

## Grain-boundary relaxation and its effect on plasticity in nanocrystalline Fe

D. Jang

*Department of Materials Science and Engineering, University of Michigan, Ann Arbor, Michigan 48109*

M. Atzmon<sup>a)</sup>

*Department of Nuclear Engineering and Radiological Sciences, University of Michigan, Ann Arbor, Michigan 48109 and Department of Materials Science and Engineering, University of Michigan, Ann Arbor, Michigan 48109*

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Grain-boundary relaxation in nanocrystalline Fe was studied by high-resolution transmission electron microscopy (HREM), and its effect on plasticity was characterized by nanoindentation. Samples with grain size of 9.8 nm were synthesized by mechanical attrition and subsequently annealed at low temperatures (80 and 100 °C) without affecting the grain size. While the hardness is not significantly affected by annealing, the strain-rate sensitivity peaks as a function of annealing time, suggesting two competing processes. HREM images show grain-boundary relaxation during annealing. Initially disconnected lattice planes were observed to evolve into a more continuous and ordered structure with regularly spaced grain-boundary dislocations. © 2006 American Institute of Physics. [DOI: 10.1063/1.2187417]

### INTRODUCTION

Nanocrystalline materials, defined as polycrystals with grain size of less than a few tens of nanometers, were reported by Birringer *et al.*,<sup>1</sup> and have been attracting theoretical and practical interests due to their unique properties. For example, contrary to the typical strengthening of conventional coarse-grained materials with decreasing grain size,<sup>2,3</sup> when the grain size is in the range of a few tens to a few nanometers, the dislocation pileup mechanism for Hall-Petch strengthening breaks down.<sup>4,5</sup> Softening with decreasing grain size is observed, termed as inverse Hall-Petch behavior.<sup>6,7</sup> For a grain size in the range of a few nanometers, the grain boundaries can occupy as much as 30% of the volume.<sup>8</sup> Therefore, understanding the structure and behavior of grain boundaries is of essential importance to the study of nanocrystalline materials.

Many researchers have attempted to determine the structure of nanocrystalline materials. In studies of nanocrystalline Pd, synthesized by inert gas condensation, by high-resolution transmission electron microscopy (HREM), Thomas *et al.*<sup>9</sup> and Straub *et al.*<sup>10</sup> reported that the lattice fringes from adjacent grains approached each other very closely at the grain boundary (at most 0.4 nm), and that they changed direction abruptly at the boundary. This observation is consistent with molecular dynamics (MD) simulation result by Van Swygenhoven *et al.*<sup>11</sup> who reported structural coherency at the grain boundary of nanocrystalline Ni and Cu, created by Voronoi construction.<sup>12</sup> However, different grain-boundary structure has also been reported. Interpreting x-ray diffraction data, Zhu *et al.*<sup>13</sup> and Haubold *et al.*<sup>14</sup> suggested an amorphous-layer-like grain-boundary structure for nanocrystalline Fe,<sup>13</sup> Cu, and Pd,<sup>14</sup> formed by inert gas condensation. MD simulation by Koblinski *et al.*<sup>15</sup> also showed

a “gluelike intergranular phase” for nanocrystalline Pd, formed by crystal growth from seeds embedded into a melt.

The fact that reports for nominally the same material are in conflict with each other suggests that the synthesis method or thermal history may affect the grain-boundary structure. Modification of the grain-boundary structure by heat treatment has been reported. Using x-ray atomic distribution function analysis, Löffler and Weissmüller<sup>16</sup> reported that in nanocrystalline Pd, formed by inert gas condensation, the disorder at or near the grain boundary decreases during annealing or aging. Using *in situ* x-ray analysis, Moelle and Fecht<sup>17</sup> observed that in ball-milled nanocrystalline Fe, the lattice strain decreased significantly by annealing below 200 °C while the grain size remained constant. They suggested that this observation was likely due to structural relaxation of the grain boundaries. Grain-boundary relaxation by annealing was also observed by Hasnaoui *et al.*<sup>18</sup> in MD simulation. They reported densification of the crystallographic planes of each grain at grain boundary by annealing at 800 K for 100 ps. Using HREM for artificially formed grain boundaries of bicrystals, Merkle<sup>19</sup> has performed extensive research on grain-boundary relaxation in metals and ceramics, classifying grain-boundary relaxation into four different modes.

Given that grain-boundary activity contributes significantly to plasticity of nanocrystalline materials,<sup>20,21</sup> their mechanical behavior is likely to be affected by grain-boundary relaxation. The strain-rate sensitivity, defined as  $m = d \ln \sigma / d \ln \dot{\epsilon}$ , where  $\sigma$  is the applied stress and  $\dot{\epsilon}$  is the corresponding strain rate, is a useful parameter for understanding the grain-boundary contribution to plastic deformation because grain-boundary activity is thermally activated, and thus strain-rate dependent.<sup>21</sup> We have previously reported<sup>7</sup> that the strain-rate sensitivity increases monotonically as the grain size decreases (Fig. 1). This observation suggests that increased grain-boundary volume fraction at smaller grain size increases the contribution of grain-

<sup>a)</sup>Electronic mail: atzmon@umich.edu

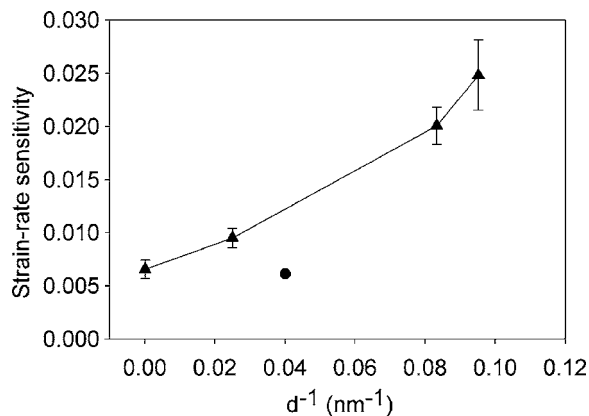


FIG. 1. Strain-rate sensitivity as a function of inverse grain size. The solid circle (●) represents the value for Fe with 24 nm grains, which had experienced heat treatment, therefore, grain growth and grain-boundary relaxation (Ref. 22).

boundary activity to deformation, making the stress more sensitive to the strain rate. Also, the strain-rate sensitivity of sintered nanocrystalline Fe,<sup>22</sup> which had undergone grain growth, was lower than that of an as-milled sample<sup>7</sup> with the same grain size (see Fig. 1), suggesting that grain-boundary relaxation, caused by heat treatment, possibly affected grain-boundary plasticity.

In the present study, relaxation of grain boundaries in nanocrystalline Fe and its effect on plasticity are investigated by differential scanning calorimetry (DSC), nanoindentation, and HREM. Nanocrystalline Fe samples with volume-averaged grain size of 9.8 nm were synthesized by ball milling and then annealed at low temperatures that did not affect the grain size. The effect of annealing on plasticity was studied by nanoindentation. The atomic structure of grain boundaries and its dependence on annealing time were observed using HREM.

## EXPERIMENTAL PROCEDURE

Fe powder of 99.9% purity was obtained from Cerac Company, with all samples originating from the same batch. Nanocrystalline Fe was prepared at room temperature by high-energy ball milling in Ar in a Spex 8000 Mixer/Mill, using a stainless steel vial and balls, for 24 h. The ball-to-sample weight ratio for high-energy ball milling was about 4:1. For comparison, one sample was prepared by low-energy ball milling in Ar using a Fritsch Pulverisette-0 vibrating-frame ball mill for 100 h. In this case, 3.2 g of Fe powder were milled with a stainless steel ball of 5 cm diameter and 500 g weight. Unless otherwise noted, results are quoted for samples formed by high-energy ball milling. We have previously ruled out a significant effect of impurities acquired during ball milling on subsequent mechanical behavior.<sup>7</sup> Samples were annealed at 80 and 100 °C from 30 min to 24 h in a Ti-gettered flowing-Ar atmosphere or in vacuum ( $<10^{-8}$  torr).

X-ray diffraction (XRD) was performed using a horizontal General Electric  $\theta$ -2 $\theta$  powder diffractometer in step-scanning mode with Mo  $K\alpha$  radiation ( $\lambda=0.07093$  nm). The peaks were fitted with a Pearson VII function, and the

volume-averaged grain size and root mean square (rms) stain were determined by the Warren-Averbach analysis.<sup>23</sup> To prepare samples for hardness and strain-rate sensitivity measurements, the powder was consolidated under 17 MPa at room temperature and then cold mounted with epoxy resin. The powder particles were at least 15  $\mu\text{m}$  in diameter. The hardness and strain-rate sensitivity were measured using Nanoinstruments NANO INDENTER® II with a Berkovich tip. Measurements were conducted in displacement-control mode at a rate of 10 nm/s to a maximum depth of 300 nm. Thus, the indent was significantly smaller than the particle, but significantly larger than the grain size. The absence of surface cracks due to indentation was confirmed by scanning electron microscopy (SEM). To determine the strain-rate sensitivity, the indenter tip was kept at constant load, 20 mN, and the displacement of the tip was monitored as a function of time.<sup>7</sup> Each reported value of hardness or strain-rate sensitivity was obtained by averaging 25–30 measured values. The standard deviation of the mean was indicated as an estimate of error due to experimental scatter. DSC measurements were performed with a Perkin-Elmer DSC 7. Approximately 40 mg of nanocrystalline Fe powder were sealed in Al pans and scanned at a heating rate of 20 K/min to the maximum temperature of 750 K. To determine the baseline, a second run was conducted without interruption, using the same scanning conditions, and its trace was subtracted from that of the first run to obtain the intrinsic heat evolution of the sample. For each sample condition, curves from two individual runs were averaged.

Samples were observed using a JEOL 3011 HREM. The grain-boundary structure is only visible by HREM when the specimen thickness is similar to the grain size and when the grain boundary is aligned edge-on. To prepare sufficiently thin nanocrystalline Fe specimen ( $\sim 10$  nm), the powder particles were embedded into a Cu foil. The resulting composite was chemically jet thinned from the Cu side until it was perforated, using an electrolyte consisting of 10% perchloric acid and 90% methanol at  $-50$  °C, with an applied voltage of  $\sim 10$  V and current density of  $\sim 7$  mA/mm<sup>2</sup>, respectively. The thinnest portion, near the perforation, was used for HREM. To avoid possible structural damage to the samples, mechanical grinding or ion milling was avoided.

## RESULTS

The average grain size of the as-milled samples, determined by x-ray diffraction,<sup>23</sup> was 9.8 nm. This agrees with our HREM observation. In Fig. 2, the grain size and rms strain of samples annealed in the Ti-gettered flowing-Ar furnace at 80 and 100 °C are displayed as a function of annealing time. The rms strain decreases slightly at the early stage of the annealing and then remains nearly unchanged. The grain size remains nearly constant for the entire range of annealing time. In order to observe changes in enthalpy content upon annealing, two samples were investigated by DSC: as-milled nanocrystalline Fe and a milled sample, subsequently annealed at 100 °C for 24 h in vacuum. In Fig. 3, the DSC traces, each obtained by averaging two curves, measured at a heating rate of 20 K/min, are displayed. The as-

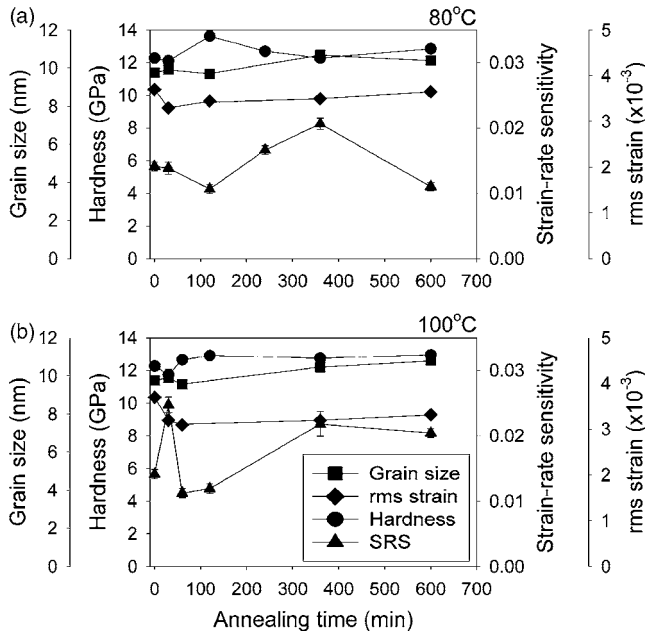


FIG. 2. Grain size, rms strain, hardness, and strain-rate sensitivity as functions of annealing time at (a) 80 °C and (b) 100 °C. The samples were annealed in Ti-gettered Ar.

milled sample shows an exothermic peak starting at around 370 K, but the annealed sample does not show a significant signal up to around 400 K. A shoulder, observed on the exothermic peak of the as-milled sample around 445 K (arrow in Fig. 3), is not present for the annealed sample. The enthalpy releases of as-milled and annealed nanocrystalline Fe in the range of 323–552 K (between two dashed lines in Fig. 3) were 950 and 630 J/mol, respectively. Thus, annealing leads to a decrease of 320 J/mol ( $\pm 20\%$ ) in stored enthalpy.

The hardness and strain-rate sensitivity of the nanocrystalline Fe samples, which were annealed at 80 and 100 °C for various times, were measured by nanoindentation. The results are displayed in Fig. 2, together with grain size and rms strain, as a function of annealing time under Ti-gettered Ar. The hardness varies just slightly with annealing time. The initial strain-rate sensitivity of the as-milled sample, 0.012, increases above 0.2 after 6 h for both temperatures. At 80 °C, the strain-rate sensitivity decreases slightly below the initial value after 10 h, but at 100 °C, it is still higher than

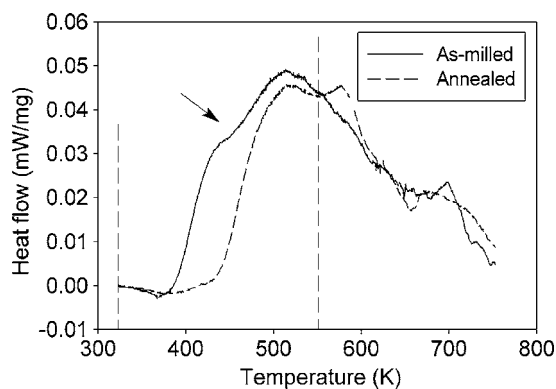


FIG. 3. DSC traces for as-milled and annealed (24 h at 100 °C) nanocrystalline Fe.

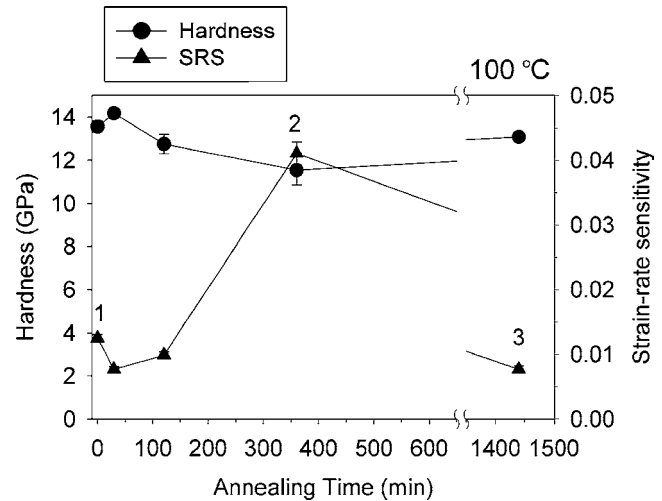


FIG. 4. Hardness and strain-rate sensitivity as functions of annealing time for samples annealed in vacuum at 100 °C.

the initial value after 10 h. Furthermore, at 100 °C, a narrow peak is observed at 30 min. A similar narrow peak was also observed in samples that were formed by low-energy ball milling and subsequently annealed at 100 °C for up to 6 h under Ti-gettered Ar. It is noted that the hardness variations follow closely those of the strain-rate sensitivity, but in the opposite direction. The same measurements were also performed on samples annealed in vacuum at 100 °C, as shown in Fig. 4. The initial strain-rate sensitivity of 0.012 increases above 0.04 after 6 h and then decreases slightly below the initial value after 24 h. However, the strain-rate sensitivity of samples annealed in vacuum does not peak at 30 min, as it does for the sample annealed under Ti-gettered Ar at the same temperature.

Changes in atomic configuration at or near grain boundaries during annealing were observed using HREM. Three samples were selected for observation: as-milled, annealed 6 h at 100 °C, which has the highest strain-rate sensitivity value, and annealed 24 h at 100 °C, both in vacuum. These are labeled as 1, 2, and 3, respectively, in Fig. 4. 16, 4, and 12 grain boundaries were imaged for each of the three samples, respectively. In Figs. 5 and 6, typical grain-boundary images are shown for an as-milled (1 in Fig. 4) and annealed (100 °C 24 h, 3 in Fig. 4) samples, respectively. In both Figs. 5 and 6, grain I is aligned along the  $[\bar{1}11]$  zone axis, and the three lattice fringes observed correspond to the (101), (110), and (01 $\bar{1}$ ) planes. The interplanar spacing in grain II is within 3% of the values in grain I, allowing us to identify the planes observed in grain II as (110). It should be noted that the instrumental line-to-line resolution, 0.14 nm, is close to the (200) interplanar spacing, 0.143 nm, allowing only (110) planes to be clearly resolved. This is the reason that no other lattice fringes are observed in any of the images. The angles between the (110) planes in neighboring grains are about 21° in both Figs. 5 and 6. It should be noted that, because of the rotational degree of freedom about the [110] axis in grain II, a determination of the misorientation angle between the grains is not possible. However, considering all (110) planes in grain II, the angle observed between

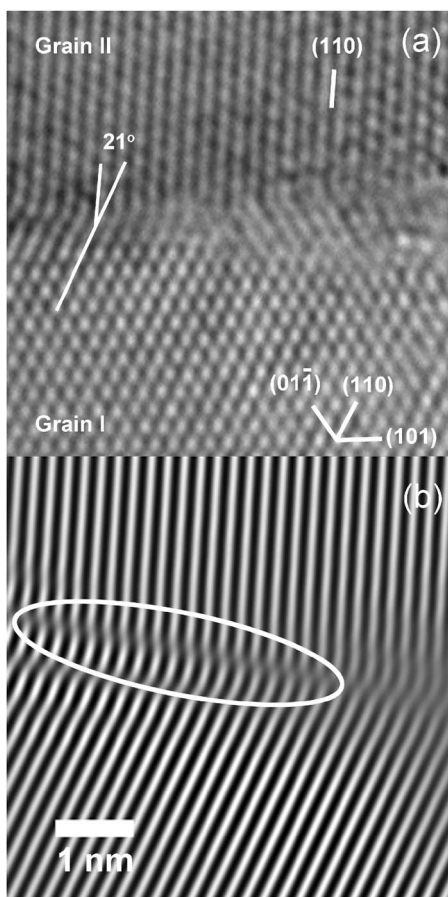


FIG. 5. (a) HREM image of a grain boundary in as-milled nanocrystalline Fe. Grain I is aligned along the  $[111]$  zone axis. The angle between the  $(110)$  planes in grain I and grain II is  $21^\circ$ . (b) Fourier-filtered image. Disconnected lattice fringes can be observed at the grain boundary (encircled).

the planes is a lower limit on the misorientation angle between the grains. Thus, the boundaries in Figs. 5 and 6 are high-angle grain boundaries.

To highlight the atomic structure at and near the grain boundary, the images were Fourier filtered, so that only one set of lattice fringes in each grain, aligned approximately perpendicular to the grain boundary, is visible [Figs. 5(b) and 6(b)]. The clearest difference *between samples* was found in the connectivity of lattice fringes at the grain boundary. In as-milled Fe, all lattice fringes are disconnected at the grain boundary, where a gap is observed [Fig. 5(b), encircled]. However, in the sample annealed at  $100^\circ\text{C}$  for 24 h, a significant portion of the lattice fringes is continuous, and the misfit of lattice planes is accommodated by regularly spaced grain-boundary dislocations [Fig. 6(b)]. Attempts to observe samples *in situ* in the HREM during heating were not successful—samples curled up instantly upon heating. Furthermore, even if such attempts were successful, the observed behavior could be severely affected by surface effects. We have confidence in our results because of their reproducibility. The observations we report above are typical—15 of 16 and 12 of 12 grains boundaries, respectively, had the same structure. For each of the two sample types, one to two of the boundaries observed were possibly low-angle boundaries [ $<15^\circ$  between the  $(110)$  planes], and even these followed

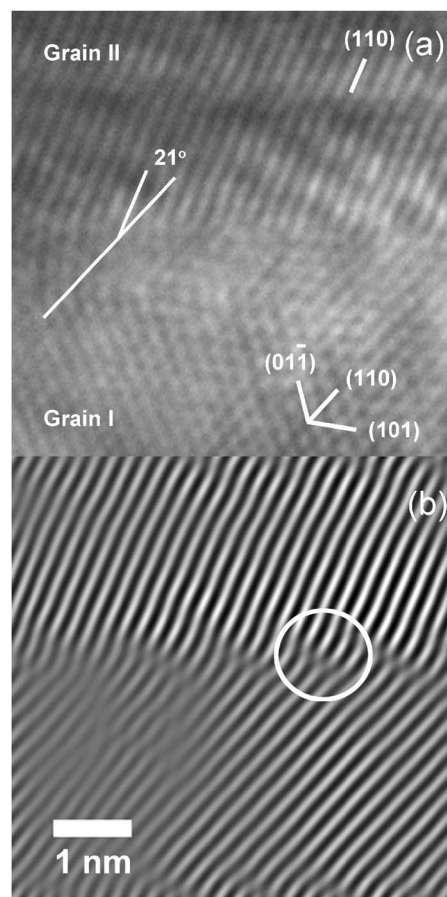


FIG. 6. (a) HREM image of a grain boundary in annealed nanocrystalline Fe at  $100^\circ\text{C}$  for 24 h. Grain I is aligned along the  $[111]$  zone axis. The angle between the  $(110)$  planes in grain I and grain II is  $21^\circ$ . (b) Fourier-filtered image. Regularly spaced grain-boundary dislocations (encircled) are observed at the grain boundary.

the same trends as the high-angle boundaries. In Figs. 7–9, grain-boundary images of a sample annealed for 6 h at  $100^\circ\text{C}$  are displayed. Unlike samples 1 and 3, whose grain-boundary structures are nearly uniform, the  $100^\circ\text{C}$  6 h annealed sample contains a mixture of the grain-boundary structures observed in Figs. 5 and 6: some grain boundaries have a structure typical of the as-milled sample (Fig. 7), and others have a structure typical of the  $100^\circ\text{C}$  24 h annealed sample (Fig. 8). Furthermore, in some boundaries, both structures are observed side by side in a single grain boundary (Fig. 9), suggesting that the grain-boundaries are in the middle of a transition between the initial and steady-state structures. It is finally noted that no lattice dislocations were found in any of the samples observed.

## DISCUSSION

The following macroscopic observations suggest relaxation of the grain boundaries: (1) the decrease in rms strain at fixed grain size, (2) the shift in onset temperature of the exothermic peak in DSC trace from  $\sim 370$  to  $\sim 400$  K (Fig. 3), and (3) the decrease in stored enthalpy by  $320$  J/mol. Since the material does not contain lattice dislocations,<sup>24</sup> the observed changes cannot be attributed to dislocations pro-

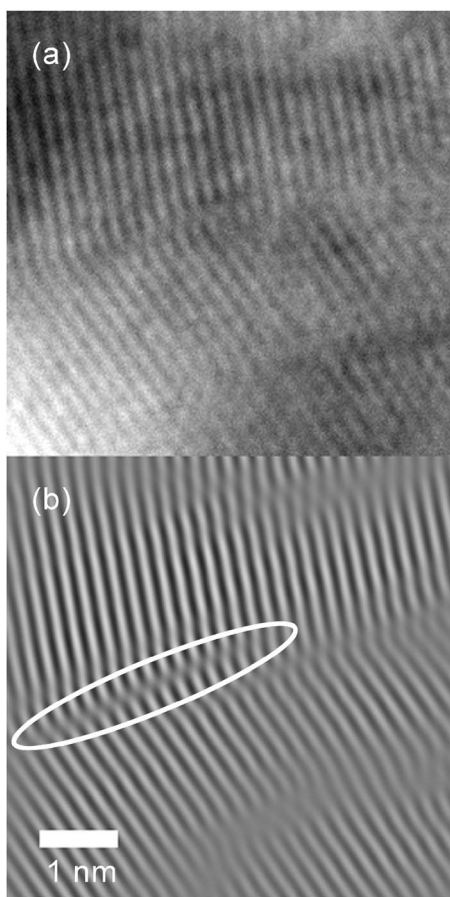


FIG. 7. (a) HREM image of a grain boundary in 100 °C 6 h annealed nanocrystalline Fe. (b) Fourier-filtered image. It has a structure similar to that of grain boundaries in as-milled nanocrystalline Fe.

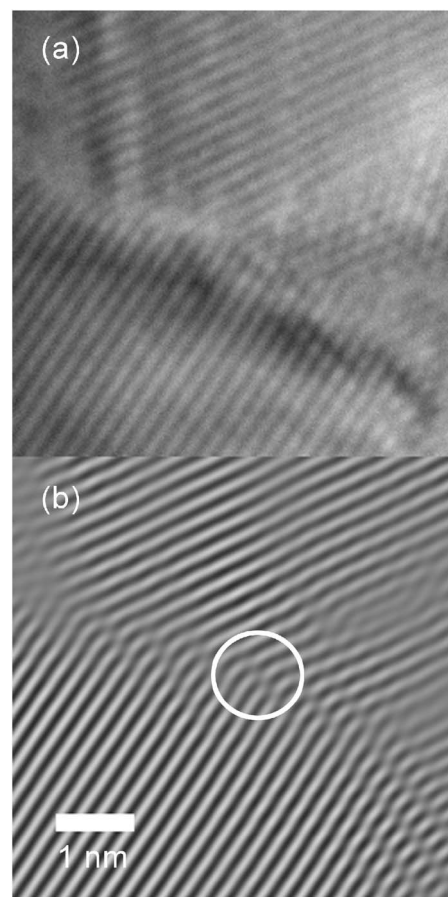


FIG. 8. (a) HREM image of a grain boundary in 100 °C 6 h annealed nanocrystalline Fe. (b) Fourier-filtered image. It has a structure similar to that of grain boundaries in 100 °C 24 h annealed nanocrystalline Fe.

cesses, such as polygonization. Consequently, it is plausible to attribute the relaxation to the grain boundaries.

Assuming the grain shape to be the regular 14-sided tetrakaidecahedron, and the grain size to be the maximum diameter of an inscribed sphere,<sup>8</sup> the grain-boundary area per unit volume is  $3.44 \times 10^8 \text{ m}^{-1}$  for an average grain diameter of 9.8 nm. Assuming the decrease in stored enthalpy, 320 J/mol, is due to grain-boundary relaxation only, it corresponds to a decrease of grain-boundary enthalpy by  $0.13 \text{ J/m}^2$ . The average grain-boundary free energy of bcc Fe is about  $1.96 \text{ J/m}^2$ ,<sup>25</sup> thus the energy decrease is about 7% of the total grain-boundary free energy. Moelle and Fecht<sup>17</sup> report for as-milled Fe powder a DSC curve that is roughly similar to ours, if one accounts for the differences in grain size and heating rates. Using the total enthalpy release due to extensive grain growth, and comparing it with theoretically calculated grain-boundary energy values, they estimate a higher contribution of relaxation to the stored enthalpy than our value. While it is possible that our annealed sample is not fully relaxed, we note that our measurement provides a more direct estimate of the enthalpy associated with relaxation. In addition, grain diameters were determined in Ref. 17 using the integral breadths of Bragg peaks, which yield results that are different from those of the Warren-Averbach method.<sup>26</sup>

The following pitfalls should be noted when interpreting

HREM images: (1) Attempts at *in situ* annealing in the HREM were not successful. The differences in the grain-boundary structures are determined only statistically based on *ex situ* observation. (2) The HREM images of nanocrystalline samples may contain artifacts such as overlap or grain boundaries that are not aligned edge-on. (3) The Fourier filtering technique, which was used to highlight the lattice fringes, may induce an artifact because the Fourier transform of a finite domain has finite-size ripples in reciprocal space. Nonetheless, the conclusion on the evolution of the grain-boundary structure is still sound because of its statistical consistency.

The grain-boundary structure is known to depend on the boundary misorientation.<sup>27</sup> Therefore, in order to study the effect of annealing on the grain-boundary structure, boundaries of similar misorientation should be compared for different samples. In the present study, two grain boundaries before and after annealing, displayed in Figs. 5 and 6, have the same lower bound on their misorientation angle,  $21^\circ$ , ruling out the possibility that the difference observed in the grain-boundary structure is due to one of them being a low-angle boundary. Furthermore, the authors observed a discontinuous grain-boundary structure in the as-milled sample for both high-angle grain boundaries and those that may be low-angle boundaries, but for neither boundary type in the sample annealed for 24 h. Because low-angle grain bound-

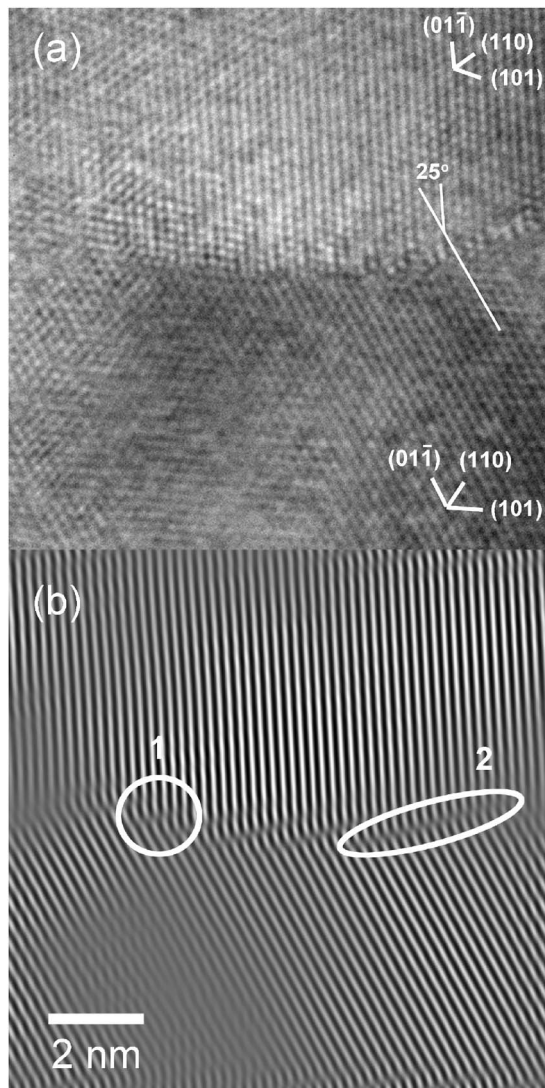


FIG. 9. (a) HREM image of a grain boundary in 100 °C 6 h annealed nanocrystalline Fe. (b) Fourier-filtered image. A single grain boundary exhibits both continuous (1) and discontinuous (2) lattice fringes.

aries generally have a more-ordered structure,<sup>27</sup> this observation strengthens the conclusion that the as-milled sample has a discontinuous grain-boundary structure, which evolves into a more-ordered structure due to annealing.

The HREM images suggest that at 100 °C, the initially discontinuous grain-boundary structure is transformed into a more-ordered structure with regularly spaced grain-boundary dislocations. The latter structure likely corresponds to a relaxed state with lower free energy. Relaxation is likely to lead to strengthening<sup>18</sup> and to a decrease of the strain-rate sensitivity. For all isothermal annealing conditions, contrary to expectation, the strain-rate sensitivity does not evolve monotonically, but goes through a peak (Figs. 2 and 4). This suggests that there is not a single process responsible for its evolution, but that at least two processes compete. Further study is required to understand which additional process may be responsible for the initial increase of the strain-rate sensitivity. Considering the effect of annealing atmosphere, we cannot rule out a role of impurities in the evolution of the strain-rate sensitivity.

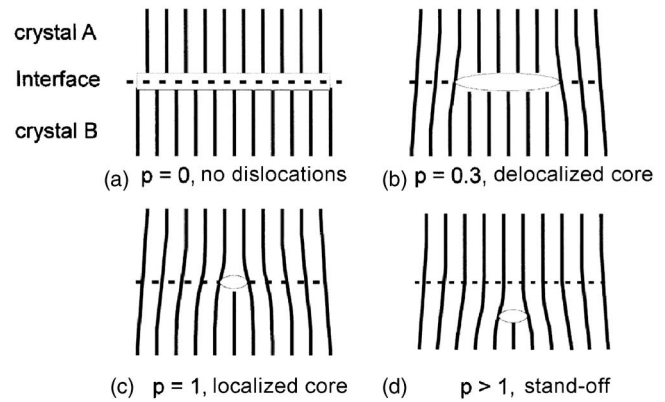


FIG. 10. The degree of matching across the grain boundary as a function of  $p$  value [after Lojkowski and Fecht (Ref. 28)]. The lattice connectivity increases from (a) to (c).  $P > 1$  is unlikely for an elemental solid.

In the presence of grain-boundary dislocations, as we report, one set of parameters describing the grain boundary are the characteristics of the dislocations, e.g., dislocation core width. The relation between the grain-boundary structure and the grain-boundary dislocation distribution was analyzed by Lojkowski and Fecht,<sup>28</sup> who introduced the concept of the localization parameter,  $p$ . It is defined as the ratio of core widths of lattice and grain-boundary dislocations or, equivalently, as the ratio of shear modulus of the grain interior and grain boundary, and provides a criterion for classifying the type of grain boundary. In Fig. 10, grain-boundary types are illustrated for different values of  $p$ .

According to Lojkowski and Fecht,<sup>28</sup> given that the localization parameter is equal to the ratio of the shear modulus of the grain interior and grain boundary, the transition of the grain boundary to a thermodynamically more-stable state can induce a grain-boundary structure change via a change in the grain-boundary shear modulus. That is, if a grain boundary relaxes by a thermal process, the increased bond strength of atoms at the grain boundary can induce an increase of lattice plane connectivity via an increase of the grain-boundary shear modulus. Observation of the grain-boundary structure of nanocrystalline Fe in the present study suggests that the transformation described above takes place during low-temperature annealing. The initially discontinuous lattice planes of as-milled Fe become continuous, with regularly spaced grain-boundary dislocations, as a result of annealing. The decreasing portion of the strain-rate sensitivity curve (Figs. 2 and 4) is possibly due to the stiffening of grain boundaries and their enhanced connectivity.

## SUMMARY

Grain-boundary relaxation and its effect on the plasticity of nanocrystalline Fe were studied using DSC, nanoindentation, and HREM. Nanocrystalline Fe samples with identical grain size, but different degrees of grain-boundary relaxation, were obtained by high-energy ball milling and subsequent annealing at low temperatures (80 and 100 °C). While the hardness was not affected significantly by annealing, the strain-rate sensitivity shows a peak value at an intermediate annealing time and decreases to slightly below its initial value, suggesting two competing processes. HREM images

suggest that grain boundaries evolve gradually from an initially discontinuous structure to a more continuous and ordered structure with regularly spaced grain-boundary dislocations. An artifact due to a bias in grain-boundary misorientation is shown to be *unlikely*. This *relaxation* process is suggested to be responsible for a decrease of strain-rate sensitivity by annealing. However, in order to closely relate the evolution of structure with that of the strain-rate sensitivity, further study is required.

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