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Dehumidification Over Tropical Continents Reduces Climate Sensitivity and Inhibits Snowball Earth Initiation

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Dehumidification over Tropical Continents Reduces Climate Sensitivity and Inhibits Snowball Earth Initiation Richard P Fiorella¹ and Christopher J Poulsen¹ 1: Department of Earth and Environmental Sciences, University of Michigan, Ann Arbor, MI Corresponding author address: Richard P. Fiorella, Department of Earth and Environmental Sciences, University of Michigan, 2534 CC Little Building, 1100 N. University Avenue, Ann Arbor, MI 48109-1005, USA Email: richf@umich.edu

Abstract

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The enigmatic Neoproterozoic geological record suggests the potential for a fully glaciated "Snowball Earth." Low-latitude continental position has been invoked as a potential Snowball Earth trigger by raising surface albedo and reducing atmospheric CO₂ concentrations through increased silicate weathering. Herein, climate response to reduction of total solar irradiance (TSI) is tested using four different land configurations (aquaplanet, modern, Neoproterozoic, and lowlatitude supercontinent) with uniform topography in the NCAR Community Atmosphere Model (CAM, version 3.1) general circulation model with a mixed-layer ocean. Despite a lower surface albedo at 100% TSI, the threshold for global glaciation decreases from 92% TSI in the aquaplanet configuration to 85% TSI with a low-latitude supercontinent. The difference in thresholds is principally due to the sensitivity of total specific humidity and therefore greenhouse forcing to reductions in TSI. Dehumidification of the troposphere over large tropical continents decreases greenhouse forcing and also increases direct heating by decreasing cloud cover. Continental heating intensifies the Walker circulation and the transport of dry air over the ocean, enhancing surface evaporation and marine tropospheric humidification, maintaining a high specific humidity and greenhouse effect over the ocean. Topography also provides an important control on Snowball Earth initiation. Modern topography in the modern continental arrangement lowers the initiation threshold by up to 2% TSI relative to a modern continental arrangement without topography. In the absence of potential silicate weathering feedbacks, large tropical landmasses raise the barrier to initiation of Snowball events. More generally, these simulations demonstrate the substantial influence of geography on climate sensitivity.

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1. Introduction

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Several lines of geologic evidence from the mid and late Neoproterozoic suggest the episodic occurrence of prolonged low-latitude glaciation and perhaps a "Snowball Earth" resulting from global glaciation. Cryogenian (850-630 Ma) glaciogenic sequences and diamictites are found on every continent and are frequently overlain by thick "cap" carbonate sequences. Large and chemostratographically correlated carbon isotope anomalies in these cap carbonates are thought to reflect the global nature of mid and late Neoproterozoic glaciations (Hoffman and Schrag 2002; Macdonald et al. 2010). Iridium anomalies found at the base of cap carbonates in Africa suggest that the Marinoan glaciation episode lasted at least 3 million years, but more likely 12 million years (Bodiselitsch et al. 2005), though these results have not been duplicated in other cap carbonate sequences. Finally, the return of banded iron formations in the Neoproterozoic after a billion year hiatus (Hoffman and Schrag 2002; Pierrehumbert et al. 2011) has been argued to be evidence of ocean anoxia during the glaciation event, though the chemical conditions facilitating their deposition in the Neoproterozoic were likely different and more directly dependent on ocean composition than during the Paleoproterozoic. Deep ocean anoxia during Neoproterozoic glaciations may have been brought about by tropical pack ice interfering with and weakening wind-driven ocean circulation or total isolation of the ocean from the atmosphere by sea ice (Kirschvink 1992). Termination of Snowball Earth is thought to have resulted from buildup of atmospheric CO₂ levels over millions of years as silicate weathering essentially stopped during the Snowball event while volcanic outgassing would have continued unabated. Supporting this suggestion, strongly negative Δ^{17} O anomalies in Neoproterozoic sulfate minerals precipitated after termination of Snowball events suggest higher CO₂ levels than at any point during the Phanerozoic (Bao et al. 2008; 2009).

However, not all geologic evidence is compatible with Snowball Earth. Biomarker evidence of eukaryotic life dates back to the Archean (Waldbauer et al. 2009) while evidence of sponges exists prior to the Marinoan glaciation (Love et al. 2008). Therefore, emergence of Metazoan life predated the Snowball events and Metazoan survival calls into question likely global temperatures during the existence of a completely glaciated ocean. Additionally, while deposition of cap carbonates is thought to reflect amplified continental weathering or overturning of a stratified ocean following deglaciation (Hoffman and Schrag 2002), critics have questioned how these mechanisms could produce widespread and synchronous carbonate deposition (Ridgwell et al. 2003; Shields 2005).

Efforts using climate models to understand Neoproterozoic climate have only stoked the Snowball debate. According to standard solar evolution models, solar luminosity during the Neoproterozoic would have been approximately 6-7% lower than present (Gough 1981; Crowley and Baum 1993). Energy balance models (EBM) developed independently by Budyko (1969) and Sellers (1969) suggest global glaciation driven by a runaway ice-albedo feedback could result from modest reductions in solar luminosity (as little as 1.6%) after ice expands equatorward beyond a critical latitude. Improvements to EBM parameterization of heat transport and thermal diffusion made by North (1975a,b) suggested global glaciation might occur with reductions in solar luminosity closer to 7%.

Snowball Earth initiation has also been explored in several general circulation models (GCMs). Larger reductions of total solar irradiance (TSI) are typically required in GCMs to simulate a Snowball state than in EBMs, suggesting that climate dynamics and feedbacks not captured in EBMs restrict ice expansion. In support of this conclusion, previous studies have identified processes including ocean heat transport (Poulsen et al. 2001; Bendtsen 2002), wind-

driven ocean circulation, and cloud-radiative forcing (Poulsen and Jacob 2004) that restrict seaice growth. For example, clouds over sea-ice restrict surface sensible heat loss and have a warming effect over sea-ice. The ice-albedo feedback is weaker than when cloud radiative forcing is excluded (Poulsen and Jacob 2004). Likewise, ocean circulation transports heat to and stabilizes the ice line. Convective mixing warms the surface ocean adjacent to the ice margin restricting ice growth, with the largest impact occurring in the winter (Poulsen et al. 2001).

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Nonetheless, a wide range of solar luminosity thresholds for Snowball initiations have been reported (Pierrehumbert et al. 2011). With Marinoan continents, the ECHAM/MPI-OM model predicts a Snowball state when solar luminosity is reduced below 96% of the modern value (Voigt et al. 2011). With comparable greenhouse gas concentrations, however, other models including FOAM (e.g., Poulsen et al. 2001; Poulsen 2003), GISS (e.g., Chandler and Sohl 2000), and GENESIS (e.g., Baum and Crowley 2001; 2003) fail to simulate Snowball states when solar luminosity is reduced below 94%. Other solutions in addition to global ice cover have also been invoked to explain the geologic evidence. "Slushball" or "waterbelt" solutions (e.g., Hyde et al. 2000; Chandler and Sohl 2000; Baum and Crowley 2001; 2003; Micheels and Montenari 2008) describe a state with glaciated continents but ice-free tropical oceans and ice margins poleward of 25° latitude. In these studies, however, a large portion of land was positioned in mid and high latitudes, at odds with paleogeographic reconstructions for the Marinoan and Sturtian glaciations (Hoffman and Li 2009). An additional solution, the "Jormungand" state—named for the serpent-like appearance of a narrow band of iceless equatorial ocean in a Hovmöller diagram—emerges when the albedo of snow-covered and bare sea-ice differ (Abbot et al. 2011). These open-ocean solutions provide an attractive explanation for glaciogenic deposits in low paleolatitudes and provide a refuge for life, but may fail to

account for the deposition of banded iron formations since the wind-driven ocean circulation would likely prevent the deep oceans from becoming anoxic.

Across these models, a myriad of paleogeographies has been used and may in part be responsible for the range of observed sensitivities. Here, we investigate the hypothesis that low latitude continents facilitate Snowball Earth initiation (Kirschvink 1992). Multiproxy reconstructions of Neoproterozoic geography place the bulk of the land area in the tropics and the subtropics (Kirschvink 1992; Li et al. 2008; Hoffman and Li 2009), a distribution that has been considered conducive to Snowball Earth initiation for three major reasons. First, land surfaces generally have higher albedo than ocean, particularly under the relatively cloud-free descending limb of the Hadley cell in the subtropics where surface albedo has a stronger impact on planetary albedo. Second, as continental glaciers form and sea-level drops, land would replace shallow seas in the tropics, further increasing albedo (Kirschvink 1992). Finally, weathering rates of silicate minerals would increase with increased land under the intertropical convergence zone, leading to an enhanced drawdown of atmospheric CO₂ under a constant volcanic degassing rate (Marshall et al. 1988; Worsley and Kidder 1991).

Previous studies provide conflicting conclusions about the impact of increased tropical landmass on global average temperatures. GCM experiments generally suggest higher global temperatures when more land is concentrated in the tropics. For example, Barron et al. (1984) show tropical continents suppress evaporative cooling and warm the tropics and Poulsen et al. (2002) suggest mid and high-latitude continents promote cooling by increasing snow coverage over continents and sensible heat loss from the ocean, leading to greater global ice coverage. In contrast, Voigt et al. (2011) argue that increased tropical landmass cools global temperatures by increasing planetary albedo and weakening the greenhouse effect. Here, we analyze how climate

sensitivity to reduction in solar luminosity varies with paleogeography and describe its implications for snowball Earth initiation.

2. Model Description and Setup

Snowball Earth initiation experiments were performed using the National Center for Atmospheric Research Community Atmosphere Model (CAM) version 3.1 (Collins et al. 2004), consisting of an atmospheric model coupled to a 50-m mixed-layer ocean model and the Common Land Model (CLM) version 3.0 (Oleson et al. 2004). A horizontal resolution of spectral T31 (~3.75°) was used in the atmospheric model, with 26 levels in the vertical using a terrain-following vertical coordinate (Collins et al. 2004). The model top is in the stratosphere at 3.5 hPa, and model levels are roughly split equally between the troposphere and stratosphere. No additional heat transport was specified for the mixed-layer ocean, and dynamic adjustments to the mixed-layer ocean heat flux were disabled. Sea-ice is modeled purely on a thermodynamic basis with no dynamics included. To generate a better match with modern sea ice distributions, the mixed-layer ocean model in CAM adds 15 W m⁻² to the Northern Hemisphere ocean and removes 10 W m⁻² from the Southern Hemisphere ocean by default; these values were set to zero everywhere in our experiments.

Neoproterozoic boundary conditions are not well constrained. Model boundary conditions were chosen to best represent the Neoproterozoic, to facilitate glaciation, and to simplify analysis of results. Eccentricity was set to zero and obliquity was held constant at 23.5°. Carbon dioxide (280 ppm), aerosol, and ozone concentrations were set to preindustrial levels. Model default values were used for methane (1714 ppb) and nitrous oxide (316 ppb); sensitivity tests indicate this choice does not fundamentally alter the results presented here. As land plants

had not evolved in the Neoproterozoic, all land vegetation was removed and the land surface specified as desert. CLM calculates the land surface albedo as a function of soil color and the moisture content of the topsoil layer. The lightest, most reflective soil color is used, having visible/near-infrared (NIR) albedos of 0.24/0.48 when the top layer of soil is dry, and decreasing to 0.12/0.24 when the top soil layer is saturated (Oleson et al. 2004). Ocean albedo is 0.06 for direct sunlight in these experiments. The albedo of sea ice increases linearly with thickness up to 1 m, at which albedos are capped at 0.67 in the visible and 0.30 for NIR wavelengths. Cold snow has an albedo of 0.91/0.63 for visible/NIR wavelengths.

We test four different continental arrangements: (a) an aquaplanet with no land (hereafter AQP), (b) an equator-centered rectangular supercontinent extending to 43°N/S and spanning 130° degrees in longitude (hereafter RECT), (c) a continental reconstruction with two large tropical continents (hereafter MAR), based on the Marinoan (~635 Ma) paleogeography used in Voigt et al. (2011), and (d) a modern continental configuration with topography removed (hereafter MOD-NT) (see Fig. 1 for configurations). A summary of all of the runs performed is provided in Table 1. Global land fraction was similar across the configurations with 32.3%, 27.8% and 27.0% land coverage for the MOD-NT, MAR, and RECT experiments respectively (Table 2). Low-latitude (≤30°) land fraction is 28.8%, 48.1% and 38.5% for the MOD-NT, MAR, and RECT experiments. Land surface elevations are 100 m in all experiments.

For the RECT, MAR, and MOD-NT configurations, atmospheric temperatures were initialized with a cosine latitudinal gradient ranging from 28°C at the equator to 12°C at the poles, and initial wind conditions were calculated using the thermal wind. Sea-surface temperatures were prescribed using results from a coupled CCSM run simulating modern climate. The model was run at 100% of the modern solar constant for 100 years, after which branch runs

were performed by immediately reducing TSI to values as low as 80% (Table 1). A 1% reduction in TSI corresponds to a global average change in incoming solar radiation at the top of the atmosphere of 3.4 W m⁻². As stable climate states lacking permanent sea ice have been reported for modern TSI (e.g., Pierrehumbert et al. 2011), additional sensitivity experiments were performed with no initial ice and initial atmosphere and ocean temperatures set to 300K everywhere. Results presented here were insensitive to this change in initial boundary conditions.

In all cases, the model was run until the surface temperature trend over a twenty year period was less than 1°C per century, signifying a quasi-equilibrium state was reached. To ensure these equilibria are stable, we integrated several simulations for additional time, particularly near the bifurcation points (Table 1). Results shown here represent averages of the final 20 years of model integration.

3. Results

a) Climate sensitivity to TSI

At 100% TSI, mean annual surface temperatures are similar for all four geographies, ranging between 272.4 and 274.0K (Table 2). In all cases, surface temperatures are highly zonal due to the lack of topography, with some asymmetries near the edges of continents (Figs. 1a-d), presumably due to land/sea thermal contrasts. An equator-to-pole temperature gradient of 70K is simulated. Equatorial temperatures average ~295K while polar temperatures average ~225K for all four geographies. Global surface albedo varies from a maximum of 0.356 for the MAR geography to a minimum of 0.310 for the AQP geography due to differing meridional distribution of land. The range in planetary albedo between the geographies at 100% TSI is more than five times smaller than the surface albedo range (0.406 for AQP and 0.397 for MOD-NT),

and as a result, surface temperatures only vary by 1.6K. The mean climatological positions of sea ice are also similar for all four geographies, ranging from a maximum equatorward extent of 39.7°N/39.3°S (RECT) to a minimum equatorward extent of 42.7°N/40.5°S (MOD-NT).

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When TSI is reduced to 95%, prominent differences between the AQP and MOD-NT, MAR, and RECT simulations emerge. Despite a surface albedo 13% lower at 100% TSI than the MAR geography, the AQP configuration shows the highest sensitivity to TSI (Figs. 1e-h). AQP equatorial temperatures decrease to 278K and polar temperatures decrease to 190K, yielding an increased equator-to-pole temperature gradient of 88K. In contrast, equatorial and polar temperatures fall to 285-290K and 200-210K for the three geographies with land. As TSI is further reduced, AQP continues to cool most rapidly with TSI, but the MOD-NT, MAR, and RECT climatologies diverge as ice expands equatorward and show a greater range of sensitivity to TSI at lower values (Fig. 2). For example, at 90% TSI the MOD-NT configuration has an average surface temperature 4.3K cooler than the MAR and 7.2K cooler than the RECT simulations. Differing sea-ice cover drives these temperature differences. MOD-NT has the largest ice extent at this TSI (Fig. 2c), particularly in the Southern Hemisphere where there is less land, followed by MAR which has slightly greater ice expansion in the Northern Hemisphere compared to RECT. As a result, planetary albedo increases more in the MOD-NT and MAR cases than in the RECT case.

As a result of these varying sensitivities, the threshold TSI for snowball state initiation differs between geographies. Though the configurations with land concentrated in the low-latitudes have slightly higher initial surface albedo at 100% TSI, both surface and planetary albedo for the AQP configuration increases more rapidly with reductions in TSI and is higher at TSI below 100% (Figs. 2c,d). All four configurations exhibit a strong bifurcation point at the

transition between partial and global ice coverage. AQP enters the snowball state at the highest solar constant, transitioning to full ice coverage between 92 and 91% TSI. The snowball transitions occur at 88, 86, and 85% of modern TSI for the MOD-NT, MAR, and RECT configurations. Across the AQP, MOD-NT, and MAR experiments, the reduction in TSI necessary to induce a snowball state increases with low-latitude land fraction. Despite having less low-latitude land area than MAR, the RECT configuration experiences global sea-ice cover at a lower TSI, suggesting that continentality also impacts snowball initiation.

Geography also impacts the global mean surface temperature at the bifurcation point, as the minimum surface temperature without global ice cover is 233.3K for AQP but 237.2K for RECT. Similarly, the latitude of maximum ice extent before transition to global ice cover for AQP is closer to the equator (10.6°N/9.4°S) than for any of the other configurations. The minimum latitude of stable ice margins for the MOD-NT and MAR experiment are 11.8°N/10.0°S and 10.5°N/13.6°S, respectively, while the minimum stable ice margin for the RECT experiments is 12.4°N/10.4°S. In the following sections, we address the energetic and dynamical differences leading to the highly varied responses in TSI reduction and snowball earth initiation to differences in geography.

b) Tropical Energy Balance: Cloud and Water Vapor Feedbacks

Low-latitude continents are thought to facilitate global ice cover by reducing absorbed shortwave radiation though increased surface albedo. Our results suggest something different: particularly in the deep tropics, the net shortwave radiation received at the surface is greater for geographies with more tropical land despite a higher surface albedo because of the near absence of clouds over large land masses. Equatorward of 10°, absorbed shortwave radiation is 30 W m⁻² and 20 W

m⁻² higher for the RECT and MAR geographies than for the AQP geography (Fig. 3a).

Longwave radiation shows an even larger disparity between the RECT and AQP geographies, with emission being ~40 W m⁻² higher in these latitudes in the RECT case (Fig. 3b). Clouds are universally cooling in the tropics at 100% TSI, but the amount of cooling varies by 20-30 W m⁻² between the geographies, with the AQP and MOD-NT geographies having the most negative cloud forcings (Fig. 3c). This difference in cloud forcing results from stark differences in tropical cloud fraction between the geographies (Fig. 4). Cloud fractions over land, particularly in the subtropics, tend to remain below 0.3 while exceeding 0.6 over the oceans at comparable latitudes. Presence of continents also reduces cloud fractions over the ocean near land (Fig. 3e, 4) increasing cloud radiative forcing over the ocean. This effect is most pronounced in the RECT case, where cloud forcing over the ocean in the deep tropics are 20 W m⁻² higher than in the AQP case at 100% TSI.

Geography alters tropical energy balance, as shown by the large disparities in surface shortwave absorption and longwave emission between cases. The higher sensitivity of AQP to TSI reduction is related to its larger dependence on greenhouse forcing to maintain energy balance (Fig 5). Specific humidity (q) and surface temperature (T) are tightly coupled through the Clausius-Clapeyron equation in all of our simulations (Fig. 6a). Reductions in TSI reduce direct solar heating linearly, while greenhouse forcing decreases nearly exponentially above the bifurcation points for all geographies (Fig 5d). Therefore, TSI reductions have the strongest impact on global temperatures in AQP resulting from its large greenhouse forcing term in the tropical energy balance. In contrast, configurations with land have a larger radiatively direct proportion of the energy budget and therefore show lower sensitivities to TSI reductions than AQP. Among configurations with land, the RECT (MOD-NT) configuration has the largest

(smallest) radiatively direct energy budget portion (Figs. 3,5), and shows the lowest (highest) sensitivity in q to changes in TSI (Fig. 6b). Tight q-T coupling suggests dynamical differences between the experiments affect energy partitioning that drives the different sensitivities to changes in TSI observed.

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c) Amplified Walker-like circulation in configurations with high tropical land area For clarity, we focus on the differences between the AQP and RECT configurations to assess the impact of large land-sea contrasts on zonal circulation patterns in the deep tropics. AQP shows strong upward motion at all longitudes to about 350 mb (Fig. 7). The atmosphere is nearly saturated at the surface and in the upper troposphere. In contrast, RECT shows strong and deeper upward motion over the ocean (to about 250 mb), but strong radiatively driven subsidence in the upper and mid troposphere over the continent, limiting upward motion to the boundary layer and leading to low relative humidity over the continent. Relative humidity drops below 5% over the western third of the continent. Temperatures over the oceans decrease with altitude following the moist adiabatic lapse rate, while dry subsiding air over the continent warms near the dry adiabatic lapse rate during adiabatic descent. As a result of this increased lapse rate as well as increased direct solar heating in the absence of clouds, temperatures at the surface are warmer than they would be were there no continent. Surface winds advect this warm, dry continental air over the oceans, increasing surface evaporation (Fig. 8). Though higher amounts of land decrease the total water flux from tropical oceans, evaporation is more efficient over the ocean in the MAR and RECT cases. Warmer global temperatures and enhanced ocean surface water fluxes yield higher specific humidity over the oceans for the RECT configuration than for the AQP at lower TSIs, particularly in the mid and upper troposphere (Fig. 7) where water vapor has a more

potent greenhouse effect relative to the boundary layer (Held and Soden 2000). The MAR and MOD-NT configurations also show a Walker cell, but both are weaker than the RECT cell as the amount of contiguous land in the deep tropics is lower and land-sea contrasts are less extreme.

Higher moisture contents over the RECT ocean imply the potential for greater latent heat release, driving deeper convection. Convective precipitation rates are double those over the ocean at 95 and 93% TSI in the RECT geography compared to AQP. Deeper convection and increased atmospheric moisture content changes the vertical distribution of clouds (Fig 9) and may be responsible for the increased cloud forcing over the oceans in the RECT case compared to the AQP case. As noted above, clouds are nearly absent over the continent but over the oceans, midlevel cloud cover decreases and high-level cloud cover increases in the tropics and subtropics (Fig. 9). With higher moisture contents, low-level clouds tend to have a stronger impact on shortwave forcing due to higher optical depth while high-level clouds tend to have a stronger warming effect through increased longwave forcing (Hartmann 1994). As a result, the shift in the vertical distribution of clouds results in 20-40 W m⁻² more shortwave warming in the RECT case while longwave warming is only reduced by 10-15 W m⁻², suggesting that the enhanced Walker circulation warms the deep tropics by 5-30 W m⁻² through altering the cloud distribution.

d) Effect of topography in the Modern Simulations

To assess the impact of topography on Snowball initiation, we performed additional simulations using the modern continental arrangement with modern topography (MOD-WT). Including topography lowers the Snowball Earth initiation threshold from 90% in the MOD-WT case to 88% in the MOD-NT case. At TSI higher than the bifurcation point, however, mean annual

surface temperature and surface and planetary albedo remain similar (Fig. 10), indicating that the change in the bifurcation point cannot be explained by a difference in global albedo.

Topography reduces the equator-to-pole temperature gradient by up to 15K, a change that is attributed to increases in subtropical albedo and decreases in midlatitude albedo (particularly in the northern hemisphere), which change the albedo distribution without changing planetary albedo (Fig. 11b). The meridional distribution of albedo changes in response to a decrease (increase) of cloud cover (Fig. 11c) and snow depth (Fig. 11d) in the midlatitudes (subtropics). Though Figure 11 shows results for 90% TSI, the trend shown is present for all TSI above the bifurcation point. The magnitude of this albedo change increases with decreasing TSI, driving an increasingly small equator-to-pole temperature gradient relative to MOD-NT in the MOD-WT simulations as TSI is reduced to the bifurcation points.

This redistribution of albedo alters the top of the atmosphere radiative balance, warming the midlatitudes and cooling the subtropics in MOD-NT. Temperatures in the northern hemisphere midlatitudes increase by 15K. Ice expands more rapidly with reduction in TSI with a decreased meridional temperature gradient as a negative global temperature perturbation plunges a larger area below freezing causing greater expansion of sea-ice. This effect is especially apparent as ice encroaches into the tropics where the meridional temperature gradient is weaker than in the subtropics and midlatitudes.

4. Discussion

a) Paleogeographic controls on Snowball Earth initiation

The forcings required to initiate a snowball state for all continental configurations tested here require reductions in TSI below expected Neoproterozoic values based on the standard solar

evolution model. Therefore, our results do not directly support the "Hard Snowball" hypothesis, though it is possible radiation receipt at the surface may have been lower through mechanisms other than reduced solar irradiance, such as an increase in volcanic aerosols or a lower concentration of greenhouse gases than used in these simulations. Our results do suggest, however, that if Neoproterozoic glaciations were global, concentration of land in the tropics was unlikely to be the facilitating factor.

Land surface albedos in these simulations are lower than those in previous studies. Desert soil albedos (0.35 and 0.51 in the visible and near-infrared) in the FOAM (e.g., Poulsen et al. 2001; Poulsen 2003; Poulsen and Jacob 2004) and GENESIS (e.g., Baum and Crowley 2001; 2003) models are based on the biosphere-atmosphere transfer scheme (Dickinson et al. 1993). Surface albedos in the GISS model (0.35 in both visible and near-infrared, Hansen et al. 1983) are slightly lower in the near-infrared than the FOAM or GENESIS models. Land albedo may play a strong role in setting the strength of the Walker circulation observed through modulating the intensity of radiatively driven subsidence over the continent as well as the size of the radiatively direct portion of the surface energy budget. Land surface albedos in the Neoproterozoic are poorly constrained, but albedos lower than those in modern deserts might be reasonable as a result of land surface coverage by ancient terrestrial microbial communities (Lenton and Watson 2004; Knauth and Kennedy 2009).

The impact of paleogeography on Snowball initiation within one model may be as large as variation between models. For example, Snowball initiation experiments for modern continents with modern topography with CCSM3 (CAM coupled to dynamic ocean and sea-ice models, Collins et al. 2006) by Yang et al. (2012a) yield a Snowball state with a reduction between 10 and 10.5% TSI; our experiments predict a Snowball state with a reduction between

10 and 11% TSI. In contrast, larger TSI reduction are required for the RECT geography in CAM than required in FOAM for a similar idealized supercontinent (Poulsen and Jacob 2004). In Poulsen and Jacob (2004), a CO₂ concentration of 140 ppm was used, reducing greenhouse radiative forcing by 3.7 W m⁻² (Myhre et al. 1998). Assuming a fairly low planetary albedo of 0.3 near the bifurcation point, however, the extra 6% reduction in TSI necessary to initiate a Snowball Earth in our experiments represents a 14.4 W m⁻² change in top of the atmosphere radiative forcing, indicating Snowball states are more difficult to simulate in CAM. Finally, much larger reductions in TSI for Snowball Earth initiation are required herein than for experiments performed in ECHAM5/MPI-OM for both the modern (Voigt and Marotzke 2009) and Marinoan (Voigt et al. 2011) continental configurations.

A large portion of the initiation threshold range can be attributed to model differences in albedo parameterizations (Pierrehumbert et al. 2011). Ice albedos in FOAM exceed ice albedos in CAM by ~0.10, suggesting a stronger ice-albedo feedback in FOAM than CAM with all else equal. Supporting this conclusion, setting both FOAM and CAM ice albedos to 0.6 yields identical initiation thresholds (Pierrehumbert et al. 2011). Comparison between ECHAM and CAM is more challenging. Ice albedo in ECHAM (0.75 below -1°C, Roeckner et al. 2003) is similar to the snow covered sea-ice albedo in CAM (broadband albedo ≈0.74 for T31, Collins et al. 2004; Pierrehumbert et al. 2011). CAM tracks snow on sea-ice and assigns a higher albedo to snow covered sea-ice, however, while ECHAM does not. As ice entering in the subtropics is more likely to be bare, model solutions for ECHAM and CAM diverge when ice encroaches on the subtropics (Abbot et al. 2011), therefore, a stronger ice albedo feedback may in part explain the difference between our results in CAM and reported ECHAM/MPI-OM results (Voigt and Marotzke 2009; Voigt et al. 2011). Pierrehumbert et al. (2011) note however, that initiation of

snowball states using ECHAM remains easier than using CAM when the snow and sea-ice albedos are all set to 0.6, suggesting differences in atmospheric circulation or cloud forcing between the models must also play a role. Further supporting this conclusion, snowball states remain easier to simulate in ECHAM in the mixed-layer ocean configuration as compared to CAM (Pierrehumbert et al. 2011).

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Differences in complexity between models used to approach the Snowball problem also certainly play a role in expanding the range of initiation thresholds. The absence of dynamic ocean circulation and sea ice are limitations of our study. Previous studies using both diffusive energy balance models (Rose and Marshall 2009) and general circulation models (Poulsen et al. 2001) find that ocean heat transport stabilizes sea-ice margins at higher latitudes. As a result, addition of ocean heat transport to our model would imply increased TSI thresholds for global glaciation. Addition of sea-ice dynamics may limit the stabilizing effect of ocean heat transport, however. Experiments by Yang et al. (2012a) show a rapid transition between a stable state with ~40% sea ice coverage and one with ~60% ice coverage when CO₂ concentration is reduced from 70 to 50 ppm at 94% TSI, a change in radiative forcing of about 1.8 W m⁻² (Myhre et al. 1998). Global ice coverage occurs when CO₂ is further lowered to 17.5 ppm, suggesting this high ice-coverage state is only stable in a small parameter space. Curiously, this rapid transition was not observed when TSI is lowered instead of CO₂. Yang et al. (2012a) attribute this to the latitudinal dependence of solar radiative forcing compared to the latitudinal independence of greenhouse forcing for well-mixed greenhouse gases, using a modified EBM to show that when the ice line expands equatorward of 20°, sensitivity due to solar forcing is stronger than sensitivity due to greenhouse gas forcing.

Jormungand states, defined by Abbot et al. (2011) as stable ice latitudes equatorward of 25°, were simulated for all four geographies. Unlike the results in Abbot et al. (Abbot et al. 2011), we do not find a rapid transition between ice-free conditions and the Jormungand state. Immediate transition between an ice-free state and the Jormungand state is attributed to model deficiencies associated with the idealized configuration of CAM (aquaplanet with no aerosols, modified greenhouse gas concentrations, no ocean heat transport, detailed further in Pierrehumbert et al. 2011). Our choice to vary TSI instead of greenhouse gas concentration may be responsible for the absence of the strong bifurcation between ice-free conditions and the Jormungand state. Radiative forcing was never high enough to simulate an ice-free state in our experiments, but states with stable ice latitudes intermediate to ice-free and Jormungand states likely result from a smaller reduction in radiative forcing in the extratropics from TSI reductions relative to the reduction in radiative forcing resulting from reduced greenhouse gas concentrations. For the same reduction in global radiative forcing, reductions in greenhouse forcing affect all latitudes similarly while the impact of TSI reductions is most pronounced in the tropics due to the latitudinal dependence of solar insolation (Yang et al. 2012a,b).

A final attractive aspect of the low-latitude continent Snowball Earth initiation theory was that as glaciers formed on land, sea level would drop exposing continental shelves and further increasing surface albedo as land replaces ocean. Instead, our results suggest as sea level drops and shallow seas in the tropics are replaced by land, Snowball Earth initiation becomes more difficult as continental interiors dry, land-sea contrast becomes more pronounced, and a Walker-style circulation intensifies.

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b) Paleogeographic impacts on climate sensitivity

The impact of paleogeography on global temperatures remains an open question. On one hand, several studies report higher global temperatures when tropical land area is increased as a result of suppressed precipitation and evaporative cooling in the tropics (Barron et al. 1984). Our results also suggest tropical continents increase global temperatures at decreased TSI, though our simulations don't include the effect of geochemical processes such as silicate weathering.

Conversely, simulations with ECHAM/MPI-OM show lower global temperatures when tropical land area is increased (Voigt et al. 2011), and a more facile Snowball Earth initiation. Initiation of a Snowball state using a modern continental arrangement and topography required a 1.5% greater reduction (to 94%) in TSI than initiation of a Snowball state using a Marinoan reconstruction (Voigt and Marotzke 2009). Voigt et al. (2011) relate this difference to elevated albedo and a weakened greenhouse effect when land is shifted equatorward. This comparison is complicated and possibly compromised by inclusion of land plants in the modern simulation, however, which may have promoted a stronger greenhouse effect and lower land albedo.

In contrast, geochemical modeling suggests lower global temperatures due to decreased greenhouse gas concentrations when land is positioned in the tropics due to enhanced silicate weathering (Marshall et al. 1988; Worsley and Kidder 1991). While no dynamic carbon cycle was modeled in these experiments, precipitation over the continent in the RECT experiment is extremely low, suggesting that weathering would also be low despite warm tropical continents. Absence of land plants and orography reduce moisture transport over large tropical continents. Precipitation and runoff increases following supercontinental breakup suggest higher weathering rates and a reduction of greenhouse forcing (Donnadieu et al. 2004; 2006). In our experiments, low latitude (≤30°) precipitation rates over land in the MAR simulation are nearly three times higher than in the RECT experiment at 95 and 93% TSI, upholding the conclusion that

continental fragmentation may be as important as latitudinal position to weathering rates. Apart from affecting precipitation patterns, the absence of land plants in the Neoproterozoic would limit the efficiency of silicate weathering, as root growth, organic acid secretion, and enhanced runoff retention by plants all enhance weathering rates (Berner 1997), though microbial communities on Neoproterozoic continents may have enhanced weathering rates to a lesser extent (Lenton and Watson 2004). Chemical weathering rates may also have been enhanced in fine-grained soils (Keller and Wood 1993) prior to land plant evolution.

Rising anthropogenic greenhouse gas emissions highlight the importance of quantifying modern climate sensitivity to changes in radiative forcing. Despite intense study, modern and paleo- climate sensitivities remain difficult to constrain (Roe and Baker 2007; Knutti and Hegerl 2008; Rohling et al. 2012). In the absence of long-term observations, paleoclimate archives have been used to estimate climate sensitivity on long time scales in response to changes in greenhouse forcing by comparing changes in proxy CO₂ with proxy temperature (e.g., Royer et al. 2007). A potential problem with this approach is that boundary conditions specific to past time periods may influence climate sensitivity. In fact, our results show the danger of this approach by suggesting geography strongly modulates climate processes and sensitivity. A high degree of land-sea contrast in the tropics increases the importance of direct radiative heating at the expense of greenhouse heating, lowering the climate sensitivity to changes in the water vapor feedback with TSI reductions.

5. Conclusions

Our simulations do not support the hypothesis that increased tropical land mass would facilitate initiation of a Snowball Earth event. While raising surface albedo, large tropical landmasses alter

the partitioning of energy in the tropics such that global climate is less sensitive to reductions in top of the atmosphere shortwave forcing. The energy budget in AQP has a stronger dependence on greenhouse forcing while the energy budget of MOD-NT, MAR, and RECT show stronger dependence on direct radiative forcing. As greenhouse forcing decreases exponentially with temperature following specific humidity scaling by the Clausius-Clapeyron equation, direct top of the atmosphere shortwave forcing decreases linearly with incremental decreases in TSI. When there is a high degree of land/sea contrast in the deep tropics, an amplified Walker circulation raises tropical temperatures, enhances evaporation over tropical oceans and maintains a stronger greenhouse effect at reduced TSI than when no land is present. Topography also provides a strong control on Snowball Earth initiation. Including topography in a modern continental configuration lowered the TSI threshold for Snowball Earth initiation by up to 2% by altering cloud and snow coverage and the meridional albedo gradient, reducing the equator-to-pole temperature gradient.

More generally, this study demonstrates that paleogeography and paleotopography can have a significant impact on climate sensitivity by repartitioning energy in the climate system. This must be considered when using paleoclimate records to estimate climate sensitivity to changing atmospheric CO₂.

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640 Tables

Configuration	TSI	Years Run	Result		
S	(% of modern)				
Aquaplanet	100	400-530	Ice latitude of 39.1°S, 39.5°N		
(AQP)	99	400-500	Ice latitude of 35.2°S, 35.7°N		
	97	400-651	Ice latitude of 20.4°S, 22.7°N		
	95	400-567	Ice latitude of 15.0°S, 16.0°N		
	93	400-553	Ice latitude of 12.2°S, 13.0°N		
	92	400-595	Ice latitude of 9.4°S, 10.6°N		
	91	400-525	Snowball by year 485		
	90	400-511	Snowball by year 431		
	88	400-575	Snowball by year 420		
	85	400-511	Snowball by year 414		
	80	400-539	Snowball by year 410		
Modern – no	100	400-507	Ice latitude of 41.1°S, 42.1°N		
topography	97	500-569	Ice latitude of 33.0°S, 36.0°N		
(MOD-NT)	95	500-653	Ice latitude of 26.5°S, 30.0°N		
	93	500-620	Ice latitude of 20.1°S, 23.7°N		
	90	500-625	Ice latitude of 14.5°S, 17.4°N		
	89	500-611	Ice latitude of 12.5°S, 15.5°N		
	88	500-697	Ice latitude of 10.0°S, 11.8°N		
	87	500-647	Snowball by year 534		
	85	500-611	Snowball by year 511		
	80	500-611	Snowball by year 505		
Marinoan	100	400-539	Ice latitude of 40.6°S, 39.3°N		
(MAR)	97	499-602	Ice latitude of 34.6°S, 32.2°N		
	95	499-655	Ice latitude of 30.4°S, 27.4°N		
	93	499-610	Ice latitude of 26.9°S, 23.2°N		
	90	499-617	Ice latitude of 21.5°S, 18.1°N		
	88	499-681	Ice latitude of 18.1°S, 14.4°N		
	87	499-723	Ice latitude of 15.8°S, 12.9°N		
	86	499-674	Ice latitude of 13.6°S, 10.5°N		
	85	499-652	Snowball by year 572		
	80	499-638	Snowball by year 504		
Rectangle	100	400-537	Ice latitude of 38.8°S, 39.1°N		
(RECT)	97	510-565	Ice latitude of 33.3°S, 33.7°N		
	95	510-635	Ice latitude of 29.8°S, 29.9°N		
	93	510-579	Ice latitude of 26.0°S, 26.4°N		
	90	510-635	Ice latitude of 20.6°S, 21.4°N		
	88	510-635	Ice latitude of 17.3°S, 18.0°N		
	85	510-663	Ice latitude of 10.2°S, 12.4°N		
	84	510-772	Snowball by year 533		

	83	510-642	Snowball by year 519
	80	510-621	Snowball by year 515
Modern –	100	1-142	Ice latitude of 41.4°S, 38.6°N
modern	97	101-209	Ice latitude of 28.1°S, 30.8°N
topography	95	101-222	Ice latitude of 22.2°S, 25.2°N
(MOD-WT)	93	101-219	Ice latitude of 18.4°S, 20.7°N
	90	101-250	Ice latitude of 11.4°S, 15.0°N
	89	101-422	Snowball by year 198
	88	101-310	Snowball by year 124
	85	101-229	Snowball by year 110
	80	101-242	Snowball by year 107

Table 1. Summary of model experiments performed and result of run when terminated after

years indicated. TSI is listed as a percentage of the modern value.

	Continental Arrangement			
Variable	AQP	MOD-NT	MAR	RECT
Land Fraction, global	0	0.323	0.278	0.270
Land Fraction, 30°N/S	0	0.288	0.481	0.385
Surface Albedo	0.310	0.326	0.356	0.353
Planetary Albedo	0.406	0.397	0.403	0.402
Mean Surface Temperature (K)	273.4	274.0	272.6	272.4
Mean Annual Ice Latitude	39.5°N,	42.1°N,	39.3°N,	39.1°N,
	39.1°S	41.1°S	40.6°S	38.8°S

Table 2. Results for AQP, MOD-NT, MAR, and RECT configurations at 100% TSI.

Figure Captions

Figure 1. Surface temperature contour plots at 100% (left column) and 95% TSI (right column) for the AQP (top), MOD-NT, MAR, and RECT (bottom) experiments. Temperature structures are highly zonal, with some asymmetries near land-ocean transitions. At 100% TSI, the four geographies have a highly similar temperature structure with mean annual tropical and polar temperatures around 295K and 225K respectively. At 95% TSI, the temperature structures of the four geographies diverge – AQP averages at least 15K colder than the MOD-NT, MAR, and RECT geographies at all latitudes.

Figure 2. Mean annual surface temperature (a), mean annual ice margin (b), and surface (c) and planetary (d) albedos as a function of TSI. A strong bifurcation notes the transition to a snowball state for each configuration as TSI is decreased. AQP enters the snowball state most readily (TSI < 92%), followed by MOD-NT, MAR, and RECT when TSI is reduced below 88%, 86%, and

659 85% respectively.

Figure 3. Components of surface energy balance at 100% TSI for all four geographies. (a) Net surface shortwave radiation absorbed, (b) net surface longwave radiation emitted, (c) total cloud forcing, (d) total greenhouse forcing (difference in upwelling clearsky longwave radiation at the top of the atmosphere and the surface), and (e) cloud forcing over the ocean. All plots are in W m⁻². Geographies with large amounts of tropical land area (i.e., MAR and RECT) have higher net shortwave in the deep tropics than geographies with little or no tropical land area (i.e., AQP) primarily due to a lower cloud fraction. This energy disparity is ameliorated by increased longwave emission and a smaller greenhouse forcing for the high land configurations, leading to similar tropical temperatures across all four geographies.

Figure 4. Total annual average cloud fraction for all four geographies at 100% TSI. Cloud fractions are highest over regions of vigorous convection over the ocean and lower in the subtropics and particularly over land. Cloud fractions are notably low over MAR and RECT geographies due to the presence of large continents in the subtropics, though cloud fractions are low over the entirety of the RECT continent, while cloud fractions remain higher in the deep tropics over the smaller continents of the MAR geography.

Figure 5. Weighted area average (from 30°N/S) components of the surface energy balance as a function of TSI for all four geographies. (a) Net surface shortwave radiation absorbed, (b) net surface longwave radiation emitted, (c) total cloud forcing, and (d) total greenhouse forcing (difference in upwelling clearsky longwave radiation at the top of the atmosphere and the surface). All plots are in W m⁻². At 100% TSI, the AQP and MOD-NT configurations are more reliant on greenhouse forcing to maintain tropical temperatures, and are most sensitive to decreases in TSI as the amount of atmospheric water vapor decreases with TSI. Cloud forcing in

the MAR and RECT configurations are more positive by ~10 W m⁻² prior to Snowball initiation,

and higher absorbed shortwave fluxes that decrease less rapidly than greenhouse forcing.

Figure 6. Vertically integrated water vapor as a function of (a) average global surface temperature (K) and (b) TSI from 30°N/S for all four geographies. At 100% TSI, vertically integrated water vapor is notably lower for high tropical land configurations MAR and RECT, with the RECT geography having 35% less tropical atmospheric water vapor than the AQP geography. Temperature and specific humidity are tightly coupled (a), as vertically integrated water vapor flows average global surface temperature as suggested by the Clausius-Clapeyron equation. As TSI is reduced however, the rate of decrease of vertically integrated water vapor (b) is higher for AQP than the other configurations with land since the energy budget has a larger greenhouse forcing term.

Figure 7. Annually and meridionally (5°N/S) averaged longitudinal/height cross-sections of vertical pressure velocities (mb day⁻¹) and relative humidity (%) for the AQP (a,c) and RECT

(b,d) geographies at 95% TSI. Black contour lines in b and d indicate specific humidity in 1 g kg⁻¹ intervals. Vertical pressure velocities in the AQP configuration are ubiquitously negative in the troposphere, suggesting vigorous upward motion at all longitudes in the annual average. In contrast, the RECT configuration shows strong upward motion over the oceans, but strong subsidence in the free troposphere limiting weaker upward motion to the boundary layer over the continent. Likewise, the troposphere is near saturation at the surface and near the tropopause in the AQP experiment at all longitudes while the troposphere is near saturation in the boundary layer and tropopause over the ocean, but is markedly drier over the continent, particularly over the continent's western half. Over the ocean, near surface specific humidity is nearly three times higher in the RECT experiment than in the AQP experiment. Additionally, higher specific humidity in the middle-troposphere of the RECT experiment results in greenhouse forcing less sensitive to TSI than in the AQP case.

Figure 8. (a) Total surface water flux (Pg s⁻¹) and (b) surface water flux normalized by ocean fraction (Pg s⁻¹ ocean fraction⁻¹) from 30°N/S for all four geographies as a function of TSI. At 100% TSI, surface water flux is highest for the AQP configuration but falls most rapidly with TSI. Surface water flux normalized by ocean fraction shows evaporation in tropical oceans is more efficient for the MAR and RECT configurations than for the AQP configurations.

Figure 9. Difference between RECT and AQP vertical cloud fractions at 95% TSI. Both zonally averaged differences (a) and longitudinal differences in the deep tropics (b, 10°N/S) are shown. Contour interval is 0.05 and negative values are shaded. The Walker circulation simulated decreases the amount of low- and midlevel clouds and increases the amount of high-level clouds,

increasing the shortwave cloud forcing in the tropics by 20-40 W m⁻² while only decreasing the 723 longwave cloud forcing by 10-15 W m⁻². 724 725 726 Figure 10. (a) Mean annual surface temperature (K) and (b) ice margin latitude for the MOD-727 WT and MOD-NT experiments as a function of TSI. Albedos are tightly coupled to temperature 728 and are not shown. Adding topography to the MOD-NT continental configuration raises the 729 snowball bifurcation point from 88% to 90% modern TSI, though above the bifurcation points, 730 little difference exists between global surface temperature and ice margin latitude. 731 732 Figure 11. Contour difference plots between MOD-WT and MOD-NT for (a) surface 733 temperature (K), (b) planetary albedo, (c) cloud fraction, and (d) snow depth (meters liquid water 734 equivalent) at 90% TSI. Adding topography reduces the equator-to-pole temperature gradient (a) 735 by raising subtropical albedo and lowering midlatitude albedo (b), particularly in the Northern 736 Hemisphere where the majority of modern land is located. Albedo changes are driven by changes 737 in cloud cover and snow depth (c and d).

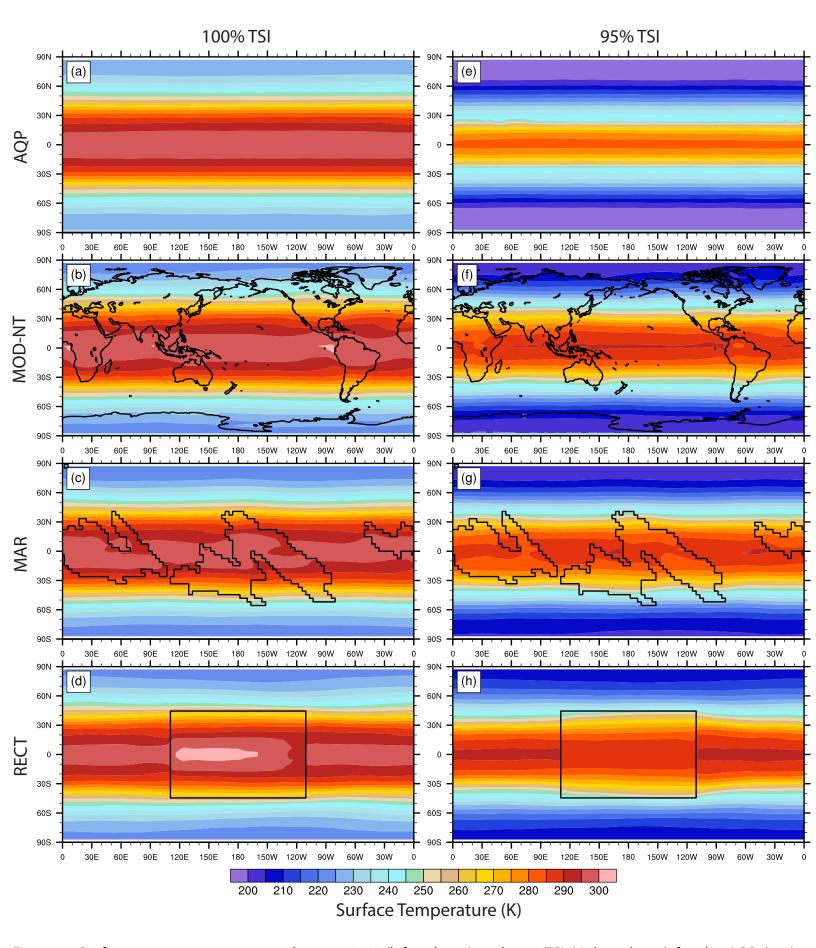


Figure 1. Surface temperature contour plots at 100% (left column) and 95% TSI (right column) for the AQP (top), MOD-NT, MAR, and RECT (bottom) experiments. Temperature structures are highly zonal, with some asymmetries near land-ocean transitions. At 100% TSI, the four geographies have a highly similar temperature structure with mean annual tropical and polar temperatures around 295K and 225K respectively. At 95% TSI, the temperature structures of the four geographies diverge – AQP averages at least 15K colder than the MOD-NT, MAR, and RECT geographies at all latitudes.

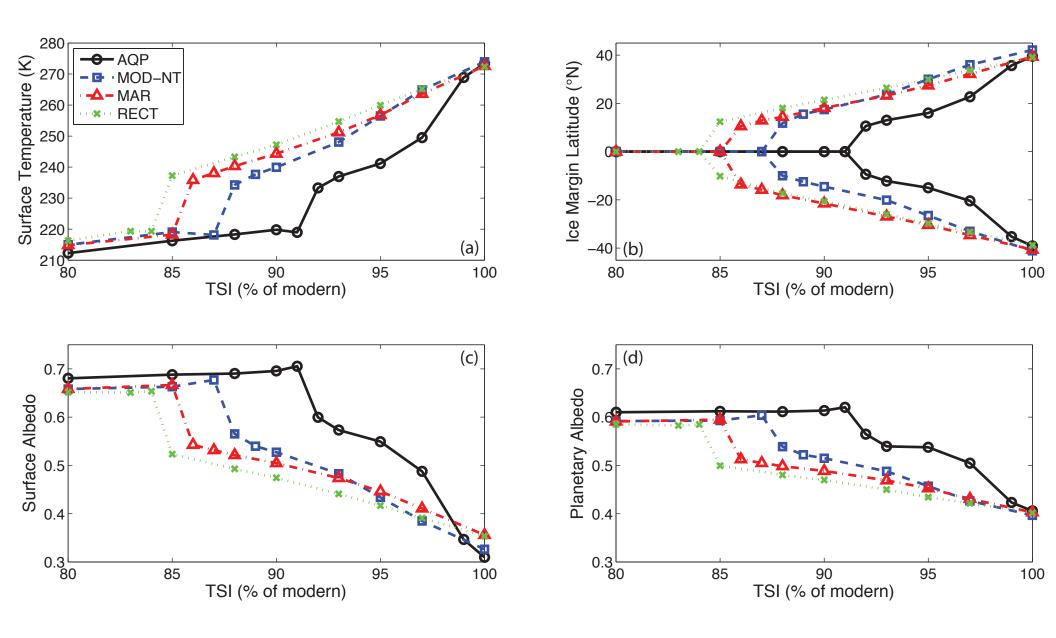
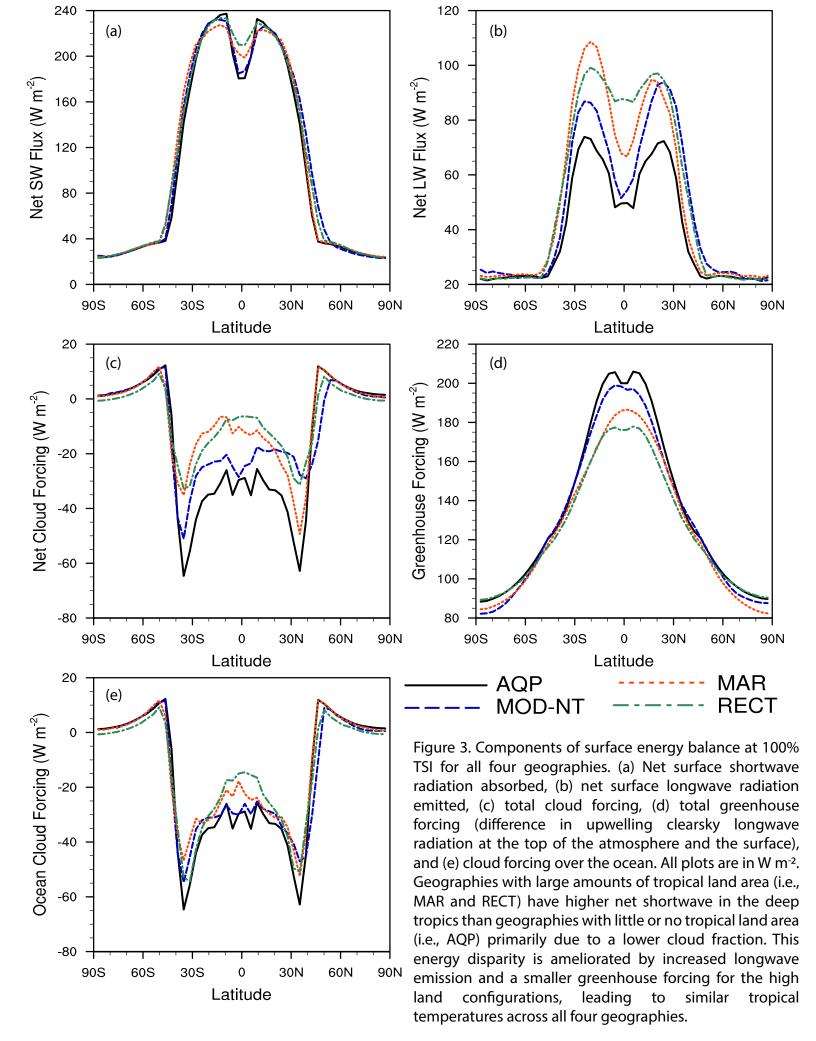


Figure 2. Mean annual surface temperature (a), mean annual ice margin (b), and surface (c) and planetary (d) albedos as a function of TSI. A strong bifurcation notes the transition to a snowball state for each configuration as TSI is decreased. AQP enters the snowball state most readily (TSI < 92%), followed by MOD-NT, MAR, and RECT when TSI is reduced below 88%, 86%, and 85% respectively.



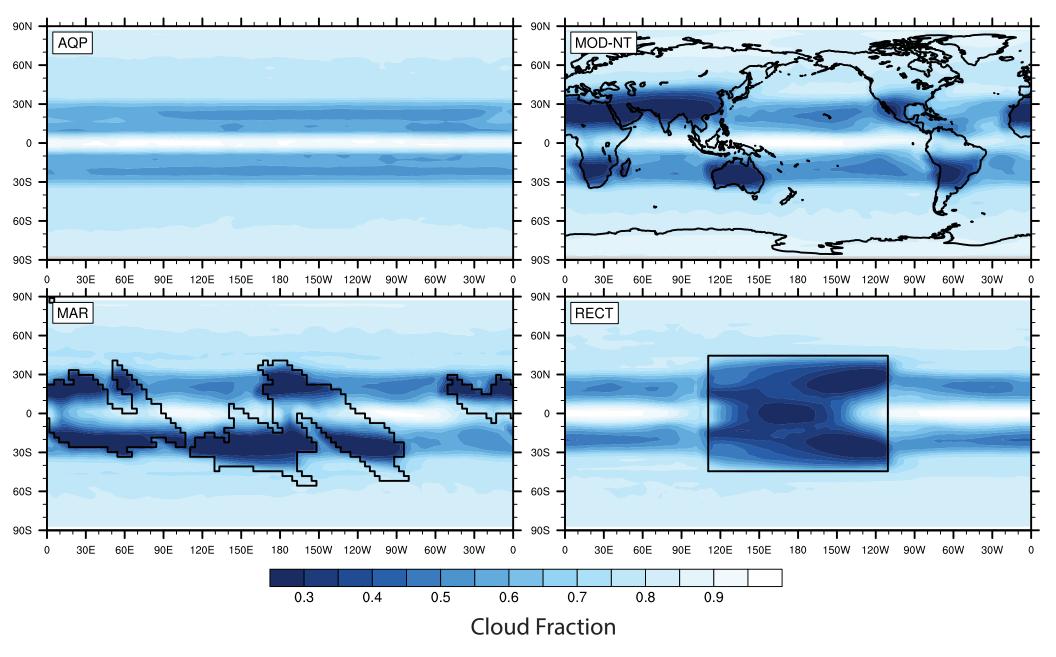


Figure 4. Total annual average cloud fraction for all four geographies at 100% TSI. Cloud fractions are highest over regions of vigorous convection over the ocean and lower in the subtropics and particularly over land. Cloud fractions are notably low over MAR and RECT geographies due to the presence of large continents in the subtropics, though cloud fractions are low over the entirety of the RECT continent, while cloud fractions remain higher in the deep tropics over the smaller continents of the MAR geography.

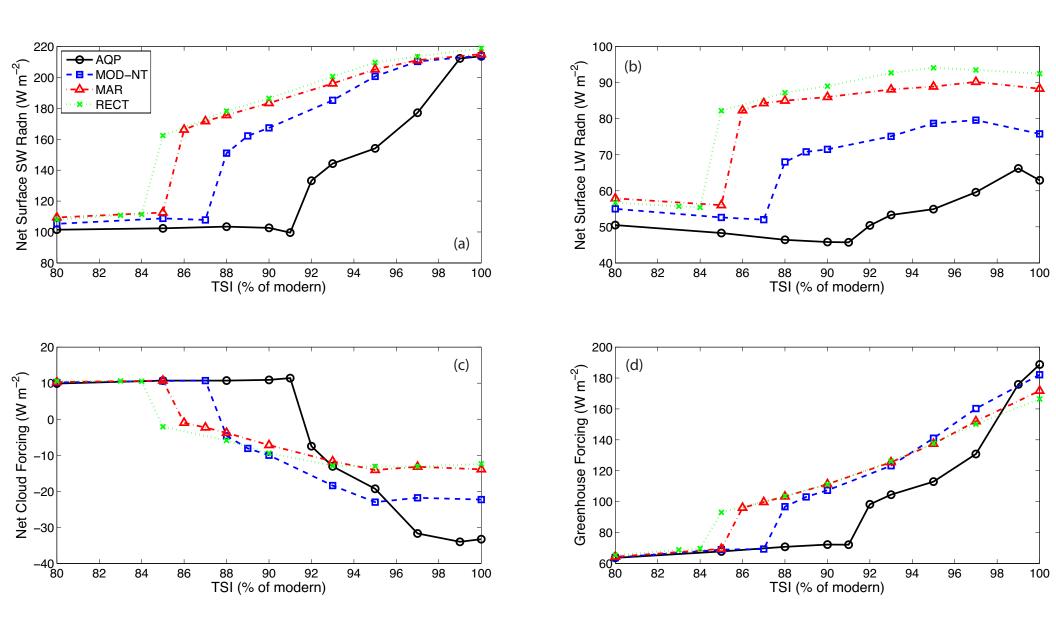


Figure 5. Weighted area average (from 30°N/S) components of the surface energy balance as a function of TSI for all four geographies. (a) Net surface shortwave radiation absorbed, (b) net surface longwave radiation emitted, (c) total cloud forcing, and (d) total greenhouse forcing (difference in upwelling clearsky longwave radiation at the top of the atmosphere and the surface). All plots are in W m⁻². At 100% TSI, the AQP and MOD-NT configurations are more reliant on greenhouse forcing to maintain tropical temperatures, and are most sensitive to decreases in TSI as the amount of atmospheric water vapor decreases with TSI. Before snowball initiation, cloud forcing increases modestly for the AQP and MOD-NT configurations as TSI is reduced, but not enough to overcome the rapid decrease in greenhouse forcing. In contrast, the MAR and RECT configurations show cloud forcings that are more positive by ~10 W m⁻² prior to snowball initiation, and higher absorbed shortwave fluxes that decrease less rapidly than greenhouse forcing.

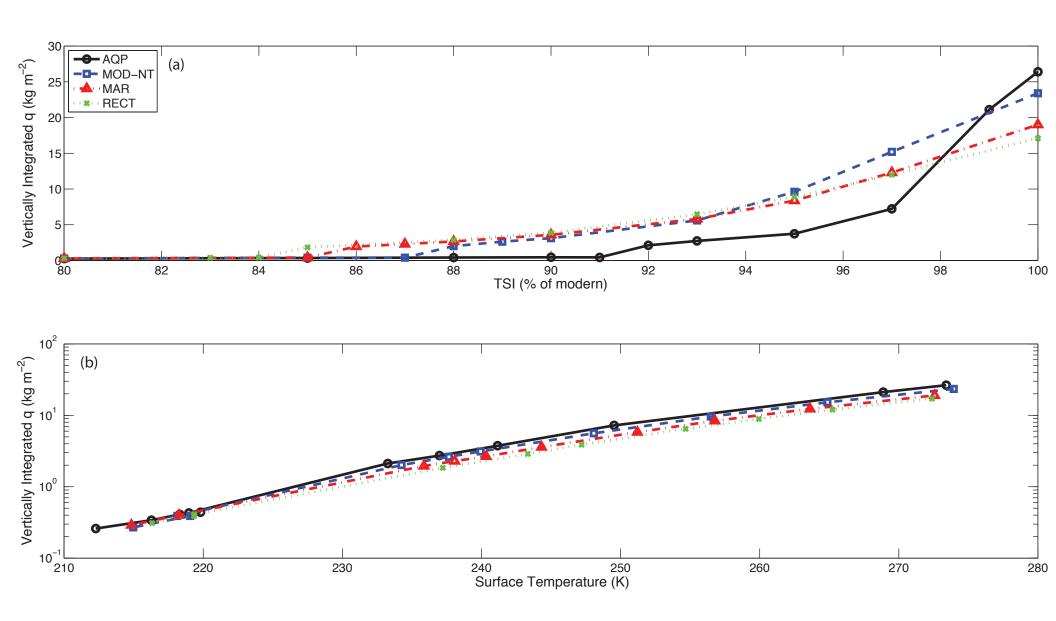


Figure 6. Vertically integrated water vapor as a function of (a) average global surface temperature (K) and (b) TSI from 30°N/S for all four geographies. At 100% TSI, vertically integrated water vapor is notably lower for high tropical land configurations MAR and RECT, with the RECT geography having 35% less tropical atmospheric water vapor than the AQP geography. Temperature and specific humidity are tightly coupled (a), as vertically integrated water vapor follows average global surface temperature as suggested by the Clausius-Clapeyron equation. As TSI is reduced however, the rate of decrease of vertically integrated water vapor (b) is higher for AQP than the other configurations with land since the energy budget has a larger greenhouse forcing term.

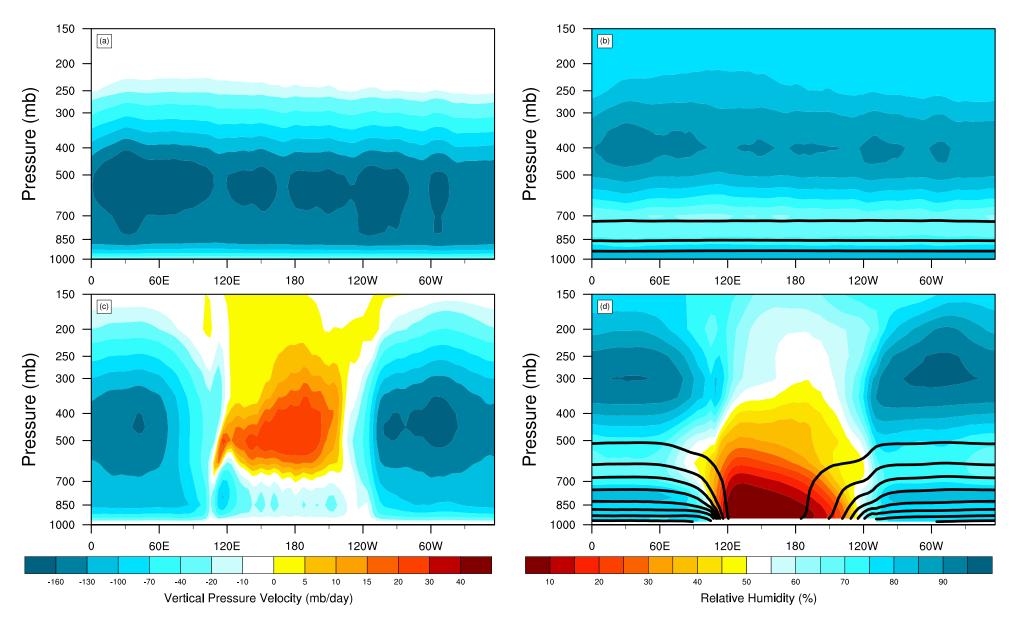


Figure 7. Annually and meridionally (5°N/S) averaged longitudinal/height cross-sections of vertical pressure velocities (mb/day) and relative humidity (%) for the AQP (a,c) and RECT (b,d) geographies at 95% TSI. Black contour lines in b and d indicate specific humidity in 1 g/kg intervals. Vertical pressure velocities in the AQP configuration are ubiquitously negative in the troposphere, suggesting vigorous upward motion at all longitudes in the annual average. In contrast, the RECT configuration shows strong upward motion over the oceans, but strong subsidence in the free troposphere limiting weaker upward motion to the boundary layer over the continent. Likewise, the troposphere is near saturation at the surface and near the tropopause in the AQP experiment at all longitudes while the troposphere is near saturation in the boundary layer and tropopause over the ocean, but is markedly drier over the continent, particularly over the continent's western half. Over the ocean, near surface specific humidity is nearly three times higher in the RECT experiment than in the AQP experiment. Additionally, higher specific humidity in the mid-troposphere of the RECT experiment results in greenhouse forcing less sensitive to TSI than in the AQP case.

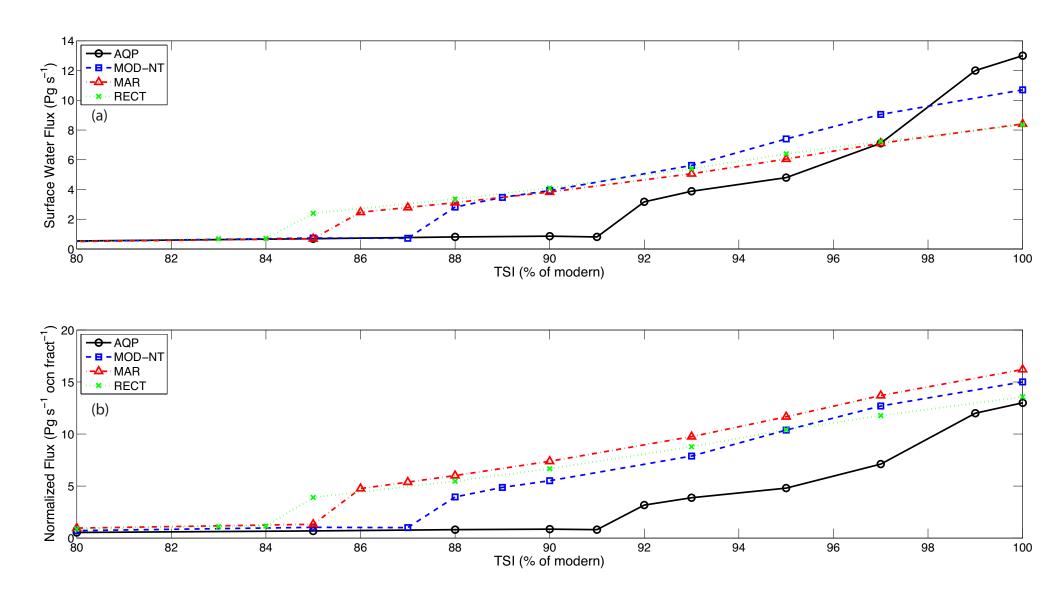


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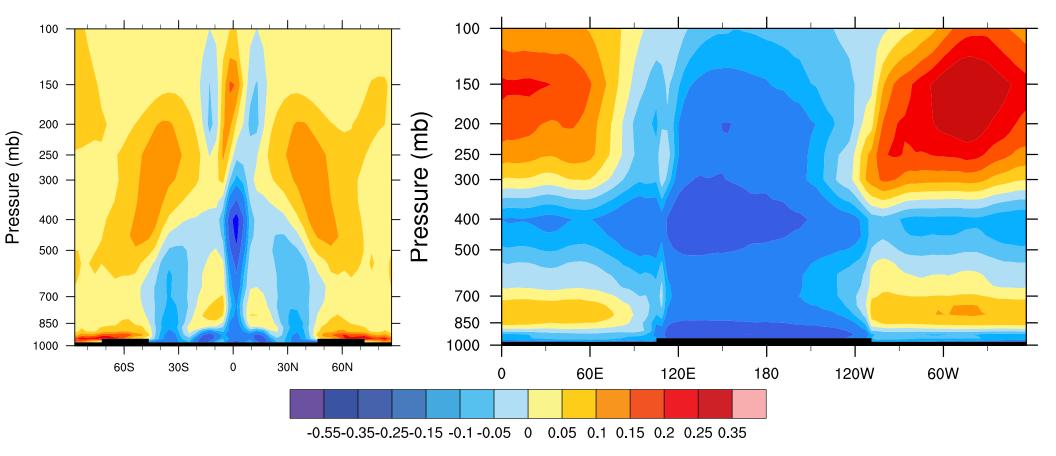


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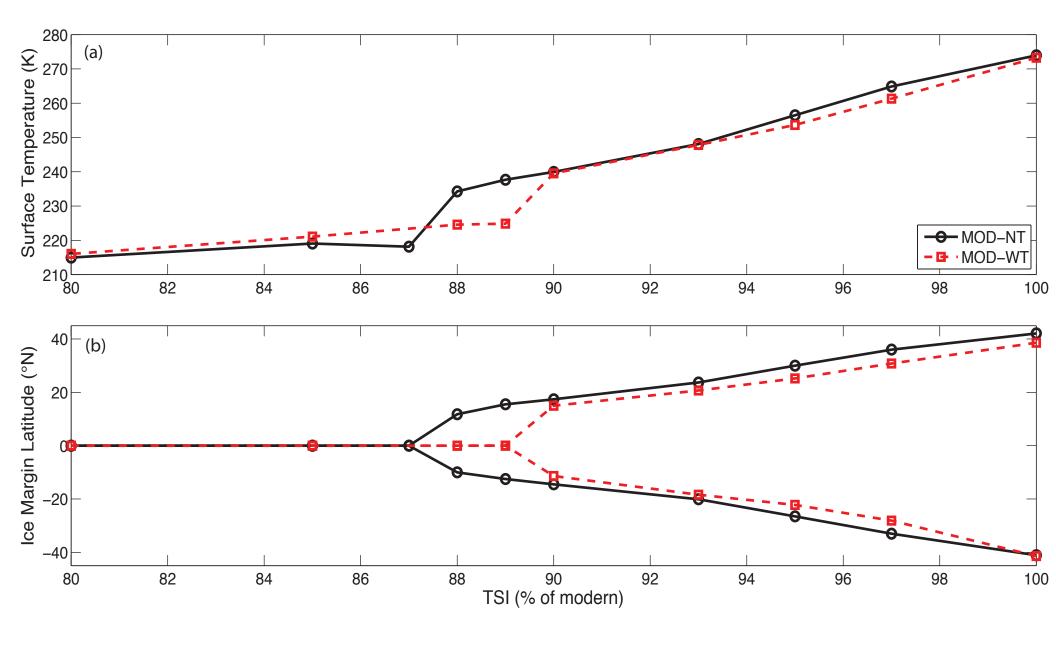


Figure 10. (a) Mean annual surface temperature and (b) ice margin latitude for the MOD-WT and MOD-NT experiments as a function of TSI. Albedos are tightly coupled to temperature and are not shown. Adding topography to the MOD-NT continental configuration raises the snowball bifurcation point from 88% to 90% modern TSI, though above the bifurcation points, little difference exists between global surface temperature and ice margin latitude.

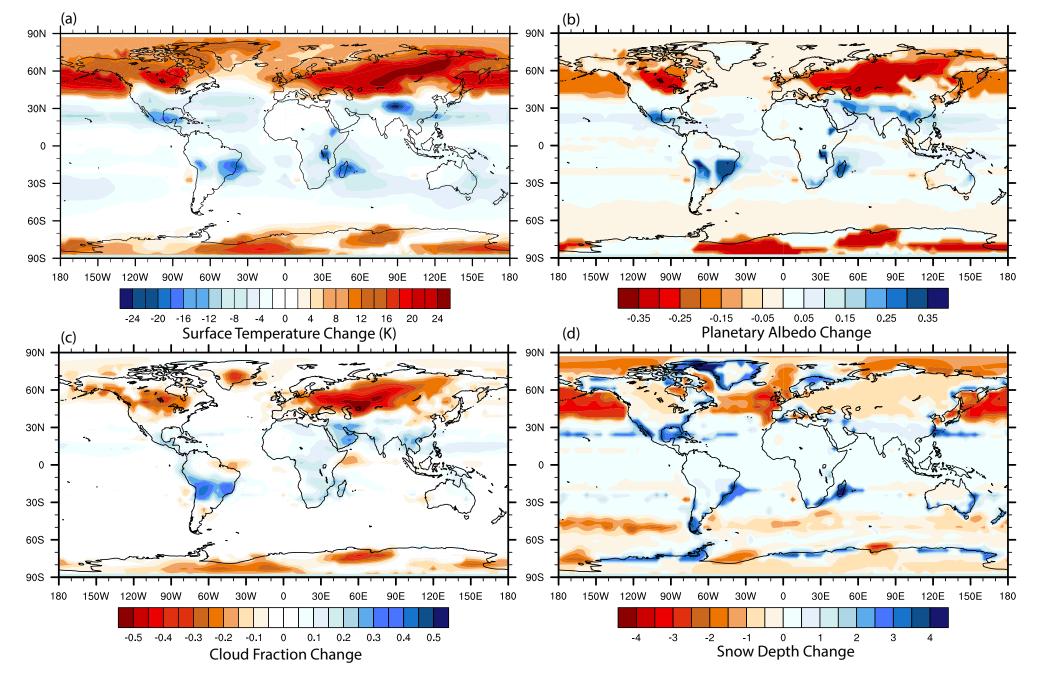


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