $N_A$  is Avogadro's number  $(6.022 \times 10^{23} \text{ mol}^{-1})$  and  $\Lambda_c^0$  is the limiting ionic conductance. For Case 1, K increases linearly with Q which is impacted by the zeta potential in Eq. (4.2). For Case 2, K is independent of Q due to counter-ion condensation effects. Increasing the amount of counter-ions adds to the condensation region, leaving the charge and potential outside that region unchanged. In this case, the electrical conductivity no longer increases linearly with volume fraction and it begins to plateau at high volume fraction. These relationships were used to explain the experimental data of the PG-based nanofluids.

Since the value of  $\Lambda_c^0$  was an unknown for the ZnO/PG system, the value of  $\lambda$  is also an unknown. The value of  $\Lambda_c^0$  is independent of the particle size and concentration since it depends on the counter-ion at the particle surface. The root-mean-square error between the experimentally measured and model predicted K, for all particle sizes and volume fractions was minimized to determine  $\Lambda_c^0$ . The value of  $\Lambda_c^0$  is substituted into Eq. (4.5) to find  $\lambda$ , which is the input to Eq. (4.3) or (4.4) to solve K.

# 4.4 RESULTS AND DISCUSSION

# 4.4.1 Zeta Potential

The results of the zeta potential measurements are shown for the PG-based ZnO nanofluids in Table 4.1. There is little difference between the zeta potential magnitudes based on particle size for each base fluid. Since the fluids were prepared using the same

process, the zeta potential was controlled and should be constant. The average zeta potential of 48.7 mV was used for modeling.

# **4.4.2 Modeling Parameters**

To determine which case was valid for the particle size and volume fraction, the values of the inequality in Eq. (4.1) were calculated and compared. Using Eq. (4.2) in Eq. (4.1), the value of the left hand side of the inequality of Eq. (4.1) was determined to be 3.78 and did not change for different particle sizes or volume fractions. The values of the right hand side of Eq. (4.1) are dependent on the volume fraction and are shown in Table 4.2. For the PG/ZnO based nanofluids, Case 1 was valid only for the 1% volume fraction suspensions of all three particle sizes, and Case 2 was valid for the rest of the suspensions (3, 5, and 7%). From the analysis, the value of  $\Lambda_c^0$  was determined to be 49.4 S cm<sup>2</sup>/mol. This is in line with common values for  $\Lambda_c^0$  of ions arising from the colloidal charge process in aqueous suspensions, which are between about 40-350 S cm<sup>2</sup>/mol [15]. Based on the  $\Lambda_c^0$  value obtained,  $\lambda = 6.26 \times 10^{-12} \text{ C}^2/(\text{S m}^2)$ .

# 4.4.3 Comparison of Experimental and Modeling Results

The experimental data is shown in Fig. 4.2 with solid symbols. The addition of nanoparticles to the suspension increased the electrical conductivity when compared to the PG base fluid ( $K = 0.1 \,\mu\text{S/cm}$ ). The electrical conductivity increased with the volume fraction for each suspension. At the same volume fraction, smaller particles had a higher electrical conductivity. The electrical conductivity of the 20 nm suspension increased from 2.72  $\mu$ S/cm at 1% volume fraction to 9.60  $\mu$ S/cm at 7%, which is an increase of nearly 100 times compared to the base fluid. The electrical conductivity of the 40 nm suspension increased from 1.35  $\mu$ S/cm to 2.35  $\mu$ S/cm over the same range of volume

fractions. While the 60 nm suspension increased from 0.76  $\mu$ S/cm to 1.99  $\mu$ S/cm. Compared to PG, the 40 nm and 60 nm particle sized suspensions increased by 23.5 and 19.9 times, respectively.

The model data for the PG-based nanofluids is also shown in Fig. 4.2 with solid lines. The model is similar to the experimental data by showing the electrical conductivity increases with increasing volume fraction and increases with decreasing particle size at the same volume fraction. The model for the 20 nm and 40 nm suspension fit quite well; however, the 60 nm under predicts the experimental data. This may be due to the measurement error associated with low values of electrical conductivity associated with the 60 nm nanofluid.

## **4.4.4 Particle Size and Concentration Effects**

The size of the ZnO particle affected the nanofluid electrical conductivity significantly. The electrical conductivity decreased with increasing particle size at the same volume fraction. However, the change in electrical conductivity was more significant from 20 nm to 40 nm than from 40 nm to 60 nm, as shown experimentally and as modeled in Fig. 4.3. In the model, the electrical conductivity is proportional to  $a^{-2}$ ; therefore, it is expected that the electrical conductivity will increase with decreasing particle size. It is noted that this difference increases with volume fraction from 20 to 40 nm, and the opposite trend is observed volume fraction increase from 40 to 60 nm for the experimental data, while the model predicts an increase with volume fraction for both.

The concentration affects the electrical conductivity as well. While the electrical conductivity increases with volume fraction, the rate of change of the electrical conductivity per volume fraction change was further investigated. The rate of change

was defined as the change in electrical conductivity divided by the change in volume fraction. There is a large increase in thermal conductivity at low volume fractions; however, at higher volume fractions, the increase in electrical conductivity with volume fraction slows, as shown in Fig. 4.4. For the 20 nm suspension, there is a rapid increase in the electrical conductivity from the base fluid to the 1% volume fraction suspension and for 1% to 3%; however, the increase from 3% to 5% and 5% to 7% is lower than the initial increase. There is also similar trend for the 40 nm and 60 nm. For volume fractions greater than 3%, the model is uses Case 2 and the electrical conductivity is proportional to  $\phi \ln(1/\phi)$ , which predicts a non-linear response to the volume fraction with a decrease in slope at higher volume fractions.

# **4.4.5** Effects of Surface Charge in Modeling

The surface charge is important for the accuracy of the modeling since it determines which case is used. For low surface charge Case 1 is relevant and for high surface charge Case 2 is relevant. For the PG/ZnO nanofluids, the model switched from Case 1 to Case 2 between 1% and 3% volume fraction. To see the difference between these two cases, Fig. 4.5 shows the K value calculated based on Case 1 (Eq. (4.3)) and Case 2 (Eq. (4.4)) for the 20 nm ZnO suspension. Case 1 over predicts the electrical conductivity data at higher volume fractions with a linear increase with volume fraction. Case 2 increases with the volume fraction at a rate of  $\phi \ln(1/\phi)$ , which better fits the data. This demonstrates the effect of surface charge on the accuracy in modeling of K. Similar trends were seen for 20 and 40 nm particle sizes. At volume fractions less than 1%, there is not a significant difference in the electrical conductivity predicted by the two cases.

## 4.4.6 Electrokinetic Radius

For suspensions with a thick EDL, the electrical conductivity is larger than that of the bulk fluid and the surface conductance increases the effective electrical conductivity [22]. The electrokinetic radius,  $\kappa a$ , is the ratio of the particle diameter to the thickness of the EDL,  $\kappa^{-1}$ , and is used to describe the coupled effects of the particle size and ion concentration. The Debye-Hückle parameter,  $\kappa$ , is defined as:

$$\kappa = \frac{1}{a} \left[ \frac{3ze\zeta}{kT(1/\phi - 1)} \right]^{1/2} \tag{4.6}$$

Therefore,

$$\kappa a = \left[\frac{3ze\zeta}{kT(1/\phi - 1)}\right]^{1/2} \tag{4.7}$$

Eq. (4.7) shows that  $\kappa a$  is independent of particle size and increases with volume fraction. For values  $\kappa a$  <1, imply a relatively thick double layer. Posner et al. [7] stated that more work is needed in predicting the conductivity for  $\kappa a \sim 1$ .

Fig. 4.6 shows the value of  $\kappa a$  is less than 1 for the suspensions studied and it increases with volume fraction according to Eq. (4.7) for the ZnO/PG-based nanofluids used in this study. The value of  $\kappa a$  ranges from 0.34 at 1% volume fraction to 0.92 for 7% volume fraction. Therefore, the ZnO/PG nanofluids have a thick EDL which affects the electrical conductivity of the suspensions.

# 4.5 CONCLUSIONS

A model developed for colloidal suspensions in a salt-free medium was used to model the large increases in the electrical conductivity of ZnO/PG nanofluids. The data for the 20 nm and 40 nm suspensions fit well, while the model under predicted the 60 nm. Both the experimental data and the model showed the electrical conductivity increases with increasing volume fraction and with decreasing particle size at the same volume fraction. Also, at higher volume fractions, the increase of electrical conductivity begins to level off, which may be attributed to ion condensation effects in the high surface charge regime. The electrokinetic radius was found to be less than 1, which implies a large double layer thickness.

Table 4.1. Zeta potential measurements

2a (nm)	$\zeta$ (mV)
20	$49.3 \pm 0.1$
40	$48.6 \pm 0.1$
60	$48.3 \pm 0.1$

Table 4.2. Values of the right hand side of the inequality in Eq. (4.1)

φ	$ln(1/\phi)$
1%	4.61
3%	3.51
5%	3.00
7%	2.66

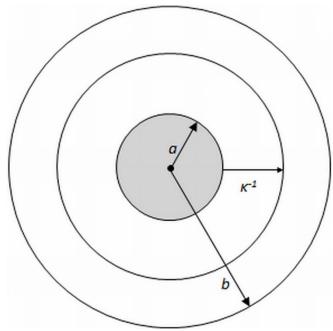


Figure 4.1. Schematic of Kuwabara's cell model with particle radius, a, layer thickness,  $\kappa$ -1, and particle space, b.

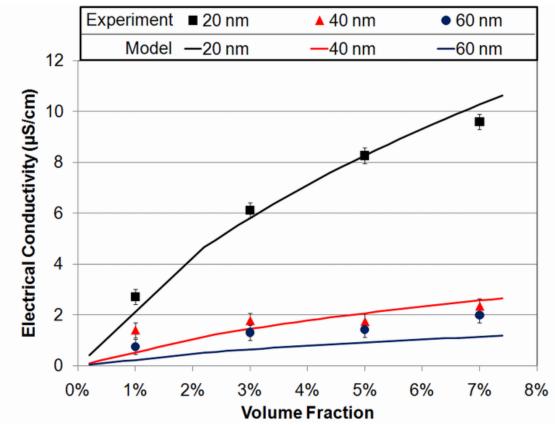


Figure 4.2. Experimental electrical conductivity of propylene glycol-based ZnO nanofluids with 20, 40, and 60 nm particles in comparison to a model for the electrical conductivity of a colloidal suspension in a salt-free medium.

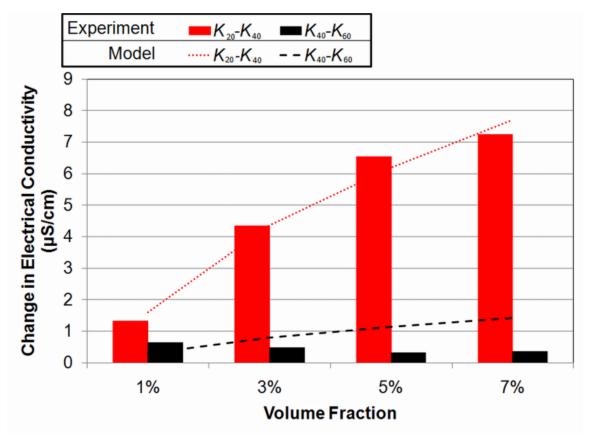


Figure 4.3. The difference in electrical conductivity between particle sizes of  $20~(K_{20})$  and  $40~\text{nm}~(K_{40})$  and  $40~\text{nm}~(K_{60})$  at the same volume fraction. The difference in electrical conductivity between the 20~nm and 40~nm is larger than the difference between 40~nm and 60~nm particles.

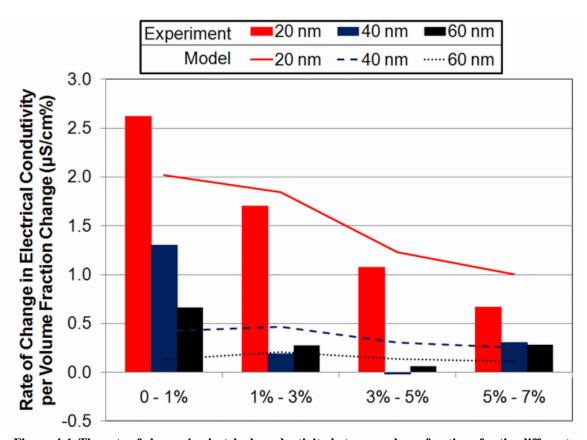


Figure 4.4. The rate of change in electrical conductivity between volume fractions for the different particle sizes. The rate of change decreases at higher volume fractions.

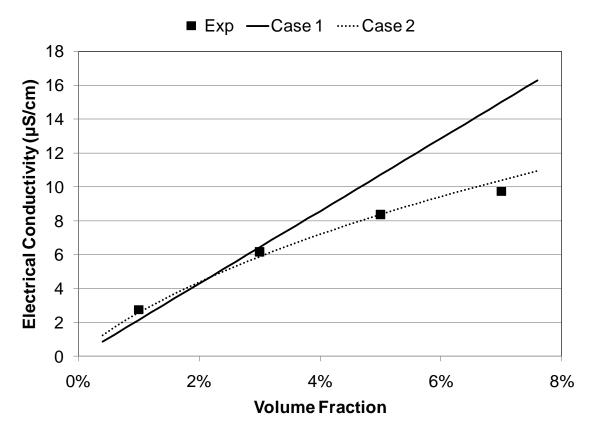


Figure 4.5. Two potential cases for the model using Eqs. (4.3) and (4.4) for the low surface charge (Case 1) and high surface charge (Case 2), respectively, for the 20 nm ZnO/PG suspension.

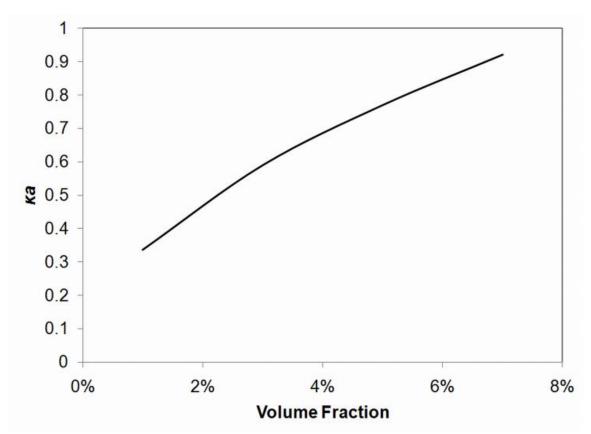


Figure 4.6. Variation of the electrokinetic radius,  $\kappa a$ , with volume fraction.

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## **CHAPTER 5**

## CONCLUSIONS AND FUTURE WORK

## **5.1 CONCLUSIONS**

This dissertation investigated nanofluids for applications involving boiling heat transfer. The impact of nanoparticle deposition on the boiling performance of a surface was studied using two different methods: particle deposition during boiling of nanofluids and electrophoretic deposition of particles from a nanofluid. Also, the particle size and concentration effects on the electrical conductivity of nanofluids were investigated.

The creation of a nanoparticle layer on the heated surface was shown to increase the boiling performance of water up to 62% and correlate well with a model that uses AFM-measured surface roughness data. Nanofluid boiling performance was shown to increase by about 25% initially (compared to water) before the deposited nanoparticle layer is completely formed. After the nanoparticle layer fully covers the heater surface, the nanofluid boiling coefficient declines, which is to a suppression of bubble nucleation and transport by the suspended nanoparticles.

High particle volume fraction nanofluids were used to create EPD coating and demonstrated a significant improvement in boiling performance. When coating times were sufficient, a nearly 200% increase in HTC was measured. The coating increased the cavities available for nucleation, resulting in an increase in boiling HTC. The 2.9 times increase in active nucleation site density calculated from the

experimental data matched well with the image analysis to determine the change in the number of cavities (2.4 time increase). This demonstrated that the increase in nucleation site density is potentially the cause of the improved boiling performance.

The experimentally measured electrical conductivity of nanofluids increased with increasing volume fraction and decreasing particle size. For a 20 nm ZnO particle in propylene glycol, the electrical conductivity increased by over 100 times, while the 40 nm and 60 nm particle sized suspensions increased nearly 25 and 20 times, respectively. This large of an increase was not predicted by conventional models, which only predict a small increase. A model developed for colloidal suspensions in a salt-free medium was used to model the electrical conductivity of ZnO/PG nanofluids. The data for the 20 nm and 40 nm suspensions fit well, while the model under predicted the 60 nm. Furthermore, at higher volume fractions, the increase of electrical conductivity begins to level off, which is attributed to ion condensation effects in the high surface charge regime.

The major contributions of this research can be summarized as follows.

Nanofluid boiling was shown to be influenced by two competing processes. First, nanofluids enhanced boiling by depositing layer of nanoparticles on the boiling surface. Second, nanofluids degraded boiling performance due to the particles in suspension which suppressed bubble growth. The increased boiling performance due to the nanoparticle layer formed during boiling was attributed to a roughening of the surface. This was corroborated with a model used to correlate the roughness with the boiling heat transfer coefficient.

- A novel method to create enhanced boiling surfaces using EPD of nanoparticles was introduced. With adequate coating times, the EPD creates a layer which increases the boiling heat transfer coefficient compared to an uncoated surface. The increase boiling heat transfer due to the EPD coating was thought to be due to an increase in nucleation site density. This was verified by comparing the calculated change in nucleation site density from the experimental data to an estimation of cavities obtained from SEM image analysis.
- The electrical conductivity of nanofluids were measured and validated with a model showing that the electrical conductivity increases with increasing volume fraction and decreasing particle size. The model used was able to predict the large increase in electrical conductivity of the nanofluids that could not be predicted with conventional models. Also, the leveling off of electrical conductivity at higher volume fractions was ascribed to ion condensation effects.

## **5.2 FUTURE WORK**

This research has identified the need for future studies in the following areas.

- Further study is needed for the mechanism of nanoparticle deposition during boiling
  of nanofluids. Many researchers have observed this phenomenon and it was used in
  this study to create enhanced boiling surfaces, but the process of particle attachment
  is not well understood. Knowledge of this process could be used to create surfaces
  designed specifically for boiling nanofluids.
- 2. This research showed that coating time played an important role during the EPD process. Therefore, optimal coating parameters of the EPD process should be further

- studied. Also, treatments after coating used in other EPD processes, such as sintering, may change the quality of the layer. Once optimized, large scale manufacturing using this coating method could be developed for specific applications that require enhanced boiling surfaces.
- 3. EPD can be used to apply coatings to surfaces of different shapes and can be used to infiltrate foams. So, the EPD of nanoparticles from nanofluids coating process could be applied to boiling surfaces with enhanced features such as micro-fins and porous foams to create hybrid surfaces for further enhancement.
- 4. To create a tunable thermal management system, an electric field could be applied to the surfaces during boiling of nanofluids. In high performance situations, a field could be applied to attract particles. In low performance situations, the field could be turned off or reversed to prevent particle deposition.
- 5. To extend the model for electrical conductivity of nanofluids, a variety of particles and base fluids of the nanofluids should be investigated. Since a fluid's thermal and electrical properties are often interrelated, information on the electrical conductivity could lend new insight into the thermal conductivity enhancements of nanofluids. This information could be used to weigh the tradeoffs between enhanced thermal performance and increased electrical conductivity.