

TESTING TECHNICOLOR THEORIES

G. L. Kane
 Randall Laboratory of Physics
 University of Michigan, Ann Arbor, MI 48109

ABSTRACT

We provide improved estimates of the masses, decay modes and widths, and production cross sections of the physical particles expected in theories with dynamical symmetry breaking. The most important results are charged pseudo-Nambu-Goldstone-bosons a_{\pm}^{\pm} with $m_{\pm} \approx 8$ GeV (and thus detectable at PETRA/PEP), two neutral pseudoscalars with $m_0 \lesssim 2.5$ GeV, and the colored technieta ($m=240$ GeV) with observable production cross sections at the Tevatron Collider and at Isabelle. The calculations were done with S. Dimopoulos and S. Raby.

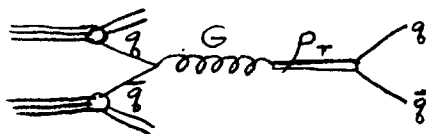
INTRODUCTION

There is not yet a dynamical symmetry breaking model which could be fully realistic. As a result, any calculations must be done in models which might not be generally applicable. Nevertheless, many features are expected to hold in any reasonable Technicolor theory,^{1,2} such as the techniquarks being colored, and the existence of interactions which couple quarks to leptons. In our calculations we use an SU(N) model which has such general features and we avoid particular assumptions which might have less generality. We give numerical results for $N=4$. Our results are given in detail in ref. 3.

There are many Goldstone bosons which arise in such a theory from breaking the original chiral symmetry. We will mention 12 of them here, the technieta color octet η_{\pm}^a , and the color-singlet light pseudoscalars ("pseudos") a_{\pm}^{\pm} , a_0^0 , \tilde{a}_0^0 . There are no light scalars in a Technicolor theory, an important prediction. The pseudo-Nambu-Goldstone bosons get mass from color, electroweak, and extended technicolor interactions^{4,5}.

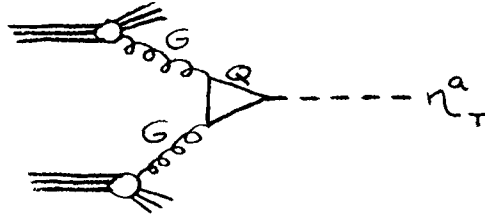
Particles come in 3 mass scales. Resonances such as the technirho, ρ_{\pm} , will occur on the mass scale of the theory, about 1 TeV. Pseudos that get mass from color, etc., will have $m^2 \sim \alpha_C m_{TC}^2 \sim (300 \text{ GeV})^2$. Those which still get no mass from color interactions will have $m^2 \sim \alpha m_Z^2 \sim (\text{few GeV})^2$.

ρ_{\pm} We find the technirho will have mass of about 900 GeV. It will be produced in $\bar{p}p$ collisions, as shown,



with $(d\sigma/dy)_{y=0} \approx 6 \times 10^{-36} \text{ cm}^2$ at $\sqrt{s} = 2000 \text{ GeV}$. This rate scales as $u(m_{\rho_T}/\sqrt{s})/m_{\rho_T}^4$ so it would increase considerably if m_{ρ_T} went down. The total width of ρ_T is $\Gamma \approx 4 \text{ GeV}$, with dominant decay modes being $\rho_T \rightarrow GG, qq, W_L^+ W_L^-$.

η_T^a A clear test of technicolor theories will come from production of η_T^a since a lower limit in its production cross section can be computed from the triangle contribution. Since chiral symmetry should be as good an approximation here as for a pion, this should be reliable. We find $M(\eta_T^a) \approx 240 \text{ GeV}$. It is produced via



and gives a large cross section because the technifermions are assumed to carry color and couple to gluons. We find

$\sqrt{s}(\text{GeV})$	$(d\sigma/dy)_{y=0}(\text{cm}^2)$
500	2×10^{-37}
800	4.5×10^{-36}
2000	44×10^{-36}

This cross section has been computed in ref. 6 also and they agree with us. Note that this result is about $8 \times N^2 = 128$ times larger than the cross section for producing a fundamental Higgs of similar mass, since η_T^a is a color octet and there are $N=4$ technifermions in the loop to sum over.

The important η_T^a decays are $\eta_T^a \rightarrow GG$ (opposite to the production) with $\Gamma(GG) \approx 60 \text{ MeV}$, and $\eta_T^a \rightarrow f\bar{f}'$ with f, f' heavy fermions. The effective coupling to the fermions is $(m_f + m_{f'})/F_T$, from models or from a Goldberger-Treiman type argument. For $m_f = 25 \text{ GeV}$, $\Gamma(f\bar{f}') \approx 1 \text{ GeV}$ and will dominate. Various modes are $\eta_T^a \rightarrow t\bar{t}, b\bar{b}, GG, G\gamma, GZ^0, GGG, \dots$. Other pseudos will have modes such as $GW^+, t\bar{b}$.

η_T^a The charged light pseudos get about 7.7 GeV of mass from electroweak interactions, while the neutrals remain massless. All get some mass from the extended technicolor interaction^{4,5}. While this contribution cannot be calculated reliably, we can put a limit on it. We make the important assumption that the same leptoquark bosons that couple technifermions to technileptons (and give mass contributions) also couple quarks to leptons (and give flavor charging neutral currents). To not violate existing limits on $K_L \rightarrow \mu e$, this implies $g^2/M^2 < (1/310 \text{ TeV})^2$. Adding this to the electroweak contribution gives charged and neutral Higgs-like particles with masses

$$m(a_T^\pm) \approx 8 \text{ GeV}$$

$$m(a_T^0) \lesssim 2.1 \text{ GeV}$$

$$m(\tilde{a}_T^0) \lesssim 2.5 \text{ GeV}$$

The a_T^\pm are charged pseudoscalars and can be found in $e^+e^- \rightarrow a_T^+ a_T^-$ at PETRA or PEP. They are produced with a β^3 threshold behavior, $1/4$ unit of R , $\sin^2\theta$ production distribution. They decay dominantly into $\tau\nu_\tau$ (about 40%), $c\bar{s}$ (like a heavy F^\pm ; about 40%), and $c\bar{b}$ (about 20%), and $\mu\nu_\mu$ (about 0.1%). About 60% of the events have 4 strange quarks.

Interestingly, if a_T^\pm exists, the mode $t \rightarrow ba_T^+$ dominates t decay since it is semiweak, and t decays are not as in the standard model.

The neutrals a_T^0, \tilde{a}_T^0 can be produced in Drell-Yan reactions (K beams are best for good signal/noise), or in decay of heavier states such as ψ, T . Their main modes are $a_T^0, \tilde{a}_T^0 \rightarrow \mu^+\mu^-$ (about 1/3), $K^*\bar{K}, \phi\phi, \Lambda\bar{\Lambda}, K\bar{K}\pi$, and perhaps the parity violating mode $K\bar{K}$. They are pseudoscalars that may have parity violating couplings to fermions but not to $\gamma\gamma, GG$.

The above predictions can be tested, and guarantee that soon (finally) there will be experimental input into understanding the origin of spontaneous symmetry breaking.

REFERENCES

1. See M. A. B. Bég, Proceedings of Orbis Scientiae, 1980, Coral Gables, (see ref. 2, for a recent review of the theory).
2. L. Susskind, these proceedings.
3. S. Dimopoulos, S. Raby, and G. L. Kane, Michigan preprint UM HE 80-22.
4. S. Dimopoulos and L. Susskind, Nuc. Phys. B1555 237 (1979).
5. E. Eichlen and K. Lane, Phys. Lett. 90B 125 (1980).
6. F. Hayot and O. Napoly, Saclay preprint.